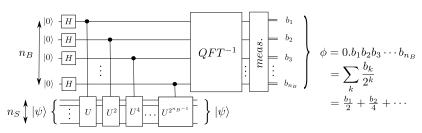
Quantum phase estimation

An example from chemistry

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quantum phase estimation: original form



With $|\psi\rangle$ eigenvector of U with eigenvalue $e^{-2i\pi\phi}$,

$$(U|\psi
angle = e^{-2i\pi\phi}|\psi
angle$$
, $0\leq\phi<1$)

The circuit above allows to compute the n_B most significant bits of ϕ in a **binary decomposition**.

Template for *all* quantum algorithms with **exponential speedup**. (i.e with quantum circuits of **polynomial** size for classically-intractable problems)

iterative quantum phase estimation: what we will use

for k in range $(n_B, 0, -1)^1$:

- 1. compute $\phi_k = \sum_{l=k+1}^{n_B} \frac{b_l}{2^{l-k+1}}$
- 2. execute the following circuit:

$$|0\rangle - H - RZ(\phi_k) - H - meas = b_k$$

$$n_S \quad |\psi\rangle \left\{ - U^{2^{k-1}} \right\} |\psi\rangle$$

With
$$H=rac{1}{\sqrt{2}}egin{pmatrix}1&1\\1&-1\end{pmatrix}$$
 and $RZ(\phi_k)=egin{pmatrix}1&0\\0&e^{i\phi_k}\end{pmatrix}$



¹Python style

$$|0\rangle - H - RZ(\phi_k) - H - meas. = b_k$$

$$n_S \quad \Big|\psi\rangle \Big\{ - U^{2^{k-1}} \Big\} |\psi\rangle$$

Recall that $\phi=0.b_1\cdots b_{n_B}$ in binary decomposition and $\phi_k=\sum_{l=k+1}^{n_B}\frac{b_l}{2^{l-k+1}}=0.0b_{k+1}\cdots b_{n_B}.$

Evolution of the system:

$$\begin{split} &|0\rangle|\psi\rangle \underset{H}{\rightarrow} \left(\frac{|0\rangle+|1\rangle}{\sqrt{2}}\right)|\psi\rangle = \left(\frac{|0\rangle|\psi\rangle+|1\rangle|\psi\rangle}{\sqrt{2}}\right) \\ &\underset{\bullet - U}{\rightarrow} \left(\frac{|0\rangle|\psi\rangle+e^{-2i\pi\cdot2^{k-1}\phi}|1\rangle|\psi\rangle}{\sqrt{2}}\right) = \left(\frac{|0\rangle+e^{-2i\pi\cdot0.b_k\cdots b_{n_B}}|1\rangle}{\sqrt{2}}\right)|\psi\rangle \\ &\underset{RZ(\phi_k)}{\leftrightarrow} \left(\frac{|0\rangle+e^{-2i\pi\cdot0.b_k}|1\rangle}{\sqrt{2}}\right)|\psi\rangle \underset{H}{\rightarrow} |b_k\rangle|\psi\rangle \end{split}$$

Our use-case: ground state energy calculation

Given \mathcal{H} , Hamiltonian of target system (e.g a molecule), we want to compute its ground state energy.

Actually, usually \mathcal{H} depends on some structural/geometric parameters (e.g H_2 dissociation curve)

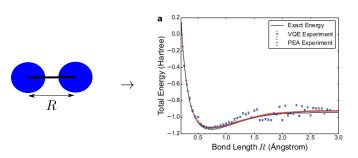


FIG. 3. Computed H₂ energy curve and errors. (a) Ene

$$|0\rangle - H - RZ(\phi_k) - H - meas - b_k$$

$$n_S \downarrow |\psi\rangle \left\{ - U^{2^{k-1}} \right\} |\psi\rangle$$

We will therefore work with:

- $V = e^{-i\mathcal{H}dt}$, for some value of dt to be chosen.
- $|\psi\rangle$ the ground state of H, whose energy E we want to compute.
- we then have $U|\psi\rangle=e^{-iEdt}|\psi\rangle=e^{-2i\pi\frac{Edt}{2\pi}}$

Problem 1: the algorithm only allows us to compute $\phi = \frac{Edt}{2\pi} \mod [1]$ ($\in [0,1]$)

Following [2], we will use an *interval guess* $[E_{min}, E_{max}]$ for E. Then $\mathcal{H} \leftarrow \mathcal{H} - E_{min} \cdot I$ ensures E > 0 and choosing $dt = \frac{2\pi}{E_{max}}$ ensures $\frac{Edt}{2\pi} < 1$



Other problems

Problem 2 In general we do not have access to the ground state.

 \rightarrow use an **ansatz**² with theoretical good overlap with ground state.

Problem 3 How do we implement $U^{2^k} = e^{-i\mathcal{H}2^k dt}$?

ightarrow usually $\mathcal H$ comes as a sum of simpler terms $\sum_l \mathcal H_l$ (e.g Pauli products) for which $e^{-i\mathcal H_l dt}$ is simple to implement.

 \rightarrow and use **Trotter's formula** [2]

$$e^{-i\left(\sum_{l}\mathcal{H}_{l}\right)dt}\simeq\left(\prod_{l}e^{-rac{i\mathcal{H}_{l}dt}{p}}
ight)^{p}$$

with p integer, as large as possible.



²easy-to-prepare approximation

Peter JJ O'Malley, Ryan Babbush, Ian D Kivlichan, Jonathan Romero, Jarrod R McClean, Rami Barends, Julian Kelly, Pedram Roushan, Andrew Tranter, Nan Ding, et al. Scalable quantum simulation of molecular energies. *Physical Review X*, 6(3):031007, 2016.

James D Whitfield, Jacob Biamonte, and Alán Aspuru-Guzik. Simulation of electronic structure hamiltonians using quantum computers.

Molecular Physics, 109(5):735–750, 2011.