

Section 1
Problem 1

GRING 2001



x marks center of mass, CM
Each orbit is an ellipse with foci at CM.
If radius is b , distance from center of
circle to CM is $b/2$
But this ought to be ae where a is semi major
(if unprojected).

Semi major axis projects to b

$$\text{therefore } \frac{ae}{a} = \frac{b/2}{b} \Rightarrow e = \frac{1}{2}.$$

1-2 SHO Hamiltonian - raising & lowering ops.

$$\frac{p^2}{2m} + kx^2 = \mathcal{H}$$

$$\left(\frac{p}{\sqrt{2m}} + i\sqrt{k}x\right)\left(\frac{p}{\sqrt{2m}} - i\sqrt{k}x\right) = a^+ a = \mathcal{H}$$

$$\left. \begin{aligned} a^+ &= \frac{p}{\sqrt{2m}} + i\sqrt{k}x \\ a &= \frac{p}{\sqrt{2m}} - i\sqrt{k}x \end{aligned} \right\} \Rightarrow a^+ - a = 2i\sqrt{k}x$$

$$a^+ |n\rangle = \kappa |n+1\rangle \quad a |0\rangle = 0$$

Some constant

$$a |n\rangle = \kappa' |n-1\rangle$$

Then $\langle n' | x | n \rangle \propto \langle n' | a^+ - a | n \rangle$

$$\neq 0 \quad \text{if} \quad n' \pm 1 = n \quad \text{if } n', n \neq 0$$

$$n' = 0, n = 1$$

$$n' = 1, n = 1$$

$$\Delta E = \hbar \omega$$

$$1.3 \quad m\ddot{x} = -mg + A\rho C_D \dot{x}$$

$$m\dot{v} = -mg + A\rho C_D v$$

$$\text{Solution: } v(t) = \alpha e^{-\lambda t} + \beta, \quad \beta = v_{\text{term}}$$

$$\dot{v} = -\lambda \alpha e^{-\lambda t} \quad r = A\rho C_D$$

$$\text{Then } -m\lambda \alpha e^{-\lambda t} = -mg + r\alpha e^{-\lambda t} + r\beta$$

$$\text{Two conditions: } -m\lambda \alpha e^{-\lambda t} = r\alpha e^{-\lambda t}$$

$$\text{and} \quad \lambda = r/m = A\rho C_D/m$$

$$-mg + r\beta = 0 \Rightarrow \beta = \frac{mg}{r} = \frac{mg}{A\rho C_D}$$

$$\text{For a person: } A \sim 1 \text{ m}^2$$

$$m \sim 100 \text{ kg}$$

$$g \sim 10 \text{ m/s}^2$$

$$C_D \sim 10^{-3} \text{ m/s}$$

$$\rho = \frac{149}{22.4 \times 10^3 \text{ cm}^3} = 0.6 \times 10^{-3} \frac{\text{g}}{\text{cm}^3} \times \frac{1 \text{ kg}}{10^3 \text{ g}} \times \frac{10^6 \text{ cm}^3}{1 \text{ m}^3} = 0.6 \frac{\text{kg}}{\text{m}^3}$$

$$r = 1 \text{ m}^2 \times 0.6 \frac{\text{kg}}{\text{m}^3} \times \frac{10 \text{ m}}{\text{s}} = 6 \text{ kg/s}$$

$$\beta = \frac{100 \text{ kg} \times 10 \text{ m/s}^2}{6 \text{ kg/s}} = 0.16 \times 10^3 \text{ m/s} = 160 \text{ m/s}$$

$$\frac{1}{\lambda} = \frac{100 \text{ kg}}{6 \text{ kg/s}} = 16 \text{ s}$$

1-4 $\psi_0 = \begin{pmatrix} 1 \\ \vdots \\ 1 \end{pmatrix}$ gives $H\psi = E_0 N\psi$

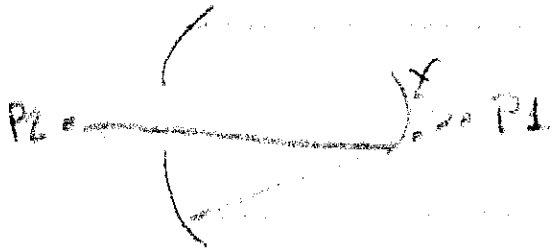
An orthogonal ev. must also be an ev. \rightarrow

$$\psi_1 = \begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix}, \psi_2 = \begin{pmatrix} -1 \\ 1 \\ 0 \end{pmatrix}, \dots$$

Eigenvalues are NE_0 and 0. Degeneracy of NE_0 is 1, 0 is $n-1$

Section 2

Problem 1 Change parabolic to paraboloidal; add word perfect
betro focus, twice



Hyperboloid. All paths to P_1 are same length. All paths to P_2 should likewise be same length. The second mirror adds a line segment P_2X and subtracts P_1X from the path length.

$\overline{P_2X} - \overline{P_1X} = \text{constant}$. But a hyperbola is the locus of points the difference of whose distances to two points is a constant. Rotate about symmetry axis & get a hyperboloid.

$$2-2 \quad Z = \sum_{l=0}^{\infty} (2l+1) e^{-l(l+1)\epsilon/kT} \quad \epsilon = \hbar\omega$$

Internal energy $U = kT^2 \frac{\partial Z}{\partial T}$

$$Z \rightarrow \int_0^{\infty} 2l e^{-l^2 \epsilon/kT} dl \quad \text{for } l \text{ large}$$

$$x = \frac{l^2 \epsilon}{kT} \rightarrow ; \quad dx = \frac{2l \epsilon}{kT} dl$$

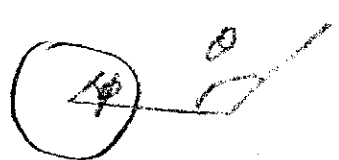
Then $Z = \int_0^{\infty} e^{-x} \frac{kT}{\epsilon} dx$

$$= -\frac{kT}{\epsilon} e^{-x} \Big|_0^{\infty} = \frac{kT}{\epsilon}$$

$$U = kT^2 \frac{\partial Z}{\partial T} = kT^2 \frac{k}{\epsilon} = \frac{k^2 T^2}{\epsilon}$$

$$C_v = \frac{\partial U}{\partial T} = \frac{2k^2 T}{\epsilon}$$

$$2-3 \quad H = \frac{p_\theta^2}{2m} + \frac{p_\phi^2}{2m} = -\frac{\hbar^2}{2m} \left(\frac{1}{a^2} \frac{\partial^2}{\partial \theta^2} + \frac{1}{b^2} \frac{\partial^2}{\partial \phi^2} \right)$$



$$\psi \propto e^{i l \phi} e^{i n \theta}$$

$$H\psi = -\frac{\hbar^2}{2m} \left(\frac{(-n^2)}{a^2} + \frac{(-l^2)}{b^2} \right) \psi = E\psi$$

$$\frac{\hbar^2}{2m} \left(\frac{n^2}{a^2} + \frac{l^2}{b^2} \right) = E_{nl}$$

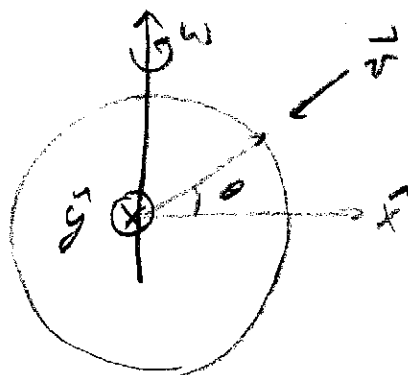
All states have infinite degeneracy

$$\psi \propto e^{i l \phi + \phi_0} e^{i n \theta + \theta_0}$$

$$2-4 \quad F_{\text{cor}} = -2m \vec{\omega} \times \vec{v}$$

$$\vec{v} = v \cos \theta \hat{x} + v \sin \theta \hat{y}$$

$$\vec{\omega} = \omega \hat{z}$$



$$\vec{\omega} \times \vec{v} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ v \cos \theta & v \sin \theta & 0 \\ 0 & 0 & \omega \end{vmatrix}$$

$$= 0 \hat{x} + -\omega v \cos \theta \hat{y} + 0 \hat{z} = -\omega v \cos \theta \hat{y}$$

$$\vec{F}_{\text{cor}} = +2m \omega v \cos \theta \hat{y} \quad (\hat{y} = \text{east})$$

$$\text{Fall time: } d = \frac{1}{2} g t_0^2 \Rightarrow t_0 = \sqrt{\frac{2d}{g}} = \sqrt{\frac{440 \text{ m}}{10 \text{ m/s}^2}} = 6.7 \text{ s}$$

$$\text{Then, } \Delta v_y = 2 \omega v \cos \theta$$

$$v_r = -gt \Rightarrow v_y(t) = 2\omega \cos \theta \left(-\frac{gt^2}{2} \right)$$

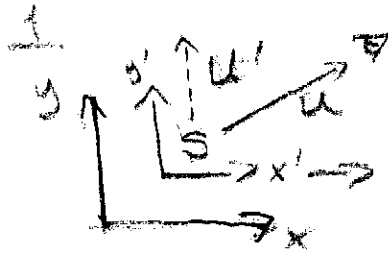
$$\Rightarrow y(t) = -\frac{\omega g t^3}{3} \cos \theta$$

$$\Rightarrow y_0 = -\frac{\omega g}{3} \left(\frac{2d}{g} \right)^{3/2} \cos \theta$$

$$= -\left(\frac{7.3 \times 10^5 / \text{s}}{3} \right) \left(\frac{10 \text{ m}}{\text{s}^2} \right) (6.7 \text{ s})^3 = 5 \times 10^{-3} \text{ m}$$

$$= 5 \text{ mm}$$

Section 3



$$x = \gamma(x' - vt')$$

$$t = \gamma(t' - \frac{vx'}{c^2})$$

$$y = y'$$

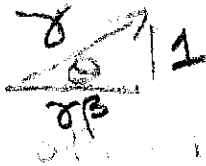
+ or - not
sure
try -

$$u' = \frac{dy'}{dt'} = c$$

$$u = \frac{dx}{dt} \hat{x} + \frac{dy}{dt} \hat{y}$$

$$\left(\frac{\tan \theta}{\text{angle}} \right) = \frac{dy/dt}{dx/dt} = \frac{dy/dt'}{dx/dt'}$$

$$= \frac{c}{\gamma v} = \boxed{\frac{1}{\gamma \beta} = \tan \theta}$$



$$\cot \theta = \gamma \beta$$

$$1 + \cot^2 \theta = \csc^2 \theta$$

$$1 + \beta^2 \gamma^2 = \gamma^2$$

$$\gamma^2 = \csc^2 \theta$$

$$1 + \frac{\beta^2}{1 - \beta^2} = \frac{1}{1 - \beta^2} = \gamma^2$$

$$\boxed{\sin \theta = \frac{1}{\gamma}}$$

3-2 Particle on a sphere

$$\hat{H} = \left(-\frac{\hbar^2}{2m} \nabla^2 + V \right) \psi = E \psi$$

$$\nabla^2 = \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2} = L^2$$

$$\rightarrow -\frac{\hbar^2}{2m} L^2 \psi = E \psi \quad \psi = \delta(r-a) Y_{lm}(\theta, \phi)$$

Eigenfunctions: $L^2 Y(\theta, \phi) = \lambda Y(\theta, \phi)$

$$\lambda = l(l+1)$$

Then, $+\frac{\hbar^2}{2ma^2} (l+1)l \psi = E \psi$

$$\Rightarrow E_l = \frac{+\hbar^2 l(l+1)}{2ma^2} \quad \frac{\text{kg m}^2}{\text{s}^2} \frac{\cancel{\text{kg}} \cancel{\text{m}^2} \cancel{\text{s}^2}}{\cancel{\text{s}^2} \cancel{\text{m}^2} \cancel{\text{kg}}}$$

$$E = \frac{\hbar^2}{2ma^2} \quad E_l = 0E, 2E, 6E, \dots$$

$$P = -\frac{1}{4\pi a^2} \frac{\partial E_l}{\partial a} = -\frac{1}{4\pi a^2} \left(\frac{+\hbar^2 l(l+1)}{2m} \right) \left(-\frac{2}{a^3} \right)$$

$$= \frac{\hbar^2 l(l+1)}{4\pi m a^5} \quad \text{Units } \frac{\hbar^2}{\text{md}^2} \cdot \frac{1}{a^3} \cdot \frac{F \cdot a}{a^3} = \frac{F}{a^2}$$

Stark effect $V = E_0 z = E_0 r \cos \theta = E_0 a \cos \theta$

$$V_{100} = \sqrt{\frac{4\pi}{3}} Y_{10} k$$

$$\Delta E = \int \psi_{lm}^* V \psi_{lm} d\Omega$$

$$= \left(\int \psi_{lm} Y_{10} \psi_{lm} d\Omega \right) \left(\frac{4\pi}{3} \right) E_0$$

$$\propto \begin{pmatrix} l & 1 & l \\ -m & 0 & m \end{pmatrix} \begin{pmatrix} l & 1 & l \\ 0 & 0 & 0 \end{pmatrix}$$

Clebsch-Gordan rules unless

$$-m + 0 + m = 0 \quad \checkmark$$

$l + 1 + l$ is even

3-3

$$Z = \sum_{n=0}^N e^{-n\varepsilon/kT}$$

Sum the series

$$\begin{aligned} Z &= 1 + \sum_{n=1}^N e^{-n\varepsilon/kT} = 1 + e^{-\varepsilon/kT} \sum_{n=0}^{N-1} e^{-n\varepsilon/kT} \\ &= 1 + e^{-\varepsilon/kT} (Z - e^{-N\varepsilon/kT}) = 1 + e^{-\varepsilon/kT} Z - e^{-(N+1)\varepsilon/kT} \end{aligned}$$

$$Z(1 - e^{-\varepsilon/kT}) = 1 - e^{-(N+1)\varepsilon/kT}$$

$$Z = \frac{1 - e^{-(N+1)\varepsilon/kT}}{1 - e^{-\varepsilon/kT}}$$

Internal energy $U = kT^2 \frac{\partial Z}{\partial T}$

$$\langle S \rangle = \frac{U}{\varepsilon} = \frac{kT^2}{\varepsilon} \frac{\partial Z}{\partial T}$$

Closed links = $N - \langle S \rangle$

3-4



$$\Phi = \pi a^2 B = \frac{\pi a^2 B_0 x}{x_0}$$

$$\mathcal{E} = -\frac{1}{c} \frac{\partial \Phi}{\partial t} = \frac{\pi a^2 B_0}{x_0} \frac{dx}{dt} = \frac{\pi a^2 B_0}{x_0} v$$

$$I = \frac{\mathcal{E}}{R} ; \quad P = I\mathcal{E} = \frac{\mathcal{E}^2}{R} = Fv = m\dot{v}v$$

$$\Rightarrow m\dot{v}x = -\left(\frac{\pi a^2 B_0}{x_0}\right)^2 vx$$

$$\dot{v} = -\frac{1}{m} \gamma v$$

$$\gamma = \left(\frac{\pi a^2 B_0}{x_0}\right)^2$$

$$v(t) = \alpha e^{-\lambda t} + \beta \quad \text{and} \quad v(0) = \alpha + \beta = v_0$$

$$v(\infty) = 0 = \beta$$

$$\Rightarrow v(t) = v_0 e^{-\lambda t} ; \quad \dot{v} = -\lambda v_0 e^{-\lambda t}$$

Then,

$$-\lambda v_0 e^{-\lambda t} = -\frac{1}{m} \gamma v_0 e^{-\lambda t} \Rightarrow \lambda = \frac{\gamma}{m}$$

Total distance is then

$$x = \int_0^{\infty} v(t) dt = \left. -\frac{mv_0}{\gamma} e^{-\gamma/m} \right|_0^{\infty} = \frac{mv_0}{\gamma} = mv_0 \left(\frac{x_0}{\pi a^2 B_0}\right)^2$$

4-1 Without GR,

$$\frac{Mv^2}{\cancel{r}} = \frac{GMm}{a^2} \quad \text{and} \quad v = a\dot{\theta}$$

$$\text{Orbit period is } \tau = \frac{2\pi a}{v} = \frac{2\pi a}{a\dot{\theta}} = \frac{2\pi}{\dot{\theta}}$$

$$\rightarrow \dot{\theta} = \frac{2\pi}{\tau}, \quad \text{Then } a^2\dot{\theta}^2 = \frac{4\pi^2 a^2}{\tau^2} = \frac{GM}{a}$$

$$\tau^2 = \frac{4\pi^2 a^3}{GM}$$

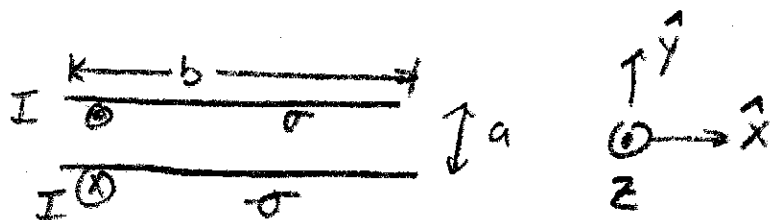
$$\text{With GR} \quad \frac{\cancel{M} a^3 4\pi^2}{\tau_g^2} = \frac{GMm}{(a - GM/c^2)^2}$$

$$\begin{aligned} \tau_g^2 &= \frac{4\pi^2 (a - GM/c^2)^2 a}{GM} \\ &= \frac{4\pi^2 a^3}{GM} \left(1 - \frac{GM}{ac^2}\right)^2 = \tau^2 \left(1 - \frac{2GM}{ac^2}\right) \end{aligned}$$

$$\frac{\tau_g}{\tau} = \left(1 - \frac{2GM}{ac^2}\right)^{1/2} = 1 - \frac{GM}{ac^2}$$

Period is shorter by a factor GM/ac^2

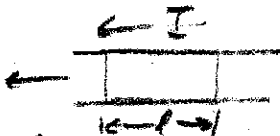
4-2



$$\vec{E} = -4\pi\sigma\hat{y} \rightarrow V = 4\pi\sigma a$$

A length l carries charge σlb , then

$$C = \frac{\sigma lb}{4\pi\sigma a} \rightarrow C_0 = \frac{C}{l} = \frac{b}{4\pi a}$$



$$\vec{B} = \frac{4\pi}{c} \frac{I}{b} \hat{x} \rightarrow \Phi = la \frac{4\pi I}{cb} = \frac{4\pi I la}{cb}$$

$$\mathcal{E} = -L\dot{I} = -\frac{1}{c} \frac{d\Phi}{dt} \Rightarrow L\dot{I} = \frac{1}{c} \frac{4\pi la}{cb} \dot{I}$$

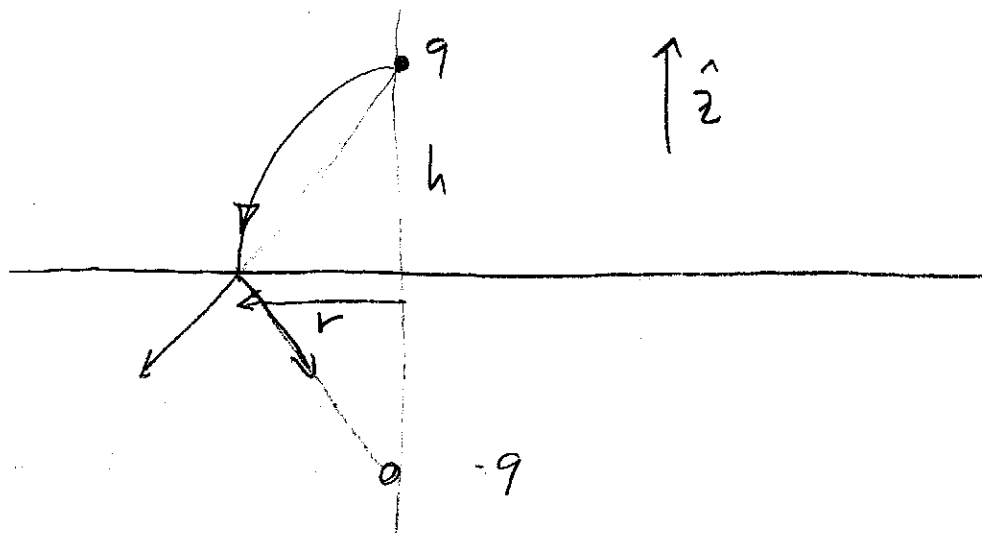
$$L_0 = L/l = \frac{4\pi a}{bc^2}$$

$$Z = \sqrt{\frac{L_0}{C_0}} = \sqrt{\frac{4\pi a}{bc^2} \frac{4\pi a}{b}} = \frac{4\pi a}{c b}$$

$$Z = \frac{12}{3 \times 10^{10}} \frac{0.5}{5} \frac{\text{s}}{\text{cm}} = \frac{1.2 \times 10^{-10} \text{s}}{3} = 0.4 \times 10^{-10} \frac{\text{s}}{\text{cm}}$$

$$= 40 \Omega$$

4-3



Brute force with an image charge
at r , $\vec{E} \perp$ to plate

$$E_z = \left[\frac{-q}{r^2 + h^2} \frac{h}{(r^2 + h^2)^{3/2}} - \frac{q}{r^2 + h^2} \frac{h}{(r^2 + h^2)^{3/2}} \right] \text{ on plate}$$

Surface charge density (w/o image charge)
is then

$$E = 4\pi\sigma \Rightarrow \sigma = \frac{-qh}{2\pi(h^2 + r^2)^{3/2}}$$

Total induced charge is

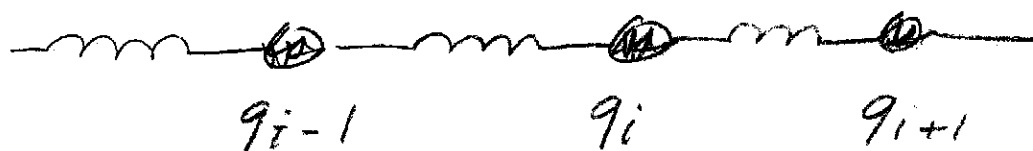
$$Q_{ind} = \int_0^\infty r dr \int_0^{2\pi} d\phi \frac{-qh}{2\pi(h^2 + r^2)^{3/2}}$$

$$= -qh \int_0^\infty \frac{r dr}{(h^2 + r^2)^{3/2}} \quad \begin{aligned} u &= h^2 + r^2 \\ du &= 2r dr \end{aligned}$$

$$q_{ind} = -gh \frac{1}{2} \int_{h^2}^{\infty} \frac{du}{u^{3/2}} = -gh \frac{\cancel{2}}{\cancel{2}} \frac{1}{u^{1/2}} \bigg|_{h^2}^{\infty} = -g$$

There are more elegant ways to do this problem.

4-4



$$m\ddot{q}_i - K(q_{i+1} - q_i) - K(q_i - q_{i-1}) = 0$$

$$m\ddot{q}_i - K(q_{i+1} - 2q_i + q_{i-1}) = 0$$

General solution $q_i = a_i e^{i\omega t}$

Then, $\ddot{q}_i = -\omega^2 a_i e^{i\omega t}$

and $-m\omega^2 a_i - K(a_{i+1} - 2a_i + a_{i-1}) = 0$

Trial solution for coefficients

$$a_j = a e^{i(j\tau - \delta)}$$

$$-m\omega^2 a e^{i(j\tau - \delta)} - K(a e^{i((j+1)\tau - \delta)} - 2a e^{i(j\tau - \delta)} + a e^{i((j-1)\tau - \delta)}) = 0$$

$e^{-i\delta}$ factors out,

$$e^{ij\tau} (-\omega^2 m - K(e^{i\tau} - 2 + e^{-i\tau})) = 0$$

$$-\omega^2 m + 2K - 2K \cos \tau = 0$$

$$m\omega^2 = 2K(1 - \cos \gamma)$$

$$\begin{aligned} \left[\sin^2 x = \frac{1}{4} (e^{2ix} - 2 + e^{-2ix}) = \frac{1}{4} (2 - 2\cos 2x) \right. \\ \left. = \frac{1}{2} (1 - \cos 2x) \right] \end{aligned}$$

$$\begin{aligned} \rightarrow m\omega^2 &= 4K \sin^2 \gamma/2 \\ \omega &= 2 \sqrt{\frac{K}{m}} \sin \gamma/2 \end{aligned}$$

Boundary Condition: $a_0 = 0 = a_N$

$$a_0 = a e^{-i\delta} = 0 \Rightarrow \delta = \pi/2$$

$$\begin{aligned} a_N = 0 = a e^{i(N\gamma - \pi/2)} &\Rightarrow N\gamma - \pi/2 = (2n+1)\pi/2 \\ N\gamma - \pi/2 &= \frac{n\pi}{2} + \pi/2 \\ N\gamma &= \frac{n\pi}{2} + \pi \\ \gamma &= \frac{\pi(n+2)}{N} \end{aligned}$$

$$\Rightarrow \omega = 2 \sqrt{\frac{K}{m}} \sin \frac{\pi(n+2)}{4N}$$

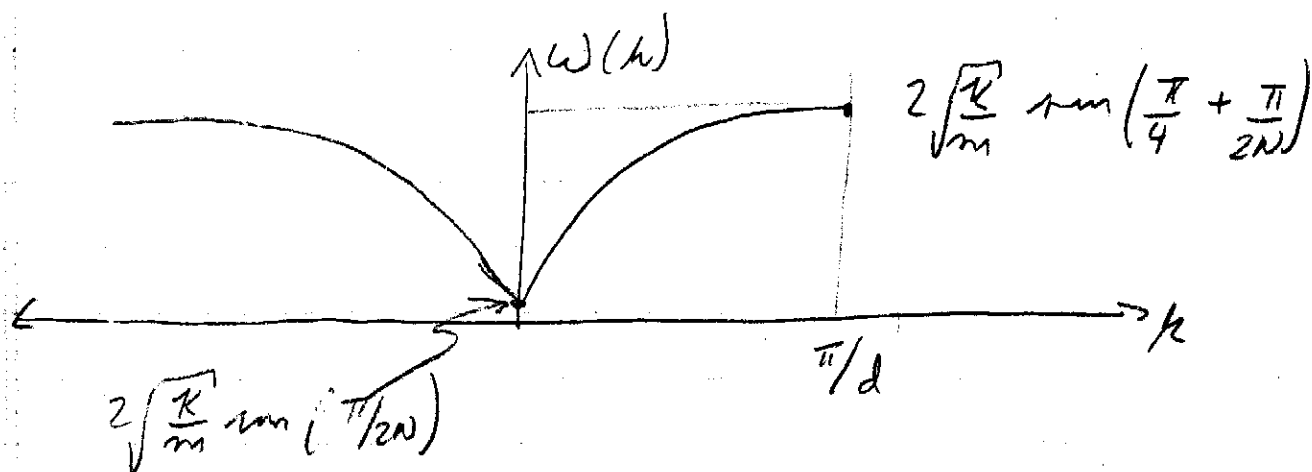
The wavelength $\lambda = 2Nd/n$

$$\Rightarrow k = 2\pi/\lambda = \frac{2\pi}{2Nd} n = \frac{\pi n}{Nd}$$

$$\text{Then } \omega = 2 \sqrt{\frac{K}{m}} \sin \left[\frac{\pi n}{4N} + \frac{\pi}{2N} \right]$$

$$\omega = 2 \sqrt{\frac{\hbar}{m}} \sin \left(\frac{k d}{4} + \frac{\pi}{2N} \right)$$

$$\text{Then, } v = \omega \hbar = \frac{2 N d}{\pi \hbar} \sqrt{\frac{\hbar}{m}} \sin \left(\frac{k d}{4} + \frac{\pi}{2N} \right)$$



5-1
Ans.

$$\beta_2 = v_2/c, \quad \gamma_2 = \frac{1}{\sqrt{1-\beta_2^2}}$$

$$p_2 = m \beta_2 \gamma_2 c$$

$$E_2 = \gamma_2 m c^2$$

F_1 is rest frame of 1,

$$E'_2 = \gamma_1 E_2 + \gamma_1 \beta_1 c p_2$$

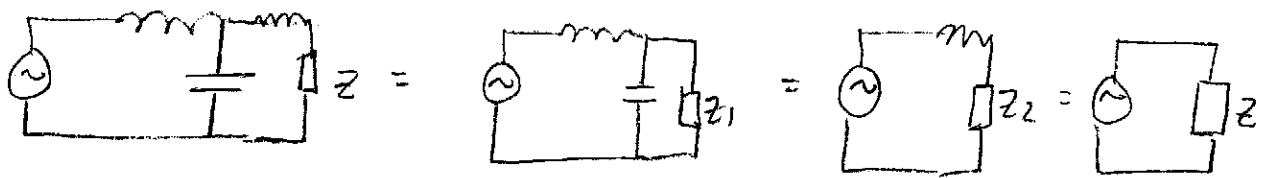
$$= \gamma_1 \gamma_2 m c^2 + \gamma_1 \beta_1 c m \beta_2 \gamma_2 c$$

$$= m c^2 \gamma_1 \gamma_2 (1 + \beta_1 \beta_2)$$

$$= \frac{m c^2 (1 + v_1/c v_2/c)}{\sqrt{(1 - v_1^2/c^2)(1 - v_2^2/c^2)}}$$

$$\beta'_2 = \frac{\beta_1 + \beta_2}{1 - \beta_1 \beta_2}$$

5-2



$$Z_1 = i\omega L + Z$$

$$Z_2 = \left(i\omega C + \frac{1}{Z_1} \right)^{-1} = \left(\frac{i\omega C Z_1 + 1}{Z_1} \right)^{-1} = \frac{Z_1}{1 + i\omega C Z_1}$$

$$Z = i\omega L + Z_2 = i\omega L + \frac{i\omega L + Z}{1 + i\omega C(i\omega L + Z)}$$

$$(Z - i\omega L)(1 - \omega^2 LC + i\omega C Z) = i\omega L + Z$$

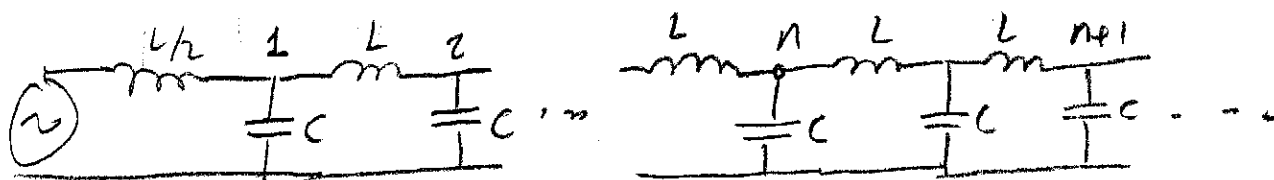
$$i\omega C Z^2 + Z(1 - \omega^2 LC + \omega^2 L^2 C - 1) - 2i\omega L + i\omega^2 L^2 C = 0$$

$$Z^2 = \frac{2iL - i\omega^2 L^2 C}{iC} = \frac{L}{C} (2 - \omega^2 LC)$$

$$Z = \sqrt{\frac{L}{C} (2 - \omega^2 LC)}$$

Wave propagates if $2 - \omega^2 LC > 0 \Rightarrow 2 > \omega^2 LC$
 $\sqrt{\frac{2}{LC}} > \omega$

For a long chain, treat n , the step number, as a continuous variable



Then, Q_n is the charge on capacitor n and

$$\dot{Q}_n = -I_{n+1} + I_n \rightarrow \frac{\partial I(n)}{\partial n} = \frac{\partial Q(n)}{\partial t} = \frac{1}{C} \frac{\partial V}{\partial t} \quad (1)$$

$$V_{n+1} - V_n = \frac{L I_{n+1} + L I_n}{2} \rightarrow \frac{\partial V}{\partial n} = L \frac{\partial I(n)}{\partial t} \quad (2)$$

Then

$$\left. \begin{aligned} \frac{1}{C} \frac{\partial^2 V}{\partial t^2} &= \frac{\partial^2 I}{\partial t \partial n} \quad (1) \\ \frac{\partial^2 V}{\partial n^2} &= L \frac{\partial^2 I}{\partial t \partial n} \quad (2) \end{aligned} \right\} \frac{1}{LC} \frac{\partial^2 V}{\partial t^2} - \frac{\partial^2 V}{\partial n^2} = 0$$

has solution $V = f(n \pm ct) = f(u)$, $u = n \pm ct$

$I = g(u)$ and

$$\frac{1}{LC} c^2 f'' - f'' = 0 \Rightarrow c^2 = LC \Rightarrow c = \sqrt{LC}$$

and $g = f$. Can also do problem discretely by using a series.

$$5-3 \ a) \quad i^i = (e^{i\pi/2})^i = e^{-\pi/2}$$

$$\frac{3.14}{2} = 1.57$$

$$b) \quad \int_{-\infty}^{\infty} \delta(f(x)) dx = \int_{-\infty}^{u_0} \delta(u) \frac{du}{f'(x_0)} = \frac{1}{f'(x_0)} = \frac{2\sqrt{b}}{c}$$

$$\text{Let } f(x) = u, \quad du = \frac{d}{dx} dx$$

$$\text{If } f(x) = a + \sqrt{b+cx} \quad - \quad 0 = a + \sqrt{b+cx}$$

$$\frac{d}{dx} : \frac{1}{2} \frac{1}{\sqrt{b+cx}} c = \frac{c}{2\sqrt{b+cx}} \quad a^2 = b+cx$$

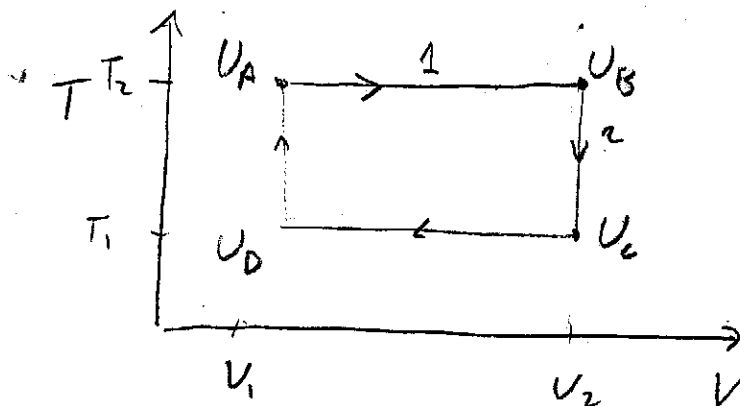
$$f'(x_0) = \frac{c}{2\sqrt{b+a^2-b}} \quad \frac{a^2-b}{c} = x_0$$

$$= \frac{2|a|}{c}$$

5-4 Cycle

1. Expand gas from V_1 to V_2 at T_2
2. Cool at V_2 from T_2 to T_1
3. Compress gas from V_2 to V_1 at T_1
4. Warm gas from T_1 to T_2 at V_1

$$du = k_B T ds - p dV$$



During 1, $Q_1 - W_1 = U_B - U_A$

2, $dw=0$ $du = U_C - U_B = Q_2 = -Q_1$

3, $Q_3 - W_3 = U_D - U_C$

4, $dw=0$ $du = U_A - U_D = Q_4$

The total work out is $W_1 + W_3 = W$

$Q_1 = -Q_3$ for an ideal gas

$$W_1 + W_3 = Q_1 + U_A - U_B + Q_3 + U_C - U_D$$

$$= (Q_1 + Q_3) + Q_4 + Q_2$$

$$\left. \begin{aligned} Q_1 &= \int k T_2 dS = k T_2 4S \\ Q_3 &= -k T_1 4S \end{aligned} \right\} \frac{Q_1}{T_2} = -\frac{Q_3}{T_1}$$

$$W = Q_1 (1 - T_1/T_2)$$

Total heat taken in is $Q_1 + Q_3 = Q_1 + \lambda Q_1$

$$\Rightarrow \eta = \frac{(1 - T_1/T_2)}{(1 + \lambda)}$$