Dynamics

Relative Motion for Rotating Frame

$$\underline{v} = \frac{d\underline{r}}{dt} = \underline{v}_{o'} + \underline{v}_r + \underline{w} \times \underline{r}$$

$$\frac{d\underline{v}}{dt} = \underline{a}_r + 2\underline{\omega} \times \underline{v}_r + \frac{d\underline{\omega}}{dt} \times \underline{r} + \underline{a}_{o'} + \underline{\omega} \times (\underline{\omega} \times \underline{r})$$

Polar Coordinates:

$$\underline{v} = \dot{r}\underline{e}_r + r\dot{\theta}\underline{e}_\theta$$

$$\underline{a} = (\ddot{r} - r\dot{\theta}^2)\underline{e}_r + (r\ddot{\theta} + 2\dot{r}\dot{\theta})\underline{e}_{\theta}$$

Binet's Formula:

For $\underline{a} \times \underline{r} = 0$ we have $a_{\theta} = 0 \implies r^2 \dot{\theta} = c$ which gives

$$a_r = -\frac{c^2}{r^2} \left[\frac{d^2}{d\theta^2} \left(\frac{1}{r} \right) + \frac{1}{r} \right]$$

Gravity

$$m\underline{a} = -\frac{GmM}{r^2}\hat{\underline{r}} = -\frac{m\mu}{r^2}\hat{\underline{r}}$$

Total Mechanical Energy

$$E = \frac{1}{2}mv^2 - \frac{m\mu}{r}$$

Angular Momentum:

$$\underline{L} = \underline{r} \times m\underline{v}$$

$$|\underline{L}| = mr^2\dot{\theta} \implies \text{conserved}$$

Keplerian Motion

Two-Body Dynamics

Gravitational Forces:

$$m_i \ddot{\underline{r}}_i = \pm \frac{Gm_i m_j}{r^2} \hat{\underline{r}}_i$$

Centre-of-Mass:

$$\underline{r}_C = \frac{m_1\underline{r}_1 + m_2\underline{r}_2}{m_1 + m_2}$$

Sum of forces acting on centre of mass $= 0 \implies \underline{r}_C = \underline{r}_0 + \underline{r}'t$ Approximate $\mu \approx Gm_1$ gives:

$$\ddot{\underline{r}} = -\frac{\mu}{r^2}\hat{\underline{r}}$$

Total energy is conserved, take $\underline{\ddot{r}}\cdot\underline{r}$ to prove (d/dt)E=0, since:

$$\frac{d}{dt}\left(\frac{1}{2}v^2\right) = \ddot{\underline{r}} \cdot \dot{\underline{r}} = \frac{d}{dt}\left(\frac{G(m_1 + m_2)}{r}\right)$$

Equations of Motion

Use chain rule and $r^2\dot{\theta} = l$ to get:

$$\frac{dr}{dt} = \frac{l}{r^2} \frac{dr}{d\theta}$$

Change variable to u = 1/r:

$$\frac{dr}{dt} = -l\frac{du}{d\theta}$$

Gives:

$$\frac{d^2r}{dt^2} = -l^2u^2\frac{d^2u}{d\theta^2} \qquad \text{and} \qquad -r\dot{\theta}^2 = -l^2u^3$$

Equations of motion in polar coordinates are:

$$\ddot{r} - r\dot{\theta}^2 = -\frac{\mu}{r^2}$$
 and $r\ddot{\theta} + 2\dot{r}\dot{\theta} = \frac{1}{r}\frac{d}{dt}(r^2\dot{\theta}) = 0$

Transforming for u:

$$-l^2 u^2 \frac{d^2 u}{d\theta^2} - l^2 u^3 = -\mu u^2$$

$$\frac{d^2u}{d\theta^2} + u = \frac{\mu}{l^2}$$

Finally change to $z = u - \mu/l^2$ to get:

$$\frac{d^2z}{d\theta^2} + z = 0 \implies z = A\cos(\theta - \theta_0)$$

Which gives polar equation of conic section for r:

$$r(\theta) = \frac{l^2/\mu}{1 + \frac{Al^2}{\mu}\cos(\theta)}$$

The energy equation is:

$$E = \frac{1}{2}m(\dot{r}^2 + r^2\dot{\theta}^2) - U(r)$$

Which can be solved by integrating w.r.t r, changing variable from t to θ , and substituting $U(r) = m\mu/r$. This gives:

$$r = \frac{p}{\pm 1 + e\cos(\theta - \theta_0)}$$

where

$$p = \frac{c^2}{|\mu|} = a(1 - e^2)$$
 and $e = \sqrt{\frac{2Ec^2}{m\mu^2} + 1}$

so r describes an ellipse with semimajor axis a and eccentricity e, pericentre at $\theta = 0$ and apocentre at $\theta = \pi$:

$$r = \frac{a(1 - e^2)}{1 + e\cos(\theta - \theta_0)}$$

Orbital Parameters

$$a = \frac{r_p + r_a}{2}$$
 and $e = \frac{r_a - r_p}{r_a + r_p}$

so

$$r_p = a(1-e)$$
 and $r_a = a(1+e)$

Comparing equations for r:

$$l^2 = \mu a (1 - e^2)$$

$$v_p = \sqrt{\frac{\mu(1+e)}{a(1-e)}}$$
 and $v_a = \sqrt{\frac{\mu(1-e)}{a(1+e)}}$

Conservation of energy:

$$E = -\frac{\mu}{2a} \implies v^2 = \mu \left(\frac{2}{r} - \frac{1}{a}\right)$$

Kepler's Laws

1st Law: Orbits are elliptical with attracting body at focus of ellipse

2nd Law: Radial vector sweeps out equal areas in equal time

$$\implies T = 2\pi \sqrt{\frac{a^3}{\mu}}$$

3rd Law: Orbital period squared \propto semi-major axis cubed:

Kepler's Equation

Mean motion

$$n = \sqrt{\frac{\mu}{a^3}}$$

Mean anomoly

$$M = n(t - \tau)$$

Kepler's equation is

$$M = E - e \sin E$$

and use trig identity to get θ from E

$$\tan\left(\frac{\theta}{2}\right) = \sqrt{\frac{1+e}{1-e}}\tan\left(\frac{E}{2}\right)$$

This must be solved iteratively, start with guess $E_0 = M$, then find correction

$$\Delta E = \frac{M - E_0 + e \sin E_0}{1 - e \cos E_0}$$

Then repeat, using $E_i = E_{i-1} + \Delta E_{i-1}$ until $\Delta E_i \approx 0$

Position and Velocity from Orbital Elements

From orbit equation

$$\dot{r} = \frac{a(1 - e^2)}{(1 + e\cos\theta)^2} e\sin\theta \dot{\theta} = r \frac{e\sin\theta}{1 + e\cos\theta} \dot{\theta}$$

Flight path angle

$$\tan \gamma = \frac{v_r}{v_\theta} = \frac{\dot{r}}{r\dot{\theta}} = \frac{e\sin\theta}{1 + e\cos\theta}$$

To get in plane position and velocity at time t, solve Kepler's equation for τ time of passage from pericentre, then compute orbit radius and thus cartesian positions.

$$x = r\cos\theta$$
 and $y = r\sin\theta$

Velocity components:

$$v_r = \frac{r^2 e \sin \theta}{p} \dot{\theta}$$
 and $v_\theta = \frac{l}{r} = r \dot{\theta}$

Gives cartesian velocities:

$$\dot{x} = -\sqrt{\frac{\mu}{p}}\sin\theta$$
 and $\dot{y} = \sqrt{\frac{\mu}{p}}(e + \cos\theta)$

Orbit Transfers

Hohmann Transfer

Minimum energy transfer between two circular, coplanar orbits is an ellipse with semimajor axis $a = \frac{1}{2}r_1 + r_2$. The change in velocities required are the difference between the velocities of the transfer ellipse at pericentre and apocentre, and the velocities of the circular orbits at r_1 and r_2

$$\Delta v_1 = \sqrt{\frac{2\mu}{r_1} - \frac{2\mu}{r_1 + r_2}} - \sqrt{\frac{\mu}{r_1}}$$

$$\Delta v_2 = \sqrt{\frac{\mu}{r_2}} - \sqrt{\frac{2\mu}{r_2} - \frac{2\mu}{r_1 + r_2}}$$

Parabolic Transfer

For a parabolic transfer where the orbiting body escapes circular orbit at r_1 to infinity, and is then captured from infinity into orbit at r_2

$$\Delta v_1 = (\sqrt{2} - 1)\sqrt{\frac{\mu}{r_1}}$$

$$\Delta v_2 = (\sqrt{2} - 1)\sqrt{\frac{\mu}{r_2}}$$

Single Impulse Manoeuvres

Raise apocentre:

$$\Delta v = \sqrt{-\frac{\mu}{a_2} + \frac{2\mu}{r_p}} - \sqrt{-\frac{\mu}{a_1} + \frac{2\mu}{r_p}}$$

Change inclination:

$$\Delta v = \sqrt{v_1^2 + v_2^2 - 2v_1v_2\cos\Delta i}$$

Multiple impulse manoeuvres can be used to reduce required delta v, by first raising apocentre and performing multiple burns.

Patched Conic

In Earth-centred frame, velocity to escape is:

$$\Delta v_{esc} = \sqrt{\frac{2\mu_{Earth}}{r_p} + v_{\infty,1}^2} - \sqrt{\frac{2\mu_{Earth}}{r_p} - \frac{\mu_{Earth}}{a_p}}$$

In sun centre frame, $v_{\infty,1}$ must match the first Δv for transfer. E.g. Δv_1 of Hohmann transfer, and similar for capture.

Rocket Equation

Specific impulse

$$I_{sp} = \frac{F}{\dot{m}g_0}$$

Effective exhaust velocity

$$c = I_{sp}g_0$$

Rocket equation

$$(m - dm)dv = -dmv_e$$

$$\Delta v = v_e \ln \frac{m_i}{m_f}$$

$$m_f = m_i e^{-\frac{\Delta v}{v_e}}$$

Aerobraking

$$D = \frac{1}{2}\rho C_D A v^2$$
$$L = \frac{1}{2}\rho C_L A v^2$$

Flight path angle γ is measured to horizontal

$$m\frac{dv}{dt} = -D - mg\sin\gamma$$

$$mv\frac{d\gamma}{dt} = L - mg\cos\gamma$$

Atmosphere density is exponential model, relative to a scale height H and base density ρ_0

$$\rho(h) = \rho_0 \exp\left[-\frac{h}{H}\right]$$

For aerocapture, assume gravity terms negligible

$$\frac{1}{v}\frac{dv}{d\gamma} = -\frac{D}{L} = -\frac{C_D}{C_L}$$

Integrate to get

$$\ln\left(\frac{v_2}{v_1}\right) = -\frac{C_D}{C_L}(\gamma_2 - \gamma_1)$$

$$\Delta v = v_1 \left(1 - \exp\left[2\frac{C_D}{C_L} \gamma_1 \right] \right)$$

Vertical Dynamics

$$\frac{dh}{dt} = v \sin \gamma = \frac{dh}{d\gamma} \frac{d\gamma}{dt} = \frac{dh}{d\gamma} \frac{L}{mv}$$

Substituting for L and using density function

$$\frac{dh}{d\gamma} = \frac{2m\sin\gamma}{\rho_0 C_L A} \exp\left[\frac{h}{H}\right]$$

Integrate to find h_{min}

$$\int_{h_1}^{h_{min}} \exp\left(-\frac{h}{H}\right) = \frac{2m}{\rho_0 C_L A} \int_{\gamma_1}^0 \sin \gamma d\gamma$$

Note that $h_1 >> H$ and so

$$h_{min} = H \ln \left(\frac{\rho_0 C_L A H}{2m(1 - \cos \gamma_1)} \right)$$