

chap-1 Energy and Potential (contd)

Date 17-Aug-19
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- (*) energy utilized in moving a point charge in an electric field

$$\vec{E} = \vec{F} \quad \text{[Point charge is moving in the electric field]}$$

Q

[Electric potential is (-)]

$$\vec{F}_E = Q \vec{E}$$

↳ force from the field

$$(\text{component}) \quad F_{EL} = \vec{F} \cdot \vec{dl}$$

= $Q \vec{E} \cdot \vec{dl}$

$$F_{app} = -Q \vec{E} \cdot \vec{q}$$

Usage of energy = $F \times \text{displacement}$
small work done

$$= -Q \vec{E} \cdot \vec{dl} \cdot dl$$

$$dW = -Q \vec{E} \cdot d\vec{l}$$

$$W = -Q \int_{\text{initial}}^{\text{final}} \vec{E} \cdot d\vec{l}$$

Example An electric field is given by $\vec{E} = 6y^2 \hat{z} + 12xy \hat{z} + 6xy^2 \hat{z}$ v/m.
An incremental path is represented by $6xy^2 \hat{z}$ v/m.
 $d\vec{l} = -3x \hat{i} + 5y \hat{j} - 2z \hat{k}$ micrometer. Find the deft work done in moving a $2 \mu C$ charge along this path if the location of charge is at

(i) P_A(0, 2, 5) (ii) P_B(1, 1, 1)

SOL.

$$\vec{E} \cdot d\vec{l} = -18y^2 z + 60xyz - 12xy^2 z$$

$$W = -Q \int_{(0,2,5)}^{(1,1,1)} (-18y^2 z + 60xyz - 12xy^2 z)$$

$$= 8 \times 10^{-6} \times 10^{-6} \times (-18 \times 4 \times 5)$$

$$= -720 \text{ pJ}$$

(11)

$$W = -Q \int_{(0,0,0)}^{(1,1,1)} (-18y^2 z + 60xyz - 12xy^2 z)$$

$$= -2 \times 10^{-12} [-18 \times 1 \times 1 + 60 \times 1 - 12 \times 1]$$

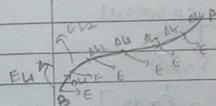
$$= -60 \text{ pJ Ans.}$$

line integral &

$$W = -Q \int_{\text{initial}}^{\text{final}} \vec{E} \cdot d\vec{l}$$

without using vector analysis

$$W = -Q \int_{\text{initial}}^{\text{final}} E_L dL$$



$$W = -Q [E_1 dL_1 + E_2 dL_2 + E_3 dL_3 + E_4 dL_4 + E_5 dL_5 + E_6 dL_6]$$

OR

$$= -Q [E_1 \cdot d\vec{l}_1 + E_2 \cdot d\vec{l}_2 + \dots + E_6 \cdot d\vec{l}_6]$$

$$E_1 = E_2 = E_3 = \dots = \vec{E}$$

$$= -Q \vec{E} (\vec{d}\vec{l}_1 + \vec{d}\vec{l}_2 + \dots + \vec{d}\vec{l}_6)$$

$$= -Q \vec{E} \cdot \vec{d}\vec{l}_{BA}$$

$$= -Q \int_A^B \vec{E} \cdot d\vec{l}$$

$$W = -Q \vec{E} \int_A^B d\vec{l}$$

- Work does not depend on path.

Example: A Non Uniform field is given by $\vec{E} = y\vec{a}_x + 2\vec{a}_y + 2\vec{a}_z$
 determine the work done in carrying 2 C charge from
 $B(1,0,1)$ to $A(0.8, 0.6, 1)$ along the shorter arc of
 the circle $x^2 + y^2 = 1$ & $z = 1$.

Soln

$$W = -q \int_A^B \vec{E} \cdot d\vec{l}$$

B

$$d\vec{l} = dx \vec{a}_x + dy \vec{a}_y + dz \vec{a}_z$$

$$q = 2 \text{ C}$$

$$\vec{E} = y \vec{a}_x + 2 \vec{a}_y + 2 \vec{a}_z$$

(0.8, 0.6, 1)

$$\text{So, } W = -q \int_{(1,0,1)}^{(0.8, 0.6, 1)} y dx + x dy + dz$$

(1,0,1)

$$= -q \left[\frac{y^2}{2} \right]$$

$$= -q \left[\int_1^{0.8} xy + xy + dz \right]_{(1,0,1)}^{(0.8, 0.6, 1)}$$

$$= (-2) \left[\int_1^{0.8} y dx + \int_0^{0.6} x dy + \int_1^{0.8} dz \right]$$

$$= (-2) \left[\int_1^{0.8} \sqrt{1-x^2} dx + \int_0^{0.6} \sqrt{1-y^2} dy + 0 \right]$$

$$= (-2) \left[\left(\frac{\pi}{2} \sqrt{1-x^2} + \frac{1}{2} \sin^{-1} x \right) \Big|_1^{0.8} + \left(\frac{y}{2} \sqrt{1-y^2} + \frac{1}{2} \sin^{-1} y \right) \Big|_0^{0.6} \right]$$

$$= (-2) \left[\frac{\pi}{2} \sqrt{1-0.8^2} + 0.54 \right]$$

$$= (-1) \left[0.2 \sqrt{1-0.8^2} + \sin^{-1}(0.8) - 1 \times 0 - \sin^{-1}(1) \right]$$

$$+ 0.6 \sqrt{1-0.6^2} + \frac{1}{2} \sin^{-1}(0.6) - \sin^{-1}(0)$$

$$= -0.96 \text{ J}$$

IIo

$$y = mx + c$$

$$y - y_B = \frac{y_A - y_B}{x_A - x_B} (x - x_B)$$

$$y = \frac{0.6 - 0}{0.8 - 1} (x - 1)$$

$$y = -3(x - 1)$$

$$x = \frac{-y + 1}{3};$$

$$\text{So, } W = (-2) \left[\int_1^{0.8} -3(x-1) dx + \int_0^{0.6} \frac{-y+1}{3} dy \right]$$

$$= (-2) \left[\left(-3x^2 + 3x \right) \Big|_1^{0.8} + \left(\frac{-y^2 + y}{6} \right) \Big|_0^{0.6} \right]$$

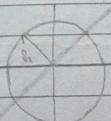
$$= (-2) \left[\frac{-0.06}{(-12.48) + 0.54} \right]$$

$$= -0.96 \text{ J}$$

19-Aug.

$$\text{Infinite line charge } \vec{E} = E_0 \vec{a}_\theta = \frac{\rho_L}{2\pi\epsilon_0 r} \vec{a}_\theta$$

(i)



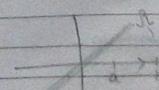
Moving along the Circle

$$dL = r d\theta \vec{a}_\theta$$

$$W = -q \int_{\text{ini}}^{\text{fin}} \vec{E} \cdot d\vec{l}$$

| $W=0$

(ii)



moving along radial path a to b

$$W = -\Phi \int_{\text{ini}}^{\text{final}} \frac{d\vec{r}}{2\pi\epsilon_0} \cdot d\vec{p} \cdot d\vec{p}$$

$$= -\Phi \int_{\text{ini}}^{\text{final}} \frac{d\vec{r}}{2\pi\epsilon_0} \cdot d\vec{p}$$

$$W = -\Phi \frac{qL}{2\pi\epsilon_0} \ln\left(\frac{b}{a}\right)$$

Potential & Potential Difference:

Potential difference is defined as workdone to move a +ve point charge from one point to another in direction of E.

$$W = -\Phi \int_{\text{ini}}^{\text{final}} \vec{E} \cdot d\vec{l}, V = W/\Phi$$

$$\text{pot. diff' } V = - \int_{\text{ini}}^{\text{final}} \vec{E} \cdot d\vec{l} [V \text{ or J/C}]$$

$$V_{AB} = - \int_B^A \vec{E} \cdot d\vec{l}$$

$$\text{let } W = \Phi \frac{qL}{2\pi\epsilon_0} \ln\left(\frac{b}{a}\right)$$

$$V_{AB} = W = \frac{\Phi L}{2\pi\epsilon_0} \ln\left(\frac{b}{a}\right)$$

$$\vec{E} = E_r \hat{r} = \frac{\Phi}{4\pi\epsilon_0 r^2} \hat{r}$$

$$dl = dr \hat{r}$$

$$V_{AB} = - \int_B^A \frac{\Phi}{4\pi\epsilon_0 r^2} \hat{r} \cdot dr \cdot \hat{r}$$

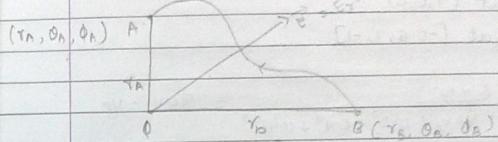
$$= - \int_B^A \frac{\Phi}{4\pi\epsilon_0 r^2} \cdot \frac{dr}{r^2}$$

$$V_{AB} = \frac{\Phi}{4\pi\epsilon_0} \left(\frac{1}{r_A} - \frac{1}{r_B} \right) = V_A - V_B$$

The potential field of a point charge

$$V_{AB} = \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{r_A} - \frac{1}{r_B} \right)$$

$$= V_A - V_B$$



$$d\vec{r} = dr \hat{r} + r d\theta \hat{\theta} + r \sin\theta d\phi \hat{\phi}$$

$$V_{AB} = - \int_{r_A}^{r_B} E_r dr = - \int_{r_A}^{r_B} \frac{\Phi}{4\pi\epsilon_0 r^2} dr$$

$$= \frac{\Phi}{4\pi\epsilon_0} \left(\frac{1}{r_A} - \frac{1}{r_B} \right)$$

: potential difference is independent of path but dependent on distance of charge to that point.

Let V=0 at ∞ point

$\Rightarrow g_1 = g_{\infty}$ is at ∞

$$\therefore V = \frac{\Phi}{4\pi\epsilon_0 r} = \frac{\Phi}{4\pi\epsilon_0 r}$$

$\therefore \Phi/4\pi\epsilon_0 r$ J of work is to be done to bring 1C charge from ∞ to a point at distance r from Q charge.

\Rightarrow Equipotential surfaces

potential difference between two points in a surface = 0 ; Hence $W=0$

Eg: A point charge of q is located at the origin in free space. Find V_p if point P is located at $[0.2, -0.4, 0.4]$ and

- (i) $V=0$ at ∞
- (ii) $V=0$ at $(1,0,0)$
- (iii) $V=90V$ at $[0.5, 1, -1]$

Sol:

$$(i) V_{p\infty} = \frac{q}{4\pi\epsilon_0 (0.2^2 + 0.4^2 + 0.4^2)^{1/2}} = 90V = V_p$$

$$(ii) V_{p0} = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{r_p} - \frac{1}{r_0} \right) = 6 \times 9 \left(\frac{5}{3} - 1 \right) = 36V = V_p$$

$$V_{p0} = V_p - V_{Q0}$$

$$(iii) V_{pQ} = 6 \times 9 \left(\frac{5}{3} - \frac{2}{3} \right) = 54V$$

$$V_{pQ} = V_p - V_Q$$

$$54 = V_p - 90$$

$$V_p = 74V$$

\Rightarrow Potential field of a system of charges - Conservative property

$$V(\vec{r}) = \frac{q_1}{4\pi\epsilon_0 |\vec{r}-\vec{r}_1|}$$

$$V(\vec{r}) = \frac{q_1}{4\pi\epsilon_0 |\vec{r}-\vec{r}_1|} + \frac{q_2}{4\pi\epsilon_0 |\vec{r}-\vec{r}_2|}$$

n-point charges

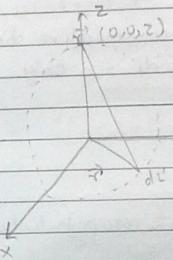
$$V(\vec{r}) = \frac{q_1}{4\pi\epsilon_0 |\vec{r}-\vec{r}_1|} + \frac{q_2}{4\pi\epsilon_0 |\vec{r}-\vec{r}_2|} + \dots + \frac{q_n}{4\pi\epsilon_0 |\vec{r}-\vec{r}_n|}$$

$$= \sum_{m=1}^n \frac{q_m}{4\pi\epsilon_0 |\vec{r}-\vec{r}_m|}$$

$$V(\vec{r}) = \frac{\delta V(\vec{r})}{4\pi\epsilon_0 |\vec{r}-\vec{r}_1|} + \frac{\delta V(\vec{r})}{4\pi\epsilon_0 |\vec{r}-\vec{r}_2|} + \dots$$

$$= \int_{\text{Vol}} \frac{\delta V(\vec{r}') dV}{4\pi\epsilon_0 |\vec{r}-\vec{r}'|}$$

\Rightarrow potential of a uniform line charge (in ring shape)



$$|\vec{r} - \vec{r}'| = \sqrt{a^2 + z^2}$$

$$dL' = a d\phi'$$

$$\theta = a; z = 0$$

$$V(\vec{r}) = \int \frac{[p_L(\vec{r}')] dL'}{4\pi\epsilon_0 |\vec{r} - \vec{r}'|}$$

$$dL' = a d\phi'; \vec{r} = z \vec{a}_z; |\vec{r}'| = a \vec{a}_\theta$$

$$|\vec{r} - \vec{r}'| = \sqrt{a^2 + z^2}$$

$$V(\vec{r}) = \int_0^{2\pi} \frac{p_L \cdot a \cdot d\phi'}{4\pi\epsilon_0 (a^2 + z^2)^{1/2}}$$

$$= \frac{p_L \cdot a}{4\pi\epsilon_0 (a^2 + z^2)^{1/2}} \cdot (2\pi)$$

$$V(\vec{r}) = \frac{p_L \cdot a}{2\pi (a^2 + z^2)^{1/2}}$$

- any field which satisfies $\oint \vec{E} \cdot d\vec{l} = 0$ is a conservative field as work done is 0

Eg: Gravitational field

Eg:- $E = \sin(\pi r) \vec{a}_\theta$

We are assuming circular path of $P = P$

$$d\vec{l} = \int d\phi \cdot \vec{a}_\theta$$

$$W = \oint \vec{E} \cdot d\vec{l}$$

$$= \Phi \int_0^{2\pi} \sin(\pi r) \cdot \vec{a}_\theta \cdot d\phi$$

$$= \Omega 2\pi P \sin \theta$$

$$\text{at } \theta = 1; \omega = 2\pi P \sin \theta$$

$$\omega = \oint \vec{E} \cdot d\vec{l} = \begin{cases} 0; & S = 0, 24 \text{ rad} \\ 1; & S = 112, 144 \end{cases}$$

above is example of non-conservative field.
in non-conservative field; $\oint \vec{E} \cdot d\vec{l} = 0$ only
possible at same points.

Potential Gradient:

rate of change of spacing between equipotential surfaces is called potential gradient.

Gradient: rate of change of one thing wrt to other.

$$V = - \oint \vec{E} \cdot d\vec{l}$$

$$d\vec{l}, \vec{E}$$

$$\Delta V = - \vec{E} \cdot \vec{dL} \quad \vec{E} \rightarrow \text{constant}$$

$$d\vec{l} = dL \hat{a}_l$$

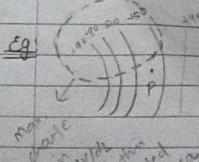
$$\Delta V = - E dL \cos \theta$$

$$\boxed{\Delta V = - E \cos \theta = \frac{dV}{dL}}$$

$\left| \frac{dV}{dL} \right|_{\max}$; when $\cos \theta = -1$; it will may possible when dV & dL are in opp direction to each other.

$$\left| \frac{dV}{dL} \right|_{\max} = E$$

$\vec{E} = - \left| \frac{dV}{dL} \right|_{\max} \hat{a}_n$ direction of E , normal to the equipotential surface



$$\text{spherical: } \frac{\partial V}{\partial r} \hat{a}_r + \frac{1}{r} \frac{\partial V}{\partial \theta} \hat{a}_{\theta} + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \hat{a}_{\phi}$$

$$\text{cylindrical: } \frac{\partial V}{\partial r} \hat{a}_r + \frac{1}{r} \frac{\partial V}{\partial \theta} \hat{a}_{\theta} + \frac{\partial V}{\partial z} \hat{a}_z$$

$$\left| \frac{dV}{dL} \right|_{\max} = \frac{dV}{dN}$$

$$\therefore \vec{E} = - \frac{dV}{dN} \hat{a}_N = - \text{grad } V$$

$$\text{Gradient of } T = \text{grad } T = \frac{dT}{dN} \hat{a}_N$$

$$\vec{E} = - \text{grad } V$$

$$dV = \frac{\partial V}{\partial x} dx + \frac{\partial V}{\partial y} dy + \frac{\partial V}{\partial z} dz \quad (1)$$

$$dV = - \vec{E} \cdot d\vec{L} = - E_x dx - E_y dy - E_z dz \quad (2)$$

Comparing (1) & (2)

$$E_x = - \frac{\partial V}{\partial x}; \quad E_y = - \frac{\partial V}{\partial y}; \quad E_z = - \frac{\partial V}{\partial z}$$

$$\text{so } \vec{E} = - \left(\frac{\partial V}{\partial x} \hat{a}_x + \frac{\partial V}{\partial y} \hat{a}_y + \frac{\partial V}{\partial z} \hat{a}_z \right)$$

$$\text{grad } V = \frac{\partial V}{\partial x} \hat{a}_x + \frac{\partial V}{\partial y} \hat{a}_y + \frac{\partial V}{\partial z} \hat{a}_z$$

$$\boxed{\nabla = \frac{\partial}{\partial x} \hat{a}_x + \frac{\partial}{\partial y} \hat{a}_y + \frac{\partial}{\partial z} \hat{a}_z}$$

Eg: given the potential field $V = 2x^2y - 5z$ and a point $P(-4, 3, 6)$. find the potential V at Point P , the electric field intensity E , the dir^c of E , electric flux density D and the volume charge density at point P .

$$V_{-4,3,6} = 2(-4)^2(3) - 5(6) = 2 \times 16 \times 3 - 30 = 66 \text{ V}$$

$$\vec{E} = - \text{grad } V = - 4xy \hat{a}_x + 2x^2 \hat{a}_y + 5 \hat{a}_z$$

$$E_{\perp} = +48 \hat{a}_x - 32 \hat{a}_y + 5 \hat{a}_z \text{ V/m}$$

$$\vec{E} = 4.8 \vec{a}_x + 3.2 \vec{a}_y + 5 \vec{a}_z$$

$$= 5.790 \text{ V/m}$$

$$= 0.82 \vec{a}_x + 0.55 \vec{a}_y + 0.86 \vec{a}_z \text{ V/m}$$

$$\vec{D} = \epsilon_0 \vec{E} = \epsilon_0 (-4.8 \vec{a}_x + 2 \vec{a}_y + 5 \vec{a}_z)$$

$$= \epsilon_0 (4.8, -3.2, 5) = 35.4 \text{ nC/m}^2 \vec{a}_x - 17.71 \text{ nC/m}^2 \vec{a}_y + 44.3 \text{ nC/m}^2 \vec{a}_z$$

Volume charge density $\rho_v = \nabla \cdot \vec{D}$

$$= \epsilon_0 [-4y + 0 + 0]$$

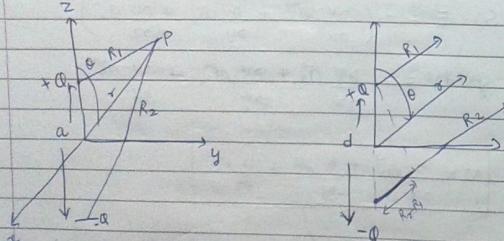
$$= -4\epsilon_0$$

$$= -4 \times 3 \times 9.75 \times 10^{-12}$$

$$= 35.4 \text{ nC}$$

$$= 106.2 \text{ PC}$$

Dipole: two opposite charges are placed at certain distance from each other.



$$V = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$$

$$= \frac{q}{4\pi\epsilon_0} \frac{d \cos\alpha}{R_1 R_2}$$

Here $R_2 - R_1 = d \cos\alpha$
 $R_1 = R_2 = r$

$$V = \frac{q \cdot d \cos\alpha}{4\pi\epsilon_0 r^2}$$

$$\vec{E} = -\nabla V$$

$$= -\left(\frac{\partial V}{\partial x} \vec{a}_x + \frac{\partial V}{\partial y} \vec{a}_y + \frac{\partial V}{\partial z} \vec{a}_z \right)$$

$$= -\frac{qd \cos\alpha}{4\pi\epsilon_0 r^3} \vec{a}_x + \frac{qd \sin\alpha}{4\pi\epsilon_0 r^3} \vec{a}_y + \frac{1}{4\pi\epsilon_0 r^3} \vec{a}_z$$

$$\vec{E} = \frac{qd \cos\alpha}{4\pi\epsilon_0 r^3} \vec{a}_x + \frac{qd \sin\alpha}{4\pi\epsilon_0 r^3} \vec{a}_y$$

$$\boxed{\vec{E} = \frac{qd}{4\pi\epsilon_0 r^3} (2 \cos\alpha \vec{a}_x + \sin\alpha \vec{a}_y)}$$

∴ we can say that $V \propto \frac{1}{r^2}$ & $E \propto \frac{1}{r^3}$

29-Aug

Dipole moment:

$$\vec{P} = q \vec{d} \text{ cm}$$

$$\vec{d} \cdot \vec{a}_r = d \cos\alpha$$

$$V = \frac{\vec{P} \cdot \vec{a}_r}{4\pi\epsilon_0 r^2}$$

$$V = \frac{1}{4\pi\epsilon_0} \frac{\vec{P} \cdot (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^2} \rightarrow$$

→ determine the dipole centered

Ex: Two dipoles with dipole moment $-5 \vec{a}_z \text{ nCm}$ & $9 \vec{a}_z \text{ nCm}$, one located at points $(0, 0, -2)$ & $(0, 0, 3)$ respectively. Find the potential at the origin.

Sol:

$$V_1 = \frac{p_1 \alpha_1}{4\pi\epsilon_0 (r_1)^2}$$

$$= \frac{-5 \vec{a}_z \cdot 2 \vec{a}_z}{4\pi\epsilon_0 (2)^2} = -10 \quad (\alpha_1 = 2)$$

$$V_2 = \frac{9 \vec{a}_z \cdot (-3 \vec{a}_z)}{4\pi\epsilon_0 (3)^2} = \frac{27}{4\pi\epsilon_0 (3)^2} \quad (\alpha_2 = (0, 0, 0) - (0, 0, 3) = -3 \vec{a}_z) \quad \alpha_2 = 3$$

$$V = -25 \text{ V}$$

$$V = \frac{1}{4\pi\epsilon_0} \left[\frac{\vec{P}_1 \cdot \vec{r}_1}{|\vec{r}_1|^3} + \frac{\vec{P}_2 \cdot \vec{r}_2}{|\vec{r}_2|^3} \right]$$

Energy density in the electrostatic field:

$$\text{Work done to position } Q_2 = Q_2 \cdot V_{2,1}$$

$$\text{Work done to position } Q_3 = Q_3 \cdot V_{3,1} + Q_3 \cdot V_{3,2}$$

(Potential at Q_2
because of Q_3)

Total Work

$$W_E = Q_2 V_{2,1} + Q_3 (V_{3,1} + V_{3,2}) \quad \text{--- (1)}$$

$$Q_3 \cdot V_{3,1} = Q_3 \cdot \frac{Q_1}{4\pi\epsilon_0 R_{3,1}} = \frac{Q_1 Q_3}{4\pi\epsilon_0 R_{3,1}} = Q_1 \cdot V_{1,3}$$

$$\text{So, } W_E = Q_1 \cdot V_{1,2} + Q_1 \cdot V_{1,3} + Q_2 \cdot V_{2,3} + Q_2 \cdot V_{2,1} \quad \text{--- (2)}$$

(1) + (2)

$$2W_E = Q_1 (V_{1,2} + V_{1,3} + V_{1,4} + \dots) + Q_2 (V_{2,1} + V_{2,2} + \dots) + Q_3 (\dots) + \dots$$

But assume $V_{1,2} + V_{1,3} + V_{1,4} = V_1$

$$\text{So, } W_E = \frac{1}{2} (Q_1 V_1 + Q_2 V_2)$$

$$W_E = \frac{1}{2} \sum_{m=1}^N Q_m V_m$$

$$\text{Continuous charge distribution } W_E = \frac{1}{2} \int_V \rho_v dV \cdot V \rightarrow \text{Potential}$$

$$\delta V = \nabla \cdot \vec{D}$$

$$\nabla \cdot (\nabla \vec{V}) = \nabla \cdot (\nabla \cdot \vec{D}) + \vec{D} \cdot \nabla \nabla$$

$$\text{So, } W_E = \frac{1}{2} \int_{\text{vol}} (\nabla \cdot \vec{D}) V dV$$

$$= \frac{1}{2} \int_{\text{vol}} [\nabla \cdot (\nabla \vec{D}) - \vec{D} \cdot \nabla^2] dV$$

Divergence theorem $\oint \vec{D} \cdot d\vec{s} = \int \nabla \cdot \vec{D} dV$
(Volume integral converted into surface integral)

$$\text{So, } W_E = -\frac{1}{2} \int_{\text{vol}} \vec{D} \cdot (\nabla V) dV + \frac{1}{2} \int_S V \vec{D} \cdot d\vec{s}$$

$$W_E = -\frac{1}{2} \int_{\text{vol}} \vec{D} \cdot (\nabla V) dV$$

$$\vec{E} = -\nabla V$$

$$\text{So, } W_E = \frac{1}{2} \int_{\text{vol}} \vec{D} \cdot \vec{E} dV$$

$$W_E = \frac{1}{2} \int_{\text{vol}} \epsilon_0 E^2 dV$$

$$D = \epsilon_0 E$$

$$\text{Energy density} = \frac{dW_E}{dV}$$

$$= \frac{1}{2} \vec{D} \cdot \vec{E}$$

$$= \frac{1}{2} \epsilon_0 E^2 / \frac{D^2}{\epsilon_0}$$

$$= \epsilon_0 E^2$$

$$\frac{1}{2} \epsilon_0 E^2 = N$$

Ex- Point charge -1nc, +1nc and 3nc are located at $(0,0,0)$, $(0,0,1)$ & $(1,0,0)$ respectively. Find the energy in the SIS.

Sol:

$$W_1 = 0$$

$$W_2 = \frac{Q_2}{4\pi\epsilon_0} V_{2,1}$$

$$= \frac{Q_2 Q_1}{4\pi\epsilon_0 R_{1,2}}$$

$$R_{1,2} = (0,0,1) - (0,0,0)$$

$$\approx (0,1)$$

$$W_3 = \frac{Q_3}{4\pi\epsilon_0} V_{3,1} + \frac{Q_3}{4\pi\epsilon_0} V_{3,2}$$

$$= \frac{Q_3 Q_1}{4\pi\epsilon_0 R_{1,3}} + \frac{Q_3 Q_2}{4\pi\epsilon_0 R_{2,3}}$$

$$R_{1,3} = (1,0,0) - (0,0,0)$$

$$\approx (1,0,0)$$

$$R_{2,3} = (1,0,0) - (0,0,1)$$

$$\approx (1,0,-1)$$

$$\text{So, } W_E = 0 + \frac{(-4) \times 10^{-9} \times 10^{-9}}{4\pi\epsilon_0 (1)} \times \frac{9 \times 10^{-18}}{4\pi\epsilon_0 (1)} + \frac{12 \times 10^{-18}}{4\pi\epsilon_0 (2)}$$

$$= -7 \times 10^{-8} + \frac{12 \times 10^{-18}}{4\pi\epsilon_0 (2)}$$

$$= \frac{1.48 \times 10^{-18}}{4\pi\epsilon_0}$$

$$= 13.36 \times 10^{-19} \text{ J}$$

Ch-5.6

(190)

CURRENT AND CONDUCTORS

Convection & conduction currents:

$$I = \frac{dQ}{dt}$$

$$\text{current density } \vec{J} = \frac{AI}{AS}$$

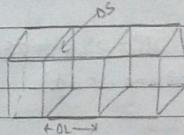
$$\Delta I = \vec{J} \cdot AS \rightarrow ($$

$$\Delta I = \vec{J} \cdot AS$$

$$I = \int_S \vec{J} \cdot dS$$

Convection current: it does not involve any conductor and not satisfy ohm's law. It occurs when current flow from insulator to insulator like vacuum, liquid etc.

Example: e⁻ beam in vacuum tube



$$\vec{J} = u_y \hat{a}_y$$

$$\Delta I = \frac{\Delta Q}{\Delta t}$$

$$\Rightarrow \vec{J} = \frac{\rho v \Delta V}{AS} = \rho v \cdot AS \frac{\Delta L}{\Delta t}$$

$$\Delta I = \rho v \Delta S u_y$$

$$J_y = \frac{\Delta I}{\Delta S} = \rho v u_y$$

$$\boxed{\vec{J} = \rho v \hat{a}_y} \quad \text{convection density}$$

Conduction current: it involves conductor and satisfies the ohm's law.

$$\vec{F} = -e\vec{E}$$

$m\ddot{\vec{u}} = -e\vec{E}$ \Rightarrow time of collision

τ

$$\ddot{\vec{u}} = -e\vec{E}$$
 (without resistance)

m

$$J_V = ne$$

$$\vec{J} = J_V \vec{U}$$

$$= (-ne) \left(-e\vec{E} \right)$$

$$\vec{J} = ne^2 \vec{E}$$

$$m$$

$$\boxed{\vec{J} = \sigma \vec{E}}$$
 → point form of ohm's law

conductivity

Num: If current den. $\vec{J} = \frac{1}{r^3} (\cos \theta \hat{r} + \sin \theta \hat{\theta}) A/m^2$, calculate current passing through

- A hemi-spherical shell of radius 20 cm.
- A spherical shell of radius 10 cm.

Sol:

$$i) I = \int \vec{J} \cdot d\vec{s}$$

$$d\vec{s} = r^2 \sin \theta d\theta d\phi \hat{\theta}$$

$$\int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} \frac{1}{r^3} (2 \cos \theta \hat{r} + \sin \theta \hat{\theta}) r^2 \sin \theta d\theta d\phi \hat{\theta}$$

$$= \frac{1}{r} \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} \cos \theta \left[\frac{-\cos \theta}{2} \right]_0^{\pi/2} d\theta d\phi$$

$$= \frac{25}{r} \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} \left[-\cos^2 \theta + \cos \theta \right] d\theta d\phi$$

$$= \frac{\pi \times 10^2}{20} [2] = \pi \times 10^2$$

$$= 31.4 A$$

ii)

$$I = \frac{2\pi}{r^2} \left[-\cos \theta \right]_0^{\pi}$$

$$= -\frac{\pi}{10 \times 10^{-2}} (1-1) = 0$$

Materials:

conductors, insulator, semiconductor (dielectric material)



$$\boxed{\vec{E} = 0}$$

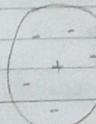
$$E = 0 \text{ bcz } E = \vec{J}/\sigma$$

$r \rightarrow \infty$

$$ii) V_{diff} = 0$$

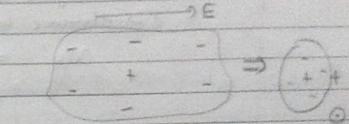
$$iii) \text{conductivity} = \infty$$

Polarization in Dielectric:



(*)

When we apply \vec{E} (electric field)



$$\vec{P}_d = \sigma \vec{E}$$

$$\vec{E} = -C \vec{E}$$

bcz (+) & (-) charges are separated because of E , then it will create dipole (align in dir. of E field)

$$\vec{P} = \sigma \vec{E}$$

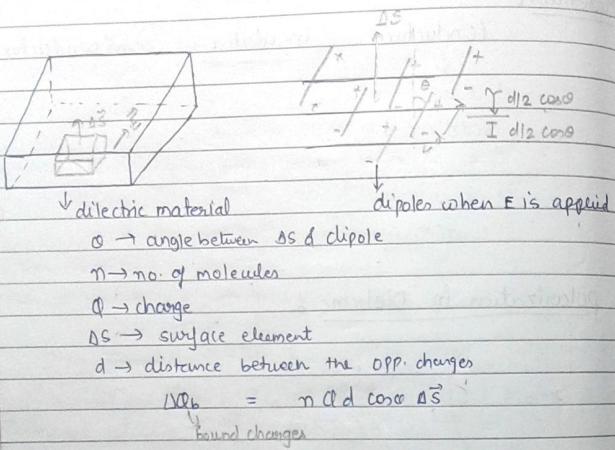
$$d_1 \vec{d}_1 + d_2 \vec{d}_2 + d_3 \vec{d}_3 + \dots + d_m \vec{d}_m = \sum_{k=1}^K d_k \vec{d}_k$$

$$\text{Polar intensity: } \vec{P} = \lim_{\Delta V \rightarrow 0} \sum_{\Delta V} \vec{d}_k \frac{dV}{\Delta V}$$

Non-polar mat: Hydrogen, Oxygen (where dipole is not present, but by applying E , then dipole is created)

Polar material: H_2O , (dipole is already available in the material)

polarization in dielectrics



$$\Delta Q_b = n Q d \cos \theta dS$$

$$\Delta Q_b = \vec{P} \cdot d\vec{S}$$

$$Q_p = - \oint \vec{P} \cdot d\vec{S}$$

↑ -ve sign just indicate the

the charges are present inside the surface, not outside the surface.

$$Q_p = \oint_s \epsilon_0 \vec{E} \cdot d\vec{S} = \oint_s \vec{D} \cdot d\vec{S}$$

$$Q_p = Q + Q_b$$

(free) (bounded)

$$Q = Q_T - Q_b$$

$$= \oint_s \epsilon_0 \vec{E} \cdot d\vec{S} + \oint_s \vec{P} \cdot d\vec{S}$$

$$Q = \oint_s (\epsilon_0 \vec{E} + \vec{P}) \cdot d\vec{S}$$

$$\text{so, } \vec{D} = \epsilon_0 \vec{E} + \vec{P} \rightarrow \text{dielectric flux density}$$

flux density of dielectric material, is dependent on electric field and polarization

$$\text{By divergence theorem } Q_p = \int_V \rho_p dv$$

$$Q = \int_V \rho_p dv$$

$$\nabla \cdot \vec{P} = -\rho_p$$

$$\nabla \cdot \epsilon_0 \vec{E} = \rho_p$$

$$\nabla \cdot \vec{D} = \rho_p$$

* (a) linear material: when D varies linearly with E .

(b) isotropic material: material for which D , B , E are in same direction. $\xrightarrow{\text{flux density}} \text{electric field intensity}$

Homogeneous material: material which E is constant.

Relationship between polarization & D for isotropic material

$$\vec{P} = \chi_e \epsilon_0 \vec{E}$$

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P}$$

$$= \epsilon_0 \vec{E} + \chi_e \epsilon_0 \vec{E}$$

$\rightarrow \epsilon_r \rightarrow \text{relative Permittivity}$

$$= \epsilon_0 \epsilon_r \vec{E}$$

$$\boxed{\vec{D} = \epsilon \vec{E}} \rightarrow \text{this relationship is true only for isotropic}$$

Ex:- Polystyrene
Electric field intensity in $\epsilon_r = 2.55$, filling the space between the plates of a parallel plate cap is 10 KV/m. The distance between the plate is 1.5 mm. Calculate flux density and polarization. ($\chi_e = 1.55$)

$$\epsilon_e = 1.55 ; \epsilon_r = 2.55$$

$$d = 1.5 \text{ mm} ; E = 10 \text{ KV/m}$$

$$\vec{D} = \epsilon_0 E + \vec{P}$$

$$\text{polarization: } \vec{P} = \rho_p \epsilon_0 \epsilon_r \vec{E}$$

$$= 1.55 \times 10^{-12} \times 2.5 \times 10^{-3}$$

$$= 1.55 \times 2.5 \times 10^{-12} \times 10 \times 10^3 = 13.75 \times 10^{-8}$$

$$= 13.75 \text{ nC/m}^2$$

Relaxation time



$$\text{flux density } D = \epsilon_0 \epsilon_r E \\ = 8.85 \times 10^{-12} \times 10 \times 10^3 \times 2.55 \\ = 22.56 \text{ n}$$

$$\text{Potential diff } V = E \cdot d \\ = (10 \times 10^3) (1.5 \text{ mm}) \\ = 15 \text{ V}$$

$$\textcircled{1} D \text{ of free charges on plates} = \vec{D} \cdot \vec{n} = D$$

Continuity Equation:

principle of charge convergence:
-ve sign bcz charge moving
out, hence decrease in the

$$J_{\text{out}} = \oint \vec{J} \cdot d\vec{s} = - \frac{dQ_{\text{in}}}{dt} \quad \text{--- (1)}$$

$$\text{divergence theorem} \quad \oint \vec{J} \cdot d\vec{s} = \int \nabla \cdot \vec{J} dV \quad \text{--- (2)}$$

$$-\frac{dQ_{\text{in}}}{dt} = -\frac{d}{dt} \int_V \rho_V dV = - \int \frac{\partial \rho_V}{\partial t} dV$$

$$\int_V \nabla \cdot \vec{J} dV = - \int \frac{\partial \rho_V}{\partial t} dV$$

$$\textcircled{2} \quad \nabla \cdot \vec{J} = - \frac{\partial \rho_V}{\partial t}$$

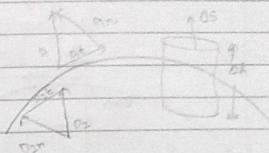
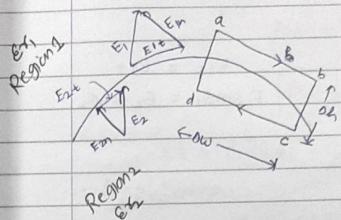
for steady current; $\nabla \cdot \vec{J} = 0$,
i.e., no movement of charges. Hence
no change in density.

Boundary Conditions:

$$\oint \vec{E} \cdot d\vec{l} = 0 \quad (\text{work done in a closed path} = 0)$$

$$\oint \vec{D} \cdot d\vec{s} = q_{\text{enclosed}}$$

(i) Dielectric - dielectric B.C.



$$\vec{E}_1 = \vec{E}_{1t} + \vec{E}_{1n} \quad (\text{long & normal component})$$

$$\vec{E}_2 = \vec{E}_{2t} + \vec{E}_{2n}$$

$$0 = E_{1t} \cdot \Delta w - E_{1n} \frac{\Delta h}{2} - E_{2t} \cdot \Delta w + E_{2n} \frac{\Delta h}{2} + E_{m} \frac{\Delta t}{2}$$

$$E_{1t} = |E_{1t}| \quad \& \quad E_{2t} = |E_{2t}|$$

$$0 = (E_{1t} - E_{2t}) \Delta w$$

$$\Delta h \rightarrow 0$$

$$\boxed{E_{1t} = E_{2t}} \quad \rightarrow \text{continuous}$$

$$\vec{D} = \epsilon \vec{E} = \vec{D}_t + \vec{D}_n$$

$$\frac{D_{1t}}{E_1} = E_{1t} = E_{2t} = \frac{D_{2t}}{E_2}$$

$$\text{OR} \quad \boxed{\frac{D_{1t}}{E_1} = \frac{D_{2t}}{E_2}} \quad \rightarrow \text{discontinuous}$$

$$\partial Q = \partial_s A_S = \partial_{in} A_S - \partial_{out} A_S$$

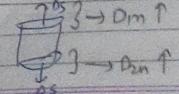
$$D_{in} - D_{out} = P_s$$

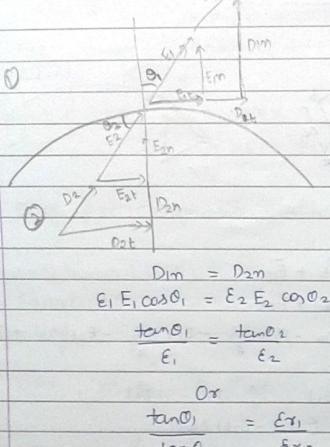
$$P_s = 0 \Rightarrow$$

$$\boxed{D_{in} = D_{out}} \quad \rightarrow \text{continuous}$$

$$\boxed{E_1 E_{1n} = E_2 E_{2n}}$$

direction of D_{in} & D_{out} are same





$$E_1 t = E_2 t$$

$$E_1 \sin \theta_1 = E_2 \sin \theta_2$$

$$D_{1m} = D_{2m}$$

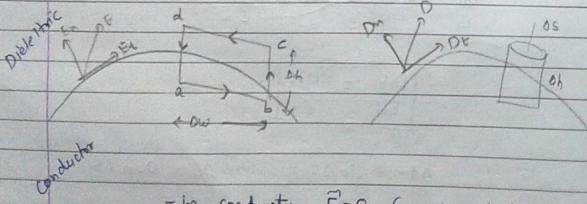
$$E_1 E_1 \cos \theta_1 = E_2 E_2 \cos \theta_2$$

$$\frac{\tan \theta_1}{E_1} = \frac{\tan \theta_2}{E_2}$$

Or

$$\frac{\tan \theta_1}{\tan \theta_2} = \frac{E_2}{E_1}$$

(iii) Conductor - Dielectric B.C.



- In conductor $E=0$ (equipotential surface)

$$0 = \frac{0 \cdot Dm}{2} + \frac{0 \cdot D2m}{2} + \frac{E_n D2m}{2} - E_t Dm - E_n D2m - \frac{0 \cdot D1m}{2}$$

$$Dm \rightarrow 0$$

$$[E_t = 0]$$

$$\Delta Q = D_m \Delta S - 0 \cdot D_m \Delta S$$

$$\vec{D} = \epsilon \vec{E} = 0$$

$$D_m = \frac{\Delta Q}{\Delta S} = \rho_s$$

$$[D_m = \rho_s]$$

$$\epsilon E_r = \rho_s$$

$$E_r = \frac{\rho_s}{\epsilon}$$

$$- D_t = \epsilon_0 \epsilon_r E_t = 0 \quad \& \quad D_m = \epsilon_0 \epsilon_r E_n = \rho_s$$

Application:

- Shielding

Example two extensive homogeneous isotropic dielectrics meet on plane $z=0$. For $z > 0$, $\epsilon_{r1}=4$ for $z < 0$, $\epsilon_{r2}=3$ a uniform field $\vec{E}_1 = 5 \hat{i} - 2 \hat{j} \text{ KV/m}$ exists for $z > 0$

(a) find \vec{E}_2 for $z < 0$

(b) angle \vec{E}_1 & \vec{E}_2 makes with interference

(c) the energy density in both the dielectric

(d) the energy within a cube of side 2m centered at $(3, 4, -5)$ \rightarrow in second region

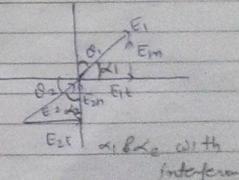
Sol:

$$(a) E_m = \vec{E}_1 \cdot \vec{a}_m$$

$$= \vec{E}_1 \cdot \vec{a}_2^2$$

$$= 3 \text{ KV/m}$$

$$\vec{E}_m = 3 \vec{a}_2^2 \text{ KV/m}$$



$$\vec{E} = \vec{E}_1 + \vec{E}_2$$

$$\vec{E}_2 = 5 \vec{a}_2 - 2 \vec{a}_2 \text{ KV/m}$$

$$\vec{E}_1 = 5 \vec{a}_1 - 2 \vec{a}_1 = \vec{E}_t$$

$$\vec{E}_m = \frac{E_m}{\epsilon_r} \vec{E}_1 = \frac{4}{3} (3) = 4 \text{ KV/m}$$

$$\vec{E}_2 = \vec{E}_{2t} + \vec{E}_n$$

$$\vec{E}_2 = 5\hat{x} - 2\hat{y} + 4\hat{z} \text{ kV/m Ans.}$$

(b)

$$\tan \theta_1 = \frac{|E_{1t}|}{|E_{1n}|} = \frac{(5^2 + 2^2)^{1/2}}{3} \Rightarrow \theta_1 = 60^\circ$$

$$\tan \theta_2 = \frac{|E_{2t}|}{|E_{2n}|} = \frac{(5^2 + 2^2)^{1/2}}{4} \Rightarrow \theta_2 = 53.7^\circ$$

$$\alpha_1 = 99.12^\circ$$

$$\alpha_2 = 36.6^\circ$$

(c)

$$\text{energy density}_1 = \frac{1}{2} \epsilon_0 E_1^2 \epsilon_r$$

$$= \frac{1}{2} \times 8.85 \times 10^{-12} \times 4 \times (5^2 - 2^2)^2$$

$$= 10.94 \mu J/m^3$$

$$\text{en. den}_2 = \frac{1}{2} \epsilon_0 \epsilon_r E_2^2$$

$$= \frac{1}{2} \times 8.85 \times 8 \times (5^2 - 2^2)^2 \times 10^{-12}$$

$$= 5.9 \mu J/m^3$$

(d) $\vec{E} \cdot \vec{dl}$

$$W = \int \vec{E} \cdot d\vec{l}$$

$$= \int_2^4 E_x dx + \int_3^5 E_y dy + \int_{-6}^4 E_z dz$$

$$= 5(2) + (-2)(5-3) + 4(-4+6)$$

$$= 10 - 4 + 8 = 14 \text{ J}$$

$$F_x = 5(2) - 2(2) + 3(2) = 10 + 6 - 4 = 12 \text{ N}$$

(iii) Conductor-free space BC

$$E_t = 0$$

$$\vec{D} = \epsilon_0 \vec{E} = 0$$

$$E_t = 1$$

$$D_n = Ps = 0.0105$$

$$P_n = Ps$$

$$D_n = \epsilon_0 E_n = Ps$$

$$E_n = Ps/\epsilon_0$$

$$E_n = Ps/E_0$$

\Rightarrow Derivation of Poisson's and Laplace eqn 8

$$\nabla \cdot \vec{D} = Pv \quad \dots \quad (1)$$

$$\vec{D} = \epsilon \vec{E} \quad \dots \quad (2)$$

$$\vec{E} = -\nabla V \quad \dots \quad (3)$$

$$\nabla \cdot (\epsilon \vec{E}) = -\nabla \cdot (\epsilon \nabla V) = Pv$$

OR

$$\text{poisson's eqn} \quad \nabla \cdot \nabla V = -\frac{Pv}{\epsilon} \quad \dots \quad (4)$$

divergence of gradient of V

rect. coord. S

$$\nabla \cdot \vec{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$$

$$\nabla V = \frac{\partial V}{\partial x} \hat{x} + \frac{\partial V}{\partial y} \hat{y} + \frac{\partial V}{\partial z} \hat{z}$$

$$\nabla \cdot \nabla V = \frac{\partial}{\partial x} \left(\frac{\partial V}{\partial x} \right) + \frac{\partial}{\partial y} \left(\frac{\partial V}{\partial y} \right) + \frac{\partial}{\partial z} \left(\frac{\partial V}{\partial z} \right)$$

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = -\frac{Pv}{\epsilon}$$

If charge & potential is given at particular boundary
then we can only used poisson & laplace eqn.

$\nabla^2 V = 0$

$\nabla^2 V = 0 \rightarrow$ it's called as Laplace eqn
 \downarrow Laplacian operator

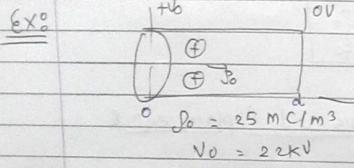


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$$\text{Rectangular: } \nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}$$

$$\text{Cylindrical: } \nabla^2 V = \frac{1}{r} \frac{\partial}{\partial r} \left(\frac{\partial V}{\partial r} \right) + \frac{1}{r^2} \left(\frac{\partial^2 V}{\partial \theta^2} \right) + \frac{\partial^2 V}{\partial z^2}$$



here we can't use Laplace eqn bcz
($\rho V + 0$ given)

$$\nabla^2 V = -\frac{\rho}{\epsilon} \quad \text{Poisson's eqn}$$

charge in z direction only \rightarrow

$$\frac{dV}{dz} = -\frac{\rho}{\epsilon}$$

$$\frac{dV}{dz} = -\frac{\rho}{\epsilon} z + A$$

$$V = -\frac{\rho}{\epsilon} z^2 + Az + B$$

$$\text{At } z=0, V=V_0 \quad \text{at } z=d, V=0$$

$$B=V_0 = 2KV \quad | \quad 0 = -\frac{\rho}{\epsilon} d^2 + A \cdot d + 2KV$$

$$B=V_0; \quad A = \frac{\rho}{\epsilon} d - \frac{V_0}{d}$$

$$\text{So } V = -\frac{\rho}{\epsilon} z^2 + \left(\frac{\rho}{\epsilon} d - \frac{V_0}{d} \right) z + V_0$$

$$\text{Field intensity } E = -\nabla V = -\frac{dV}{dz} \hat{a}_z$$

$$= \left(\frac{\rho}{\epsilon} z - A \right) \hat{a}_z$$

$$= \left(\frac{\rho}{\epsilon} z - \frac{\rho}{\epsilon} d + \frac{V_0}{d} \right) \hat{a}_z$$

net force

$$F = qE$$

$$F = \int \rho V E dV = \rho \int ds \int E dz$$

$$= \rho \cdot S \left[\frac{V_0 z}{d} + \frac{\rho}{\epsilon} (z^2 - dz) \right] \hat{a}_z$$

$$F = \rho V \cdot S \cdot V_0 \hat{a}_z \quad N = \rho V \cdot S \left[\frac{V_0 + \frac{\rho}{\epsilon} z^2 - \rho}{d} \right] \hat{a}_z$$

force per unit area or pressure is

$$P = F/A$$

$$= \rho V \cdot S = 25 \times 10^{-3} \times 22 \times 10^3$$

$$= 550 \text{ N/m}^2$$

Uniqueness Theorem

Solution of Laplace & Poisson's eqn is V , is that unique or not?

↪ its unique

V_1 & V_2 are soln of Laplace or Poisson's eqn

$$\nabla^2 V_1 = 0 \quad \text{and} \quad \nabla^2 V_2 = 0$$

$$\nabla^2 (V_1 - V_2) = 0$$

$V_b \rightarrow$ potential at boundary

$$V_1 \text{ at the boundary} = V_{1b}$$

$$V_2 \text{ " " " } = V_{2b}$$

$$V_{1b} = V_{2b} = V_b \quad (\text{bcuz at band, potentials are same})$$

or

$$V_{1b} - V_{2b} = 0$$

$$\nabla \cdot (V \vec{B}) \equiv V (\nabla \cdot \vec{B}) + \vec{B} \cdot (\nabla V)$$

$$V \rightarrow V_1 - V_2 \quad | \quad V_1 - V_2 = \text{scalar} \quad (\text{potential is scalar})$$

$$\vec{B} \rightarrow \nabla (V_1 - V_2) \quad | \quad \nabla (V_1 - V_2) = \text{vector}$$

$$\nabla \cdot ((V_1 - V_2) \nabla (V_1 - V_2)) \approx (V_1 - V_2) \nabla \cdot (\nabla V_1 - \nabla V_2) + (\nabla V_1 - \nabla V_2) \cdot \nabla (V_1 - V_2)$$

integrating the eqⁿ
 (By divergence theorem) $\oint A \cdot d\vec{s} = \int_V \nabla \cdot A \, dv$

$$\text{Left side} = \oint (V_{1b} - V_{2b}) \nabla (V_{1b} - V_{2b}) \, d\vec{s}$$

now

$V_{1b} = V_{2b}$ (By the boundary condition)

$$= 0$$

Right side

$$\nabla \cdot \nabla (V_1 - V_2)$$

$$\text{but } V_{1b} = V_{2b}$$

$$\therefore \nabla^2 (V_1 - V_2) = 0$$

$$\int_{\text{vol}} [\nabla (V_1 - V_2)]^2 \, dv = 0$$

$$[\nabla (V_1 - V_2)]^2 = 0$$

$$\nabla (V_1 - V_2) = 0$$

$$V_1 - V_2 = 0$$

$$V_1 = V_2$$

- importance of the theorem ①

and how we can say that the soln of laplace & poisson's eqⁿ are unique

Capacitance

- Capacitor holds charges and stored electric energy.

$$Q \propto V \quad \text{--- ①}$$

$$\begin{array}{c} + \\ | \rightarrow | \\ + \end{array}$$

$$C = Q/V$$

\hookrightarrow the measure of capacity
 capacitor, capacitance

parallel plate spherical cylindrical

$$\vec{E} = \sigma / \epsilon_0$$

$$E = \frac{Q}{A \epsilon_0}$$

$$V = \vec{E} \cdot d$$

$$= \frac{Qd}{A \epsilon_0}$$

So in eqⁿ ①

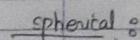
$$C = \frac{A \epsilon_0}{d}$$

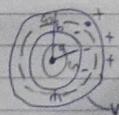
$A \rightarrow$ area of plates

$d \rightarrow$ distance between plates

Here, we have to maintain the min distance between capacitor plates, so that plates can't attract each other

-  \rightarrow Pseudo cap.
 ↓ dielectric material

spherical 



$$\frac{Q}{\epsilon_0} = \oint E \cdot dA$$

due to $C_1 \Rightarrow d = 0$
 due to $C_2 \Rightarrow \epsilon_r \neq 1$