

A Space-Time Approach to Analyzing Stochastic Gradient Descent

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Abstract

We present a diagrammatic calculus for reasoning about the behavior, at small learning rates, of SGD and its variants. We interpret the diagrams as histories of scattering events, thus offering a new physical analogy for descent. Illustrating this technique, we construct a regularizing term that causes large-batch GD to emulate small-batch SGD, present a model-selection heuristic that depends only on statistics measured before optimization, and exhibit a counter-intuitive loss landscape wherein SGD eternally cycles counterclockwise around a circle of minima.

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1. Introduction

Stochastic gradient descent (SGD) decreases an unknown objective l by performing discrete-time steepest descent on noisy estimates of l . A key question is how the noise affects the final objective value. We connect SGD dynamics to physical scattering theory, thus providing a quantitative and qualitative toolkit for answering this question.

Specifically, we derive a diagram-based formalism for reasoning about SGD via a path integral over possible interactions between weights and data. The formalism permits perturbative analysis, leading to predictions of learning curves for small η . Unlike the continuous-time limits of previous work, this framework models discrete time, and with it, the potential **non-Gaussianity** of noise. We thus obtain new results quantifying the **effect of epoch number, batch size, and momentum** on SGD test loss. We also contrast SGD against popular continuous-time approximations such as ordinary or stochastic differential equations (ODE, SDE).

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Path integrals offer not only quantitative predictions but also an exciting new viewpoint — that of iterative optimization as a **scattering process**. Much as individual Feynman diagrams (see [Dyson \(1949a\)](#)) depict how local particle interactions compose into global outcomes, our diagrams depict how individual SGD updates influence each other before affecting a final test loss. In fact, we import from physics tools such as **crossing symmetries** (see [Dyson \(1949b\)](#)) and **re-normalization** (see [Gell-Mann & Goldberger \(1954\)](#)) to simplify our calculations and refine our estimates. The diagrams’ combinatorial properties immediately yield several precise qualitative conclusions as well, for instance that to order η^2 , **inter-epoch** shuffling does not affect expected test loss.

1.1. Related Work

It was [Kiefer & Wolfowitz \(1952\)](#) who, in uniting gradient descent ([Cauchy, 1847](#)) with stochastic approximation ([Robbins & Monro, 1951](#)), invented SGD. Since the development of back-propagation for efficient differentiation ([Werbos, 1974](#)), SGD and its variants have been used to train connectionist models including neural networks ([Bottou, 1991](#)), in recent years to remarkable success ([LeCun et al., 2015](#)).

Several lines of work quantify the overfitting of SGD-trained networks ([Neyshabur et al., 2017a](#)). For instance, [Bartlett et al. \(2017\)](#) controls the Rademacher complexity of deep hypothesis classes, leading to generalization bounds that are post hoc or optimizer-agnostic. However, since deep networks trained via SGD generalize despite their seeming ability to shatter large sets ([Zhang et al., 2017](#)), one infers that generalization arises from the aptness to data of not only architecture but also optimization ([Neyshabur et al., 2017b](#)). Others have focused on the implicit regularization of SGD itself, for instance by modeling descent via stochastic differential equations (SDEs) (e.g. [Chaudhari & Soatto \(2018\)](#)). However, as explained by [Yaida \(2019\)](#), such continuous-time analyses cannot treat covariance correctly, and so they err when interpreting results about SDEs as results about SGD for finite trainset sizes.

Following [Roberts \(2018\)](#), we avoid making a continuous-time approximation by instead Taylor-expanding around the learning rate $\eta = 0$. In fact, we develop a diagrammatic

method for evaluating each Taylor term that is inspired by the field theory methods popularized by Dyson (1949a). Using this technique, we quantify the overfitting effects of batch size and epoch number, and based on this analysis, propose a regularizing term that causes large-batch GD to emulate small-batch SGD, thus establishing a precise version of the Covariance-BatchSize-Generalization relationship conjectured in Jastrzębski et al. (2018).

While we make rigorous, architecture-agnostic predictions of learning curves, these predictions become vacuous for large η . In particular, while our work does not assume convexity of the loss landscape, it also is blind to large- ηT convergence of SGD. Other discrete-time dynamical analyses allow large η by treating deep generalization phenomenologically, whether by fitting to an empirically-determined correlate of Rademacher bounds (Liao et al., 2018), by exhibiting generalization of local minima **flat** with respect to the standard metric (see Hoffer et al. (2017), Keskar et al. (2017), citetwa18), or by exhibiting generalization of local minima **sharp** with respect to the standard metric (see Stein (1956), Dinh et al. (2017), Wu et al. (2018)). Our work, which makes explicit the dependence of generalization on the underlying metric and on the form of gradient noise, reconciles those latter, seemingly clashing claims.

Others have imported the perturbative methods of physics to analyze descent dynamics: Dyer & Gur-Ari (2019) perturb in inverse network width, employing 't Hooft diagrams to compute deviations of non-infinitely-wide deep learning from Gaussian processes. Meanwhile, (Chaudhari & Soatto, 2018) and Li et al. (2017) perturb in learning rate to second order by approximating noise between updates as gaussian and uncorrelated. This approach does not generalize to higher orders, and, because correlations and heavy tails are essential obstacles to concentration of measure and hence of generalization, it does not model the generalization behavior of SGD. By contrast, we use Penrose diagrams to compute test and train losses to arbitrary order in learning rate, quantifying the effect of non-gaussian and correlated noise. We hence extend Roberts (2018) beyond leading order and beyond 2 time steps, allowing us to compare, for instance, the expected test losses of multi-epoch and single-epoch SGD.

2. Background and Notation

2.1. A Smooth Stage: Tensor Conventions

We adopt summation notation for Greek but not Roman indices, suppressing indices when clear. To expedite dimensional analysis, we follow (Bonnabel, 2013) in regarding the learning rate as an inverse metric $\eta^{\mu\nu}$ that converts a gradient into a displacement. We will use η to raise indices; for example, with C denoting the covariance of gradients, its “trace” will be $C^\mu_\mu = \eta^{\mu\nu} C_{\mu\nu}$. Standard syntactic constraints

then make manifest which expressions transform naturally with respect to optimization dynamics.

We assume that all polynomials of the 0th and higher derivatives of the losses l_n , considered as random functions on weight space, have infinitely differentiable expectations.

Kolář gives a careful introduction to these differential geometric ideas (1993).

2.2. Combinatorial Costumes: Structure Sets

We make use of *structure sets*, i.e. sets S equipped with a preorder \leq and an equivalence relation \sim . The morphisms of structure sets are strictly increasing maps that preserve \sim and its negation. A structure set is *pointed* if it has a unique maximum element and this element forms a singleton \sim -class. The categories \mathcal{S} of structure sets and \mathcal{P} of pointed structure sets enjoy a free-forgetful adjunction \mathcal{F}, \mathcal{G} .

A *diagram* is a rooted tree equipped with an equivalence relation on nodes. We draw the tree of \leq by thin edges, with the root at the far right, and we draw the equivalence relation \sim by fuzzy ties. By reading the tree as a Hasse graph, we see that each diagram D induces a structure set, by abuse of notation also named D . An \mathcal{S} -map from D to $[P] = (\mathcal{G} \circ \mathcal{F})^P(\text{empty diagram})$ is an *ordering* of D , where P counts D ’s equivalence classes. Let, $o(D)$ count orderings of D .

Fong gives a swift introduction to these category theoretic and diagrammatic ideas (2019).

2.3. The Parameterized *Personae*: Forms of SGD

SGD decreases an objective l by updating on smooth, unbiased i.i.d. estimates ($l_n : 0 \leq n < N$) of l . The pattern of updates is determined by a structure set \mathcal{S} whose preorder is a total preorder: for a map $\pi : S \rightarrow [N]$ that induces \sim , we define SGD inductively as $\text{SGD}_{\mathcal{S}}(\theta) = \theta$ when S is empty and otherwise

$$\text{SGD}_{\mathcal{S}}(\theta) = \text{SGD}_{\mathcal{S} \setminus M}(\theta^\mu - \eta^{\mu\nu} \nabla_\nu l_M(\theta))$$

where $M = \min S \subseteq S$ specifies a batch and $l_M = \frac{1}{M} \sum_{m \in M} l_{\pi(m)}$ is a batch average. Since the distribution of l_n is permutation invariant, the non-canonical choice of π does not affect the distribution of output θ s.

Of special interest are structure sets that divide into $M \times B$ many *epochs* each with N/B many disjoint *batches* of size B . An SGD instance is then determined by N, B, M , and an *inter-epoch shuffling scheme*. The cases $B = 1$ and $B = N$ we call *pure SGD* and *pure GD*.

2.4. The Tempting Tool: Taylor Series

Intuitively, each descent step displaces θ by $-\eta \nabla l$ and hence decreases the loss $l(\theta)$ by $\eta(\nabla l)^2$; thus, we expect after T

steps a net decrease of $T\eta(\nabla l)^2$:

$$l(\theta_T) \approx l(\theta_0) - T \cdot \eta \cdot (\nabla l(\theta_0))^2 \quad (1)$$

This intuition fails to capture two crucial facts: **curvature** — that as θ changes during training, so may $\nabla l(\theta)$ — and **noise** — that l_n and l may differ.

To account for noise, we should replace each $(\nabla l_t)(\nabla l)$ by an expectation. If we are interested in train instead of test loss, We get some expectations of the form $(\nabla l_t)(\nabla l_t)$, and hence obtain a different result than for test loss.

To account for curvature, **FILL IN**

3. Diagram Calculus for SGD

3.1. Role of Diagrams

Suppose s is smooth on weight space; e.g. s may be a test loss. We track $s(\theta)$ as θ evolves by SGD:

Key Lemma 1. *The Maclaurin series of $s(\theta_T)$ with respect to η is:*

$$\sum_{(d_t: 0 \leq t < T)} (-\eta)^{\sum_t d_t} \left(\prod_{0 \leq t < T} \frac{(g \nabla)^{d_t}}{d_t!} \Big|_{g=\nabla l_t(\theta)} \right) (s)(\theta_0) \quad (2)$$

In averaging over training sets ($l_t : 0 \leq t < T$) we may factor the expectation of the above product according to independence relations between the l_t . We view various training procedures (e.g. pure GD, pure SGD) as **prescribing different independence relations** that lead to different factorizations and hence to potentially different generalization behavior at each order of η .

An instance of the above product (for $s = l_a$ drawn from a test set and $0 \leq c \leq b < T$) is $-\eta^3 (\nabla l_c \nabla)^2 (\nabla l_b \nabla) l_a$, which is

$$\begin{aligned} & -(\nabla^4 l_c)(\nabla^\mu l_c)(\nabla_\lambda \nabla_\mu \nabla^\nu l_b)(\nabla_\nu l_a) - (\nabla^4 l_c)(\nabla^\mu l_c)(\nabla_\lambda \nabla^\nu l_b)(\nabla_\mu \nabla_\nu l_a) \\ & -(\nabla^4 l_c)(\nabla^\mu l_c)(\nabla_\mu \nabla^\nu l_b)(\nabla_\lambda \nabla_\nu l_a) - (\nabla^4 l_c)(\nabla^\mu l_c)(\nabla^\nu l_b)(\nabla_\lambda \nabla_\mu \nabla_\nu l_a) \end{aligned}$$

To reduce clutter, we adapt the string notation of [Penrose \(1971\)](#). Then, in expectation over (l_c, l_b, l_a) drawn i.i.d.:

$$\dots = \text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \text{diagram 4} \quad (3)$$

$$= \underbrace{2 \text{diagram 1}}_{-2 \mathbb{E}[(\nabla l)(\nabla l)] \mathbb{E}[\nabla \nabla l] \mathbb{E}[\nabla l]} + \underbrace{2 \text{diagram 2}}_{-2 \mathbb{E}[(\nabla l)(\nabla l)] \mathbb{E}[\nabla \nabla l] \mathbb{E}[\nabla \nabla l]} \quad (4)$$

Above, each node corresponds to an l_n (here, red for l_c , green for l_b , blue for l_a), differentiated g times for a degree- g node (for instance, l_b is differentiated thrice in the first diagram and twice in the second). Thin *edges* mark contractions by $-\eta$. Fuzzy *ties* denote correlations by connecting identical loss functions (here, l_c with l_c). The colors are redundant with the fuzzy ties. The value $v(D)$ of a diagram D is the expected value of the corresponding tensor expression.

Crucially, for a fixed, i.i.d. distribution over (l_c, l_b, l_a) , **the topology of a diagram determines its value**. For instance,

$$\text{diagram 1} = \text{diagram 2}$$

Thus follows the simplification of equation 4. We may convert back to explicit tensor expressions, invoking independence between untied nodes to factor the expression. However, as we will see, the diagrams offer physical intuition, streamline computations, and determine useful unbiased estimators of the statistics they represent.

We define a diagram with fuzzy outlines to be the difference between the fuzzy tied and untied versions : $\text{diagram 1} - \text{diagram 2}$.

The recipes for writing down test (or train) losses of SGD and its variants are straight-forward in the diagram notation because they reduce the problem of evaluating the previous dynamical expressions to the problem of counting isomorphic graphs. The more complicated the direct computation, the greater the savings of using diagrams. An appendix provides details and proofs for a variety of situations. For now, we focus on the test loss of SGD.

3.2. Recipe for the Test Loss of SGD

Our results all follow from this theorem and its analogues. Throughout, the d -edged diagrams give the order η^d terms.

Theorem 1. *SGD's expected test loss has a Maclaurin series given as a weighted sum of diagrams:*

$$\sum_{D \in \text{im}(\mathcal{F})} \left(\sum_{f: D \rightarrow \mathcal{F}(S)} \prod_{i \in S} \frac{1}{|f^{-1}(i)|!} \right) v(D) \quad (5)$$

Here, D is (an isomorphism class of) a diagram of form $\mathcal{F}(T)$, $e(D)$ counts D 's edges, and f is a morphism in \mathcal{P} .

Theorem 2. *SGD's expected generalization gap (test loss minus train loss) is formula 5 with $v(D)$ replaced by $\sum_{p \in D/\sim_D} v(D_p)/N$. Here, p ranges through equivalence classes of D , and D_p is D with a fuzzy outline connecting D 's maximal node to p , e.g. $(\text{diagram})_{p= \cdot} = \text{diagram}$.*

In the special case of $B = 1, M = 1$:

Proposition 1. *The order η^d contribution to the expected test loss of one-epoch SGD with singleton batches is:*

$$\frac{(-1)^d}{d!} \sum_{D \in \text{im}(\mathcal{F})} o(D) \binom{N}{P-1} \binom{d}{d_0, \dots, d_{P-1}} v(D) \quad (6)$$

where D ranges over d -edged diagrams whose equivalence classes have sizes $d_p : 0 \leq p \leq P$, with $d_P = 1$ and, without loss, are each antichains.

A P -part, d -edged diagram then contributes $\Theta((\eta N)^d N^{P-d-1})$ to the loss. For example, there are six

diagrams to third order, and they have $(4+2)+(2+2+3)+(1)$ many orderings — see Table 1. Intuitively, ηN measures the **physical time** of descent, and $1/N$ measures **coarseness** of time discretization. So we have a double-series in $(\eta N)^d N^{P-d-1}$, where d counts thin edges and $d+1-P$ counts fuzzy ties; the $P = d+1$ terms correspond to a discretization-agnostic (hence continuous-time, noiseless) ODE approximation to SGD, while $P \leq d$ gives correction terms modeling time-discretization and hence noise.

Corollary 1. *For one-epoch SGD on singleton batches through fixed physical time T : the order N^{-1} deviation of SGD’s test loss from ODE’s is $(T^2 N^{-1}/2)$. The order N^{-2} deviation of SGD’s test loss due to non-gaussian noise is $-(T^3 N^{-2}/6)$.*

For finite N , these effects make SDE different from SGD. SDE also fails to model the correlations between updates in multiepoch SGD. On the other hand, in the $N = \infty$ limit for which SDE matches SGD, optimization and generalization become computationally intractable and trivial, respectively.

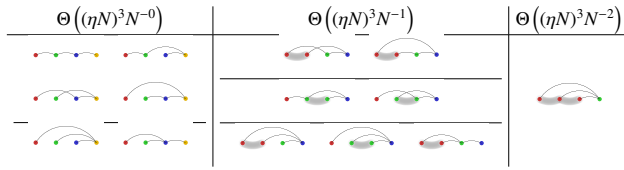


Table 1. Degree-3 scattering diagrams for $B = M = 1$ SGD’s test loss. **Left:** $(d, P) = (3, 3)$. Diagrams for ODE behavior. **Center:** $(d, P) = (3, 2)$. 1st order deviation of SGD away from ODE. **Right:** $(d, P) = (3, 1)$. 2nd order deviation of SGD from ODE with appearance of non-Gaussian statistics.

A quick combinatorial argument shows:

Corollary 2. *To order η^2 , inter-epoch shuffling doesn’t affect SGD’s expected test loss.*

Indeed, for any inter-epoch shuffling scheme:

Proposition 2. *To order η^2 , the test loss of SGD — on N samples for M epochs with batch size B dividing N and with any shuffling scheme — has expectation*

$$\begin{aligned} & + MN \text{ (diagram)} + MN \left(MN - \frac{1}{2} \right) \text{ (diagram)} \\ & + MN \left(\frac{M}{2} \right) \text{ (diagram)} + MN \left(\frac{M - \frac{1}{B}}{2} \right) \text{ (diagram)} \end{aligned}$$

Corollary 3. *To order η^2 , one-epoch SGD has $\left(\frac{M-1}{M} \right) \left(\frac{B+1}{B} \right) \left(\frac{N}{2} \right)$ less test loss than M -epoch SGD with learning rate η/M .*

Given an unbiased estimator \hat{C} of gradient covariance, we may get GD to mimic SGD:

Corollary 4. *The expected test loss of pure SGD is, to order η^2 , less than that of pure GD by $\left(\frac{M}{2} \right) \left(\frac{N-1}{N} \right)$. Moreover, GD on a modified loss $\tilde{l}_n = l_n + \left(\frac{N-1}{4N^2} \right) \hat{C}_v^v(\theta)$ has an expected test loss that agrees with SGD’s to second order.*

3.3. Renormalization

An important idea is that of renormalization, i.e. the summarization of myriad small-scale interactions into an effective large-scale theory. We can use this two ways: (A) to refine our computations if we know the hessian; (B) to refine our computations if we know the “effective propagator”. Lorem ipsum dolor sit amet, consectetur adipiscing elit...

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3.4. Descent as Scattering

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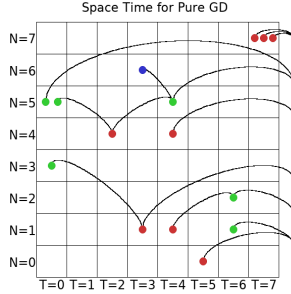


Figure 1. **a** Some diagrams embedded in pure GD's spacetime. Top three give order η^3 , η^5 , and η^1 contributions. Bottom three all give identical order η^2 contributions to test loss, including the bottom-most, by crossing symmetry. **b c d**

4. Consequences and Applications

4.1. Vanilla SGD

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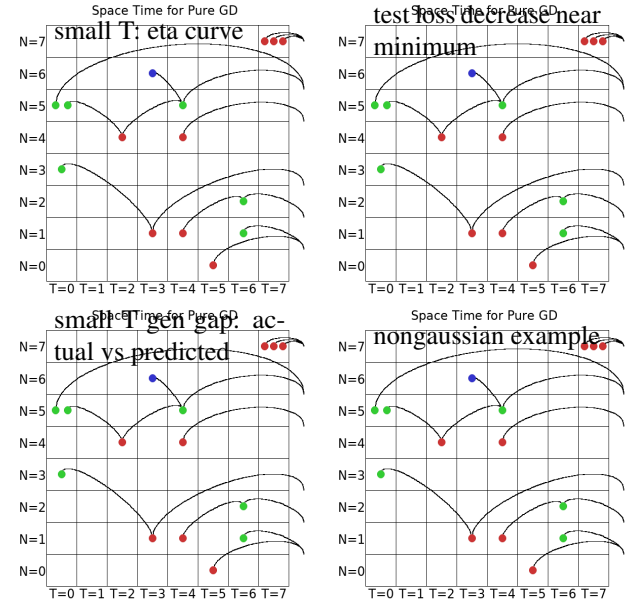


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4.2. Emulating Small Batches with Large Ones

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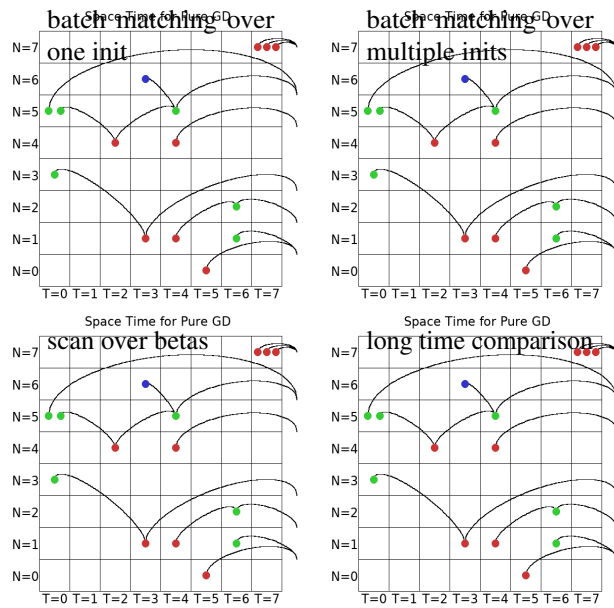


Figure 3. Lorem ipsum dolor sit amet, consectetur adipiscing elit...

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4.3. Analyzing Second Order Methods

We demonstrate how our approach extends to more sophisticated optimizers by analyzing momentum and a hessian-based method.

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Now consider a hessian-based update parameterized by a scalar λ :

$$\theta \leftarrow \theta - (\eta^{-1} + \lambda \nabla \nabla l_t(\theta))^{-1} \nabla l_t(\theta)$$

invhess

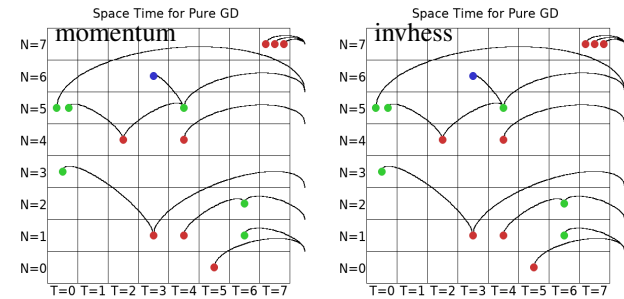


Figure 4. Lorem ipsum dolor sit amet, consectetur adipiscing elit...

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4.4. Epochs and Overfitting

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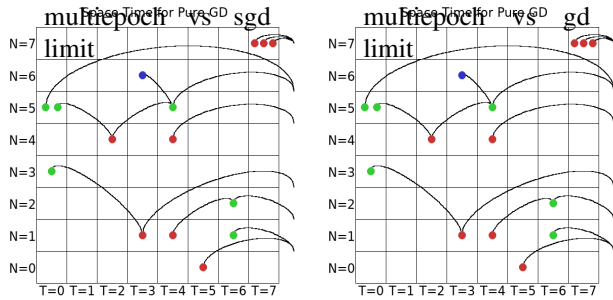


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4.5. Myopic Model Selection

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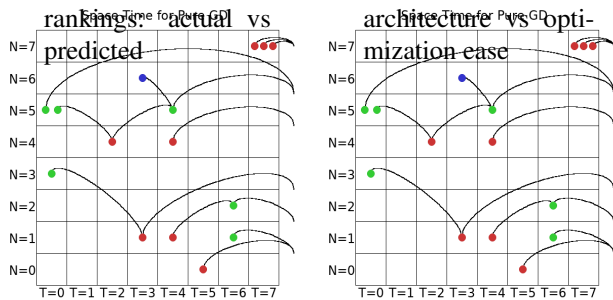


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4.6. Comparison to Continuous Time

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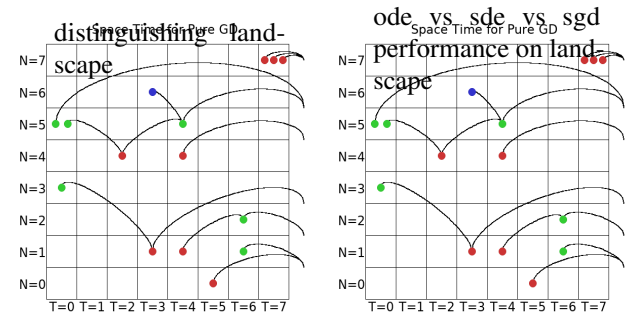


Figure 7. Lorem ipsum dolor sit amet, consectetur adipiscing elit...

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4.7. Thermodynamic Engine

We clarify Lorem ipsum dolor sit amet, consectetur adipiscing elit...

...sunt in culpa qui officia deserunt mollit anim id est laborum. We constructed a counter-intuitive loss landscape

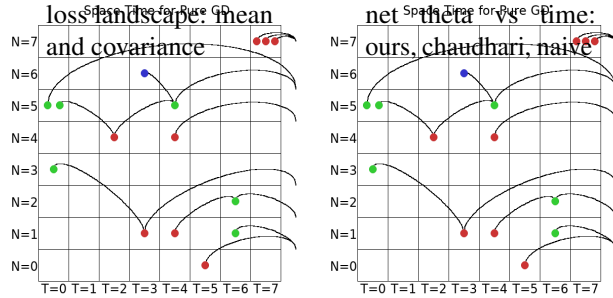


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wherein, for arbitrarily small learning rates, SGD cycles counterclockwise around a circle of minima. Our mechanism differs from that discovered by Chaudhari & Soatto (2018) discuss the thermodynamic significance of both

5. Conclusion

We presented a novel diagrammatic tool for analyzing gradient-based descent. Via a new regularizing term, we showed that **large-batch GD can be made to emulate small-batch SGD**, thus completing a project suggested by Roberts (2018). This is significant because, while small batch sizes can lead to better generalization (Bottou, 1991), modern infrastructure increasingly rewards large batch sizes (Goyal et al., 2018). We showed also that in multi-epoch SGD, inter-epoch shuffling induces only a 3rd order effect on test loss. Intuitively, we proved that **the hessian matters asymptotically more than shuffling order**.

The diagram method is also a rich source of intuitions and physical analogies. For example, it offers a clearer understanding of the empirically verified limit cycles found in Chaudhari. As our physical analogy emphasizes the underlying metric, it reconciles competing views of whether

sharp or flat minima generalize. Further exploration of this bridge to particle physics, especially within the framework of renormalization theory, pose a promising direction for future research.

Variances

5.1. Acknowledgements

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A. Derivation of Diagram Rules

5.2. Dyson Series for Iterative Optimizers

If a density ρ governs a point θ in weight space, then after a sequence of updates $\theta \mapsto \theta - \eta^{\mu\nu} \nabla_{\mu} l(\theta)$ on losses ($l_t : 0 \leq t < T$), the following density (up to an error term whose Maclaurin series vanishes; all perturbative results will implicitly carry such terms) will govern the new point:

$$\exp\left(+\eta^{\mu\nu} \nabla_{\mu} l_{T-1}(\theta) \nabla_{\nu}\right) \cdots \exp\left(+\eta^{\mu\nu} \nabla_{\mu} l_0(\theta) \nabla_{\nu}\right) \rho \quad (7)$$

or $\prod \exp(+\eta \nabla l / \nabla) \rho$ for short. The exponent above is a linear operator that acts on a space of sufficiently smooth maps; in particular, the ∇_{ν} does not act on the $\nabla_{\mu} l(\theta)$ with which it pairs. Integrating by parts, we write the expectation over initial values after T steps of a function s of weight space (e.g. s may be test loss) as:

$$\int_{\theta} \rho(\theta) \left(\prod_{0 \leq t \leq T} \exp\left(-\eta^{\mu\nu} \nabla_{\mu} l(\theta) \nabla_{\nu}\right) s \right) (\theta) \quad (8)$$

Since the exponentials above might not commute, we may not compose the product of exponentials into an exponential of a sum. We instead compute an expansion in powers of η .

Setting the initialization $\rho(\theta) = \delta(\theta - \theta_0)$ to be deterministic, and labeling as θ_t the weight after t steps, we find:

$$s(\theta_T) = \sum_{0 \leq d < \infty} (-\eta)^d \sum_{\substack{(d_t: 0 \leq t < T) \\ \sum_t d_t = d}} \left(\prod_{0 \leq t < T} \frac{(\nabla l_t(\theta) \nabla)^{d_t}}{d_t!} \right) s(\theta_0) \quad (9)$$

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B. Tutorial on Diagram Rules

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C. Derivations of Perturbative Results

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D. Diagram Rules vs Direct Perturbation

Diagram methods from Stueckelberg to Peierls have flourished in physics because they enable swift computations and immediate intuition that would otherwise require laborious algebraic manipulation. We exhibit how our diagram formalism likewise streamlines analysis of descent by comparing direct perturbation and the new formalism on some sample problems.

One-Epoch, Singleton-Batch SGD Test Loss to 3rd Order

For single-epoch SGD with singleton batches, we sum all relevant diagrams through order 3; the coefficients 4, 2; 2, 2, 3; 1 come from counting the elements of Table 1, and the other coefficients come from analogous tables. This yields:

$$\begin{aligned} \mathcal{L}_{\text{test}}^{\text{SGD}}(T, \eta) \in & \\ & + \frac{\eta}{1!} \binom{T}{1} \left(\text{diagram 1} \right) \\ & + \frac{\eta^2}{1!1!} \binom{T}{2} \left(2 \cdot \text{diagram 2} \right) + \frac{\eta^2}{2!} \binom{T}{1} \left(\text{diagram 3} \right) \\ & + \frac{\eta^3}{1!1!1!} \binom{T}{3} \left(4 \cdot \text{diagram 4} + 2 \cdot \text{diagram 5} \right) \\ & + \frac{\eta^3}{2!1!} \binom{T}{2} \left(2 \cdot \text{diagram 6} + 2 \cdot \text{diagram 7} + 3 \cdot \text{diagram 8} \right) \\ & + \frac{\eta^3}{3!} \binom{T}{1} \left(\text{diagram 9} \right) + o(\eta^3)^1 \end{aligned}$$

Test Loss of One-Epoch SGD vs Multi-Epoch SGD

Generalization Gap of a Hessian-Based Method

Test Loss of SGD vs GD

The generalization gap $\mathcal{L}_{\text{gen}}^{\text{SGD}} = \mathcal{L}_{\text{test}}^{\text{SGD}} - \mathcal{L}_{\text{train}}^{\text{SGD}}$ is suppressed by a factor $1/N$ ($N \leq T$):

$$\begin{aligned} N \cdot \mathcal{L}_{\text{gen}}^{\text{SGD}}(T, \eta) \in & \\ & + \eta \binom{T}{1} \left(\text{diagram 10} \right) + \eta^2 \binom{T}{2} \left(\text{diagram 11} + \text{diagram 12} \right) \\ & + \frac{\eta^2}{2!} \binom{T}{1} \left(\text{diagram 13} - \text{diagram 14} \right) + o(\eta^2) \\ = & + \eta^{\lambda\mu} \binom{T}{1} C_{\lambda\mu} + \eta^{\lambda\mu} \eta^{\nu\xi} \binom{T}{2} \left(\frac{1}{2} G_{\lambda} \nabla_{\mu} C_{\nu\xi} + C_{\mu\nu} H_{\xi\lambda} \right) \\ & + \frac{\eta^{\lambda\mu} \eta^{\nu\xi}}{2!} \binom{T}{1} \left(\langle l_{\lambda} l_{\nu} l_{\mu\xi} \rangle - (G_{\mu} G_{\nu} + C_{\mu\nu}) H_{\xi\lambda} \right) + o(\eta^2) \end{aligned}$$

The leading order term is $N \cdot \mathcal{L}_{\text{gen}}^{\text{SGD}}(T, \eta) \approx \eta T \left(\text{diagram 15} - \text{diagram 16} \right) = T \cdot \eta^{\lambda\mu} C_{\lambda\mu}$, where C is the covariance of gradients. We thus recover a main result of Roberts (2018).

¹ We use little- $o(\eta^d)$ instead of big- $O(\eta^{d+1})$ to avoid specializing to analytic loss landscapes. Error terms depend on the loss landscape and on T . When gradients are uniformly bounded, the T dependence is at most linear.

E. The η -Series' Domain of Convergence

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G. Generalized Bessel Factors

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F. Autonomous ODE Fitting

We fit a . These have the benefit In particular:

$$y'(t) = a \rightarrow$$

$$y(t) = y(0) + at \quad (10)$$

$$y'(t) = by + a \rightarrow y(t) = (y(0) - (a/b)) \exp(bt) + (a/b) \quad (11)$$

$$y'(t) = cy^2 + by + a \rightarrow \quad (12)$$

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