

1. *Forward problem* (initial boundary value problem for wave equation

$$\begin{aligned}\frac{1}{c^2(x)}u_{tt} &= \Delta u \text{ in } \Omega \times [0, 2T], \\ u|_{t=0} &= 0, \quad u_t|_{t=0} = 0 \text{ in } \Omega, \\ u_n|_{\partial\Omega} &= f,\end{aligned}$$

where  $\Omega$  is a disk in  $\mathbb{R}^2$

$$\Omega = \{x \in \mathbb{R}^2 \mid |x| < r\},$$

$c$  is the speed of sound, and  $f$  is a Neumann boundary source (control). Control  $f$  runs over control space  $F$ . It is enough assume that  $F = C_0^\infty(S \times [0, T])$ , where  $S = \{x \in \mathbb{R}^2 \mid |x| = r\}$ . The solution will be denoted by  $u^f$ .

For solving forward problem we use Matlab, and we choose  $f$  as a rectangular pulse (depend on time only) or as Ricker pulse.

2. *Inverse problem* is finding  $c$  by the inverse data:

$$u^f|_{S \times [0, 2T]}, \quad f \in F.$$

We will consider a bit another problem: to find eikonal function

$$\tau(x) = \text{dist}(x, S),$$

where distance means the travel time from the point  $x$  to the boundary  $S$ . Notice, that function  $\tau$  determines  $c$  since it satisfies eiconal equation

$$|\nabla \tau| = 1/c.$$

3. *Control problem.*

Let  $u^h$  be a solution to the forward problem (corresponding to the control  $h$ ) and  $v$  be an arbitrary solution to the wave equation. Then one can get differential equality

$$\frac{1}{c^2}(vu_t^f - u^f v_t)_t = \text{div}(v \nabla u^f - u^f \nabla v).$$

After integration over  $\Omega \times [0, s]$  we have integral equality

$$\int_{\Omega^s} \frac{1}{c^2}(vu_t^f - u^f v_t)dx = \int_S (vf - u^f v_n)dSdt,$$

where  $S = \partial\Omega = \{x \in \mathbb{R}^2 \mid |x| = r\}$ ,  $v_n$  is the normal derivative and  $\Omega^s$  is the set filled with wave  $u$  up to the moment  $s$ . More precisely

$$\Omega^s = \{x \in \Omega \mid \tau(x) \leq s\},$$

where  $\tau$  is the eikonal,

$$\tau(x) = \text{dist}(x, S),$$

where distance means the travel time from the point  $x$  to the boundary  $S$ . If we put  $v$  the even part of  $u^g$  with respect to  $t = T$

$$v(x, t) = u_+^g(x, t) = \frac{1}{2}(u^g(x, t) + u^g(x, 2T - t))$$

then

$$\int_{\Omega^s} \frac{1}{c^2} u_t^h u^g(., T) dx = \int_S [u_+^g h - u^h(g_n)_+] dS dt. \quad (1)$$

Now, if we put  $v = \varphi$ , where  $\varphi$  is arbitrary harmonic function (does not depend on time), then we have

$$\int_{\Omega^s} \frac{1}{c^2} u_t^h(., T) \varphi dx = \int_S [\varphi h - u^h \varphi_n] dS dt. \quad (2)$$

Let

$$f(., t) = \int_0^t h(., t') dt' \stackrel{def}{=} Ih(t).$$

Then 1 and 2 and can be rewritten as

$$C(f, g) \stackrel{def}{=} \int_{\Omega^s} \frac{1}{c^2} u^f u^g(., T) dx = \int_{S \times [0, T]} [u_+^g I f - I u^f(g_n)_+] dS dt, \quad (3)$$

$$\Phi(f) \stackrel{def}{=} \int_{\Omega^s} \frac{1}{c^2} u^f(., T) \varphi dx = \int_S [\varphi I f - I u^f \varphi_n] dS dt. \quad (4)$$

Thus we have bilinear functional  $C$  and linear functional  $\Phi$ . Let us consider equation

$$A(f, g) = \Phi(f), \quad \forall f \in F^s \quad (5)$$

with respect to  $g \in F^s$ . By definition this is equivalent to the equation

$$\int_{\Omega^s} \frac{1}{c^2} u^f u^g(., T) dx = \int_{\Omega^s} \frac{1}{c^2} u^f(., T) dx.$$

But the last equality holds if and only if

$$u^g(., T) = \varphi.$$

It follows from the approximate boundary controllability, that is any function from  $L^2(\Omega^s)$  can be approximated by a wave  $u^f(., T)$ . Another words the following equality holds

$$\overline{U^s} = L^2(\Omega^s)$$

where

$$U^s = \{u^f(., T) \mid f \in F^s\} \subset L^2(\Omega^s)$$

is the reachable set and  $F^s \subset F$  is the subspace of the control space that consist of controls acting only during period  $[T - s, T]$ . Such controls generate waves which fill the domain  $\Omega^s$  up to the moment  $T$ . Notice, that equation  $A(f, g) = \Phi(f)$  is not solvable in general, but it may be shown it is dense solvable that

is there exist a sequence of controls  $\{g_n\}$  such that  $A(f, g_n) \rightarrow \Phi(f)$ . Formally one can construct a wider control space in which the equation  $A(f, g) = \Phi(f)$  will be solvable. So, let  $g$  is the solution of this equation. It implies the equality  $u^g(., T) = \varphi$ . Let us put

$$\begin{aligned}\varphi(x) &= \frac{1}{|x - x_0|}, \quad n = 3, \\ \varphi(x) &= \ln|x - x_0|, \quad n = 2.\end{aligned}\tag{6}$$

These functions are fundamental solutions to the Laplace equations (up to some constants) and if  $c$  then they are harmonic in  $\Omega^s$  if  $x_0 \notin \Omega^s$ .

How we are going to find  $\tau(x_0)$ ? Let us fix a point  $x_0$  and choose  $\varphi$  as in 6. If  $x_0 \notin \Omega^s$  then after solving 5 we get control  $g$  that generates

$$u^g(., T) \approx \varphi \text{ in } \Omega^s.$$

Then the value of  $\Phi(g)$  equals

$$\Phi(g)(s) = \int_{\Omega^s} \frac{1}{c^2} \varphi^2 dx.$$

This value tends to  $\infty$  when  $s$  tends to  $\tau(x_0)$ . At this moment we get blow-up

$$\Phi(g)(\tau(x_0)) = \infty.$$