

Following the promising results of the α_T jet-balancing method previously described for the all-hadronic SUSY searches, it is a natural progression to look for extensions of this approach to the single-electron search, where a significant presence of QCD multi-jet backgrounds is expected.

The α_T variable is now redefined as an N-object system where the set of objects is 1 electron and N-1 jets. The new definition reproduces the kinematics of a di-jet system by construction two pseudo-jets, which balance one another in H_T . Pseudo-jets are formed from the combination of the objects by minimising the variable $\Delta H_T = |H_{T,1} - H_{T,2}|$, where the definition of the α_T variable is now written as:

$$\alpha_T = \frac{1}{2} \frac{H_T - \Delta H_T}{M_T} = \frac{1}{2} \frac{H_T - \Delta H_T}{\sqrt{H_T^2 - M H_T^2}} \quad (1)$$

The basic selection for single electron events in the leptonic α_T approach has the following requirements:

- Exactly one electron passing requirements as follows:
 - $p_T > 20$ GeV
 - $|\eta| < 2.4$
 - Passes the Cut Based ID formed by simple, yet robust, variables (these are the H/E , the super-cluster (SC) - track matching variables $\Delta\phi, \Delta\eta$, and shower shape variable $\sigma_{i\eta i\eta}$, Tracker Isolation, ECAL Isolation and HCAL Isolation¹).
- The event is vetoed if there are any muons passing the following requirements:
 - $p_T > 15$ GeV
 - $|\eta| < 2.1$
 - Passes ID requirement: IsGlobalMuon
- A jet is rejected if it is found close to a tight and isolated electron within $\Delta R=0.5$. The jet selection is as follows:
 - $p_T > 20$ GeV
 - $|\eta| < 5$
 - EMF < 0.9

The analysis uses data samples for the QCD multi-jet background processes produced with the full simulation production for Physics at 7TeV with CMS. The Monte Carlo datasets used are detailed in Table 1. The available luminosity is shown, although the plots are normalised to 1pb^{-1} . We use the EMenriched and BCtoE samples which are designed to be enriched in electrons, in order to enhance statistics. In addition, as these samples only cover the region $20 < \hat{p}_T < 170$ GeV/c, we include the inclusive QCD Jets sample for the region $\hat{p}_T > 170$ GeV/c.

Data Set	N events	σ (pb)	Equivalent luminosity (pb^{-1})
QCD BCtoE [$20 < \hat{p}_T < 30$]	432380	108330	3.99
QCD BCtoE [$30 < \hat{p}_T < 80$]	840100	138762	6.05
QCD BCtoE [$80 < \hat{p}_T < 170$]	682720	9422.4	72.4
QCD EM Enriched [$20 < \hat{p}_T < 30$]	6169999	1719150	3.59
QCD EM Enriched [$30 < \hat{p}_T < 80$]	9054696	3498700	2.59
QCD EM Enriched [$80 < \hat{p}_T < 170$]	2492814	134088	18.59
QCDJets $\hat{p}_T > 170$	3132800	25470	122.99

Table 1: *The Monte Carlo datasets used to investigate the Delta ID Inversion method in QCD backgrounds. The available Lumniosity is shown, although plots produced are normalised to 1pb^{-1} for the purpose of understanding near-reach of CMS.*

¹The cut based ID selection cuts are taken such so that they correspond to an 80% efficiency in the $W \rightarrow e\nu$ analysis.

This session is dedicated to a first approach of commissioning the α_T observable and study its behavior in pure fake electron events. It is therefore desirable to collect a suitable control sample which will be dominated by fake electrons whereas eliminate sources of prompt electron events (like W events). One way to obtain such a sample is using the anti-selection method on electron ID variables which are less correlated with the missing transverse energy. In this section, we investigate the possibility of inverting the $\Delta\eta(\text{trk-SC})$ and $\Delta\phi(\text{trk-SC})$ id cuts in the electron selection. The anti-selected events in this method are collected with electrons that pass all selection id criteria except the $\Delta\phi(\text{trk-SC})$ and $\Delta\eta(\text{trk-SC})$ ones.

In order to establish the anti-selection method above, we compare next the performance of the leptonic α_T as obtained from the control sample and the actual QCD events passing the electron criteria defined in the “signal” region.

As mentioned previously, SUSY events are expected to have high H_T . Therefore it is critical to understand the behaviour in leptonic α_T as a function of H_T . The plots in 1 show the normalised shape of distributions, both before H_T cut and the evolution with H_T . The selected and anti-selected distributions agree very well.

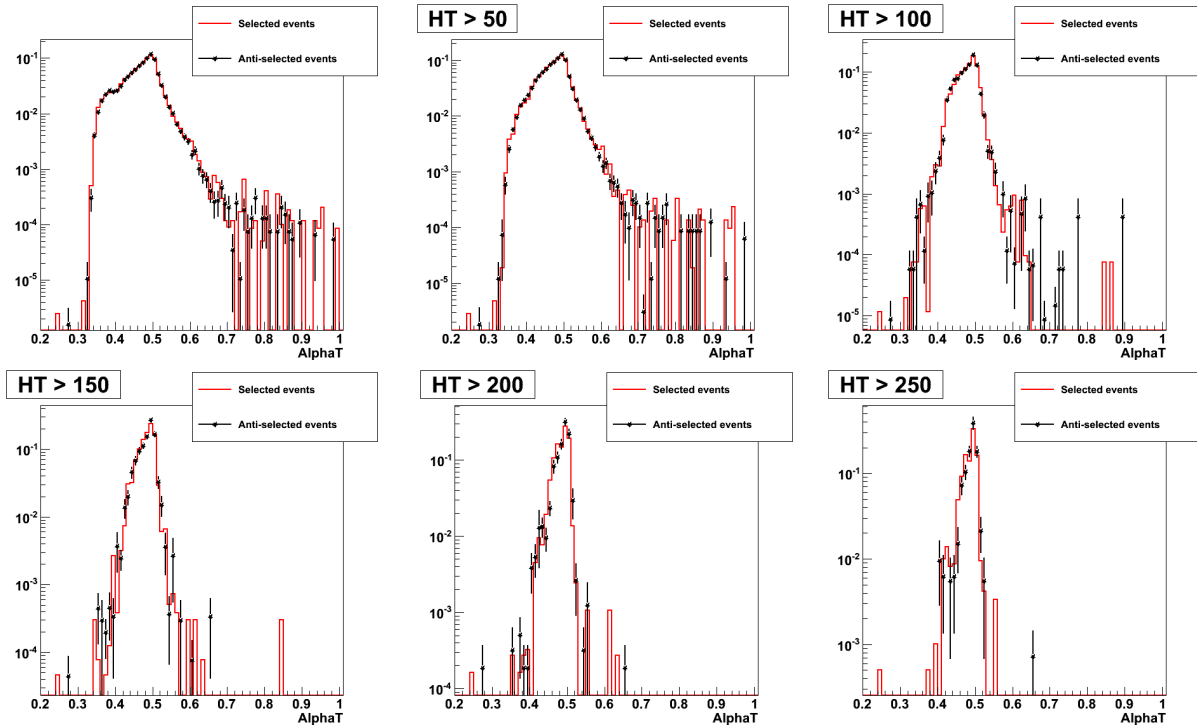


Figure 1: The α_T distributions for selected (red) and anti-selected events (blue) from inversion of the $\Delta\phi$ and $\Delta\eta$ ID Cutts, shown without H_T cut (Top Left) and with progressive H_T cuts (left-right, top-bottom). These distributions are normalised to unity for shape comparison. There is good agreement between the selected and anti-selected regardless of H_T requirement, and the high α_T tails reduce as expected when moving to higher H_T cuts.

In order to demonstrate the power of H_T in α_T tail-reduction, we introduce the variable R_{α_T} which is defined as the ratio of the number of events passing the α_T cut over the number of events failing it:

$$R_{\alpha_T} = \frac{N(\alpha_T > 0.55)}{N(\alpha_T < 0.55)} \quad (2)$$

The “default” cut value here is the value prompted from the all-hadronic analysis, 0.55. Figure 2 shows a plot of R_{α_T} as a function of the H_T cut applied. As the H_T requirement increases, R_{α_T} decreases in an exponential manner. The selected and anti-selected events from the electron id inversion method are in good agreement.

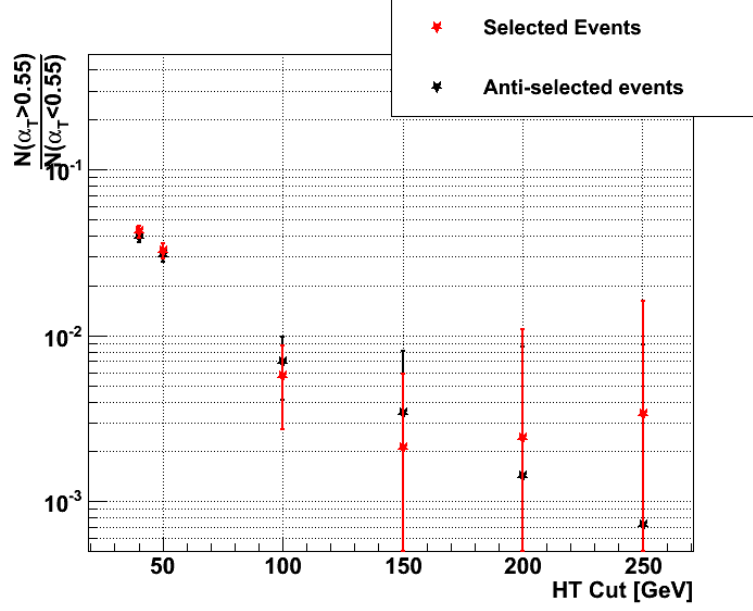


Figure 2: The R_{α_T} versus the H_T cut applied for the QCD multi-jet background, shown for both selected and anti-selected events in the Delta ID Inversion method.

In order to verify that the tail is not reduced proportionally to the peak, an additional set of ratios may be defined. Defining a reference istribution for low H_T region $[0,100]$, we can investigate distributions in higher HT bins ($[100,150]$, $[150,200]$, $[200,250]$) by dividing them by the reference distribution. These plots are shown in Figure 3. As the region of H_T is increased, the difference in ratio between the peak and the tail becomes more pronounced.

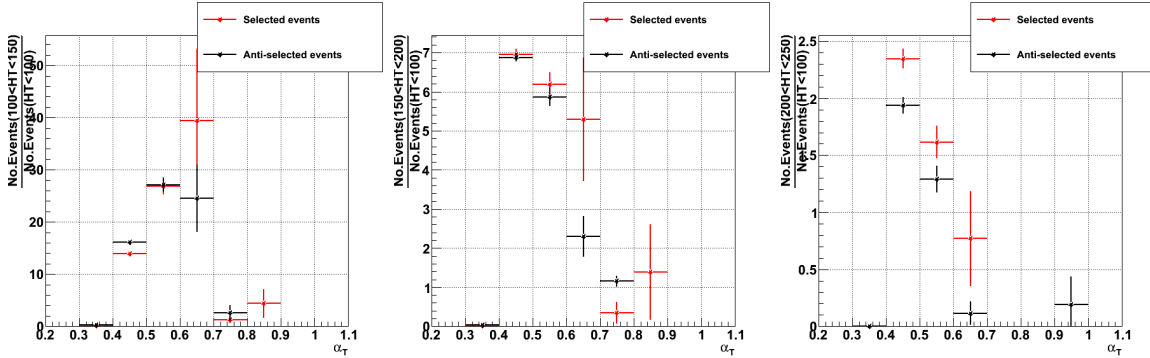


Figure 3: The ratio of α_T distributions in high H_T regions (Left: $[100,150]$ GeV, Middle: $[150,200]$ GeV, Right: $[200,250]$ GeV) to the α_T distribution in a low H_T region ($[0,100]$ GeV). Both selected (red) and anti-selected events (blue) from inversion of the $\Delta\phi$ and $\Delta\eta$ ID Cuts are shown.