Particle Physics Phenomenology

Part 10

March 20, 2018

The Quark-Gluon Plasma and heavy-ion collisions: an experimental overview

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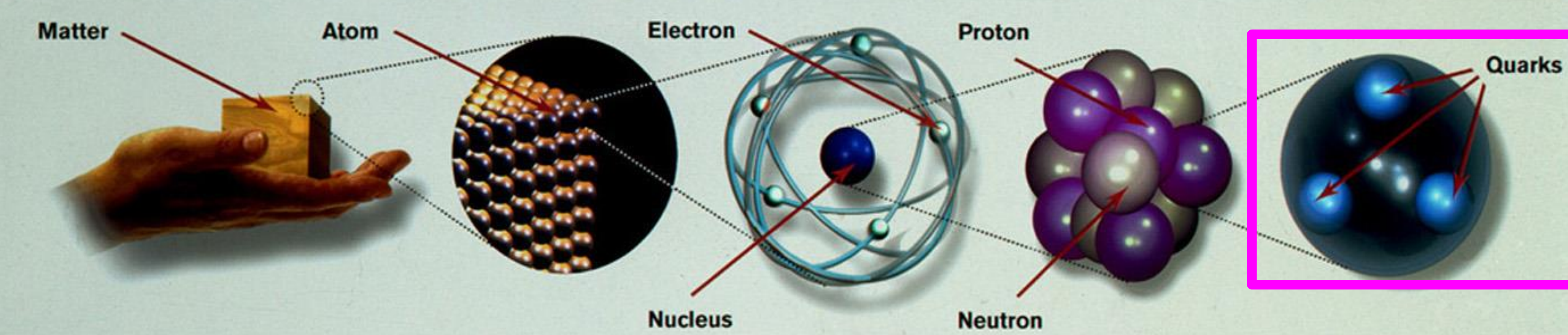
Outline

- Part 1
  - The medium temperature
  - Hard probes
    - The standard candles
    - Jets and high p_T particles
    - Quarkonium
- Part 2
  - Soft medium properties
    - Collective flow
  - Quark Gluon Plasma (QGP) in small systems?



Matter particles

All ordinary particles belong to this group

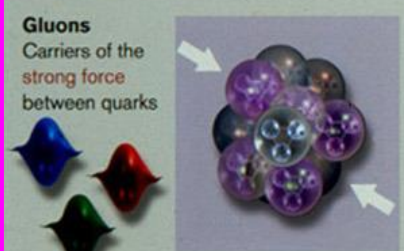
	LEPTONS		QUARKS	
FIRST FAMILY	**Electron** Responsible for electricity and chemical reactions; it has a charge of -1	**Electron neutrino** Particle with no electric charge, and possibly no mass; billions fly through your body every second	**Up** Has an electric charge of plus two-thirds; protons contain two, neutrons contain one	**Down** Has an electric charge of minus one-third; protons contain one, neutrons contain two
SECOND FAMILY	**Muon** A heavier relative of the electron; it lives for two-millionths of a second	**Muon neutrino** Created along with muons when some particles decay	**Charm** A heavier relative of the up; found in 1974	**Strange** A heavier relative of the down; found in 1964
THIRD FAMILY	**Tau** Heavier still; it is extremely unstable. It was discovered in 1975	**Tau neutrino** not yet discovered but believed to exist	**Top** Heavier still	**Bottom** Heavier still; measuring bottom quarks is an important test of electroweak theory

These particles existed just after the Big Bang. Now they are found only in cosmic rays and accelerators

Force particles

These particles transmit the four fundamental forces of nature although gravitons have so far not been discovered

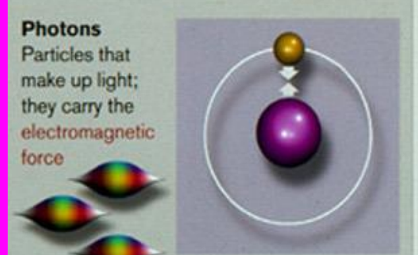
Gluons
Carriers of the strong force between quarks



Felt by:
quarks

The explosive release of nuclear energy is the result of the strong force

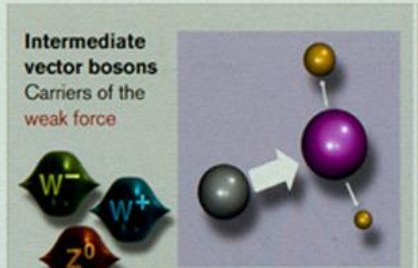
Photons
Particles that make up light; they carry the electromagnetic force



Felt by:
quarks and charged leptons

Electricity, magnetism and chemistry are all the results of electro-magnetic force

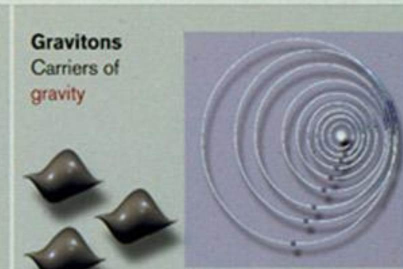
Intermediate vector bosons
Carriers of the weak force



Felt by:
quarks and leptons

Some forms of radio-activity are the result of the weak force

Gravitons
Carriers of gravity



Felt by:
all particles with mass

All the weight we experience is the result of the gravitational force

GRAPHICS: PETER CROWTHER

What is the Quark-Gluon Plasma? (1/2)

- Today quarks and gluons are confined inside hadrons
- Question: has hadrons always existed?
- Answer: not in the very early Universe
- Question: what was there before hadrons?
- Answer: the quarks and gluons were deconfined (a Plasma). This state of matter is called the Quark-Gluon Plasma

And we recreate this phase of matter at LHC!

What is the Quark-Gluon Plasma? (2/2)

“`QGP' is not a new particle but a paradigm-shift of how we understand matter in extreme conditions. Thus there is no clear criterion by which one can claim an experimental discovery. For this reason QGP was rediscovered again with new experimental results obtained at the order of magnitude higher RHIC collision energies. In addition to confirming the CERN results, RHIC produced new puzzling phenomena; some will be discussed below. The circumstance repeats for the third time today: LHC data confirms SPS and RHIC results, and is offering another rich field of new experimental results.

Since no one plans to announce the QGP discovery at the LHC, we conclude that QGP has gained considerable acceptance as a new form of matter.”

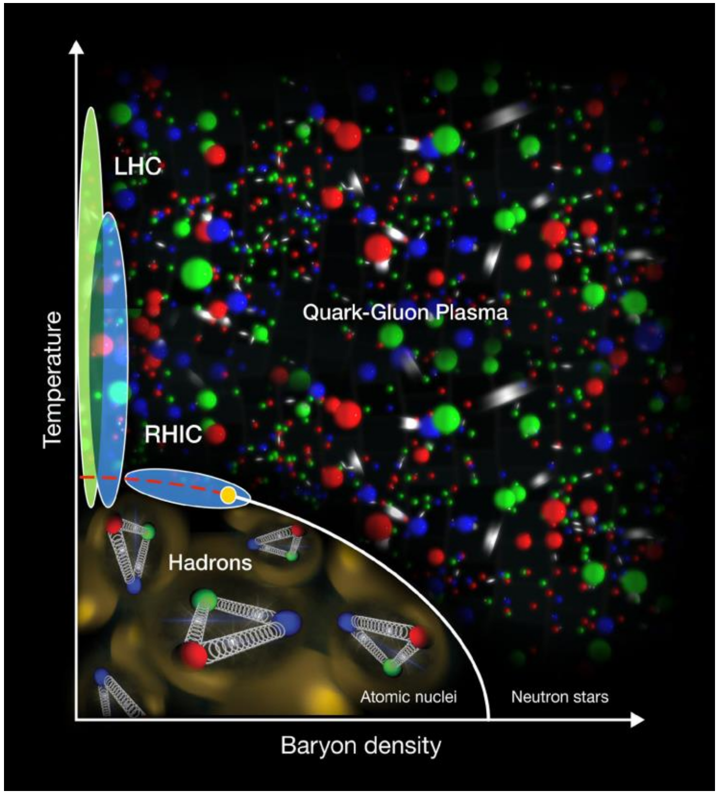
Johann Rafelski and Jeremey Birrell, arXiv:1311.0075.

They point out a big problem: our main understanding is experimental and very little of the theory is on a firm basis.

When we will have a better theoretical picture we can then really point out what the QGP properties are and when it was discovered.

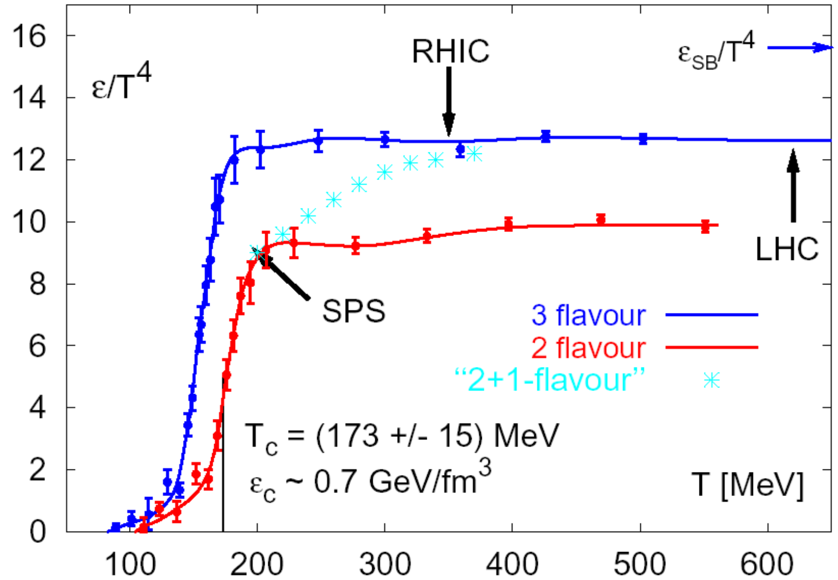
Deconfinement at high energy densities

Tc ~ 160-170 MeV
(2.000.000.000.000 K)



I will only talk about the high temperature transition which is what we probe in heavy-ion collisions

Lattice QCD calculation of the energy density



$$\varepsilon_{Quark-Gluon\ gas} = \frac{\pi^2}{30}\left(2\times 8 + \frac{7}{8}2\times 2\times 3\times 3\right)T^4$$

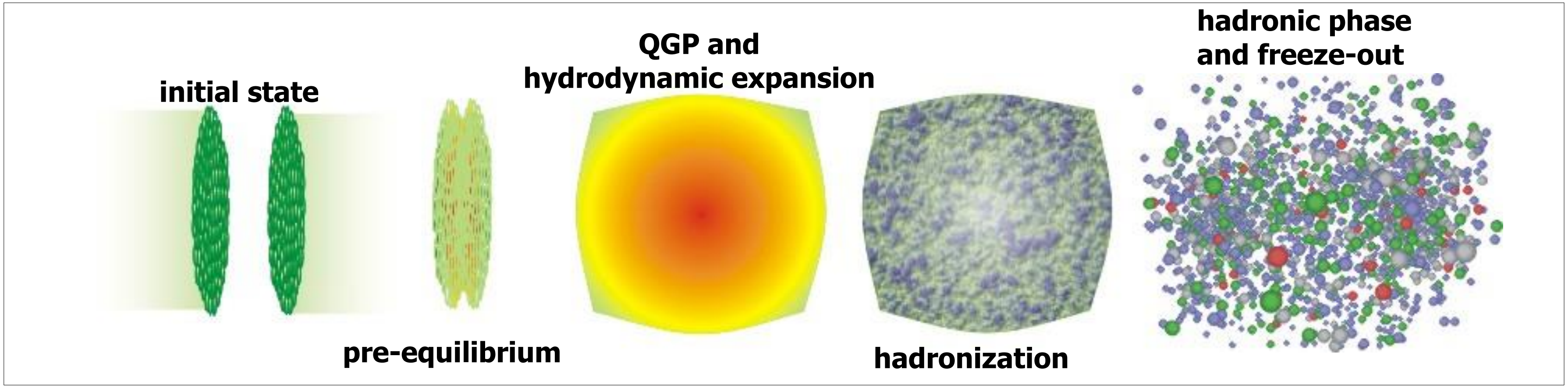
Gluon spin and color

(Anti+)quark spin, color and flavor

With lattice QCD one can study some aspects of QCD numerically. In this case the energy density shows that at a temperature of ~170 MeV there is a phase transition.

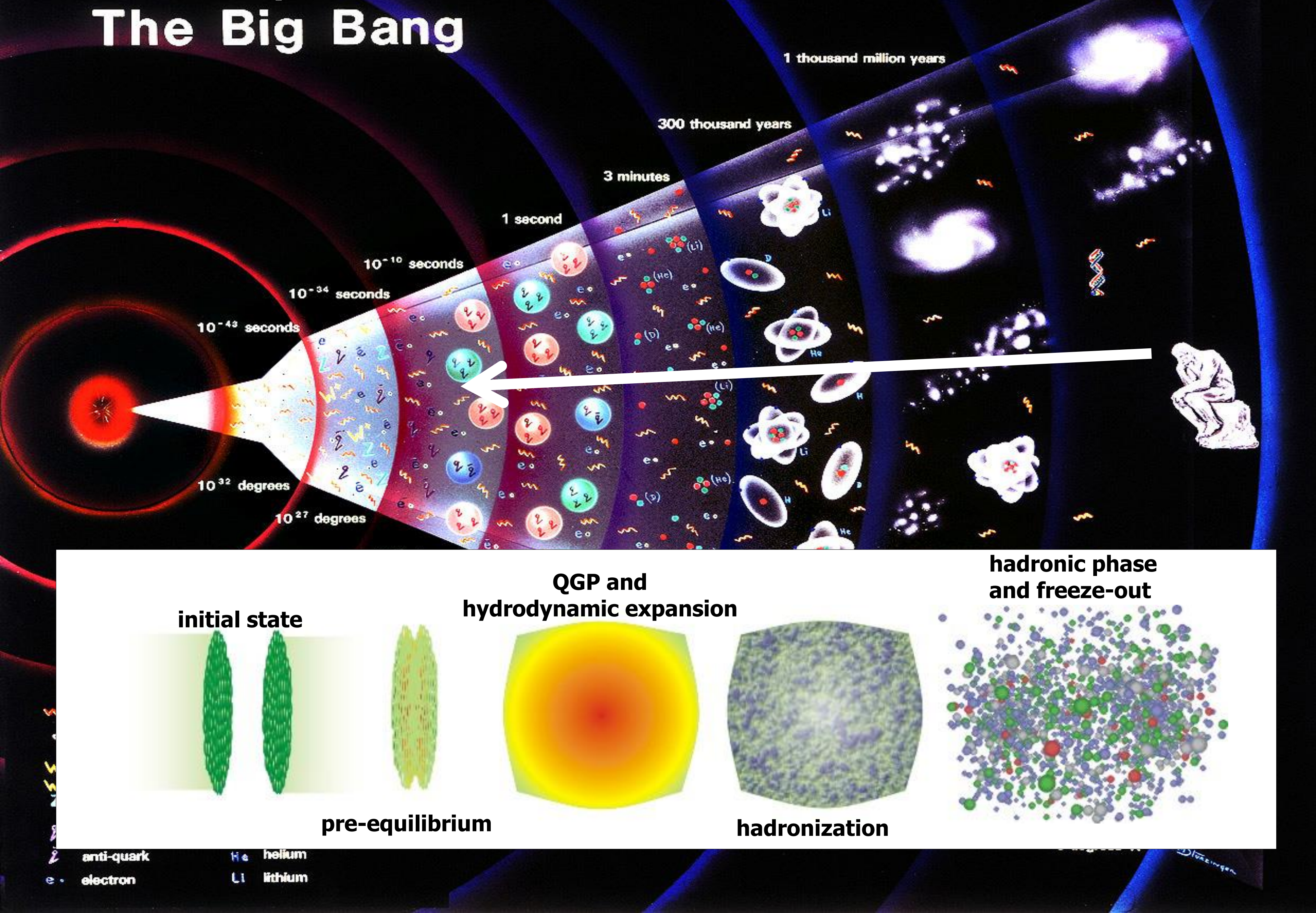
The phase transition is believed to be a Xover, meaning that for temperatures around Tc the hadronic and QGP phases coexist and that no entropy is produced in the phase transition.

Heavy ion collisions



- The only way we can create the QGP in the laboratory!
- By colliding heavy ions it is possible to create a large (»1fm^3) zone of hot and dense QCD matter
- Goal is to create and study the properties of the Quark Gluon Plasma
- Experimentally mainly the final state particles are observed, so the conclusions have to be inferred via models

QGP – the phase of the universe a few micro seconds after



History of ultra-relativistic heavy-ion physics

- 1^{st} generation $\sqrt{s_{NN}}$<20GeV (fixed target) AGS (BNL, US), SPS (CERN) late 80s and 90s
  - Experiments: NA61, NA49, NA60, NA50, NA44, ……
- 2^{nd} generation $\sqrt{s_{NN}}$<200GeV (collider) RHIC (BNL, US) 2000-Now
  - Experiments: PHENIX, STAR, PHOBOS, BRAHMS
- 3^{rd} generation $\sqrt{s_{NN}}$=2760GeV in run 1 (collider), $\sqrt{s_{NN}}$=5020GeV in run 2 (2015-) LHC (CERN) 2010-Now
  - Experiments: ALICE, CMS, and ATLAS.

I will mainly show results from LHC!

Establishing the QGP

HOT

COLORED

DECONFINED

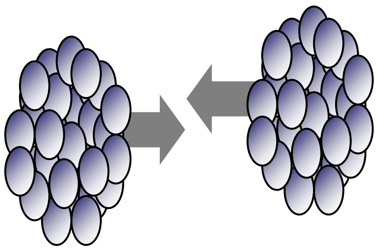
MEDIUM

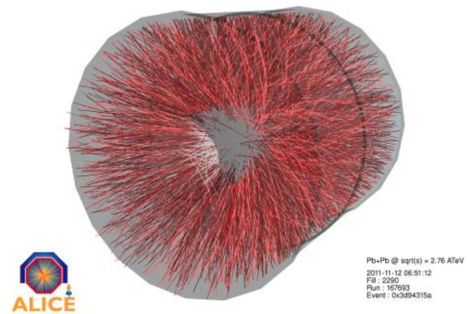
THERMALIZED

And then we have to also measure its properties!

The three systems (understanding before 2012)

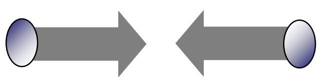
Pb-Pb

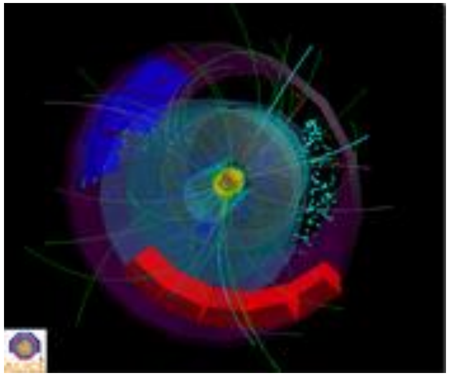




Hot QCD matter:
This is where we expect the QGP to be created in central collisions.

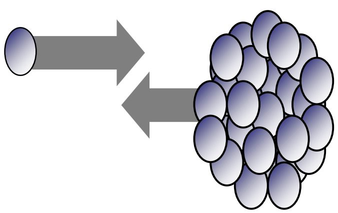
pp

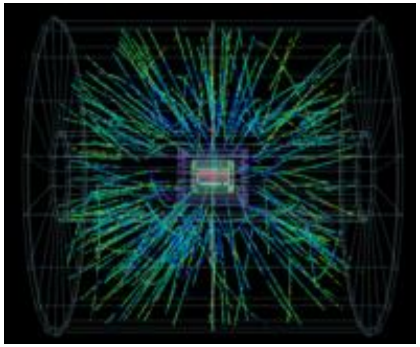




QCD baseline:
This is the baseline for "standard" QCD phenomena.

p-Pb





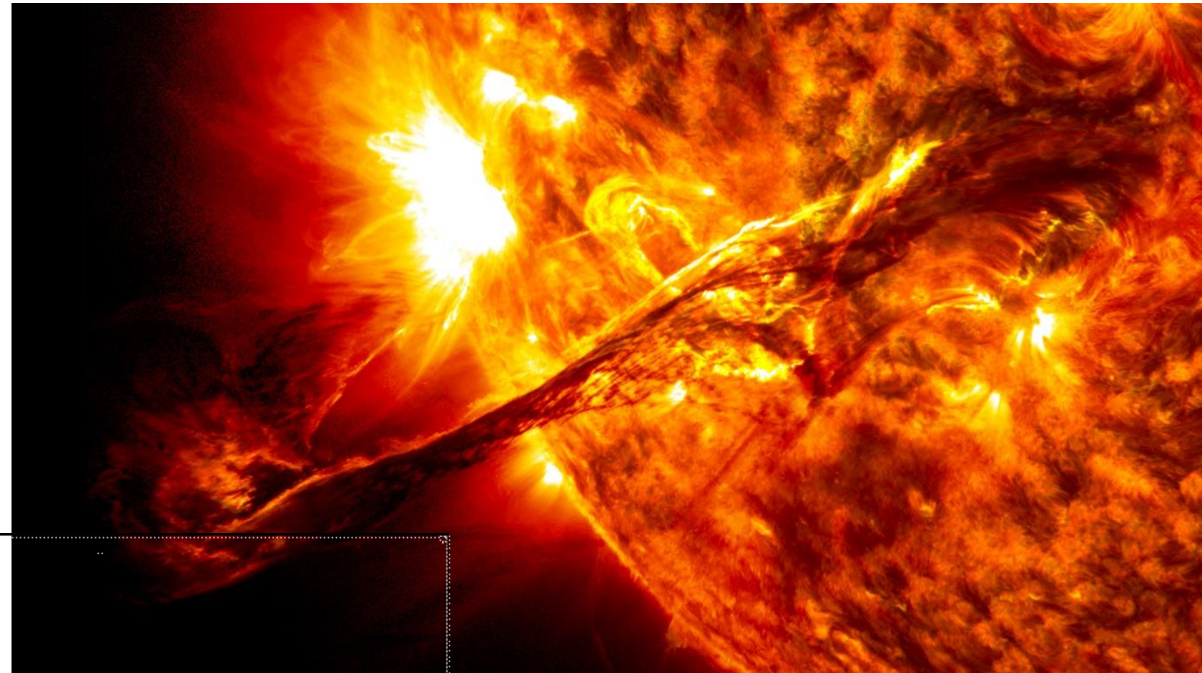
Cold QCD matter:
This is to isolate nuclear effects, e.g. nuclear pdfs.

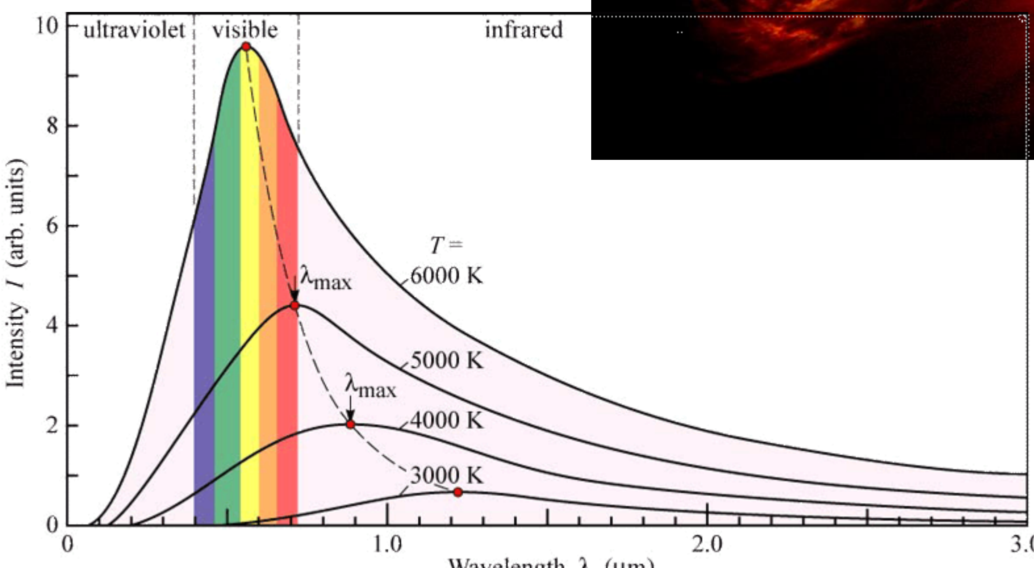
THERMAL PHOTONS

Measuring the medium temperature

- To establish the necessary conditions for a medium one would like to measure the temperature/energy density
- There are in general three possibilities
  - Extracting the final temperature from the p_T slope and particle yields (this is the final temperature)
  - Bjorken has made a famous relation between the initial energy density and the final transverse energy (using hydrodynamics and assumptions on when a medium is formed) Phys. Rev. D27 (1983) 140-151 (2300 citations)
  - Measuring the thermal photon spectrum
    - Photons only interacts weakly with the QGP and so one can probe the early times directly (one measures the time integrated spectrum)

In a similar way as we measure the surface temperature of the sun…

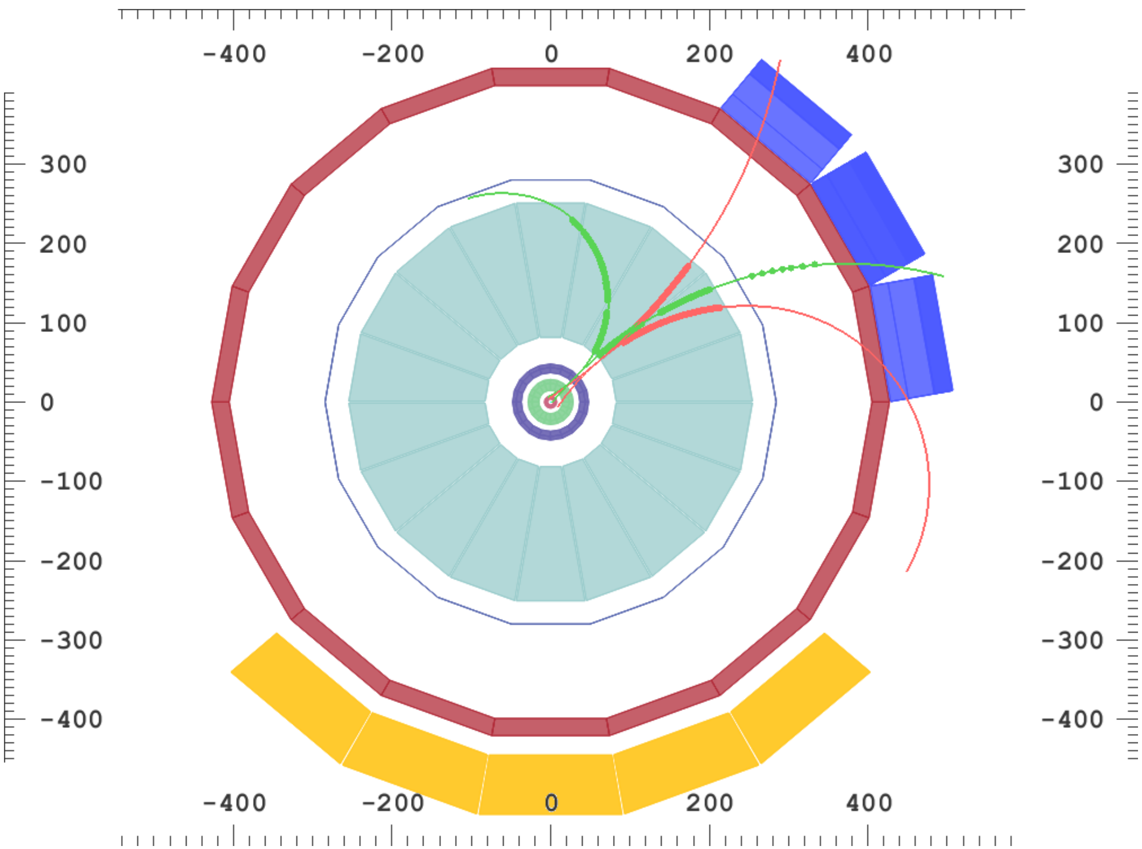


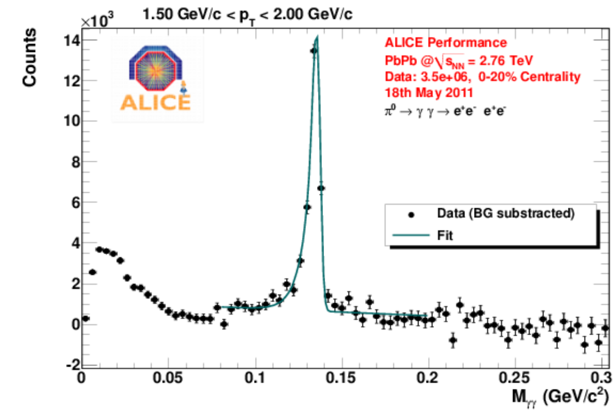


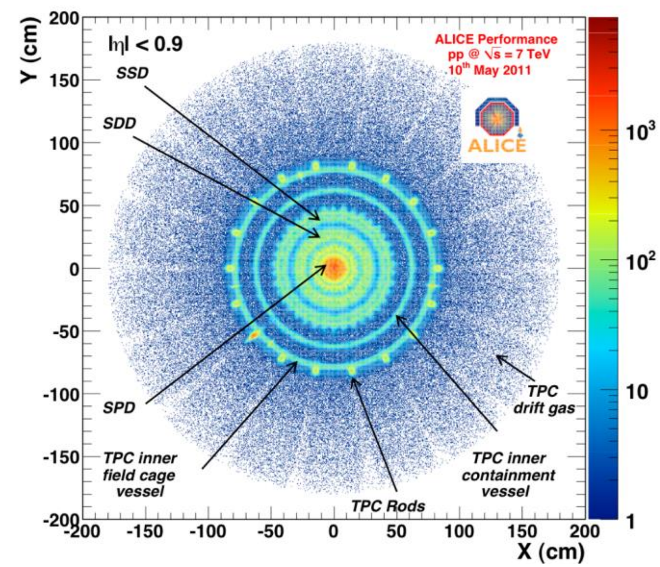
… we can measure the temperature of the QGP

Photon identification in the TPC via conversion

$\pi^0 \rightarrow 2\gamma$ converts to $2*(e^- + e^+)$ = background for signal



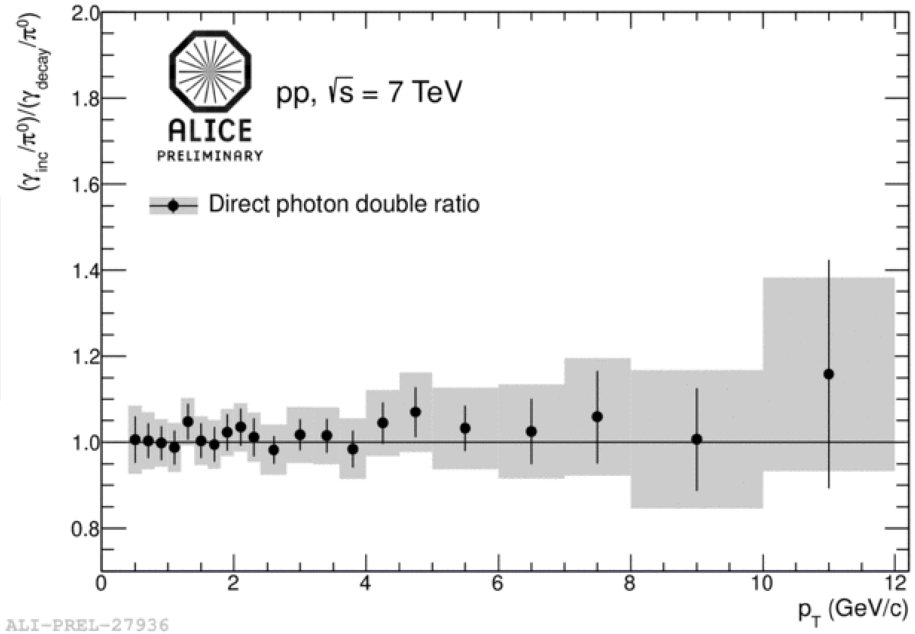




Direct photons in pp collisions

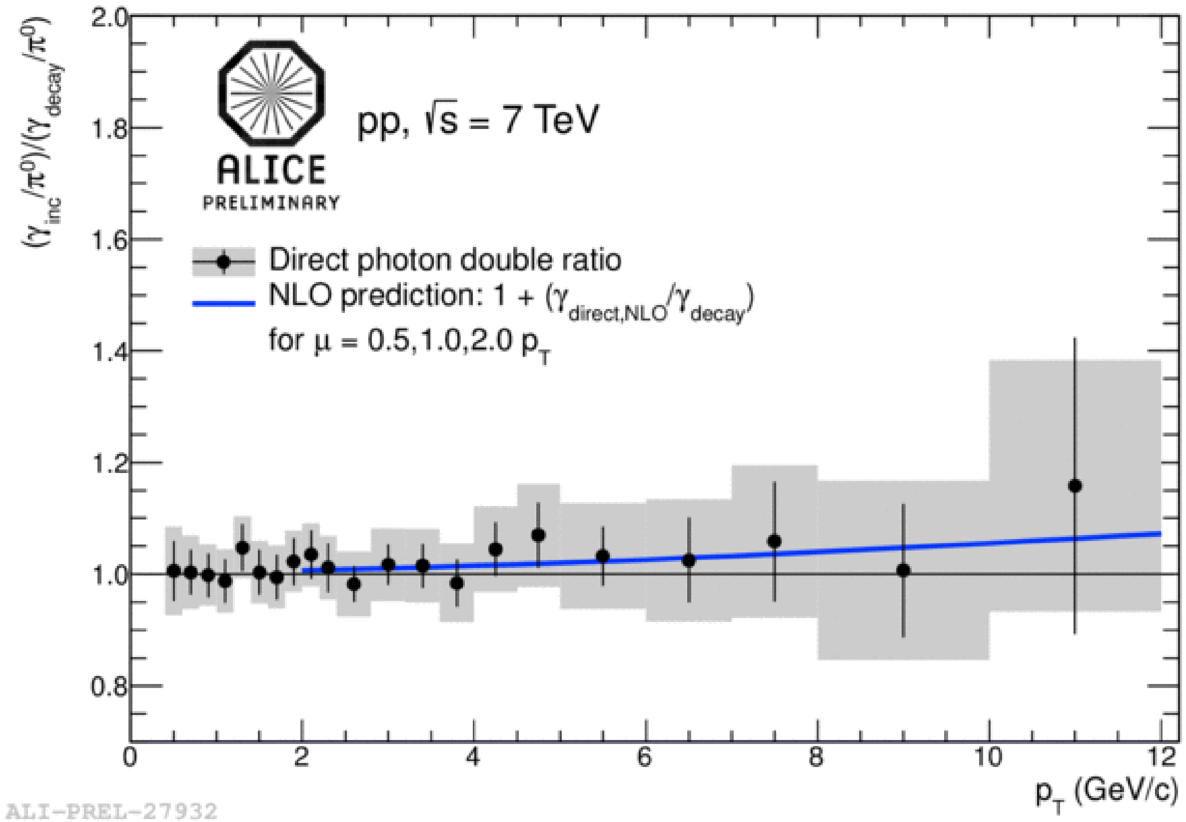
Double Ratio: $\frac{\gamma_{inc}}{\pi^0} \Big/ \frac{\gamma_{decay}}{\pi^0_{param}} \approx \frac{\gamma_{inc}}{\gamma_{decay}}$

$\rightarrow$ cancellation of uncertainties

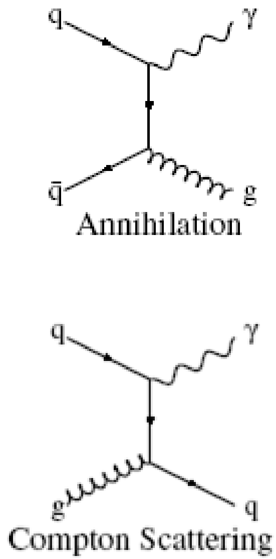


- Construct double ratio to elliminate/reduce systematics
- Numerator is the actual measurement
- Denominator is from a cocktail calculation

Direct photons in pp collisions

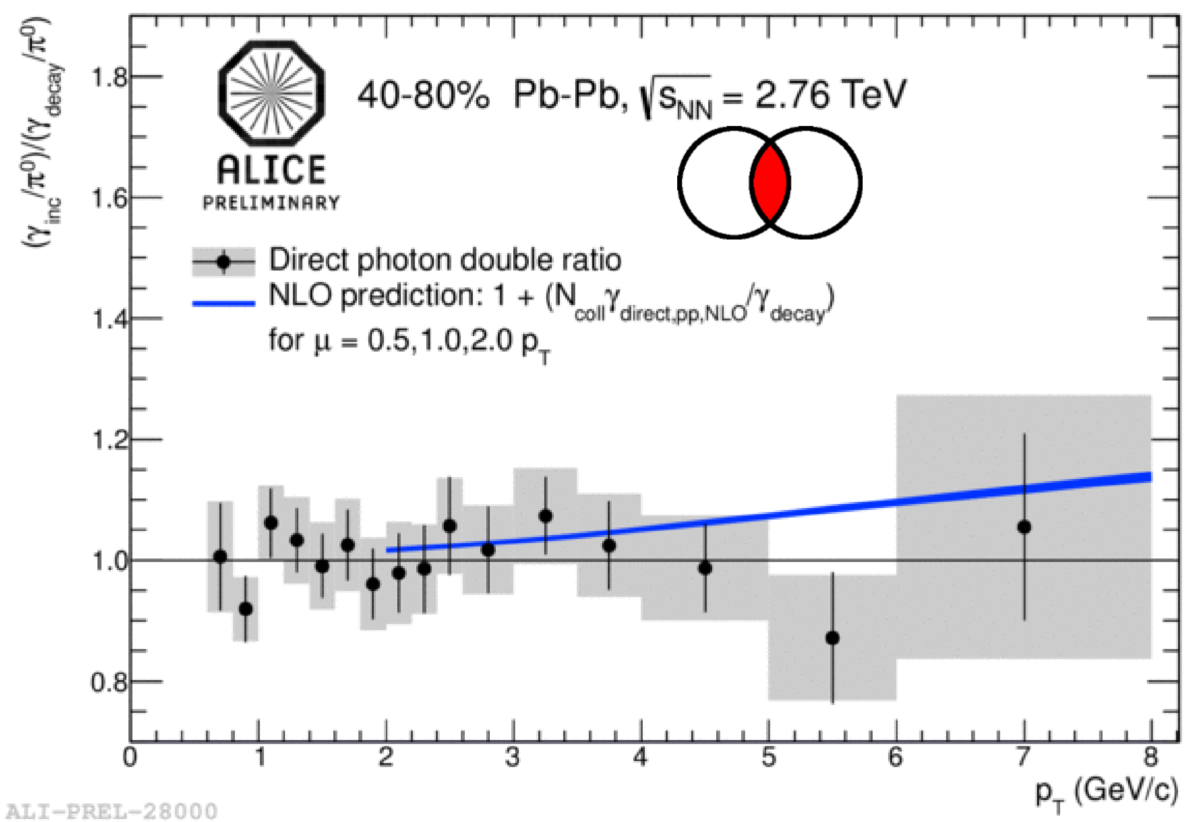


Sources of direct photons



- Comparison to pQCD NLO calculation

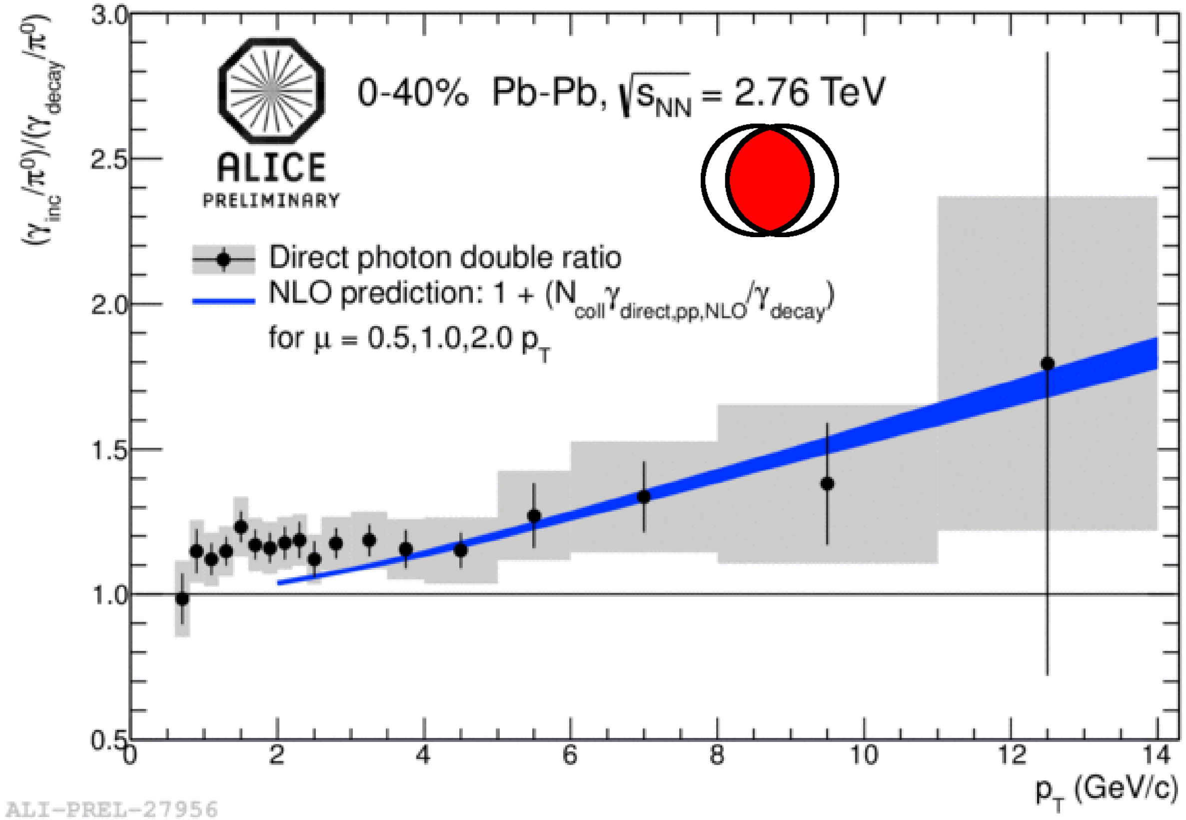
Direct photons in peripheral Pb-Pb collisions



- Peripheral Pb-Pb

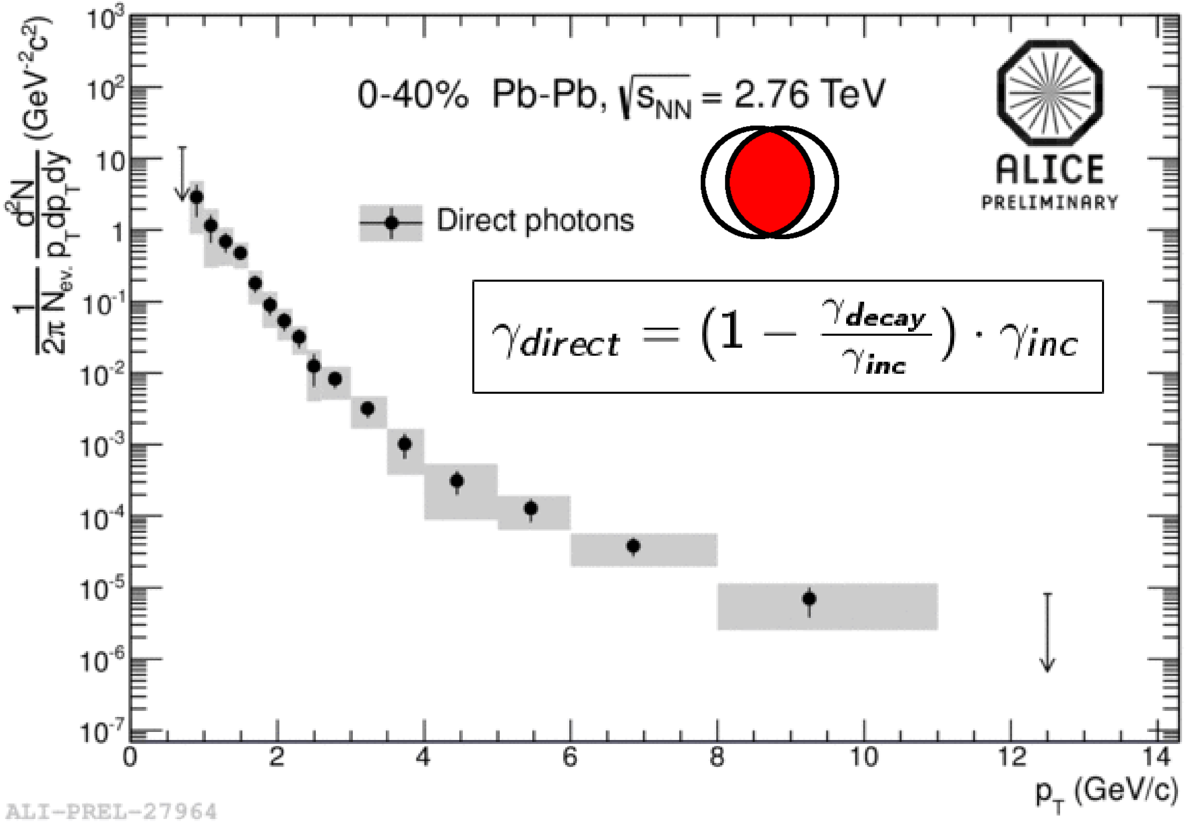
  Consistent with only direct and decay photons

Direct photons in central Pb-Pb collisions



- Central Pb-Pb
  - Surplus of photons at high pT is expected from hard production
  - Surplus of photons at low pT are from thermal radiation from the QGP

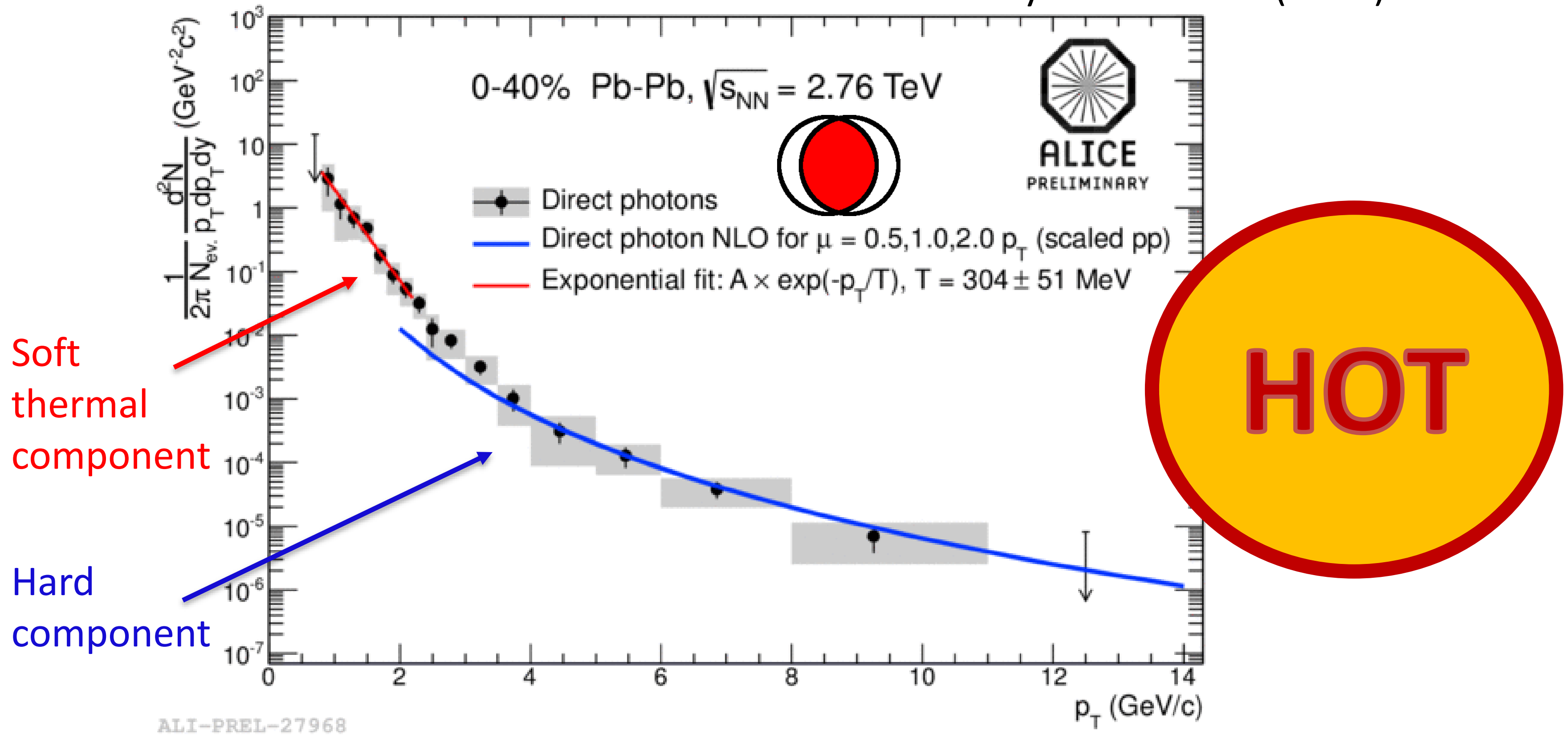
Direct photon spectrum in central Pb-Pb



- Obtain the direct γ spectrum by scaling with the inclusive γ spectrum

Direct photon spectrum in central Pb-Pb

Phys. Lett. B 754 (2016) 235-248



- The temperature of the low p_T direct γ spectrum is of order 300 MeV (recall that this is a time average over the lifetime of the medium)

The highest man-made temperature

Highest man-made temperature

Share

Who

CERN, LARGE HADRON COLLIDER

What

5X10^12 DEGREE(S) KELVIN

Where

SWITZERLAND

When

13 AUGUST 2012

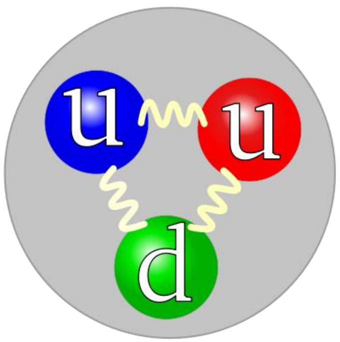
On 13 August 2012 scientists at CERN's Large Hadron Collider, Geneva, Switzerland, announced that they had achieved temperatures of over 5 trillion K and perhaps as high as 5.5 trillion K. The team had been using the ALICE experiment to smash together lead ions at 99% of the speed of light to create a quark gluon plasma – an exotic state of matter believed to have filled the universe just after the Big Bang.

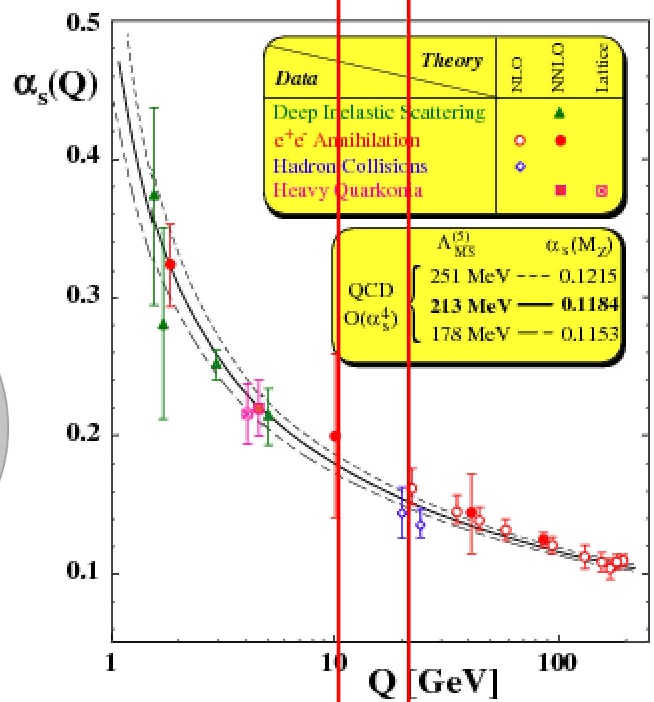
HARD PROBES

What happens in pp and Pb-Pb collisions – a simple picture

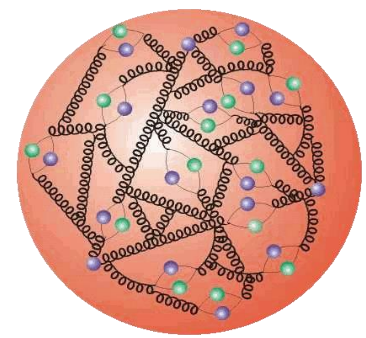
- 2 limits

SOFT





HARD



Non-perturbative physics
(knows the equations but not how to solve them)
Bulk properties (=QGP)

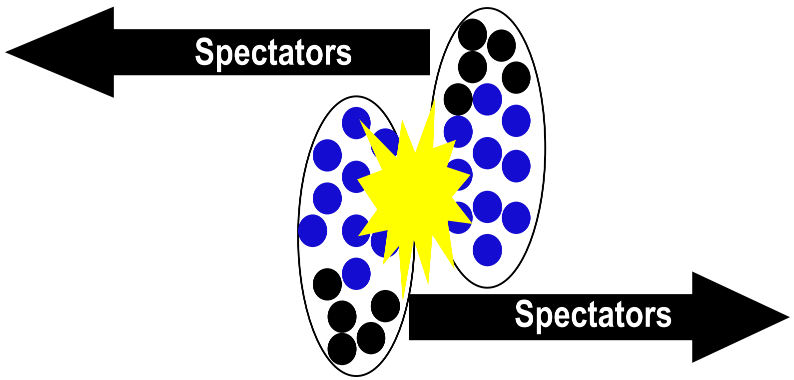
Perturbative physics
(theoretical predictions)
Rare jets (=probes)

Hard probes in heavy-ion collisions

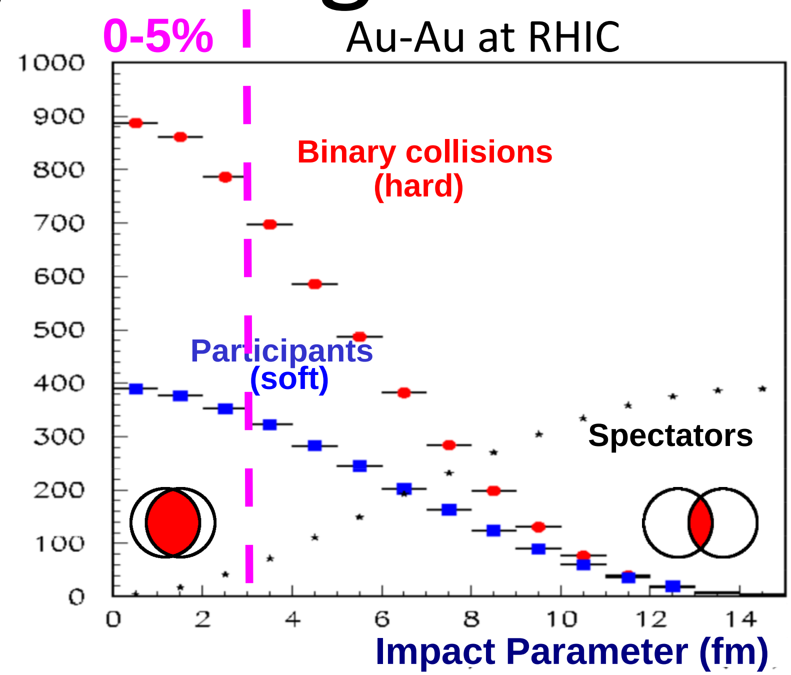
- The cross sections for Hard probes (HP) can in principle be calculated using pQCD (+ pdfs and FFs). In practice they need to be measured to achieve the best systematic precision.
- The cross sections can easily be extrapolated from pp to Pb-Pb collisions using binary scaling (next slide) unless there are "nuclear" effects
  - HP allow the precise study of nuclear effects
- Penetrating probes
  - Photons, W, Z, used to benchmark the binary scaling
- Strongly interacting probes
  - Heavy quarks (quarkonia), jets

The collision geometry and binary scaling

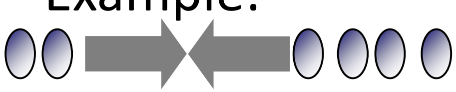
- Centrality (ex. for Au+Au):



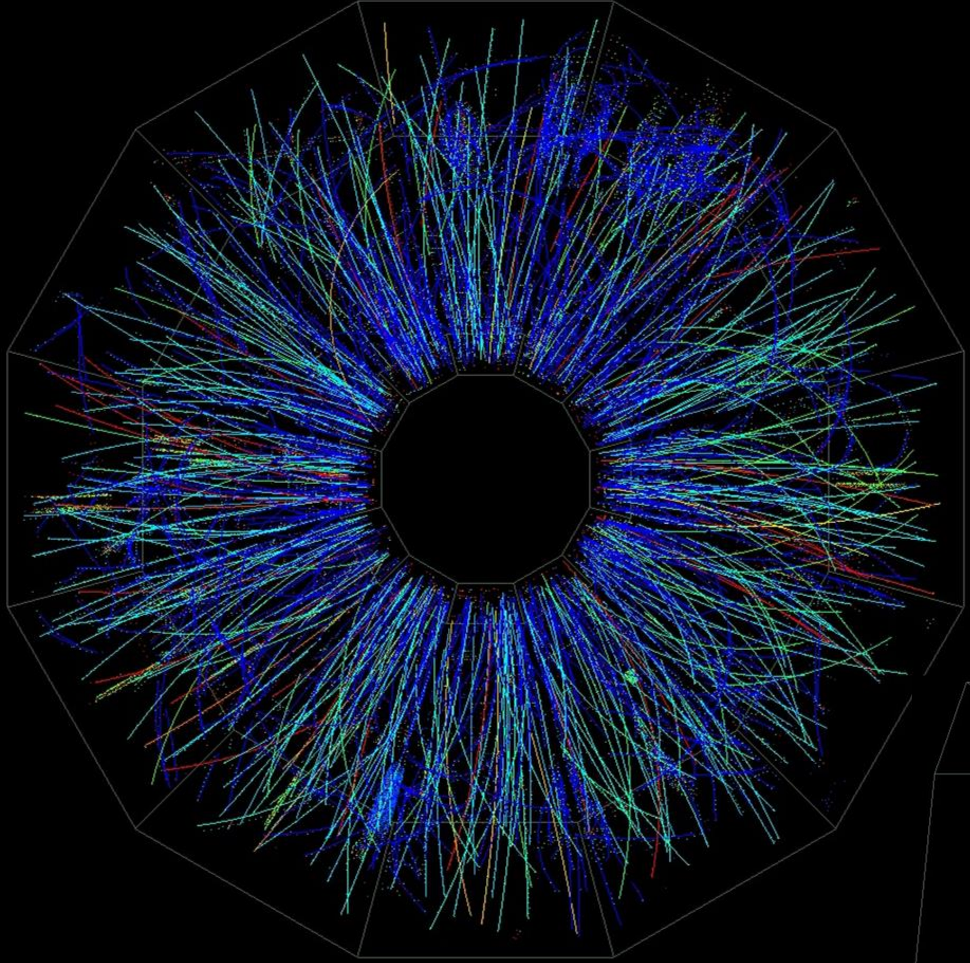
Participants = 2*197 - Spectators



- The "medium" energy is proportional to the # of participant (Npart)
- The number of parton-parton (quark-quark, quark-gluon, gluon-gluon) is proportional to the # of binary collisions (Nbin)
- Example:

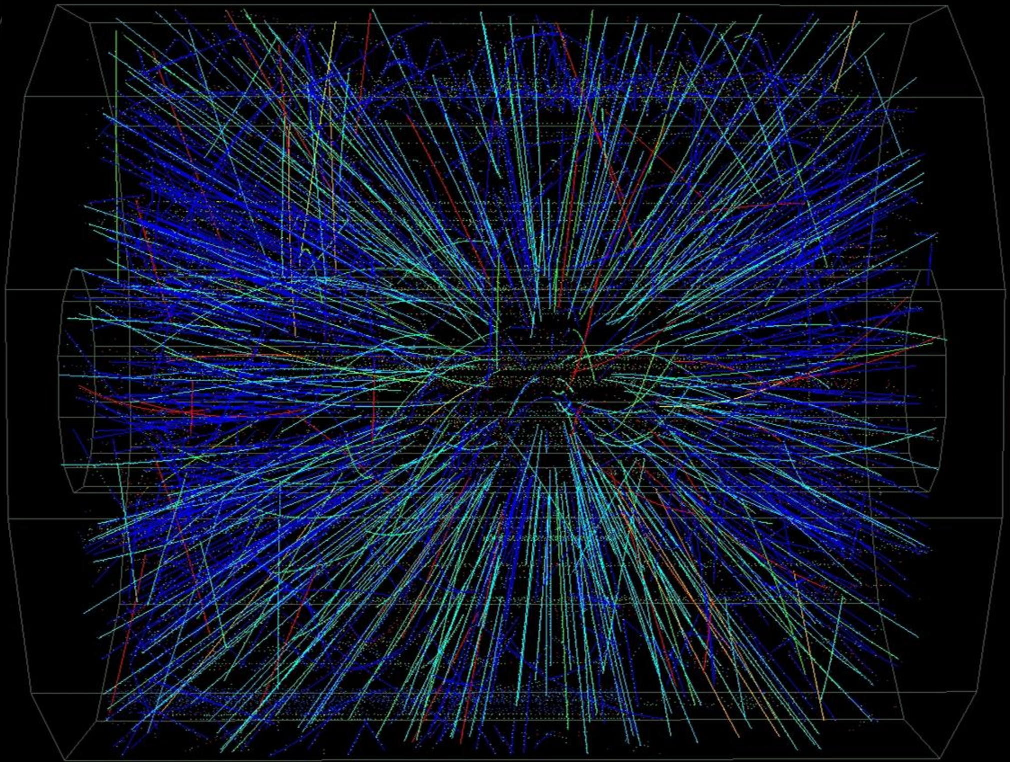


6 participant
8 binary collisions
(pp has 2 participant and 1 binary collision)



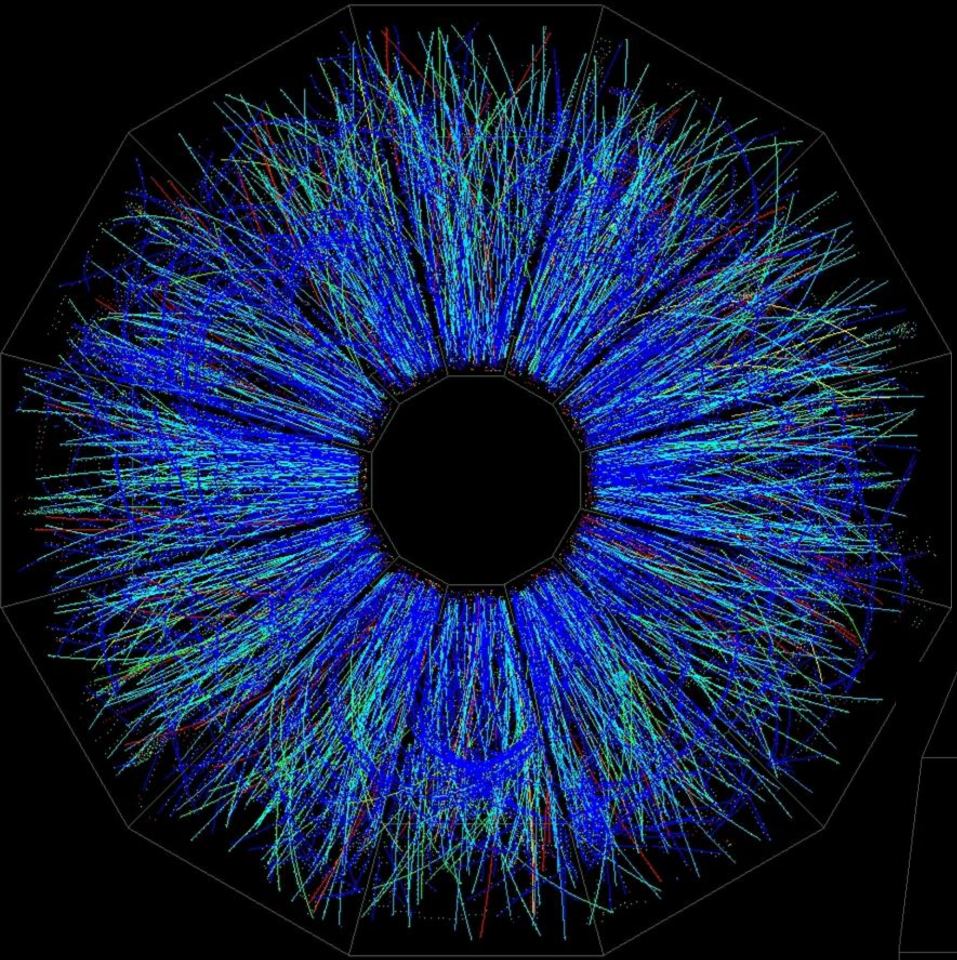
Peripheral Event

From real-time Level 3 display.



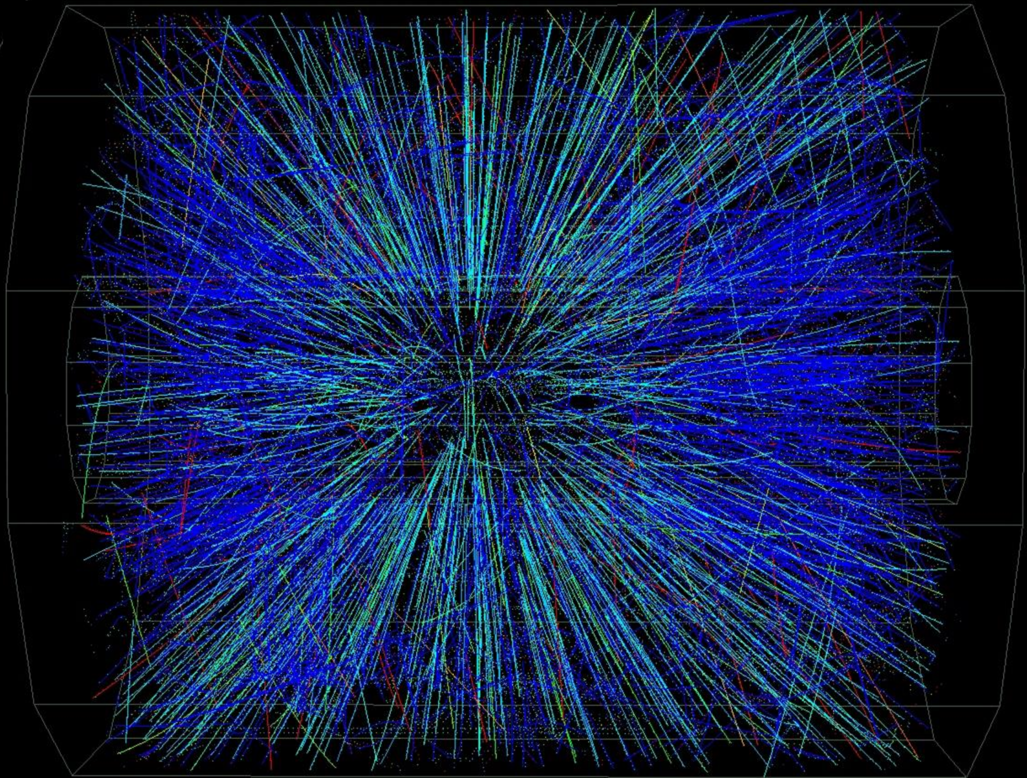
color code $\Rightarrow$ energy loss

STAR

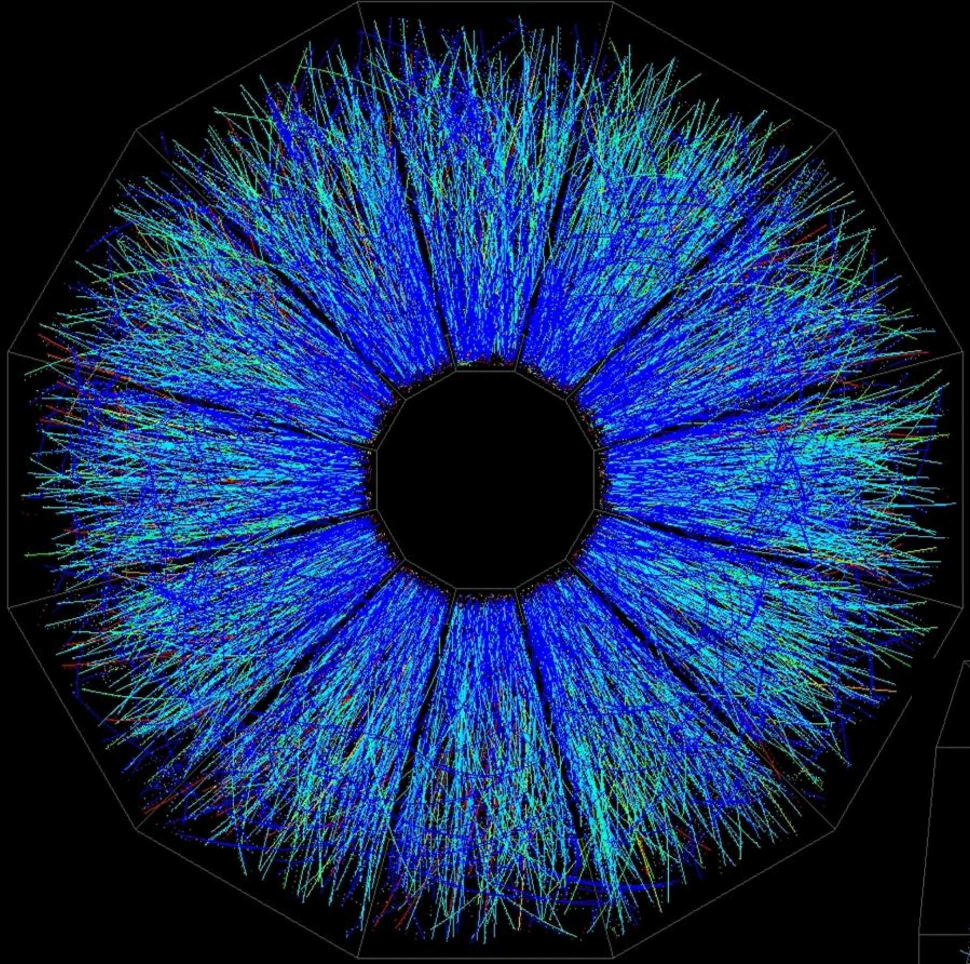


Mid-Central Event

From real-time Level 3 display.

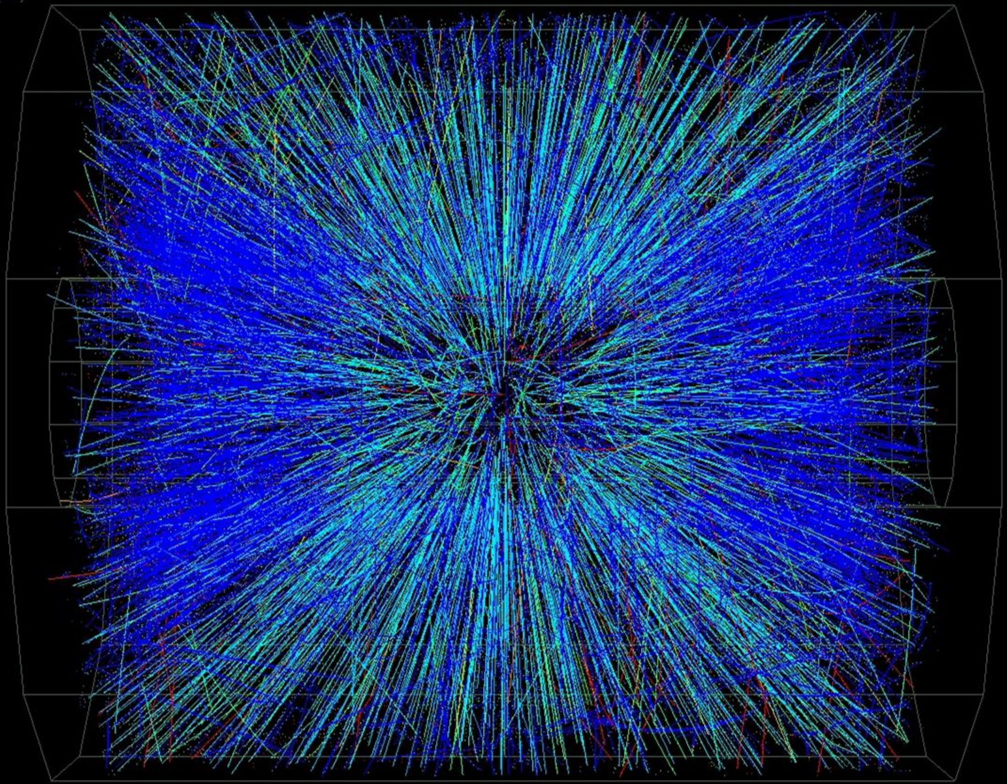


STAR



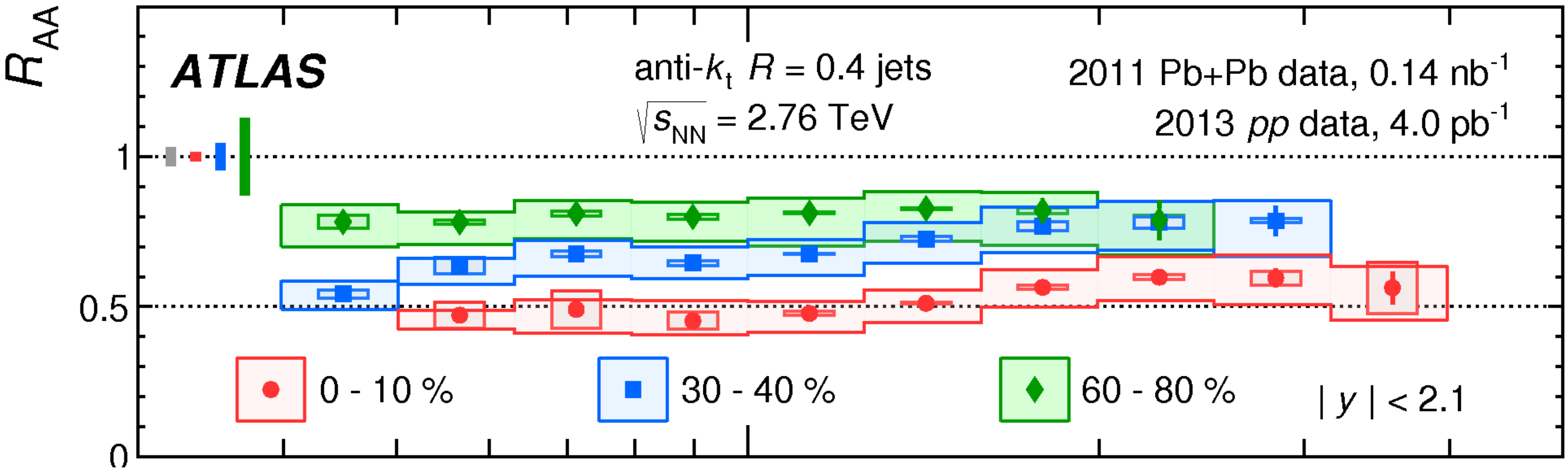
Central Event

From real-time Level 3 display.



STAR

The nuclear modification factor R_{AA}



$$R_{AA} = \frac{d^2N^{AA}/dp_T d\eta}{\langle T_{AA}\rangle d^2\sigma^{pp}/dp_T d\eta}$$

$<T_{AA}>\sigma^{pp} = <N_{coll}>$

N_{coll} is the # of binary collisions

For perturbative QCD processes:

R_{AA}<1: suppression

R_{AA}=1: no nuclear effects

R_{AA}>1: enhancement

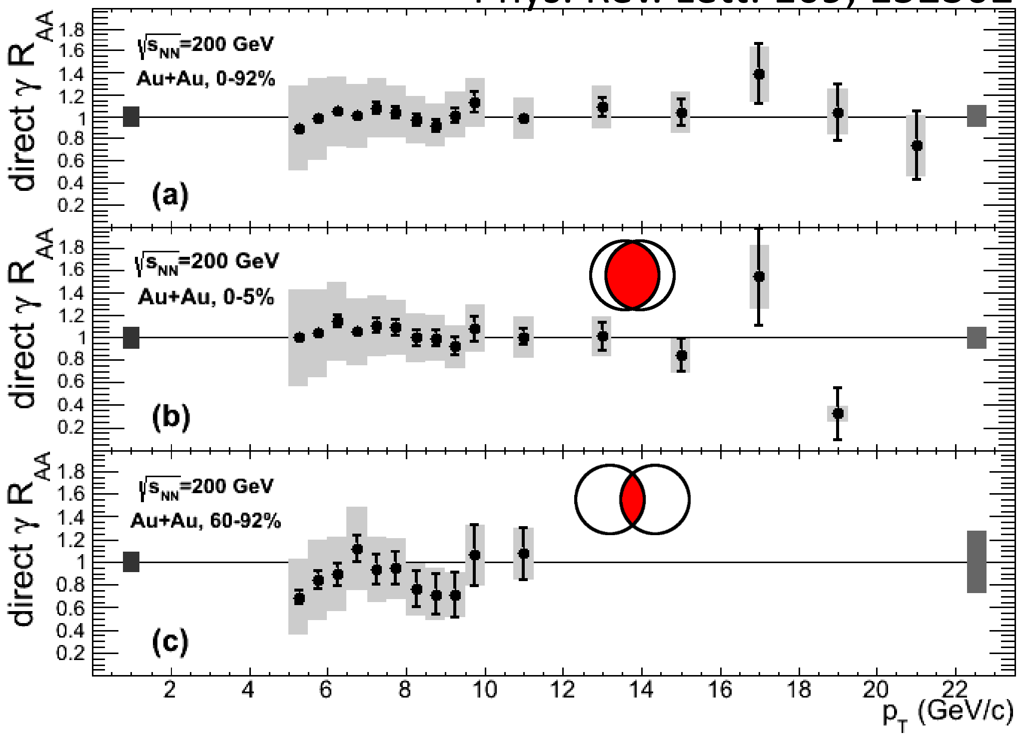
So R_{AA} is a way to quantify the nuclear effects!

Standard candles

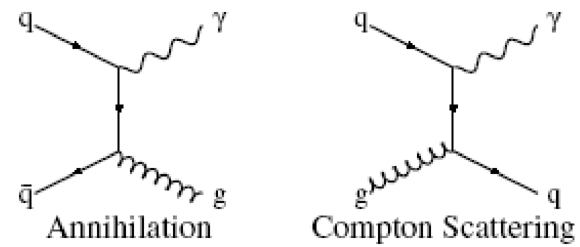
We need penetrating hard probes to validate the binary scaling

Ncoll scaling for direct photons

Phys. Rev. Lett. 109, 152302



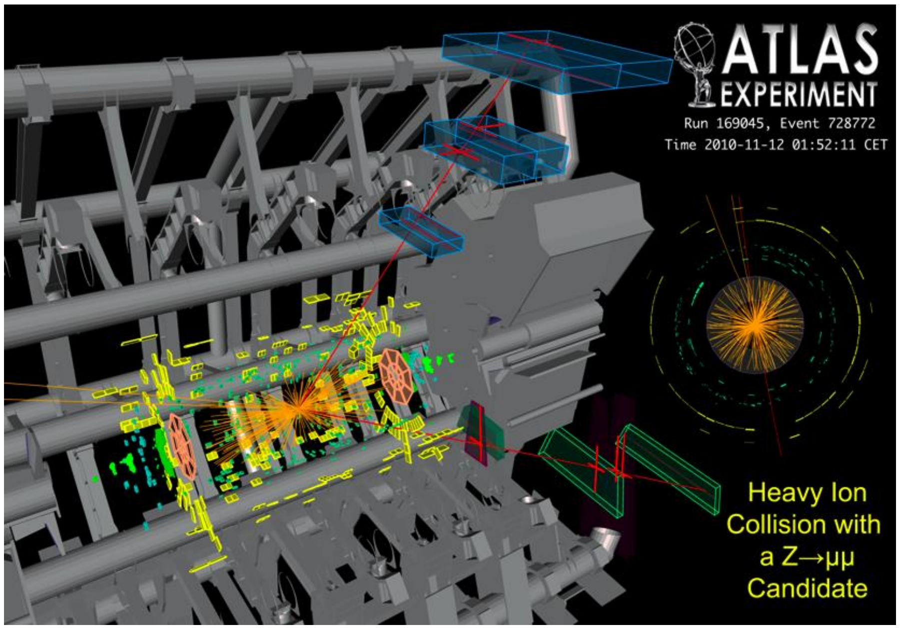
Source of direct photons

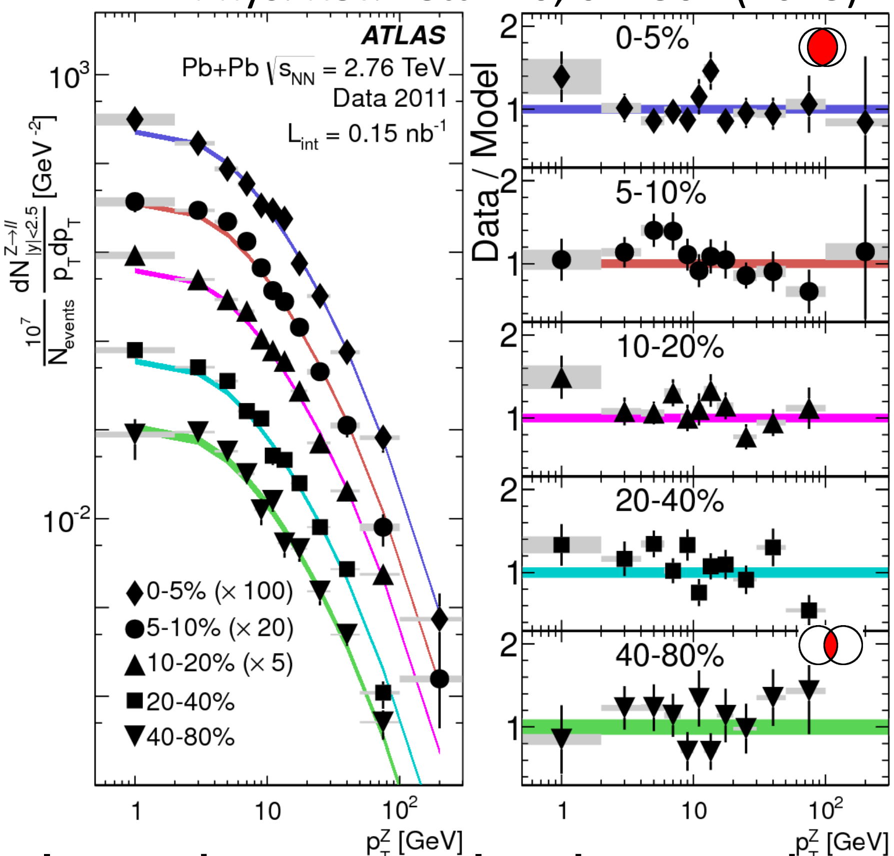


- Direct photons does not interact with final state hadronic matter and the results confirm binary scaling of hard processes!
- The best standard candle at RHIC

New "standard candle" at LHC: ATLAS measures Z bosons

Phys. Rev. Lett 110, 022301 (2013)

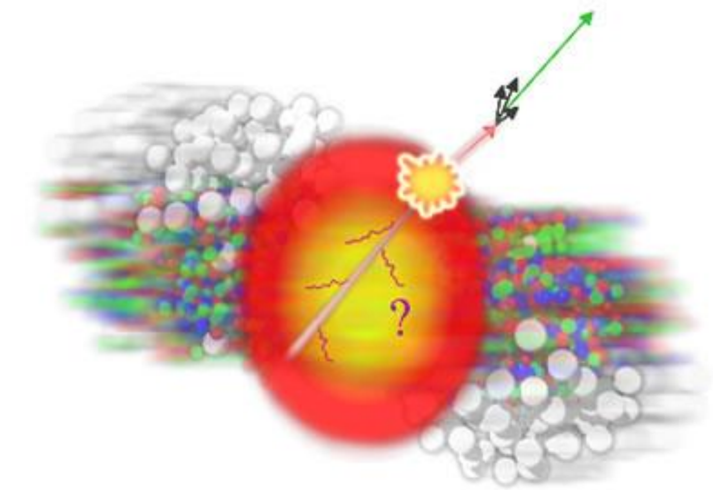




- The Z does not interact strongly and so can also be used to check binary scaling at LHC
- Also the W have been used at LHC

HARD PROBES: JETS

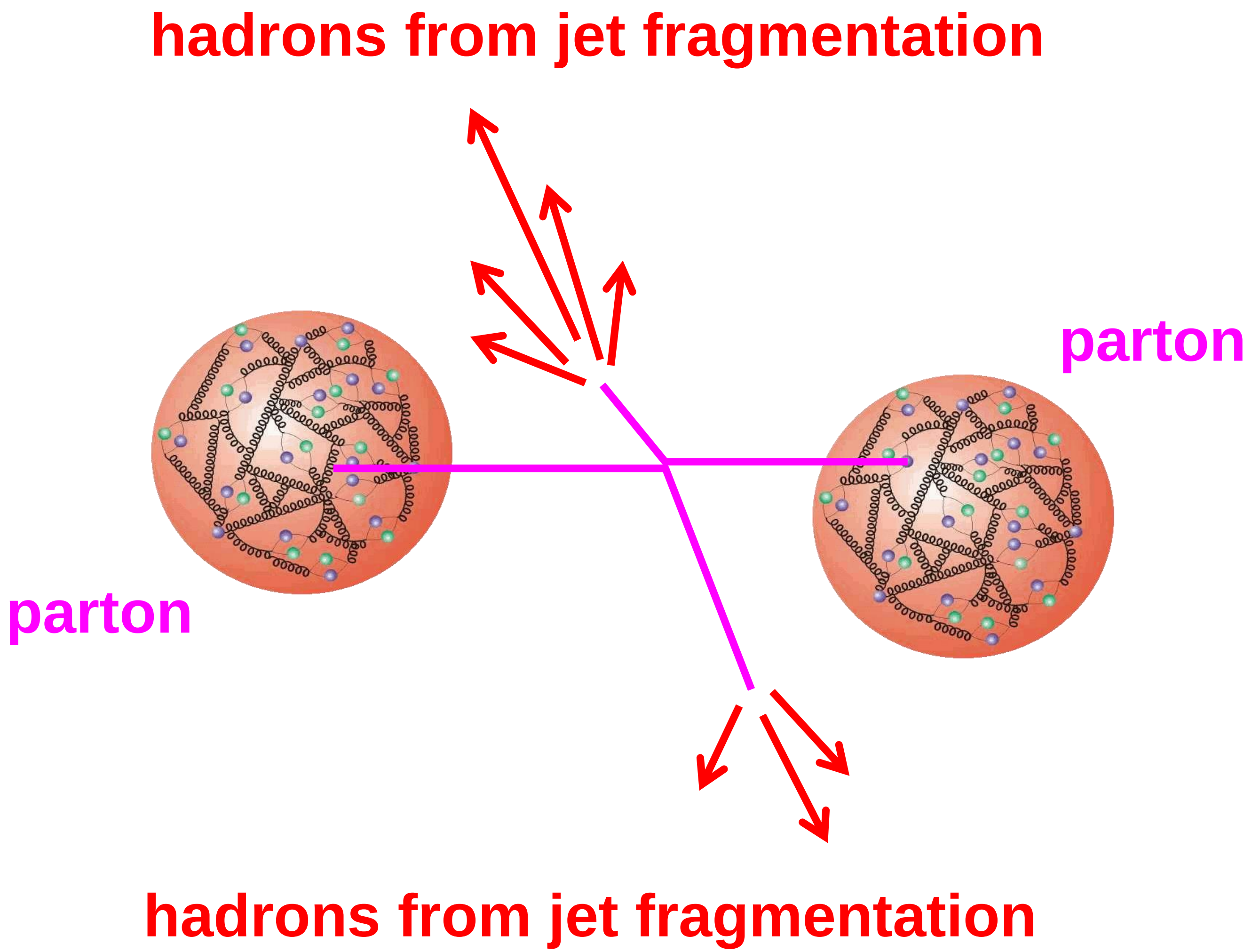
Jets as probes of the QGP



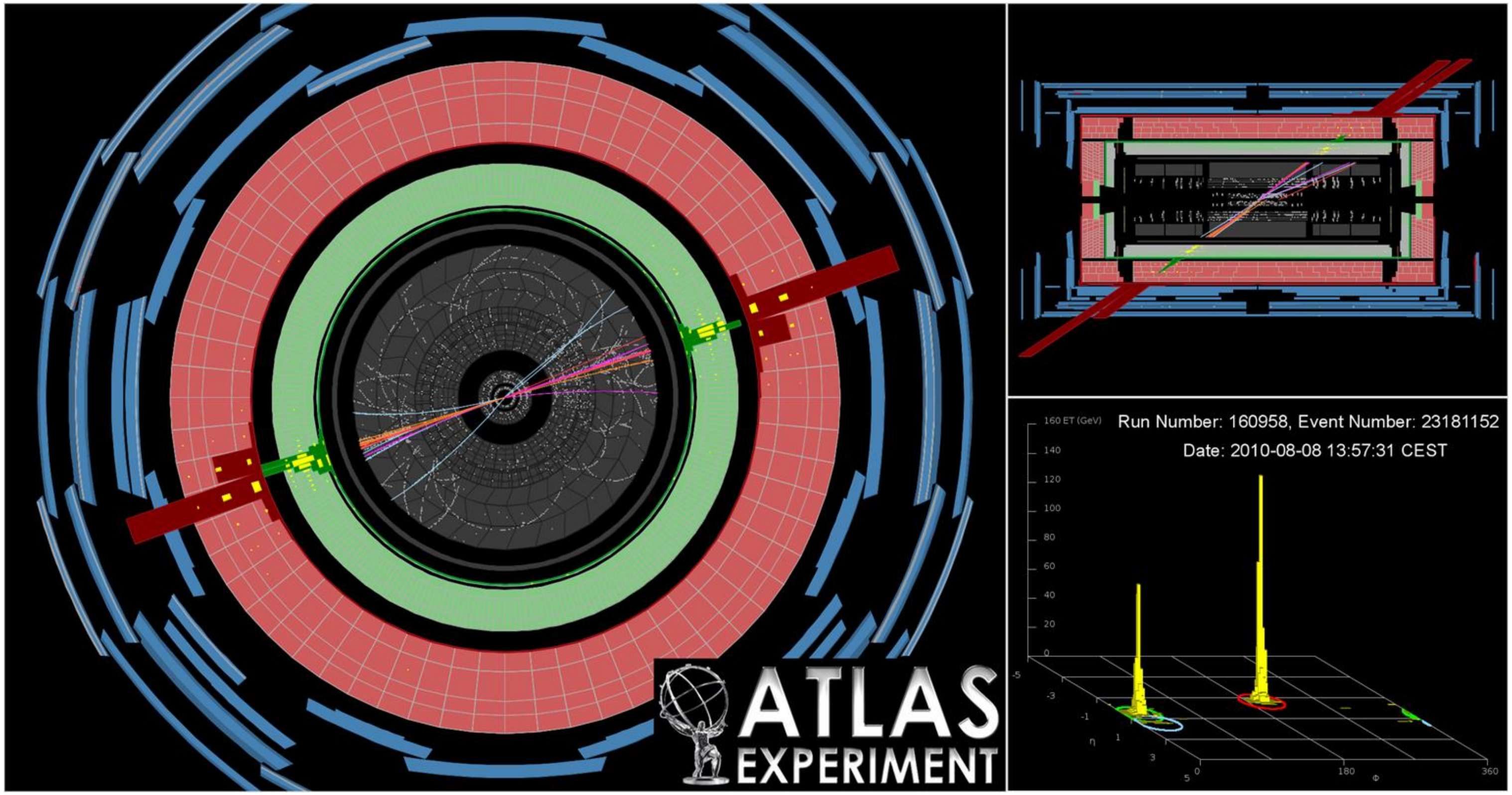
- The idea to use jets to probe the QGP goes back to Bjorken. He made a first study of collisional energy loss but never published this as he realized that radiative energy loss was much more important

- The phenomenology is quite difficult and has many variants so I focus on the experimental results

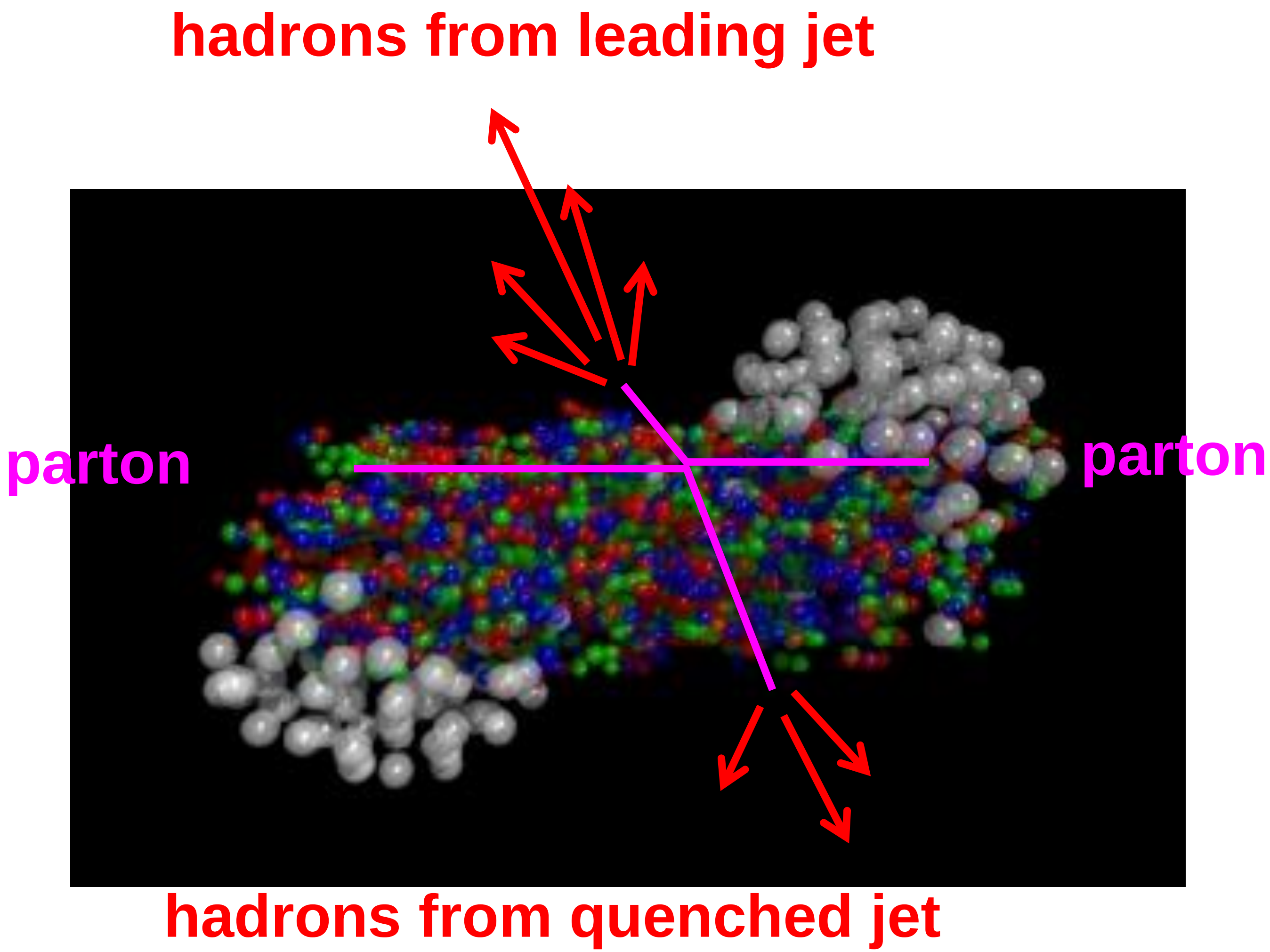
Jets in pp


hadrons from jet fragmentation
parton
parton
hadrons from jet fragmentation

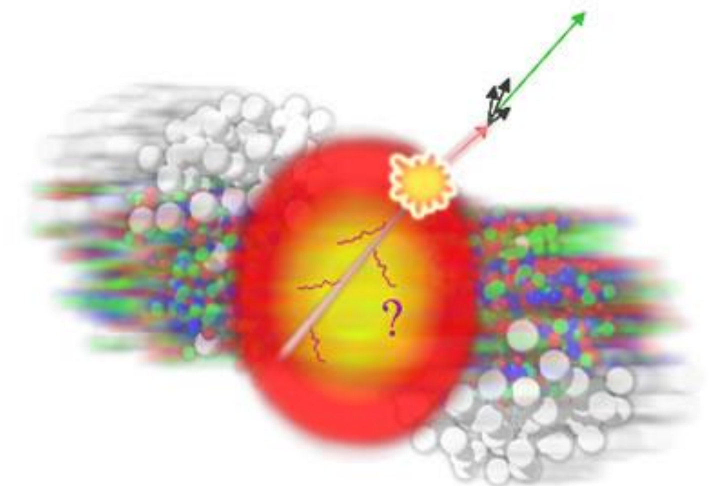
Jets in pp


ATLAS
EXPERIMENT
Run Number: 160958, Event Number: 23181152
Date: 2010-08-08 13:57:31 CEST
160 ET (GeV)

Jets in Pb-Pb

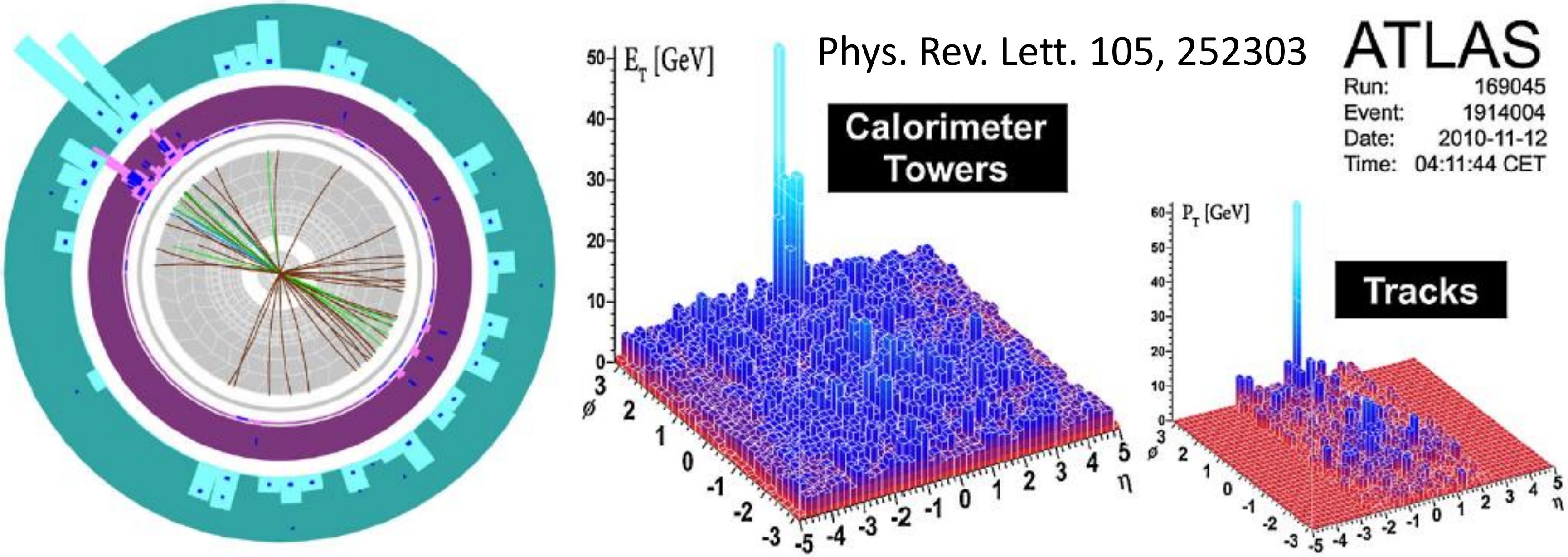

hadrons from leading jet
parton
parton
hadrons from quenched jet

Jet quenching in the QGP

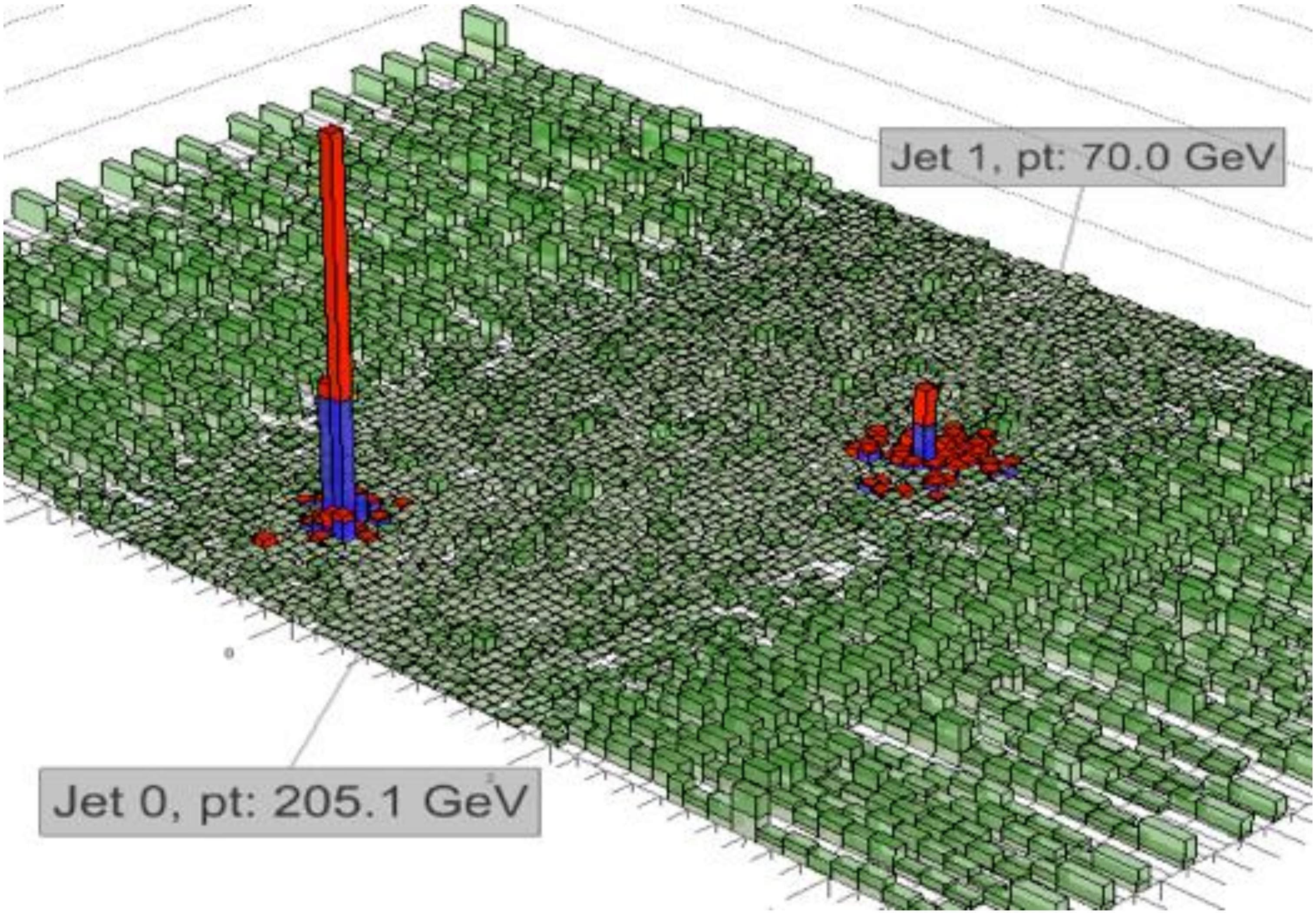


For color charges one expects energy losses of GeV/fm vs MeV/cm for electric charges in normal atomic matter. 10,000,000,000,000,000 (10^{16}) times larger energy loss.

Phys. Rev. Lett. 105, 252303

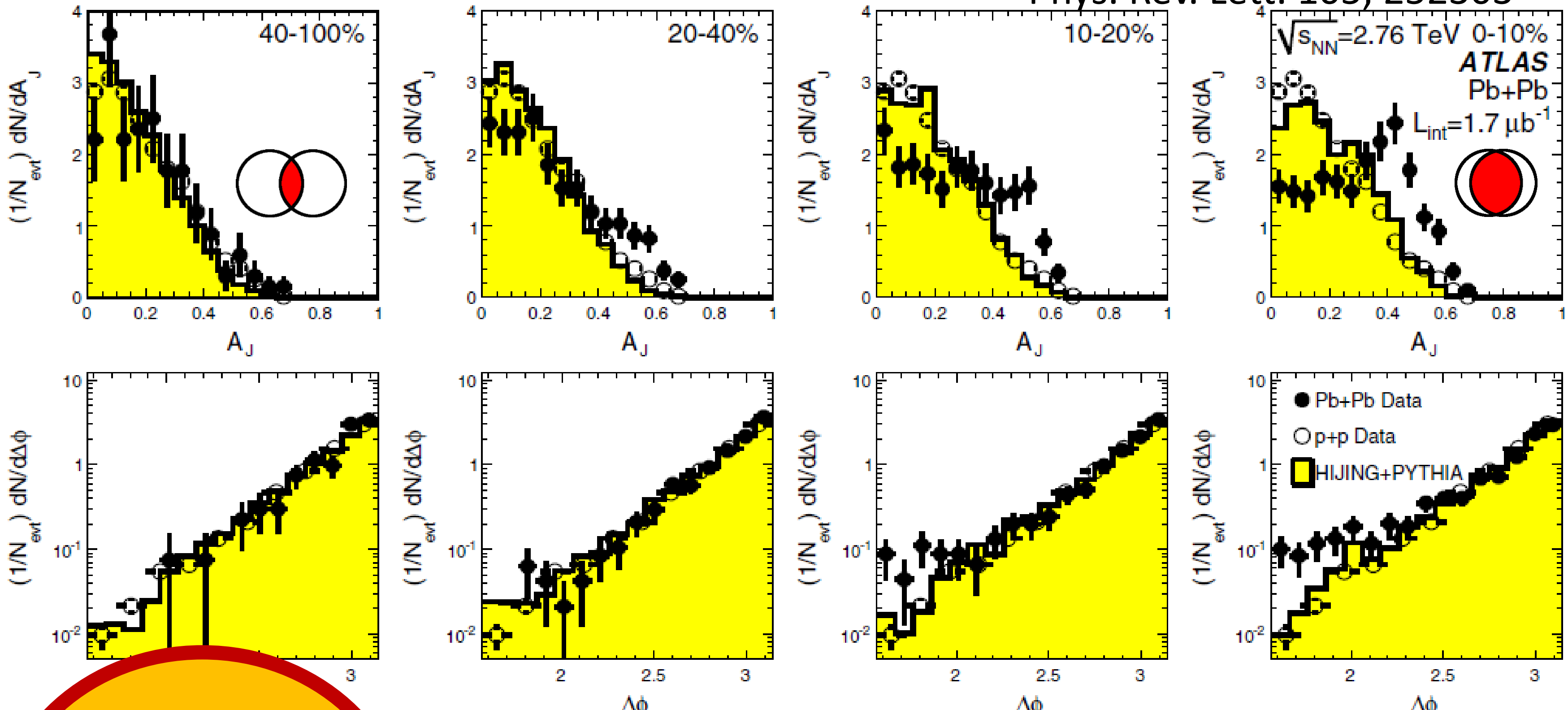


Jets in Pb-Pb (CMS)



Quantifying the dijet asymmetry

Phys. Rev. Lett. 105, 252303



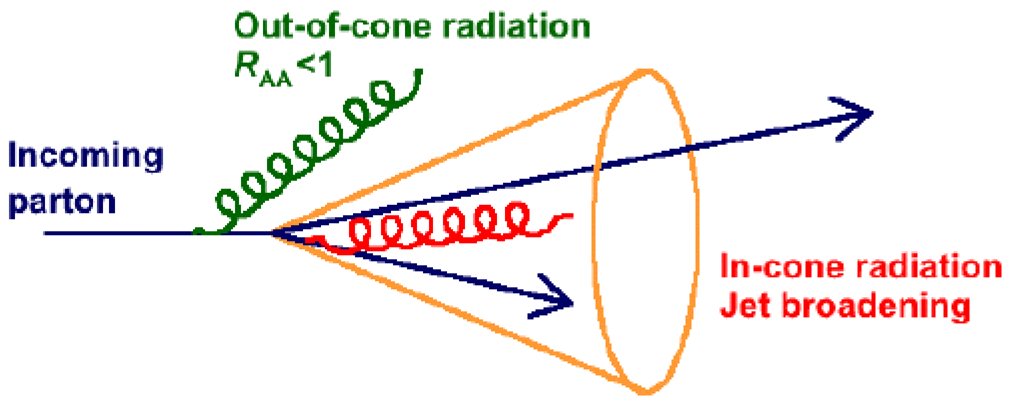
COLORED

$$A_J = \frac{E_{T1} - E_{T2}}{E_{T1} + E_{T2}}$$

Where E_{T1} (E_{T2}) is the transverse energy of the leading (subleading) jet (E_{T1}>100 GeV and E_{T2}>25 GeV).

Notice that the jets are still back-to-back!

Looking into the jet structure

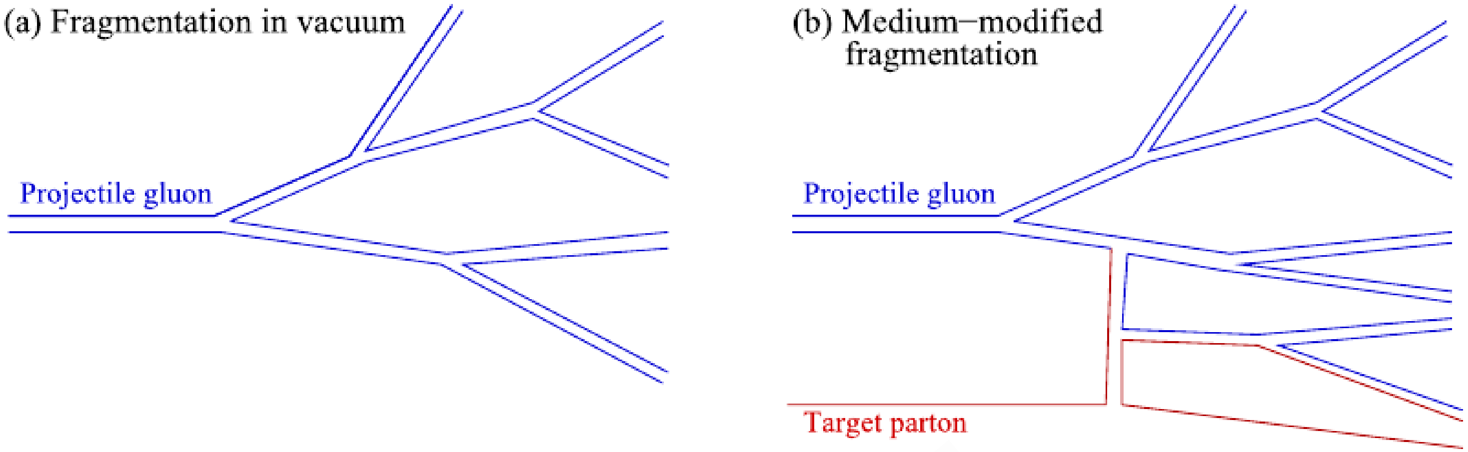


The motivation for these studies was:

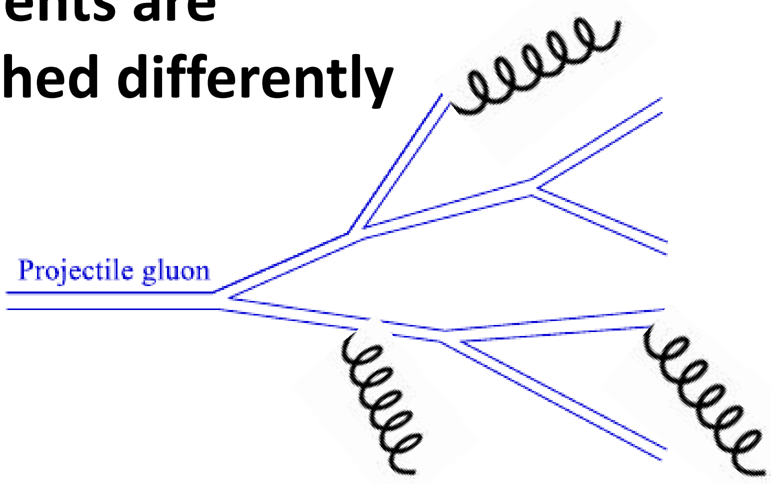
- to recover some of the radiated energy
  (in principle jets did not have to suppressed!)
- to study how the FF is modified

Two examples of ideas for modified FFs

Direct interaction with medium

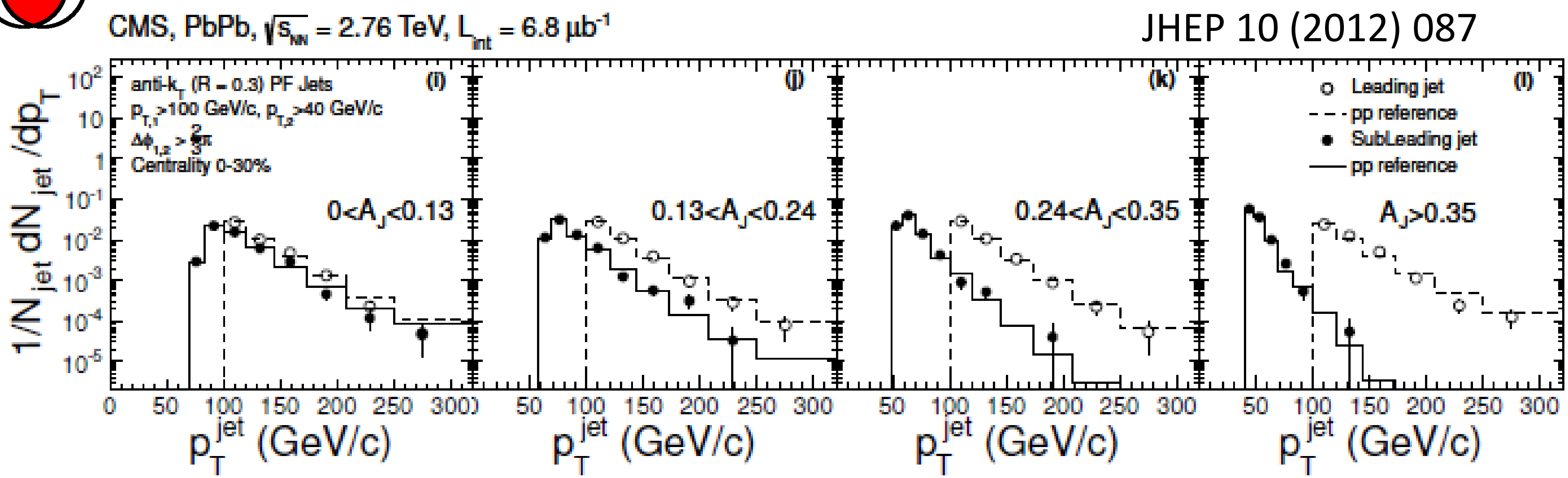


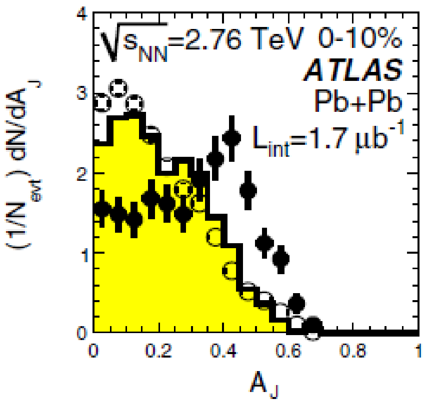
Fragments are quenched differently



Selecting dijet events and comparing to pp

JHEP 10 (2012) 087



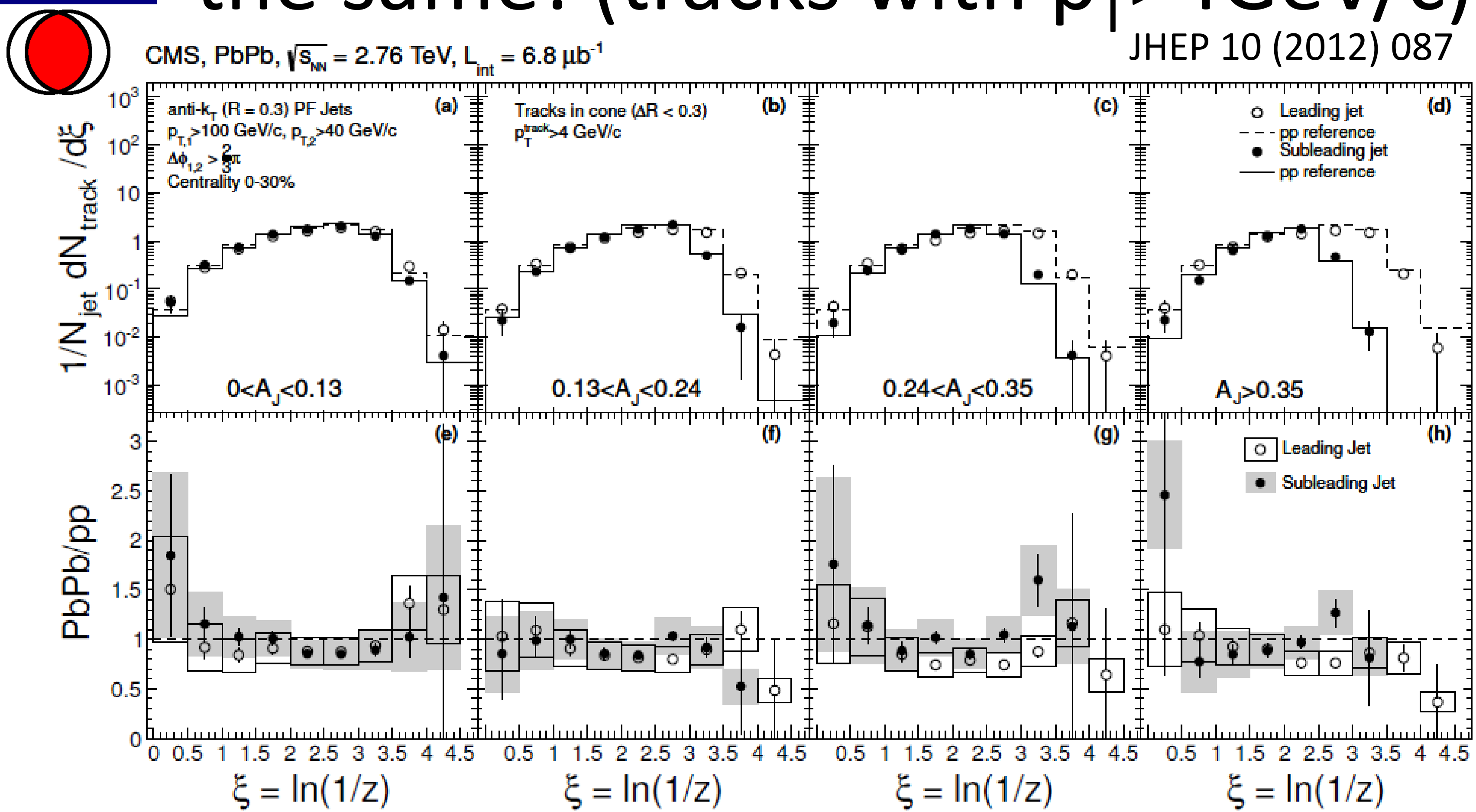


Tracks with p_T>4GeV/c.

The A_J selection introduces the same bias on the dijet samples

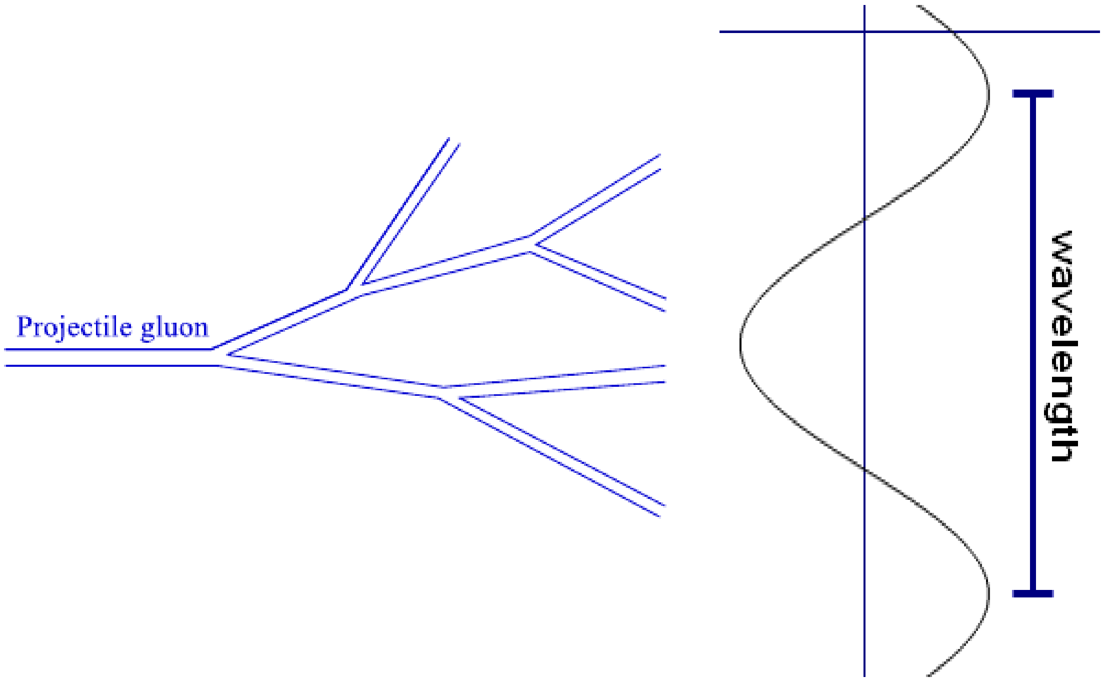
Surprisingly the jet structure is the same! (tracks with p_T>4GeV/c)

JHEP 10 (2012) 087



The result shows that quenched jets looks like pp (vacuum) jets!
Even in the case where A_J is large and for the subleading jet!

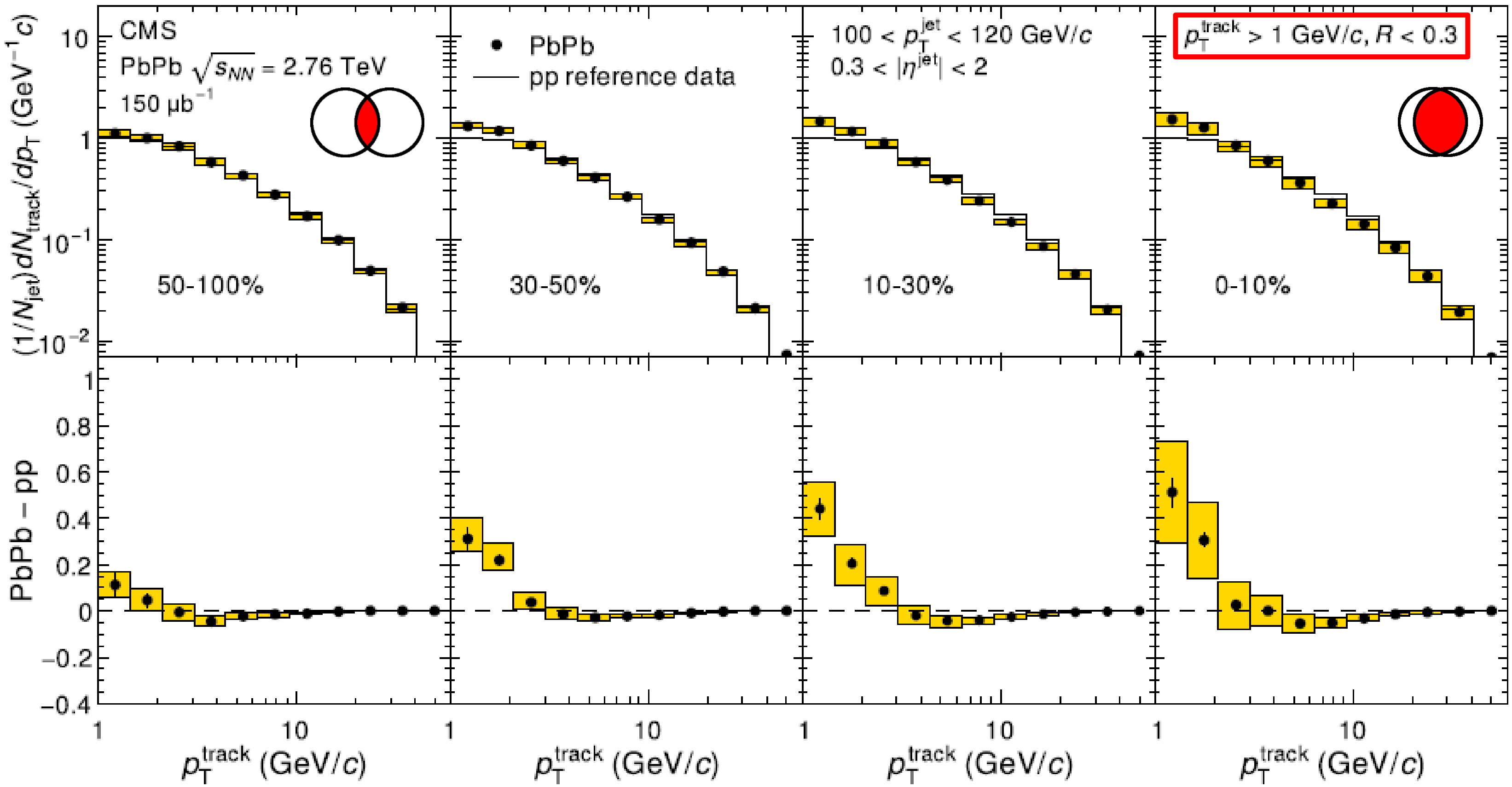
The results show that quenching is coherent!?



One proposed solution is that the medium cannot resolve the jet constituents (Phys.Lett. B725 (2013) 357-360).
So for the medium the jet will look essentially as a single parton and so all fragments are quenched coherently!

The finer structure of the fragmentation function

Phys. Rev. C 90, 024908

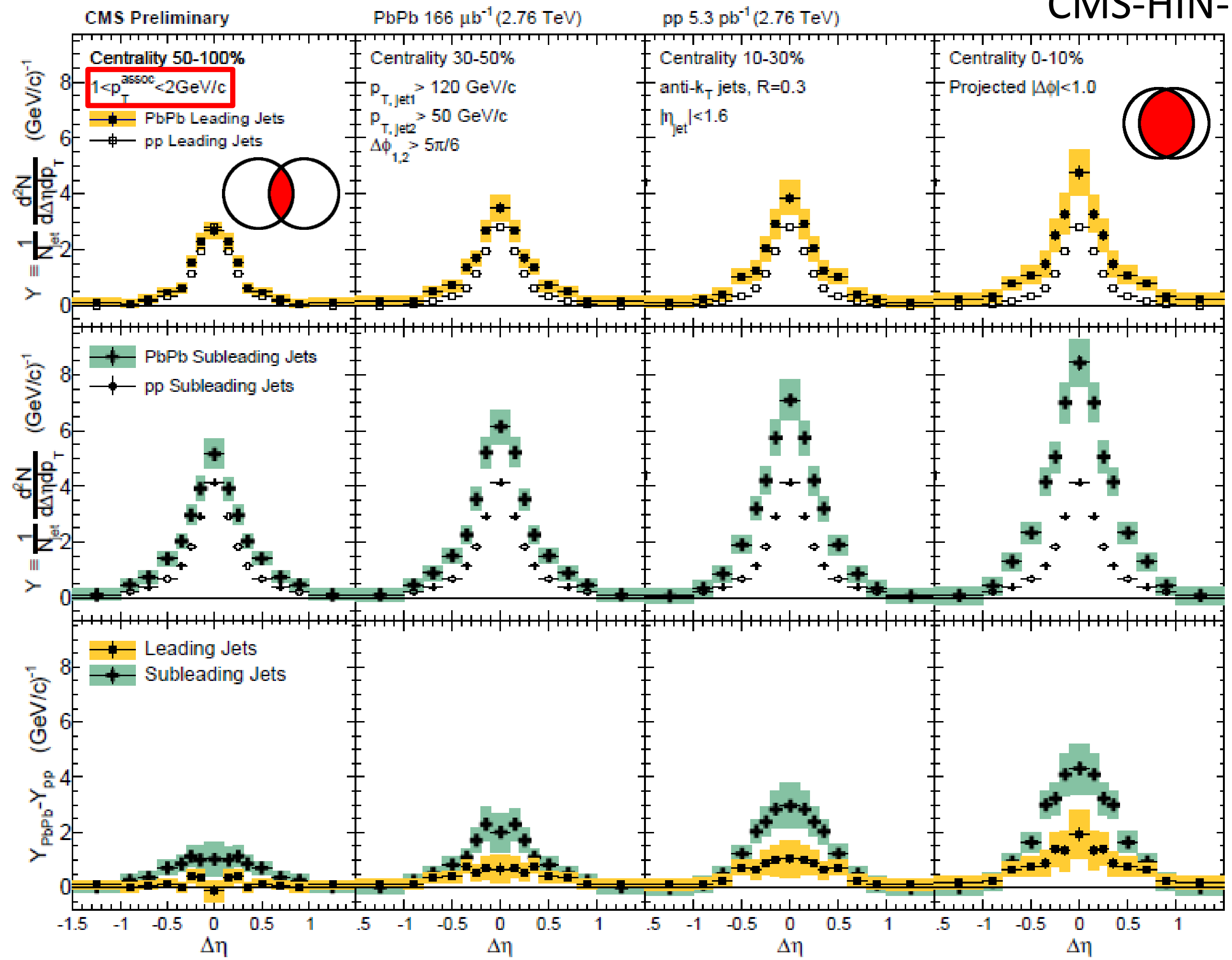


This analysis considers tracks down to p_T = 1 GeV/c

The modifications are sitting at low p_T, mainly p_T<3GeV/c

Tracking the energy loss

CMS-HIN-14-016



CMS have shown that the energy loss can be recovered at low p_T at large Δη and large Δϕ

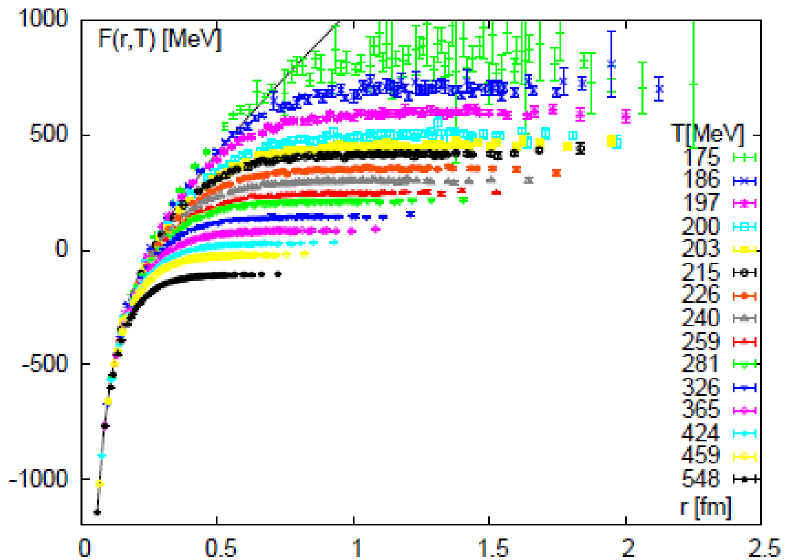
Conclusions on jets

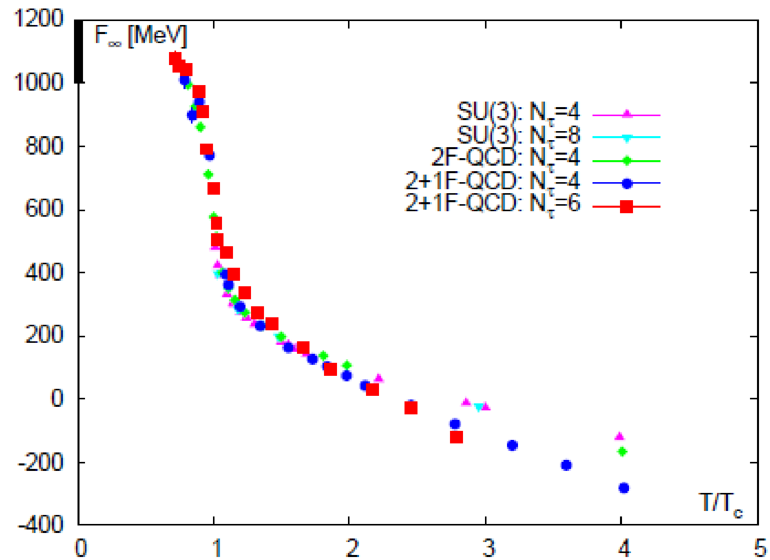
- High p_T particles (and jets) were one of the hottest topics at RHIC where they were difficult to measure
  - And it seemed at RHIC that jets were modified in a spectacular way giving rise to exotic effects
- At LHC the picture we have of jet quenching has turned out to be surprisingly simple
  - Experimentally the jets seem to first loose energy in the QGP and then afterwards to fragment as vacuum jets
    The energy loss is observed as soft particles at large angles

HARD PROBES: HEAVY QUARKS

Heavy quarks:
I will only focus on Quarkonia

Lattice QCD results for the heavy quark potential (free energy): arXiv:0710.0498

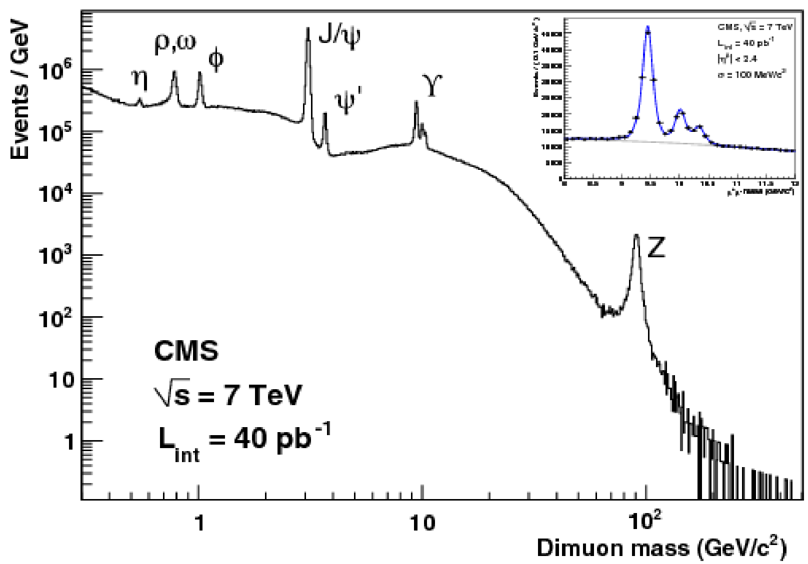




Lattice QCD predicts that the long range force will be screened in the plasma
Bound state properties will change and some states will disappear / melt

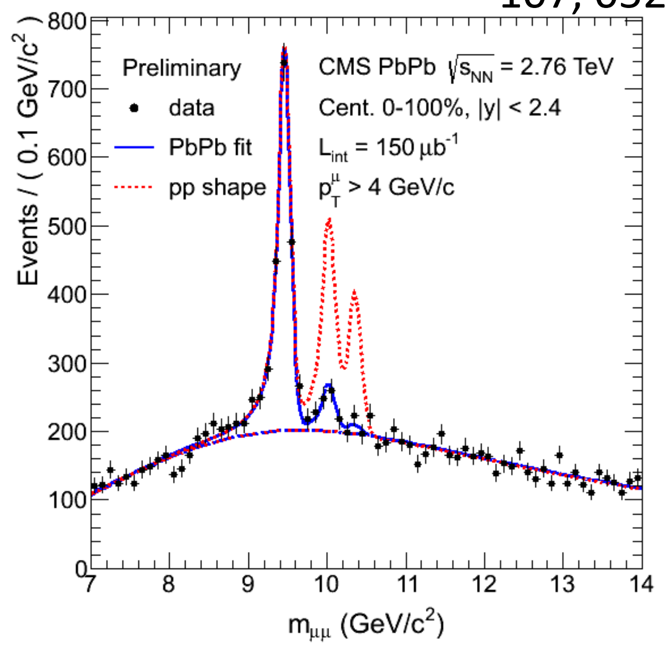
Original idea:
"J/ψ Suppression by Quark-Gluon Plasma Formation" by T. Matsui and H. Satz
Phys.Lett. B178 (1986) 416
Cited by 2108 records

LHC has delivered



Υ: 1S 2S 3S

Published in PRL 107, 052302

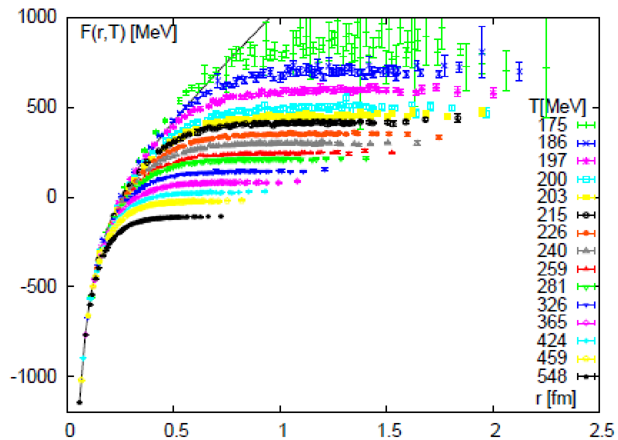


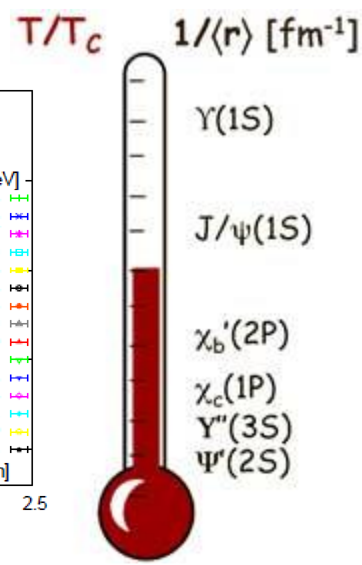
CMS has proven to be a marvelous detector for bottomonium
ALICE can complement by going to lower p_T for charmonium

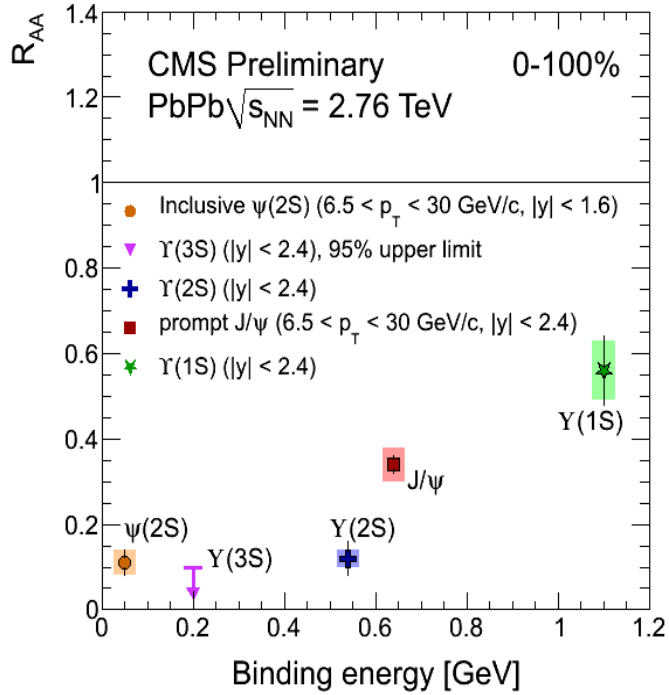
These textbook results impressed at Quark Matter 2011

Suppression of heavy quarkonia can work as a thermometer

Note: $6.5<p_T<30$ GeV for J/ψ and ψ(2s)

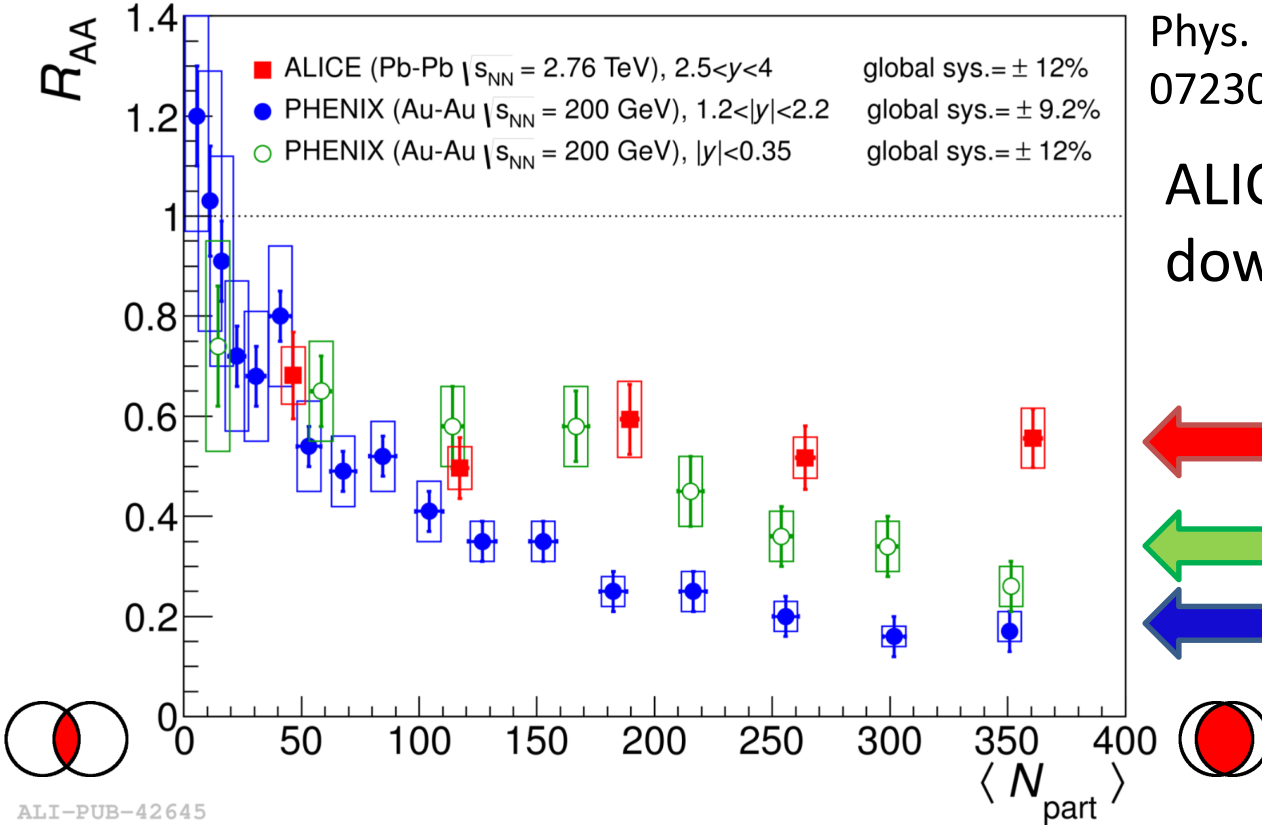






Suppression qualitatively depends on binding energy as predicted

J/ψ results from ALICE: Evidence for regeneration



Phys. Rev. Lett. 109, 072301 (2012)

ALICE can go down to p_T~0.

LHC

RHIC (mid)

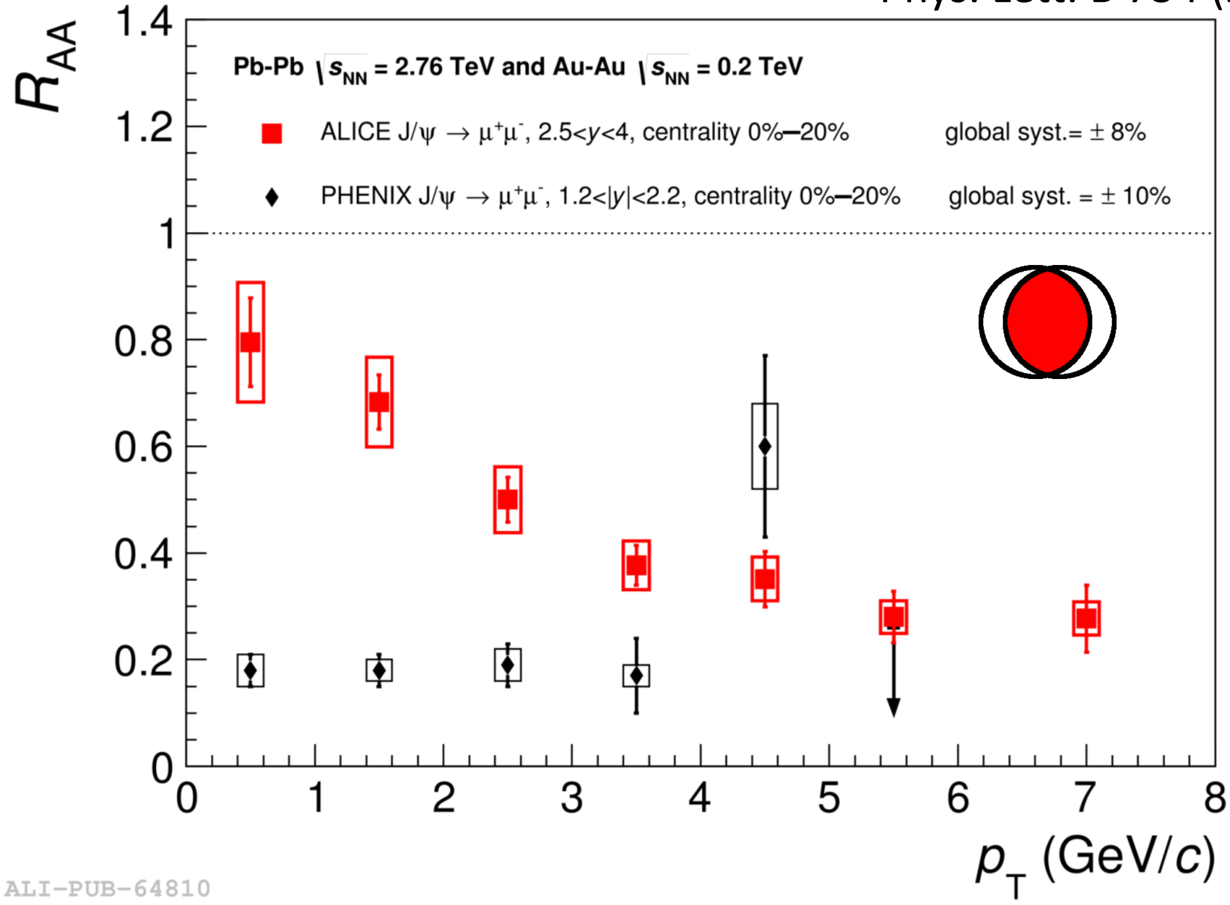
RHIC (fwd)

The suppression of J/ψ at LHC is less than at RHIC!

Understanding: due to the large cross section "random" charm and anti-charm quarks combine (see next slides)

J/Psi differential results

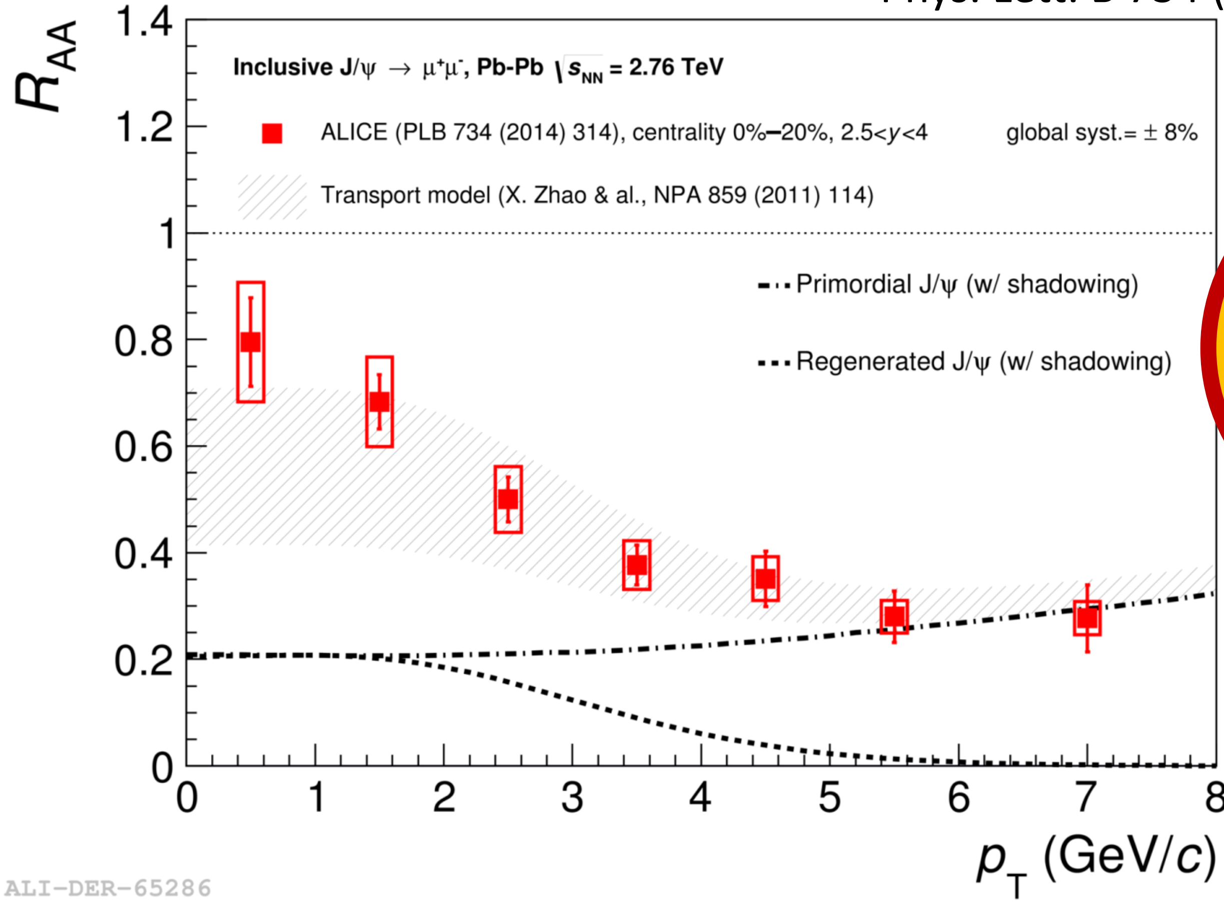
Phys. Lett. B 734 (2014) 314-327



The difference in suppression is visible at low p_T

J/Psi differential results

Phys. Lett. B 734 (2014) 314-327

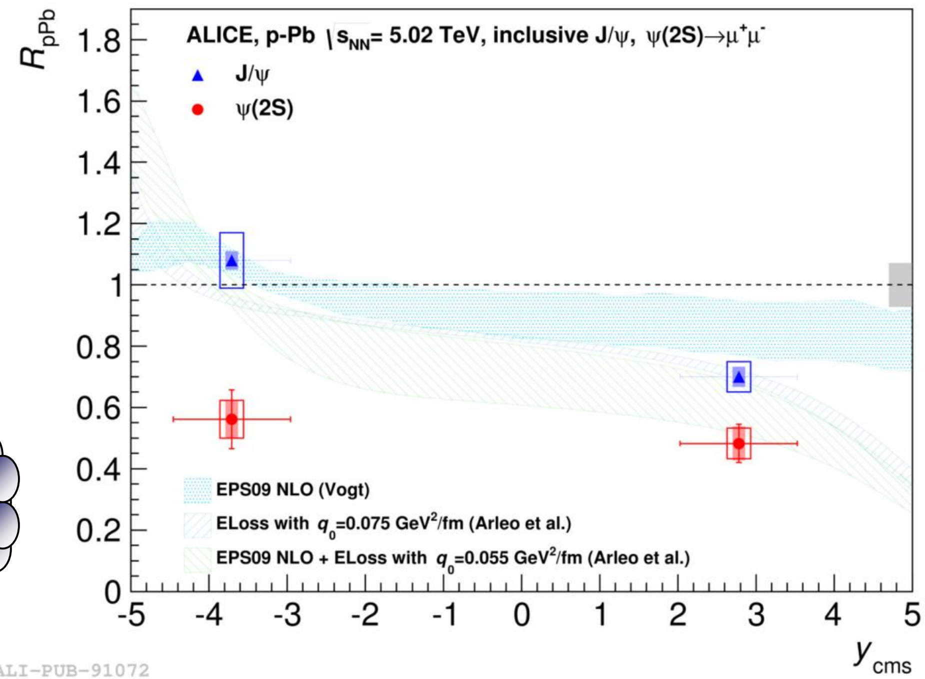


DECONFINED

The difference can be explained by J/ψ regeneration (much larger charm Xsection at LHC)

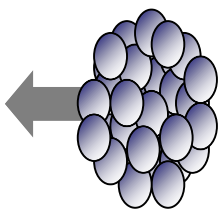
(This effect was already proposed before the RHIC results came out.)

Caveat: cold nuclear matter effects are not understood



JHEP12(2014)073

lead rapidity



proton rapidity

For J/ψ we have reasonbale good description: energy loss really just means that due to collisions with other nucleons the rapidity of the J/ψ is shifted towards the Pb-nuclei.

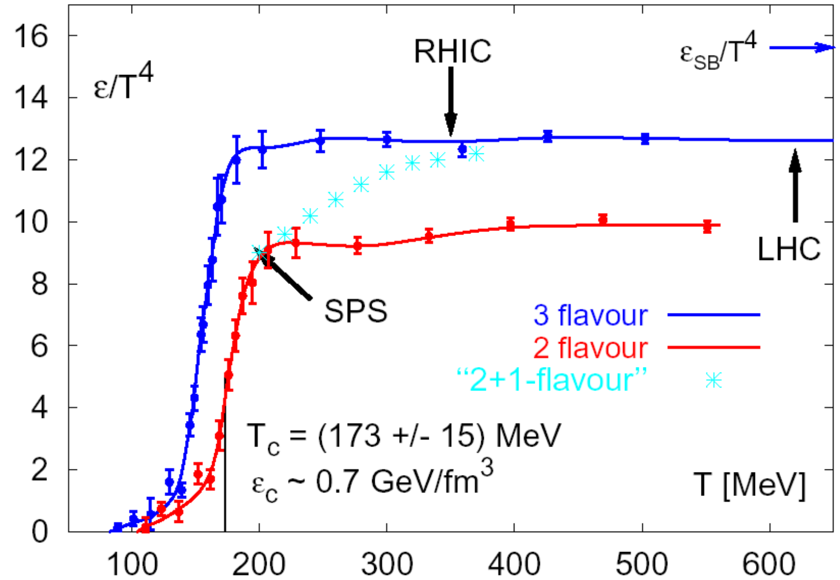
But why is the 2S state extremely suppressed in p-Pb (at the rapidity) where the 1S state is a bit enhanced?

Summary of quarkonia results

- The qualitative results are in agreement with what one expects for a Quark Gluon Plasma
- The advantage of LHC is that all quarkonia states have been measured making the measurements less sensitive to assumptions about feed-down corrections
- The quantitative picture is complicated and there are many issues that are even interesting for pp phenomenology

SOFT MEDIUM PROPERTIES

Lattice QCD calculation of the energy density



$$\varepsilon_{Quark-Gluon\ gas} = \frac{\pi^2}{30}\left(2\times 8+\frac{7}{8}2\times 2\times 3\times 3\right)T^4$$

Gluon spin and color

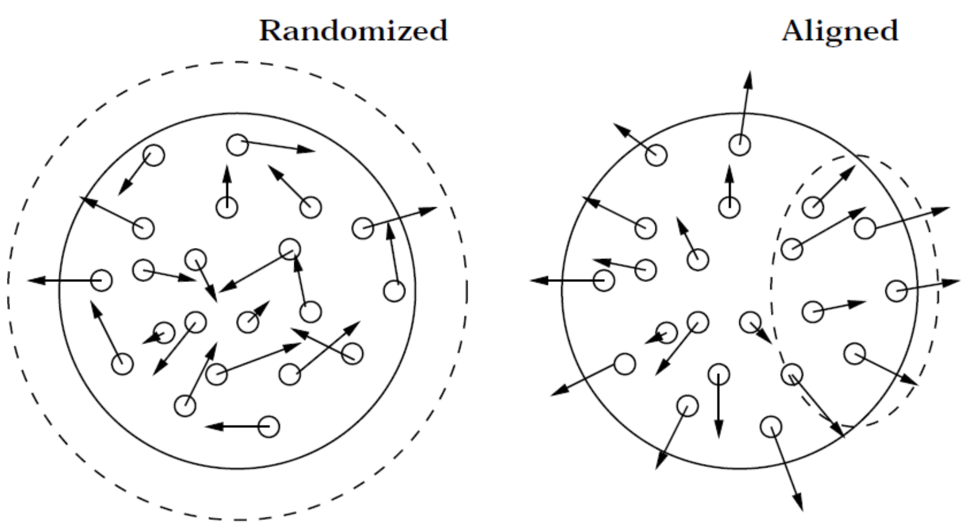
(Anti+)quark spin, color and flavor

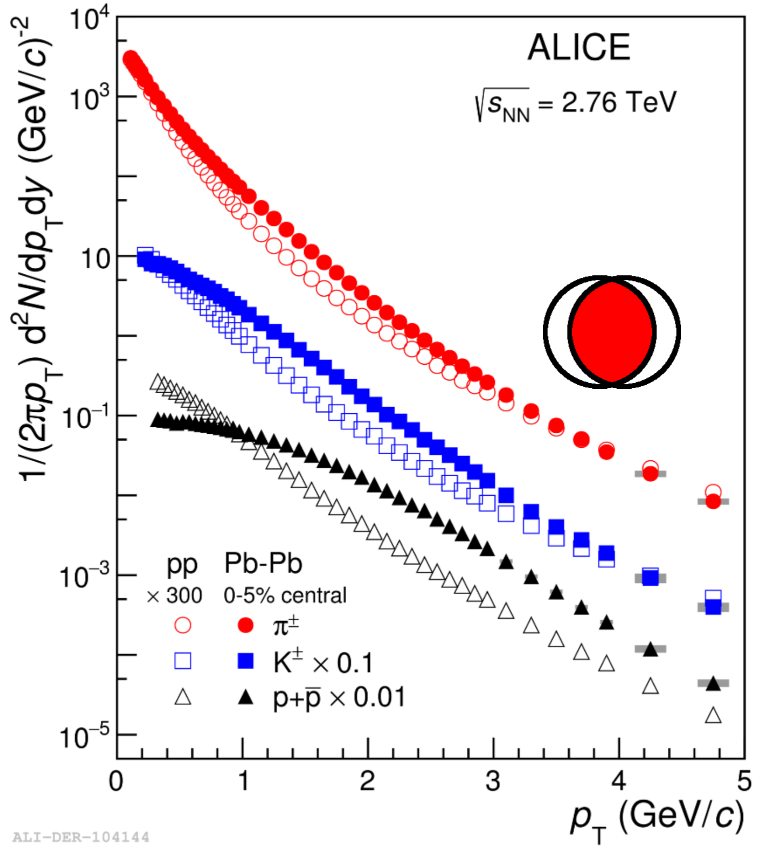
Because of the similarity with Stefan-Boltzmann energy density for a quark-gluon gas:

QGP should be weakly coupled

At a deeper level this is also what we expect from asymptotic freedom

Radial flow



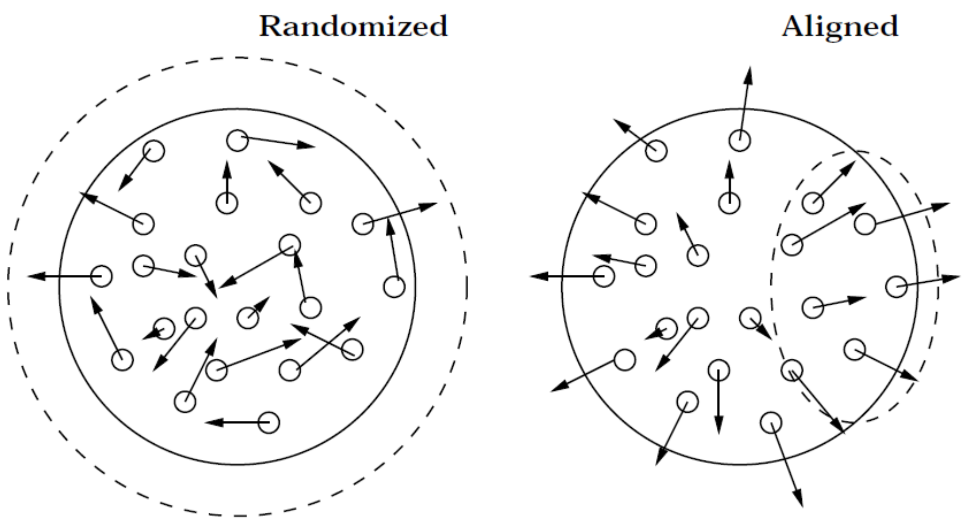


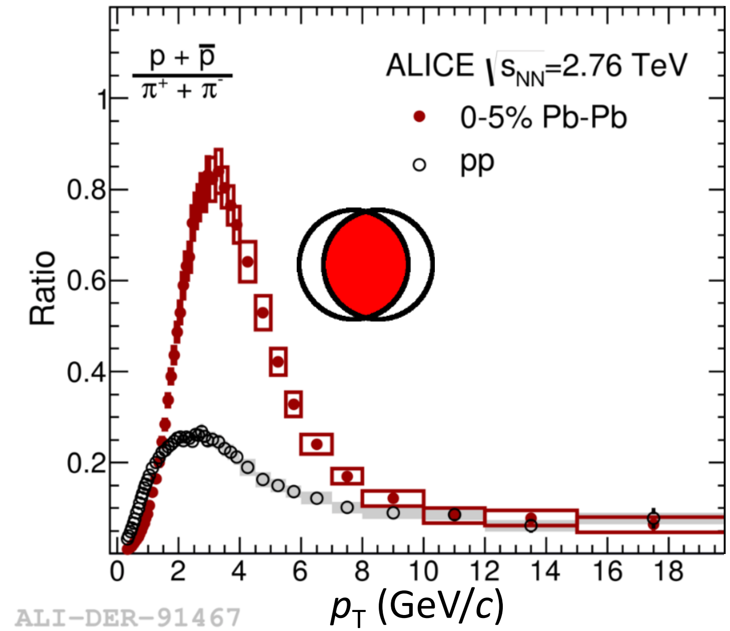
- Flow in general plays a very important role in heavy-ion collisions.
- We believe that flow in the partonic phase is imprinted on the final state hadrons at freeze out.

Flow velocity β_r → mass dependent boost:

$p_T \sim \gamma\beta_r m$ (for particle initially at rest)

Radial flow

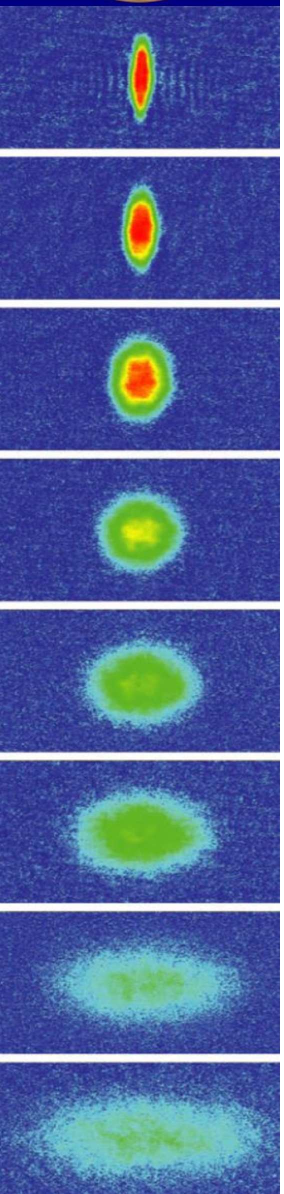


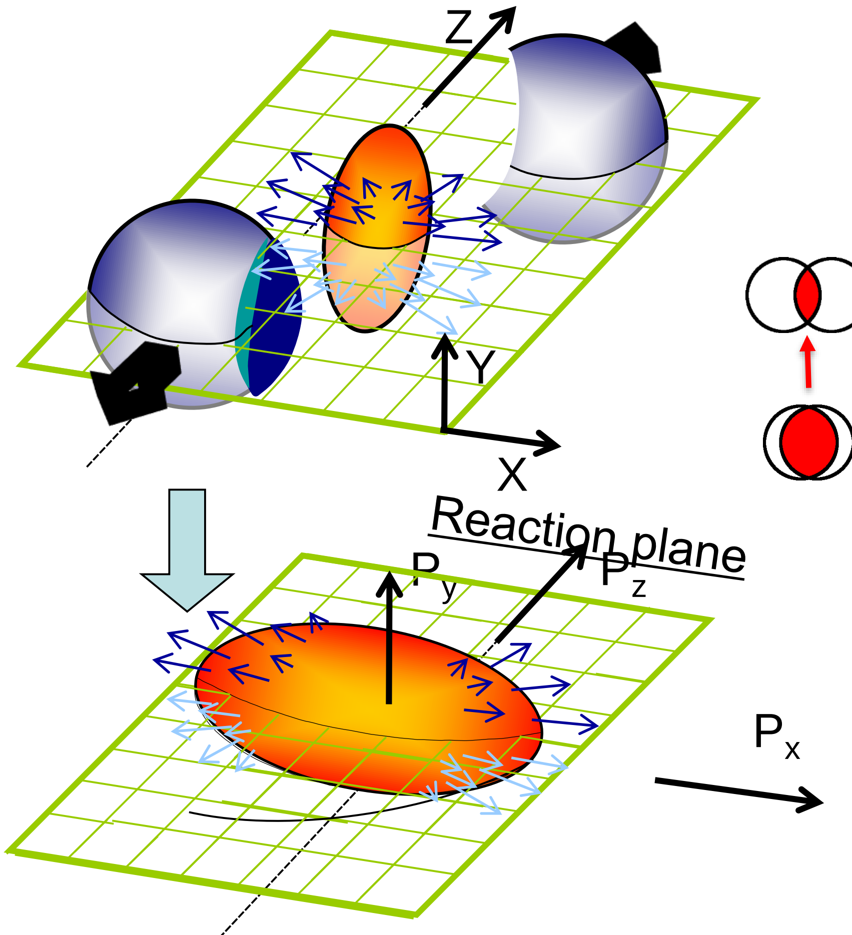


- Flow in general plays a very important role in heavy-ion collisions.
- We believe that flow in the partonic phase is imprinted on the final state hadrons at freeze out.

Flow velocity β_r $\rightarrow$ mass dependent boost:
$p_T \sim \gamma\beta_r m$ (for particle initially at rest)

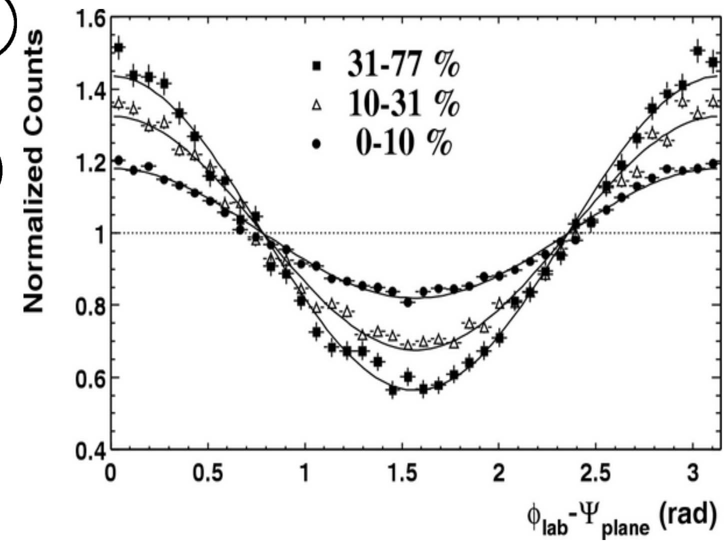
Elliptic flow (v_2)





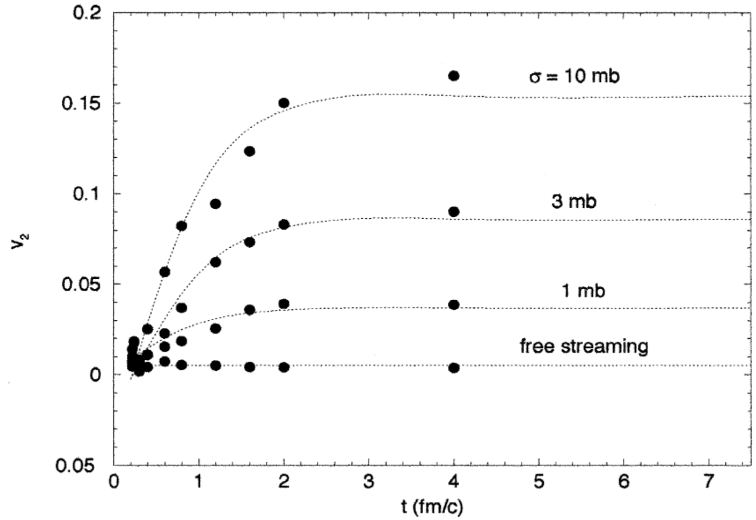
Fourier decomposition:

$dN/d\Delta\phi = 1 + 2\, v_2 \cos(2\, \Delta\phi)$



Initial spatial anisotropy → Strong pressure gradients → v_2 Azimuthal anisotropy → **Sensitivity to early expansion**

Elliptic flow requires early strong interactions to form

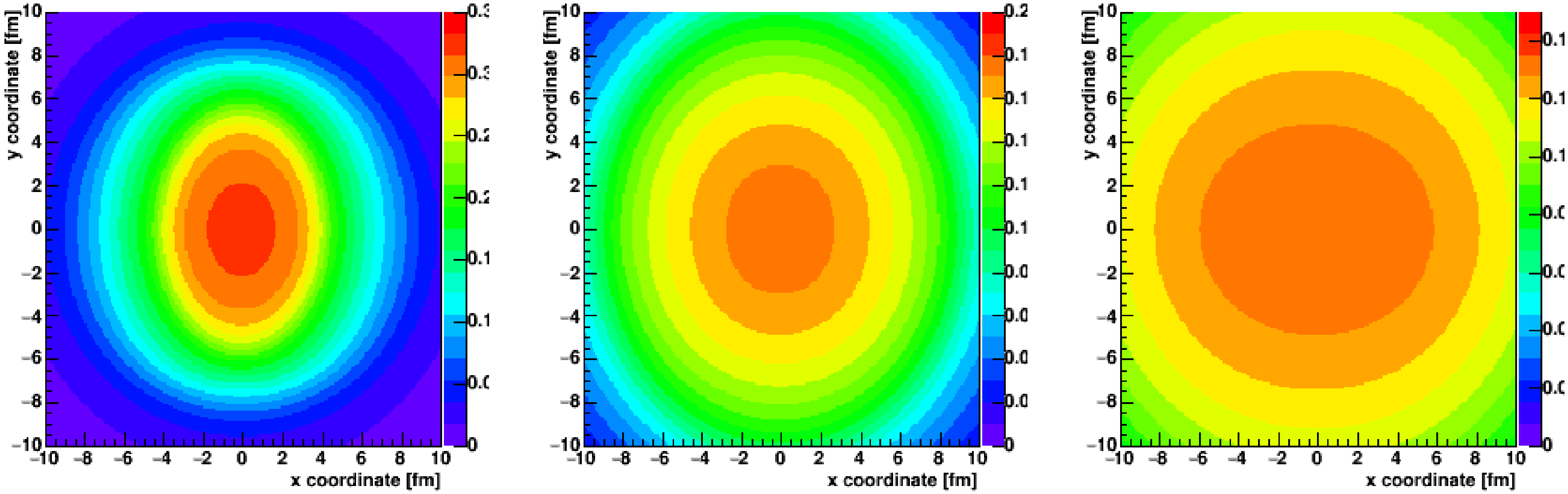


Zhang, Gyulassy, Ko, Phys. Lett. B455 (1999) 45

- Each nucleon-nucleon interaction produces on average a spherical symmetric distribution.
  Only by interacting elliptic flow is generated

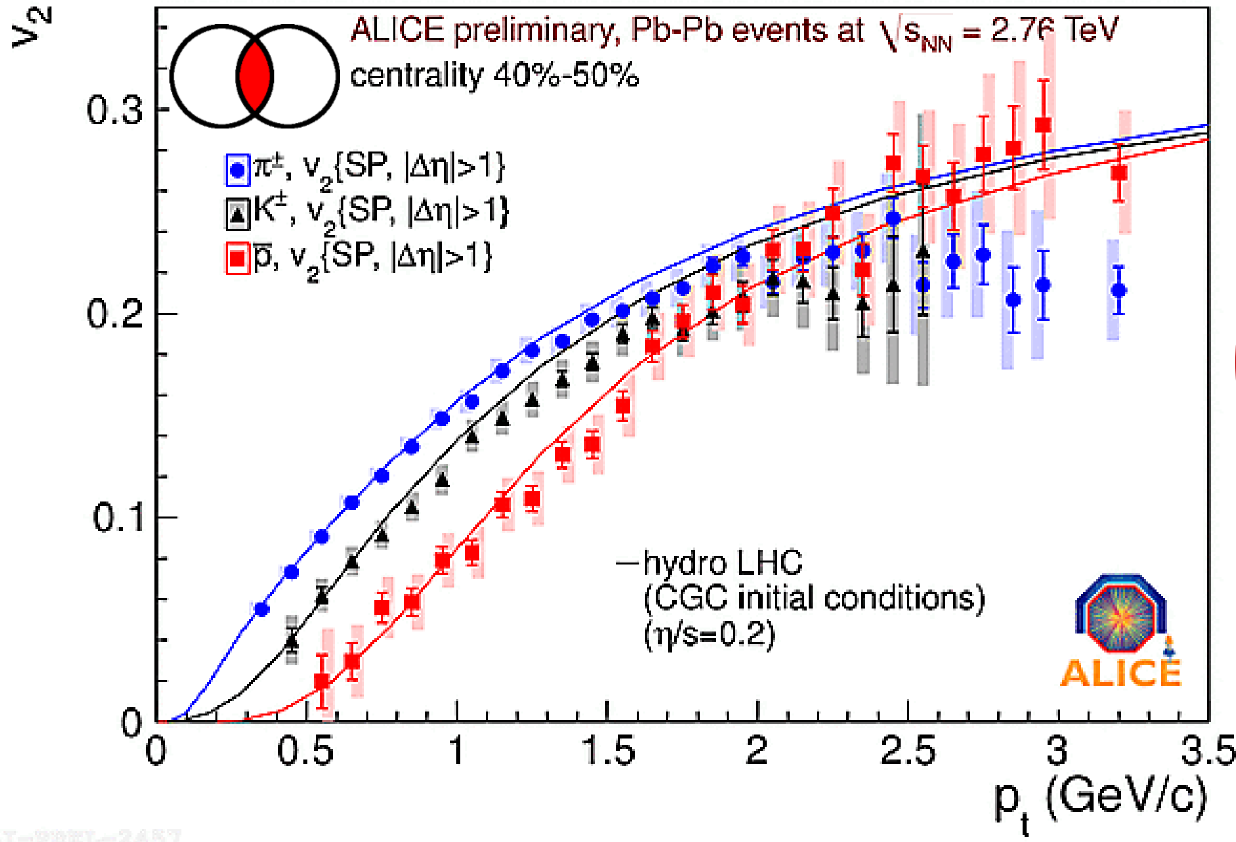
A hydro expansion

t = 1 fm/c t = 6 fm/c t = 12 fm/c



- In hydrodynamics there are no particles. These are first created at freeze-out (when T~160 MeV)
- Calculation by J. Nagle using P. Romatschke's SONIC for Au+Au @ 200 GeV collision with b = 6.5 fm

Mass ordering of the elliptic flow

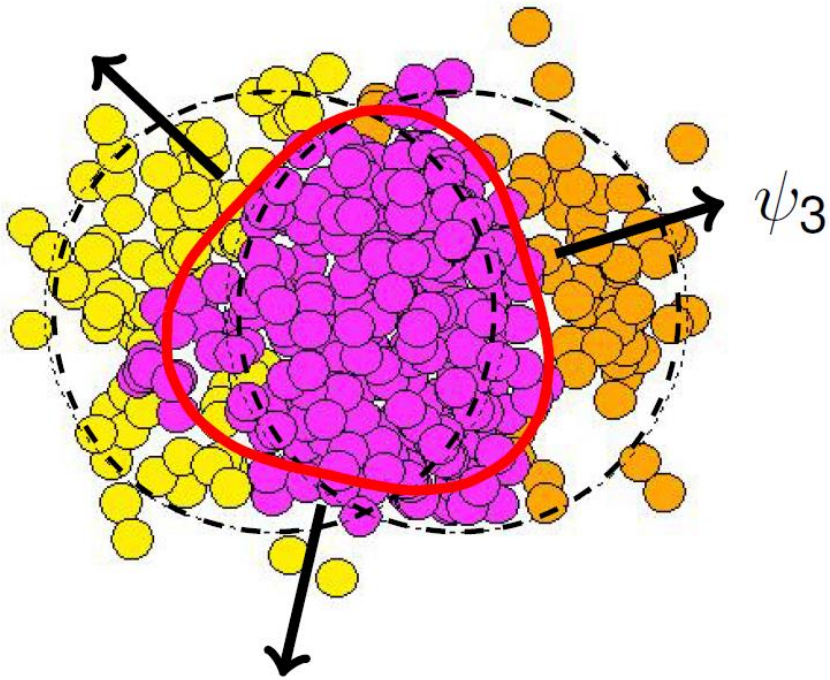


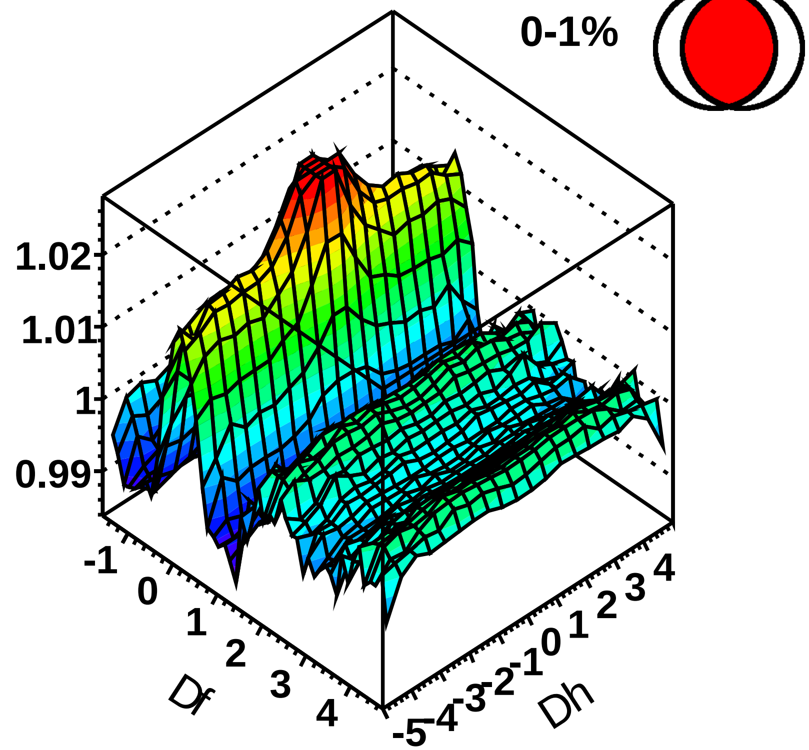
MEDIUM

The mass ordering is characteristic of flow since heavier particles are pushed out to higher p_T

Surprisingly this is well described by nearly ideal hydrodynamics

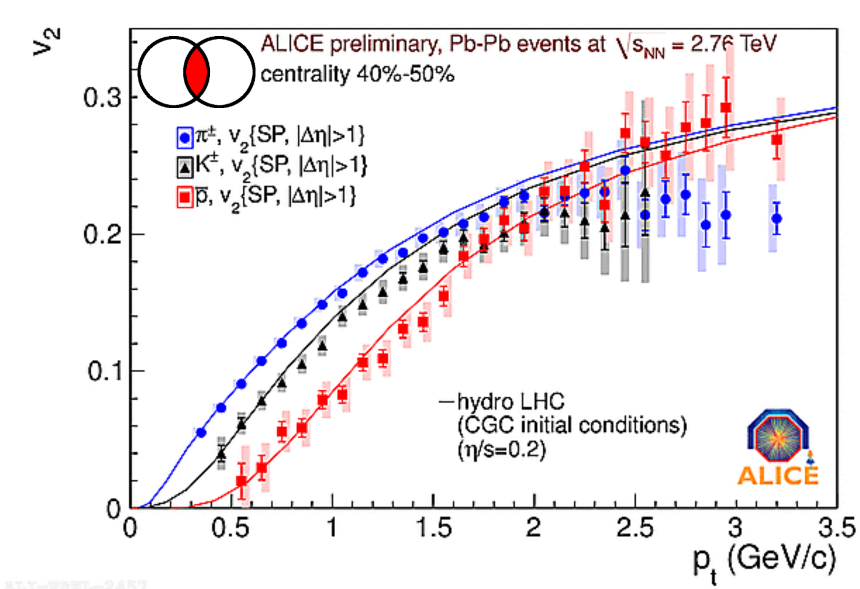
Flow fluctuations

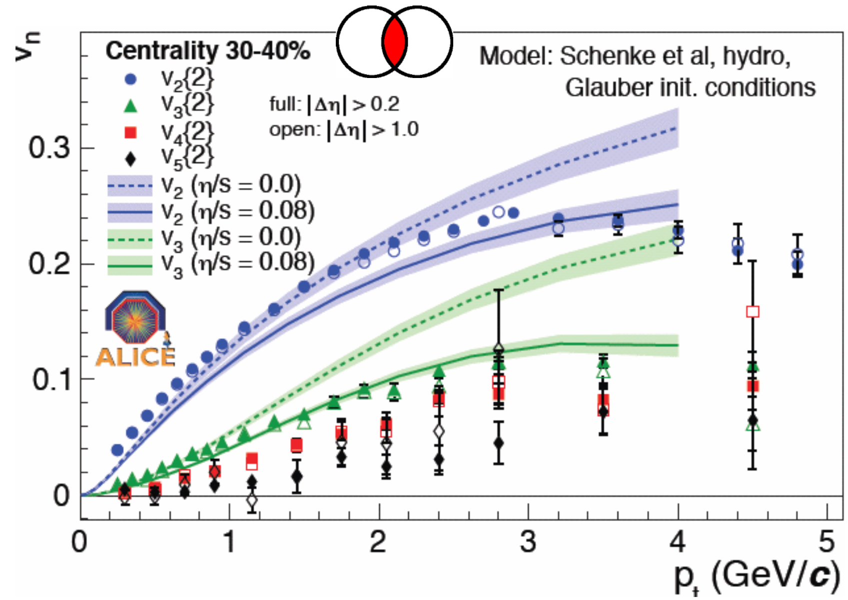




Because the nuclei are not homogenous the initial state can be quite asymmetric giving rise to e.g. triangular flow!
Famous paper: B. Alver, G. Roland, Phys.Rev. C81 (2010) 054905 (420 citations)

Elliptic flow and triangular flow is almost ideal!





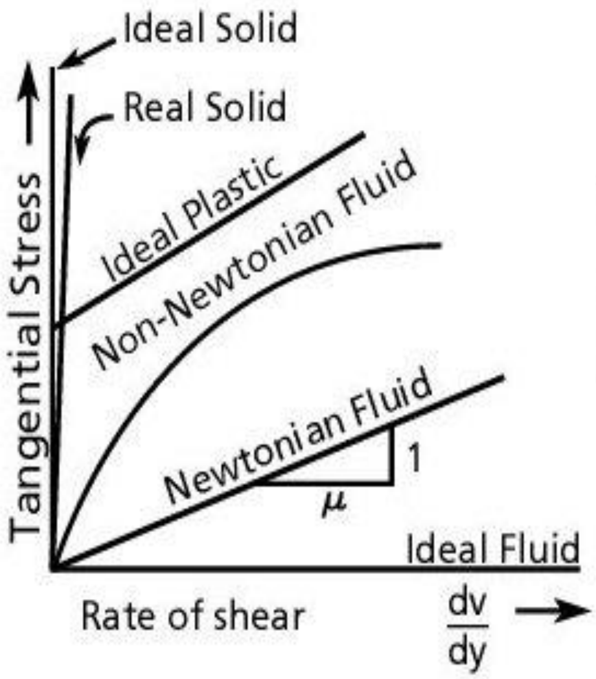
- Huge flow at intermediate p_T:

  2 times more particles in plane than out

  Nearly ideal fluid

- Significant higher order flow caused by fluctuations – also described by nearly ideal hydro + initial state

Shear viscosity



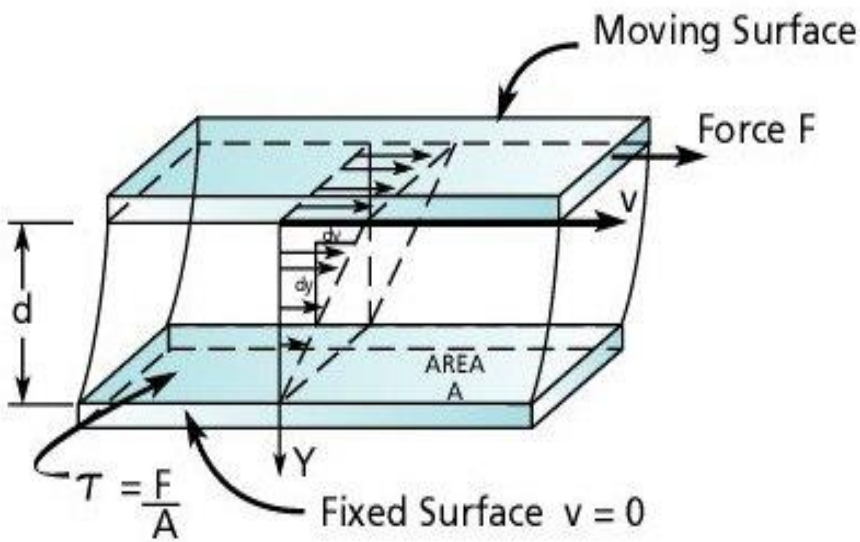


Figure taken from http://www.pumpfundamentals.com/about_fluids.htm

The shear force is given as F=ηAv/d

The shear vicosity-to-entropy density ratio, η/s, is a unitless quantity for characterizing fluids

The QGP behaves like a synchro team!

- Big theoretical challenges:
  - Why is the QGP behaving like a liquid? (hints from string theory)
  - How to go from initial random collisions to organized state in a VERY short time (<1fm/c$\sim10^{-23}$s). This remains to be understood

How to understand this? AdS-CFT

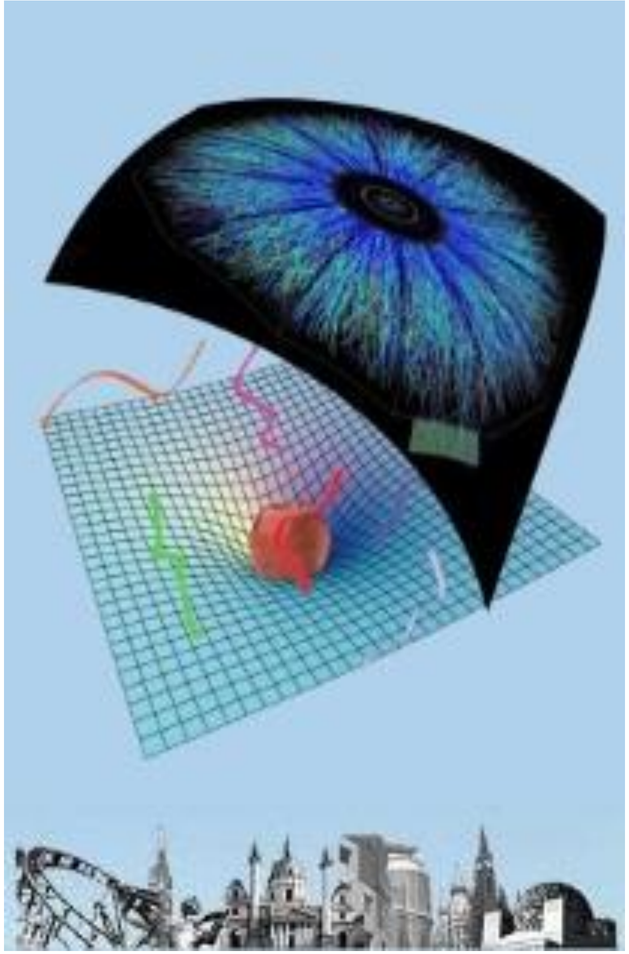
- How to reconcile nearly ideal fluid with energy density like a relativistic gas?
- AdS-CFT correspondence (conjecture)

  J.M. Maldacena, Adv.Theor.Math.Phys.2:231-252, 1998, ~10,000 citations on inspire=most cited
- Duality between ***weakly*** coupled gravity like theory (AdS) and ***strongly*** coupled QCD like theory (CFT)
- QCD like theory, but

  conformal (no confinement, no running coupling) (like QGP!)

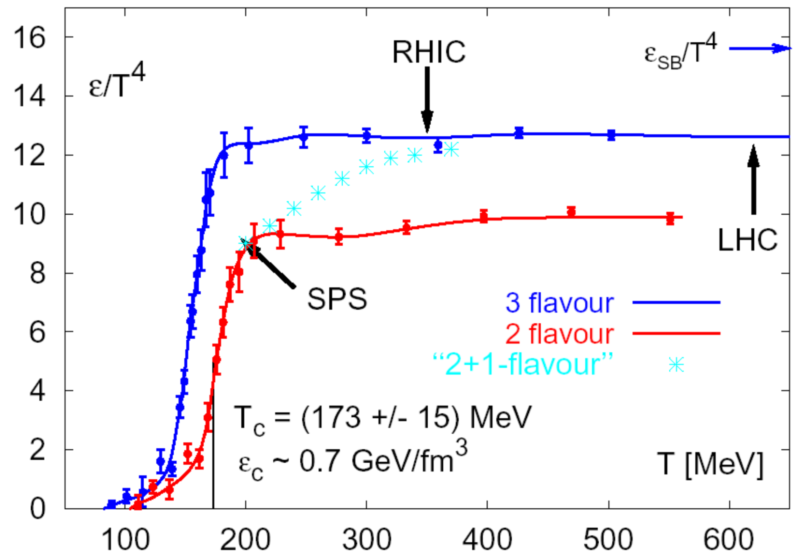
  infinite Ncolors

  SUSY

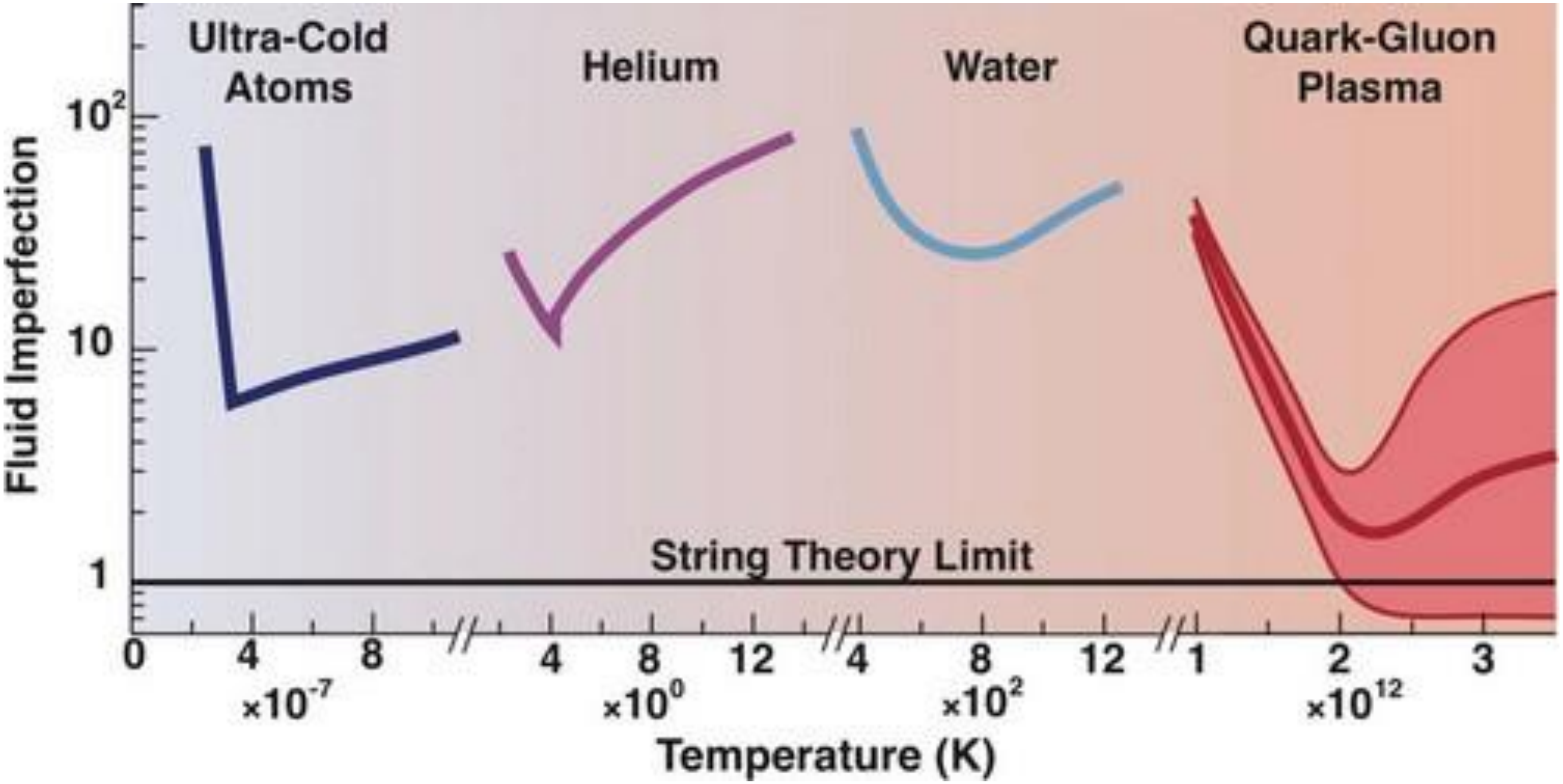


AdS-CFT

- Two very important results:
  - Conjectured bound on shear viscosity: η/s≥1/(4π)~0.08
    - **Viscosity in strongly interacting quantum field theories from black hole physics, P. Kovtun, D.T. Son, A.O. Starinets, Phys.Rev.Lett. 94 (2005) 111601. (~1400 citations on inspire.)**
  - Possibility of infinitely strong coupling at energy density of ¾ SB gas

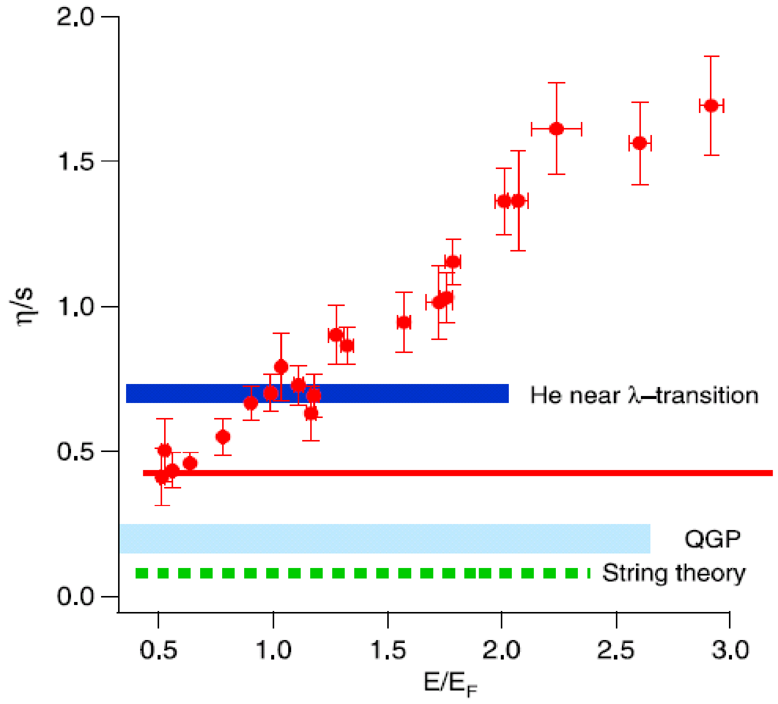


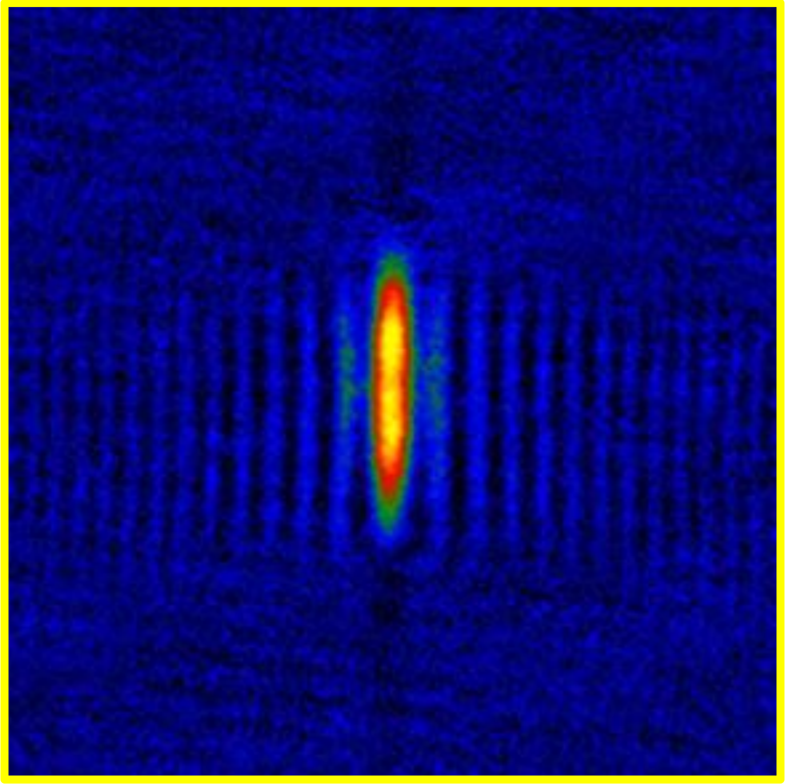
The QGP fluid compared to other fluids



The QGP fluid compared to other fluids

Strongly interacting ultra-cold Li atoms released from a trap





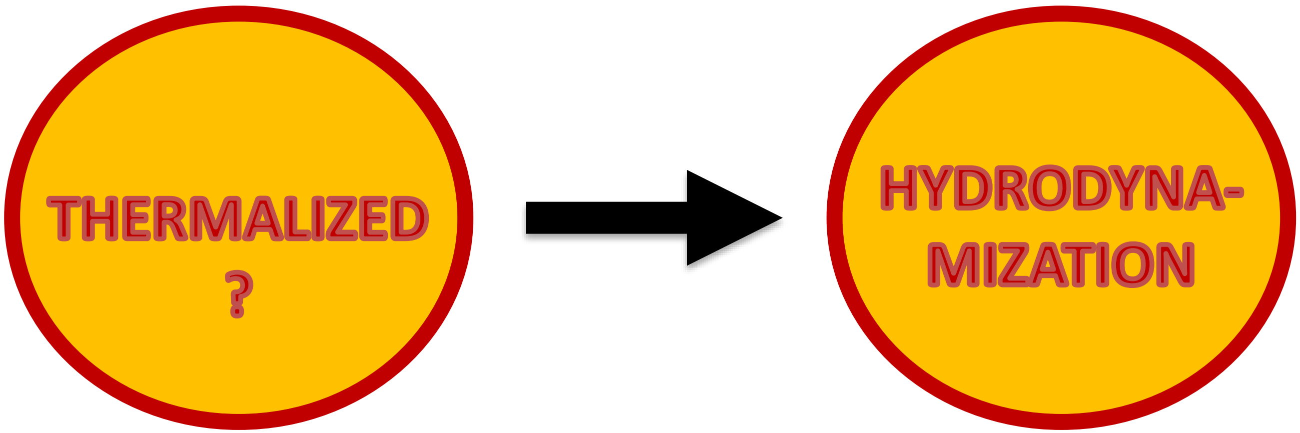
$\eta/s \sim 7 \times 1/4\pi$

http://www.physics.ncsu.edu/jet/index.html

Conclusions about soft physics

- The medium produced in heavy-ion collisions behaves like a nearly perfect liquid. In fact like the most perfect liquid we know!
- This was completely unexpected based on lattice QCD results
- We can get some insight into the liquid nature from AdS-CFT but so far this does not give a full picture e.g. it does not yet describe jet quenching

What about thermalization?

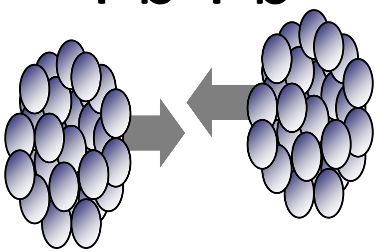


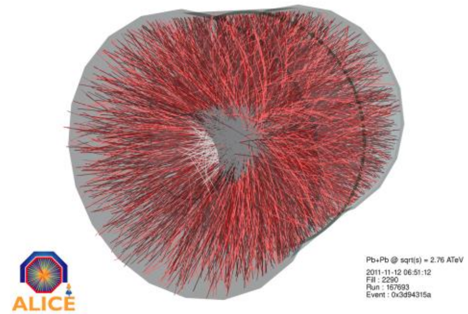
- At some point one thought that hydrodynamics implied thermalization. Now one has found that hydrodynamics can also describe non-equilibrium systems so one talks about hydrodynamization (= when the system can be described by hydrodynamics)
- Thermalization is still being debated

COLLECTIVITY IN SMALL SYSTEMS

The effect of system size:
Macroscopic effects in small systems?

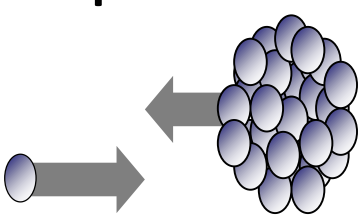
Pb-Pb

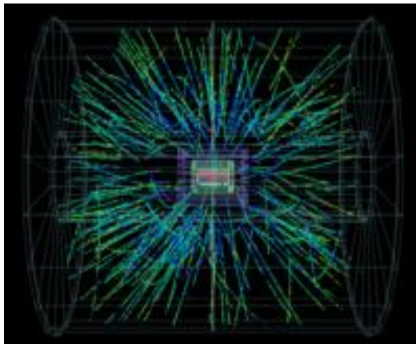




Hot nuclear matter

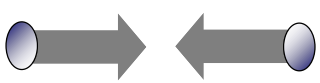
p-Pb

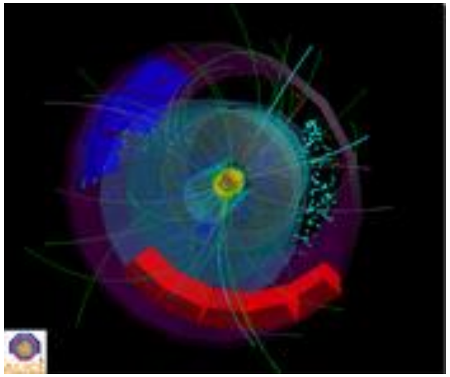




Cold nuclear matter

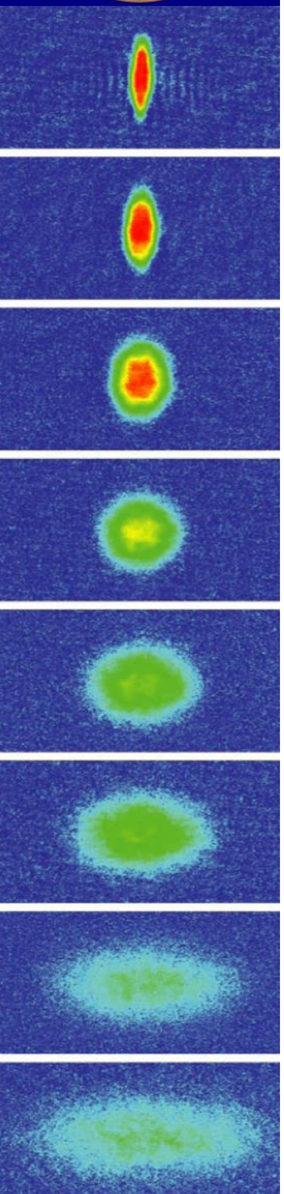
pp

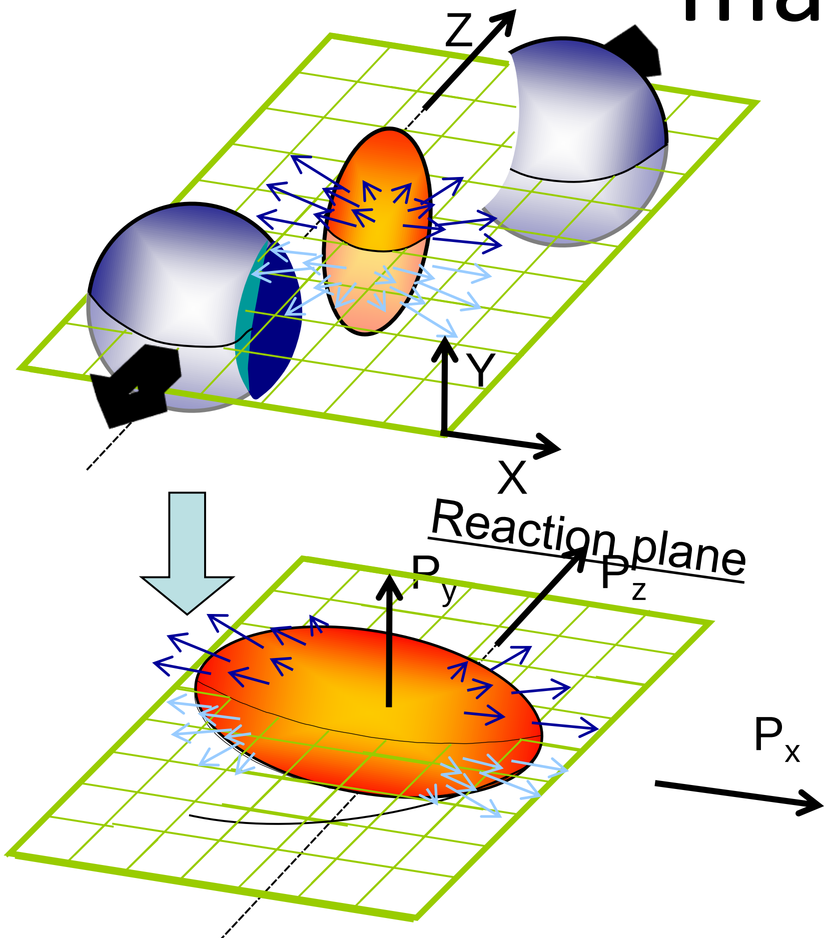


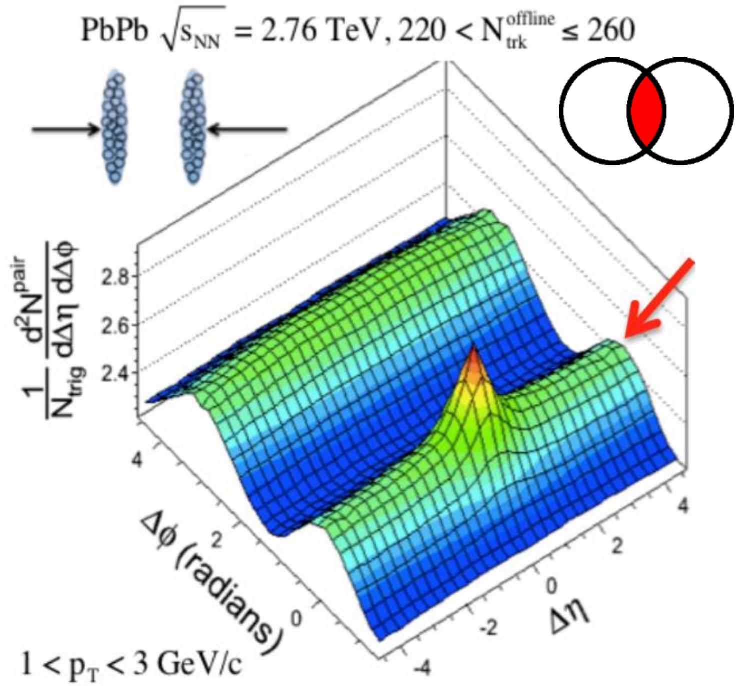


QCD baseline

The ridge: a fingerprint of the macroscopic QGP

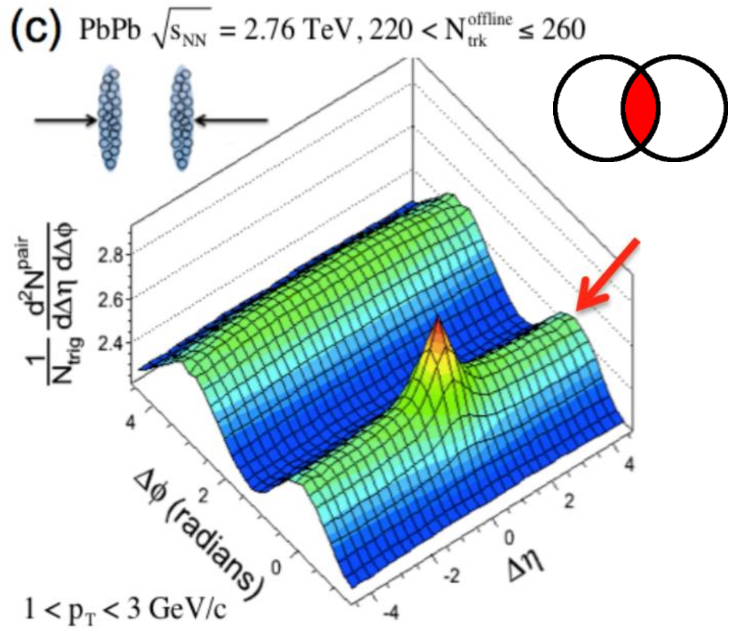






Initial spatial anisotropy → Strong pressure gradients → v_2 Azimuthal anisotropy → **Sensitivity to early expansion**

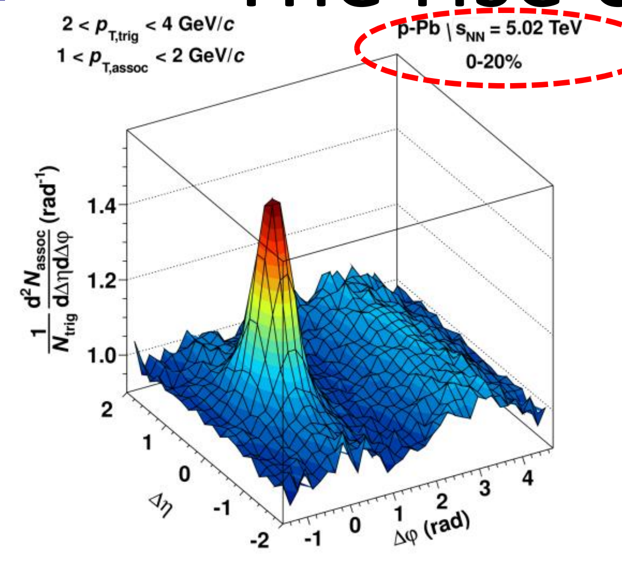
Elliptic flow from the ridge



Singles: $$\frac{dN}{d\varphi} \propto 1 + \sum_n 2vn \cos(n[\varphi - \Psi_n])$$

Pairs(a,b): $$\frac{dN}{d\Delta\varphi} \propto 1 + \sum_n 2v_{n,a}v_{n,b} \cos(n\Delta\varphi)$$

Collectivity in p-Pb collisions? The rise of the double ridge



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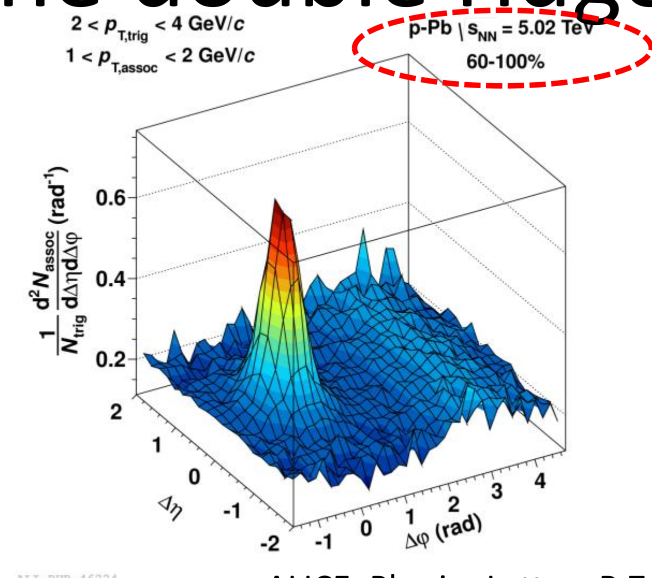

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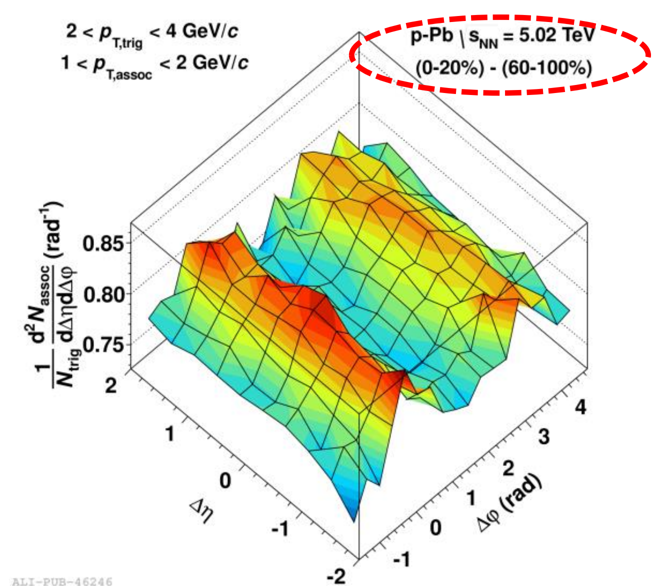
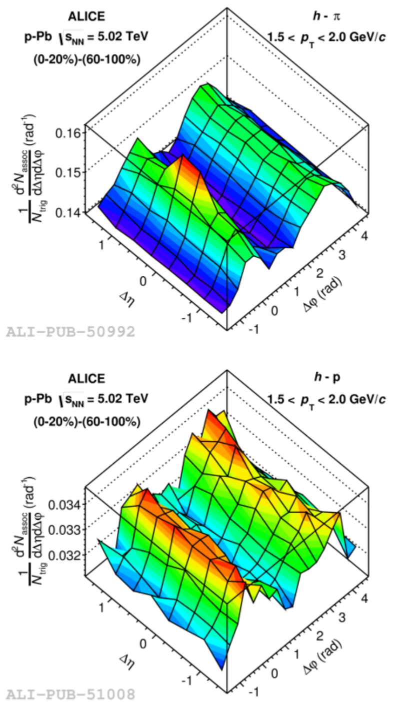
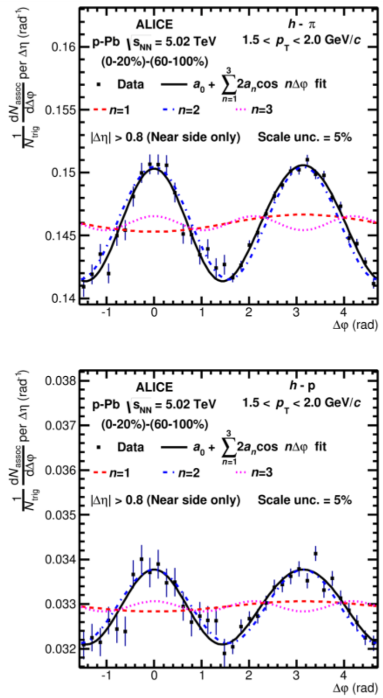


ALICE: Physics Letters B 719 (2013)

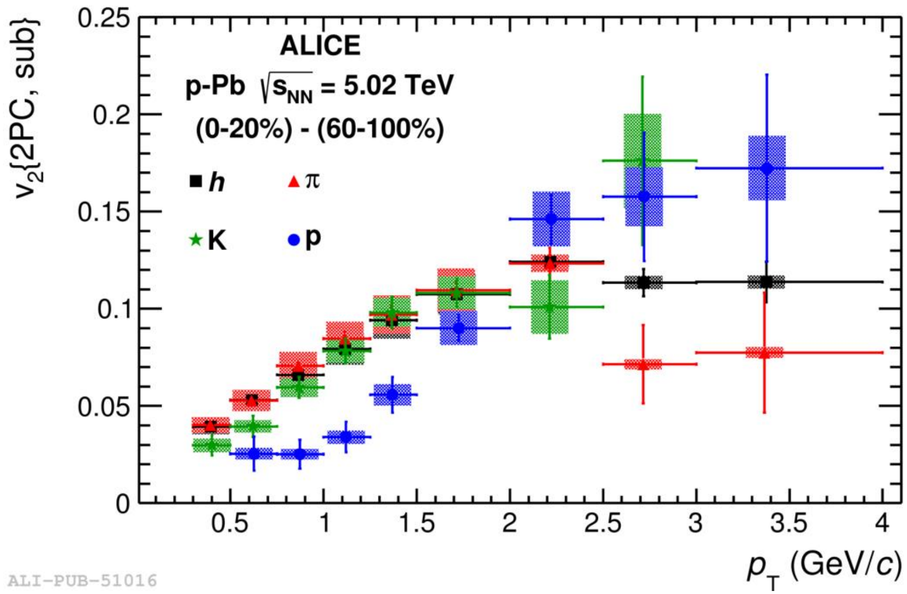
- Double ridge structure reminiscent of azimuthal flow in Pb-Pb collisions

Using Particle IDentification to study the double ridge





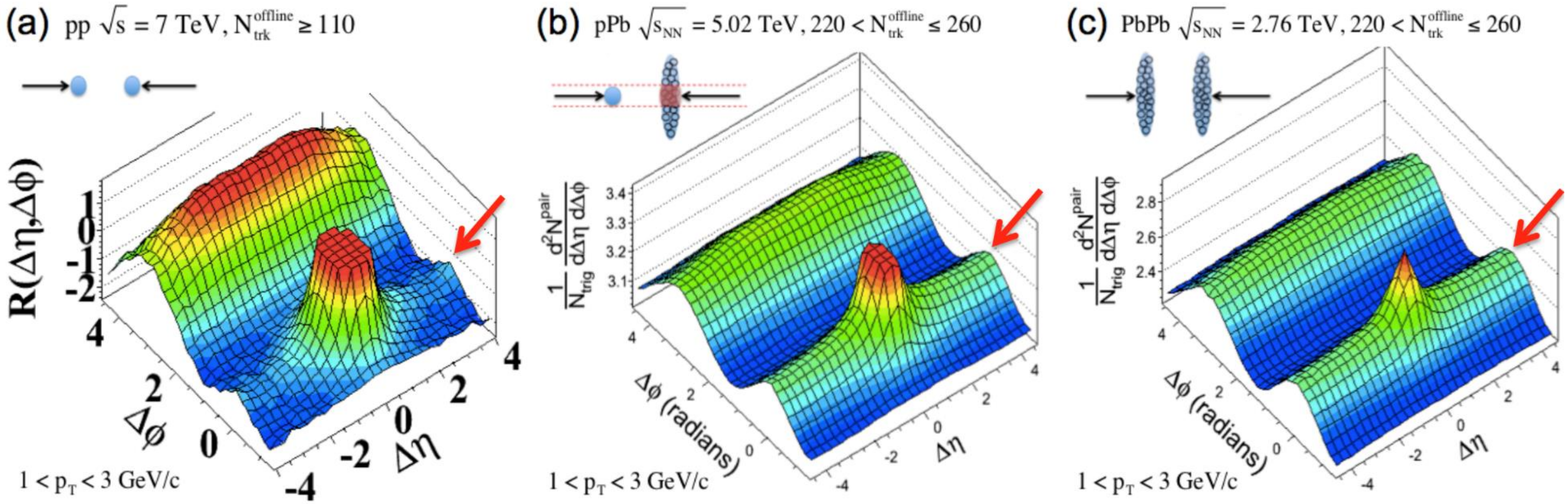
Fourier coefficients:



Phys. Lett. B 726 (2013) 164–177

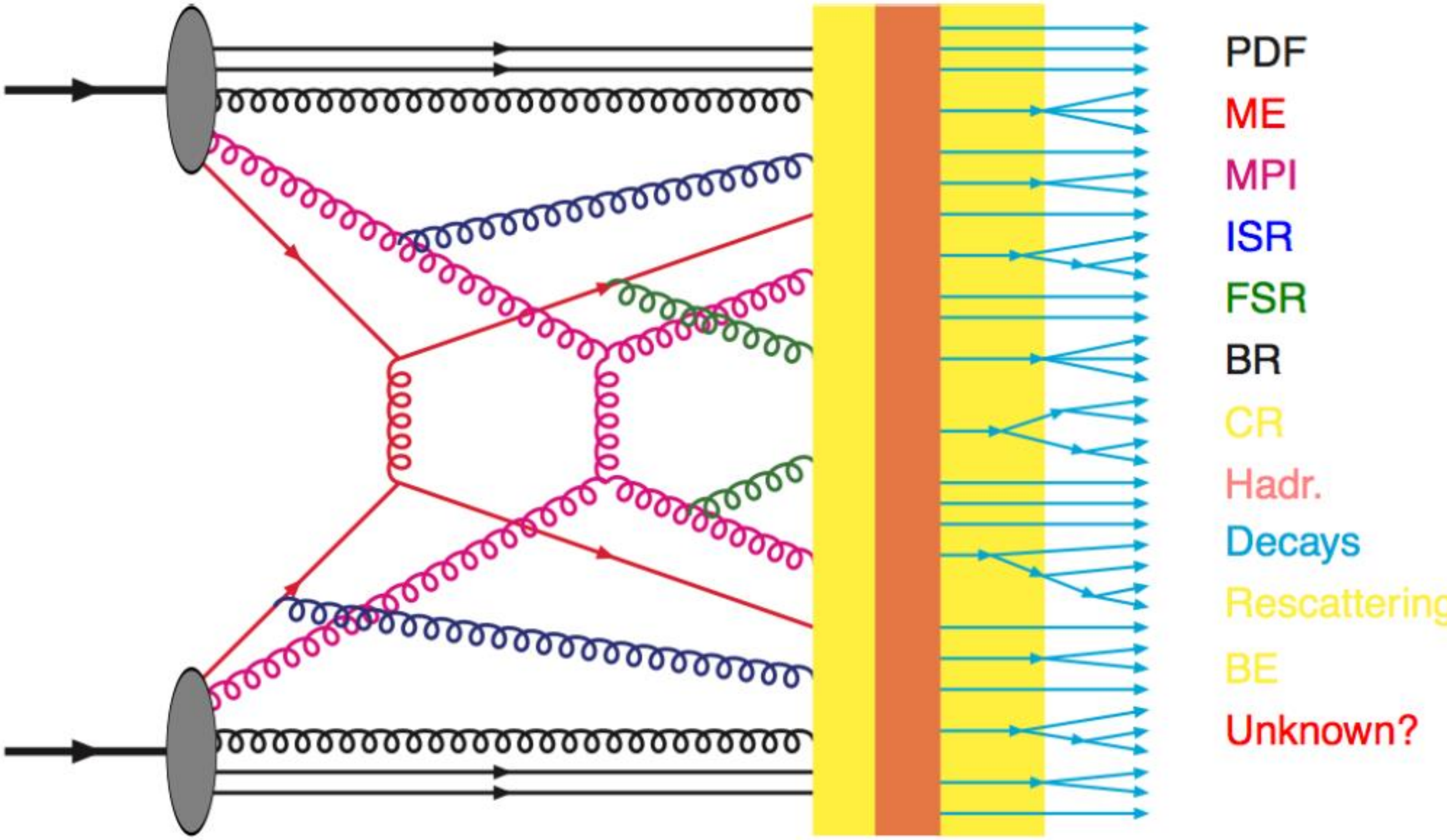
- Clear mass ordering suggests flow

Ridges in all systems



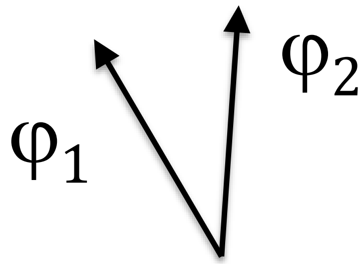
The perfect liquid is produced in all systems suggesting that small QCD systems produce “macroscopic” matter

What about other effects in small collisional systems?



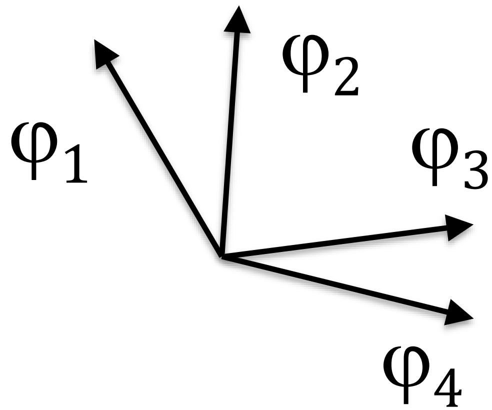
Need to use more advanced methods

2 particle correlations



- Measure $\langle 2\rangle = \langle\cos(2[\;{}_2 - \varphi_1])\rangle = \langle v_2^2\rangle$ (see exercise 1.2)
- Define $v_2\{2\} = \sqrt{\langle 2\rangle}$
- Caveat: has contributions from, e.g., resonance decays

4 particle correlations



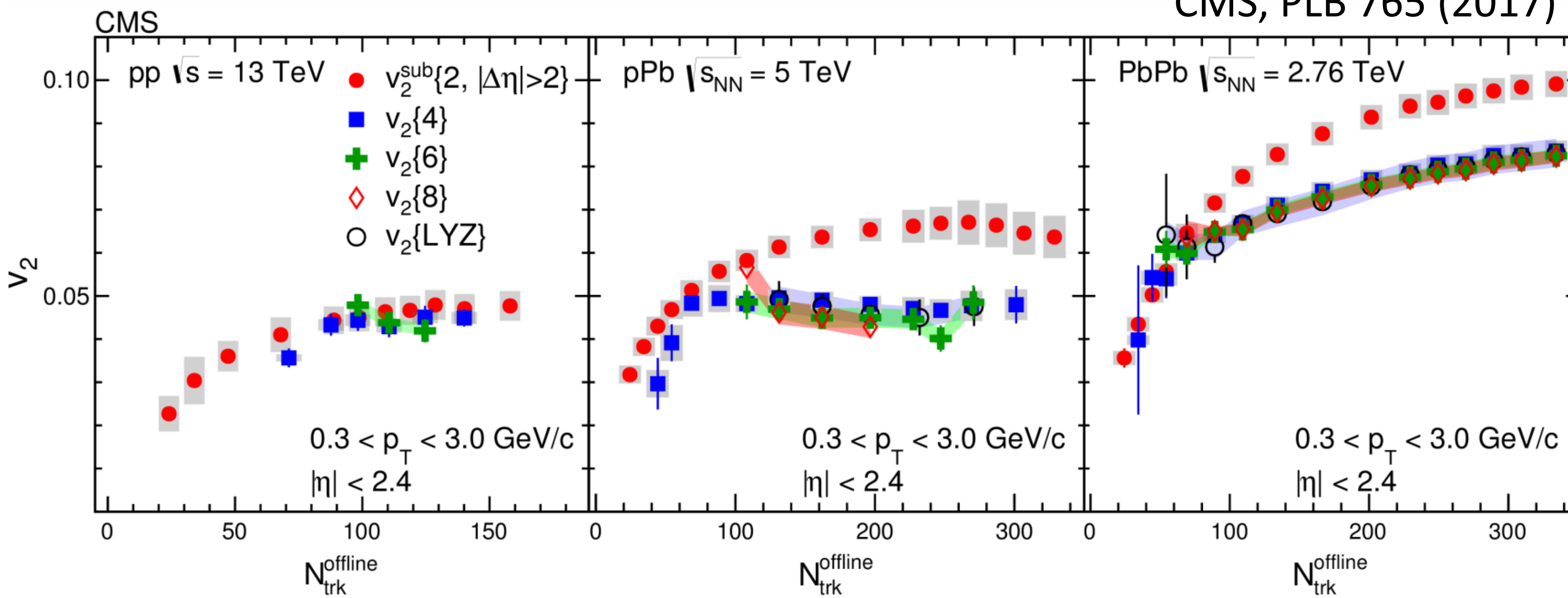
- Measure $\langle 4\rangle = \langle\cos(2[\varphi_1 + \varphi_2 - \varphi_3 - \varphi_4])\rangle$
- Define cumulant:

$$c_2\{4\} = \langle 4\rangle - 2\langle 2\rangle^2$$

- Now $v_2\{4\} = \sqrt[4]{-c_2\{4\}}$
  will not have contributions from 2-particle correlations such as resonance decays!

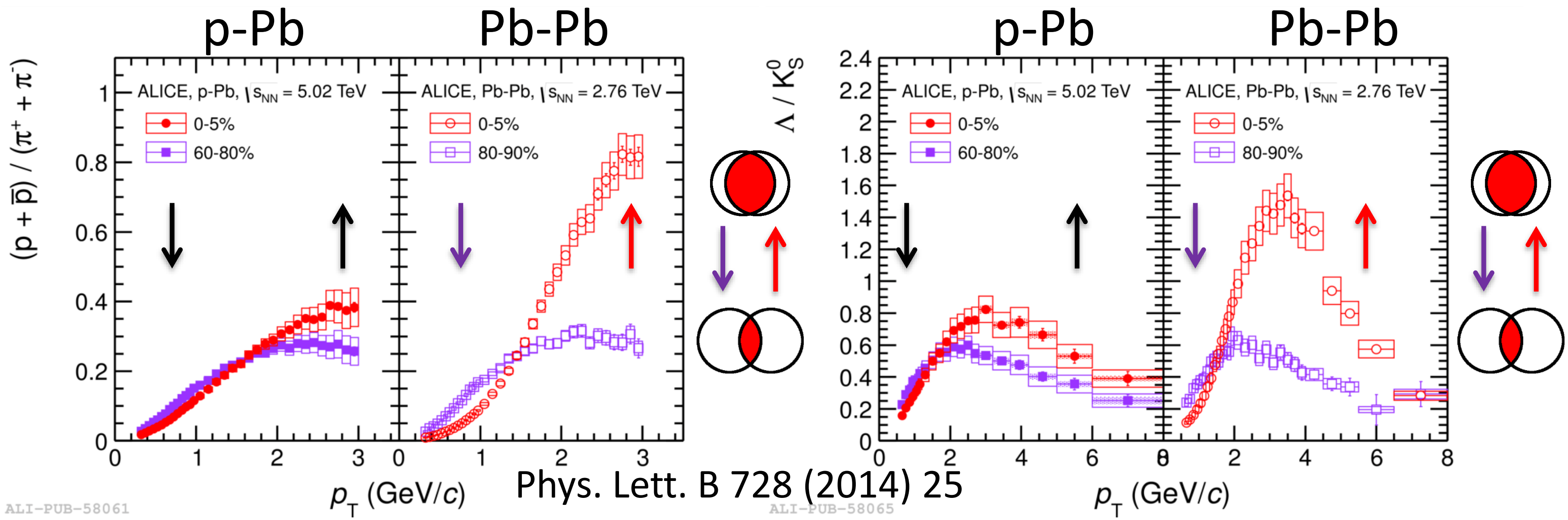
Collectivity in small systems?

CMS, PLB 765 (2017) 193



- Also higher-order flow coefficients are all consistent with flow
- Note that $v_2\{4\}$ can only be calculated for $c_2\{4\} < 0$, which is considered a strong indication for collectivity

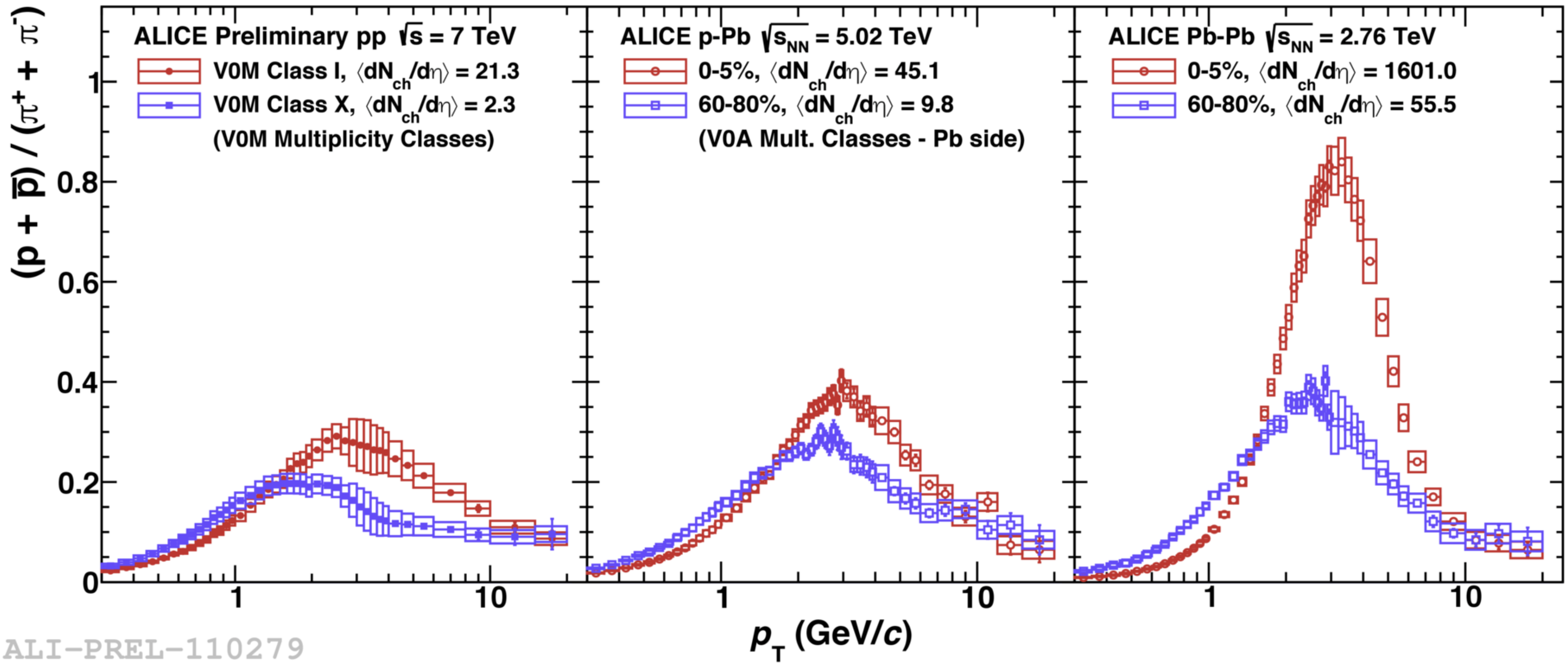
Particle ratios in p-Pb and Pb-Pb show similar features



Phys. Lett. B 728 (2014) 25

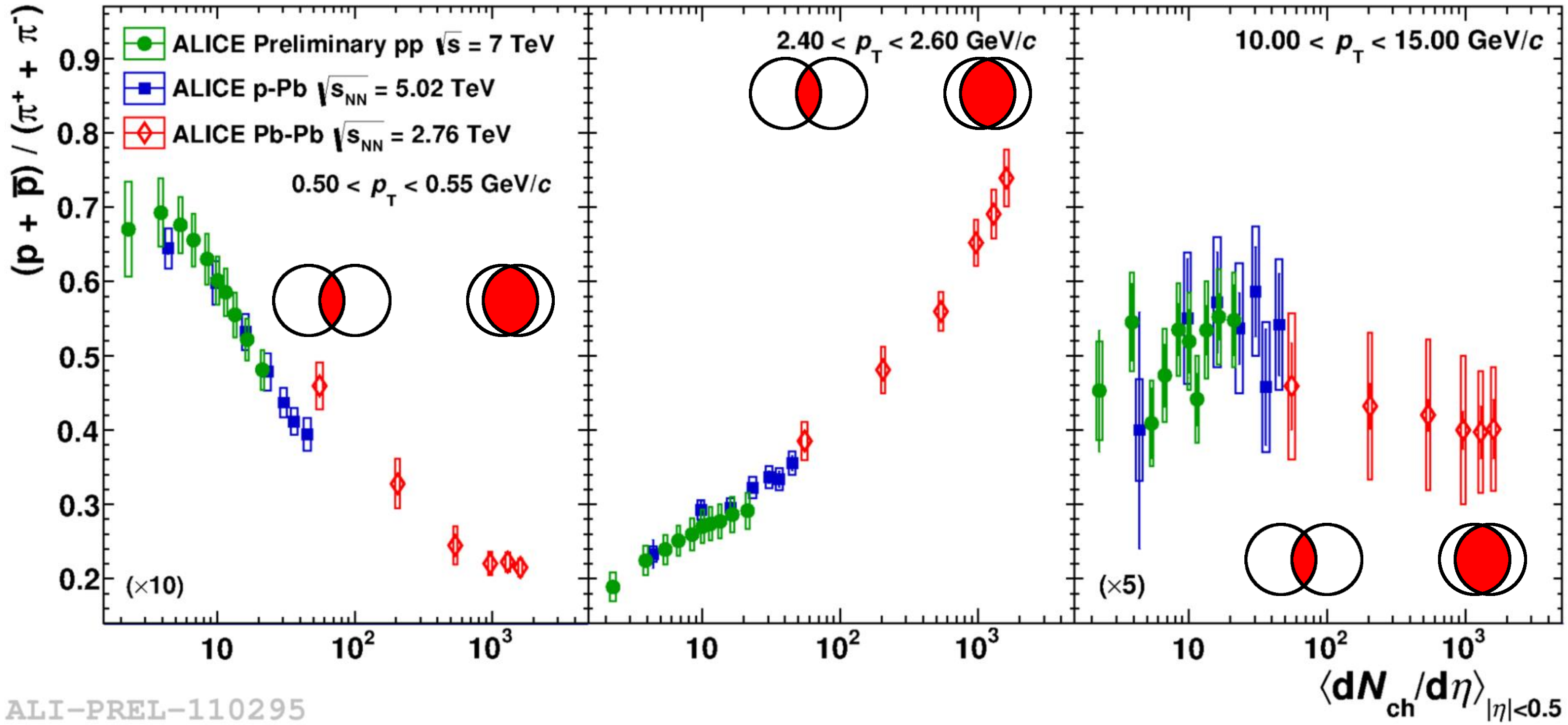
- Characteristic evolution of p/π and Λ/ K^0_S with multiplicity is reminiscent of Pb-Pb where it is believed to be due to radial flow
- NB! The solid boxes for p-Pb ratios indicate the uncorrelated systematic error
  $\Rightarrow$ the relative trend can be measured rather precisely

Also seen in pp



- The proton-to-pion ratio shows the qualitative same features, indicative of radial flow, in all 3 systems

Universal trend vs <dN/dη>



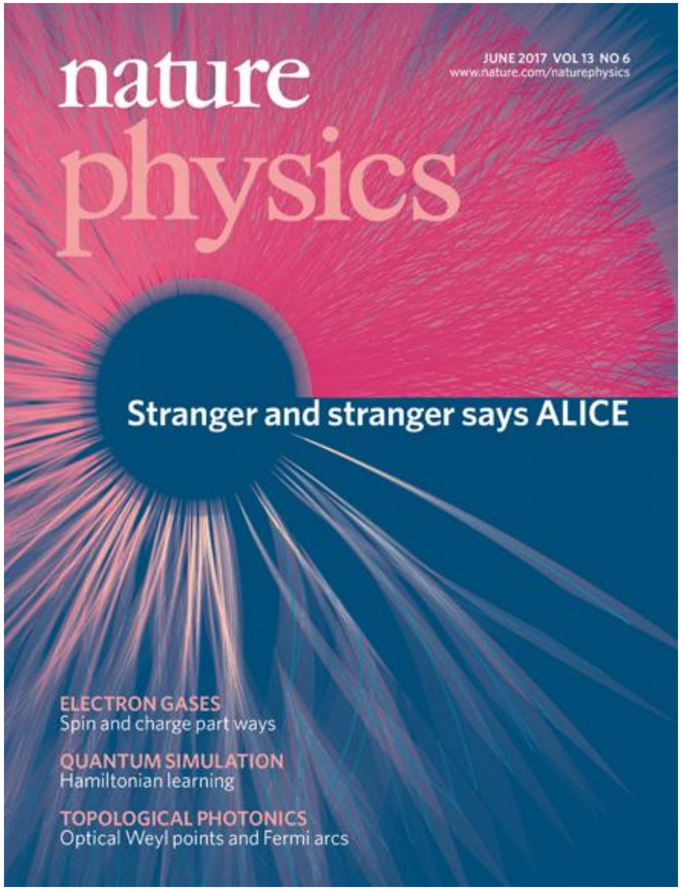
Low p_T:
Decrease

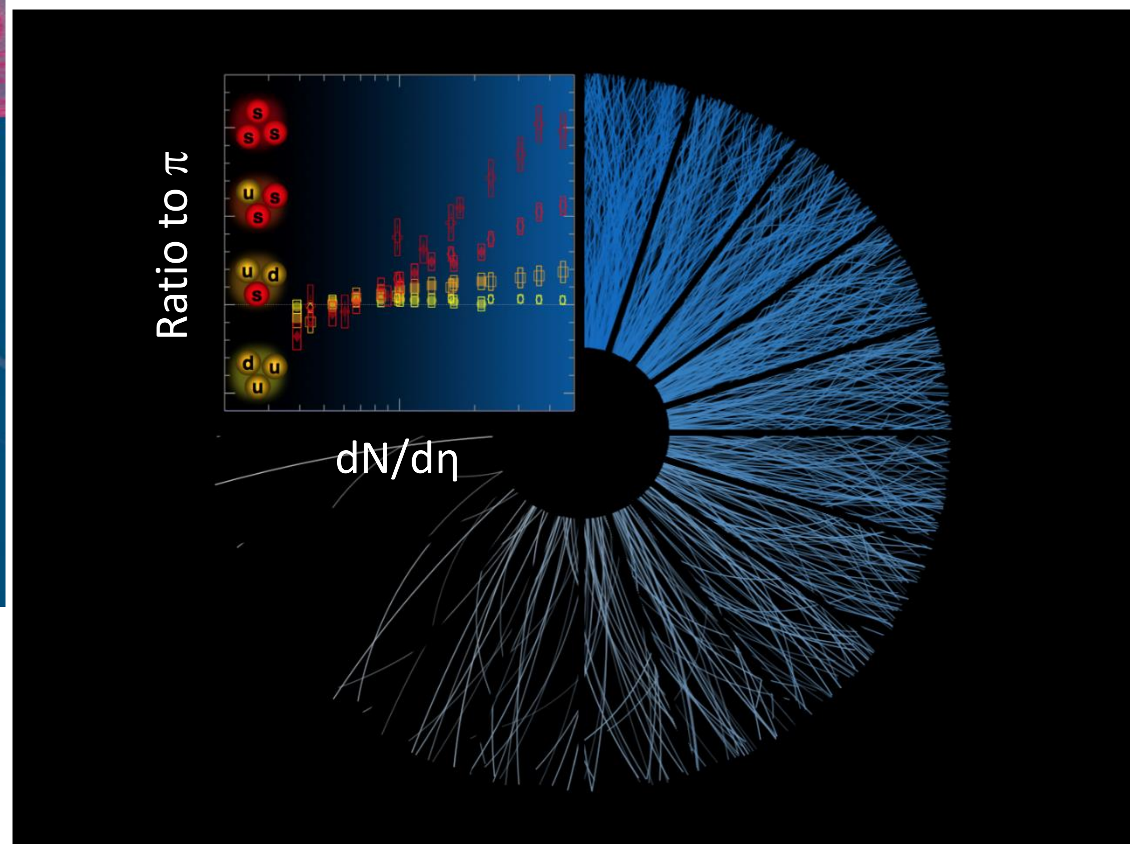
Intermediate p_T:
Increase

High p_T:
Unchanged

Note that <dN/dη> is measured at forward rapidity to avoid auto-correlations.

Strangeness production in small systems

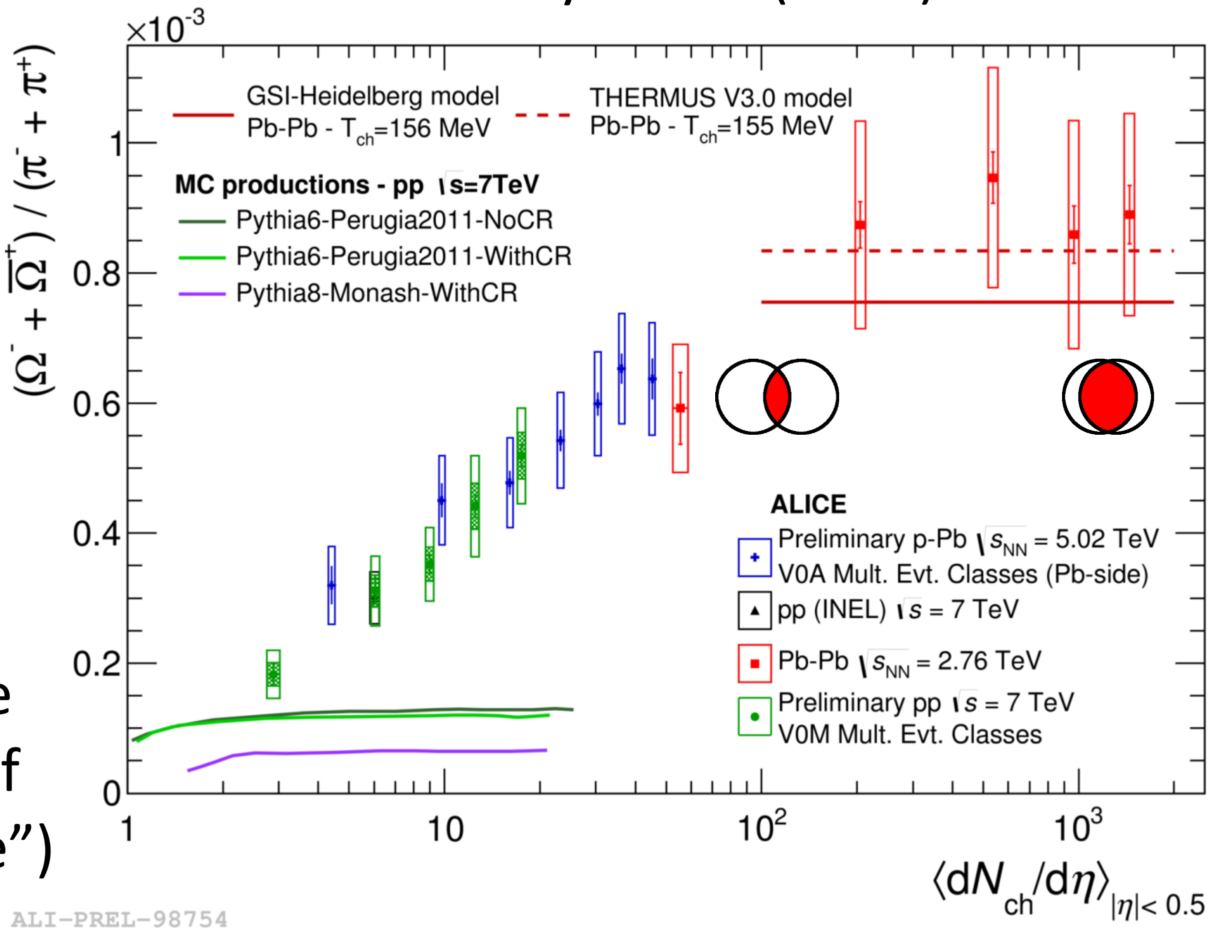




pp phenomenologists' favorite figure from ICHEP 2016

Nature Physics 13 (2017) 535

(sss)

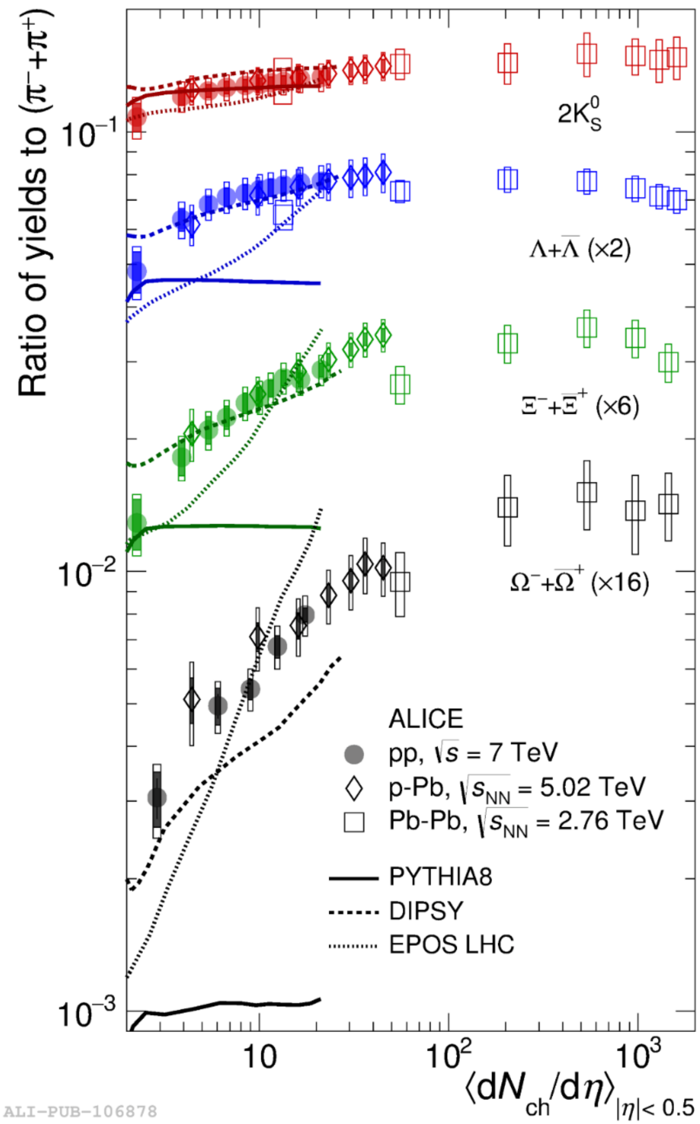


AA limit

e⁺+e⁻like ("more of the same")

Need bulk physics (medium) to describe increase!

Integrated particle ratios



DIPSY Color rope model:
C. Bierlich, G. Gustafson, L. Lönnblad, A. Tarasov (Jefferson Lab), JHEP 1503 (2015) 148

Nature Physics 13 (2017) 535

Macroscopic vs microscopic models

- Stat. thermal model
  - Canonical
  - Grand-canonical
- Hydrodynamics
  - Radial flow
  - Azimuthal anisotropic

- Tunneling of $q\bar{q}$-pairs
  - Strings
  - Ropes
- String interactions
  - Color reconnection
  - Shoving

Project: “**Pinning down the origin of collective effects in small collision systems**”

Grant: SEK 26 200 000 over five years

Principal investigator: **Associate Professor Peter Christiansen, Lund University**

P. Christiansen (Partikelfysik)

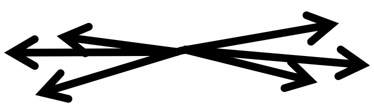
L. Lönnblad (Teoretisk högenergifysik)

Discovery → Control: Transverse spherocity

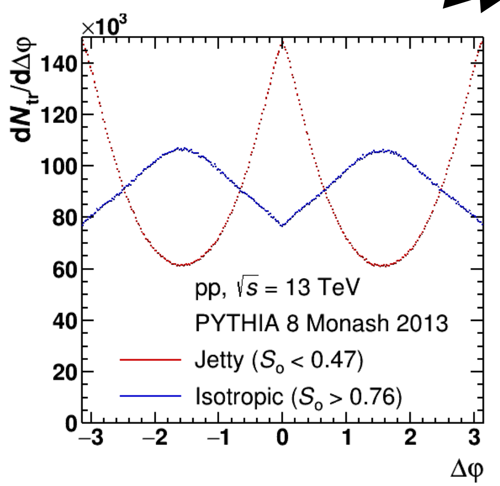
- In the following analyses, transverse spherocity, S_O, is used to characterize event shapes:
  - Low spherocity: "jetty" events
  - High spherocity: "isotropic" events

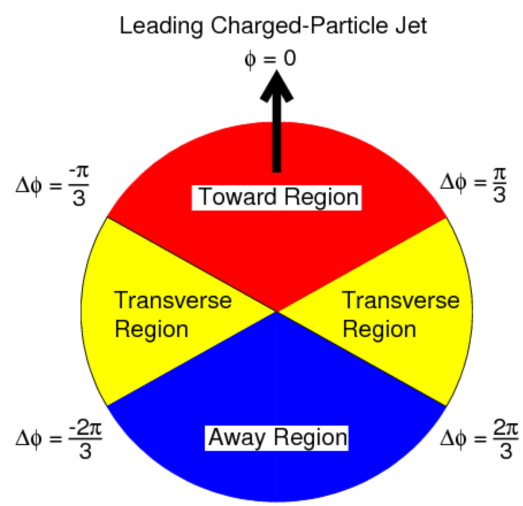
$$S_O = \frac{\pi^2}{4} \min_{\hat{n}} \left(\frac{\sum_i \vec{p}_{T,i} \times \hat{n}}{\sum_i p_{T,i}} \right)^2$$

Jetty: $S_O \to 0$



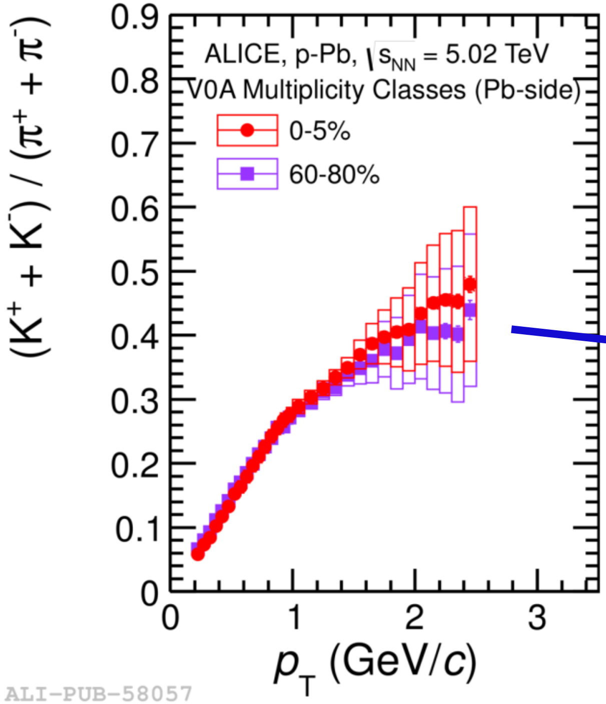
Isotropic: $S_O \to 1$



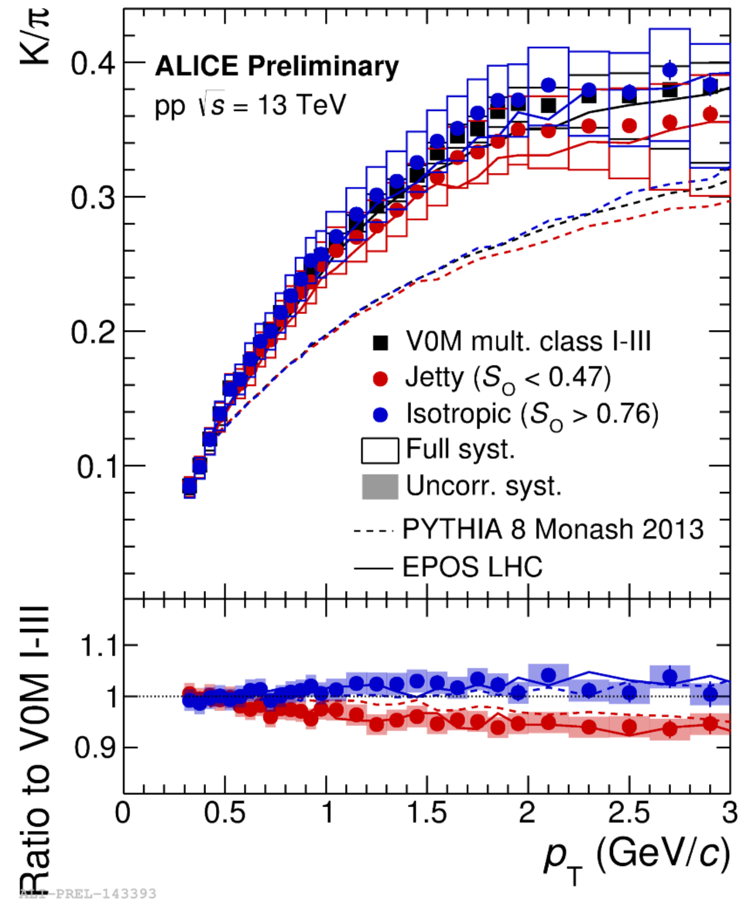


p_T-differential K/π ratios as a function of S_O

- Isotropic:
  - No strong modifications of ratios expected (à la p-Pb)



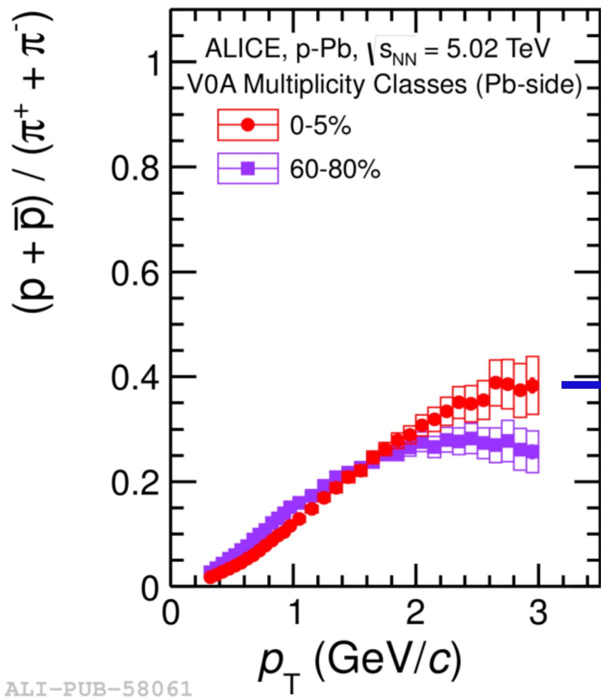
Phys. Lett. B 728 (2014) 25-38



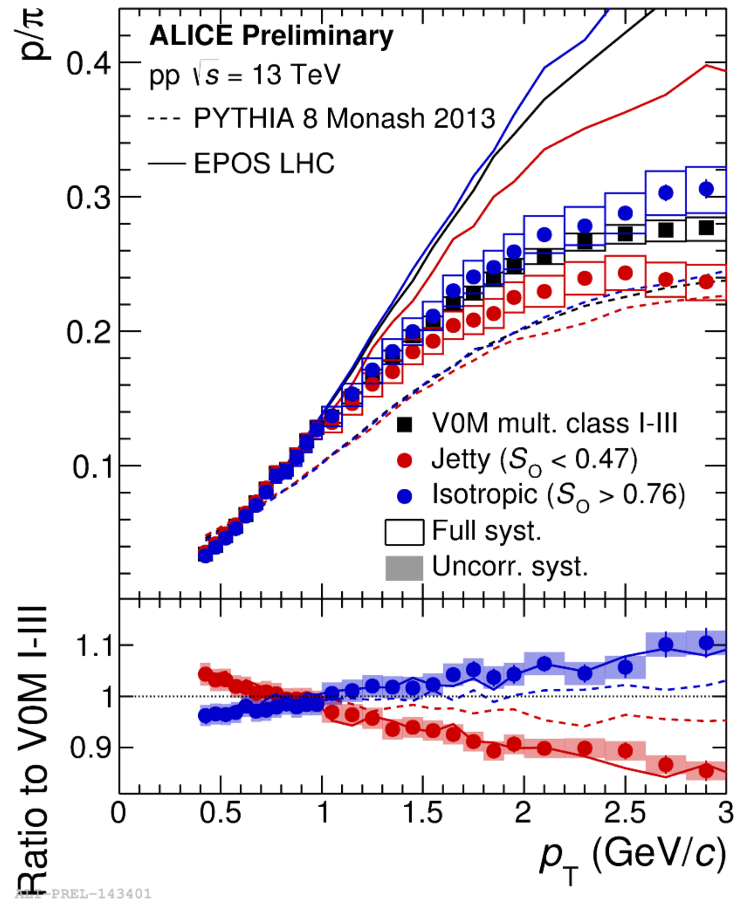
Analysis by
Vytautas Vislavicius

p_T-differential p/π ratios as a function of S_O

- Isotropic:
  - In p-Pb, flow manifests via boosted proton spectra



Phys. Lett. B 728 (2014) 25-38



Analysis by
Vytautas Vislavicius

What next?

- More sensitive observables
  - Multi-strange baryons are more sensitive to QGP-like effects – will they show bigger effect?
  - The phi has hidden strangeness and can differentiate between some QGP and some string-like models
- Different estimators
  - Can we find an estimator that both enhances radial and elliptic flow?

- Collectivity in small systems challenges two paradigms at once!
  1. How far down in systems size does the "SM of heavy ions" remain?
  2. Can the standard tools for min bias pp remain standard?

C. Bierlich

Conclusions

- The reason we believe in a QGP
  - We form a medium with strong collective properties that behaves as a nearly ideal fluid with a "temperature" in the range required by LQCD
  - It quenches jets and gives rise to energy losses of order GeV/fm which is as expected from color charges
  - It dissolves quarkonia as expected indicating that the long range confining potential is screened and indeed it seems that deconfined charm quarks regenerate to form J/ψ mesons
- The challenge to this paradigm is that we now observe medium like effects in small systems
  - We need to understand this better!

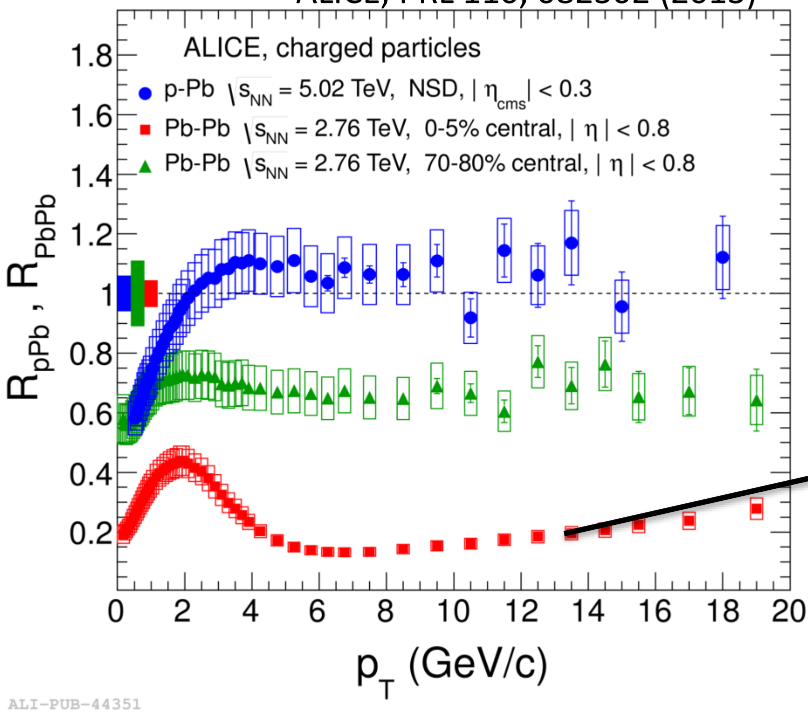
Backup slides

What I did not cover

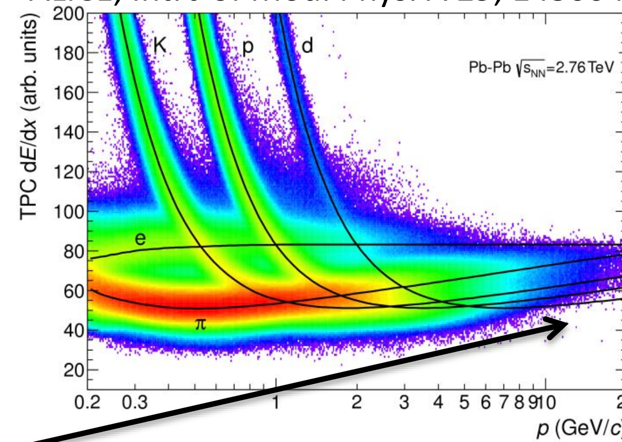
- Charged multiplicity and hadronization
  - Color Glass Condensate and Statistical Model
- Quark scaling of elliptic flow
  - Recombination of quark like degrees of freedom (picture is more complicated at LHC)
- Quenching and flow of heavy quarks
  - Suggests that there are also large collisional energy loss

Extending the R_{AA} to identified particles

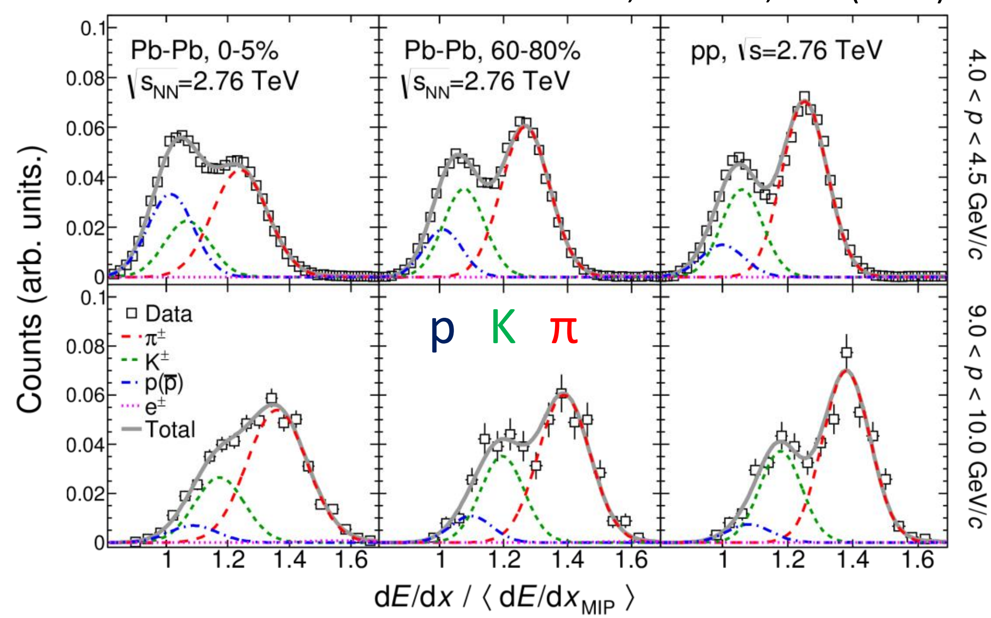
ALICE, PRL 110, 082302 (2013)



ALICE, Int. J of Mod. Phys. A 29, 1430044 (2014)



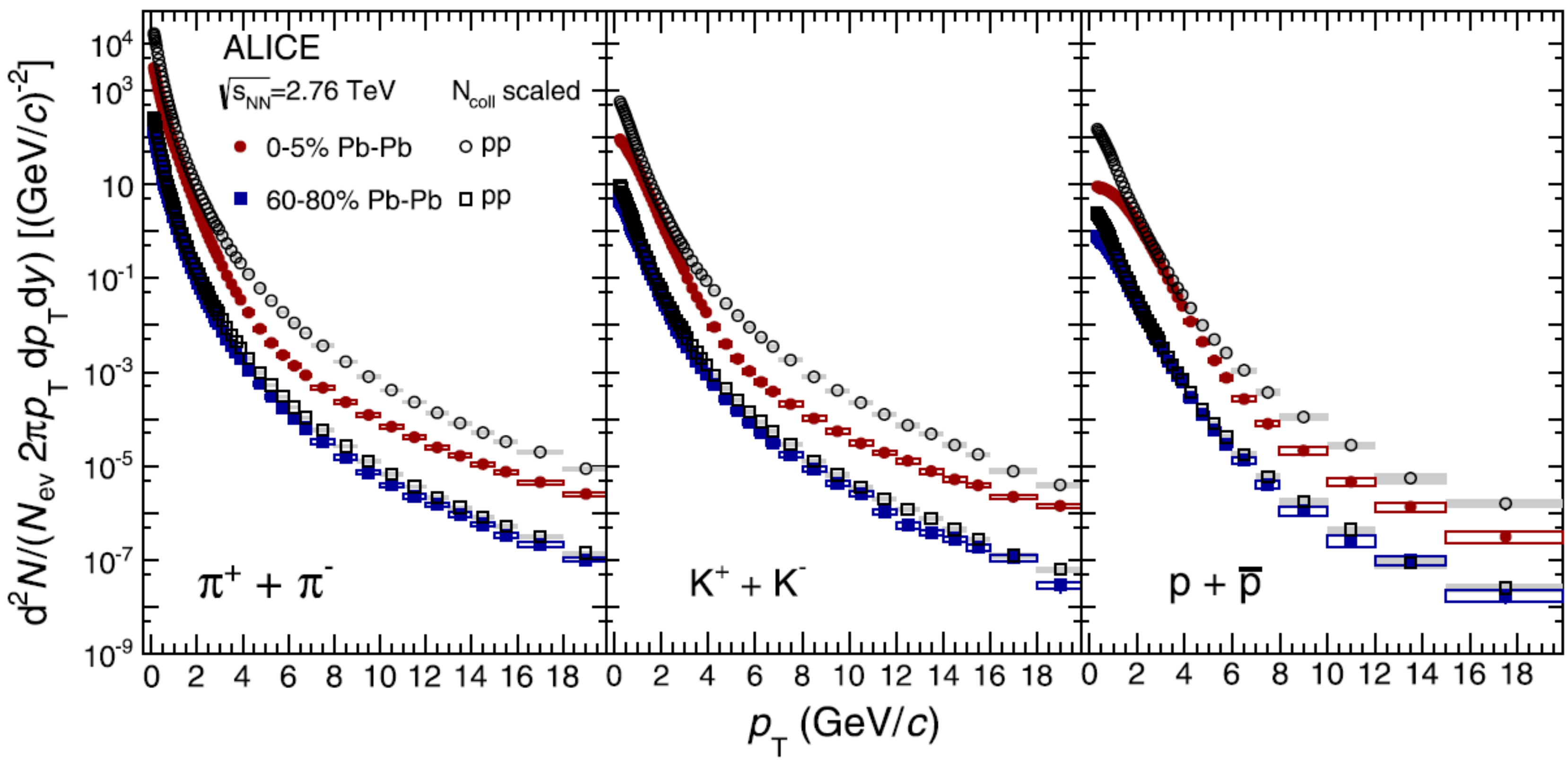
ALICE , PLB 736, 196 (2014)



Each TPC track also has an associated dE/dx that can be used for ID on the relativistic rise

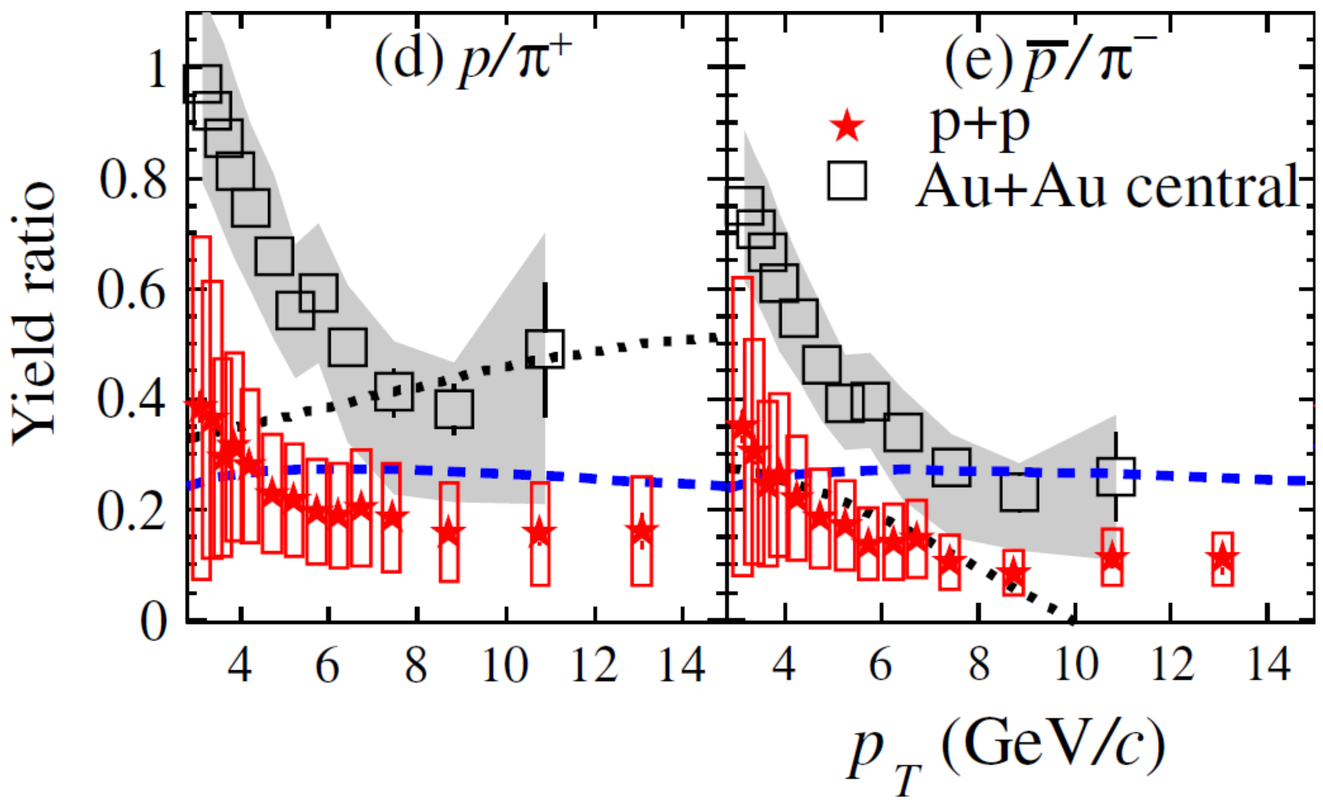
Charged π, K, and p spectra in pp and Pb-Pb collisions

ALICE , PLB 736, 196 (2014)



High p_T p/π ratio in pp and central Au-Au collisions at $\sqrt{s_{NN}}$=200 GeV

STAR, PRL 108, 072302 (2012)

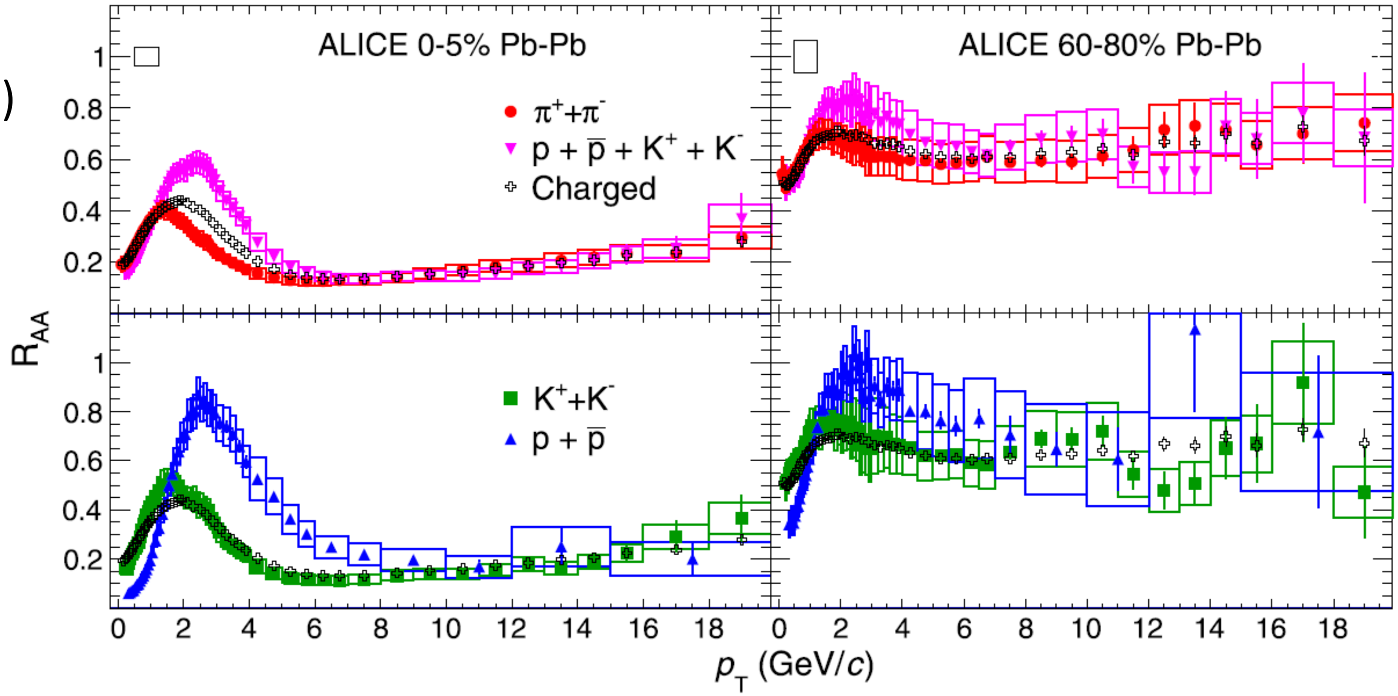


FFs:
DSS
AKK2008

NB! At this low energy there is significant difference between baryon and anti-baryon production.

The nuclear modification factor R_{AA} for charged π, K, and p

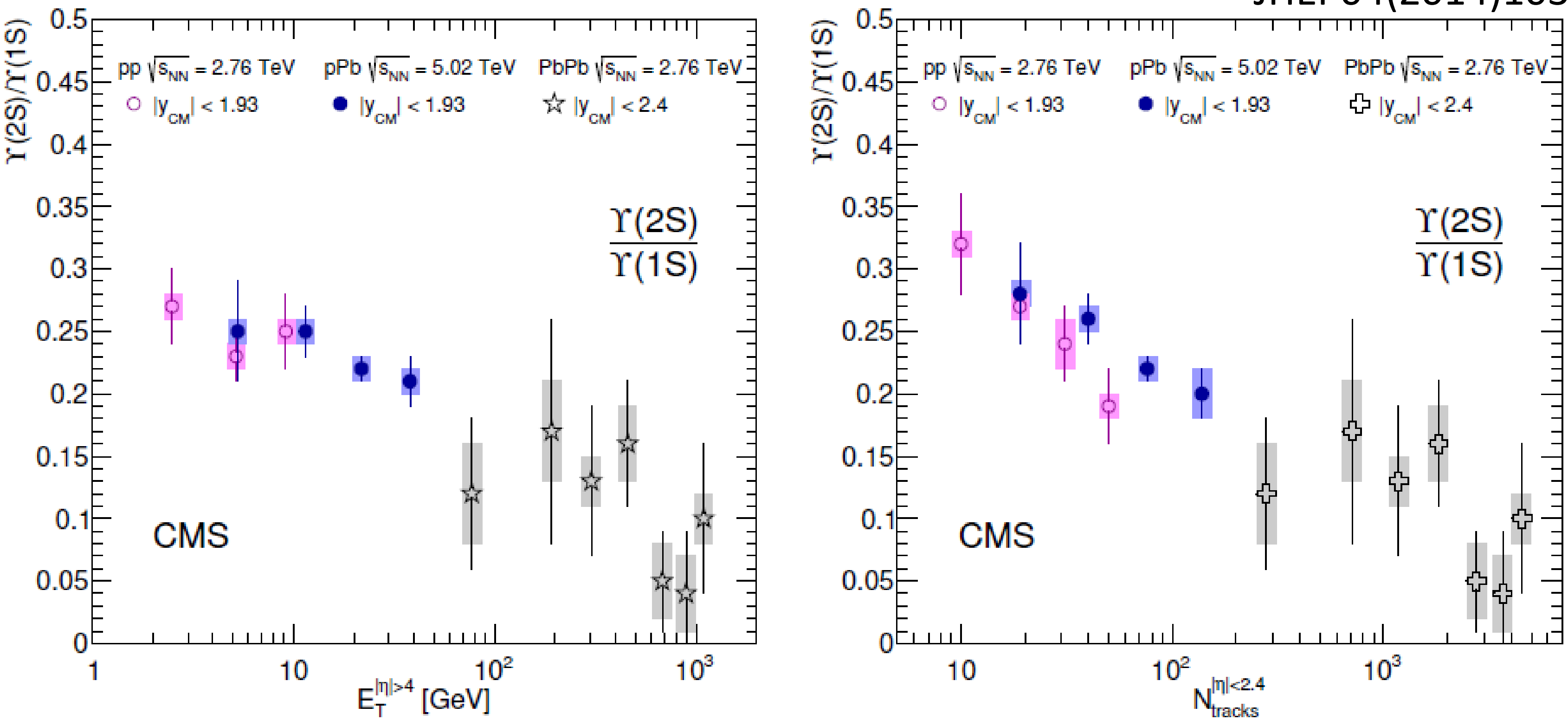
PLB 736, 196 (2014)



The nuclear modification factors are consistent for π, K, p for p_T > 10 GeV/*c*.
For 0-5% Pb-Pb collisions: K and π to within ~10% (syst), p and π to within ~20% (syst)
The results rules out popular ideas of jet quenching where the large energyloss is coupled to large leading order particle species dependent effects.

The relative production of ϒ states is affected by event acivity (?)

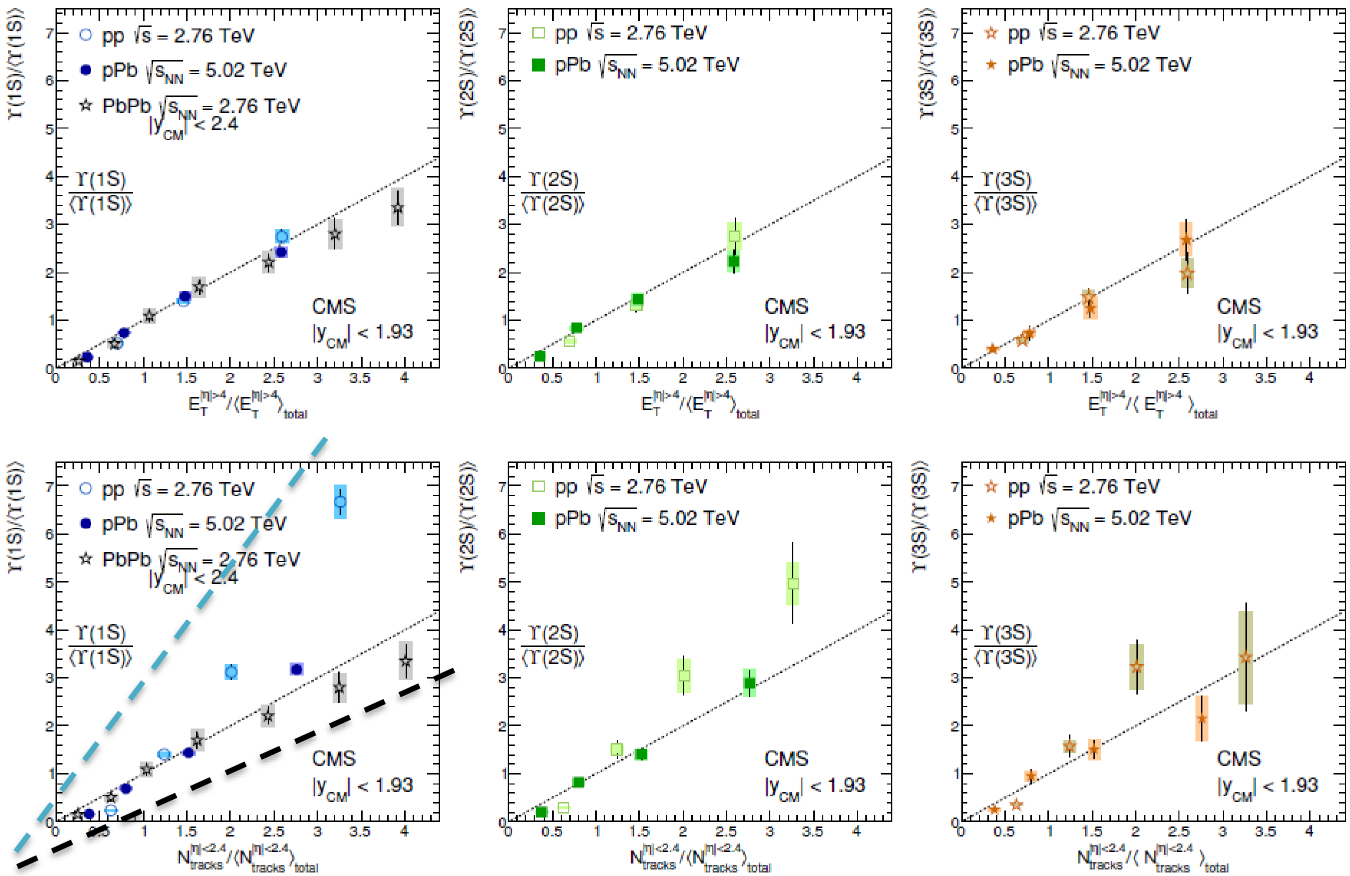
JHEP04(2014)103



Also for ϒ we observe varying trends for the different states even for pp collisions!
(More info on next slide)

Heavy-ion physics and the QGP (P. Christiansen, Lund)

In small systems we observe enhancement with event activity



JHEP04 (2014)103

In this case we observe different trends with event activity

Caveat: should study binary scaling (but this does not work for pp)