### **Newtonian Mechanics**

 $\mathbf{N1}$  - If no forces act on a body, it remains at rest or moves with constant velocity:  $\dot{\underline{v}} = 0$ 

$$\mathbf{N2} - \underline{\dot{p}} = \underline{F}$$

$$\mathbf{N3} - F_{ab} = -F_{ba}$$

$$\underline{L} \equiv \underline{r} \times p$$

$$W_{BA} \equiv \int_{A}^{B} \underline{F} \cdot d\underline{r} = T_{B} - T_{A}$$

Orbits (cylindrical polars):

$$\begin{split} \underline{e_r} &= \cos \phi \underline{i} + \sin \phi \underline{j} \\ \underline{e_\phi} &= -\sin \phi \underline{i} + \cos \phi \underline{j} \\ \dot{\underline{r}} &= \dot{r} \underline{e_r} + r \dot{\phi} \underline{e_\phi} \\ E &= \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\phi}^2) + V(r) \end{split}$$

## Newton to Lagrange

Holonomic constraint is an algebraic relation between coordinates:

$$f(r_a, r_b, ..., r_N; t) = 0$$

For system with N cartesian coordinates  $x_i$ , M constraints, and 3N-M generalised coordinates  $q_i$ , and  $x_i=x_i(\{q\},t)$ 

Virtual displacement:

$$\delta x_i = \sum_i \frac{\partial x_i}{\partial q_j} \delta q_j + 0$$

Generalised forces:

$$Q_j = \sum_i F_i \frac{\partial x_i}{\partial q_j}$$

for a function  $f = f(\{q\}, \{\dot{q}\}, t)$ 

$$df = \sum_{j} \frac{\partial f}{\partial q_{j}} + \sum_{j} \frac{\partial f}{\partial \dot{q}_{j}} d\dot{q}_{j} + \frac{\partial f}{\partial t} dt$$

for a function  $f = f(\{q\}, \{\not q\}, t)$  cancellation of dots

$$\frac{\partial \dot{f}}{\partial \dot{q}_j} = \frac{\partial f}{\partial q_j}$$

commuting derivatives:

$$\frac{d}{dt}\left(\frac{\partial f}{\partial q_j}\right) = \frac{\partial}{\partial q_j}\left(\frac{df}{dt}\right)$$

Lagrange's equations (general form):

$$\frac{d}{dt} \left( \frac{\partial T}{\partial \dot{q}_j} \right) - \frac{\partial T}{\partial q_j} = Q_j$$

or

$$\frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial q_j} \right) - \frac{\partial \mathcal{L}}{\partial q_j} = 0$$

$$(\hat{c}) t - T((c), (\hat{c}), t) \quad V((c), t) = 0$$

where  $\mathcal{L}(\{q\}, \{\dot{q}\}, t) = T(\{q\}, \{\dot{q}\}, t) - V(\{q\}, t)$ 

### Calculus of Variations

Euler's Equation: for functional of the form

$$I[y(s)] = \int_a^b F(y(s), y'(s), s) ds$$

$$\frac{\partial F}{\partial y} - \frac{d}{ds} \left( \frac{\partial F}{\partial y'} \right) = 0$$

with first integrals

$$\frac{\partial F}{\partial v'} = \text{constant}$$
 if F does not depend on y

$$y'\frac{\partial F}{\partial y'} - F = \text{constant}$$
 if F does not depend on s

Hamilton's Principle:

$$\delta S = 0$$
 for the action  $S = \int_{t_1}^{t_2} \mathcal{L}dt$ 

### **Energy Function**

The energy function (h) is usually conserved, whereas the energy (E = T + V) might not be. h is given by:

$$h = \sum_{i} \frac{\partial L}{\partial \dot{q}_{i}} \dot{q}_{i} - L = \sum_{i} p_{i} \dot{q}_{i} - L$$

h and L are the legendre transforms of each other on the variables  $p_i$  and  $\dot{q}_i$ :

$$p_i = \frac{\partial L}{\partial \dot{q}_i}, \dot{q}_i = \frac{\partial h}{\partial p_i}$$

#### Relativistic Particle

The lagrangian for a free relativistic particle can be found by minimising the action S, then requiring  $\frac{\partial L}{\partial \dot{q}_i} = \gamma m \dot{q}_i$ 

$$S = \int_{\tau_1}^{\tau_2} \epsilon d\tau = \epsilon \int_{t_1}^{t_2} \frac{1}{dt/d\tau} dt = \epsilon \int_{t_1}^{t_2} \frac{dt}{\gamma(\dot{q})}$$

Finding  $\epsilon$  from relativistic momentum:

$$p_i = \gamma m \dot{q}_i = \frac{\partial L}{\partial \dot{q}_i} = -\frac{\epsilon \gamma \dot{q}_i}{c^2}$$

And we have  $\epsilon = -mc^2$ , so the lagrangian is:

$$L = -\frac{mc^2}{\gamma}$$

# Lagrangian of Rotation

The lagrangian for a rotating symmetric top is given by:

$$\mathcal{L} = \frac{1}{2} [A(\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) + C(\dot{\phi} \cos \theta + \dot{\psi})^2] - Mgl \cos \theta$$

With  $A=I_1=I_2$  and  $C=I_3$ . This won't need to be derived in an exam - it'll be given or will need to be quoted.