

# **Detectors and Physics at a Future Linear Collider**

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## Abstract

An electron-positron linear collider is an option for future large particle accelerator projects. Such a collider would focus on precision tests of the higgs boson properties. This thesis describes several studies related to the optimisation of high granular calorimeters. Three main areas were covered.

The performance of photon reconstruction is improved. Photon reconstruction algorithms were developed within PandoraPFA, a world-leading pattern-recognition software for particle flow calorimetry. A sophisticated pattern recognition algorithm was implemented, which uses the topological properties of electromagnetic showers to identify photon candidates and separate them from nearby particles. It performs clustering of the energy deposits in the detector, followed by topological characterisation of the clusters, with the results being considered by a multivariate likelihood analysis. This algorithm leads to a significant improvement in the reconstruction of both single photons and multiple photons in high energy jets.

Reconstruction and classification of tau lepton decay modes were studied. Tau decay products, such as photons, were reconstructed as separate entities. Utilising high granular calorimeters, the resolution of energy and invariant mass of the tau decay products is improved. A hypothesis test was performed for expected decay final states. A multivariate analysis was trained to classify decay final states with a data-driven machine learning method. The performance of tau decay classification is used for the electromagnetic calorimeter optimisation at the ILC or CLIC.

Sensitivity of higgs couplings at CLIC was studied, using simulated double Higgs boson production. Algorithms were developed to identify

isolated high energy leptons, and results were fed into a multivariate analysis. The study was done for two CLIC energy scenarios. This sensitivity study of Higgs trilinear self-coupling and quartic coupling is a part of scientific cases for CLIC. This work provides further motivation for high granular particle flow calorimetry for a future electron-positron linear collider.

## Declaration

This dissertation is the result of my own work, except where explicit reference is made to the work of others, and has not been submitted for another qualification to this or any other university. This dissertation does not exceed the word limit for the respective Degree Committee.

Boruo Xu



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*'You cannot open a book without learning something.'*

— Confucius, 551 BC – 479 BC



# Chapter 1

## Introduction

*‘The journey of a thousand miles begins with a single step.’*

— Lao Zi, 604 BC – 531 BC

Since twenty years ago, the high energy physics community has been considering a next-generation electron–positron collider after the Large Hadron Collider (LHC). Measurements from the LHC helped to establish the Standard Model of particle physics. Yet there are issues that the Standard Model could not explain. For example, the origin of the masses of neutrinos and the particles accounted for cosmic dark matter are questions that need to be addressed. Precision measurements from a next-generation electron–positron collider will hopefully provide answers to some of these questions.

The advantage of an  $e^+e^-$  linear collider over a hadron collider include: few background events from photon–photon collisions; a similar rate of pairs production of all particles as photon couples to all particles equally; smaller theoretical uncertainties on the electroweak interaction; and full reconstruction of events.

The International Linear Collider (ILC) [1], and the Compact Linear Collider (CLIC) [2], are two most promising candidates of the next-generation electron–positron collider. The ILC is capable of operating at centre-of-mass energies from 250 GeV to 500 GeV. CLIC can reach centre-of-mass energies from 350 GeV to 3 TeV. Both colliders will be able to measure Higgs couplings precisely via processes like  $e^+e^- \rightarrow ZH$  and measure top quark mass and couplings via processes such as  $e^+e^- \rightarrow t\bar{t}$ .

The optimisation of the design of the detectors for the future linear colliders is crucial to improve the ability to reconstruct events. By reconstructing the event to the

per-particle level, an event can be studied in detail. At the same time, physics simulation studies are important to demonstrate the physics reach of the future linear collider.

Chapter 2 starts with an overview of Standard Model of particle physics, including brief discussions on the quantum electrodynamics, quantum chromodynamics, and the electroweak interaction. The focus of the Standard Model discussion is on the Higgs mechanism and the Higgs boson in the Standard Model. The discussion then moves on to theories beyond the Standard Model, with an example of a general parametrisation of the Higgs theory. Last part of the chapter is dedicated to the discussion on studying the correlation between the polarisations of the tau pair from a boson decay to determine statistically if the parent boson is a scalar or a vector.

In chapter 3, the detector designs currently considered for two future electron-positron linear colliders, the ILC and CLIC, are described. After a short introduction of the two colliders, physics programme for these future colliders is discussed, followed by impact of physics requirements on the detector design. Afterwards, the International Large Detector (ILD), one detector option for the ILC, is discussed in detail, followed by overviews on each sub-detector in the ILD. The chapter finishes with a discussion on the modified ILD detector concept for CLIC.

In chapter 4, softwares for event simulation, event reconstruction, and event analysis are discussed. PandoraPFA, a world-leading pattern-recognition software for particle flow calorimetry, is presented. A discussion on jet algorithms is provided, followed by a discussion on multivariate analysis, where different fitting models, optimisation, and overfitting are discussed. The multivariate analysis is used extensively in chapter 5, chapter 6, and chapter 8.

Chapter 5 describes several new PandoraPFA algorithms for photon reconstruction. Sophisticated pattern recognition algorithms were implemented, which uses the topological properties of electromagnetic showers to identify photon candidates and separate them from nearby particles. It performs clustering of the energy deposits in the detector, followed by topological characterisation of the clusters, with the results being considered by a multivariate likelihood analysis. The algorithms lead to a significant improvement in the reconstruction of both single photons and multiple photons in high energy jets, which in return improves the jet energy resolution at high energies.

In chapter 6, a classification of tau lepton decay modes is presented. The analysis contains the event generation, simulation, reconstruction, and the use of the multivariate classifier for the classification. Utilising highly granular calorimeters, the resolutions of

energy and invariant mass of the tau decay products are improved. A hypothesis test was performed for expected decay final states. The performance of the tau decay mode classification is given, followed by an electromagnetic calorimeter optimisation study of the ILD detector based on the tau decay mode classification. The tau decay mode classification is further used in a proof-of-principle analysis in chapter 7 to demonstrate the ability to use the tau pair polarisation correlation as a signature for Higgs boson using the tau pair decay process, where both  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ .

In chapter 8, a full CLIC\_ILD detector simulation study is performed for the double Higgs production channel,  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$ , via  $W^+W^-$  fusion. Event generation and simulation will be discussed first. An overview of the analysis, including lepton finding and jet reconstruction, is presented, followed by an optimised multivariate analysis to distinguish signal from background processes. The optimised event selection is used to derive an estimate of the uncertainty on the cross section of double Higgs production at CLIC. The event selection is further used to provide an estimate of the uncertainty on the measurements of trilinear Higgs self-coupling and quartic coupling at CLIC.



# Chapter 2

## Theoretical overview

*'I believe it is impossible to be sure of anything.'*

— Han Fei Zi, 280 BC – 233 BC

This chapter provides a review of the Standard Model of Particle Physics, with an emphasis on the Higgs mechanism and the Higgs boson. A general parametrisation of the Higgs theory is discussed, which supplies the theoretical background for the physics analysis in chapter 8. Lastly a discussion of the usage of the tau pair polarisation correlations as a signature of Higgs boson is presented, which motivates the study in chapter 7.

### 2.1 Overview of the Standard Model

The Standard Model (SM) [3–6] is a quantum field theory concerning the fundamental particles and three fundamental interactions of nature: the electromagnetic; the weak; and the strong interactions. The fundamental particles in the SM consist of bosons and fermions. The bosons mediate the fundamental forces between particles: the photon is the force carrier of the electromagnetic force;  $W^+$ ,  $W^-$ , and  $Z$  bosons are the force carriers of the weak force; and the gluon,  $g$ , is the force carrier of the strong force. The properties of the force-exchange bosons and Higgs boson are listed in table 2.1.

The other fundamental particles are spin- $\frac{1}{2}$  fermions. For each fermion in the SM, there is an anti-fermion with the same mass and spin, but opposite charge. These fermions have three generations. Each generation of fermions has the same interaction

Force	Boson	Mass	Spin	Charge / $e$
Electromagnetic	photon	0	1	0
	$W^+$	80.385(15) GeV	1	1
	$W^-$	80.385(15) GeV	1	-1
Weak	Z	91.1876(21) GeV	1	0
	gluon	0	1	0
Strong	Higgs	125.1(3) GeV	0	0
-				

**Table 2.1:** Masses, spins, and charges of fundamental bosons in the SM. Values are taken from [3].

property, but different masses. Experimental evidences of three generations include the measurements of the Z boson decay-width, which strongly suggested three generations of neutrinos [7].

These fermions come in two distinct categories: leptons and quarks. Neutral leptons (the neutrinos) only experience the weak forces. Charged leptons ( $e^\pm, \mu^\pm, \tau^\pm$ ) experience the weak forces and the electromagnetic forces. Quarks experience all three fundamental forces described by the SM. Properties of these fermions are listed in table 2.2.

Many SM predictions have been experimentally verified. Some recent highlights include the discovery of the top quark in 1995 [8], the tau neutrino in 2000 [9], and the Higgs boson in 2012 [10, 11]. However, there are observations which are not explained by the SM. One issue is that the SM does not incorporate the gravitational force. Another issue is that the SM does not natively allow neutrino masses and mixings. The SM also does not explain the existence of the dark matter. There are many theories beyond the Standard Model (BSM) trying to provide an explanation for these issues. One example is the generalisation of the Higgs theory to allow non-SM coupling strengths [12, 13].

## 2.2 Quantum electrodynamics

QED is a quantum gauge field theory explaining electromagnetic interactions. Quantum field theory (QFT) is the theoretical framework for constructing quantum mechanical models of fundamental particles. Particles are treated as excited states of the underlying physical field in the QFT. A gauge theory is a type of field theory in which the Lagrangian is invariant under a continuous group of local transformations. Gauge invariance or gauge symmetry refers to when a field is transformed, but the Lagrangian is not.

Type	Generation	Fermion	Mass	Charge / $e$
Lepton	1	$e^-$	$0.548579909070(16) \text{ MeV}$	-1
		$\nu_e$	-	0
	2	$\mu^-$	$105.6583745(24) \text{ MeV}$	-1
		$\nu_\mu$	-	0
	3	$\tau^-$	$1776.86(12) \text{ MeV}$	-1
		$\nu_\tau$	-	0
Quark	1	u	$2.2^{+0.6}_{-0.4} \text{ MeV}$	$+\frac{2}{3}$
		d	$4.7^{+0.5}_{-0.4} \text{ MeV}$	$-\frac{1}{2}$
	2	c	$1270 \pm 30 \text{ MeV}$	$+\frac{2}{3}$
		s	$98^{+8}_{-4} \text{ MeV}$	$-\frac{1}{3}$
	3	t	$173210 \pm 510 \pm 710 \text{ MeV}$	$+\frac{2}{3}$
		b	$4180^{+40}_{-30} \text{ MeV}$	$-\frac{1}{3}$

**Table 2.2:** Masses and charges of the fundamental fermions in the SM. All fermions are spin- $\frac{1}{2}$  particles. For each fermion in the SM, there is an anti-fermion with the same mass and spin, but opposite charge. Neutrinos are known to have non-zero masses from the observation of neutrino flavour oscillations. The upper bound on the neutrino mass is 2 eV. For the top quark mass, the statistical uncertainty is listed first, followed by the systematic uncertainty. Values are taken from [3].

QED is an abelian gauge theory with the U(1) symmetry group. The gauge field, which mediates the interaction between the charged spin- $\frac{1}{2}$  fields, is the electromagnetic field, denoted  $A^\mu$ . The QED Lagrangian [14] for a spin- $\frac{1}{2}$  field interacting with the electromagnetic field is given by:

$$\mathcal{L}_{QED} = \bar{\psi} (i\gamma^\mu D_\mu - m) \psi - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}, \quad (2.1)$$

where  $\psi$  is the spin- $\frac{1}{2}$  Dirac field satisfying the Dirac equation, given by the Lagrangian density:

$$\mathcal{L} = \bar{\psi} (i\gamma^\mu D_\mu - m) \psi, \quad (2.2)$$

where the  $\gamma^\mu$  are the Dirac gamma matrices with  $\mu \in \{0, 1, 2, 3\}$ ;  $\bar{\psi}$  is defined as  $\psi^\dagger \gamma^0$ ;  $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$  is the electromagnetic field tensor;  $e$  is the coupling constant, which is equal to the electric charge;  $m$  is the mass of the electron; and the gauge covariant

derivative is given by:

$$D_\mu \equiv \partial_\mu + ieA_\mu, \quad (2.3)$$

where  $A_\mu$  is the covariant four-vector potential of the electromagnetic field generated by the electron itself.

## 2.3 Quantum chromodynamics

Quantum chromodynamics (QCD) is the quantum field theory of strong interactions. QCD theory is invariant under local non-Abelian SU(3) transformations. There are eight gauge bosons, the gluons, corresponding to the eight ( $8 = 3^2 - 1$ ) generators of the SU(3) symmetry group. Gluons carry colour charges. There are three types of colour charges, sometimes labelled as red, green, and blue. Anti-particles carry anticolour. Quarks are associated with a single colour. Gluons are made up of a colour and an anticolour (or superposition of colour–anticolour pair). The QCD Lagrangian is given by:

$$\mathcal{L}_{QCD} = \sum_{f \in u, d, s, c, b, t} \bar{\psi}_i \left( \left( i\gamma^\mu \partial_\mu - g_s \gamma^\mu G_\mu^a \frac{\lambda^a}{2} \right)_{ij} - m_f \delta_{ij} \right) \psi_j - \frac{1}{4} G_{\mu\nu}^a G^{a\mu\nu}, \quad (2.4)$$

where  $\psi$  represents a quark with a colour charge, indicated by  $i$  or  $j$ ;  $m$  controls the mass of the quark;  $g_s$  is the strong coupling constant;  $a$  is the colour charge;  $\lambda^a$  represents one of the eight Gell-Mann matrices; and  $G_{\mu\nu}^a$  represents the gauge invariant gluon field strength tensor, given by:

$$G_{\mu\nu}^a = \partial_\mu \gamma_\nu^a - \partial_\nu \gamma_\mu^a - g_s f_{abc} G_\mu^b G_\nu^c, \quad (2.5)$$

where  $G_\mu^b$  is the gluon field with colour charge  $b$ .

## 2.4 The electroweak interaction

The electroweak interaction can be thought as an extension to QED to incorporate the weak force, the force describing nuclear radioactive decay. The unification of the electromagnetic and the weak force is accomplished under an  $SU(2)_L \times U(1)$  gauge symmetry group. The corresponding gauge bosons are three  $W$  bosons ( $W^1$ ,  $W^2$ , and  $W^3$ ) from  $SU(2)_L$  gauge symmetry, and  $B$  boson from  $U(1)$  gauge symmetry. All gauge

bosons are initially massless. Fermion mass terms are forbidden under  $SU(2)_L$  gauge symmetry.

The electroweak Lagrangian can be written as

$$\mathcal{L}_{Electroweak} = \mathcal{L}_{Boson} + \mathcal{L}_{Fermion} + \mathcal{L}_{Higgs} + \mathcal{L}_{Yukawa}. \quad (2.6)$$

First consider  $\mathcal{L}_{Boson}$ , given by:

$$\mathcal{L}_{Boson} = -\frac{1}{4}W_{\mu\nu}^i W^{i\mu\nu} - \frac{1}{4}B_{\mu\nu} B^{\mu\nu}, \quad (2.7)$$

$$W_{\mu\nu}^i = \partial_\nu W_\mu^i - \partial_\mu W_\nu^i - g\varepsilon^{ijk}W_\mu^j W_\nu^k, \quad (2.8)$$

$$B_{\mu\nu} = \partial_\nu B_\mu - \partial_\mu B_\nu, \quad (2.9)$$

where  $B$  field is invariant under  $U(1)$  transformations;  $W$  field is invariant under non-Abelian  $SU(2)$  transformations; and the indices,  $i$ ,  $j$ , and  $k$ , indicate three  $W$  fields.

The term  $\mathcal{L}_{Fermion}$  describes the massless fermion fields coupling to the fermions, and the propagation of the fermion fields. The left-handed ( $\psi_L$ ) and the right-handed fermions ( $\psi_R$ ) are treated differently. The right-handed fermions are singlets. The left-handed fermions are in doublets with the corresponding fermions of the same generation. The term  $\mathcal{L}_{Fermion}$  is given by:

$$\mathcal{L}_{Fermion} = \sum_{\psi \in fermions} \bar{\psi}_L \gamma^\mu D_\mu^L \psi_L + \bar{\psi}_R \gamma^\mu D_\mu^R \psi_R, \quad (2.10)$$

where covariant derivatives  $D_\mu^L$  and  $D_\mu^R$  are defined as

$$D_\mu^L = \partial_\mu + ig\frac{\tau_i}{2}W_\mu^i + ig'Y_\psi B_\mu, \quad (2.11)$$

$$D_\mu^R = \partial_\mu + ig'Y_\psi B_\mu. \quad (2.12)$$

The structure of this Lagrangian allows  $W$  and  $B$  fields to couple with left-handed fermions, but only allows the  $B$  field to couple with right-handed fermions. The  $\tau_i$  matrices are the generators of  $SU(2)$  and  $Y_\psi$  is the hypercharge associated with the fermion field  $\psi$ . The  $W$  field couples with strength  $g$  to the fermion field. The  $B$  field couples with strength  $g'$  to the particles carrying weak hypercharge  $Y$ .

The term  $\mathcal{L}_{Higgs}$  describes the Higgs field. After electroweak symmetry breaking of the Higgs field, the mass terms of the gauge bosons are introduced. The term  $\mathcal{L}_{Yukawa}$  produces the mass terms of the quarks and charged leptons. Firstly a general spontaneous symmetry breaking mechanism is provided, followed by a description of the electroweak symmetry breaking.

### 2.4.1 Spontaneous symmetry breaking

Consider a complex scalar field, with the Klein-Gordon Lagrangian:

$$\mathcal{L} = \partial^\mu \psi^* \partial_\mu \psi - m^2 |\psi|^2 = \partial^\mu \psi^* \partial_\mu \psi - V(\psi), \quad (2.13)$$

where  $m$  is the mass term and  $V(\psi)$  is the potential of the field  $\psi$ . This Lagrangian has a global symmetry  $\psi \rightarrow e^{i\phi}\psi$ . The potential can be modified to add an interaction term without breaking the invariance of the global symmetry:

$$V(\psi) = m^2 |\psi|^2 + \lambda |\psi|^4, \quad (2.14)$$

where  $\lambda$  controls the interaction strength. This modified potential has a minimum at  $|\psi| = 0$  for  $m^2 > 0$ . However, if  $m^2 < 0$ , the minimum of the potential occurs when:

$$\frac{\partial V(\psi)}{\partial |\psi|} = 2m^2 |\psi| + 4\lambda |\psi|^3 = 0, \quad (2.15)$$

leading to a non-negative expectation value for the field:

$$|\psi| = \sqrt{\frac{-m^2}{2\lambda}} \equiv \frac{\nu}{\sqrt{2}}, \quad (2.16)$$

where  $\nu = \sqrt{-m^2/\lambda}$ . The solution that minimises the potential is not unique; it corresponds to a circle of points in the complex  $\psi$  plane. By choosing any one of these points, which are degenerate in energy, the symmetry of  $\psi \rightarrow e^{i\phi}\psi$  is broken. This phenomenon is known as the spontaneous symmetry breaking.

A consequence of spontaneous symmetry breaking is that the perturbation of the field along the degenerate energy direction, which is the circle in complex  $\psi$  plane, have no associated potential energy. This is formalised as Goldstone's theorem [15, 16]. The theorem states that spontaneous symmetry breaking always implies the existence of a massless particle.

To demonstrate Goldstone's theorem, using the Lagrangian in equation 2.13 as an example, after the spontaneous symmetry breaking of the field, the perturbation of the field  $\psi$  near the field minimum point can be written as

$$\psi = \frac{1}{\sqrt{2}}(\nu + \psi_1 + i\psi_2), \quad (2.17)$$

where  $\nu = \sqrt{-m^2/\lambda}$  refers to the minimum point in the potential, and  $\psi_1$  and  $\psi_2$  are real scalar fields. Substituting  $\psi$  in the Lagrangian in equation 2.13 gives:

$$\mathcal{L} = \frac{1}{2}\partial^\mu\psi_1\partial_\mu\psi_1 + \frac{1}{2}\partial^\mu\psi_2\partial_\mu\psi_2 - m^2\psi_1^2 + \dots \quad (2.18)$$

The mass term for the  $\psi_1$  field is  $\sqrt{-m^2}$  whereas there is no mass term for the  $\psi_2$  field, as stated by Goldstone's theorem.

In the previous example, the Lagrangian in equation 2.13 possesses the global symmetry of  $\psi \rightarrow e^{i\phi}\psi$ . Instead, if there is a local U(1) gauge symmetry of  $\psi \rightarrow e^{i\phi(x)}\psi$ , this implies a corresponding field  $A_\mu$ , which transforms as  $A_\mu \rightarrow A_\mu - \partial_\mu\phi(x)$ . For the gauge invariance, the covariant derivative becomes  $D_\mu = \partial_\mu + ieA_\mu$ . Hence the Lagrangian in equation 2.13 becomes:

$$\mathcal{L} = (D^\mu\psi)^*(D_\mu\psi) - m^2|\psi|^2 - \lambda|\psi|^4. \quad (2.19)$$

When the field is expanded around the minimum of the potential,  $\nu = \sqrt{-m^2/\lambda}$ , with  $m^2 < 0$ , the gauge boson mass term

$$+ \frac{e^2\nu^2}{2}A^\mu A_\mu, \quad (2.20)$$

is obtained from the  $(D^\mu\psi)^*(D_\mu\psi)$  term in the Lagrangian. Therefore the spontaneous symmetry breaking of a gauge field gives rise to a gauge boson mass.

## 2.5 Higgs mechanism

The Higgs mechanism is an extension of the example of the spontaneous symmetry breaking introduced in the previous section. It can provide mass terms for bosons and fermions that are compatible with the gauge invariance of the SM. Consider a complex scalar Higgs field,  $\Phi_H$ , transforms as a doublet of SU(2) with hypercharge  $Y = \frac{1}{2}$ . The

Higgs Lagrangian is given by:

$$\mathcal{L}_{Higgs} = \left( D_\mu \Phi_H \right)^\dagger \left( D^\mu \Phi_H \right) - \mu^2 \Phi_H^\dagger \Phi_H + \lambda \left( \Phi_H^\dagger \Phi_H \right)^2, \quad (2.21)$$

where  $\lambda$  and  $\mu$  are constants. The  $SU(2)_L \times U(1)$  symmetry of the electroweak Lagrangian demands that the covariant derivative of the Higgs field takes the form

$$D_\mu = \left( \partial_\mu + ig \frac{\tau_i}{2} W_\mu^i + ig' \frac{1}{2} B_\mu \right), \quad (2.22)$$

where  $g$  is the coupling constant of the  $SU(2)_L$  gauge symmetry;  $g'$  is the coupling constant of the  $U(1)$  gauge symmetry; and the  $\tau_i$  are Pauli matrices. The Higgs potential is given by:

$$V(H) = \mu^2 \Phi_H^\dagger \Phi_H - \lambda \left( \Phi_H^\dagger \Phi_H \right)^2. \quad (2.23)$$

The Higgs potential is minimised when

$$\sqrt{\Phi_H^\dagger \Phi_H} = \frac{\nu}{\sqrt{2}} = \sqrt{\frac{\mu^2}{2\lambda}}. \quad (2.24)$$

By expanding the Higgs field about the minimum point of the potential, the non-zero vacuum expectation value (VEV) can be written as:

$$\langle \Phi_H \rangle = \begin{pmatrix} 0 \\ \frac{\nu}{\sqrt{2}} \end{pmatrix}, \quad (2.25)$$

with a real  $\nu$ . Substituting the Higgs VEV into the  $\mathcal{L}_{Higgs}$  in equation 2.21, the  $\left( D_\mu \Phi_H \right)^\dagger \left( D^\mu \Phi_H \right)$  term becomes

$$-\frac{1}{8} \begin{pmatrix} 0 & \nu \end{pmatrix} \begin{pmatrix} gW_\mu^3 + g'B_\mu & g(W_\mu^1 - iW_\mu^2) \\ g(W_\mu^1 + iW_\mu^2) & -gW_\mu^3 + g'B_\mu \end{pmatrix} \begin{pmatrix} gW_\mu^3 + g'B_\mu & g(W_\mu^1 - iW_\mu^2) \\ g(W_\mu^1 + iW_\mu^2) & -gW_\mu^3 + g'B_\mu \end{pmatrix} \begin{pmatrix} 0 \\ \nu \end{pmatrix}. \quad (2.26)$$

Ignoring the negative sign, equation 2.26 simplifies to

$$\frac{\nu^2 g^2}{8} (W_\mu^1 - iW_\mu^2)(W_\mu^1 + iW_\mu^2) + \frac{\nu^2}{8} (gW_\mu^3 - g'B_\mu)^2. \quad (2.27)$$

The physical fields  $W_\mu^+$  and  $W_\mu^-$  can be identified with the first part of equation 2.27, as

$$W_\mu^+ = \frac{1}{\sqrt{2}}(W_\mu^1 - iW^2\mu), \quad (2.28)$$

$$W_\mu^- = \frac{1}{\sqrt{2}}(W_\mu^1 + iW^2\mu). \quad (2.29)$$

The physical fields  $Z_\mu$  and  $A_\mu$  are associated with  $W_\mu^3$  and  $B_\mu$ . Since the Z boson is massive and the photon is massless, the second part of equation 2.27 should give rise to Z boson mass term only, with no mass term for photon. This can be achieved by rearranging the second part of the equation 2.27:

$$\frac{\nu^2}{8}(gW_\mu^3 - g'B_\mu)^2 = \frac{\nu^2(g^2 + g'^2)}{8} \left( \frac{g}{\sqrt{g^2 + g'^2}}W_\mu^3 - \frac{g'}{\sqrt{g^2 + g'^2}}B_\mu \right)^2. \quad (2.30)$$

A convenient way to connect  $g$  and  $g'$  is to use the Weinberg mixing angle [17], denoted as  $\theta_W$ . The Weinberg mixing angle is defined as

$$\cos(\theta_W) = \frac{g}{\sqrt{g^2 + g'^2}}, \quad (2.31)$$

$$\sin(\theta_W) = \frac{g'}{\sqrt{g^2 + g'^2}}. \quad (2.32)$$

Equation 2.30 can be rewritten using the Weinberg mixing angle:

$$\frac{\nu^2(g^2 + g'^2)}{8}(\cos(\theta_W)W_\mu^3 - \sin(\theta_W)B_\mu)^2. \quad (2.33)$$

The physical field  $Z_\mu$  can be immediately identified as:

$$Z_\mu = \cos(\theta_W)W_\mu^3 - \sin(\theta_W)B_\mu. \quad (2.34)$$

Consequently, the physical field  $A_\mu$  with associated massless photon can be written as:

$$A_\mu = \sin(\theta_W)W_\mu^3 + \cos(\theta_W)B_\mu. \quad (2.35)$$

Equation 2.27 can be written in terms of the physical fields  $W_\mu^+$ ,  $W_\mu^-$ ,  $Z_\mu$ , and  $A_\mu$ :

$$\frac{(g\nu)^2}{4}W_\mu^+W^{-\mu} + \frac{(g^2 + g'^2)\nu^2}{8}Z_\mu Z^\mu. \quad (2.36)$$

The first term gives mass of the  $W^+$  and  $W^-$  vector bosons. The second term gives mass of the  $Z$  vector boson. There is no mass term for the photon. The spontaneous symmetry breaking of the Higgs field breaks the electroweak  $SU(2)_L \times U(1)$  gauge symmetry to the  $U(1)$  gauge symmetry of the electromagnetism. The masses of the  $W^+$ ,  $W^-$  and  $Z$  bosons are given by:

$$m_{W^\pm} = m_{W^-} = \frac{g\nu}{2}, \quad m_Z = \frac{\nu\sqrt{g^2 + g'^2}}{2} = \frac{m_W}{\cos(\theta_W)}. \quad (2.37)$$

## 2.6 Higgs boson

For the Higgs doublet complex field in the SM, there are four real scalar degrees of freedom. Three degrees of freedom are “eaten” to form the longitudinal polarisations of the  $W_\mu^\pm$  and  $Z_\mu$  fields. The remaining one real scalar degree of freedom forms the Higgs boson. The properties of the Higgs bosons can be shown in the unitary gauge, where three degrees of freedom are manifestly eaten. The Higgs field is given by:

$$H(x) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ \nu + h(x) \end{pmatrix}, \quad (2.38)$$

where  $h(x)$  is the real scalar field of the Higgs boson and  $\nu$  the Higgs vacuum expectation value. The Higgs boson is not charged under electromagnetism as the field is real. The coupling of the Higgs boson to other fields can be calculated out by replacing  $\nu$  with  $\nu + h(x)$  in equation 2.36:

$$m_W^2 \left( \frac{2h}{\nu} + \frac{h^2}{\nu^2} \right) W_\mu^+ W^{-\mu} + \frac{m_Z^2}{2} \left( \frac{2h}{\nu} + \frac{h^2}{\nu^2} \right) Z_\mu Z^\mu. \quad (2.39)$$

The Higgs boson self-interaction terms are obtained by replacing  $\nu$  with  $\nu + h(x)$  in the Higgs field potential in equation 2.23:

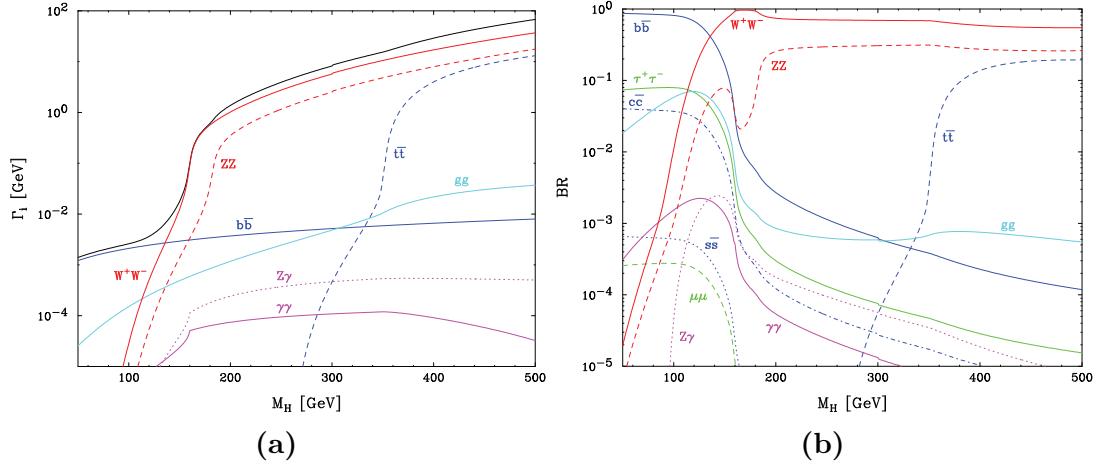
$$\frac{\mu^2}{2} (\nu + h)^2 - \frac{\lambda}{4} (\nu + h)^4 \supset -\lambda\nu^2 h^2 - \lambda\nu h^3 - \frac{\lambda}{4} h^4 \quad (2.40)$$

The quadratic term,  $-\lambda\nu^2 h^2$ , is the Higgs boson mass term,  $m_H = \sqrt{2\lambda}\nu$ . The terms in  $h^3$  and  $h^4$  give trilinear and quadlinear Higgs self-interaction terms.

Once the Higgs boson mass is known,  $\lambda$  can be determined and the Higgs boson decay widths and branching fractions can be calculated. Figure 2.1 shows the Higgs boson

partial decay widths and the branching ratios as a function of the Higgs boson mass for different Higgs decay modes.

The Higgs boson mass is measured to be  $125.09 \pm 0.24$  GeV [3]. Because the Higgs boson is lighter than a pair of heavier particles such as  $W^+W^-$  or  $ZZ$ , the processes  $H \rightarrow W^+W^-$  and  $H \rightarrow ZZ$  are forbidden kinematically. However, the quantum field theory allows such processes to happen, if one of the decay products is virtual and not on the mass shell. The virtual gauge boson subsequently decays to real on-mass-shell particles.



**Figure 2.1:** a) The Higgs boson partial decay widths, and b) Higgs boson branching ratios, plotted as a function of the Higgs boson mass,  $m_H$ . In a), the black curve shows the total decay width. Both figures are taken from [18].

## 2.7 Yukawa couplings

The Yukawa sector of the electroweak Lagrangian provides mass terms for quarks and charged leptons after the spontaneous symmetry breaking of the Higgs field. The corresponding term in the Lagrangian is:

$$\mathcal{L}_{Yukawa} = -\lambda^u \overline{q}_L \Phi_H^c u_R - \lambda^d \overline{q}_L \Phi_H d_R - \lambda^e \overline{l}_L \Phi_H e_R + h.c., \quad (2.41)$$

where  $q_L$  is the left-handed quark doublet field;  $u_R$  is the up-type right-handed quark singlet field;  $d_R$  is the down-type right-handed quark singlet field;  $l_L$  is the left-handed lepton doublet field;  $e_R$  is the right-handed charged lepton singlet field;  $\lambda$  is a constant;  $\Phi_H^c \equiv i\sigma^2 H^*$  is a SU(2) doublet field with hypercharge  $Y = -\frac{1}{2}$ ;  $h.c.$  indicates the Hermitian conjugate terms; and the Lagrangian is summed over all possible quarks and leptons. When the Higgs vacuum expectation value is substituted into  $\mathcal{L}_{Yukawa}$ , the

Yukawa interaction terms give the fermion mass terms:

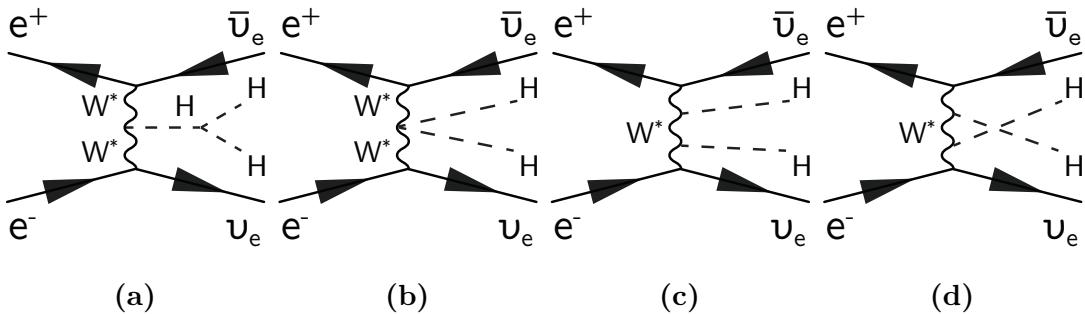
$$m_u = \frac{\lambda^u \nu}{\sqrt{2}}, \quad m_d = \frac{\lambda^d \nu}{\sqrt{2}}, \quad m_e = \frac{\lambda^e \nu}{\sqrt{2}}. \quad (2.42)$$

## 2.8 Beyond the Standard Model Higgs models

A number of BSM Higgs theories have been proposed. For example, the light Higgs could be a composite bound state of a new strongly-interacting sector at the TeV scale. If the composite Higgs is pseudo Nambu-Goldstone boson from spontaneous global symmetry breaking, the Higgs can be naturally light [12]. In this model, the couplings of the Higgs would deviate from those in the SM for Higgs interactions at the TeV scale.

An important physics process for testing the Higgs theory is the double Higgs production via vector boson fusion at the TeV scale [19–21]. For the composite Higgs scenario, the scattering amplitude for this process increases with energy. It is difficult to measure the double Higgs production at the LHC due to the large SM background rate [20]. However, a multi-TeV electron–position linear collider, such as the Compact Linear Collider, would be able to measure the cross section for this process [22].

The study of double Higgs production via  $W^+W^-$  fusion can probe the Higgs trilinear self coupling,  $g_{HHH}$ , and quartic coupling,  $g_{WWHH}$ . The coupling  $g_{HHH}$  is associated with the terms in  $h^3$  in Higgs potential in equation 2.40. The coupling  $g_{WWHH}$  is associated with the terms in  $h^2$  in Higgs interaction with other fields in equation 2.39. Leading-order Feynman diagrams for double Higgs production via  $W^+W^-$  fusion are shown in figure 2.2. The diagram shown in figure 2.2a contains the triple Higgs vertex, which is sensitive to the Higgs trilinear self coupling  $g_{HHH}$ . The diagram in figure 2.2b is sensitive to the quartic coupling  $g_{WWHH}$ . Figures 2.2c and 2.2d show Feynman diagrams for irreducible background processes containing two  $HW^+W^-$  vertices.



**Figure 2.2:** The main Feynman diagrams for the leading-order  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  processes.

Following the assumption made in [20, 21], the self-interaction of the light scalar Higgs,  $h$ , and its coupling to other SM bosons can be described by a Lagrangian using the notation in [21]. In this description, after the electroweak symmetry breaking, the Lagrangian is given by:

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - V(h) + \left(m_W^2 W_\mu^+ W^{-\mu} + \frac{m_Z^2}{2} Z_\mu Z^\mu\right) \left[1 + 2a\frac{h}{\nu} + b\frac{h^2}{\nu^2} + \dots\right], \quad (2.43)$$

where  $V(h)$  is the  $h$  field potential

$$V(h) = \frac{1}{2}m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2\nu}\right) h^3 + d_4 \left(\frac{m_h^2}{8\nu^2}\right) h^4 + \dots, \quad (2.44)$$

and  $a$ ,  $b$ ,  $d_3$  and  $d_4$  are dimensionless parameters. Higher-order terms in  $h$  are omitted. The parameters  $a$  and  $b$  are proportional to the coupling strengths of the  $VVh$  and  $VVhh$  vertices, where  $V$  represents a vector boson, and the parameters  $d_3$  and  $d_4$  are proportional to the trilinear and quadlinear  $h$  self-coupling strengths respectively. Comparing with the  $\mathcal{L}_{Higgs}$  in the SM (see equation 2.39 and equation 2.40), it can be seen that  $a = b = d_3 = d_4 = 1$  in the SM, and all higher order terms vanish. However, BSM Higgs models allow  $a, b, d_3, d_4$  to take different values.

Consider a pair of the longitudinal polarised vector bosons ( $V_L$ ) coupling to two  $h$  fields. The scattering amplitude for  $V_L V_L \rightarrow hh$  can be written as:

$$A = a^2(A_{SM} + A_1\delta_b + A_2\delta_{d_3}), \quad (2.45)$$

where  $A_{SM}$  is the SM amplitude and:

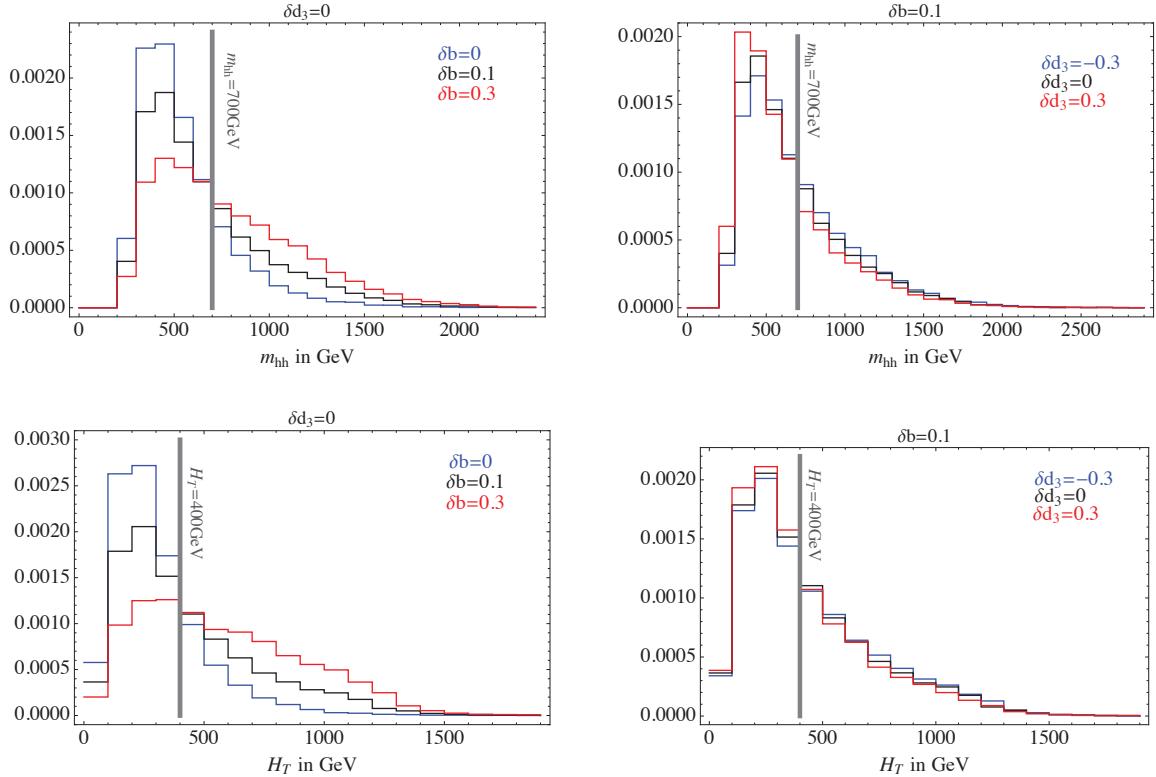
$$\delta_b \equiv 1 - \frac{b}{a^2}, \quad (2.46)$$

$$\delta_{d_3} \equiv 1 - \frac{d_3}{a}. \quad (2.47)$$

The term  $A_1$  grows like the square of energy at a large center-of-mass energy,  $E \gg m_V$ . The terms  $A_{SM}$  and  $A_2$  have no energy dependence. Therefore, the parameter  $\delta_b$  controls the magnitude of the increasing of the scattering amplitude as a function of energy. In an electron–positron collider, this scattering process can be studied via the double Higgs production  $e^+e^- \rightarrow \nu\bar{\nu}hh$  channel, where the cross section can be written as

$$\sigma = a^4 \sigma_{SM} \left(1 + A\delta_b + B\delta_{d_3} + C\delta_b\delta_{d_3} + D\delta_b^2 + E\delta_{d_3}^2\right), \quad (2.48)$$

where  $\sigma_{SM}$  is the SM cross section. Variables that increase with the increasing of the centre-of-mass energies are suitable for studying the cross section dependence on parameters  $\delta_b$  and  $\delta_{d_3}$ . Two examples of such variables are the invariant mass of the two Higgs system,  $m_{hh}$ , and the scalar sum of two Higgs transverse momenta,  $H_T$ . Figure 2.3 shows that the  $m_{hh}$  and  $H_T$  distributions are sensitive to the values of  $\delta_b$  and  $\delta_{d_3}$  [21]. The changes in the  $m_{hh}$  and  $H_T$  distributions can be related to the change in  $\delta_b$  and  $\delta_{d_3}$ . Therefore deviations of  $\delta_b$  and  $\delta_{d_3}$  from those SM values, 1, could be established using the  $m_{hh}$  and  $H_T$  distributions. It should be noted that figure 2.3 shows a generator-level study; detector effect will affect the distributions because of, for example, the loss of the reconstruction efficiency in the barrel/endcap overlap region.



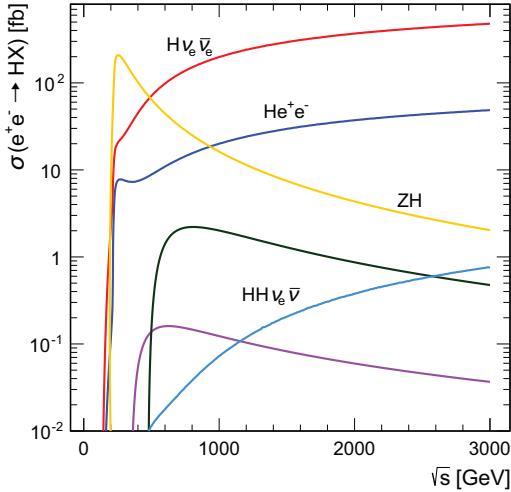
**Figure 2.3:** Normalised differential cross sections  $d\sigma/dm_{hh}$  and  $d\sigma/dH_T$  for  $e^+e^- \rightarrow \nu\bar{\nu}hh$  process for CLIC at  $\sqrt{s} = 3$  TeV after applying generator-level identification cuts, for several values of  $\delta_b$  and  $\delta_{d_3}$ . Figures are taken from [21].

In the expression of the cross section of the double Higgs production via  $e^+e^- \rightarrow \nu\bar{\nu}hh$  in equation 2.48, the parameter  $a$ , which is proportional to  $g_{VVH}$ , enters as an overall factor. Figure 2.4 shows the comparison of cross sections as a function of the centre-of-mass energy, for different the Higgs production modes. Up to a centre-of-mass energy of  $\sqrt{s} = 3$  TeV, the cross sections of the single Higgs production are two orders of magnitude larger than the cross sections of the double higgs production. The cross section of

$e^+e^- \rightarrow \nu\bar{\nu}h$  channel is given by:

$$\sigma = \sigma_{SM} (1 + A\Delta a + B\Delta a^2), \quad (2.49)$$

where  $\Delta a \equiv 1 - a$  is the change in  $a$ , and  $A$  and  $B$  are two dimensionless coefficients. The measurement of the parameter  $a$ , using  $e^+e^- \rightarrow \nu\bar{\nu}h$  channel, would be performed before the measurement of the  $\delta_b$  and  $\delta_{d_3}$  for the double Higgs production.



**Figure 2.4:** Cross sections as a function of centre-of-mass energy for Higgs production processes at an electron-positron collider for a Higgs mass of 126 GeV. The cross section values correspond to unpolarised beams and do not include the effect of beamstrahlung. The plot is taken from [23].

Therefore, for the purpose of measuring  $g_{VVHH}$  and  $g_{HHH}$  via double Higgs production, it is sufficient to treat the parameter  $a$  as a known constant. Hence only a two-dimensional fit of the parameters  $\delta_b$  and  $\delta_{d_3}$  would be performed to extract values of  $\delta_b$  and  $\delta_{d_3}$ .

## 2.9 Tau pair polarisation correlations as a signature of Higgs boson

The tau lepton has been studied extensively in the past at the Large Electron–Positron Collider (LEP) [24]. The tau lepton is a fundamental particle, with a negative electric charge and a spin of  $\frac{1}{2}$ . It has the same fundamental interaction property as an electron, but a much larger mass. Unlike the stable electron, because the tau lepton is massive, it decays via the weak interaction with a mean decay lifetime of  $(290.3 \pm 0.5) \times 10^{-15}$  s [25]. The tau lepton has many decay modes. The decay modes with branching ratio above 2% are listed in table 2.3.

Decay modes	Final states	Branching ratio
$e^- \bar{\nu}_e \nu_\tau$	$e^- \bar{\nu}_e \nu_\tau$	$17.83 \pm 0.04\%$
$\mu^- \bar{\nu}_\mu \nu_\tau$	$\mu^- \bar{\nu}_\mu \nu_\tau$	$17.41 \pm 0.04\%$
$\pi^- \nu_\tau$	$\pi^- \nu_\tau$	$10.83 \pm 0.06\%$
$\rho \nu_\tau$	$\pi^- \pi^0 \nu_\tau$	$25.52 \pm 0.09\%$
$a_1 \nu_\tau$ neutral	$\pi^- \pi^0 \pi^0 \nu_\tau$	$9.30 \pm 0.11\%$
$a_1 \nu_\tau$ charged	$\pi^+ \pi^- \pi^- \nu_\tau$	$8.99 \pm 0.06\%$
$\pi^+ \pi^- \pi^- \pi^0 \nu_\tau$	$\pi^+ \pi^- \pi^- \pi^0 \nu_\tau$	$2.70 \pm 0.08\%$

**Table 2.3:** Decay modes, final state particles, and branching ratios of the seven major  $\tau^-$  decays, taken from [3].

A scalar Higgs boson with spin-0 can decay to  $\tau_L^+ \tau_L^-$  or  $\tau_R^+ \tau_R^-$ , whereas a vector boson Z with spin-1 can decay to  $\tau_L^+ \tau_R^-$  or  $\tau_R^+ \tau_L^-$ , where L, R denotes the tau lepton helicity. Therefore, by studying the correlation between the polarisations of the tau pair from a boson decay, one can determine statistically if the parent boson is a scalar or a vector.

The tau pair polarisation correlation can be studied using various tau decay modes. Here reference [26] is followed and the  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$  decay mode is used as the example. The Higgs and Z boson decay to a tau pair, where both tau leptons subsequently decay via  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ , can be represented as:

$$X \rightarrow \tau_\alpha^+ \tau_\beta^- \rightarrow \pi^+ \pi^- + \nu s, \quad (2.50)$$

where  $X$  is either H or Z, and  $\alpha, \beta$  are the tau lepton helicities, L or R. In the collinear limit where  $m_\tau^2/m_X^2 \ll 1$ , the appropriate kinematic variables are the energy fractions:

$$z = \frac{E_{\pi^-}}{E_{\tau^-}}, \quad (2.51)$$

$$\bar{z} = \frac{E_{\pi^+}}{E_{\tau^+}}. \quad (2.52)$$

For a single tau decay, the differential cross section distribution can be written as:

$$\frac{1}{\Gamma_\tau} \frac{d\Gamma}{dz} = Br(\tau^- \rightarrow \pi^- \nu_\tau) f(\tau_\alpha^- \rightarrow \pi^-; z), \quad (2.53)$$

where  $Br(\tau^- \rightarrow \pi^- \nu_\tau)$  is the branching fraction of  $\tau^- \rightarrow \pi^- \nu_\tau$  decay mode. The form  $f$  can be obtained by working out the matrix element from the Feynman diagram and

integrating the square of the matrix element over the phase space [27]:

$$f(\tau_\alpha^- \rightarrow \pi^-; z) = 1 + P_\alpha(2z - 1), \quad (2.54)$$

where  $P_L = -1$  and  $P_R = +1$ . Hence for the tau pair decay, the differential cross section distribution is of the form:

$$\frac{d^2 N(X \rightarrow \tau^+ \tau^- \rightarrow \pi^+ \pi^- + \nu' s)}{dz d\bar{z}} = \left( Br(\tau^- \rightarrow \pi^- \nu_\tau) \right)^2 \sum_{\alpha, \beta} C_{\alpha\beta}^X f(\tau_\alpha^- \rightarrow \pi^-; z) f(\tau_\beta^+ \rightarrow \pi^+; \bar{z}), \quad (2.55)$$

where the only non-zero correlation coefficients  $C_{\alpha\beta}$  for the parity-conserving  $H \rightarrow \tau^+ \tau^-$  are:

$$C_{LL}^H = C_{RR}^H = \frac{1}{2}. \quad (2.56)$$

In contrast, the non-zero correlation coefficients for the  $Z \rightarrow \tau^+ \tau^-$  are:

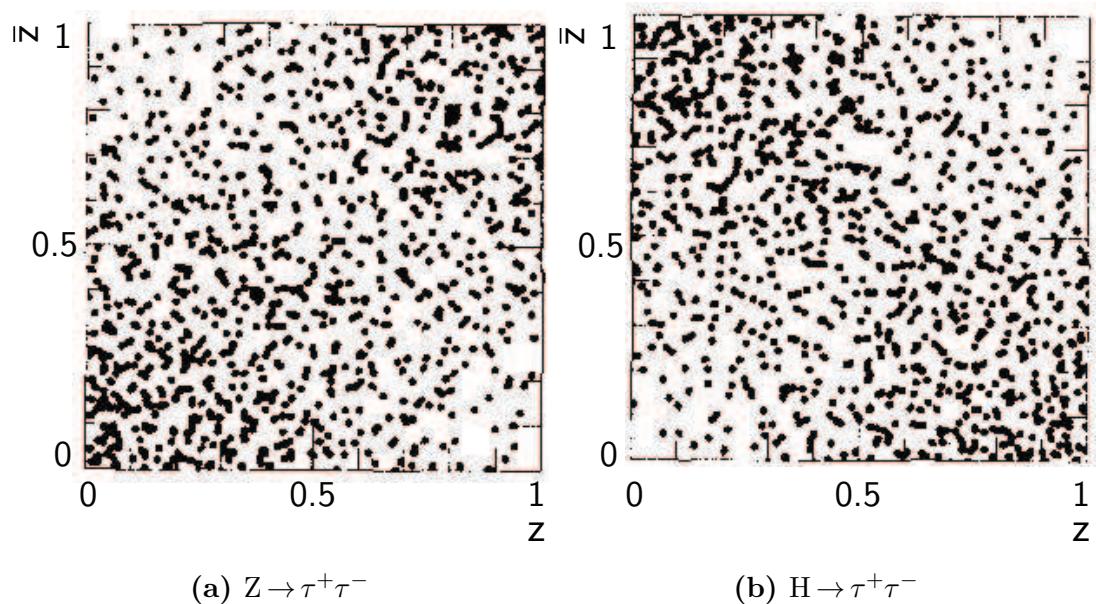
$$C_{LR}^Z = \frac{1}{2}(1 - P_\tau), \quad C_{RL}^Z = \frac{1}{2}(1 + P_\tau), \quad (2.57)$$

where  $P_\tau$  is the mean non-zero tau polarisation of  $Z$  decays. The tau polarisation is not zero because the process  $Z \rightarrow \tau^+ \tau^-$  is not parity-conserving. In the SM:

$$P_\tau = \frac{-2va}{v^2 + a^2}, \quad (2.58)$$

where the parameter  $v = -\frac{1}{2} + \sin^2 \theta_W$  and  $a = -\frac{1}{2}$  are the respective vector and axial-vector  $Z\tau^+ \tau^-$  couplings.

Figure 2.5 shows the resulting two-dimensional distributions of  $\bar{z} = \frac{E_{\pi^+}}{E_{\tau^+}}$  versus  $z = \frac{E_{\pi^-}}{E_{\tau^-}}$  for  $Z \rightarrow \tau^+ \tau^-$  and  $H \rightarrow \tau^+ \tau^-$  channels, where both tau leptons decay via  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ . The difference of the tau pair polarisation correlation between  $Z$  and  $H$  is clear. The energy distribution of the charged pion from  $Z \rightarrow \tau^+ \tau^-$  has the form of  $\bar{z} \sim z$ , whilst the distribution from  $H \rightarrow \tau^+ \tau^-$  has the form of  $\bar{z} \sim (1 - z)$ . Therefore, in  $Z \rightarrow \tau^+ \tau^-$  process, a high-energy  $\pi^\pm$  is likely to be associated with a high-energy  $\pi^\mp$ . In  $H \rightarrow \tau^+ \tau^-$  process, the opposite is favoured. If the tau pair decay from Higgs boson is observed, the decay can be recognised in the  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$  mode as a high-energy  $\pi^\pm$  with a low-energy  $\pi^\mp$ . Hence, the tau decay product energy distribution can be used as a signature for  $H \rightarrow \tau^+ \tau^-$ .



**Figure 2.5:** Two-dimensional distributions of  $\bar{z} = E_{\pi^+}/E_{\tau^+}$  plotted against  $z = E_{\pi^-}/E_{\tau^-}$  for a)  $Z \rightarrow \tau^+ \tau^-$ , and b)  $H \rightarrow \tau^+ \tau^-$  processes, where both tau leptons decay via  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ , adapted from reference [27].

# Chapter 3

## Detectors for Future Electron–Positron Linear Colliders

*‘The person attempting to travel two roads at once will get nowhere.’*

— Xun Kuang, 313 BC – 238 BC

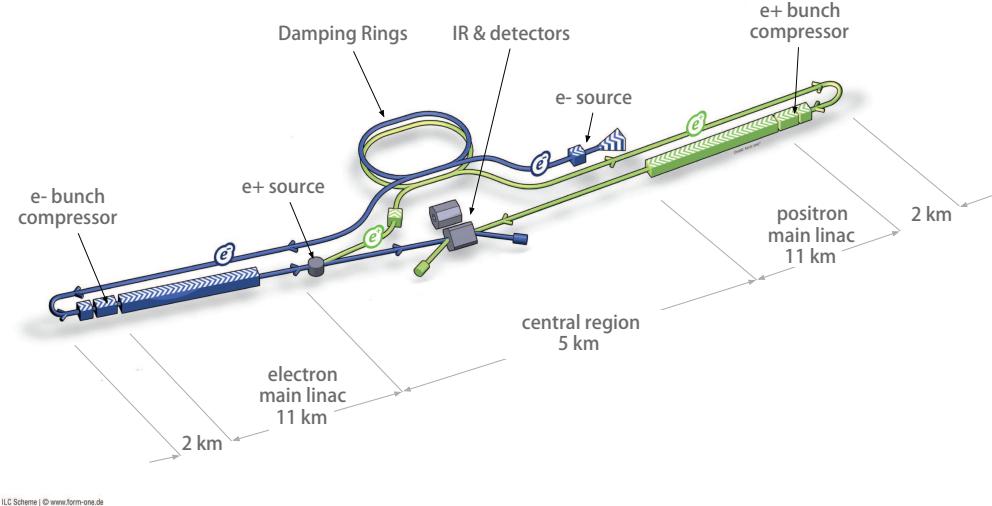
Two leading candidates for next-generation electron-positron linear particle colliders are the International Linear Collider (ILC) [1], and the Compact Linear Collider (CLIC) [2]. This chapter provides an overview of the two colliders, followed by the physics programme at these colliders, the detectors requirements, and the description of detectors for the ILC and CLIC.

### 3.1 International Linear Collider

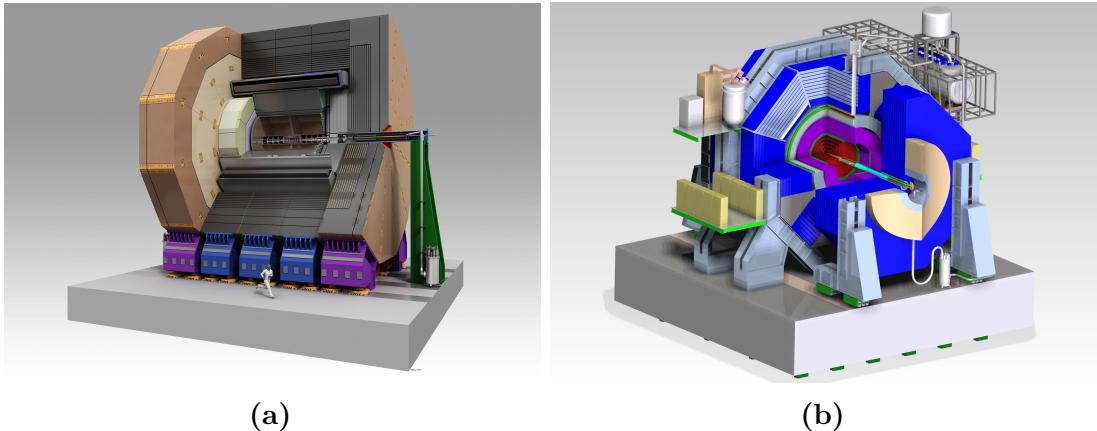
The ILC is a high-luminosity future electron-positron linear particle collider. The machine will be built in two stages. The first stage would have a centre-of-mass energy of 250/350 GeV. The second stage would have a centre-of-mass energy of 500 GeV with a possible upgrade to 1 TeV. The layout of the collider complex is shown in figure 3.1. Two detector concepts have been developed for the ILC: the International Large Detector (ILD) [28] and the Silicon Detector (SiD) [29]. Both detectors are shown in figure 3.2.

### 3.2 Compact Linear Collider

CLIC is a potential next-generation electron-positron linear particle collider at CERN [2]. CLIC is designed to be built in three stages: a first stage of a centre-of-mass energy of



**Figure 3.1:** Schematic layout of the International Linear Collider, indicating all the major subsystems (not to scale), taken from [30].

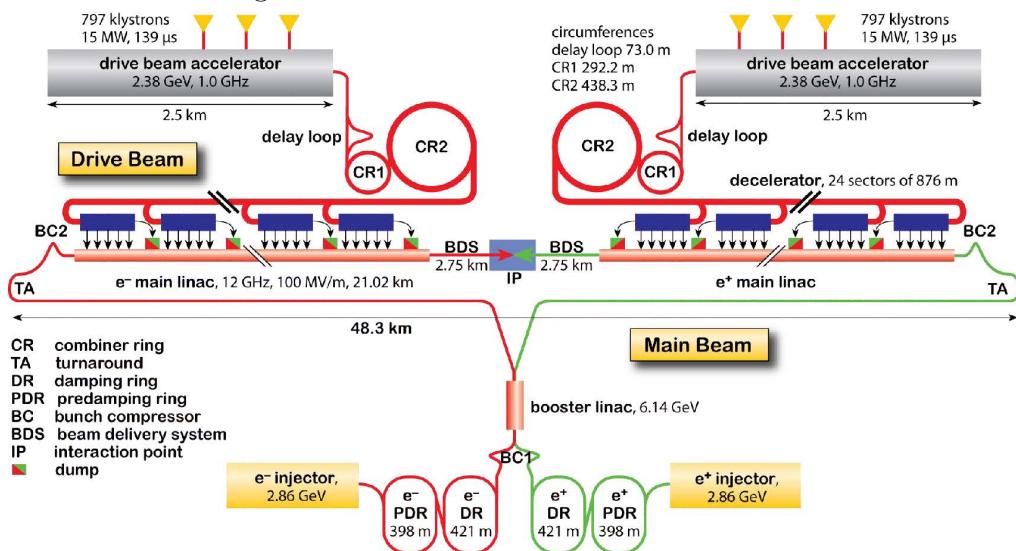


**Figure 3.2:** a) The International Large Detector, and b) the Silicon Detector. Both detector concepts are developed for the International Linear Collider. Both figures are taken from [30].

380 GeV; a second stage of a centre-of-mass energy of 1.4 TeV; and the final stage of a centre-of-mass energy of 3 TeV. The layout of the CLIC complex at the final stage is shown in figure 3.3.

The ILC and CLIC share some common features. Both colliders will be linear colliders as opposed to circular colliders, like the Large Hadron Collider (LHC). Detectors for both colliders will use the high granularity particle flow calorimetry [2, 30]. One major difference between the two colliders is the operating energy. Due to a higher centre-of-mass energy at CLIC, there are significant beam related backgrounds. The  $e^+e^-$  incoherent pair background has a major influence on the design of the inner region and

the forward region of the detectors [2]. The pile-up of  $3.2\ \gamma\gamma \rightarrow \text{hadrons}$  background events per bunch, also integrating over 60 bunch crossings, need to be mitigated for physics analyses [2]. Another difference between the ILC and CLIC is that the timing separation between bunches is much shorter at CLIC. The CLIC beam contains 312 bunch trains with a train repetition rate of 50 Hz, separated by 0.5 ns between each bunch train. This short timing separation suggests that the detector will integrate over a number of bunch crossings.



**Figure 3.3:** The layout of the Compact Linear Collider at the final stage of a centre-of-mass of energy of 3 TeV, taken from [31].

### 3.3 Physics at future linear colliders

An  $e^+e^-$  linear collider has advantages over the current hadron collider, the LHC. Those advantages include:

- i) Events in the  $e^+e^-$  collider will be cleaner than those in the hadron collider. In the LHC, many proton–proton collisions per bunch crossing are expected [32], generating hundreds of particles from parton collisions. In the  $e^+e^-$  collider, the main source of background comes from photon–photon collisions [1, 2]. Depending on the operating energy and scheme, there will be only a few of these photon–photon collisions per bunch crossing. Particles produced from these collisions are mainly in the forward direction, which can be identified relatively easily.
  - ii) Electroweak interactions in the  $e^+e^-$  collider will be democratic as the photon couples to all particles in and beyond the Standard Model equally [3–6]. The production of pairs of all particles will be at a similar rate. In the LHC, the non-perturbative strong

interaction is the main channel for the particle production. Heavy particles have lower production rates than light ones, as the parton distributions fall sharply for a composite object like a proton [33].

- iii) In the LHC, calculation of the cross section depends on quantum chromodynamics and the proton structure function, which have large systematic errors. In an  $e^+e^-$  collider, the initial particles,  $e^+$  and  $e^-$ , are point-like particles, interacting through electroweak forces only. Consequently, theoretical uncertainties are smaller.
- iv) The physics at the  $e^+e^-$  collider can be studied in detail. Without complicated underlying events, complete events can be reconstructed. Polarised beams of electrons and positrons with known initial and final polarisation states could also be used to enhance the production of certain interactions, for example, electron–positron annihilation with opposite helicities.

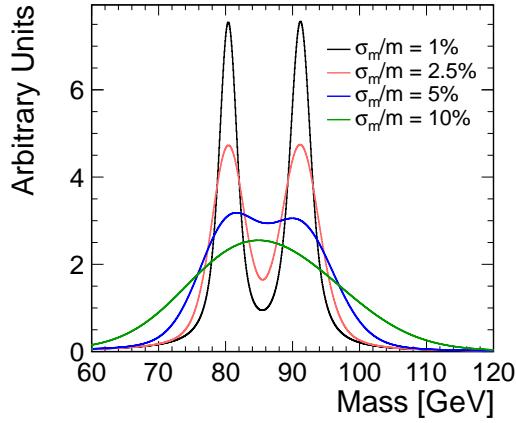
The physics programmes for the ILC and CLIC, which are the driving forces behind the detector design, share some common features. At a centre-of-mass energy of 250 GeV, the collider could operate as a Higgs factory, allowing precise measurements of Higgs couplings via channels like  $e^+e^- \rightarrow ZH$ . At a centre-of-mass energy of 350 GeV, the collider can continue to measure Higgs couplings, as well as to measure top quark mass and couplings via channels such as  $e^+e^- \rightarrow t\bar{t}$ . At a centre-of-mass energy of 1 TeV and beyond, the collider would be able to produce rare Higgs decays, allowing measurements of Higgs self-couplings and probing composite Higgs sector, and to search for supersymmetric particles [34].

### 3.4 Detector requirements

Many physics processes at a future linear collider can be characterised by multi-jet final states, often with charged leptons or missing momentum associated with neutrinos. The reconstruction of the invariant masses of two or more jets is crucial for event reconstruction and event selection. At the Large Electron–Positron Collider (LEP), kinematic fitting [35] allowed precise invariant mass reconstruction. At a future linear collider, reconstructing invariant mass of multiple jets for final states with missing momentum will rely heavily on the intrinsic jet energy resolution of the detector.

One goal of jet energy resolution at a future linear collider is to be able to separate  $W$  and  $Z$  bosons, by reconstructing their invariant masses via quark-jets using the hadronic decay channel. The idealised reconstructed  $W$  and  $Z$  boson mass distributions

for different jet mass resolutions are shown in figure 3.4. As the invariant mass resolution is comparable to the gauge boson widths, i.e.  $\sigma_m/m \approx \Gamma_W/m_W \approx \Gamma_Z/m_Z$ , a separation of  $2.5\sigma$  in the mass distributions implies a jet energy resolution of 3.5% [36] for a range of jet energies from 50 GeV to 1 TeV.



**Figure 3.4:** Ideal W/Z boson mass separation for different jet mass resolutions obtained using a Gaussian smearing of Breit–Wigner distribution, taken from [2].

### 3.5 Particle flow calorimetry

A jet energy resolution of 3.5% is unlikely to be achieved with a traditional calorimeter design. Traditionally, jet energies are measured as a sum of energies deposited in the electromagnetic (ECAL) and hadronic calorimeter (HCAL), giving a jet energy resolution of the form

$$\frac{\sigma_E}{E} = \frac{\alpha}{\sqrt{E(\text{GeV})}} \oplus \beta. \quad (3.1)$$

The stochastic term  $\alpha$  is typically greater than 60% [28, 30], and the constant term  $\beta$  is a few percent [28, 30]. To achieve a jet energy resolution of 3.5% or better, the stochastic term should be less than 30% with a small constant term, which is unlikely to be achieved by a traditional calorimeter.

In a typical jet, about 62% of the jet energy is from charged particles, 27% from photons, 10% from long-lived neutral hadrons, and 1.5% from neutrinos [37, 38]. In a traditional approach to calorimetry, the jet energy resolution is limited by the relatively poor energy resolution of the hadronic calorimeters.

The particle flow approach to calorimetry improves the jet energy resolution by fully reconstructing all visible particles in the detector. The jet energy is the sum of energies of individual particles, where the energies of the charged particles are measured in the tracking detectors, and the energies of neutral particles are measured in calorimeters. Hence the hadronic calorimeter only measures about 10% of the jet energy.

As shown in table 3.1, assuming 30% of the jet energy (photon energy) is measured with  $\sigma_E/E = 15\%/\sqrt{E(\text{GeV})}$ , and 10% of the jet energy (hadron energy) is measured with  $\sigma_E/E = 55\%/\sqrt{E(\text{GeV})}$  [30], a jet energy resolution of  $\sigma_E/E = 19\%/\sqrt{E(\text{GeV})}$  can be obtained. This satisfies the jet energy resolution requirement for separating W and Z bosons via their hadronic decays. In reality, this level of performance is unattainable due to incorrect association of energy deposits to particles. At jet energies beyond tens of GeVs, results from imperfect reconstruction rather than the intrinsic detector performance limit the particle flow performance [36].

Component	Detector	Energy fraction	Energy resolution	Jet energy resolution
Charged particles (C)	Tracker	$\sim 0.6E_j$	$10^{-4}E_C^2$	$< 3.6 \times 10^{-5}E_j^2$
Photons ( $\gamma$ )	ECAL	$\sim 0.3E_j$	$0.15\sqrt{E_\gamma}$	$0.08\sqrt{E_j}$
Neutral hadrons(N)	HCAL	$\sim 0.1E_j$	$0.55\sqrt{E_N}$	$0.17\sqrt{E_j}$

**Table 3.1:** Contributions from different particle components to the jet energy resolutions (all energies in GeV). The table lists the approximate fractions of charged particles, photons, and neutral hadrons in a jet of energy  $E_j$ , and the assumed single particle energy resolutions. The table is adapted from [36].

In the particle flow approach to calorimetry, the sum of calorimeter energies is replaced by a complex pattern-recognition problem, which is solved by the Particle Flow reconstruction Algorithm (PFA). Detailed simulations of the ILC and the CLIC detector concepts using the PandoraPFA [36, 39] particle flow reconstruction algorithms have demonstrated that a jet energy resolution of approximately 3% can be achieved for jet energies in the range of 100 GeV to 1 TeV.

Particle flow calorimetry works by fully reconstructing particles and associating calorimeter hits to tracks in tracking detectors. This places stringent requirements on the calorimeter designs. The ECAL and the HCAL need to be highly granular for an excellent spatial resolution to correctly associate calorimeter hits to the inner detector tracks. The tracking system needs to have an excellent momentum resolution for the momentum measurements of the charged particles.

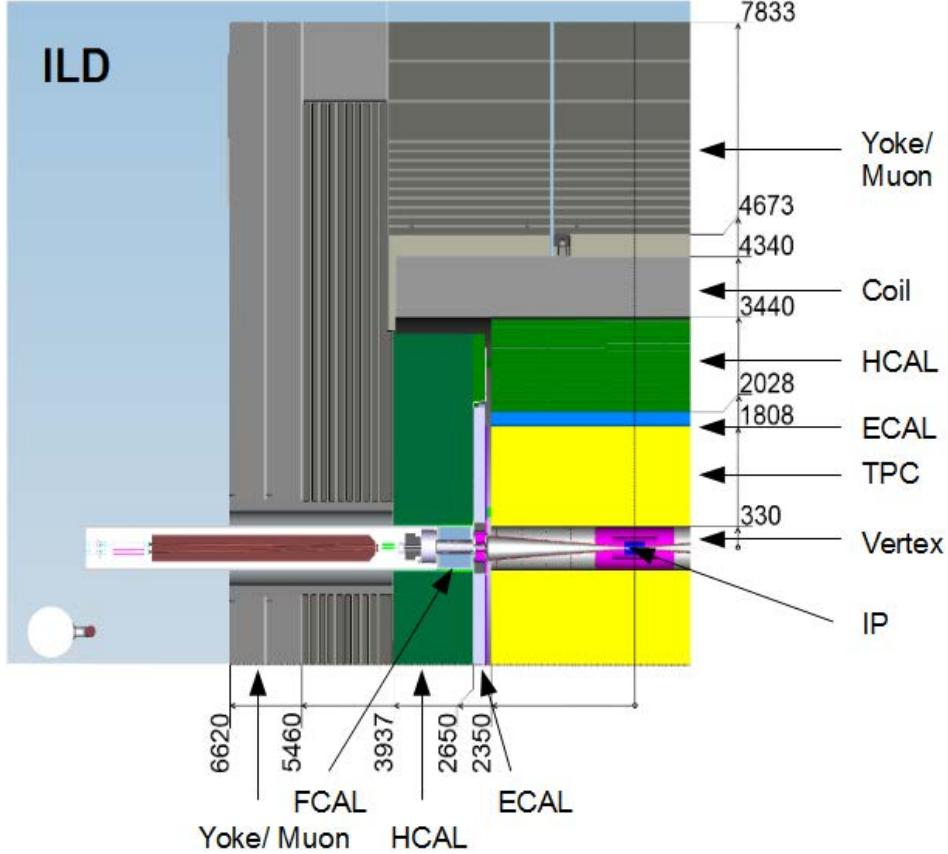
## 3.6 International Large Detector

Two detector concepts have been developed for the ILC. The motivation for having two detectors is to have multiple independent measurements within one collider for cross-checking, complementarity, and competition between collaborations. The two detectors are both designed to be general purpose detectors. The Silicon Detector, SiD [29], is a compact detector with silicon tracking modules and a magnetic field of 5 T. The other detector, the International Large Detector, ILD [28], is a larger detector with a time projection chamber as the main tracking detector.

The ILD detector concept has been optimised for particle flow techniques. Figure 3.5 shows the longitudinal cross section of top quadrant of the ILD detector concept. From the interaction point (IP) outwards, there is: a tracking system comprising a large time projection chamber (TPC) augmented with silicon tungsten layers; highly granular electromagnetic calorimeters (ECAL) and hadronic calorimeters (HCAL); forward calorimeters (FCAL); a superconducting solenoid; and muon chambers embedded within the iron return yokes. The key parameters of the ILD are listed in table 3.2. The section below describes the sub-detectors of the ILD detector concept referred to as the ILD\_o1\_v05 option in the MOKKA detector simulation [40], used for the ILD technical design report [30].

Component	ILD	CLIC_ILD
Tracker	TPC; Silicon	TPC; Silicon
Solenoid Field	3.5 T	4 T
Solenoid Field Bore	3.3 m	3.4 m
Solenoid Length	8.0 m	8.3 m
VTX Inner Radius	16 mm	31 mm
ECAL $r_{min}$	1.8 m	1.8 m
ECAL $\Delta r$	172 mm	172 mm
HCAL Absorber Barrel / Endcap	Fe / Fe	Fe / W
HCAL Interaction Length	$5.5 \lambda_I$	$7.5 \lambda_I$
Overall Height	14.0 m	14.0 m
Overall Length	13.2 m	12.8 m

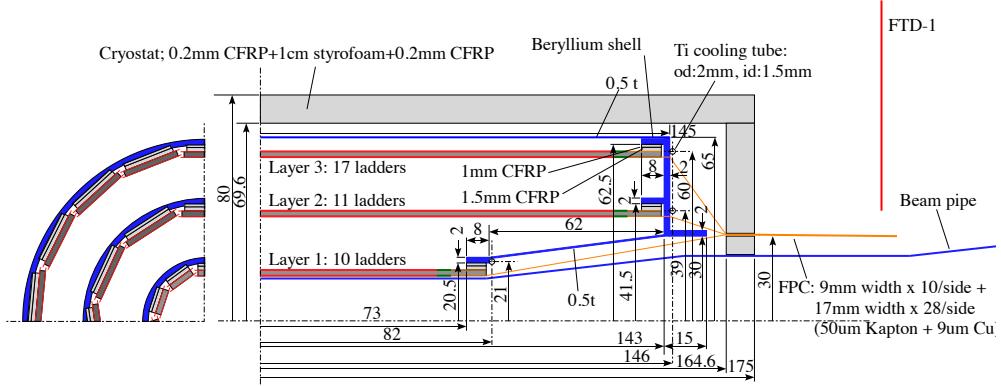
**Table 3.2:** A comparison of key parameters of the ILD and CLIC\_ILD detector concepts. ECAL  $r_{min}$  is the smallest distance from the calorimeter to the main detector axis. The table is adapted from [2].



**Figure 3.5:** The longitudinal cross section of top quadrant of the ILD, taken from [30]. From interaction point (IP) outwards, there is: a tracking system comprising a large time projection chamber (TPC) augmented with silicon tungsten layers; highly granular electromagnetic calorimeters (ECAL) and hadronic calorimeters (HCAL); forward calorimeters (FCAL); a superconducting solenoid; and muon chambers embedded within the iron return yokes. Dimensions are in units of mm.

### 3.6.1 Vertex detector

The pixel-vertex detector (VTX) needs to be close to the interaction point to reconstruct secondary vertices. As the TPC is the main tracking detector, the VTX mainly measures the impact parameter of tracks. Figure 3.6 shows the structure of the vertex detector. The structure is of three, nearly cylindrical, concentric layers of double-sided ladders. Each ladder contains pixel sensors on both sides at 2 mm separation between two layers. This results in six measured positions for each charged particle traversing the detector. The first double layer is of half the length of the other two, to avoid the high occupancy region of direct low-momentum hits from the incoherent pair background. The baseline geometry of the vertex detector can be found in table 3.3. The radii covered by the detector range from 16 mm to 60 mm.



**Figure 3.6:** Structure of the ILD vertex detector, taken from [30].

	R	$ z $	$ \cos(\theta) $
Layer 1	16 mm	62.5 mm	0.97
Layer 2	18 mm	62.5 mm	0.96
Layer 3	37 mm	125 mm	0.96
Layer 4	39 mm	125 mm	0.95
Layer 5	58 mm	125 mm	0.91
Layer 6	60 mm	125 mm	0.90

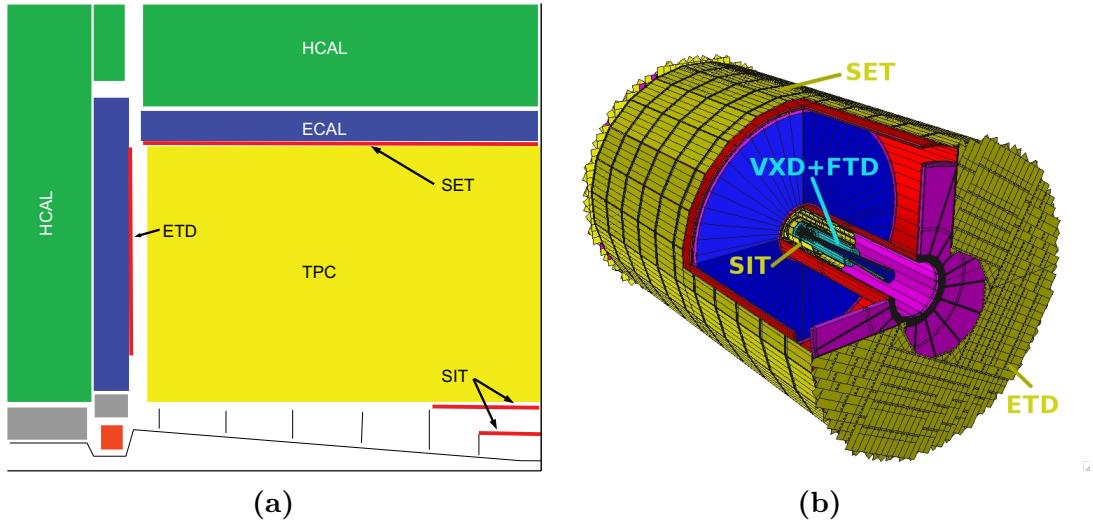
**Table 3.3:** Key parameters of vertex detector in the ILC. The table is adapted from [30].

### 3.6.2 Tracking detectors

The hybrid tracking system consists of a large time projection chamber (TPC), a Silicon Inner Tracker (SIT), a Silicon External Tracker (SET) in the barrel region, a silicon endcap tracking component (ETD) behind the endplate of the TPC, and a silicon forward tracker (FTD) in the forward region. A top quadrant view of the ILD silicon envelope system with the TPC is shown in figure 3.7. The SIT, SET, and ETD are made up of two single-sided strip layers tilted by a small angle. The FTD is a system of two silicon-pixel disks and five silicon-strip disks. The main parameters of the silicon system and the TPC can be found in table 3.4.

A TPC tracking detector has several advantages: a) tracks can be measured with a large number of three-dimensional ( $r, \phi, z$ ) spatial points; b) the continuous tracking allows precise reconstruction of tracks; and c) the TPC uses a minimum amount of material, which minimises the photon to electron pair conversion.

The silicon intermediate tracker (SIT) and the silicon envelope tracker (SET) provide spatial point measurements before and after the TPC in the barrel region. This helps to improve the overall momentum resolution by providing points to link the vertex detector with the TPC, and to extrapolate tracks from the TPC to the calorimeters. The FTD improves the low angle coverage of the tracking system, where the low angle is not covered by the TPC.



**Figure 3.7:** a) A top quadrant view of the ILD silicon envelope system, SIT, SET, FTD, and ETD, with TPC, ECAL, and HCAL, and b) a 3D detailed GEANT4 simulation description of the silicon system as sketched in the quadrant view in a). Both plots are adapted from figures in [30].

	R	z	$\cos(\theta)$
SIT	153 mm	368 mm	0.910
SIT	300 mm	644 mm	0.902
SET	1811 mm	2350 mm	0.789
ETD	419 - 1822.7 mm	2420 mm	0.985 - 0.799
TPC	329 - 1808 mm	$\pm 2350$ mm	up to 0.98

**Table 3.4:** Main parameters of the central silicon tracking systems (SIT, SET, and ETD) and the TPC. The table is adapted from [30].

### 3.6.3 Electromagnetic calorimeter

The silicon–tungsten sampling electromagnetic calorimeters in the ILD consist of an octagonal barrel and two endcap systems. The fine granularly ECAL is located inside the HCAL. Both ECAL and HCAL are inside the superconducting solenoid. Figure

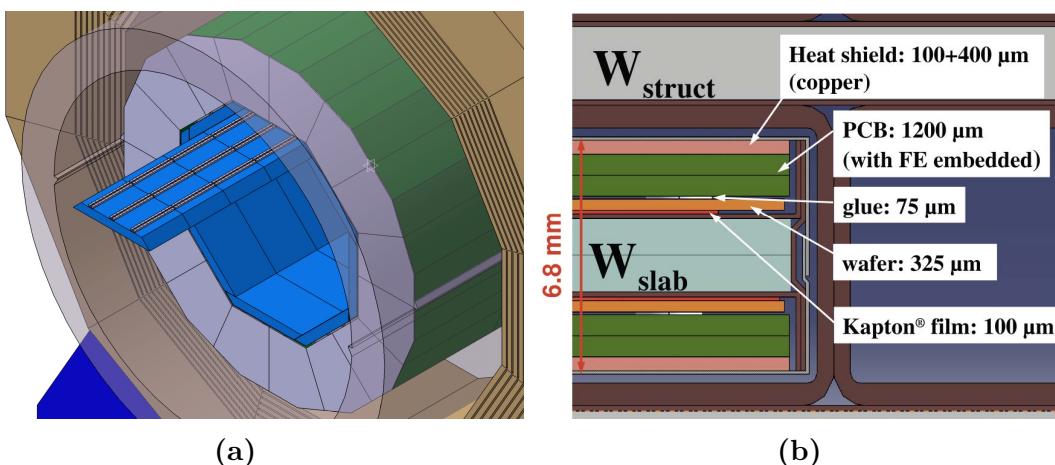
[3.8a](#) shows the position of the electromagnetic calorimeter in the ILD detector and the trapezoidal form of the modules.

The particle flow paradigm has a large impact on the ECAL design. In addition to measuring and separating photons, the ECAL needs to allow the reconstruction of detailed shower profiles to separate electromagnetic showers from hadronic showers, as approximately 50% of hadronic showers start in the ECAL.

From test beam data and simulation studies [41–43], a sampling calorimeter with a longitudinal segmentation below one radiation length and the transverse segmentation below one Molière radius is required. A compact design is realised with tungsten as the absorber material and silicon pad diodes as the active material. A cross section of an ECAL layer is shown in figure [3.8b](#). Tungsten is a dense material with a large ratio of interaction length to radiation length. A small radiation length will promote the start of the electromagnetic shower earlier in the calorimeter, whilst a large interaction length will reduce the fraction of hadronic showers starting in the ECAL.

The ECAL, which is about 20 cm thick, has 30 longitudinal layers providing about 24 radiation lengths. The first 20 layers use 2.1 mm thick absorber plates and the last 10 layers have 4.2 mm thick absorber plates.

The choice of thin silicon layers offers a great spatial resolution. The chosen size of  $5.1 \times 5.1 \text{ mm}^2$  silicon pads provides enough segmentation to meet the requirements of the particle flow paradigm [30].



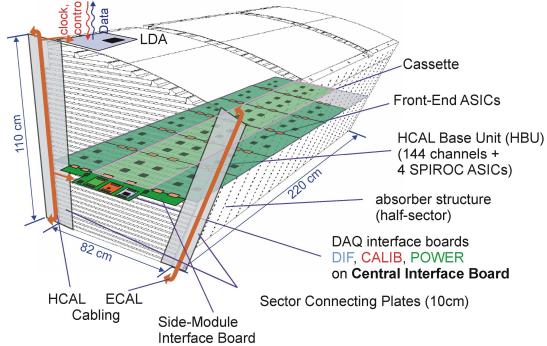
**Figure 3.8:** a) The electromagnetic calorimeter (in blue) within the ILD detector. b) A cross section through the electromagnetic calorimeter layers. Both plots are taken from [30].

### 3.6.4 Hadronic calorimeter

The principal role of the HCAL is to separate neutral hadron showers from other particles, and to measure (neutral) hadron energies. The ILD HCAL is a sampling calorimeter with steel absorber and scintillator tiles as the active medium. The layout of the HCAL is 48 longitudinal layers with  $3 \times 3 \text{ cm}^2$  scintillator tiles, using an analogue read out system. The layout of a technological prototype, the "EUDET prototype" [44] is shown in figure 3.9.

For the absorber material, stainless steel is chosen for mechanical and calorimetric reasons. Steel allows a self-supporting structure without auxiliary supports. At the same time, iron has a moderate ratio of hadronic interaction length ( $\lambda_I = 17 \text{ cm}$ ) to electromagnetic radiation length ( $X_0 = 1.8 \text{ cm}$ ), which allows a fine longitudinal sampling in  $X_0$  with a reasonable number of layers in a given total hadronic absorption length. The longitudinal system, including the ECAL, provides about 6 interaction lengths, which is sufficient to contain the hadronic showers.

The scintillator tiles provide both energy and position measurements. The transverse segmentation, chosen by optimisation studies [36] is about  $3 \times 3 \text{ cm}^2$ . This level of segmentation is sufficient to meet the requirement of the particle flow paradigm [30].

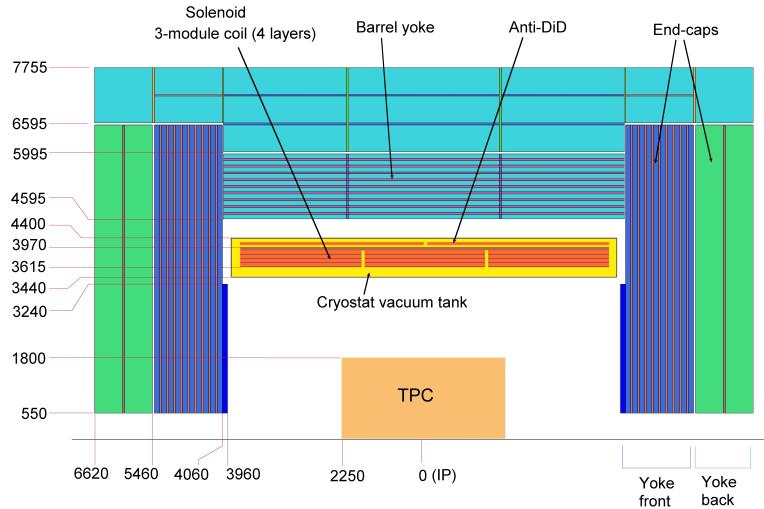


**Figure 3.9:** The schematic view of a CALICE analogue HCAL technological prototype module, taken from [30].

### 3.6.5 Solenoid, yoke, and muon system

A large superconducting solenoid, outside the calorimeters, produces a nominal 3.5 T magnetic field. Figure 3.10 shows the cross section of the ILD magnet.

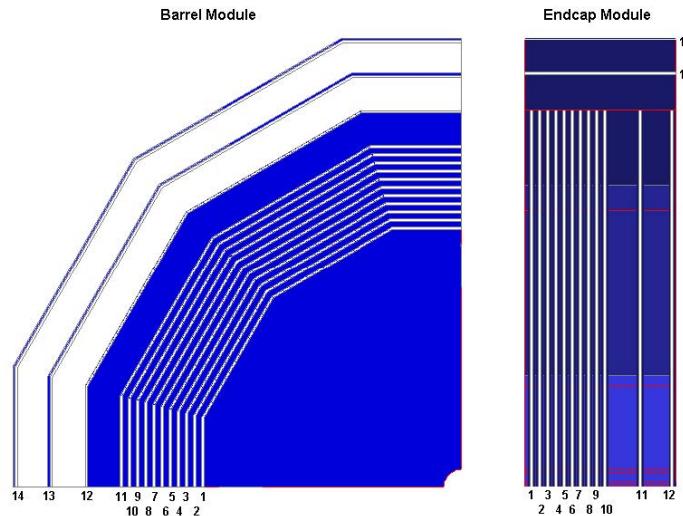
The iron yoke returns the magnetic flux. The yoke is designed to ensure that the maximum magnetic field at 15 m radial distance from the detector is 50 Gauss to ensure safety [45].



**Figure 3.10:** The ILD magnet cross section. Dimensions are in mm. Figure is taken from [30].

The iron yoke is also instrumented with scintillator strips as active layers to act as a muon detector. A highly efficient muon detector is provided by  $3 \times 3 \text{ cm}^2$  scintillator tiles. The layout of the muon detector is shown in figure 3.11.

The first layer of the muon detector, also acting as a tail catcher calorimeter, catches the energy leakage from the HCAL and the ECAL. It has been shown that a 10% improvement of single particle energy resolution is possible with the tail catcher [46].



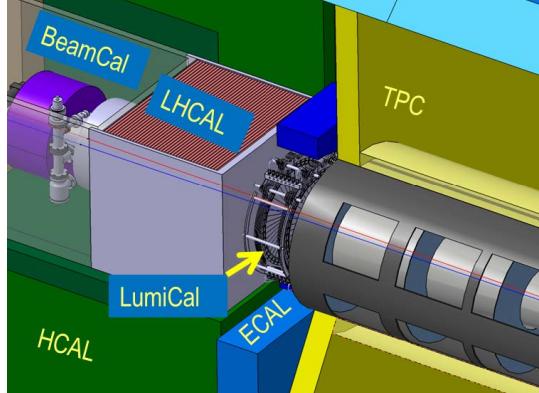
**Figure 3.11:** Sensitive layers of the ILD muon system, taken from [30].

### 3.6.6 Very forward calorimeters

The detectors in the forward region provide luminosity measurements and forward coverage of calorimeters. A system of precision and radiation resistant calorimeters is required. Figure 3.12 shows the forward calorimeters of the ILD.

The luminosity calorimeter (LumiCAL) counts Bhabha scattering to measure the luminosity to a precision of  $10^{-3}$  at a centre-of-mass energy of 500 GeV [47]. The beam calorimeter (BeamCAL), which is hit by many beamstrahlung pairs after each bunch crossing, extends the forward coverage. The BeamCAL also provides a measurement of the bunch-by-bunch luminosity. An additional hadron calorimeter in the forward region, LHCAL, extends the angular coverage of the HCAL to that of the LumiCAL.

The calorimeters in the forward region also provide enough information for high-energy electron tagging [48], which aids event reconstruction at a high centre-of-mass energy. Table 3.5 lists the key parameters of the LumiCAL and the BeamCAL in the ILD.



**Figure 3.12:** The calorimeters in the forward region of the ILD, taken from [30]. The LumiCAL, the BeamCAL, and the LHCAL are the luminosity calorimeter, the beam calorimeter, and the forward hadronic calorimeter, respectively.

## 3.7 Detector optimisation

Detector optimisation studies were performed to choose the optimal parameters of the ILD sub-detectors [30]. Here the optimisation studies for the ECAL and the HCAL are presented.

		ILD	CLIC_ILD
LumiCAL	Geometrical acceptance	31 - 77 mrad	38 - 110 mrad
	Fiducial acceptance	41 - 67 mrad	44 - 80 mrad
	z (start)	2450 mm	2654 mm
	Number of layers (W + Si)	30	40
BeamCAL	Geometrical acceptance	5 - 40 mrad	10 - 40 mrad
	z (start)	3600 mm	3281 mm
	Number of layers (W + sensor)	30	40
	Graphite layer thickness	100 mm	100 mm

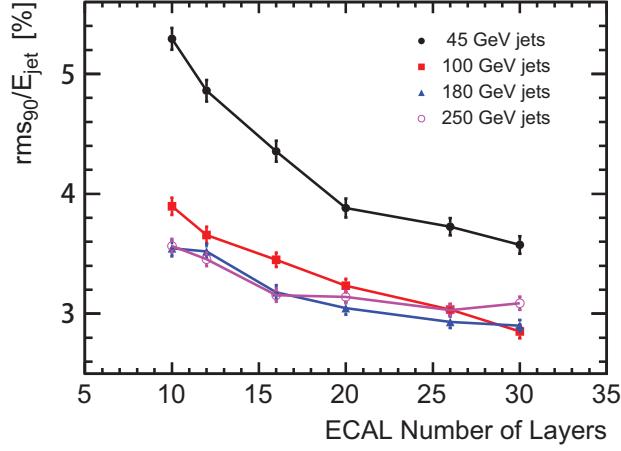
**Table 3.5:** Comparison of the key parameters of the LumiCAL and the BeamCAL at the ILD and the CLIC\_ILD detector concepts. The table is adapted from [2].

### 3.7.1 Electromagnetic calorimeter optimisation

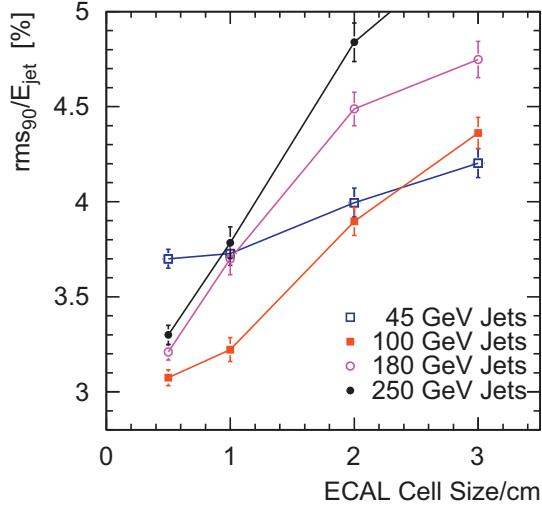
The metric for optimisation is the jet energy resolution, which is defined as the root-mean-square divided by the mean for the smallest width of distribution that contains 90% of entries, using  $e^+e^- \rightarrow Z'Z'$  events, where  $Z' \rightarrow u\bar{u}/d\bar{d}/s\bar{s}$ , at barrel region. The angular cut is to avoid the barrel/endcap overlap region. The light quark decay of the  $Z'$  is used to avoid the complication of missing momentum from semi-leptonic decay of heavy quarks. Using 90% of the entries is robust and focuses on the Gaussian part of the distribution.

The optimisation of the ILD ECAL design was performed as a function of the number of longitudinal layers, whilst keeping other geometry constant. Figure 3.13 shows the jet energy resolution for a single jet as a function of the number of longitudinal layers for four different jet energies. For a 45 GeV jet, a degradation of 10% in the jet energy resolution is observed when the number of layers decreases from 30 to 20. The degradation in the jet energy resolution is significant for number of layers fewer than 20, although the impact is smaller for high energy jets. Therefore, 30 longitudinal layers is chosen for the ECAL.

The ILD ECAL design was optimised as a function of the transverse cell sizes, whilst keeping other geometry constant. Figure 3.14 shows the jet energy resolution for a single jet, plotted as a function of transverse scintillator cell sizes for four different jet energies. The  $10 \times 10 \text{ mm}^2$  cell size is needed to meet the jet energy requirement of  $\sigma_E/E < 3.8\%$  for the jet energies relevant at  $\sqrt{s} = 1 \text{ TeV}$ , with  $5 \times 5 \text{ mm}^2$  cell size being preferable.



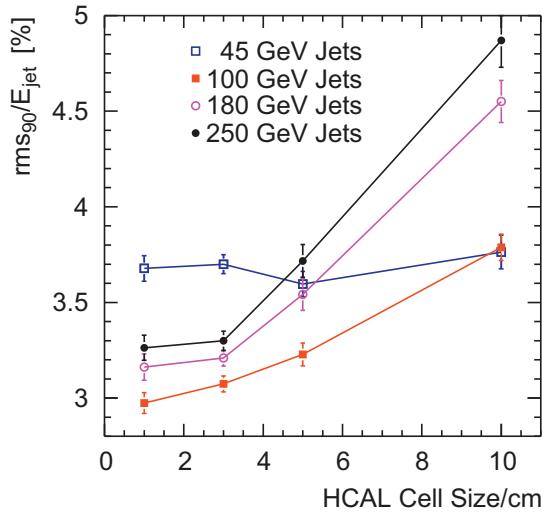
**Figure 3.13:** The single jet energy resolution as a function of the number of longitudinal ECAL layers, adapted from [30].



**Figure 3.14:** The single jet energy resolution as a function of the ECAL transverse cell sizes, adapted from [36].

### 3.7.2 Hadronic calorimeter optimisation

For the ILD HCAL design, the transverse cell sizes has been optimised using the jet energy resolution as the metric. The jet energy resolutions as a function of HCAL scintillator square cell sizes for four different jet energies are shown in figure 3.15. There is no substantial gain in the jet energy resolution for cell sizes below 3 cm. However, the jet energy resolution degrades for cell sizes above 3 cm. Hence  $3 \times 3 \text{ cm}^2$  scintillator cell size is chosen for the HCAL design.



**Figure 3.15:** The single jet energy resolution as a function of the hadronic calorimeter scintillator cell sizes, adapted from [30].

### 3.8 CLIC\_ILD detector concept

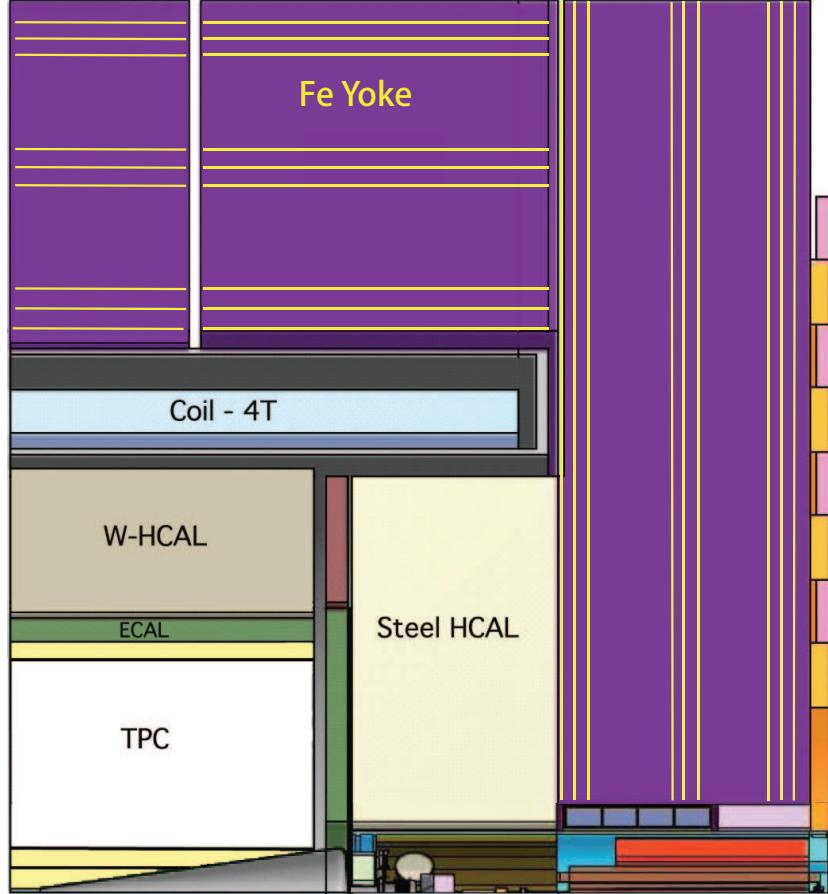
There are two detector concepts studied in the CLIC conceptual design report [2], the CLIC\_ILD and the CLIC\_SiD concepts. The CLIC\_ILD detector concept is based on the ILD detector concept. Figure 3.16 shows the longitudinal cross section of the CLIC\_ILD detector. A comparison of key parameters of the ILD and the CLIC\_ILD detector concepts is shown in table 3.2.

For the CLIC\_ILD vertex detector, the first layer is moved outwards by 15 mm due to a larger high occupancy region with a higher centre-of-mass energy and a smaller beam jet. The detector is also required to provide time stamping at the nanosecond level.

For the CLIC\_ILD tracking detector, the same silicon–TPC hybrid structure is used. At CLIC, it is challenging to use a TPC to separate two tracks in high-energy jets and to identify events in the collection of 312 bunch crossings in 156 ns. The outer silicon tracking system is important to achieve a high momentum resolution at high centre-of-mass energy. The solid angle coverage of the tracking detector is  $12^\circ \lesssim \theta \lesssim 168^\circ$ .

For the CLIC\_ILD design, the same ECAL as the ILD is assumed, as the requirements of a CLIC detector are satisfied by the ECAL design at the ILD.

For the CLIC\_ILD HCAL, extra layers are added to contain the hadronic shower for the higher centre-of-mass energies of CLIC. The increased thickness is justified by the simulation studies [2], where the jet energy resolution degrades quickly for a thinner



**Figure 3.16:** The longitudinal cross section of top quadrant of the CLIC\_ILD detector concept, taken from [2]. From interaction point (IP) outwards, there is: a tracking system comprising a large time projection chamber (TPC) augmented with silicon tungsten layers; highly granular electromagnetic calorimeters (ECAL) and hadronic calorimeters (HCAL); forward calorimeters (FCAL); a superconducting solenoid; and muon chambers embedded within the iron return yokes.

HCAL. To sustain the same inner bore radius as the ILD detector solenoid, a more dense material, tungsten, is chosen as the absorber material in the HCAL barrel.

The magnetic field is increased to 4 T for a better jet energy resolution [36] at a higher centre-of-mass energy. Due to the different magnetic field strength, the iron yoke thickness is increased to 230 cm.

The CLIC\_ILD adopts a similar very forward calorimetry system as that of the ILD. The dimensions of the elements are changed due to a difference in the beam crossing angles (20 mrad for CLIC and 14 mrad for the ILC). A comparison of the key parameters of the LumiCAL and the BeamCAL at the ILD and the CLIC\_ILD is shown in table 3.5.

# Chapter 4

## Event Generation, Simulation, Reconstruction, and Analysis

*‘When I walk along with two others, from at least one I will be able to learn.’*

— Confucius, 551 BC – 479 BC

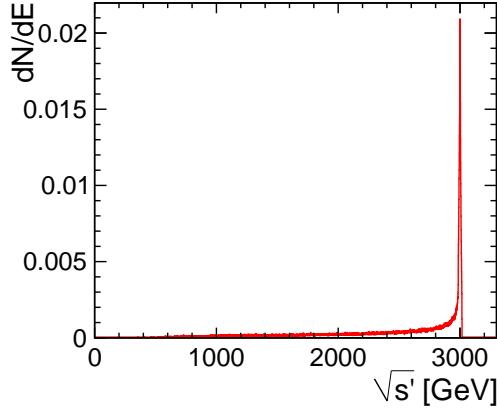
In this chapter, event generation, simulation, reconstruction, and analysis software for the future linear colliders is discussed. Reconstruction of particle flow objects with PandoraPFA is presented, which is the framework used for the photon reconstruction algorithms developed in chapter 5. The Boosted Decision Tree multivariate analysis technique, which is used in a number of places in this thesis, is described in detail.

### 4.1 Event generation

Most of the Monte Carlo (MC) samples used in this thesis were generated using WHIZARD generator software [49, 50]. The Initial State Radiation (ISR) effect is simulated in WHIZARD with the ISR photons collinear with the beam direction. Events were generated with head-on  $e^+e^-$  collisions. PYTHIA [51] was used to describe parton showering, hadronisation and fragmentation. The fragmentation parameters of PYTHIA were tuned to OPAL data [52] from the Large Electron-Positron Collider (LEP). The Final State Radiation (FSR) is treated by PYTHIA with its default parameters. TAUOLA [53] was used to describe the tau lepton decay with correct spin correlations of the tau decay products.

### 4.1.1 CLIC luminosity spectrum

The small CLIC beam size at the interaction point, which indicates a high bunch charge density, results in electrons and positrons radiating strongly in the electromagnetic field of the other beam. This effect is known as beamstrahlung. Consequently, the centre-of-mass energies of the  $e^+e^-$  collisions have a long low energy tail towards lower values than the nominal centre-of-mass energy. The luminosity spectrum for CLIC operating at  $\sqrt{s} = 3$  TeV is shown in figure 4.1, generated with GUINEAPIG [54]. Only 35% of the effective luminosity falls within 1% of the nominal centre-of-mass energy, for CLIC operating at  $\sqrt{s} = 3$  TeV.



**Figure 4.1:** The luminosity spectrum for CLIC operating at  $\sqrt{s} = 3$  TeV, taken from [2].

Due to the presence of a large amount of beamstrahlung photons, the photon–photon and electron/positron–photon background must be accounted for. The instantaneous luminosities for the electron–positron, electron/positron–photon, and photon–photon interactions at a  $\sqrt{s} = 3$  TeV CLIC machine are listed in table 4.1, obtained from GUINEAPIG.

Instantaneous luminosity / $10^{34}\text{cm}^{-2}\text{s}^{-1}$ $\sqrt{s} = 3\text{ TeV}$	
$e^+e^-$	6.7
$e^+\gamma$	5.3
$e^-\gamma$	5.3
$\gamma\gamma$	4.6

**Table 4.1:** Instantaneous luminosity for the electron–positron, electron/positron–photon, and photon–photon interactions at a  $\sqrt{s} = 3$  TeV CLIC machine. Values are taken from [55].

The GUINEA $\text{PIG}$  prediction overestimates the luminosities for interactions that involve beamstrahlung photons [55]. Hence luminosity corrections are applied to the interactions that involves Beamstrahlung photons according to table 4.2.

Luminosity correction	$\sqrt{s} = 1.4 \text{ TeV}$	$\sqrt{s} = 3 \text{ TeV}$
$L(e^+\gamma) / L(e^+e^-)$	0.75	0.79
$L(e^-\gamma) / L(e^+e^-)$	0.75	0.79
$L(\gamma\gamma) / L(e^+e^-)$	0.64	0.69

**Table 4.2:** Correction to integrated luminosities of the positron–photon, electron–photon, and photon–photon interactions where photons are from the Beamstrahlung process for CLIC at  $\sqrt{s} = 1.4 \text{ TeV}$  and  $3 \text{ TeV}$ . The corrections, normalised to the electron–positron luminosity, are taken from [55].

## 4.2 Event simulation

Studies presented in this thesis are based on event samples, fully simulated and reconstructed in ILD or CLIC\\_ILD detector models. GEANT4 [56] is used to simulate the interactions of particles through the detector material. The ILD and CLIC\\_ILD detector geometries are implemented in the MOKKA package [40]. The QGSP\\_BERT physics list from GEANT4 is used to simulate the detailed development of hadronic showers in the detector. The beam crossing angle (20 mrad for CLIC and 14 mrad for the ILC) is introduced by applying a corresponding Lorentz boost to all generated particles in the events prior to detector simulation.

In chapter 8, which presents a double Higgs production analysis using the CLIC\\_ILD detector, it is necessary to account for the pile up from beam induced background.

### 4.2.1 CLIC beam induced background

At CLIC, pile-up from beam induced background needs to be considered. The two most significant types of background at CLIC are  $\gamma\gamma \rightarrow$  hadrons and incoherent  $e^+e^-$  pairs [2]. The  $\gamma\gamma \rightarrow$  hadrons background is produced when the interaction of real and virtual photons from the colliding beams leads to two-photon interactions, resulting in hadronic final states [57, 58]. The incoherent  $e^+e^-$  pairs are produced with interactions of both real or virtual beamstrahlung photons with individual particles of the other beam, producing  $e^+e^-$  pairs in the strong electromagnetic field of the other beam [59].

Simulation of  $\gamma\gamma \rightarrow$  hadrons uses the photon spectrum from GUINEA<sub>PIG</sub> and a parametrisation of the total cross section of the  $\gamma\gamma \rightarrow$  hadrons process [60]:

$$\sigma_{\gamma\gamma}(s_{\gamma\gamma}) = 211 \text{ nb} \left( \frac{s_{\gamma\gamma}}{\text{GeV}^2} \right)^{0.0808} + 215 \text{ nb} \left( \frac{s_{\gamma\gamma}}{\text{GeV}^2} \right)^{-0.4525}, \quad (4.1)$$

where  $s_{\gamma\gamma}$  is the Mandelstam  $s$  variable for the two photons. On average, there are 3.2  $\gamma\gamma \rightarrow$  hadrons events per bunch crossing within the detector acceptance at  $\sqrt{s} = 3 \text{ TeV}$  with a  $\gamma\gamma$  centre-of-mass energy greater than 2 GeV [61]. PYTHIA is used to simulate the hard interaction and the hadronisation of these  $\gamma\gamma \rightarrow$  hadrons events. Table 4.3 shows the average energy deposited from  $\gamma\gamma \rightarrow$  hadrons and the incoherent  $e^+e^-$  pairs in different parts of the CLIC\_ILD detector. The energies in the calorimeter are integrated over the 300 ns from the start of the bunch train. The incoherent  $e^+e^-$  pairs are the dominant background in the HCAL endcaps due to the interactions of the large incoherent  $e^+e^-$  pairs in the BeamCAL, resulting in low-energy neutrons depositing energies in the HCAL endcaps. Except in the HCAL endcaps, the  $\gamma\gamma \rightarrow$  hadrons is the dominant background in all calorimeters.

Subdetector	Incoherent Pairs (TeV)	$\gamma\gamma \rightarrow$ hadrons (TeV)
ECAL Endcaps	2	11
ECAL Barrel	-	1.5
HCAL Endcaps	16	6
HCAL Barrel	-	0.3
Total Calorimeter	18	19
Central Tracker	-	7

**Table 4.3:** Average energy deposited from  $\gamma\gamma \rightarrow$  hadrons and the incoherent  $e^+e^-$  pairs in different parts of the CLIC\_ILD subdetectors. Numbers correspond to the background for an entire CLIC bunch train. The reconstructed calorimeter energies are integrated over 300 ns from the start of the bunch train. The table is adapted from [2].

For the study presented in chapter 8, only the  $\gamma\gamma \rightarrow$  hadrons background is included in the simulation. The hits from simulated  $\gamma\gamma \rightarrow$  hadrons events from 60 bunch crossings are superimposed on simulated  $e^+e^-$ ,  $e^\pm\gamma$ , and  $\gamma\gamma$  collisions before the event reconstruction. The included background corresponds to an integration time window of  $-5 \text{ ns}$  to  $+25 \text{ ns}$  around the generated physics event. Bunch trains are separated by 0.5 ns to mimic the CLIC train structure [2]. For each bunch crossing, the number of  $\gamma\gamma \rightarrow$  hadrons events included is drawn from a Poisson distribution with a mean of 3.2.

## 4.3 Event reconstruction

The linear collider reconstruction software runs in the MARLIN framework [62]. The event reconstruction consists of three main steps: digitisation of simulated tracks and calorimeter hits, reconstruction of tracks in the tracking system [63], and reconstruction of particle flow objects (PFOs) with PandoraPFA [36, 39].

Several different MARLIN processors are used to reconstruct tracks: CLUPATRA [63] is used to reconstruct tracks in the TPC; FORWARDTRACKING [63] is used to reconstruct tracks in the FTD; and SILICONTRACKING [63] is used to reconstruct tracks in other silicon tracking detectors. A final MARLIN tracking processor, FULLLDCTRACKING [63], is used to combine tracks segments produced from individual processors.

### 4.3.1 PandoraPFA

PandoraPFA was developed for the ILD detector concept [36] at the ILC and at CLIC [2, 39]. It has also been used for the SiD detector concept [29] and the studies for the CMS endcap calorimeter upgrade [64]. The latest improvement to PandoraPFA, including a better internal memory management, is summarised in [65].

PandoraPFA adopts a multi-algorithm approach to the particle flow object reconstruction. There are over sixty  $e^+e^-$  linear collider specific reconstruction algorithms. Each algorithm aims to address a particular topological issue in the reconstruction. There are nine main steps in the PandoraPFA reconstruction, listed below:

- Track processing: identifies special topologies,  $V^0$ s and kinks, of tracks and prepare them for subsequent reconstruction;
- Calorimeter hit processing: calculates the properties of a calorimeter hit, such as its position and its energy response from the calorimeter digitiser. The calorimeter hits are used in subsequent reconstruction;
- Particle identification algorithms find calorimeter hits associated with neutral particles, such as photons. Chapter 5 describes photon reconstruction algorithms in detail;
- An initial clustering step group calorimeter hits into clusters. This step is discussed further in section 4.3.1.1, since it is used in photon reconstruction algorithms described in Chapter 5;

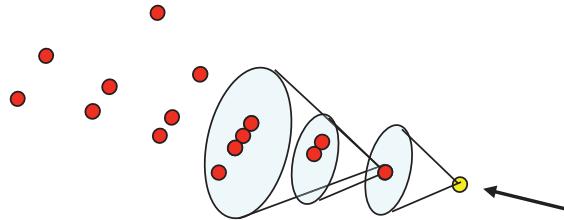
- Topological cluster association merges clusters based on clear topological signatures. Merging signatures include combining track segments, connecting track segments with gaps, connecting track segments to hadronic showers, and merging clusters when they are within close proximity.
- Track–cluster matching step associates clusters to tracks obtained from the tracking detectors;
- Re-clustering step improves the compatibility of the cluster energy and the associated track momentum on a statistical basis. This step is important for events with a dense jet environment and jets above 50 GeV;
- Fragment removal step focuses on merging clusters that are likely to be fragments of other particles. Algorithms for photon fragment merging are described in chapter 5;
- Particle Flow Object creation is the last step of the PandoraPFA reconstruction. It creates the output objects, Particle Flow Objects (PFOs). The reconstructed particles, PFOs, are the bases of all subsequent analyses.

#### 4.3.1.1 Initial clustering

In PandoraPFA, cone-based clustering algorithms are used to group calorimeter hits into clusters, illustrate in figure 4.2. The cone-based clustering algorithm is used because the direction of the particle flow is largely unchanged from the original particle. The seed for the cone clustering can be a projection of a track onto the front of the ECAL, and the initial cone direction is the direction of the track projection. Alternatively, the seed can be a calorimeter hit, and the initial cone direction is the direction from the IP to the calorimeter hit. Afterwards, a cone with a specified opening angle is formed around the direction of the seed. The building of the cone is iterated from the inner layer of the ECAL to the outer layer. At each layer, possible associations with the cone are made by considering calorimeters hits in previous layers and the same layer. If a calorimeter hit is not associated with any cone, the hit is used to seed a new cluster.

#### 4.3.2 CLIC beam induced background suppression

The analysis chapter 8 is performed for the CLIC\_ILD detector, where the CLIC beam induce background must be considered. The output of the PandoraPFA is a list of reconstructed particles. Two packages have been developed to suppress the background from pile up of  $\gamma\gamma \rightarrow$  hadrons: TRACKSELECTOR and PFOSELECTOR [39].



**Figure 4.2:** Illustration of the cone-based clustering algorithm used in PandoraPFA, taken from [66]

The TRACKSELECTOR [39] package removes poor quality and fake tracks that are likely to be from the beam induced background. It examines the number of track hits in individual tracking subdetectors and imposes track-quality cuts. It also places a cut on the arrival time of the track onto the front of the ECAL. If the arrival time of the track onto the front of the ECAL, using the helical fit of the track trajectory, differs more than 50 ns from using a straight line fit, the track will be rejected.

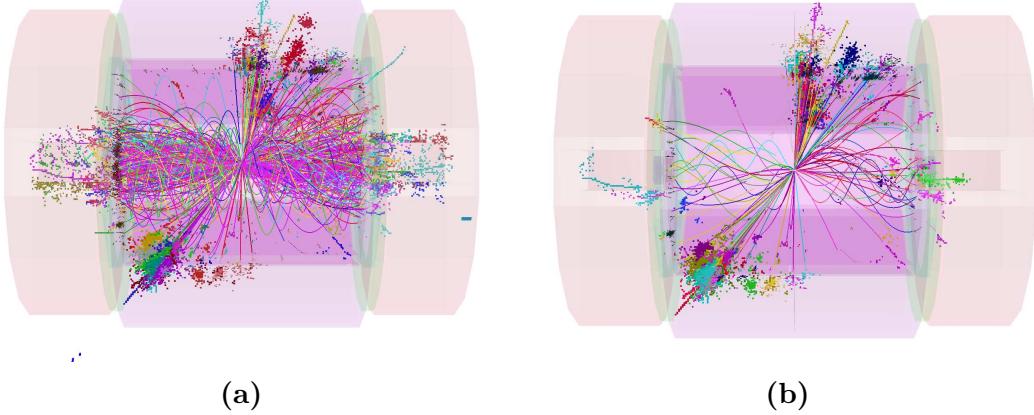
The PFOSELECTOR [39] package discards reconstructed particles that are likely to be from the beam induced background from the event reconstruction, based on the transverse momentum ( $p_T$ ) and time information of the reconstructed particles. The reconstructed particles from  $\gamma\gamma \rightarrow \text{hadrons}$  often have low  $p_T$  and are distributed in time across the 10 ns reconstruction integration timing window. In contrast, the PFOs from physics processes have a range of  $p_T$ , and have times close to the time of the brunch crossing that contains the event. The time of the brunch crossing containing the event is determined by the high-level software trigger.

PFOSELECTOR uses different  $p_T$  and timing cuts for the central part and the forward part of the detector. PFOSELECTOR also uses different  $p_T$  and timing cuts for different types of particles: photons, neutral PFOs, and charged PFOs.

Three configurations of these cuts were developed: “loose”, “normal”, and “tight” PFO selections. As the name suggested, “loose” PFO selection corresponds to a looser cut of  $p_T$  and time, preserving PFOs with a larger range of  $p_T$  and a larger range of times than the “tight” PFO selection.

Figure 4.3 shows the effect of the background suppression with the tight PFO selection. Figure 4.3a shows reconstructed particles in a simulated  $e^+e^- \rightarrow HH \rightarrow t\bar{b}b\bar{t}$  event in the CLIC\_ILD detector model, assuming an integration time window of 10 ns (100 ns in HCAL barrel), with 60 bunch crossings of  $\gamma\gamma \rightarrow \text{hadrons}$  background overlaid. 60 bunch crossings corresponds to an integration time window of 30 ns assuming a CLIC beam

structure, which is sufficient to account for main effect of the background. The effect of applying tight PFO selection cuts is shown in figure 4.3b. The energy deposited in the detector by the background is reduced from 1.2 TeV to the level of 100 GeV.



**Figure 4.3:** Reconstructed particles in a simulated  $e^+e^- \rightarrow HH \rightarrow t\bar{b}b\bar{t}$  event, integrated over a time window of 10 ns (100 ns in HCAL barrel) in the CLIC\_ILD detector model, with 60 bunch crossings of  $\gamma\gamma \rightarrow \text{hadrons}$  background overlaid in a). The effect of applying tight PFO section cuts is shown in b). The energy deposited in the detector by the background is reduced from 1.2 TeV to the level of 100 GeV. Figures are taken from [39].

## 4.4 Analysis software

### 4.4.1 Monte Carlo truth linker

It is important to be able to associate reconstructed particles to the Monte Carlo particles for the purpose of algorithm development and event selection optimisation. The MC truth linker processor provides the link between a MC particle and a reconstructed calorimeter hit. From the link, the MC particle contributing the most to a reconstructed particle can be determined, based on highest sum of the energies of calorimeter hits with the link to the same MC particle. Similarly, the MC particle contributing the most to a group of reconstructed particles (a jet) can be determined.

### 4.4.2 Jet algorithms

Jets are the results of hadronisation processes from high energy quarks or gluons. A jet is typically a visually obvious structure in an event display. The momentum and the direction of a jet are largely the same as the original particle. Despite the relative simplicity of identifying jets visually, it is a challenge for a pattern recognition program

to identify jets effectively and efficiently. Early work on jet finding started in 1977 [67], where descriptions on later developments can be found in reviews [68–70].

There are two large families of jet finding algorithms: cone based algorithms, and sequential combination algorithms. Here the focus is on the sequential combination algorithms.

Sequential combination algorithms typically calculate a pair-wise distance metric between a seed and a particle. The particle with the smallest metric is combined with the seed and into the jet. The distance metric is updated after a combination. This procedure is repeated until some stopping criteria are satisfied. The different jet algorithms typically differ in the definitions of distance metrics and stopping criteria.

The jet algorithm implementation used in this thesis is the FastJet C++ software package [71, 72]. The notations in the subsequent discussion follow the convention in [71].

#### 4.4.2.1 Longitudinally invariant $k_t$ algorithm

The longitudinally invariant  $k_t$  algorithm [73, 74] is one of the common sequential combination algorithms used in the pp collider experiments. There are two variants of the algorithm: inclusive and exclusive. In the inclusive variant, the symmetrical pair-wise distance metric between particle  $i$  and  $j$ ,  $d_{ij}$  or  $d_{ji}$ , and the beam distance,  $d_{iB}$ , are defined as

$$d_{ij} = d_{ji} = \min(p_{Ti}^2, p_{Tj}^2) \frac{\Delta R_{ij}^2}{R^2}, \quad (4.2)$$

$$d_{iB} = p_{Ti}^2, \quad (4.3)$$

where  $p_{Ti}$  is the transverse momentum of particle  $i$  with respect to the beam ( $z$ ) direction, and  $\Delta R_{ij}^2$  is the measurement of angular separation of particle  $i$  and  $j$ , defined as  $\Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$ , where  $y_i = \frac{1}{2} \ln \frac{E_i + p_{zi}}{E_i - p_{zi}}$  and  $\phi_i$  are the rapidity and azimuthal angle of particle  $i$  respectively. The free parameter  $R$  controls the jet radius.

The distance metric and the beam distance are calculated for all pairs of particles. If the minimum value of all distance metrics and beam distances is a distance metric,  $d_{ij}$ , particle  $i$  and  $j$  are merged and the four momentum of particle  $i$  is updated as the sum of the two particles. If the minimum value is a beam distance,  $d_{iB}$ , particle  $i$  becomes an output jet and is removed from the list of particles. The above procedure is repeated until no particles are left.

The exclusive variant is similar to the inclusive variant. First difference is that when the minimum value of all distance metrics and beam distances is a beam distance,  $d_{iB}$ , particle  $i$  forms part of the beam jet. The beam jet is discarded at the end of the jet clustering. The second difference is that when all distance metrics and beam distances are above a threshold,  $d_{cut}$ , the jet clustering will stop.

The exclusive mode allows a specified number of jets to be found, where the  $d_{cut}$  is automatically determined. In contrast, the inclusive mode would find as many jets as the algorithm allows.

#### 4.4.2.2 Durham algorithm

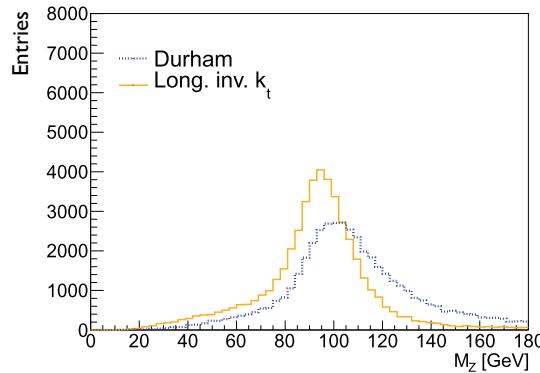
The Durham algorithm [75], also known as  $e^+e^- k_t$  algorithm, is commonly used in  $e^+e^-$  collider experiments. It has one pair-wise distance metric:

$$d_{ij} = 2 \min(E_i^2, E_j^2) (1 - \cos(\theta_{ij})), \quad (4.4)$$

where  $E_i$  is the energy of particle  $i$  and  $\theta_{ij}$  is the angle between particle  $i$  and  $j$ . The Durham algorithm can only be run in exclusive mode, which means that the clustering will stop when all distance metrics are above a threshold,  $d_{cut}$ .

#### 4.4.2.3 Jet algorithms for CLIC

Although CLIC is an  $e^+e^-$  collider, the beam induced background is significant and deposits a large amount of energy in the detector. Therefore, traditional  $e^+e^-$  jet algorithms, like the Durham algorithm, are not suitable for the CLIC colliding environment. Figure 4.4 shows the reconstructed Z boson mass distribution for  $ZZ \rightarrow q\bar{q}q'\bar{q}'$  at a 500 GeV CLIC [76], using the Durham jet algorithm and the longitudinally invariant  $k_t$  jet algorithm. The  $\gamma\gamma \rightarrow$  hadrons background corresponding to 300 bunch crossings was overlaid on the event, where each bunch crossing contains approximately 0.3  $\gamma\gamma \rightarrow$  hadrons events. Longitudinally invariant  $k_t$  algorithm gives a better reconstructed Z boson mass distribution. Other studies [2, 77] have also shown that jet algorithms for the pp colliders, such as longitudinally invariant  $k_t$  algorithm, give better reconstructed mass and energies resolutions at CLIC. Therefore, longitudinally invariant  $k_t$  algorithm is often used in analyses at CLIC.



**Figure 4.4:** The reconstructed Z boson mass distribution for  $ZZ \rightarrow q\bar{q}q'\bar{q}'$  at a 500 GeV CLIC, using the Durham jet algorithm and the longitudinally invariant  $k_t$  jet algorithm. The  $\gamma\gamma \rightarrow$  hadrons background corresponding to 300 bunch crossings was overlaid on the event, where each bunch crossing contains approximately 0.3  $\gamma\gamma \rightarrow$  hadrons events. The figure is adapted from [76].

#### 4.4.2.4 The $y$ parameter

The  $y$  parameter is a measure of the number of jets in an event. It describes the transition of going from  $N$  clustered jets to  $N+1$  clustered jets using an exclusive jet algorithm. For example,  $y_{23}$  would be the  $d_{cut}$  value for an exclusive jet algorithm, above which the jet algorithm returns 2 jets, below which the jet algorithm returns 3 jets. Numerically the  $y$  parameter is often much smaller than 1. A typical way to convert a small number to a machine acceptable range is to take negative logarithm of the number.

## 4.5 Multivariate analysis

Multivariate analysis (MVA) has become increasingly important in high energy physics. The implementation of the machine learning based MVA used in this thesis is provided by TMVA package [78]. MVA can be used for classification or regression. Classification classifies an event into one of several classes. Regression of an event gives an output in a continuous numerical range. In a typical physics analysis, MVA is often used to classify one type of event from other types.

A typical MVA classification involves two classes, sometimes referring to as the signal class and the background class. Before using the MVA classification, a machine learning model (classifier) needs to be trained with training data. The model uses a set of discriminant variables as inputs, which have different distributions for the signal and the background. To use the MVA classification, the trained model will be applied onto the testing data. The classification response is to classify a testing event into the signal class or the background class.

This two-class classification scheme can be easily extended to multiple classes, as implemented in TMVA with the `MULTICLASS` class. For example, The `MULTICLASS` approach is used in the tau decay mode classification in section 6.5 and in the flavour tagging classifier in section 8.5.

### 4.5.1 Optimisation and overfitting

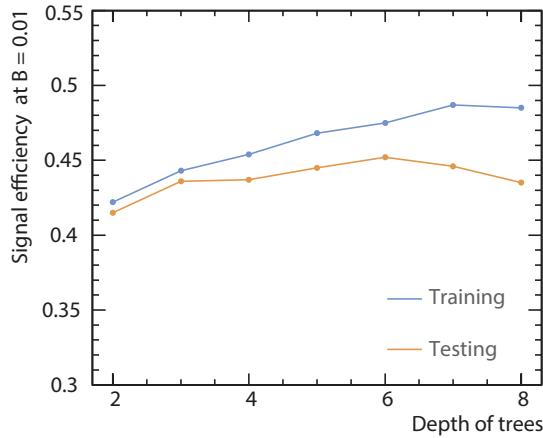
Two important concepts with MVAs are optimisation and overfitting of a model. The optimisation of a model refers to selecting the optimal free parameters of the model. One could build a complex model which fits the training samples extremely well, but the model may not be optimal for a testing sample. A simple model is less prone to statistical fluctuations of samples. However, the model might be too simple to achieve the optimal modelling. The former case is known as overfitting, or overtraining. The latter case is called underfitting, or undertraining.

An optimal model lies between an overfitted model and an underfitted model. In practice, this involves building a model with increasing complexities, and identifying the point where overfitting occurs.

One definition of overfitting is when the efficiency of the signal selection in the training samples increases, but the efficiency of the signal selection in the testing sample decreases, with the increase of the model complexity. Figure 4.5 shows the signal selection efficiency as a function of the model complexity, using an example from the double Higgs analysis of chapter 8 and the Boosted Decision Tree model. The efficiency of the signal selection is defined as the fraction of the signal selected when the fraction of background is 1%, reported by the TMVA training process. The depth of the tree reflects the complexity of the model. From a tree depth of two to six, the efficiency for both testing and training samples increases. From a tree depth of six onwards, overfitting occurs. In this particular example, one should choose a tree depth fewer than seven to avoid overfitting.

### 4.5.2 Choice of models

The model to fit the data can be as simple as a cut-based model, a likelihood estimator, or a linear regression model. The model can also be as complicated as a non-linear tree, a non-linear neural network, or a support vector machine. Regardless of the model complexity, the choice of the most optimal classifier is often data driven to match the nature of the sample. For example, a non-linear model is the best to model a non-linear response to the input variables.



**Figure 4.5:** Example of the signal selection efficiency as a function of the model complexity. The example is chosen from the double Higgs analysis at  $\sqrt{s} = 3 \text{ TeV}$ , using the Boosted Decision Tree model. The efficiency of the signal selection is defined as the fraction of the signal selected when the fraction of background is 1%, reported by the TMVA training process. The depth of the tree reflects the complexity of the model. From a tree depth of six onwards, overfitting occurs.

To rigorously identify the best model, individual optimisations of models are required, which is computationally very expensive. However, as researchers in the machine learning field suggested, the boosted decision tree is probably the best out-of-the-box machine learning model [79]. A neural network model could potentially perform better than the boosted decision tree model, but it requires more tuning, and is less intuitive to interpret. For these reasons, the boost decision tree model (BDT) is chosen for physics analyses in this thesis. Before describing the BDT model in detail, some simpler models are discussed.

### 4.5.3 Rectangular Cut model

The rectangular cut method optimises cuts to maximise some pre-defined metrics. The metric could be the signal efficiency for a given background efficiency. Alternatively, the metric can be the significance,  $\frac{S}{\sqrt{S+B}}$ , where  $S$  and  $B$  are respective numbers of signal and background events passing the rectangular cuts.

Discriminant variables give better separation power when they are gaussian-like and statistically independent [79]. Therefore it is preferable to decorrelate the variables and gaussian transform them before using the rectangular cut MVA. One way to decorrelate variables is to find linear combinations from the variables, such that the correlation between the linear combinations is less than the correlation between the variables.

#### 4.5.4 Projective Likelihood model

The projective likelihood model with probability density estimators (PDE) is used in PandoraPFA for the photon ID test, due to its simplicity and low requirement on computing resources. The PandoraPFA implementation of the projective likelihood model is discussed in section 5.3.3.

The likelihood classifier calculates the probability density for each discriminant variable, for the signal class and the background class. The overall signal and background likelihood are defined as products of the individual probability densities of each variable, for the respective signal class and background class. The likelihood ratio,  $R$ , can be then defined as the signal likelihood divided by the sum of the signal likelihood and the background likelihood.

Discriminant variables give better separation power when they are gaussian-like and statistically independent [79]. Hence similar to the rectangular cut method, the likelihood model works better with decorrelated, gaussian-like variables [79].

#### 4.5.5 Decision Tree model

The decision tree is a non-linear tree based model. Its rather complex nature requires a careful explanation of many concepts.

The decision tree is a binary tree, where each splitting node (splitting point) uses a cut on a single discriminant variable to decide whether an event is signal-like (“goes down by a layer to the left”), or background-like (“goes down by a layer to the right”), depending on whether the event passes the cut. At each splitting node, samples are divided into two sub-samples: signal-like and background-like sub-samples. This splitting process (tree growing) starts at the root node. For each sub-sample, the splitting process stops after certain criteria are met. The stopping criteria could be the minimum number of events in a node, the maximum number of layers of the tree, or a minimum/maximum signal purity of the end nodes. The end nodes, where the tree stops growing and the sub-samples are not split, contains signal and/or background events. If there are more signal than background events in an end node, it is called a signal-like end node. The opposite is called a background-like end node.

The training of the decision tree refers to finding the optimal cut at each splitting node by minimising a given metric. Assuming the probability of the cut producing the signal is  $p$ , three commonly used metrics for two-class classification are:

1. misclassification error:  $1 - \max(p, 1-p)$ ;
2. Gini index:  $2p(1-p)$ ;
3. cross-entropy or deviance:  $-p \log p - (1-p) \log (1-p)$ .

The application of a trained decision tree is performed by traversing the tree from the root node to the end node. The event is classified as signal or background, depending on whether it falls in a signal-like or background-like end node.

Figure 4.6 illustrates a simple example of a trained decision tree. PhD student is the signal class and undergraduate student is the background class. The top diamond box in figure 4.6, “Leave party before 1am”, is the root node. All diamond boxes are splitting nodes. Rectangular boxes are end nodes. The signal-like end node is represented by the red rectangle and the background-like end nodes are represented by blue rectangles. The depth of this decision tree is two. The metric used to find the optimal cut at the splitting node is the Gini index.

The attributes of signal and background classes are listed table 4.4. In this example, there are ten PhD students who leave parties before 1 am and know where a free pizza is located. In contrast, five undergraduate students leave parties after 1 am and know where a free pizza is located. Three undergraduate students leave parties before 1 am and do not know where a free pizza is located. Two undergraduate students leave parties after 1 am and do not know where a free pizza is located.

PhD students	Leave party before 1 am	Leave party after 1 am
Know where free pizza is	10	0
Not know where free pizza is	0	0
Undergraduates	Leave party before 1 am	Leave party after 1 am
Know where free pizza is	0	5
Not know where free pizza is	3	2

**Table 4.4:** The attributes of the PhD students and undergraduate students for the decision tree example shown in figure 4.6.

To demonstrate the first step of the training of the model, details of finding the optimal cut at the root node are outlined. There are two possible cuts for the root node, “Leave party before (after) 1am” and “(Not) Know where free pizza is”. If the cut at the root node is “Leave party before 1am”, the probability of the cut producing the signal,  $p$ , is  $\frac{10}{13}$ , as there are 10 PhD students and 3 undergraduate students who leave parties

before 1 am. The Gini index gives

$$2p(1-p) \simeq 0.36. \quad (4.5)$$

If the cut at the root node is “Know where free pizza is”,  $p = \frac{10}{15}$ , as there are 10 PhD students and 5 undergraduate students who know where a free pizza is located. The Gini index gives

$$2p(1-p) \simeq 0.44. \quad (4.6)$$

Therefore, by choosing the cut that minimises the Gini Index, the optimal cut for the root node is “Leave party before 1am”.

The simple tree in figure 4.6 is grown fully as each end node contains signal or background only. An example of applying the trained decision tree is provided: if there is a student who leaves parties before 1 am and knows where a free pizza is located, then the student is classified as a PhD student.

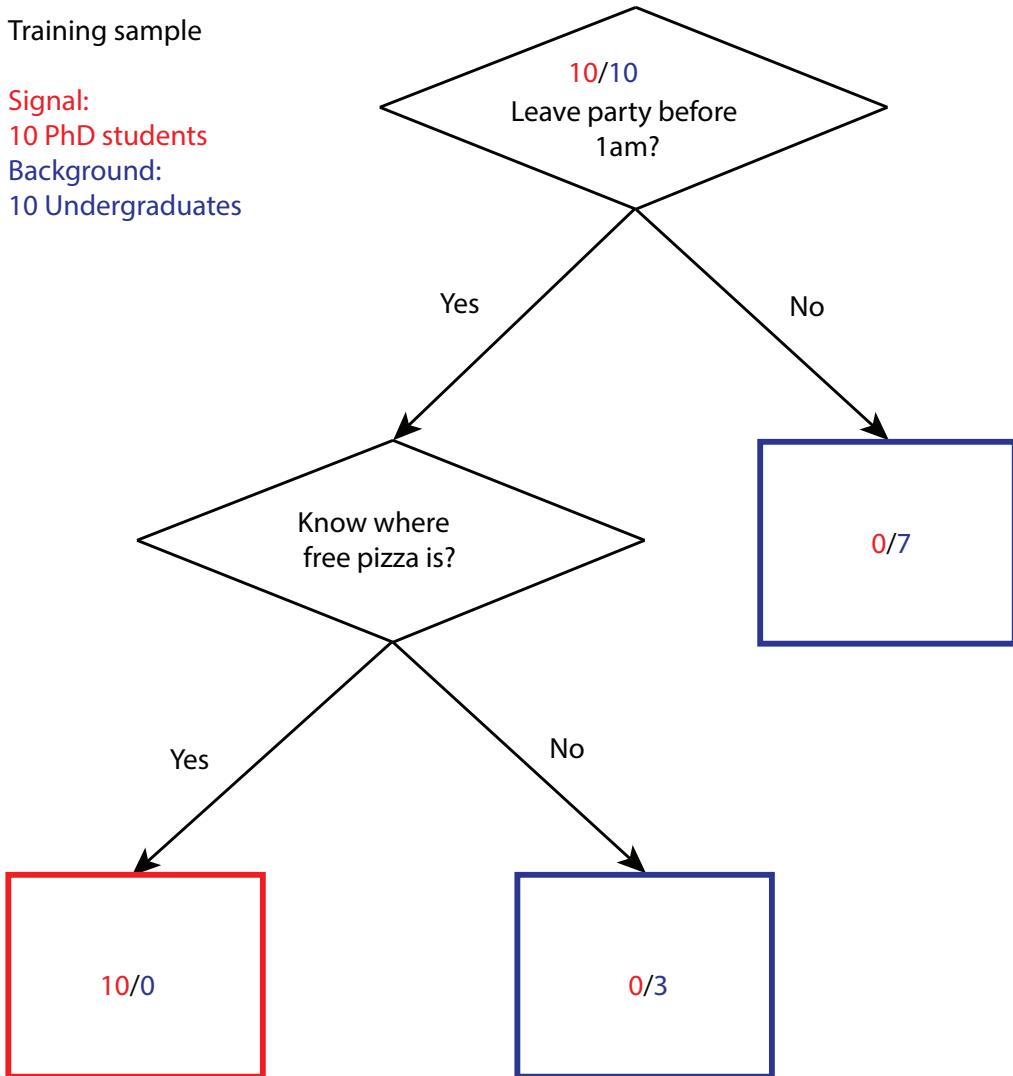
#### 4.5.5.1 Improving decision tree

It is very easy to construct a decision tree that fits the training data very well, but the tree would not be optimal for the testing sample (because of overfitting). To overcome the instability of the decision tree model, many methods have been developed. Some of the most successful ones are boosting, bagging, and random forest.

- Boosting: the basic idea of boosting is that the tree growing procedure focuses on events which are difficult to classify correctly. By assigning a weight to each event, after each tree growing iteration, the weights for misclassified events are gradually increased. Therefore misclassified events get more attention in the next iteration.
- Bagging: also known as boot-strap, it is a method that selects a random subset of the training sample, and uses the subset in the training stage.
- Random Forest: when a tree is grown, a randomly selected subset of discriminant variables are used to grow the tree.

#### 4.5.6 Boosted Decision Tree model

A Boosted Decision Tree (BDT) contains many decision trees, where the boosting is used to grow trees. There are two common boosting methods: adaptive boosting and gradient



**Figure 4.6:** Example of a decision tree. Numbers in each node represent the number of PhD students (red) and the number of undergraduate students (blue). Diamond boxes represent splitting nodes. Rectangular boxes represent end nodes. Blue boxes are background-like end nodes. The red box is a signal-like end-node.

boosting. Adaptive boosting, first introduced in [80], is discussed in further detail, as it is simpler to understand than the gradient boosting. The adaptive boosting algorithm, adapted from [79], is outlined below:

- At the initialisation stage, the event weight is initialised to  $w = 1/N$  for every event, for  $N$  total events.
- Iterate  $M$  times, where  $M$  is the total number of trees. For iteration  $m$ :
  - Create/grow a  $m^{th}$  tree with weighted samples obtained from  $(m-1)^{th}$  iteration.

- Update the  $m^{th}$  tree error function,  $err_m = \frac{\sum_{i=1}^N w_{i,m-1} B_{i,m}}{\sum_{i=1}^N w_{i,m-1}}$ .
- Update the  $m^{th}$  tree weight,  $\alpha_m = \log\left(\frac{1-err_m}{err_m}\right)$
- Update the  $i^{th}$  event weight in  $m^{th}$  tree,  $w_{i,m} = w_{i,m-1} e^{\alpha_m B_{i,m}}$ .
- The output,  $G(x)$ , for a testing event  $x$ , is a weighted vote from all  $M$  trees:

$$G(x) = \begin{cases} -1, & \text{if } \sum_{m=1}^M \alpha_m G_m(x) < 0, \\ 1, & \text{otherwise.} \end{cases} \quad (4.7)$$

The tree classifier output,  $G$ , is denoted as  $-1$  or  $+1$ . One can think of  $-1$  as background and  $+1$  as signal. There are  $N$  events and  $M$  iterations (trees). The parameter  $B$  represents whether an event is misclassified. For the  $i^{th}$  event in the  $m^{th}$  tree,  $B_{i,m} = 1$  if the event is misclassified, or  $0$  if the event is correctly classified. The parameter  $w_{i,m}$  represents the event weight for  $i^{th}$  event in  $m^{th}$  tree. In each iteration, if  $i^{th}$  event is misclassified in  $m^{th}$  tree, the event weight increases by a factor of  $(1 - err_m)/(err_m)$ . Otherwise, the event weight does not change.

The power of adaptive boosting is to dramatically improve the performance of a weak classifier. A weak classifier is a classifier which gives a predictive performance only slightly better than a random guess. A decision tree with one or two layers would be a weak classifier. By sequentially applying many weak classifiers with weighted samples, the final “forest” is very robust with a very good performance at selecting signals.

#### 4.5.6.1 Optimisation of Boosted Decision Tree

The most important parameter is the depth of a tree, which determines how many end nodes the tree has. It also affects the complexity of the BDT model. If the depth of a tree is set to a large value, it could lead to the overfitting of the model.

The number of trees is another important parameter. Past studies on BDTs show that using many small trees yields the best result [79]. It has been shown that a large number of trees does not lead to overfitting [79]. Therefore there is a debate on the metric to determine the optimal number of trees. The minimum number of events in a node, which is a stopping criterion for tree growing, affects the size of the tree, but is less important than the depth of the tree parameter.

The boosting algorithm has two variants in the TMVA implementation: adaptive boost and gradient boost. The learning rate of the adaptive boost controls how fast the event weight changes in each boosting iteration. Studies on BDT show that a small learning rate ( $\sim 0.1$ ) with many trees works better than a large learning rate with fewer trees [79].

The shrinkage rate in the gradient boost is similar to the learning rate parameter in the adaptive boost. The shrinkage rate controls how fast the weight changes for events in each boosting iteration. Again a small value ( $\sim 0.1$ ) is preferable [79].

The usual choice of the metric to optimise cuts for tree growing is either the Gini index or the cross-entropy. The two metrics make little differences to the performance.

The number of bins per variable is a necessary parameter to make tree growing efficient, because it is faster to compute the optimal cut at splitting nodes for discretely binned variables than continuous variables. This parameter, however, has little impact on the model performance. But because variables are binned, variables should be pre-processed before feeding into the training model. For example, a variable should be limited to a range to avoid the extreme values that distort the shape of the variable distribution. If the original distribution of a variable is highly skewed, the variable should be transformed to obtain a more uniform distribution.

For the end node, the output can either be signal-like or background-like, based on the majority of the training events in the end node. Numerically, it can correspond to 1/0. However, the end node could also use signal purity as the output, resulting in a continues range of [0,1].

The bagging fraction determines the fraction of randomly selected events used in each boosting iteration. By choosing a small value, events between each boosting iteration are less correlated. Hence the overall model performance improves.

The `DO_PRESLECTION` flag in the BDT of TMVA allows the classifier to identify phase spaces where there are only background events and apply cuts to discard these background events.

#### 4.5.7 Multiple classes

The above discussion assumes exactly two classes, the signal class and the background class. The classification can be extended to multiple classes. There are two ways to train

a classifier for multiple classes. The “one versus one” scheme trains each class against each other class. The second way is called “one versus all”, when each class is trained against all other classes combined.

Using a three-class example, A, B, and C, the “one versus one” scheme trains class A against class B; class B against class C; and class C against class A. “One versus all” scheme would train class A against non-A classes; class B against non-B classes; and class C against non-C classes.

The TMVA implementation of the multiple classes classifier, **MULTICLASS**, uses the “one versus all” scheme. For each class, the **MULTICLASS** classifier will train the class against all other classes. This process is repeated for each class, resulting in multiple classifiers. The overall classifier output for a single event is a normalised response using all trained classifiers, where the sum of the classifier outputs for a single event is one. The individual response of a trained classifier for an event can be treated as the likelihood. In the application stage, the event is classified into a class if the classifier for that class gives the highest output response amongst all classifiers for that event.

The advantage of using the **MULTICLASS** classifier instead of a two-class classifier for samples with multiple classes is that the classifier outputs are correctly adjusted for multiple classes. Hence one event can be unambiguously classified into only one class. The issue with the **MULTICLASS** classifier is that powerful discriminant variables for each individual class need to enter the training stage simultaneously, resulting in a large number of discriminant variables in the **MULTICLASS** classifier.

# Chapter 5

## Photon Reconstruction in PandoraPFA

*'I dreamed I was a butterfly, flitting around in the sky; then I awoke.  
Now I wonder: Am I a man who dreamt of being a butterfly, or am I a  
butterfly dreaming that I am a man?'*

— Zhuang Zi, 369 BC – 286 BC

A good single photon energy resolution and the ability to reconstruct two spatially close photons are necessary to reconstruct particles using decay processes involving photons, such as  $\pi^0 \rightarrow \gamma\gamma$  decays. Furthermore, the ability to correctly reconstruct photons in a dense jet environment improves the charged particle reconstruction by removing the calorimeter hits belonged to the photons and simplifying the pattern recognition problem for the charged particle reconstruction.

The photon reconstruction algorithms presented in this chapter have benefited a number of physics analyses. The most recent example is the  $H \rightarrow \gamma\gamma$  simulation study at CLIC [81].

This chapter starts with an overview of the electromagnetic shower produced by photons passing through a thick absorber. It then discusses photon reconstruction algorithms within the PandoraPFA framework, followed by a description of the performance of these algorithms. Part of this chapter has been published in the proceedings of 2015 International Workshop on Future Linear Colliders [82].

## 5.1 Electromagnetic showers

Electromagnetic (EM) showers develop through processes of the pair production and bremsstrahlung when a high energy photon or electron passing though a thick absorber. Many low-energy photons and electrons are generated, producing shower-like structures in the detector. Two suitable length scales to describe the EM shower growth are the radiation length and the Molière radius [83, 84].

The radiation length of a material describes the EM longitudinal shower profile, defined as the mean distance travelled by an electron for its energy to be reduced by a factor of  $1/e$  via bremsstrahlung. It is also defined as the  $7/9$  of the mean free path for pair production by a high energy photon [85].

Figure 5.1 shows the simulated longitudinal electromagnetic shower profiles as a function of the longitudinal shower depth for electrons and photons. The mean EM longitudinal shower profile can be described by the following function [86]:

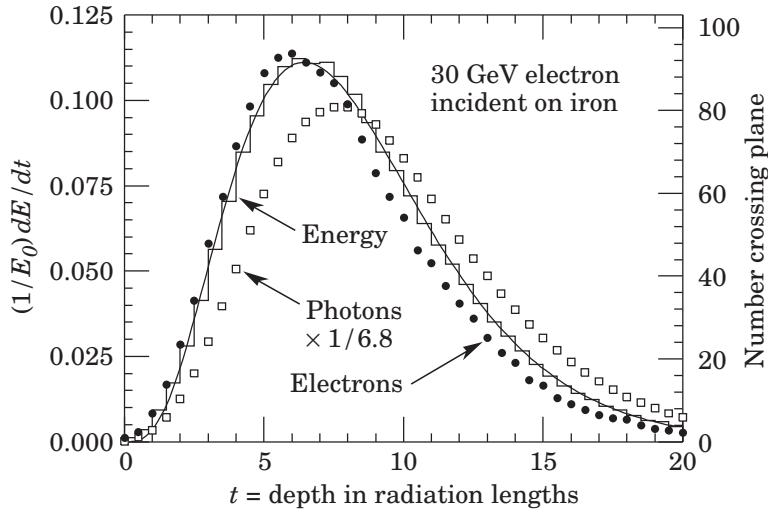
$$\frac{dE}{dt} = E_0 b \frac{(bt)^{a-1} e^{-bt}}{\Gamma(a)}, \quad (5.1)$$

where  $t$  is the number of radiation lengths; the parameter  $E_0$  is the initial energy of the photon/electron; the parameter  $b$  takes the value of 0.5, which is sufficient for the purpose of photon reconstruction [3]; and the parameter  $a$  is given by [36]:

$$a = 1.25 + 0.5 \ln \left( \frac{E_0}{E_c} \right), \quad (5.2)$$

where  $E_c$  is the critical energy. The critical energy is defined as the energy of the electron at which the rate of losing energy by bremsstrahlung is the same as the rate of losing energy by ionisation [87]. The alternative definition of the critical energy is the energy at which the energy loss by ionisation per radiation length is the same as the particle energy [88]. This parametrisation of the EM longitudinal shower profile should only be used to describe an average behaviour of the EM shower; fluctuations in individual EM shower profiles are significant.

The EM transverse shower profile can be described as a narrow core, widening as the shower develops. 90% of the shower energy is contained in a cylinder with a radius of one Molière radius. About 99% of the shower energy is contained inside of 3.5 Molière radius [83, 84].



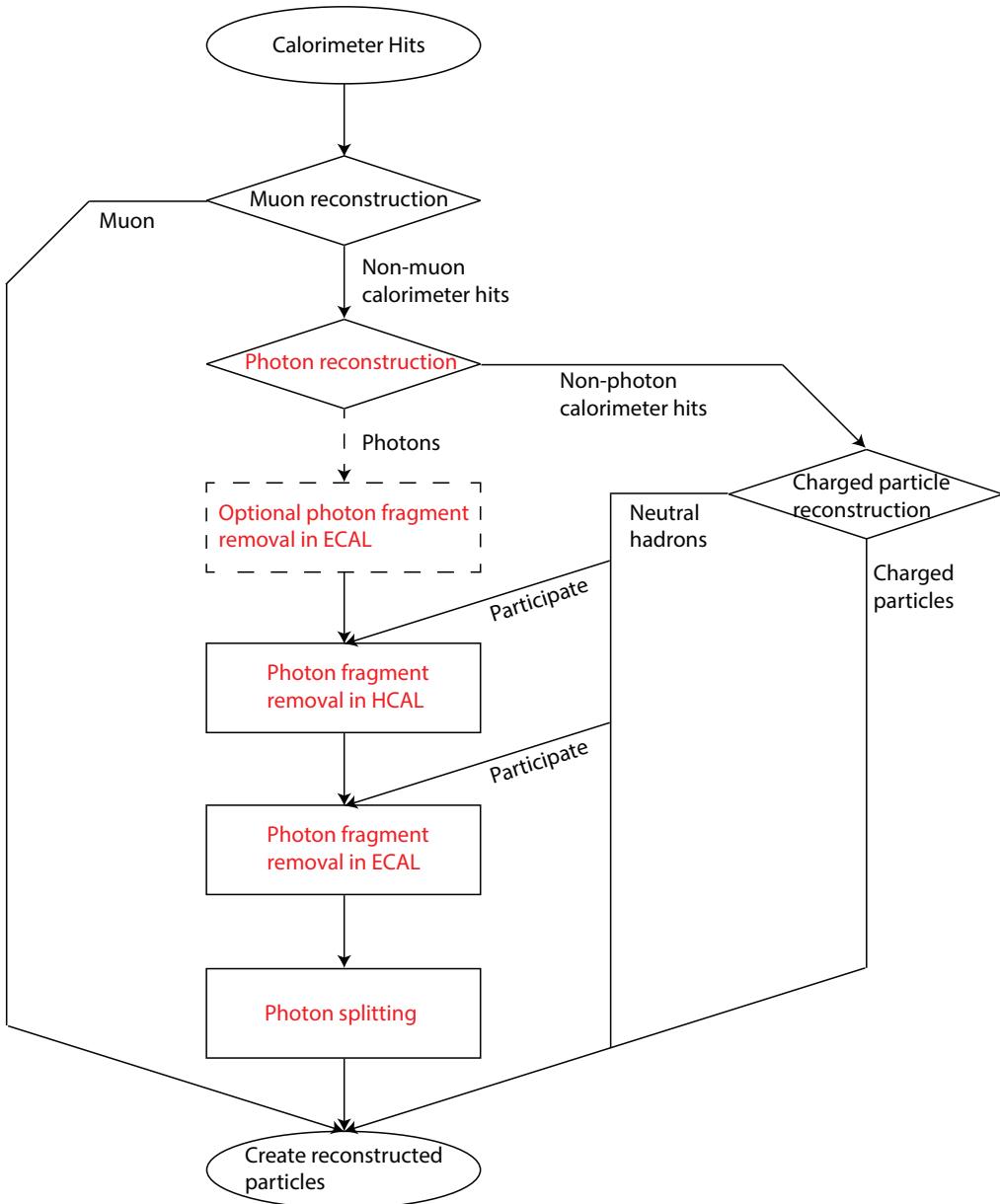
**Figure 5.1:** An EGS4 [89, 90] simulation of a 30 GeV electron-induced electromagnetic shower in iron. The histogram shows the fractional energy deposition as a function of radiation length, and the curve is a gamma-function fit to the distribution. Circles and squares are the numbers of electrons and photons respectively with total energy greater than 1.5 MeV crossing planes with scale on right. Plot is taken from [3].

## 5.2 Overview of photon reconstruction in PandoraPFA

Five algorithms are developed to tackle different issues in photon reconstruction in PandoraPFA:

- PHOTON RECONSTRUCTION algorithm reconstructs photons from calorimeter hits in the ECAL, including forming a photon candidate and applying a photon ID test, with special treatments for photons close to charged particles.
- Two photon fragment removal algorithms remove fragments in the ECAL.
- One algorithm removes fragments in the HCAL.
- A photon splitting algorithm separates accidentally merged photons.

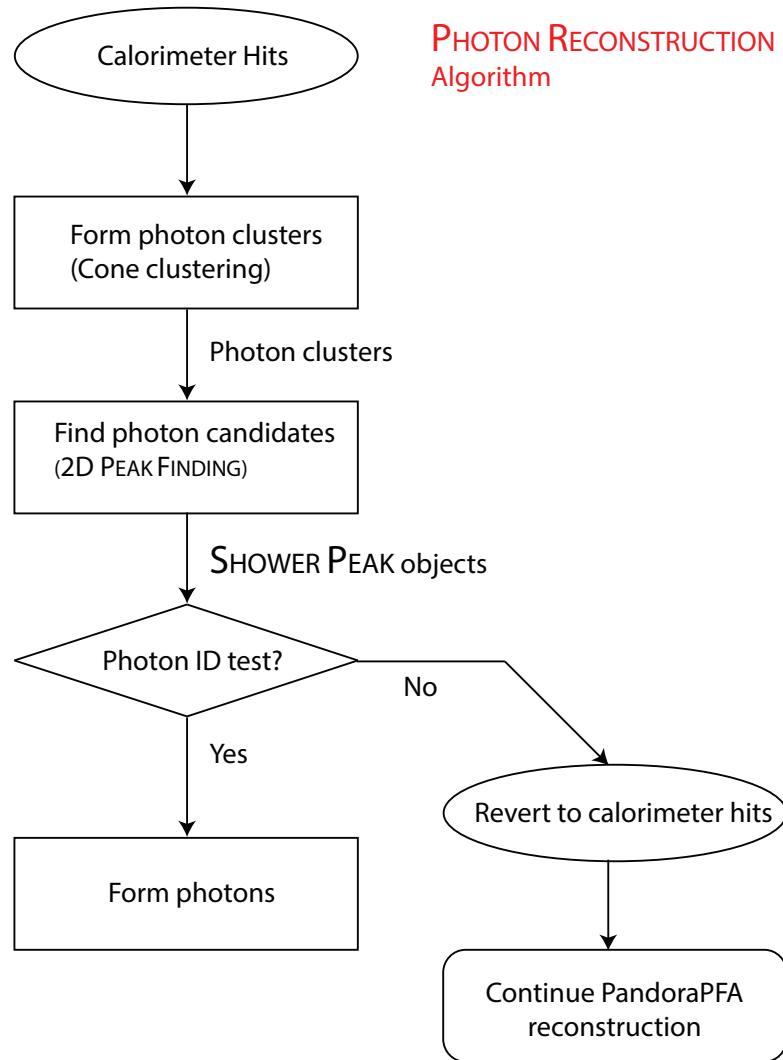
Positions of the photon algorithms used in the PandoraPFA are shown in figure 5.2. The five photon algorithms are highlighted in red. Samples used to produce figures in this chapter were simulated with the ILD detector model.



**Figure 5.2:** Positions of the photon algorithms used in the PandoraPFA. The five photon algorithms are highlighted in red.

### 5.3 Photon Reconstruction algorithm

The PHOTON RECONSTRUCTION algorithm runs at an early stage of the overall reconstruction, before the charged particle reconstruction. Main steps of the PHOTON RECONSTRUCTION algorithm, shown in figure 5.3, are: forming photon clusters; finding photon candidates; and photon ID test.



**Figure 5.3:** Main steps of the PHOTON RECONSTRUCTION algorithm: forming photon clusters; finding photon candidates; and photon ID test.

### 5.3.1 Forming photon clusters

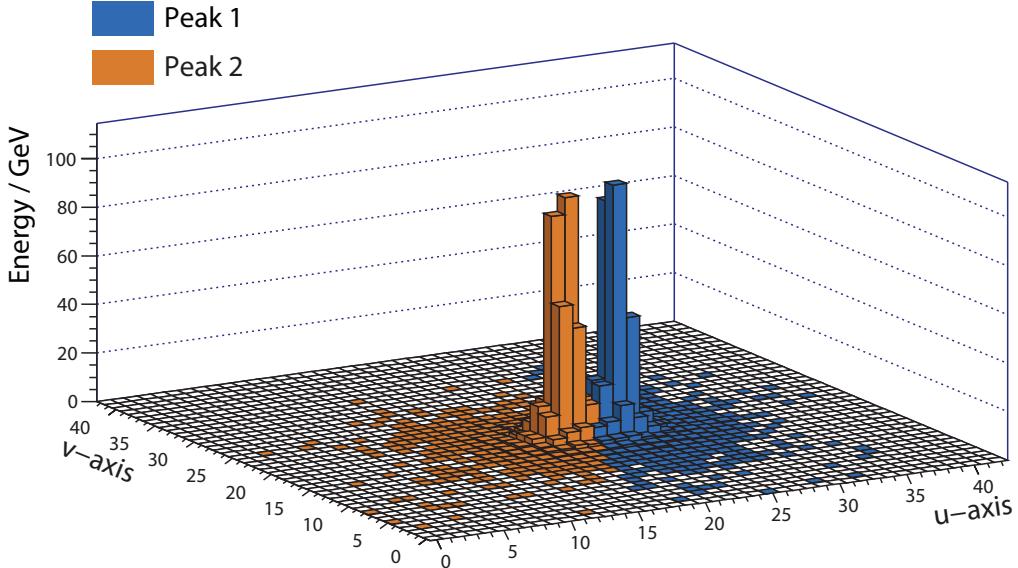
Since the PHOTON RECONSTRUCTION algorithm runs after the muon reconstruction algorithm, inputs of the PHOTON RECONSTRUCTION algorithm are calorimeter hits in the ECAL that are not associated with reconstructed muons.

This step forms clusters from calorimeter hits in the ECAL using the cone clustering algorithm. As the target for reconstruction is the neutral photon, the cone clustering algorithm uses high-energy calorimeter hits in the ECAL as initial seeds in the order of descending energies, instead of using track projections as initial seeds. Parameters for large cones are used to form clusters such that it is unlikely that one photon is split into two clusters, but one cluster may contain calorimeter hits from multiple photons.

### 5.3.2 Finding photon candidates and 2D Peak Finding algorithm

If a cluster contains calorimeter hits from several photons, this step aims to split the three-dimensional cluster into several smaller clusters (photon candidates). Each photon candidate should contain calorimeter hits from one photon only.

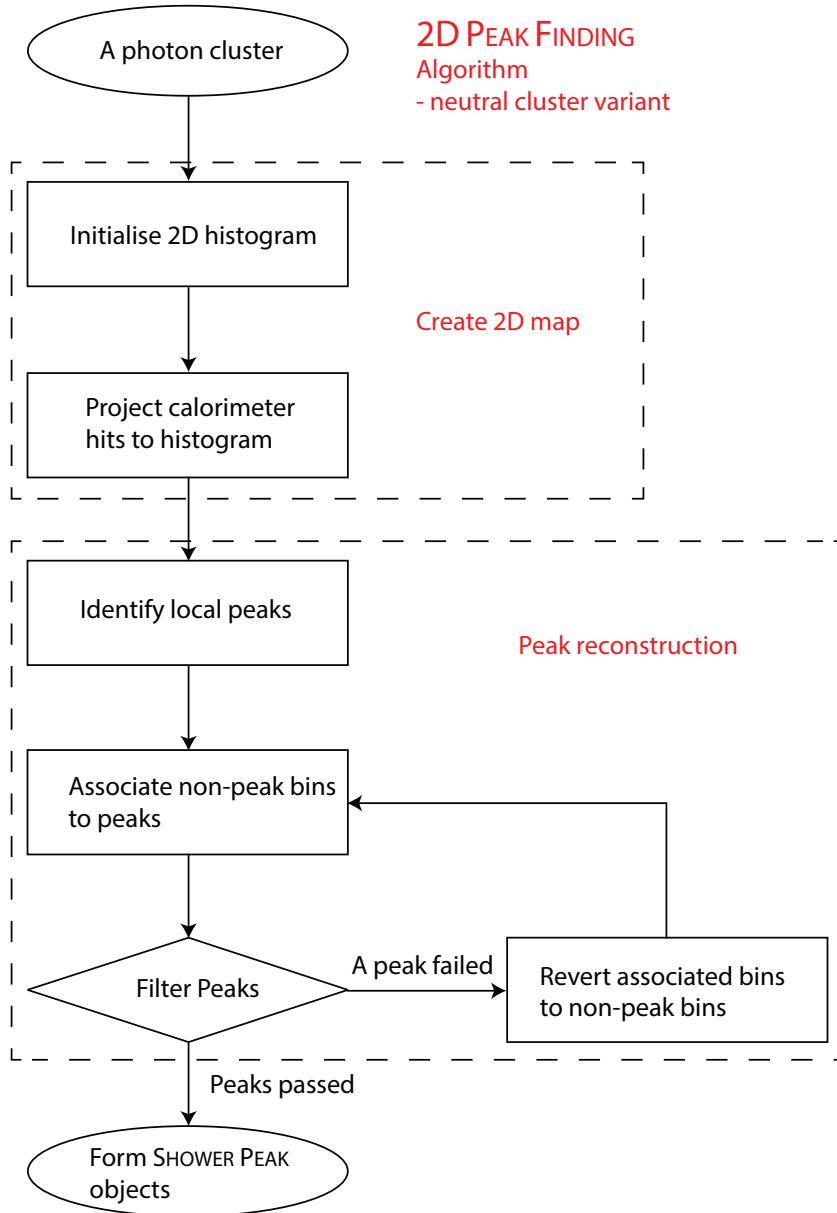
The three-dimensional splitting problem is harder than a two-dimensional one. Therefore, a translation is needed to map the three-dimensional problem to a more manageable two-dimensional problem. This translation relies on the characteristic EM transverse shower profile. When the energies of the calorimeter hits of the cluster are projected onto a two-dimensional plane, an EM shower core would appear as a mountain-like structure in the plane. Figure 5.4 shows an example of a cluster projected onto a two-dimensional plane. By identifying a peak in the two-dimensional plane, an EM shower core is identified.



**Figure 5.4:** Two 500 GeV photons (yellow and blue) within a cluster, just resolved in a transverse plane orthogonal to the direction of the cluster. The axes U and V are orthogonal axes in units of the ECAL cell sizes. The height of a bin in the histogram is the sum of the calorimeter hit energy associated with the bin.

Identifying photon candidates inside a cluster is translated into identifying peaks in a two-dimensional plane, using a two-dimensional peak-finding algorithm (2D PEAK FINDING algorithm). The 2D PEAK FINDING algorithm aims to correctly identify peak positions in a two-dimensional histogram and to associate non-peak bins to identified peaks.

There are two variants of the 2D PEAK FINDING algorithm: the neutral cluster variant and the charged cluster variant. The main steps of the neutral cluster variant are shown in figure 5.5: creating the 2D map; peak reconstruction; and forming SHOWER PEAK objects.



**Figure 5.5:** Main steps of the neutral cluster variant of the 2D PEAK FINDING algorithm: creating the 2D map; peak reconstruction; and forming SHOWER PEAK objects.

### 5.3.2.1 Creating the 2D map

A two-dimensional (2D) histogram is initialised to host the projection of the calorimeter hits of the cluster. For the best resolving power of EM showers, the projection direction

is chosen to be the direction of the cluster. Two axes of the two-dimensional histogram are chosen such that the axes and the direction of the cluster form an orthogonal basis in the three-dimensional space.

The calorimeter hits associated with the cluster are then projected onto the 2D histogram. The projection is chosen such that the cluster centroid position is projected onto the centre of the histogram. The distance between the calorimeter hit position and the cluster centroid position is then converted into a distance vector to be used to project the calorimeter hit. The distance vector,  $\vec{s}_i$ , of a calorimeter hit  $i$ , is defined as:

$$\vec{s}_i = \frac{\vec{a}_i - \langle \vec{a} \rangle}{d_{cell}}, \quad (5.3)$$

where  $\vec{a}$  is the three-dimensional position of the calorimeter hit  $i$ ;  $\langle \vec{a} \rangle$  is the centroid position of cluster  $a$ ; and  $d_{cell}$  is the ECAL square cell length. The coordinate of the calorimeter hit projection onto the histogram is calculated from the scalar products of the distance vector ( $\vec{s}_i$ ) with the axes vectors.

The height of a bin in the 2D histogram is the sum of the energies associated with the calorimeter hits that fall in that particular bin. Each bin contains calorimeter hits that projected onto the bin. One bin size along either axes on the 2D histogram corresponds to one ECAL square cell length.

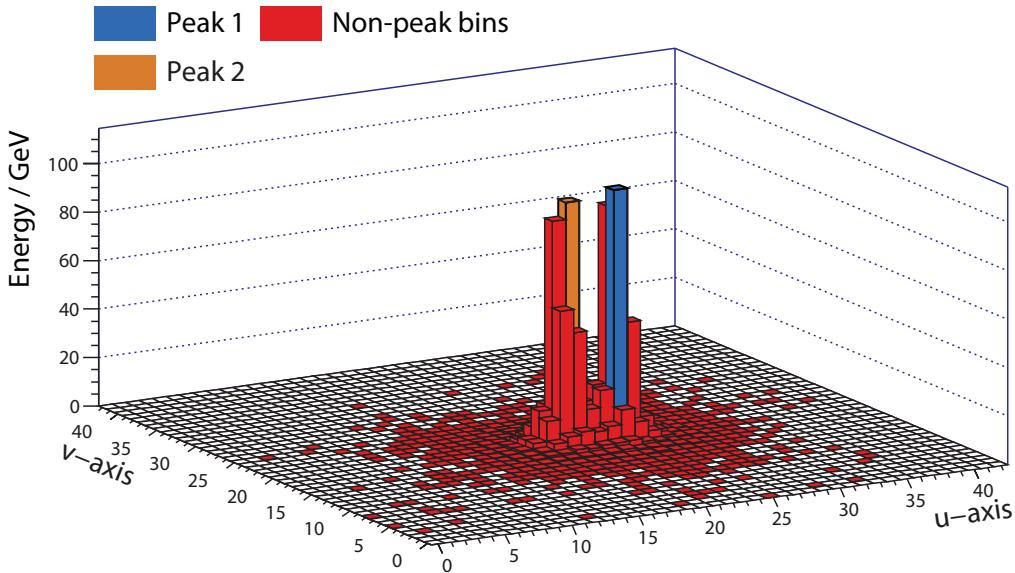
### 5.3.2.2 Peak reconstruction

All local peaks are identified in the 2D histogram. A local peak is defined as a bin where its height is above all eight neighbouring bins. All bins in the 2D histogram are considered to identify all local peaks. Figure 5.6 shows an example of a 2D histogram with two local peak bins (orange and blue) identified. Red bins are non-peak bins.

Having identified all local peaks, non-peak bins are associated to a particular peak based on the energy of the peak and the distance of the non-peak bin to the peak bin. A non-peak bin should be associated to a high-energy peak bin that is close to the non-peak bin.

A non-peak bin is associated with the peak bin that gives the smallest value of the metric:

$$\frac{d_i}{\sqrt{E_i}} \quad (5.4)$$



**Figure 5.6:** Two peak bins, orange and blue, are identified. Red bins are non-peak bins.

where  $d_i$  is the Euclidean distance between a non-peak bin and a peak bin  $i$  in the 2D histogram, and  $E_i$  is the height (energy) of the peak bin  $i$ . For each non-peak bin, the metric is considered for all peak bins to find the peak bin that produces the smallest metric. In the example of the 2D histogram in figure 5.6 with two peaks identified, the result after associating non-peak bins to peak bins is shown in figure 5.4.

The performance of the 2D PEAK FINDING algorithm is improved by peak filtering. In a 2D histogram, major peaks with many associated non-peak bins most likely correspond to physical photons, while minor peaks with a few associated non-peak bins are more likely from fluctuations in the energy deposition of the EM shower. To select only major peaks and to discard minor peaks, every time after all non-peak bins are associated with peak bins, peaks with fewer than three bins associated (including the peak bin) are discarded. These discarded bins are re-associated with other peak bins. This process iterates until all peak bins have at least three bins associated.

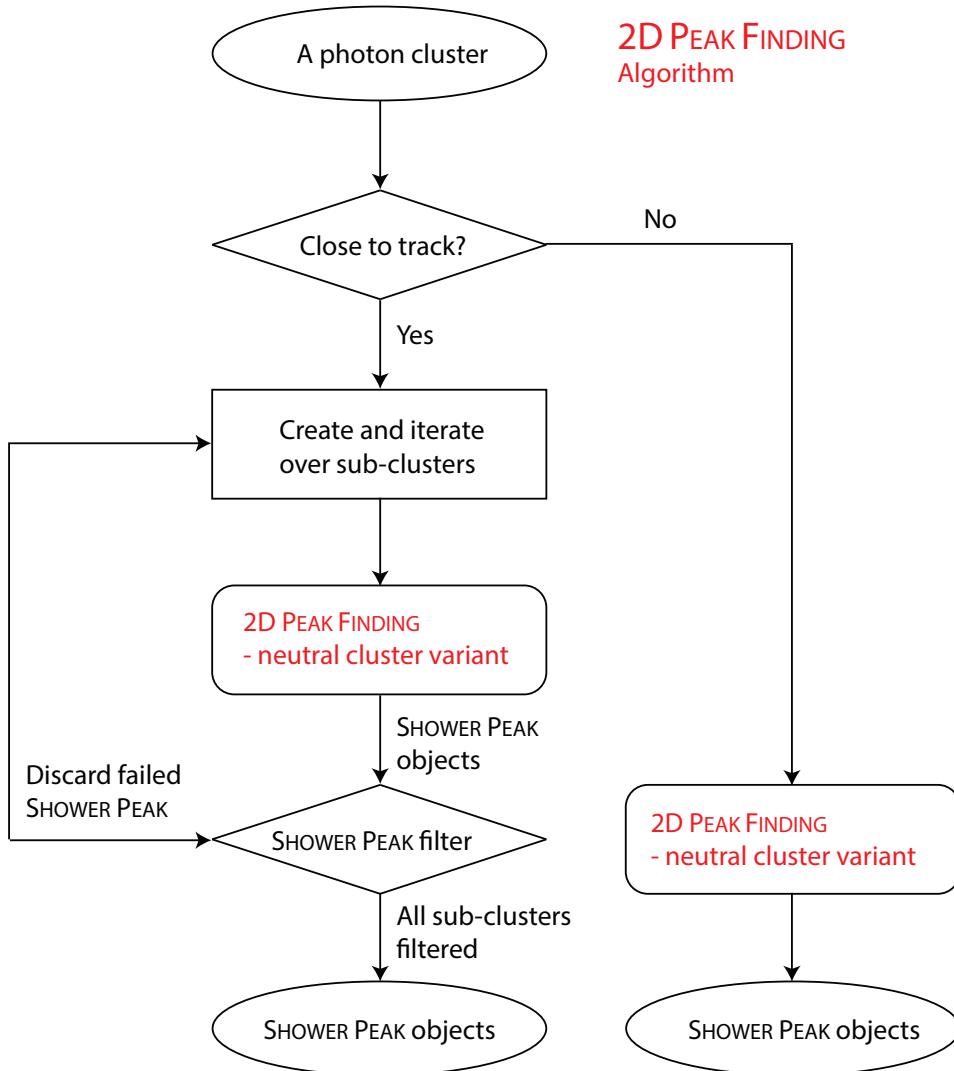
After filtering peaks, SHOWER PEAK objects are created. One SHOWER PEAK object contains one peak bin and associated non-peak bins. The associated calorimeter hits within the bins are attached to the SHOWER PEAK object as well. If multiple peaks are identified in a cluster, multiple SHOWER PEAK objects are created as outputs.

### 5.3.2.3 Charged cluster variant

In a dense jet environment, if a photon next to a charged hadron is carefully reconstructed, the charged particle reconstruction is improved. This step aims to carefully identify

photon candidates next to charged hadrons, by using track information and features of EM showers. An EM shower typically starts in the first few layers of the ECAL with direction of the EM shower largely unchanged when the shower develops.

Figure 5.7 shows the main steps in the full 2D PEAK FINDING algorithm, including the treatment of clusters close to tracks. The first step of the algorithm, "Close to track", determines if a cluster is close to a track. If the distance between a cluster and the closest track projection onto the front of the ECAL is fewer than 3 mm, the charged cluster variant of the 2D PEAK FINDING algorithm is applied to the cluster.



**Figure 5.7:** Main steps of the 2D PEAK FINDING algorithm, including the charged cluster variant: identifying whether the cluster is close to a track; creating and iterating over sub-clusters; applying 2D PEAK FINDING algorithm neutral cluster variant to sub-clusters; filtering SHOWER PEAK objects in sub-clusters; creating final SHOWER PEAK objects.

The "Create and iterate over sub-clusters" step performs the following. The ECAL is sliced longitudinally to create fiducial volumes. For example, the default three slices will result in three fiducial volumes in the ECAL. Each fiducial volume covers the space from the front of the ECAL to a third, to two thirds, and to the back of the ECAL. Three sub-clusters are created from the calorimeter hits of the cluster that are contained in each fiducial volume. In the example of the ILD detector model, the first sub-cluster is formed with the calorimeter hits of the cluster in the first 10 layers of the ECAL. The second and the third sub-cluster are formed with the calorimeter hits of the cluster in the first 20 layers the ECAL and the entire ECAL respectively.

After creating sub-clusters, the neutral cluster variant of the 2D PEAK FINDING algorithm is applied to each sub-cluster to find peaks. SHOWER PEAK objects are created from the 2D PEAK FINDING algorithm.

The SHOWER PEAK objects created from each sub-cluster undergo the "SHOWER PEAK filter" step. Peaks in the first sub-cluster are preserved. Peaks in the second sub-cluster are preserved if the peak bin position is the same as a preserved peak bin position in the first sub-cluster, allowing a shift in the peak bin position by no more than one neighbouring bin. Similarly, peaks in the third sub-cluster are preserved if the peak bin position is the same as a preserved peak bin position in the second sub-cluster, allowing a shift in the peak bin position by no more than one neighbouring bin. Only preserved peaks in the third sub-cluster are used to form the final SHOWER PEAK objects.

Furthermore, if a peak bin is within one neighbouring bin of a track projection bin, the peak is discarded. The track projection bin is the bin where the track projection onto the front of the ECAL projects onto the 2D histogram.

Figure 5.8 illustrates an example of three sub-clusters created during the charged variant of the 2D PEAK FINDING algorithm. Peaks with associated bins and track projection bins are labelled. Figure 5.8a shows the first sub-cluster, created using calorimeter hits in the first 10 layers of the ECAL. One peak is identified. Figure 5.8b shows the second sub-cluster, created using calorimeter hits in the first 20 layers of the ECAL. Two peaks are identified. Only the blue peak is in the same position of the blue peak in the first sub-cluster. Hence, the blue peak in the second cluster is preserved. Figure 5.8c shows the third sub-cluster, created using calorimeter hits in the entire ECAL. Three peaks are identified. Only blue peak is in the same position of the blue peak in the second sub-cluster. Hence, only blue peak in the third sub-cluster is preserved. The

preserved blue peak and associated bins in the third sub-cluster are then used to create one SHOWER PEAK object.

#### 5.3.2.4 Inclusive mode

The time complexity of iterating a 2D histogram is  $O(n^2)$  for a  $n$  bins by  $n$  bins sized histogram (default  $n = 41$ ). Therefore, for the purpose of speed, it is undesirable to have a large number of bins. Having a small finite-sized histogram speeds up the computation. However, because of the finite size of the histogram, only calorimeter hits projected onto the histogram would be considered by the peak finding algorithm. Calorimeter hits projected outside the histogram would not be used when SHOWER PEAK objects are constructed. This behaviour is suitable if the algorithm is only interested in finding the EM shower cores, for example, the PHOTON RECONSTRUCTION algorithm. However, for the purpose of photon splitting, all calorimeter hits from the parent photon should be used to form daughter photons. Hence the inclusive mode of the 2D PEAK FINDING algorithm is developed, and allows calorimeter hits projected outside the histogram to be associated with identified peaks.

### 5.3.3 Photon ID test

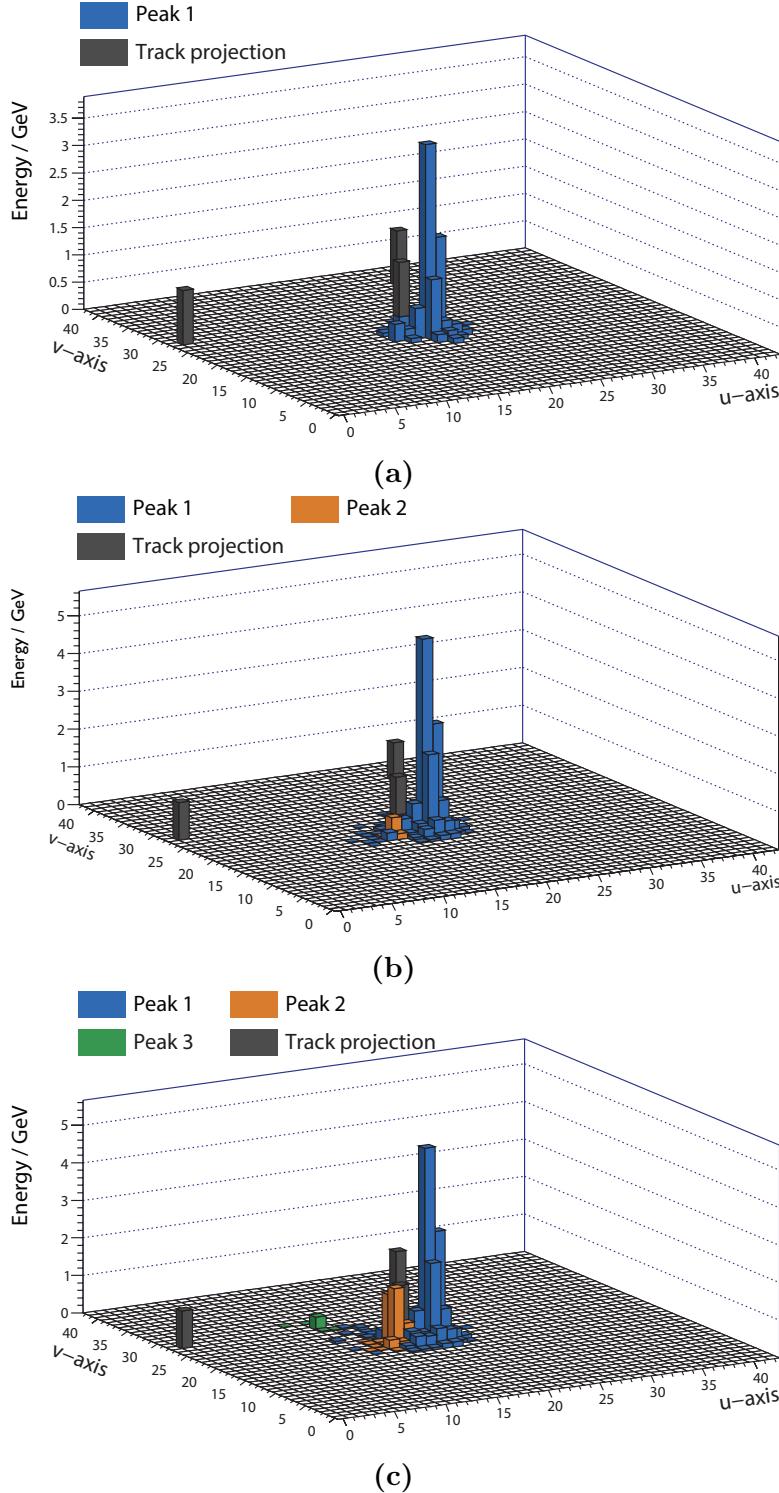
This step applies the photon ID test on the SHOWER PEAK object. The photon ID test uses a multi-dimensional likelihood classifier.

#### 5.3.3.1 Variables used in likelihood classifier

Variables used in the likelihood classifier exploit the differences between a characteristic electromagnetic shower and a hadronic shower, and the fact that a photon is less likely to be close to track projections onto the front of the ECAL, than a cluster of a charged particle. Variables used in the classifier are listed in table 5.1. All plots in this section are produced from simulated  $e^+e^- \rightarrow Z'Z'$  events, where  $Z' \rightarrow u\bar{u}/d\bar{d}/s\bar{s}$ , at  $\sqrt{s} = 500$  GeV.

Two variables are obtained from the EM longitudinal shower profile:  $t_0$  is the start layer in the ECAL of the fitted EM longitudinal shower profile, shown in figure 5.9a; and  $\delta l$  is fractional difference of the observed EM longitudinal shower profile to the expected EM longitudinal shower profile described in equation 5.1:

$$\delta l = \frac{1}{E_0} \sum_i |\Delta E_{obs}^i - \Delta E_{EM}^i|, \quad (5.5)$$



**Figure 5.8:** An illustration of sub-clusters created during the charged cluster variant of the 2D PEAK FINDING algorithm. Sub-clusters are created using calorimeter hits in: a) first 10 layers of the ECAL, b) first 20 layers of the ECAL, and c) the entire ECAL. Peaks with associated bins and track projection bins are labelled.

where  $E_0$  is the energy of the EM shower;  $\Delta E_{EM}^i$  is the energy of the expected EM longitudinal shower profile in bin  $i$ ;  $\Delta E_{obs}^i$  is the energy of the observed EM longitudinal shower profile in bin  $i$ ; the index  $i$  is summed over the ECAL layers as the EM longitudinal shower profile is binned according to the ECAL layers; and the quantity  $\delta l$  is minimised as a function of  $\Delta E_{EM}^i$ , which is a function of  $t_0$  via equation 5.1. The  $\delta l$  distributions for photons and non-photon particles are shown in figure 5.9b. For a true photon,  $t_0$  and  $\delta l$  are expected to be small, as an EM shower should start in the first few layers of the ECAL and the observed EM longitudinal shower profile should be similar to an expected EM longitudinal shower profile.

Three variables are obtained from the EM transverse shower profile: the variable  $\langle w \rangle$  is the energy weighted root-mean-square distance of all bins in a SHOWER PEAK to its peak bin, a measure of the transverse shower size, shown in figure 5.9c; the variable  $\langle w_{UV} \rangle$  is the smallest ratio of the two energy weighted root-mean-square distances of all bins in a SHOWER PEAK to its peak bin in each of the U, V axis direction, a measure of the circularity of the transverse shower; the variable,  $\tilde{E}_{cluster}$ , is the ratio of the energy of the SHOWER PEAK object to the cluster energy, a measure of the dominance of the SHOWER PEAK in a cluster.

The last variable used in the classifier,  $d$ , is the distance between the candidate and the closest track projection onto the front of the ECAL. The SHOWER PEAK object is less likely to be a photon if it is close to a track. The distributions of  $d$  for photons and non-photon particles are shown in figure 5.9d.

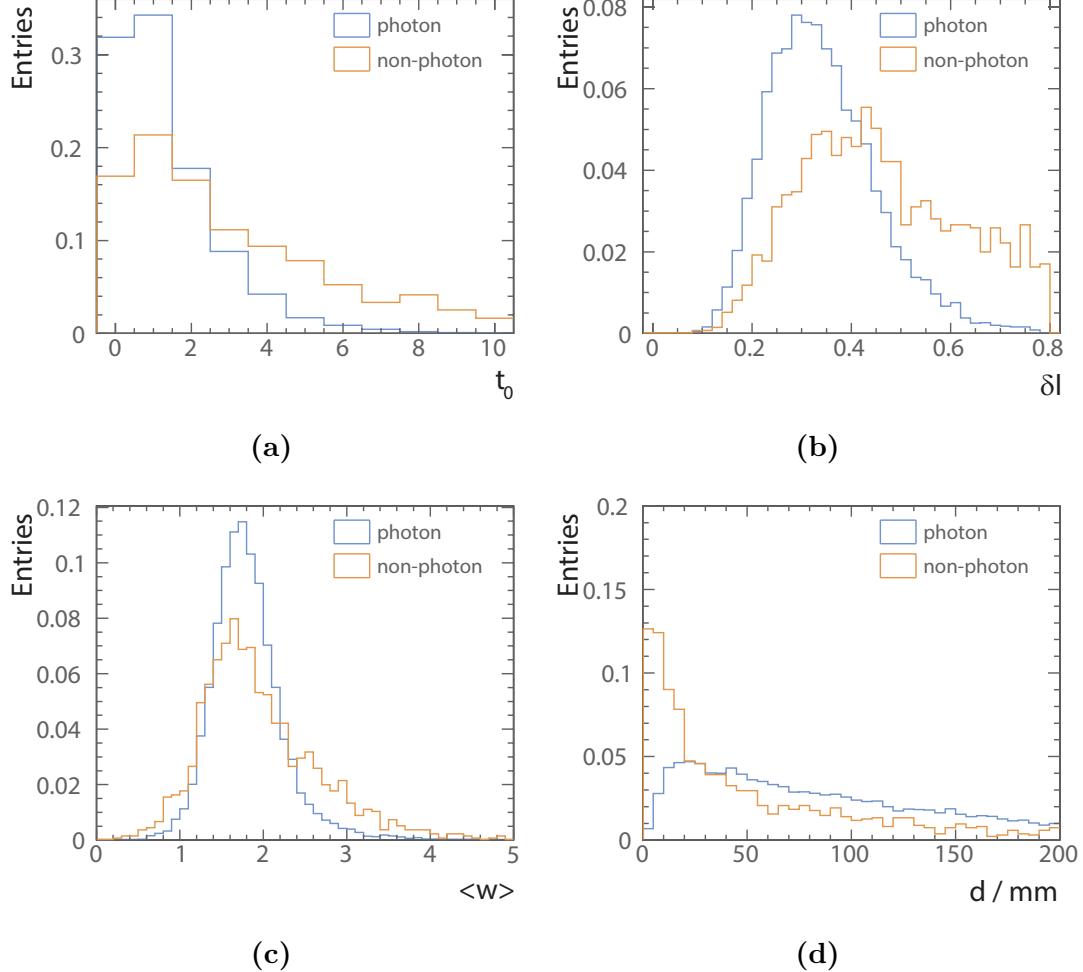
Categories	Variables
EM longitudinal shower profile	$\delta l, t_0$
EM transverse shower profile	$\langle w \rangle, \langle w_{UV} \rangle, \tilde{E}_{cluster}$
Distance to track	$d$

**Table 5.1:** Variables used in the likelihood classifier for photon ID test.

### 5.3.3.2 Projective Likelihood classifier

Projective likelihood classifier is used for the photon ID test due to its low requirement on computing resources, comparing to a Boosted Decision Tree classifier or a Neural Network classifier.

The probability distributions for each variable for photons and non-photon particles are obtained in the training stage. The distributions of these variables are normalised



**Figure 5.9:** Distributions of: a) the start layer from the longitudinal shower profile ( $t_0$ ); b) the fractional difference of the observed shower profile to the expected EM shower profile ( $\delta l$ ); c) the energy weighted root-mean-square distance of all bins in a SHOWER PEAK to its peak bin ( $\langle w \rangle$ ); and d) the distance between the photon candidate and the closest track projection onto the front of the ECAL ( $d$ ). The area under each curve is normalised to unity.

to unity, stored in binned histograms. The classifier is improved by realising that the variable distributions depend on photon energy. Thus the variables distributions are stored separately for different photon energy ranges. Eight photon energy ranges are used by binning the distribution of photon energies at 0.2, 0.5, 1, 1.5, 2.5, 5, 10, 20 GeV.

The training stage of the classifier uses simulated  $e^+e^- \rightarrow Z'Z'$  events, where  $Z' \rightarrow u\bar{u}/d\bar{d}/s\bar{s}$ , at a centre-of-mass energy of 500 GeV. Events at the centre-of-mass energy of 500 GeV allow the training of photon with energies greater than 20 GeV.

In the applying stage of the classifier, for a given particle with the energy in the energy bin  $\alpha$ , the classifier output is given by

$$\text{PID}_\alpha = \frac{N \prod_i^6 P_i}{N \prod_i^6 P_i + N' \prod_i^6 P'_i} \quad (5.6)$$

where  $P_i$  and  $P'_i$  are the values of the bins that the particle fallen in the  $i^{th}$  variable probability distributions of the respective photon and non-photon particles in the energy bin  $\alpha$ ; the variables  $N$  and  $N'$  are the number of respective photons and non-photon particles in the energy bin  $\alpha$  in the training samples.

A SHOWER PEAK object passes the photon ID test if

$$\begin{cases} \text{PID} > 0.6, & \text{if } 0.2 < E < 0.5 \text{ GeV} \\ \text{PID} > 0.4, & \text{if } E \geq 0.5 \text{ GeV} \end{cases} \quad (5.7)$$

where  $E$  is the energy of the particle. Two values of the cuts on PID are motivated by the fact that it is more likely to misidentify a low-energy particle as a photon. A low-energy EM shower does not have a dense shower core, and is more difficult to identify. Hence for candidates with energies between 0.2 and 0.5 GeV,  $\text{PID} > 0.6$  is required instead of  $\text{PID} > 0.4$ .

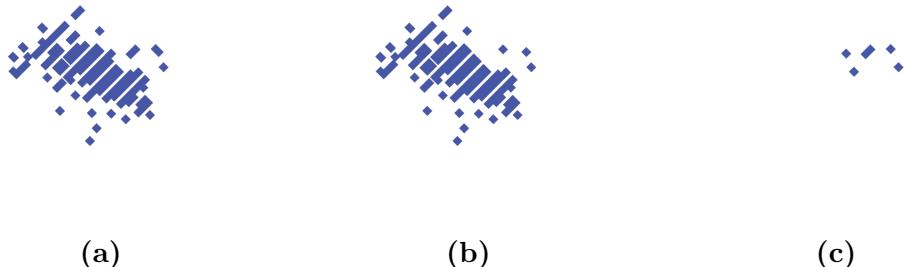
If a SHOWER PEAK object passes the photon ID test, the SHOWER PEAK object would be tagged as a photon. If a SHOWER PEAK object fails the photon ID test, the SHOWER PEAK object will be discarded. Calorimeter hits associated with the discarded SHOWER PEAK object will be passed onto the next stage of the reconstruction.

This step marks the end of the PHOTON RECONSTRUCTION algorithm. The outputs are reconstructed photons, separated from calorimeter hits that are not associated with photons.

## 5.4 Photon fragment removal algorithms in the ECAL

During the reconstruction, it is possible that a core of the photon electromagnetic shower is identified as a photon (the main photon), but the outer part of the shower is reconstructed as a separate particle (the fragment), and identified as a photon or a neutral hadron. Figure 5.10 shows a typical creation of such a photon fragment. A

fragment typically does not have the electromagnetic shower structure, and has a much lower energy than a main photon.

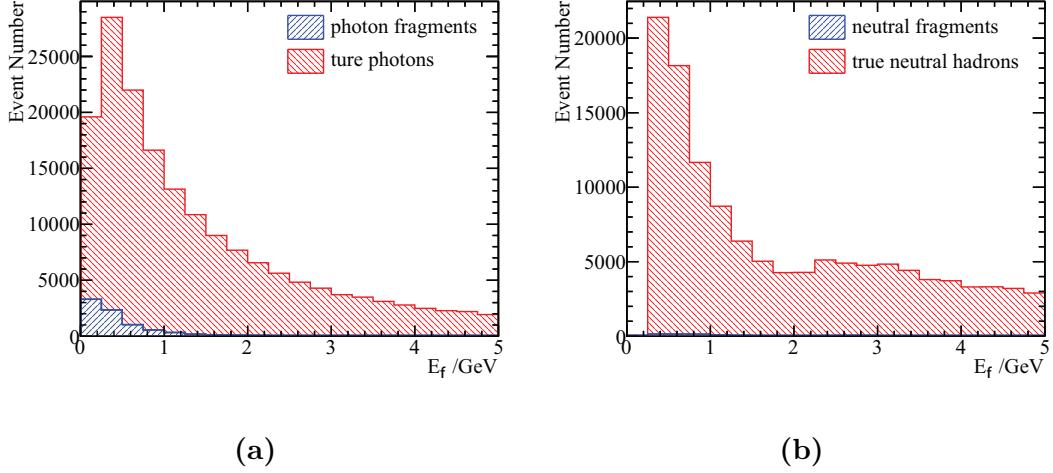


**Figure 5.10:** An event display of a) a typical 10 GeV photon, reconstructed into b) a main photon, and c) a photon fragment.

A photon and a fragment form a photon–fragment pair. Depending on whether the fragment is reconstructed as a photon or a neutral hadron, the photon–fragment pairs are further classified into photon–photon-fragment pairs and photon–neutral-fragment pairs. The neutral fragment refers to the fragment reconstructed as neutral hadrons. The photon–fragment pairs are subsequently divided into low-energy and high-energy pairs, depending on whether the fragment energy ( $E_f$ ) is above 1 GeV. Figure 5.11 shows the energies of the second most energetic reconstructed photon in the photon–photon-fragment pairs, the true photon–photon pairs, photon–neutral-fragment pairs, and true photon–neutral-hadron pairs. Most photon and neutral hadron fragments have energies below than 1 GeV. Hence the energy sub-division was chosen to be at 1 GeV. Plots in this section are obtained with simulated  $e^+e^- \rightarrow Z'Z'$  events, where  $Z' \rightarrow u\bar{u}/d\bar{d}/s\bar{s}$ , at  $\sqrt{s} = 500$  GeV, reconstructed with the PandoraPFA version 1.

There are two variants of the photon fragment removal algorithms: one immediately after the PHOTON RECONSTRUCTION algorithm, and the other one after the charged particle reconstruction (see figure 5.2). Since two algorithms share similar logics for fragment removal, the algorithm used after the charged particle reconstruction will be discussed in detail here. Table B.1 lists cuts for merging photon–photon-fragment pairs and photon–neutral-fragment pairs for both low-energy and high-energy fragments. Using the cuts for photon–photon-fragment pairs with low-energy fragments as an example, each set of logics for merging fragments is discussed. Photon–fragment pair passing any one set of cuts will be merged.

The transverse shower comparison cut merges fragments when the photon–fragment pair looks like one EM shower in the two-dimensional energy deposition projection. The transverse shower comparison requires  $\frac{E_{p1}}{E_m+E_f} > 0.9$ , demanding most energy of the



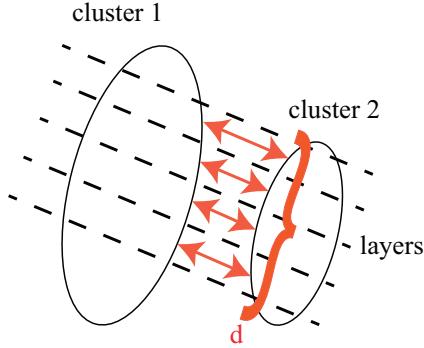
**Figure 5.11:** The energies of the second most energetic reconstructed photon in a) the photon–photon-fragment pairs, and the true photon–photon pairs , and in b) the photon–neutral-fragment pairs, and the true photon–neutral-hadron pairs.

cluster contains in the most energetic peak found by the 2D PEAK FINDING algorithm. It also demands that the second energetic peak should have less than half of the energy in the fragment,  $\frac{E_{p2}}{E_f} < 0.5$ . And the most energetic peak should have more energy than the main photon,  $E_{p1} > E_m$ , where  $E_m$  is the energy of the main photon;  $E_f$  is the energy of the fragment; and  $E_{p1}$  and  $E_{p2}$  are the respective energies of the two most energetic EM showers, identified by the 2D PEAK FINDING algorithm, ordered by descending energy, using the photon–fragment pair as input. Lastly the fragment should be close to the main photon,  $d < 30 \text{ mm}$ , where  $d$  is the average energy weighted intra-layer distance between the photon and the fragment, illustrated schematically in figure 5.12:

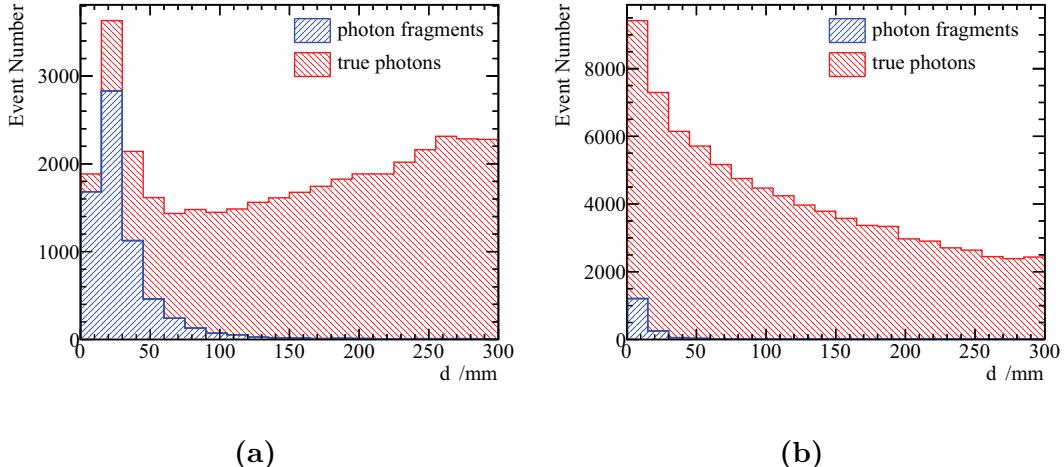
$$d = \frac{\sum_i^{\text{layers}} d_l^i E_f^i}{\sum_i^{\text{layers}} E_f^i} \quad (5.8)$$

where index  $i$  indicates  $i^{\text{th}}$  layer of the ECAL; the parameter  $d_l^i$  is the minimum distance between calorimeter hits of the photon and the fragment in the  $i^{\text{th}}$  layer; and  $E_f^i$  is the total energy of calorimeter hits of the fragment in the  $i^{\text{th}}$  layer of the ECAL. Figure 5.13a and figure 5.13b show the average energy weighted intra-layer distance ( $d$ ) for photon–photon-fragment pairs and the true photon–photon pairs, for low-energy and high-energy fragments respectively. Photon–fragment pairs typically have a small distance separation between the photon and the fragment.

The close proximity cut merges fragments when the fragment has a low energy and is spatially close to the main photon:  $d < 20 \text{ mm}$  and  $E_f < 0.2 \text{ GeV}$ .



**Figure 5.12:** An illustration of the average energy weighted intra-layer distance between the photon and the fragment,  $d$ .

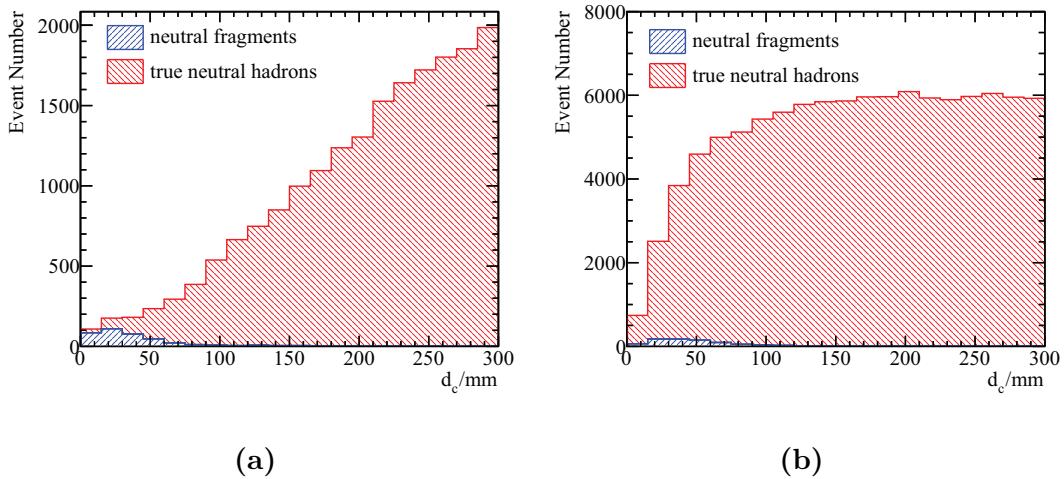


**Figure 5.13:** Distributions of average energy weighted intra-layer distances,  $d$ , for: a) photon–photon-fragment pairs and the true photon–photon pairs for low-energy fragments; and b) photon–photon-fragment pairs and the true photon–photon pairs for high-energy fragments.

The third set of logics, small fragment cuts, is when fragments that are spatially close to the main photon and have very few number of associated calorimeter hits. Either the photon–fragment pair satisfies:  $d < 30$  mm;  $d_c < 50$  mm; and  $N_{calo} < 40$ . Or the photon–fragment pair satisfies:  $d < 30$  mm, and  $N_{calo} < 50$ , where  $N_{calo}$  is the number of the calorimeter hits in the fragment. The multiple sets of cuts allow the merging of a fragment with a fewer number of calorimeter hits with a slightly larger distance separation to the main photon, or the merging of a fragment with a slightly larger number of calorimeter hits with a smaller distance separation to the main photon.

The fourth set of logics merges low-energy fragment in the endcap region of the detector. The cut demands:  $d_c < 60$  mm;  $|\cos(\theta_Z)| > 0.7$ ;  $E_f < 0.6$  GeV; and  $N_{calo} < 40$ .  $|\cos(\theta_Z)|$  is the absolute value of the cosine of the polar angle of the main photon with

respect to the beam direction.  $d_c$  gives the distance between centroids of the photon and the fragment, which is a computationally quick measurement. Figure 5.14a and figure 5.14b shows the distance between centroids ( $d_c$ ) for photon–neutral-fragment pairs and the true photon–neutral-hadron pairs, for low-energy and high-energy fragments respectively. Photon–fragment pairs typically have a small distance separation between the photon and the fragment.



**Figure 5.14:** Distributions of distances between centroids,  $d_c$ , for: a) photon–neutral-fragment pairs and the true photon–neutral-hadron pairs for low-energy fragments; and b) photon–neutral-fragment pairs and the true photon–neutral-hadron pairs for high-energy fragments.

The last set of logics is that the merged fragment should be relatively low energetic. The ratio of the fragment energy to the main photon energy should be less than 0.01. The distance between the pair should satisfies  $d < 40$  mm and  $d_h < 20$  mm, where  $d_h$  is the minimum distance between calorimeter hits of the photon and the fragment.

Cuts for high-energy fragments ( $E_f > 1$  GeV) only has logics for transverse shower comparison and relative low energy fragment, as the cut on the absolute energy of the low-energy fragments is not applicable for the high-energy fragments.

This merging test is considered for all possible photon–fragment pairs. If multiple photon–fragment pairs with the same photon pass the merging test, the pair with the smallest distance metric,  $d$ , will be merged.

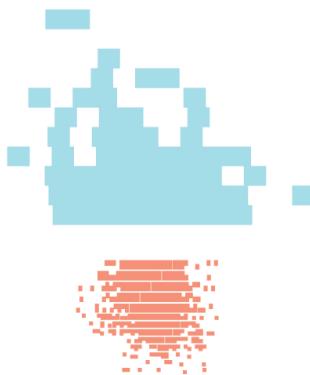
Since all possible photon–fragment pairs are considered, this is a costly cooperation with  $O(n^2)$  time complexity for  $n$  particles. The speed is improved by considering only pairs with  $d < 80$  mm.

### 5.4.1 Photon fragment removal algorithm after the Photon Reconstruction algorithm

The photon fragment removal algorithm immediately after the PHOTON RECONSTRUCTION algorithm shares similar logics as the stated above. The cuts for merging fragments are listed in table B.2.

## 5.5 Photon fragment removal algorithm in the HCAL

There is another type of fragments originated from the leakage effect of the ECAL. When a high-energy EM shower is not fully contained in the ECAL, the shower deposits energy in the HCAL, which often forms a neutral hadron in the HCAL. An example of a 500 GeV photon reconstructed into a main photon in the ECAL (red) and a neutral hadron fragment in the HCAL (blue) is shown in figure 5.15. This section presents an algorithm to merge fragments in the HCAL to the main photon.



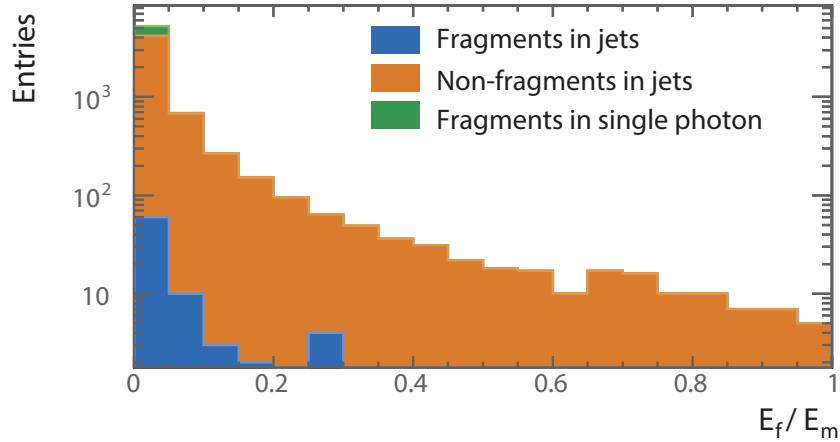
**Figure 5.15:** An event display of a typical 500 GeV photon, reconstructed into a main photon in the ECAL (red) and a neutral hadron fragment in the HCAL (blue).

Photon fragments in the HCAL are spatially close to the main photon. A cone obtained from fitting the main photon, if extended to the HCAL, should contain most of the calorimeter hits of the fragment. These features allow a set of cuts developed to merge fragments in the HCAL, which are listed in table 5.2.

This algorithm uses photons in the ECAL and neutral hadrons in the HCAL as inputs. It then considers all pairs of reconstructed photons and neutral hadrons. Photon–fragment pairs passing all the cuts will be merged.

The adjacent in layers cut demands that photon cluster deposit energies in the last outer layer of the ECAL and the fragment deposits energies in the first inner layer of the HCAL.

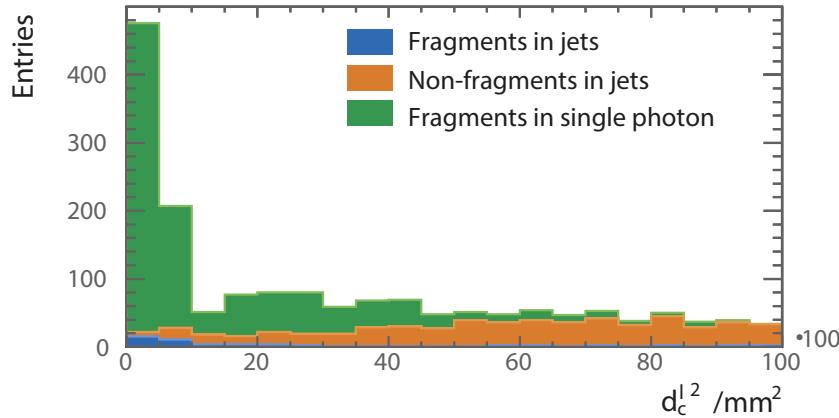
Another requirement for merging, energy comparison cut, is that the fragment should have a low energy relative to the main photon. The variables  $E_m$  and  $E_f$  are the energy of the main photon and the energy of the fragment respectively. The ratio,  $\frac{E_f}{E_m}$ , has to be less than 0.1 for merging. Figure 5.16 shows the distributions of the energy fractions ( $\frac{E_f}{E_m}$ ) after passing the adjacent in layers cut, for photon fragments in jet samples (blue), non-fragments in jet samples (orange), and photon fragments in single-photon samples (green). Jet samples are  $e^+e^- \rightarrow Z'Z'$  events, where  $Z' \rightarrow u\bar{u}/d\bar{d}/s\bar{s}$ , at  $\sqrt{s} = 500$  GeV, reconstructed with the PandoraPFA version 1. Single-photon samples are single 500 GeV photon events, reconstructed with the PandoraPFA version 1. The cut  $\frac{E_f}{E_m} < 0.1$  contains most of the fragments.



**Figure 5.16:** The stacked distributions of the energy fractions ( $\frac{E_f}{E_m}$ ) after passing the adjacent in layers cuts, for photon fragments in jet samples (blue), non-fragments in jet samples (orange), and photon fragments in single-photon samples (green).

The distance comparison cuts requires that fragment in the HCAL to be merged to the main photon should be spatially close to the main photon, measured by three distance metrics: the variable  $d_c^l$  is the distance between the centroid position of the calorimeter hits of the main photon in the last outer layer in the ECAL, and the centroid position of the calorimeter hits of the fragment in the first inner layer of the HCAL; the variable  $d_{fit}^l$  is the shortest distance between the direction fitted with the calorimeter hits of the main photon in the last outer layer in the ECAL, and the direction fitted with the calorimeter hits of the fragment in the first inner layer of the HCAL; and  $d_{fit}$  is the shortest distance between the direction fitted with the main photon, and the direction fitted with the fragment. These three distances should be small for merging.

The cuts demand:  $d_c^l \leq 173$  mm;  $d_{fit}^l \leq 100$  mm; and  $d_{fit}^l \leq 100$  mm. Figure 5.17 shows the distributions of  $d_c^{l^2}$  after passing the adjacent in layers cut and the energy comparison cut, for photon fragments in jet samples (blue), non-fragments in jet samples (orange), and photon fragments in single-photon samples (green). The cut at  $d_c^l \leq 173$  mm ( $d_c^{l^2} \leq 3000$  mm $^2$ ) covers most of the fragments.



**Figure 5.17:** The stacked distributions of  $d_c^{l^2}$  after passing the adjacent in layers cuts and the energy comparison cuts, for photon fragments in jet samples (blue), non-fragments in jet samples (orange), and photon fragments in single-photon samples (green).

The projection comparison cut states that the direction of the fragment to be merged to the main photon should be similar to the direction of the main photon. The variable  $r_f$  is the energy weighted root-mean-square distance of a calorimeter hit in the fragment to the direction fitted with the main photon. The cut requires  $r_f \leq 45$  mm.

The shower width comparison cut requires that the shower width of the fragment to be merged to the main photon and the shower width of the main photon are similar. Variable  $w_m^l$  is the root-mean-square distance of the calorimeter hits of the main photon in last outer layer in the ECAL to the centroid of the calorimeter hits in the same layer. Variable  $w_f^l$  is the root-mean-square distance of the calorimeter hits of the fragment in the first inner layer in the HCAL to the centroid of the calorimeter hits in the same layer. The ratio  $\frac{w_f^l}{w_m^l}$  needs to be in the range from 0.3 to 5 to pass the cut. The generous upper bound is because the HCAL cell size is much larger than the cell size of the ECAL.

The last cut, the cone comparison cut, demands that when a cone obtained by fitting the main photon in the ECAL is extended to the fragment in the HCAL, the cone should contain a significant amount of the fragment. The variable,  $\frac{N_{\text{cone}}}{N_f}$ , the fraction of the calorimeter hits in the fragment in the cone comparing to the calorimeter hits in the fragment, has to be greater than 0.5 for merging.

Photon fragment recovery	Cuts
Adjacent in layers	yes
Energy comparison	$\frac{E_f}{E_m} \leq 0.1$
Distance comparison	$d_c^l \leq 173 \text{ mm}; d_{fit}^l \leq 100 \text{ mm}; d_{fit} \leq 100 \text{ mm}$
Projection comparison	$r_f \leq 45 \text{ mm}$
Shower width comparison	$0.3 \leq \frac{w_f^l}{w_m^l} \leq 5$
Cone comparison	$\frac{N_{cone}}{N_f} \geq 0.5$

**Table 5.2:** The cuts for merging photon fragment in the HCAL to the main photon in the ECAL.

If multiple photon–fragment pairs pass the cuts with the same fragment, the pair with highest  $\frac{N_{cone}}{N_f}$  will be merged.

## 5.6 Photon splitting algorithm

Another aspect of photon reconstruction is to split accidentally merged photons. During the event reconstruction, it is possible that photons are accidentally merged if they are spatially close. Hence this algorithm addresses this issue and tries to split merged photons.

If a photon has the topology of multiple spatially closed photons, the parent photon will be split into several daughter photons. Extra care are taken if the parent photon is close to a track projection onto the front of the ECAL. Table 5.3 lists the cuts used in the algorithm.

The algorithm works as follows. If an energetic photon is identified, the 2D PEAK FINDING algorithm will be used to identify EM showers in the parent photon. If energy of the parent photon is bigger than  $E_{c1}$ , and the energy of the 2<sup>nd</sup> energetic EM shower is bigger than  $E_{c2}$ , the parent photon will be split into daughter photons according to the number of EM showers identified by the 2D PEAK FINDING algorithm.

The values of  $E_{c1}$  and  $E_{c2}$  depend on whether the parent photon is close to a track projection onto the front of the ECAL. The algorithm demands higher values of  $E_{c1}$  and  $E_{c2}$ , if the photon is close to the track projection. The number of nearby charged tracks is counted as number of tracks with the track projection onto the front of the ECAL fewer than 100 mm to the parent photon centroid position. If there is no nearby tracks to the parent photon,  $E_{c1}$  is set to 10 GeV and  $E_{c2}$  is set to 1 GeV. If there is one nearby

track,  $E_{c1}$  is set to 10 GeV and  $E_{c2}$  is set to 5 GeV. If there is more than one nearby track,  $E_{c1}$  is set to 20 GeV and  $E_{c2}$  is set to 10 GeV.

The constraint on  $N_p$ , the number of EM showers identified in the parent photon, should be fewer than five, as one reconstructed photon is unlikely to be accidentally merged from more than four photons.

Photon splitting	Cuts
Cuts	$E > E_{c1}, E_{p2} > E_{c2}, N_p < 5$
$E_{c1}$ and $E_{c2}$ values	
0 track nearby	$E_{c1} = 10 \text{ GeV}, E_{c2} = 1 \text{ GeV}$
1 track nearby	$E_{c1} = 10 \text{ GeV}, E_{c2} = 5 \text{ GeV}$
$> 1$ tracks nearby	$E_{c1} = 20 \text{ GeV}, E_{c2} = 10 \text{ GeV}$

**Table 5.3:** Cuts used in the photon splitting algorithm.

## 5.7 Photon reconstruction performance

Three different versions of the PandoraPFA are used to demonstrate the improvement of the photon reconstruction performance:

- with no stand-alone photon reconstruction algorithms;
- with a stand-alone photon reconstruction algorithm from PandoraPFA version 1;
- with full photon algorithms described above, incorporated in PandoraPFA version 3;

Without photon reconstruction algorithms, PandoraPFA applies a simple photon ID at the end of the reconstruction. In PandoraPFA version 1, there is a rudimentary photon reconstruction algorithm. In PandoraPFA version 3 contains all the photon algorithms presented in this chapter. The photon algorithms in PandoraPFA version 3 have replaced the photon reconstruction algorithm in PandoraPFA version 1.

Firstly the photon reconstruction performance with full photon algorithms implemented in PandoraPFA version 3 is compared with the performance with no stand-alone photon algorithms. Afterwards, the photon reconstruction performance is compared with the performance obtained from PandoraPFA version 1. The photon reconstruction performances of individual photon algorithms are then characterised, followed by the characterisation of the performance of the photon algorithms in PandoraPFA version 3.

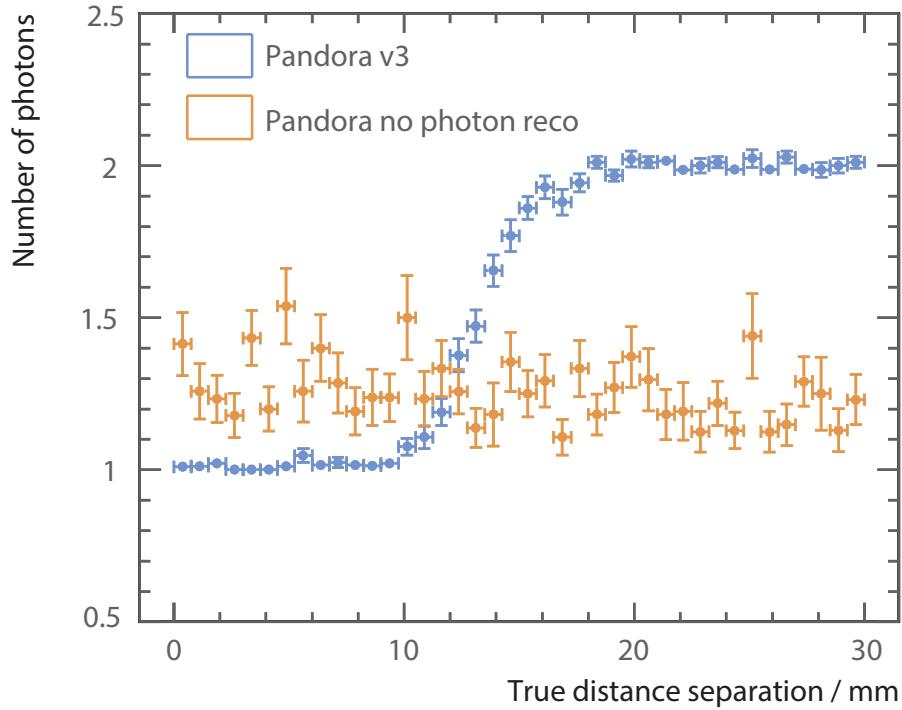
### 5.7.1 Improvement over no stand-alone photon algorithms

The improvement in the photon reconstruction is demonstrated using MC samples with two photons. The two-photon samples were generated with an uniform distribution in the solid angle of the first photon, and an uniform distribution in the opening angle between the photon pair. Events are discarded if there is a photon converting to electron pairs in the tracking detector or a photon escapes the detector undetected. The events are further restricted to the photon depositing energies in barrel and endcap regions only to avoid the barrel/endcap overlap region. Events were reconstructed using the nominal ILD detector model.

Figure 5.18 shows the average number of reconstructed photons as a function of true distance separation between two photons, using two-photon samples with photon energies of 500 GeV and 50 GeV, reconstructed with and without photon algorithms. For the reconstruction without the photon algorithms, the number of photon fluctuates between 1 and 1.5 for a distance separation of 0 to 30 mm between two photons. For the reconstruction with the photon algorithms, two photons start to be resolved at a distance separation of 10 mm between two photons, and fully resolved at a distance separation of 20 mm. The average number of reconstructed photon is 2 at a distance separation of 20 mm.

The improvement in photon reconstruction leads to a considerable improvement in the jet energy resolution. Jet energy resolution is defined as the root-mean-square divided by the mean for the smallest width of distribution that contains 90% of entries, using  $e^+e^- \rightarrow Z'Z'$  events, where  $Z' \rightarrow u\bar{u}/d\bar{d}/s\bar{s}$ , at barrel region. The angular cut is to avoid the barrel/endcap overlap region. The light quark decay of the  $Z'$  is used to avoid the complication of missing momentum from semi-leptonic decay of heavy quarks. Using 90% of the entries is robust and focus on the Gaussian part of the jet energy distribution. The total jet energies are sampled at centre-of-mass energies of 91, 200, 360 and 500 GeV.

As shown in figure 5.19, jet energy resolutions are much better at  $\sqrt{s} = 360$  GeV and 500 GeV for the reconstruction with photon algorithms. By identifying photons before reconstructing charged particles in a dense jet environment, there are fewer calorimeter hits left for the charged particle reconstruction. However, at  $\sqrt{s} = 91$  GeV and 200 GeV, jet energy resolutions are worse for the reconstruction with photon algorithms, because photon algorithms are developed and optimised with jet environments at a high centre-of-mass energy of 500 GeV.

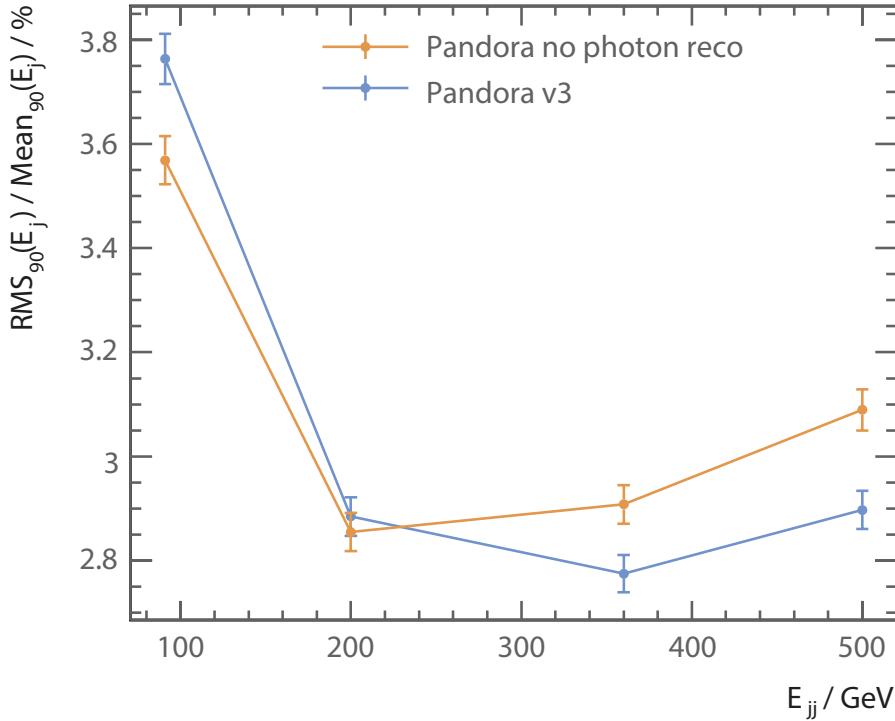


**Figure 5.18:** Average number of reconstructed photons using two-photon samples with photon energies of 500 GeV and 50 GeV, without (orange) and with (blue) photon algorithms, as a function of the true distance separation between two photons.

The impact of photon algorithms on the jet energy resolution was studied using the same jet samples with the perfect photon reconstruction, which identifies photons by associating calorimeter hits using the truth information. The photon confusion term, which are defined as the quadrature differences of the jet energy resolutions between a non-cheated reconstruction and a perfect photon reconstruction, is a measure of the failure in the photon reconstruction. Table 5.4 lists the photon confusion terms as a function of the centre-of-mass energies for the jet sample. The photon confusion terms, except at  $\sqrt{s} = 91$  GeV, have been reduced to 0.9% for the reconstruction with photon algorithms.

Photon confusion	$\sqrt{s} = 91$ GeV	200 GeV	360 GeV	500 GeV
PandoraPFA without photon algorithms	0.7%	0.9%	1.3%	1.4%
PandoraPFA with full photon algorithms	1.4%	0.9%	0.9%	0.9%

**Table 5.4:** Photon confusion terms as a function of total jet energies in the  $e^+e^- \rightarrow Z'Z'$  events, where  $Z' \rightarrow u\bar{u}/d\bar{d}/s\bar{s}$  at barrel region for reconstruction with and without photon algorithms.



**Figure 5.19:** Jet energy resolutions as a function of the total jet energy using  $e^+e^- \rightarrow Z'Z'$  events, where  $Z' \rightarrow u\bar{u}/d\bar{d}/s\bar{s}$  at barrel region. The orange and bottom points represent the reconstruction without and with photon algorithms, respectively.

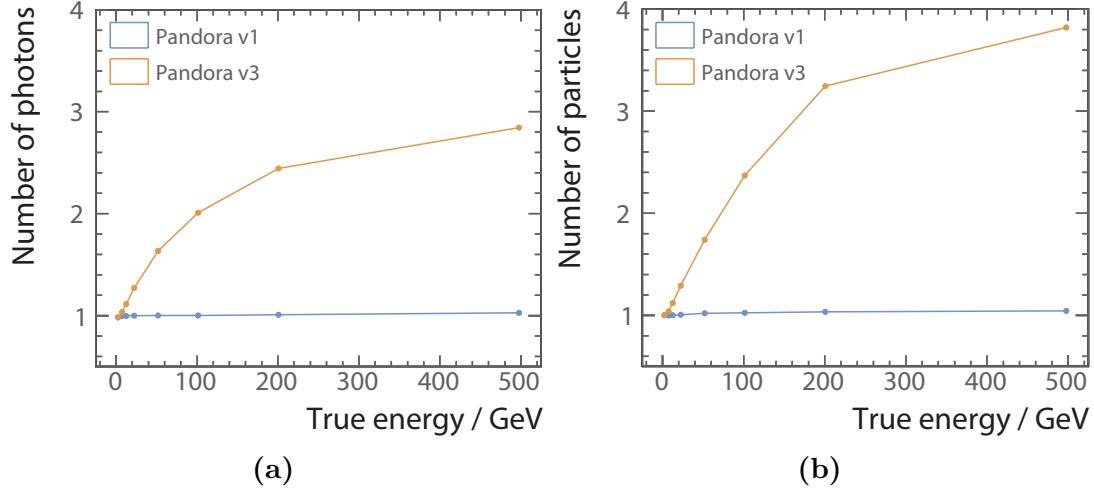
### 5.7.2 Improvement over PandoraPFA version 1

This section reviews the photon reconstruction improvement from PandoraPFA version 1 to version 3, using single-photon, two-photon, and jet samples.

The single-photon MC samples were generated with an uniform distribution in the solid angle of the photon. Other samples were generated and simulated in the same way as previously. The same pre-selection as previously was applied to the single-photon and two-photon samples .

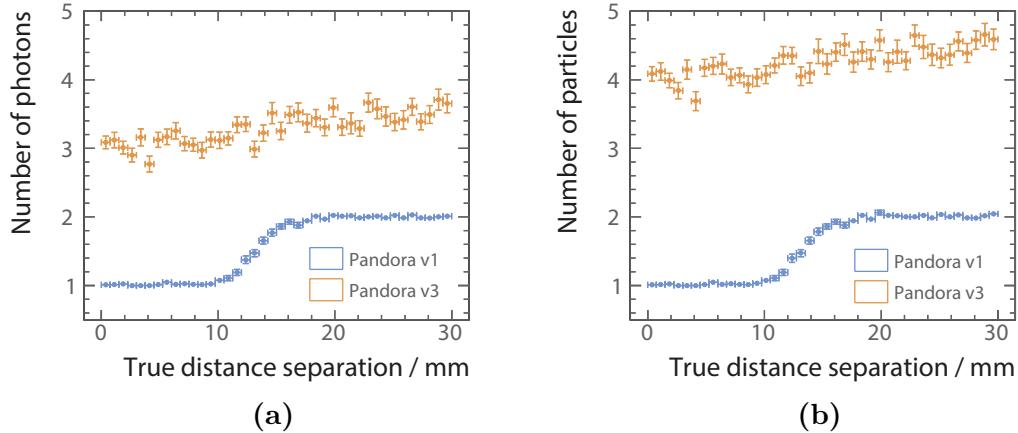
Figure 5.20a shows the average number of reconstructed photons as a function of the true photon energies, using single-photon samples. Figure 5.20b shows the average number of reconstructed particles as a function of the true photon energies. Drastic decrease in the number of fragments can be seen in both plots for photons up to 500 GeV.

Figure 5.21 shows the numbers of reconstructed photons and particles as a function of the true distance separation of the two photons, using a two-photon sample with photon energies of 500 GeV and 50 GeV. The average numbers of photons and particles for reconstruction with PandoraPFA version 3 are both below 2.05 at a distance separation of 30 mm, which is significantly lower than the numbers for reconstruction with PandoraPFA



**Figure 5.20:** Average numbers of: a) photons; and b) particles, as a function of their true energies using single-photon samples. For both figures, the top orange and bottom blue points are reconstructed with PandoraPFA version 1 and version 3, respectively.

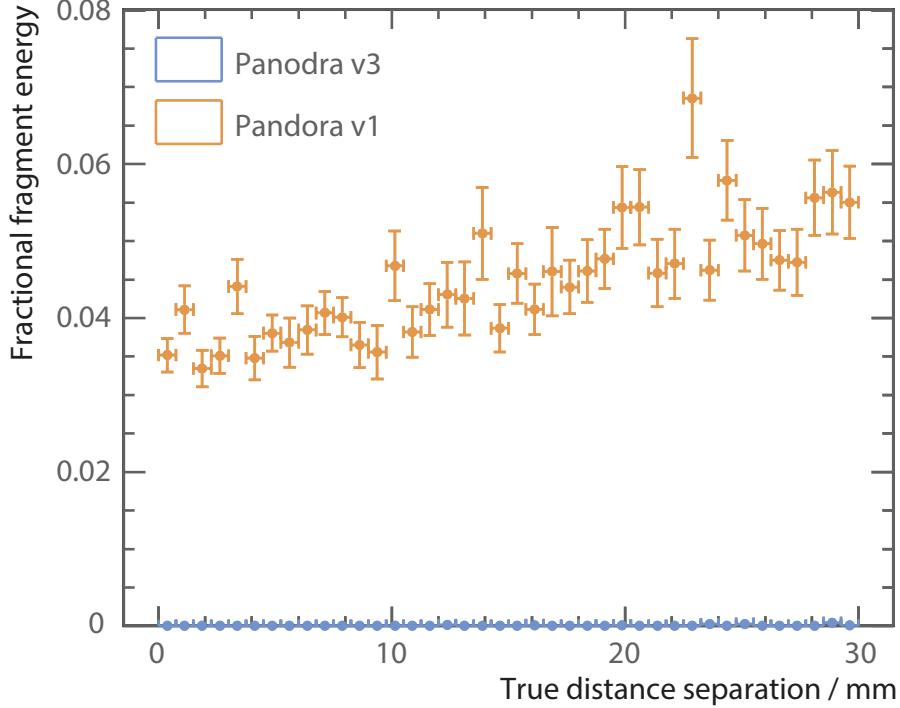
version 1. For reconstruction with PandoraPFA version 3, two photons start to be resolved at a distance separation of 10 mm, and fully resolved at a distance separation of 20 mm.



**Figure 5.21:** Average numbers of: a) photons; and b) particles, as a function of the true distance separation between two photons, using two-photon samples with photon energies of 500 GeV and 50 GeV. For both figures, the top orange and bottom blue points represent the reconstruction with PandoraPFA version 1 and version 3, respectively.

Another metric to reflect the improvement in photon reconstruction is the fraction of the fragment energy to the total energy in an event. In a two-photon sample, the fragment energy is defined as the total energy of particles excluding the two most energetic photons. Shown in figure 5.22, using two-photon sample with photon energies of 500 GeV and 50 GeV, a reduction in fragment energy can be seen clearly going from PandoraPFA version 1 to version 3. For the photon reconstruction in PandoraPFA version 3, the

average fragment energy fraction is below 0.1% up to a distance separation of 30 mm, whilst around 5% energy would be in fragments for the reconstruction with PandoraPFA version 1.

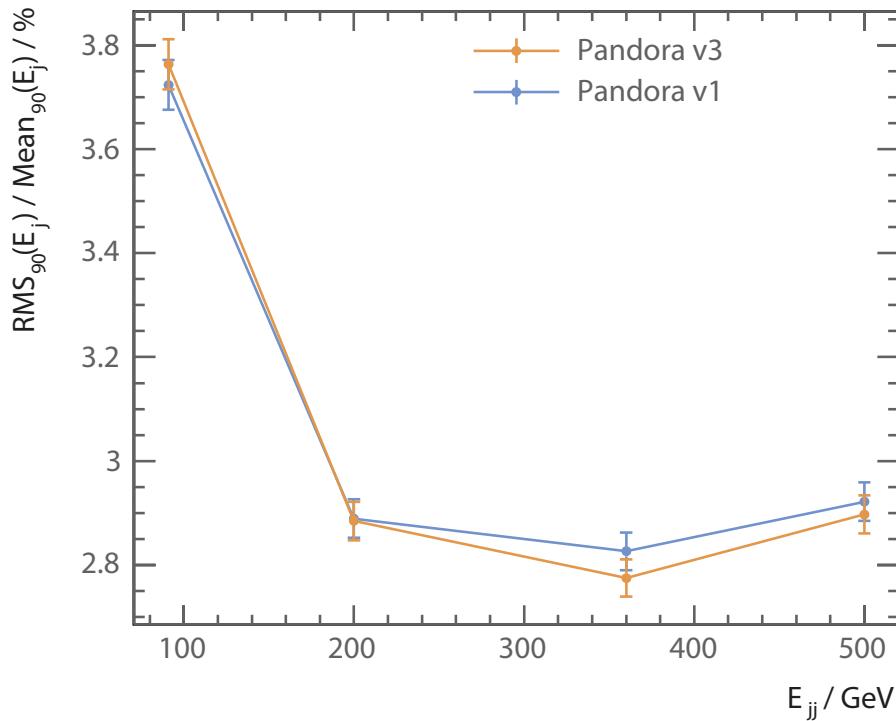


**Figure 5.22:** Average fraction of fragments energy to the total energy in the event, as a function of the true distance separation between two photons, using a two-photon sample with photon energies of 500 GeV and 50 GeV. The top orange and bottom blue points represent the reconstruction with PandoraPFA version 1 and version 3 respectively.

The reduction in the fragments, as shown in the reconstruction of the single-photon and two-photon samples, leads to a small improvement in the jet energy resolutions at a high jet energy. Using the same jet sample as in the previous section, the jet energy resolutions are better at total jet energies of 360 and 500 GeV with the photon reconstruction in PandoraPFA version 3, as shown in figure 5.23.

### 5.7.3 Performance of individual photon algorithms

Two-photon events with photon energies of 500 GeV and 500 GeV are used to show the incremental improvement of the performance of individual photon algorithms. Figure 5.24 shows the average number of reconstructed particles as a function of true distance separation between two photons, reconstructed with full photon algorithms with PandoraPFA version 3 (blue), reconstructed with only fragment removal algorithms in the

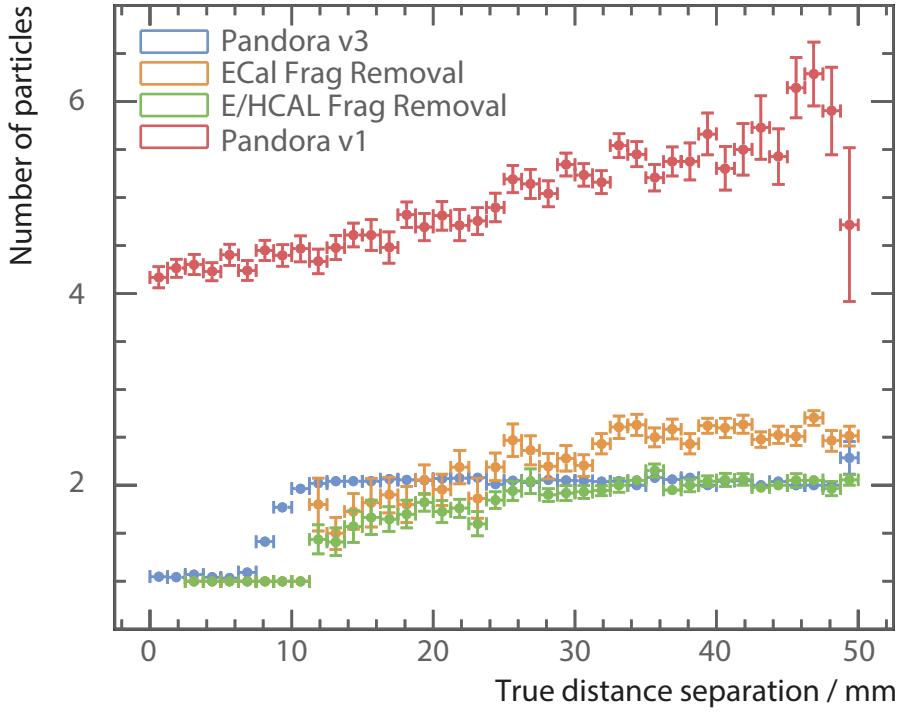


**Figure 5.23:** Jet energy resolutions as a function of the total jet energy using  $e^+e^- \rightarrow Z'Z'$  events, where  $Z' \rightarrow u\bar{u}/d\bar{d}/s\bar{s}$ , at barrel region. The top orange and bottom blue points represent the reconstruction with PandoraPFA version 1 and version 3.

ECAL and photon reconstruction in PandoraPFA version 1 (orange), reconstructed with fragment removal algorithms in the ECAL and the HCAL and photon reconstruction in PandoraPFA version 1 (green), and reconstructed with PandoraPFA version 1 (red).

For the reconstruction with fragment removal algorithm in the ECAL (orange), the number of fragments is reduced significantly, when it is compared with photon reconstruction in PandoraPFA version 1 (red). With the additional fragment removal algorithm in the HCAL (green), the number of fragments is reduced further. At a distance separation of 40 mm, there is on average less than 0.05 fragment per photon pair for the reconstruction with fragment removal algorithms in the ECAL and the HCAL (green).

The introduction of the photon reconstruction and photon splitting algorithm (blue) makes the photon pair resolve at a much shorter distance separation between two photons. Photon pairs start to be resolved at a distance separation of 5 mm, and fully resolved at a distance separation of 15 mm when reconstructed with full photon algorithms in PandoraPFA version 3.

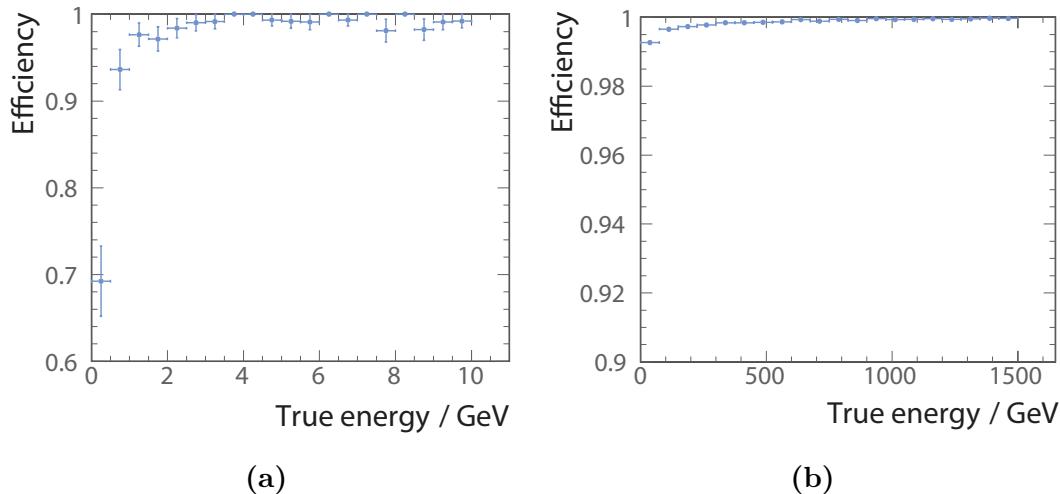


**Figure 5.24:** Average numbers of photons, as a function of the true distance separation between two photons, using a two-photon sample with photon energies of 500 GeV and 500 GeV. The blue, orange, green, and red points represent the reconstruction with PandoraPFA version 3, the reconstruction with fragment removal in the ECAL and photon reconstruction in PandoraPFA version 1, the reconstruction with fragment removal in the ECAL and the HCAL and photon reconstruction in PandoraPFA version 1, the reconstruction with PandoraPFA version 1, respectively.

#### 5.7.4 Photon reconstruction performance with PandoraPFA version 3

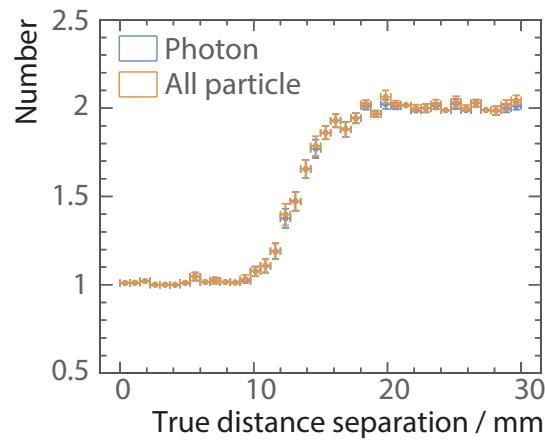
The photon reconstruction performance with PandoraPFA version 3 is demonstrated with single-photon and two-photon samples.

Figure 5.25 shows the average single photon reconstruction efficiency as a function of the true photon energies, using single-photon samples. In a single-photon sample, an event can have an efficiency of 1 or 0, depending on whether there is a reconstructed photon corresponding to the true photon. The average single photon reconstruction efficiency is above 98% for photons with energies above 2 GeV, and above 99.5% for photons with energies above 100 GeV. The low efficiency in the first bin in figure 5.25a, for photon energies in the range from 0 to 0.25 GeV, is because photon reconstruction algorithms do not attempt to reconstruct photons with energies below 0.2 GeV.



**Figure 5.25:** Single photon reconstruction efficiency as a function of true photon energies, using single-photon samples for: a) the low photon energy regime; and b) the high photon energy regime.

Figure 5.26 shows the average numbers of reconstructed photons and particles as a function of the true distance separation between two photons, using a two-photon sample with photon energies of 500 GeV and 500 GeV. A good match between the number of photons and the number of particles is achieved. The average numbers of photons and particles are both fewer than 2.05 for a distance separation beyond 20 mm, less than 1 fragment produced per 20 events.

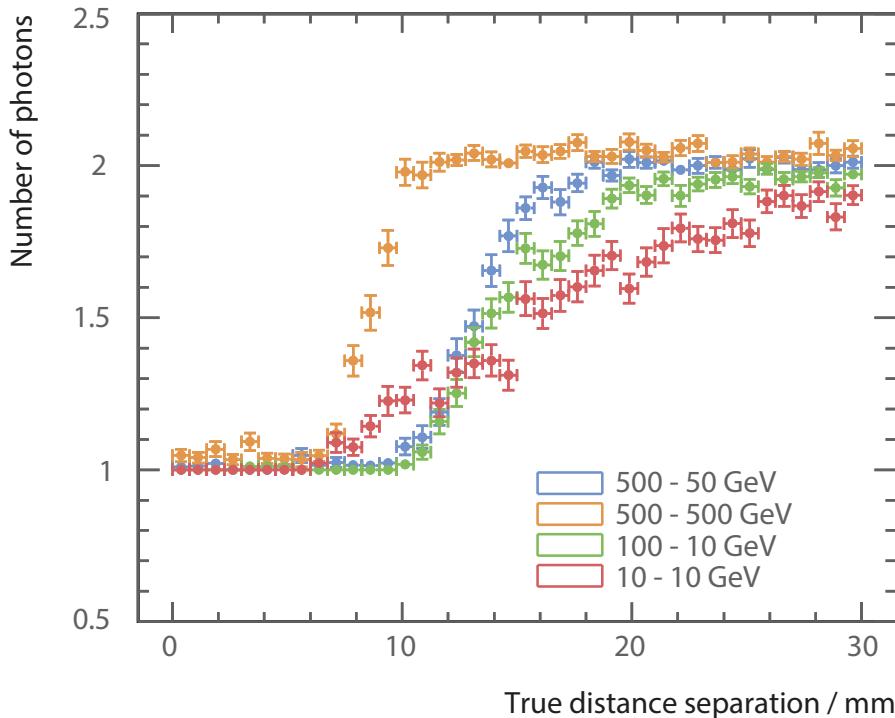


**Figure 5.26:** Average numbers of reconstructed photon (blue) and particle (orange), as a function of the true distance separation between two photons, using two photons of 500 GeV and 50 GeV per event samples.

The ability to resolve of a photon pair depends on energies of two photons. Figure 5.27 shows the average number of photon reconstructed using two-photon samples as a function of the true distance separation between two photons, for different photon

energies. When the energies of two photons are similar, the distance of two photons starting to be resolved is shorter. This is because that when the two photon showers have similar sizes, the 2D PEAK FINDING algorithm can exploit the symmetry in the size of the EM showers. For example, 500 GeV–500 GeV photon pair and 10 GeV–10 GeV photon pair start to be resolved at a distance separation of 6 mm, which is about one ECAL cell length in the simulated nominal ILD detector. In contrast, photon pairs with different energies, for example 500 GeV–50 GeV and 100 GeV–10 GeV pairs, start to be resolved at a distance separation of 10 mm, which is about two ECAL cells length.

For an energetic photon, it is easier to identify the photon, because the electromagnetic shower core is denser and contains more energies than the peripheral calorimeter hits. Therefore separating two energetic photons is easier than separating two low-energy photons. As shown in figure 5.27, at a distance separation of 20 mm, 500 GeV–500 GeV photon pairs are fully resolved, whereas approximately only 60% of 10 GeV–10 GeV photon pairs are resolved.



**Figure 5.27:** Average numbers of reconstructed photons for four different photon pairs: 500 GeV–50 GeV (blue), 500 GeV–500 GeV (orange), 100 GeV–10 GeV (green), and 10 GeV–10 GeV (red), as a function of the true distance separation between two photons.

## 5.8 Summary

Using the ILD detector model, the single photon reconstruction efficiency is above 98% for photons with energies above 2GeV, and above 99.5% for photons with energies above 100 GeV.

The number of photon fragments produced have been greatly reduced. Using a two-photon sample with photon energies of 500 GeV and 50 GeV, the average numbers of photons and particles beyond a distance separation of 20 mm are both less than 2.05, where the true value is 2.

The distance of the minimal resolved photon pairs is reduced to 6 mm for two photons with same energies, and 10 mm for two photons with different energies.

The jet energy resolution has been improved for high centre-of-mass energies. The photon confusion terms, except at  $\sqrt{s} = 91$  GeV, have been reduced to 0.9%.



# Chapter 6

## Tau Lepton Decay Mode Classification

*'I once tried standing up on my toes to see far out in the distance, but I found that I could see much farther by climbing to a high place.'*

— Xun Kuang, 313 BC – 238 BC

The tau pair polarisation correlation from a boson decay can be used to determine statistically if the parent boson is a scalar or a vector, for example, to differentiate a H boson from a Z boson [26]. It can also be used to measure the CP (the product of charge conjugation and parity symmetries) of the Higgs, via  $H \rightarrow \tau^+ \tau^-$  decay process [91].

Since the tau lepton has a very short mean decay lifetime of 290 fs [25], only tau decay products can be detected in the calorimeters and tracking detectors of the ILD detector. Therefore, the performance of the calorimetric and tracking systems determines the ability to reconstruct tau lepton decay products and to classify different tau decay modes.

The main challenge in classifying tau lepton hadronic decay modes is the reconstruction and separation of spatially close photons. For tau leptons with energies above tens of GeVs, visible decay products are often highly boosted. Consequently electromagnetic showers from photons from  $\pi^0$  decays often overlap in the ECAL. Reconstructing these photons as separate entities requires a good photon reconstruction. Hence the photon reconstruction algorithms described in chapter 5 are used in this study. This chapter

presents a study of the classification of tau decay modes in a highly granular linear collider detector.

## 6.1 Event generation and simulation

Two million  $e^+e^- \rightarrow \tau^+\tau^-$  events at a centre-of-mass energy of 100 GeV were generated with WHIZARD [49]. TAUOLA [53] was used to describe the tau lepton decays with correct spin correlations of the tau decay products. The study was focused on separating tau decay modes. Hence beam effects, such as the initial state radiation and the beam induced background, were not included. The  $e^+e^- \rightarrow \tau^+\tau^-$  events were simulated using the ILD detector model as described in chapter 4.

## 6.2 Event reconstruction

Events were reconstructed with iLCSoft version v01-17-07 [92] and PandoraPFA version 3 [65], using the photon reconstruction algorithms described in chapter 5. An event display of  $e^+e^- \rightarrow \tau^+\tau^-$  interaction reconstructed in the ILD detector is shown in figure 6.1. The top half of the event shows a tau lepton decaying into  $\pi^-\pi^0\nu_\tau$ . The bottom half of the event shows a tau lepton decaying into  $\pi^+\pi^-\pi^-\pi^0\nu_\tau$ .

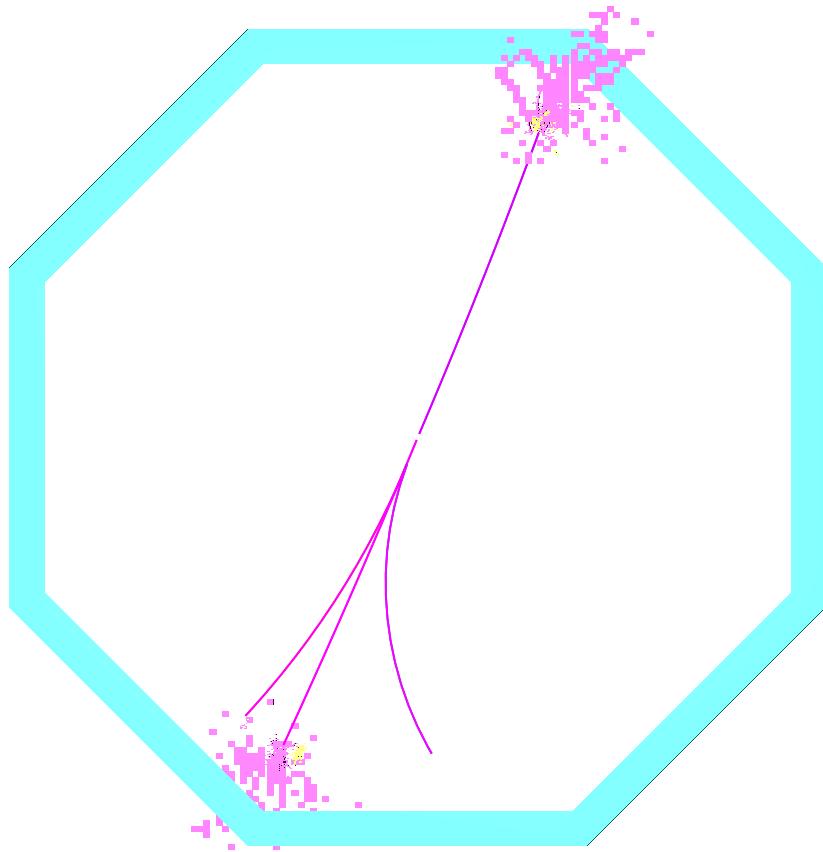
### 6.2.1 Tau decay modes

To study the main decay modes of the tau lepton, decay modes with branching ratios above 2% are classified, which results in seven decay modes, covering 92.58 % of the tau decay branching fraction [3]. The seven tau decay modes, their branching ratios, and detectable final states are listed in table 6.1.

In the  $\rho\nu_\tau$  decay mode, the  $\rho$  meson subsequently decays into  $\pi^-\pi^0$ . In the  $a_1\nu_\tau$  neutral and charged decay modes, the  $a_1$  meson subsequently decays into  $\pi^-\pi^0\pi^0$ , and  $\pi^+\pi^-\pi^-$ , respectively. The invariant masses of the  $\rho$  meson and  $a_1$  meson are  $775.11 \pm 0.34$  MeV and  $1230 \pm 40$  MeV, respectively [3].

### 6.2.2 Tau selection

The simulated  $e^+e^- \rightarrow \tau^+\tau^-$  event contains two tau leptons. Since the tau decay mode classification is applied on a per-tau-lepton basis, the decay products of the two tau



**Figure 6.1:** An event display of a simulated  $e^+e^- \rightarrow \tau^+\tau^-$  event using the ILD detector model. The top half of the event shows a tau lepton decaying into  $\pi^-\pi^0\nu_\tau$ . The bottom half of the event shows a tau lepton decaying into  $\pi^+\pi^-\pi^-\pi^0\nu_\tau$ . Purple lines represent  $\pi^\pm$  tracks in the tracking detectors. Purple squares represent calorimeter hits of the  $\pi^\pm$  hadronic showers in the ECAL and the HCAL. Yellow squares represent calorimeter hits of EM showers of photons from  $\pi^0 \rightarrow \gamma\gamma$ . The blue region is the transverse cross section of the ECAL barrel part.

Decay mode	Detectable final state	Branching ratio
$e^-\bar{\nu}_e\nu_\tau$	$e^-$	$17.83 \pm 0.04\%$
$\mu^-\bar{\nu}_\mu\nu_\tau$	$\mu^-$	$17.41 \pm 0.04\%$
$\pi^-\nu_\tau$	$\pi^-$	$10.83 \pm 0.06\%$
$\rho\nu_\tau$	$\pi^-\pi^0$	$25.52 \pm 0.09\%$
$a_1\nu_\tau$ neutral	$\pi^-\pi^0\pi^0$	$9.30 \pm 0.11\%$
$a_1\nu_\tau$ charged	$\pi^+\pi^-\pi^-$	$8.99 \pm 0.06\%$
$\pi^+\pi^-\pi^-\pi^0\nu_\tau$	$\pi^+\pi^-\pi^-\pi^0$	$2.70 \pm 0.08\%$

**Table 6.1:** Decay modes, detectable final state particles, and branching ratios of the seven major tau decay modes with the largest branching ratios. Values are taken from [3].

leptons in one event are divided into two sets for individual tau decay mode classification. By identifying the axis of the back-to-back taus in an event, the detector space can be separated in two hemispheres, where particles in each hemisphere correspond to the decay products of one tau lepton.

Separating reconstructed particles in an event into two sets is achieved using the principle thrust axis vector of the event, which is the axis that most particles are aligned to. The principle thrust axis vector,  $\hat{t}$ , is determined by maximising the thrust [93],  $T$ :

$$T = \max_{\hat{t}} \frac{\sum_i |\hat{t} \cdot \vec{p}_i|}{\sum_i |\vec{p}_i|}, \quad (6.1)$$

where  $\vec{p}_i$  is the momentum vector of particle  $i$ ; vector  $\hat{t}$  is the unit principle thrust axis vector; and index  $i$  is summed over all particles in an event. Two sets of particles are obtained based on the sign of the scalar product between the principle thrust axis vector and the momentum vector of a particle; particles with a positive sign of the scalar product are in one set and particles with a negative sign of the scalar product are in another set.

### 6.3 Pre-selection

For the purpose of this study, three pre-selection cuts, based on the MC information of the particles, are used. The cuts and the fractions of tau decays passing the cuts are listed in table 6.2.

Since this study is focused on photon reconstruction in the ECAL to classify tau decay modes, the tau decays with photon converting to electron pairs in the tracking detector are not considered.

The focus of the study is on high energy tau decays. Tau decays with the total visible energy (i.e. not accounting for neutrinos) of tau decay products,  $E_{vis}^{MC}$ , less than 5 GeV are not considered.

Lastly, tau decays are discarded when the tau decay products are in the gap region between barrel and endcap parts of the calorimeters, because there is a degradation in the particle reconstruction efficiency in this region. Tau decays with the generated polar angle of the tau lepton in the region  $0.6 < |\theta_\tau^{MC}| < 0.9$  rad are not considered.

Table 6.2 shows fractions of tau decays passing successive cuts for different tau decay final states. As expected, the cut on photon conversions only affects tau decay modes

with photons in the final states. The cut on the total visible energy of the tau decay products has the greatest effect on the leptonic decay modes with two neutrinos in the final states. The cut on the tau polar angle affects different tau decay modes equally.

Final state	No photon conversion	$E_{vis}^{MC} > 5 \text{ GeV}$	$ \theta_\tau^{MC} $
$e^- \bar{\nu}_e \nu_\tau$	100.0%	84.7%	66.2%
$\mu^- \bar{\nu}_\mu \nu_\tau$	100.0%	85.2%	66.7%
$\pi^- \nu_\tau$	100.0%	88.3%	60.9%
$\pi^- \pi^0 \nu_\tau$	77.1%	76.9%	61.9%
$\pi^- \pi^0 \pi^0 \nu_\tau$	61.3%	61.2%	50.5%
$\pi^+ \pi^- \pi^- \nu_\tau$	100.0%	100.0%	78.0%
$\pi^+ \pi^- \pi^- \pi^0 \nu_\tau$	77.0%	77.0%	61.8%

**Table 6.2:** Fractions of tau decays passing successive pre-selection cuts for different tau decay final states.

## 6.4 Multivariate classification

The classification of different tau decays uses a MVA classifier based on twenty-seven discriminant variables, listed in table 6.3. The particle ID information comes from the output of the PandoraPFA reconstruction.

Category	Variable
Particle numbers	$N_C, N_\mu, N_e, N_\gamma, N_{\pi^-}$
Invariant masses	$m_{vis}, m_C, m_N, m_\gamma, m_{\pi^-}$
Energy variables	$\tilde{E}_{vis}, \tilde{E}_C, \tilde{E}_\mu, \tilde{E}_e, \tilde{E}_\gamma, \tilde{E}_{\pi^-}$
Calorimetric energy	$E_C^{ECAL}/E_C, E^{ECAL}/E$
$\rho(\pi^- \pi^0)$ reconstruction	$m_{\pi^0}^{(\rho)}, m_\rho^{reco}$
$a_1(\pi^- \pi^0 \pi^0)$ reconstruction	$m_{\pi^0}^{(a_1)}, m_{\pi^0}^{*(a_1)}, m_{a_1}^{reco}$
EM shower profile	$\delta l, t_0, \langle w \rangle$
Calorimeter hit information	$\bar{E}_{hit}, MIP$
Track information	$E/p$

**Table 6.3:** Variables used in the MVA classification for the tau lepton decay mode classification.

### 6.4.1 Particle number variables

The most crucial variables for classifying tau decay modes are the number of different types of final state particles. There are five particle number variables used in the MVA classification: the number of charged particles ( $N_C$ ); the number of muons ( $N_\mu$ ); the number of electrons ( $N_e$ ); the number of photons ( $N_\gamma$ ); and the number of charged pions ( $N_{\pi^-}$ ). Here muon and electron IDs are provided by the PandoraPFA outputs.

Figure 6.2a shows distributions of numbers of reconstructed charged particles for different tau decay modes. Over 98% of  $\tau^- \rightarrow \pi^- \nu_\tau$  decays have exactly one reconstructed charged particle, and approximately 95% of  $a_1(\pi^+ \pi^- \pi^-)$  decays give exactly three reconstructed charged particles. Figure 6.2b and figure 6.2c show distributions of numbers of reconstructed muons and electrons respectively for different tau decay modes. Here 99% of  $\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$  decays produce exactly one reconstructed muon, and 99% of  $\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau$  decays have one reconstructed electron. Figure 6.2d shows distributions of numbers of reconstructed photons for different tau decay modes, which distinguishes hadronic tau decay final states with different numbers of  $\pi^0$ . Nearly 75% of  $\rho(\pi^- \pi^0)$  decays give exactly two reconstructed photons, and over 60% of  $a_1(\pi^- \pi^0 \pi^0)$  decays have exactly four photons.

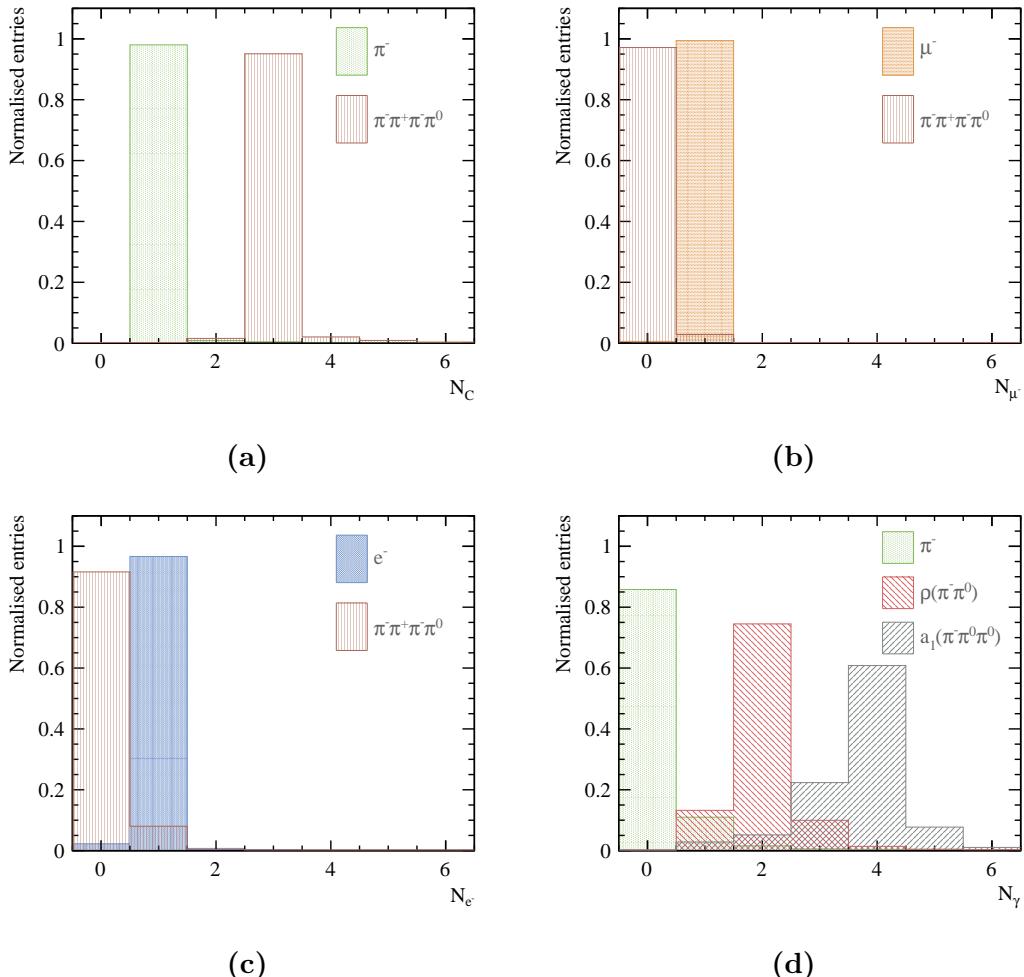
### 6.4.2 Invariant mass variables

Five invariant mass variables are used in the MVA classification: the invariant mass of all reconstructed particles ( $m_{vis}$ ); the invariant mass of all reconstructed charged particles ( $m_C$ ); the invariant mass of all reconstructed neutral particles ( $m_N$ ); the invariant mass of all reconstructed photons ( $m_\gamma$ ); and the invariant mass of all reconstructed charged pions ( $m_{\pi^-}$ ).

Figure 6.3a shows distributions of invariant masses of all reconstructed particles for different tau decay modes. Peaks in the invariant mass distributions can be seen for the  $\rho$  and  $a_1$  decay modes. Figure 6.3b shows distributions of invariant masses of all reconstructed neutral particles for different tau decay modes. Differences in the distributions for the  $\rho$  and  $a_1$  decay modes can be seen.

### 6.4.3 Energy variables

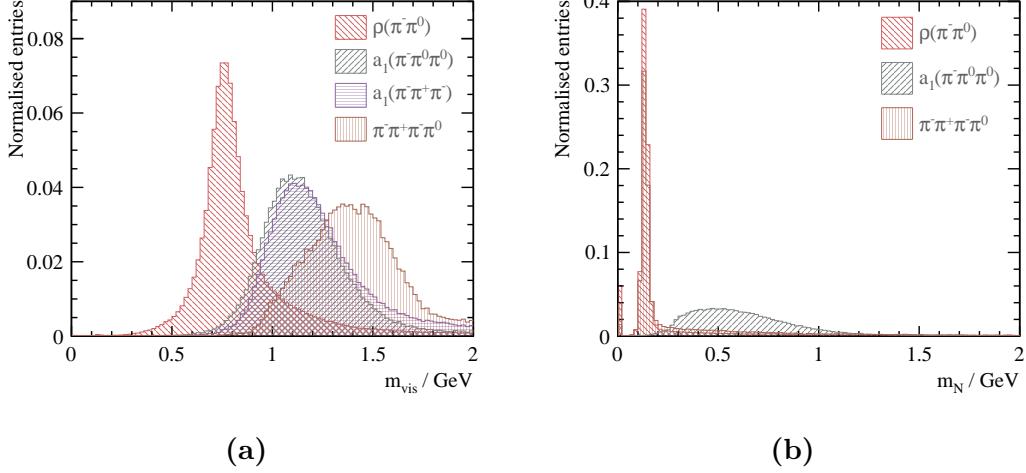
Energy information helps to further separate different tau decay modes. Six energy variables are used in the MVA classification: the normalised total energy of all recon-



**Figure 6.2:** Distributions of the number of reconstructed particles of different types: a) charged particles ( $N_C$ ); b) muons ( $N_\mu$ ); c) electrons ( $N_e$ ); and d) photons ( $N_\gamma$ ). The particle ID information comes from the output of the PandoraPFA reconstruction. The area under the curve for each decay mode is normalised to unity.

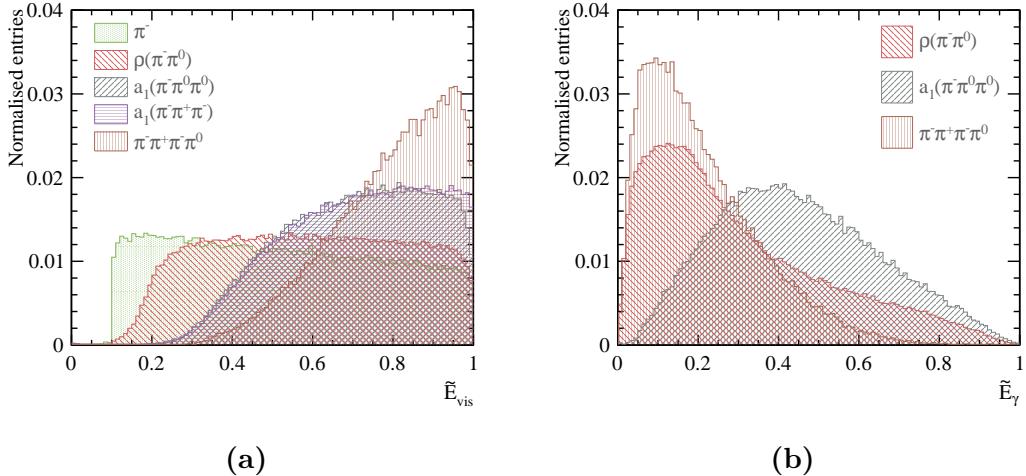
structed particles ( $\tilde{E}_{vis}$ ); the normalised total energy of charged particles ( $\tilde{E}_C$ ); the normalised total energy of muons ( $\tilde{E}_\mu$ ); the normalised total energy of electrons ( $\tilde{E}_e$ ); the normalised total energy of photons ( $\tilde{E}_\gamma$ ); and the normalised total energy of charged pions ( $\tilde{E}_{\pi^-}$ ). All variables are normalised relative to the energy of the associated tau lepton, i.e.  $\tilde{E}_{vis} = E_{vis}/E_\tau$ , where  $E_{vis}$  is the total energy of all reconstructed particles and  $E_\tau$  is the energy of the associated tau lepton.  $E_\tau$  is obtained from the generated value, i.e. 50 GeV.

Figure 6.4a and figure 6.4b show distributions of normalised energies of all reconstructed particles and photons respectively for different tau decay modes. Differences in the distributions for different tau decay modes can be seen. The cut-off at 0.1 for



**Figure 6.3:** Distributions of the invariant mass of a) all particles ( $m_{vis}$ ); and b) all neutral particles ( $m_N$ ). The particle ID information comes from the output of the PandoraPFA reconstruction. The area under the curve for each decay mode is normalised to unity.

the  $\tilde{E}_{vis}$  distribution is due to the pre-selection of  $E_{vis}^{MC} > 5$  GeV, which approximately corresponds to  $\tilde{E}_{vis} > 0.1$ .

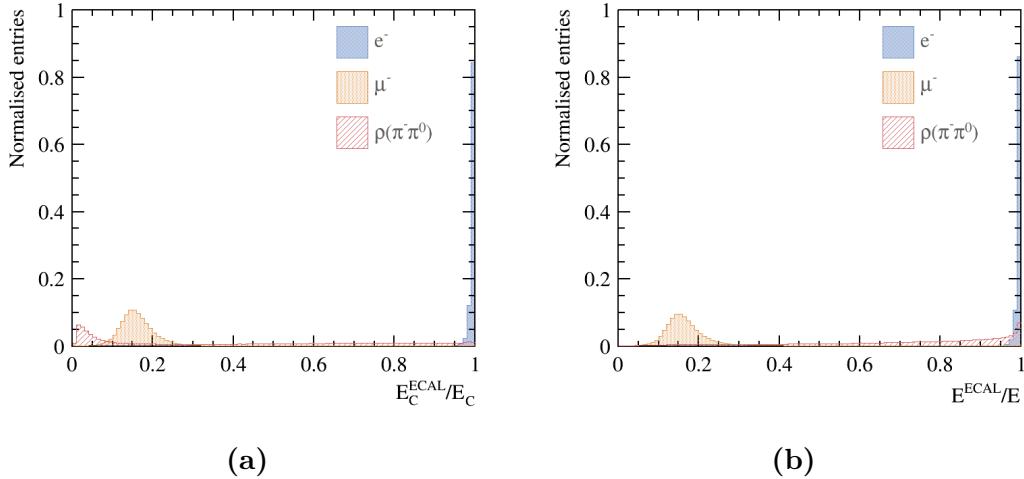


**Figure 6.4:** Distributions of the normalised energies of: a) all reconstructed particles ( $\tilde{E}_{vis}$ ); and b) all reconstructed photons ( $\tilde{E}_\gamma$ ). The particle ID information comes from the output of the PandoraPFA reconstruction. The area under the curve for each decay mode is normalised to unity.

#### 6.4.4 Calorimetric energy variables

Two calorimetric energy variables are used in the MVA classification: the fraction of the energy deposited in the ECAL divided by the energy deposited in the ECAL and HCAL, where only calorimetric deposits associated with charged particles are considered

$(E_C^{ECAL}/E_C)$ ; and the fraction of the energy deposited in the ECAL divided by the energy deposited in the ECAL and HCAL for all particles  $(E^{ECAL}/E)$ . These two variables help to improve the identification of electron and muon decay modes. Figure 6.5 show distributions of  $E_C^{ECAL}/E_C$  and  $E^{ECAL}/E$  for different decay modes. An electron typically deposits over 95% of its energy in the ECAL, and a muon typically deposits 5% to 25% of its energy in the ECAL. The difference between  $E_C^{ECAL}/E_C$  and  $E^{ECAL}/E$  is that photons and neutral hadrons, which deposit most of their energies in the ECAL and in the HCAL respectively, are not included in the calculation of  $E_C^{ECAL}/E_C$ .



**Figure 6.5:** Distributions of the fractions of the energy deposited in the ECAL divided by the energy deposited in the ECAL and HCAL: a) where only calorimetric deposits associated with charged particles are considered ( $E_C^{ECAL}/E_C$ ); and b) for all reconstructed particles ( $E^{ECAL}/E$ ). The particle ID information comes from the output of the PandoraPFA reconstruction. The area under the curve for each decay mode is normalised to unity.

#### 6.4.5 $\rho(\pi^-\pi^0)$ and $a_1(\pi^-\pi^0\pi^0)$ resonances variables

By utilising the photon identification potential of the highly granular ECAL, the identification of the  $\rho(\pi^-\pi^0)$  and  $a_1(\pi^-\pi^0\pi^0)$  decay modes is enhanced by reconstructing the  $\rho$  and  $a_1$  invariant masses. For decays with at least one charged pion and one photon, the reconstruction selects the combination of charged pions and photons that have a invariant mass most consistent with the  $\rho$  or  $a_1$  mass.

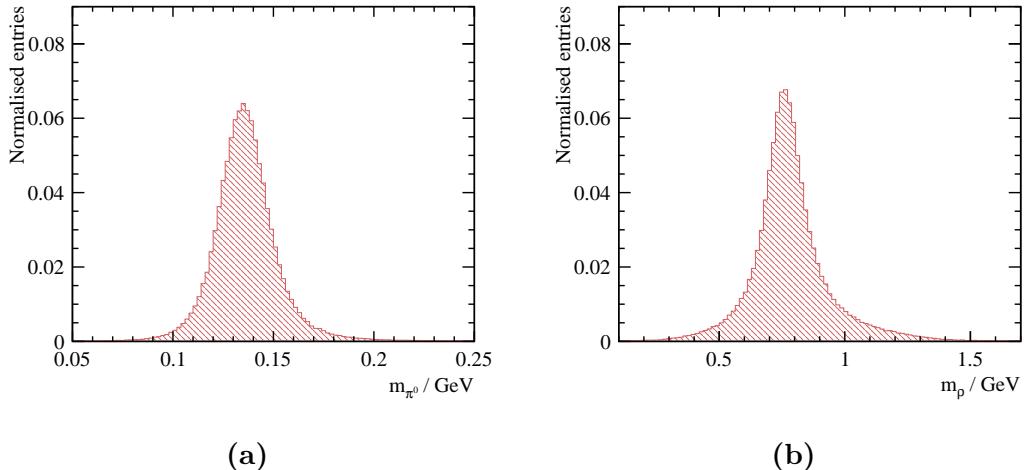
For example, the final state of the  $\rho(\pi^-\pi^0)$  decay mode contains a  $\pi^-$  and a  $\pi^0$ , where  $\pi^0 \rightarrow \gamma\gamma$ . The  $\rho(\pi^-\pi^0)$  decay mode hypothesis test is performed by selecting the combination of the charged pion and photons that gives the smallest value of a  $\chi^2$

function:

$$\chi^2 = \left( \frac{m_{tot} - m_\rho}{\sigma_\rho} \right)^2 + \left( \frac{m_{\gamma_1\gamma_2} - m_{\pi^0}}{\sigma_{\pi^0}} \right)^2, \quad (6.2)$$

where  $m_{\gamma_1\gamma_2}$  is the invariant mass of two photons; the variable  $m_{tot}$  is the total invariant mass of the two photons and one  $\pi^-$ ;  $m_\rho$  and  $m_{\pi^0}$  are the respective true masses of  $\rho$  and  $\pi^0$ ; and  $\sigma_\rho$  and  $\sigma_{\pi^0}$  are the assumed mass resolutions. Figure 6.6 shows the reconstructed invariant mass distributions for  $\pi^0$  and  $\rho$  in the  $\rho(\pi^-\pi^0)$  decay mode, obtained by selecting reconstructed particles using the MC truth information. A mass resolution of 20% is a good approximation for the invariant masses of  $\pi^0$  and  $\rho$ , and it is used for  $\sigma_\rho$  and  $\sigma_{\pi^0}$ .

The particle ID of charged pions and photons comes from the output of the PandoraPFA reconstruction. The  $\chi^2$  function works naturally if there are two reconstructed photons in a decay. If there are more than two photons in a decay, all combinations of two photons are considered and the combination with the smallest value of  $\chi^2$  is chosen. If there is only one photon in a decay, the second term in the equation 6.2 is dropped and  $m_{tot}$  is the total invariant mass of one photon and one  $\pi^-$ .



**Figure 6.6:** Reconstructed invariant mass distributions of the  $\pi^0$  and  $\rho$  in the  $\rho(\pi^-\pi^0)$  decay mode. The reconstructed masses are obtained using the MC truth information to find the corresponding reconstructed particles. The area under the curve is normalised to unity.

The  $\chi^2$  function of equation 6.2 is modified for the  $a_1(\pi^-\pi^0\pi^0)$  decay mode hypothesis test:

$$\chi^2 = \left( \frac{m_{tot} - m_{a_1}}{\sigma_{a_1}} \right)^2 + \left( \frac{m_{\gamma_1\gamma_2} - m_{\pi^0}}{\sigma_{\pi^0}} \right)^2 + \left( \frac{m_{\gamma_3\gamma_4} - m_{\pi^0}}{\sigma_{\pi^0}} \right)^2, \quad (6.3)$$

where the  $\rho$  mass has been replaced by the  $a_1$  mass and other variables are defined in the same way as previously. Four photons and one  $\pi^-$  are required for this  $\chi^2$  function. To resolve the degeneracy between two photon pairs, the requirement of  $|m_{\gamma_1\gamma_2} - m_{\pi^0}| < |m_{\gamma_3\gamma_4} - m_{\pi^0}|$  is imposed.

If there are at least four photons in a decay, all combinations of four photons are considered and the combination with the smallest value of  $\chi^2$  is chosen. If there are three photons in a decay, the last term in the equation 6.3 is dropped and the  $\chi^2$  function becomes:

$$\chi^2 = \left( \frac{m_{tot} - m_{a_1}}{\sigma_{a_1}} \right)^2 + \left( \frac{m_{\gamma_1\gamma_2} - m_{\pi^0}}{\sigma_{\pi^0}} \right)^2, \quad (6.4)$$

where  $m_{tot}$  is the invariant mass of the charged pion and three photons and  $m_{\gamma_1\gamma_2}$  is the invariant mass of two photons. Combinations of photons are iterated.

If there are only two photons in a decay, either the reconstruction failed to reconstruct one photon pair or the reconstruction fails to reconstruct both photon pair. Hence two  $\chi^2$  functions are used and the one with the smallest value is chosen. The first function is

$$\chi^2 = \frac{1}{2} \left( \frac{m_{tot} - m_{a_1}}{\sigma_{a_1}} \right)^2 + \frac{1}{2} \left( \frac{m_{\gamma_1\gamma_2} - m_{\pi^0}}{\sigma_{\pi^0}} \right)^2, \quad (6.5)$$

where  $m_{tot}$  is the invariant mass of the charged pion and two photons and  $m_{\gamma_1\gamma_2}$  is the invariant mass of two photons. The second function is

$$\chi^2 = \left( \frac{m_{tot} - m_{a_1}}{\sigma_{a_1}} \right)^2, \quad (6.6)$$

where  $m_{tot}$  is the invariant mass of the charged pion and two photons.

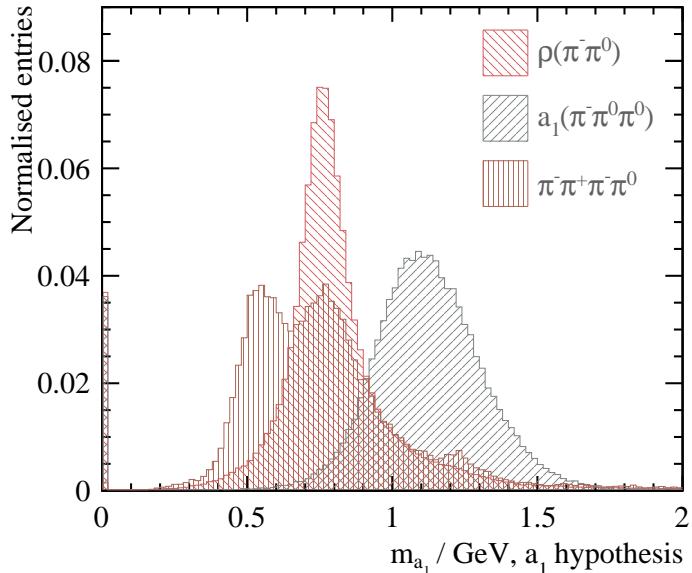
If there is only one photon in a decay,  $m_{tot}$  is the invariant mass of the charged pion and the photon.

From the  $\rho$  invariant mass reconstruction of the  $\rho(\pi^-\pi^0)$  decay hypothesis test, two variables are obtained and used in the MVA classification to help to identify  $\rho(\pi^-\pi^0)$  decay mode: the  $\rho$  mass ( $m_\rho^{reco} \equiv m_{tot}$  in equation 6.2) and the  $\pi^0$  mass ( $m_{\pi^0}^{(\rho)} \equiv m_{\gamma_1\gamma_2}$  in equation 6.2). If there is only one photon,  $m_{\pi^0}^{(\rho)}$  is set to 0.

From the  $a_1$  invariant mass resonance reconstruction, three variables are obtained and used in the MVA classification: the  $a_1$  mass ( $m_{a_1}^{reco} \equiv m_{tot}$  in equation 6.3), the first

$\pi^0$  mass ( $m_{\pi^0}^{(a_1)} \equiv m_{\gamma_1\gamma_2}$  in equation 6.3), and the second  $\pi^0$  mass ( $m_{\pi^0}^{*(a_1)} \equiv m_{\gamma_3\gamma_4}$  in equation 6.3). If there are three reconstructed photons,  $m_{\pi^0}^{*(a_1)}$  is set to 0. If there are two reconstructed photons, depending on whether equation 6.5 or equation 6.6 gives the smallest  $\chi^2$  value, either  $m_{\pi^0}^{*(a_1)}$  is set to 0 or both  $m_{\pi^0}^{(a_1)}$  and  $m_{\pi^0}^{*(a_1)}$  are 0. If there is only one reconstructed photon, both  $m_{\pi^0}^{(a_1)}$  and  $m_{\pi^0}^{*(a_1)}$  are 0.

Figure 6.7 shows the distributions of  $m_{a_1}^{reco}$  under  $a_1(\pi^-\pi^0\pi^0)$  decay mode hypothesis test for three different tau decay modes. Only the distribution for  $a_1(\pi^-\pi^0\pi^0)$  decay mode has a resonance peak at  $a_1$  mass position.



**Figure 6.7:** Reconstructed invariant mass distributions for  $a_1$  ( $m_{a_1}^{reco}$ ), reconstructed under the  $a_1(\pi^-\pi^0\pi^0)$  decay mode hypothesis for three different tau decay modes. The area under the curve is normalised to unity.

#### 6.4.6 Separating electrons from charged pions

Variables are used in this analysis to help further separate electrons from charged pions, obtained from a modified private version of PandoraPFA.

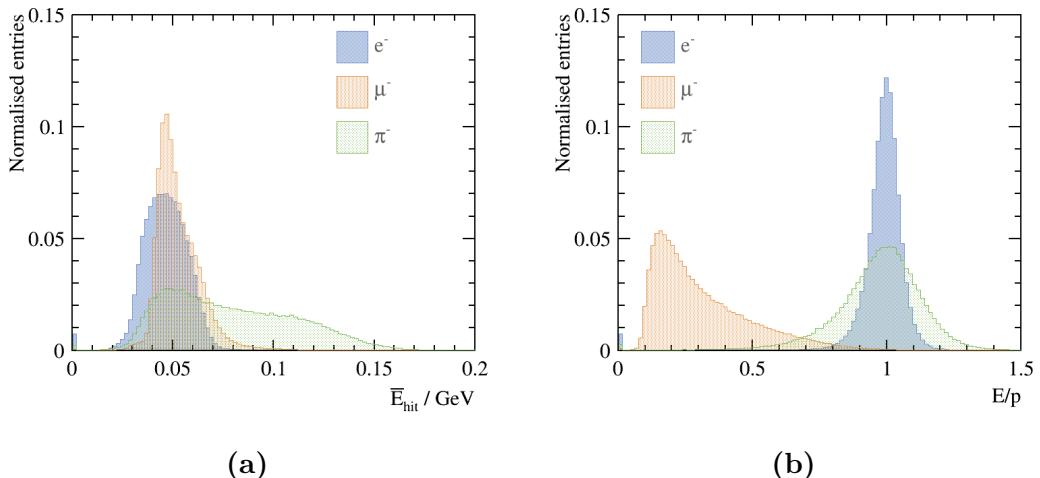
An electron develops a characteristic EM shower in the ECAL, whilst a charged pion develops a hadronic shower. Variables characterising the EM shower help to identify an electron. Three variables are used in the MVA classification: the start layer of the longitudinal shower ( $t_0$ ); the fractional difference between observed and expected longitudinal EM shower profile ( $\delta l$ ); and  $\langle w \rangle$ , a measure of the EM shower transverse width. These variables are defined in the same way as the variables used in the photon

likelihood classifier in the photon reconstruction in PandoraPFA, described in section 5.3.3.

The calorimeter hit information is also used to differentiate an EM shower from a hadronic shower. Two variables used in the MVA classification are: the average energy of a calorimeter hit ( $\bar{E}_{hit}$ ), which is the total energy deposited in the ECAL and HCAL divided by the number of the ECAL and HCAL calorimeter hits, and the average fraction of minimum ionising calorimeter hits ( $MIP$ ), which is the number of calorimeter hits in the ECAL and HCAL flagged as minimum ionising particles by the PandoraPFA reconstruction, divided by the total number of calorimeter hits in the ECAL and HCAL.

Finally track momentum is used to provide the consistency check of the track momentum with the total energies in the ECAL and HCAL for charged particles. The variable used in the MVA classification is the energy in the ECAL and HCAL divided by the track momentum, averaged over all charged particles ( $E/p$ ).

Figure 6.8 show distributions of the average energy of a calorimeter hit ( $\bar{E}_{hit}$ ), and the energy in the ECAL and HCAL divided by the track momentum, averaged over all charged particles ( $E/p$ ). Differences between the  $e^- \bar{\nu}_e \nu_\tau$ ,  $\mu^- \bar{\nu}_\mu \nu_\tau$ , and  $\pi^- \nu_\tau$  decay modes can be seen.



**Figure 6.8:** Distributions of: a) the average energy of a calorimeter hit ( $\bar{E}_{hit}$ ); and b) the energy in the ECAL and HCAL divided by the track momentum, averaged over all charged particles ( $E/p$ ). The area under the curve is normalised to unity.

## 6.5 Classification

The MULTICLASS class of the TMVA package [94] was used to perform a multiple-class classification, which classifies seven tau lepton decay final states simultaneously. The MULTICLASS classification is an extension of a standard two-class signal-background classification. The Boosted Decision Tree classifier with Gradient boost (BDTG) is used. Half of the events, randomly selected, were used in the training process and the other half were used for testing. The optimisation of the BDTG classifier followed the strategy outlined in section 4.5.1. The optimised parameters of the classifier are listed in table 6.4, where an explanation of the parameters can be found in section 4.5.6.1.

Parameter	Value
Depth of tree	5
Number of trees	3000
Boosting	gradient boost
Learning rate of the gradient boost	0.1
Metric for the optimal cuts	Gini Index
Bagging fraction	0.5
Number of bins per variables	100
End node output	yes/no

**Table 6.4:** Optimised parameters of the Boosted Decision Tree with Gradient boost MULTICLASS classifier used for the tau decay mode classification.

## 6.6 Tau decay mode classification efficiency

Two million  $e^+e^- \rightarrow \tau^+\tau^-$  events at a centre-of-mass energy of 100 GeV were used in the tau decay modes classification. For tau decays passing pre-selection cuts, the correct classification and misidentification efficiencies for the seven tau decay modes are shown in table 6.5. The correct classification efficiencies (bold numbers in table 6.5) are defined as:

$$\varepsilon_i = \frac{N_i^{correct}}{N_i^{MC}}, \quad (6.7)$$

where  $N_i^{correct}$  is the number of correctly classified tau decays for decay mode  $i$  and the  $N_i^{MC}$  is the total number of true tau decays for decay mode  $i$ .

Reco↓ Truth →	$e^-$	$\mu^-$	$\pi^-$	$\rho(\pi^-\pi^0)$	$a_1(\pi^-\pi^0\pi^0)$	$a_1(\pi^+\pi^-\pi^-)$	$\pi^+\pi^-\pi^-\pi^0$
$e^-$	<b>99.7%</b>	-	0.9%	0.6%	0.4%	-	-
$\mu^-$	-	<b>99.5%</b>	0.6%	-	-	-	-
$\pi^-$	-	0.3%	<b>94.0%</b>	0.8%	-	0.4%	-
$\rho(\pi^-\pi^0)$	-	-	3.4%	<b>93.6%</b>	9.5%	0.6%	2.3%
$a_1(\pi^-\pi^0\pi^0)$	-	-	-	4.5%	<b>89.7%</b>	-	0.6%
$a_1(\pi^+\pi^-\pi^-)$	-	-	0.9%	-	-	<b>96.8%</b>	6.4%
$\pi^+\pi^-\pi^-\pi^0$	-	-	-	0.3%	-	2.0%	<b>90.6%</b>

**Table 6.5:** Classification efficiencies for the seven tau decay modes considered here. Bold numbers represent the correct classification efficiencies. Boxes highlight one-prong and three-prong tau hadronic decay modes. The entries marked with “-” represent numbers below 0.25%. The absolute statistical uncertainty for each entry is less than 0.25%.

The particle ID from the PandoraPFA reconstruction is effective, resulting in the correct classification efficiencies for  $\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau$  and  $\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$  decays being 99.8% and 99.5% respectively. For the  $\tau^- \rightarrow \pi^- \nu_\tau$  decays, only 0.9% of decays are misclassified as  $\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau$  decays, due to the additional variables (section 6.4.6) dedicated to the separation between  $e^-$  and  $\pi^-$ .

For the separation of tau hadronic decay modes, photon reconstruction is important as the number of photons is an essential variable to distinguish different hadronic decay modes. Failure to reconstruct photons in the  $\tau^- \rightarrow a_1(\pi^-\pi^0\pi^0)\nu_\tau$  decays or extra reconstructed photons in the  $\tau^- \rightarrow \rho(\pi^-\pi^0)\nu_\tau$  decays leads to the misclassification between the  $\tau^- \rightarrow a_1(\pi^-\pi^0\pi^0)\nu_\tau$  and  $\tau^- \rightarrow \rho(\pi^-\pi^0)\nu_\tau$  decays. Similarly, failure to reconstruct photons in the  $\tau^- \rightarrow \rho(\pi^-\pi^0)\nu_\tau$  decays or extra reconstructed photons in the  $\tau^- \rightarrow \pi^- \nu_\tau$  decays can lead to the misclassification between the  $\tau^- \rightarrow \rho(\pi^-\pi^0)\nu_\tau$  and  $\tau^- \rightarrow \pi^- \nu_\tau$  decays. The misclassification between one-prong decays, as well as between three-prong decays, is highlighted in table 6.5.

A high correct classification rate is achieved for all seven classified tau decay modes. The leptonic decay modes have correct classification rates over 99.5%. For the hadronic tau decay modes, classification rates of 89.7% or above are achieved.

## 6.7 Electromagnetic calorimeter optimisation

The performance of photon reconstruction in a highly granular ECAL is an important metric for the ECAL performance. Since the classification of the tau hadronic decay modes

depends on the ability to reconstruct photons, tau hadronic decay modes classification is used as a metric to optimise the ECAL design. The tau decay mode classification was studied with ECAL square cell sizes of 3, 5, 7, 10, 15 and 20 mm, and at four centre-of-mass energies of 100, 200, 500, 1000 GeV. The other ECAL dimensions are kept the same as for the ILD nominal detector. The multivariate classifier was trained individually for each ECAL cell size and each centre-of-mass energy.

PandoraPFA was optimised for the nominal ILD detector. Therefore a re-optimisation is required for detector models with different ECAL cell sizes. In particular, the parameters used in `PHOTONFRAGMENTREMOVAL` algorithm need to be optimised for different ECAL cell sizes. The optimal `CLOSESTHITDISTANCE` parameter in `PHOTONFRAGMENTREMOVAL` algorithm, which is a distance metric controlling the merging of the fragment, was chosen by selecting the value that gives the highest overall tau hadronic decay classification rate,  $\varepsilon_{had}$  using  $e^+e^- \rightarrow \tau^+\tau^-$  samples at a centre-of-mass energy of 100 GeV. `CLOSESTHITDISTANCE` is varied amongst values of 5, 10, 20, 30, 40, and 50 mm. The overall tau hadronic decay correct classification rate,  $\varepsilon_{had}$ , is the weighted average correct classification efficiency, defined as:

$$\varepsilon_{had} = \frac{\sum_i^5 B_i \varepsilon_i}{\sum_i^5 B_i}, \quad (6.8)$$

where  $B_i$  is the branching fraction of the tau hadronic decay mode  $i$ ;  $\varepsilon_i$  is the correct classification efficiency of tau decay mode  $i$  (defined in equation 6.7); and the index  $i$  is summed over five tau hadronic decay modes considered here:  $\tau^- \rightarrow \pi^-\nu_\tau$ ;  $\tau^- \rightarrow \rho(\pi^-\pi^0)\nu_\tau$ ;  $\tau^- \rightarrow a_1(\pi^-\pi^0\pi^0)\nu_\tau$ ;  $\tau^- \rightarrow a_1(\pi^+\pi^-\pi^-)\nu_\tau$ ; and  $\tau^- \rightarrow \pi^+\pi^-\pi^-\pi^0\nu_\tau$ .

Table 6.6 shows the optimised values of `CLOSESTHITDISTANCE` parameter in `PHOTONFRAGMENTREMOVAL` algorithm as a function of the ECAL square cell sizes. As expected, for larger cell sizes, the distance metric for merging photons becomes larger.

Figure 6.9 shows  $\varepsilon_{had}$  as a function of ECAL cell sizes for four different centre-of-mass energies. The efficiency  $\varepsilon_{had}$  decreases with an increase of the centre-of-mass energy. As the centre-of-mass energy increases, the tau decay products become more boosted, making it increasingly difficult to separate tau decay products, for example, the photon pair from  $\pi^0$  decay. The reduction in the ability to separate photon pairs leads to a degradation of the classification performance.

ECAL cell length / mm	CLOSESTHITDISTANCE / mm
3	5
5	10
7	10
10	10
15	20
20	20

**Table 6.6:** Optimised values of CLOSESTHITDISTANCE parameters in PHOTONFRAGMENTREMOVAL algorithm as a function of the ECAL square cell sizes.

The efficiency  $\varepsilon_{had}$  decreases with the increasing ECAL cell sizes. The change in the ECAL cell size will change the ECAL transverse spatial resolution. Hence a large cell size will result in a low transverse spatial resolution, leading to a reduction in the ability to separate a pair of photons. Consequently, a worse classification performance is expected for a larger ECAL cell size.

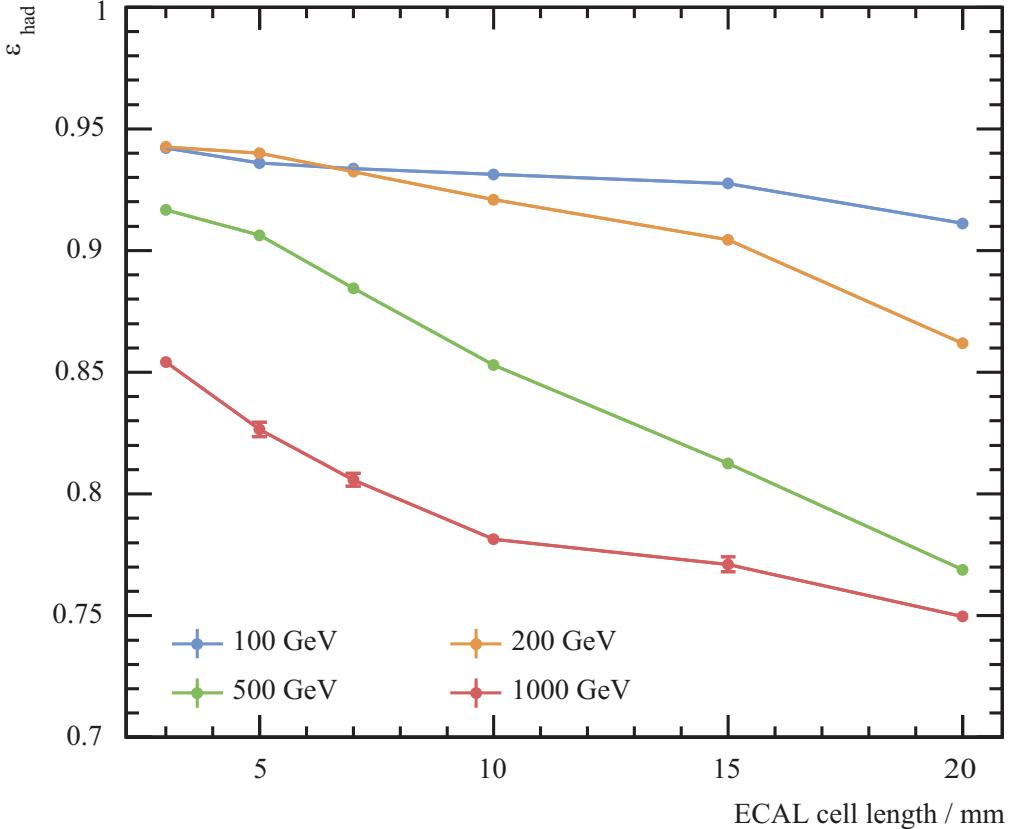
Table 6.7 lists the achieved hadronic decay mode separation, as measured by  $\varepsilon_{had}$ , with 3 mm and 20 mm ECAL cell lengths for four different centre-of-mass energies. The sensitivity of  $\varepsilon_{had}$  to different cell sizes is stronger at high centre-of-mass energies. With decay products being spatially close at high centre-of-mass energies, it is more beneficial to have a small ECAL cell size to reconstruct individual particles.

$\varepsilon_{had}$	3 mm	20 mm
100 GeV	94%	91%
200 GeV	94%	86%
500 GeV	92%	78%
1000 GeV	85%	75%

**Table 6.7:**  $\varepsilon_{had}$  with 3 mm and 20 mm ECAL cell lengths for four different centre-of-mass energies.

Figure 6.10 shows the correct classification efficiencies ( $\varepsilon_i$ ) for five tau hadronic decay modes as a function of the ECAL square cell sizes for four different centre-of-mass energies. For the ECAL square cells, the cell size is the squared of the cell length. The tau decay mode correct classification efficiencies generally decrease with an increase of centre-of-mass energies and an increase of ECAL cell sizes.

For the  $\tau^- \rightarrow \rho(\pi^-\pi^0)\nu_\tau$  decay mode, the efficiency at  $\sqrt{s} = 1000$  GeV increases as the cell size increases. This is because the multivariate classifier optimises for the overall

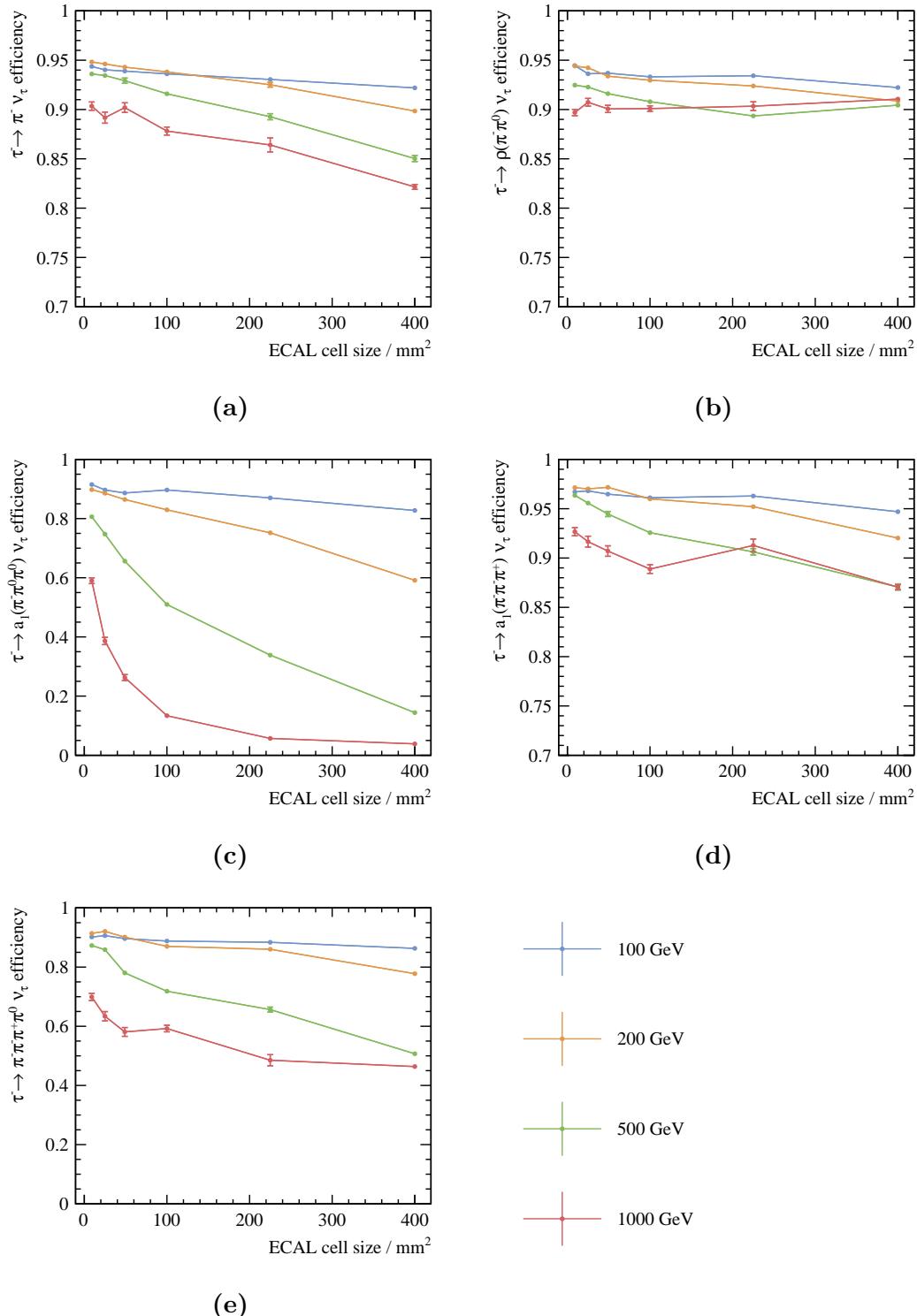


**Figure 6.9:** The weighted average tau hadronic decay correct classification efficiency,  $\epsilon_{had}$ , as a function of the ECAL cell sizes for four different centre-of-mass energies. The blue, orange, green, and red points show  $\epsilon_{had}$  at centre-of-mass energies of 100, 200, 500, and 1000 GeV, respectively.

classification efficiency, which may balance the decrease of the efficiency of one decay mode by the increase of the efficiency of another decay mode. In this case, the small increase in the efficiency for  $\tau^- \rightarrow \rho(\pi^-\pi^0)\nu_\tau$  decay mode at  $\sqrt{s} = 1000$  GeV is compensated by the drastic decrease in the efficiency for  $\tau^- \rightarrow a_1(\pi^-\pi^0\pi^0)\nu_\tau$  decay mode at the same centre-of-mass energy. For this reason,  $\epsilon_{had}$  gives a better picture of true performance.

For the  $\tau^- \rightarrow a_1(\pi^-\pi^0\pi^0)\nu_\tau$  decay mode, the loss of efficiency with an increasing ECAL cell size and an increasing centre-of-mass energy is most significant comparing to other decay modes. With more photons in the final state, it is the most challenging decay mode to reconstruct and thus most sensitive to the change in cell sizes and centre-of-mass energies.

For the  $\tau^- \rightarrow a_1(\pi^+\pi^-\pi^-)\nu_\tau$  decay mode, the efficiencies are similar to that of the  $\tau^- \rightarrow \pi^-\nu_\tau$  decay mode. Both final states contain charged particles only. Therefore it is most sensitive to the tracking detector performance, which is not affected by different ECAL cell sizes.



**Figure 6.10:** The correct classification efficiencies as a function of the ECAL square cell sizes for: a)  $\tau^- \rightarrow \pi^- \nu_\tau$  decays; b)  $\tau^- \rightarrow \rho(\pi^-\pi^0)\nu_\tau$  decays; c)  $\tau^- \rightarrow a_1(\pi^-\pi^0\pi^0)\nu_\tau$  decays; d)  $\tau^- \rightarrow a_1(\pi^+\pi^-\pi^-)\nu_\tau$  decays; and e)  $\tau^- \rightarrow \pi^+\pi^-\pi^-\pi^0\nu_\tau$  decays. Results are shown for centre-of-mass energies of 100, 200, 500, and 1000 GeV.

A previous study [95] on the tau decay mode classification was performed using the ILD detector on  $\pi^-\nu_\tau$ ,  $\rho\nu_\tau$ , and  $a_1(\pi^-\pi^0\pi^0)\nu_\tau$  decay modes. Samples used were  $e^+e^- \rightarrow Z^0 \rightarrow \tau^+\tau^-(\gamma)$  at 250 GeV. Pre-selection required that events with photon converted to electron pairs in the tracking detector were discarded. GARLIC [96, 97] photon reconstruction was used to reconstruct photons in the events. The main differences between the previous analysis and the current one are the pre-selection cuts, the photon reconstruction algorithm, and the number of classified decay modes. Table 6.8 lists the correct classification rates of  $\pi^-\nu_\tau$ ,  $\rho\nu_\tau$ , and  $a_1(\pi^-\pi^0\pi^0)\nu_\tau$  decay modes for current analysis with PandoraPFA photon reconstruction at a centre-of-mass energies of 200 GeV and the previous analysis using GARLIC photon reconstruction. Similar correct classification rates are achieved for  $\pi^-\nu_\tau$ ,  $\rho\nu_\tau$ , and  $a_1(\pi^-\pi^0\pi^0)\nu_\tau$  decay modes.

Decay mode	PandoraPFA $\sqrt{s} = 200$ GeV	GARLIC $\sqrt{s} = 250$ GeV
$\pi^-\nu_\tau$	$94.6\% \pm 0.1\%$	$96.8\% \pm 0.2\%$
$\rho\nu_\tau$	$94.2\% \pm 0.1\%$	$90.5\% \pm 0.2\%$
$a_1(\pi^-\pi^0\pi^0)\nu_\tau$	$88.6\% \pm 0.2\%$	$91.1\% \pm 0.4\%$

**Table 6.8:** Correct classification rates of  $\pi^-\nu_\tau$ ,  $\rho\nu_\tau$ , and  $a_1(\pi^-\pi^0\pi^0)\nu_\tau$  decay modes for current analysis with PandoraPFA photon reconstruction and previous analysis using GARLIC photon reconstruction. Values for the previous analysis are taken from [95].

## 6.8 Summary

For the ILC at  $\sqrt{s} = 250$  GeV or CLIC at  $\sqrt{s} = 350$  GeV, an ECAL size of 10 mm or fewer is sufficient to achieve a  $\varepsilon_{had}$  of 92%. For a linear collider operating at a centre-of-mass energy above a few hundred GeVs, such as the ILC at  $\sqrt{s} = 500$  GeV or CLIC at  $\sqrt{s} = 1.4$  TeV or 3 TeV, it is preferable to have a small ECAL cell size, i.e. 3 mm, for the best tau decay mode separation, as  $\varepsilon_{had}$  decreases drastically with an increasing ECAL cell size.

# Chapter 7

## Tau Pair Polarisation Correlation

*‘Where can I find a man who has forgotten words so I can have a word with him?’*

— Zhuang Zi, 369 BC – 286 BC

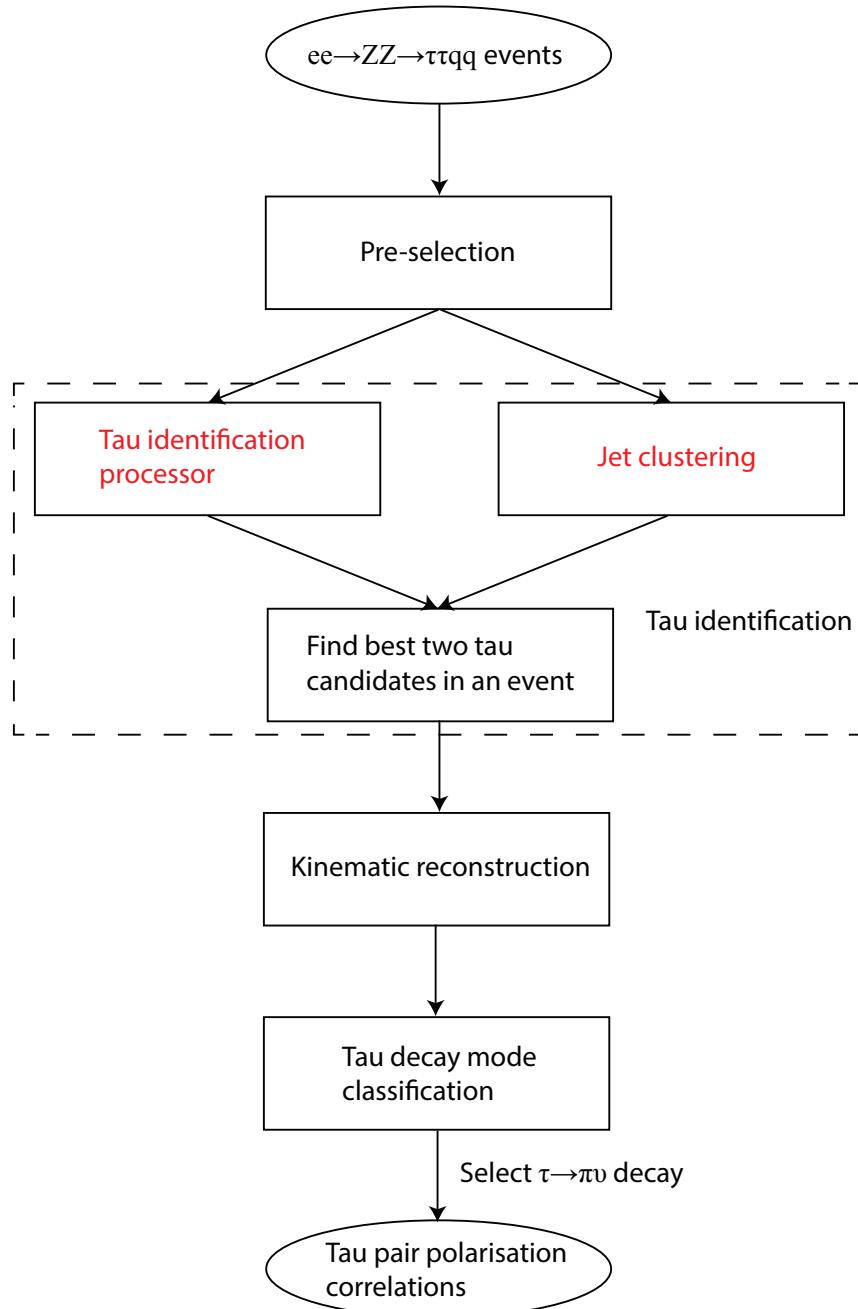
This chapter follows the theoretical discussion in section 2.9 on using the correlation between the polarisations of the tau pair from a boson decay as a signature to differentiate the Higgs boson from the Z boson.

A spin-0 scalar Higgs boson can decay to  $\tau_L^+ \tau_L^-$  or  $\tau_R^+ \tau_R^-$ , whereas the spin-1 Z boson can decay to  $\tau_L^+ \tau_R^-$  or  $\tau_R^+ \tau_L^-$ , where L, R denote the tau lepton helicities. Therefore, by studying the tau pair polarisation correlation from a boson decay, one can determine statistically if the parent boson is a scalar or a vector.

Here a proof-of-principle analysis is performed to reconstruct the polarisation correlation of the tau pairs in the  $Z \rightarrow \tau^+ \tau^-$  process, where both  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ .  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$  decay mode is selected using the tau decay mode classifier developed in chapter 6, which utilises the photon reconstruction in a highly granular calorimeter developed in chapter 5. The charged pion decay mode is chosen because the correlation between  $E_{\pi^+}/E_{\tau^+}$  and  $E_{\pi^-}/E_{\tau^-}$  is very different if the parent boson is a Z or a H boson, as suggested in figure 2.5.

The analysis starts with the event generation and simulation, followed by identifying the tau decay products in the events. Afterwards, the tau decay mode classification is used to identify  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$  decays. Lastly the tau pair polarisation correlation is presented and compared to the tau pair polarisation correlation obtained with generator-level Monte

Carlo particles. Figure 7.1 shows main steps in this proof-of-principle demonstration of the tau pair polarisation correlations using  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  events, where both  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ .



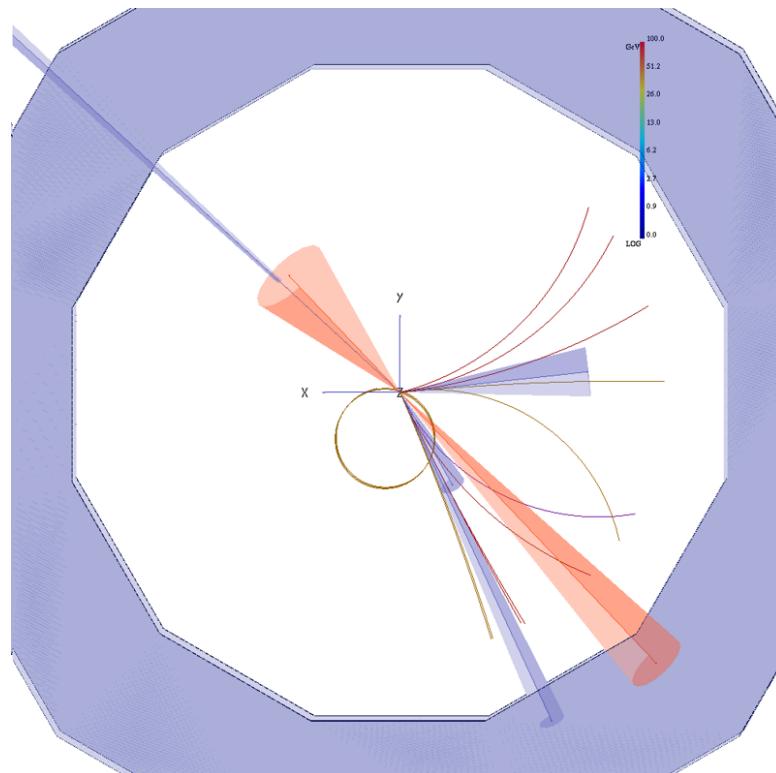
**Figure 7.1:** Main steps in the proof-of-principle demonstration of the tau pair polarisation correlations.

## 7.1 Event generation and simulation

For this proof-of-principle study,  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  events were generated at a centre-of-mass energy of 350 GeV, using WHIZARD [49] generator without ISR. TAUOLA [53] was used to describe the tau lepton decay with correct spin correlations of the tau decay products. Beam effects, such as the initial state radiation and the beam induced background, were not included. The  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  events were simulated using the ILD detector model as described in chapter 4.

## 7.2 Event reconstruction

Events were reconstructed with iLCSoft version v01-17-07 [92] and PandoraPFA version 3 [65], using the photon reconstruction algorithms described in chapter 5. Figure 7.2 shows an event display of a  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  event. The two brown cones indicate the tau decay products found by the tau identification processor. The four blue cones indicate the four jets found by the jet clustering algorithm.



**Figure 7.2:** An event display of a  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  event. The two brown cones indicate the tau decay products found by the tau identification processor. The four blue cones indicate the four jets found by the jet clustering algorithm. The blue outer region shows the HCAL ring.

Seven tau decay modes defined in section 6.1 were considered in this analysis:  $\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau$ ;  $\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$ ;  $\tau^- \rightarrow \pi^- \nu_\tau$ ;  $\tau^- \rightarrow \rho(\pi^-\pi^0)\nu_\tau$ ;  $\tau^- \rightarrow a_1(\pi^-\pi^0\pi^0)\nu_\tau$ ;  $\tau^- \rightarrow a_1(\pi^+\pi^-\pi^-)\nu_\tau$ ; and  $\tau^- \rightarrow \pi^+\pi^-\pi^-\pi^0\nu_\tau$ .

## 7.3 Pre-selection

Two pre-selection cuts defined in section 6.3 are used: the tau decay products with photon conversion to electron pairs in the tracking detector are not considered; and tau decays with the generated polar angle of the tau lepton in the region  $0.6 < |\theta_\tau^{MC}| < 0.9$  rad are not considered.

## 7.4 Tau identification

The  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  final state contains two tau leptons and two quark jets. Identifying the tau decay products in  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  events is challenging as a low multiplicity quark-jet could be topologically similar to a tau hadronic decay. Hence the tau decay product identification processor and the jet clustering algorithm are both used to find tau decay products. In the event shown in figure 7.2, particles associated with the tau decay products found by the tau identification processor are different to the particles associated with jets found by the jet clustering algorithm.

### 7.4.1 Tau identification processor

The ISOLATED TAU IDENTIFIER processor is a modified version of the tau decay product identification software used in the double Higgs analysis in section 8.3.2.2. The processor identifies high transverse momentum ( $p_T$ ) particles as tau seeds. Particles are iteratively added to a cone in the order of the ascending opening angle to the seed. The cone is called search cone, which contains potential tau decay products. After each particle addition, the temporary search cone is then considered as a temporary tau candidate and tested for isolation and consistency with a tau hadronic decay signature. The number of charged particles in the temporary tau candidate,  $N_{X+}$ , should be one or three. The invariant mass of the temporary tau candidate,  $m_c$ , should be less than 3 GeV. The temporary tau candidate also needs to pass the isolation condition to be identified as a tau candidate, which requires the opening angle between the temporary search cone and the 2<sup>nd</sup> closest charged particle,  $\theta_{X,+}^c$ , is larger than 0.6 rad.

The iterative particle addition procedure stops when the cone opening angle,  $\theta_S$ , is larger than  $\cos^{-1}(0.99)$ . If multiple temporary tau candidates of the same tau seed pass the isolation condition, the one with the smallest opening angle is chosen to form the final tau candidate. Table 7.1 lists the parameters of the ISOLATEDTAUIDENTIFER processor.

Modified ISOLATEDTAUIDENTIFER	Selection
Veto low $p_T$	$p_T < 0.5 \text{ GeV}$
Seed particle	$p_T > 1 \text{ GeV}$
Maximum search cone opening angle	$\theta_S \leq \cos^{-1}(0.99)$
Tau candidate rejection	$N_{X^+} \neq 1 \text{ or } 3; m_c > 3 \text{ GeV}$
Isolation	$\theta_{X'^+}^c > 0.6 \text{ rad}$

**Table 7.1:** Optimised parameters of the modified ISOLATEDTAUIDENTIFER.

The event is discarded if the ISOLATEDTAUIDENTIFER processor finds fewer than two tau candidates. If more than two tau candidates are found, the best two are selected by choosing the tau candidates that gives the smallest value of the  $\chi^2$  function, resulting in well reconstructed  $Z \rightarrow q\bar{q}$  decays:

$$\chi^2 = \left( m_{qq} - m_Z \right)^2 + \left( E_{qq} - \frac{\sqrt{s}}{2} \right)^2, \quad (7.1)$$

where  $\sqrt{s}$  is the centre-of-mass energy; variable  $m_Z$  is the mass of  $Z$  boson from reference [3];  $m_{qq}$  is the invariant mass of particles that do not belong to the two tau candidates; and  $E_{qq}$  is the total energy of particles that do not belong to the two tau candidates. This function is considered for all pairs of tau candidates. In the generated  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  events, the energy of each  $Z$  boson is half of the centre-of-mass energy. The invariant mass of two quarks from the  $Z$  decay should be close to  $Z$  mass.

## 7.4.2 Jet clustering

The Durham algorithm (see section 4.4.2.2) was run in the exclusive mode to force the reconstructed particles in  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  events into exactly four jets. The two tau candidate jets are identified by selecting two jets that gives the smallest value of the  $\chi^2$  function in equation 7.1. Here  $m_{qq}$  is the invariant mass of particles that are not in the two tau candidates jets and  $E_{qq}$  is the total energy of particles that are not in the two tau candidates jets. Other variables are defined in the same way as in section 7.4.1.

### 7.4.3 Selecting best tau candidates in an event

If both ISOLATEDTAUIDENTIFER processor and the jet clustering method find two tau candidates, the best pair of tau candidates should result in well reconstructed  $Z \rightarrow q\bar{q}$  decays, defined by:

$$\left| m_{qq} - m_Z \right| < 10 \text{ GeV}, \quad \left| E_{qq} - \frac{\sqrt{s}}{2} \right| < 10 \text{ GeV}. \quad (7.2)$$

The selection of best pair of tau candidates in an event proceeds as follows:

1. if pairs of tau candidates from both ISOLATEDTAUIDENTIFER processor and the jet clustering method satisfy equation 7.2, the pair with the smallest  $\chi^2$  from equation 7.1 is selected;
2. otherwise, the pair of tau candidates that satisfies equation 7.2 is selected;
3. otherwise, if one jet from the jet clustering is close to the beam pipe and there are exactly two tau candidates obtained from ISOLATEDTAUIDENTIFER, then the two tau candidates from ISOLATEDTAUIDENTIFER are selected. This choice is motivated by the fact that if one jet is close to the beam pipe, it is likely that some particles close to the beam pipe are undetected, which leads to a failure in the jet reconstruction;
4. otherwise, the two jets with the fewest number of particles are selected.

Table 7.2 lists the numbers of events with unmatched and matched taus identified in each of four steps in  $107 \text{ e}^+ \text{e}^- \rightarrow ZZ \rightarrow \tau^+ \tau^- qq$  events. An event with matched tau requires that MC particles contributing the most to identified best tau candidates are both true taus. The opposite is an event with unmatched taus. In 107 events, 93 events have matched taus and 14 events have unmatched taus.

## 7.5 Kinematic reconstruction of tau energy

The pion energy fractions,  $E_{\pi^+}/E_{\tau^+}$  and  $E_{\pi^-}/E_{\tau^-}$ , are the appropriate kinetic variables to illustrate the tau pair polarisation correlation in  $\text{e}^+ \text{e}^- \rightarrow ZZ \rightarrow \tau^+ \tau^- qq$  events, where both taus decay  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ , motivated in section 2.9. To obtain  $E_{\pi^+}/E_{\tau^+}$  and  $E_{\pi^-}/E_{\tau^-}$ , the energies of the taus,  $E_{\tau^+}$  and  $E_{\tau^-}$ , are required. Energies of the taus are also required for the calculation of the energy variables used in the tau decay mode classification.

Step	Unmatched taus	Matched taus
1	0	37
2	2	26
3	0	5
4	12	25
total	14	93

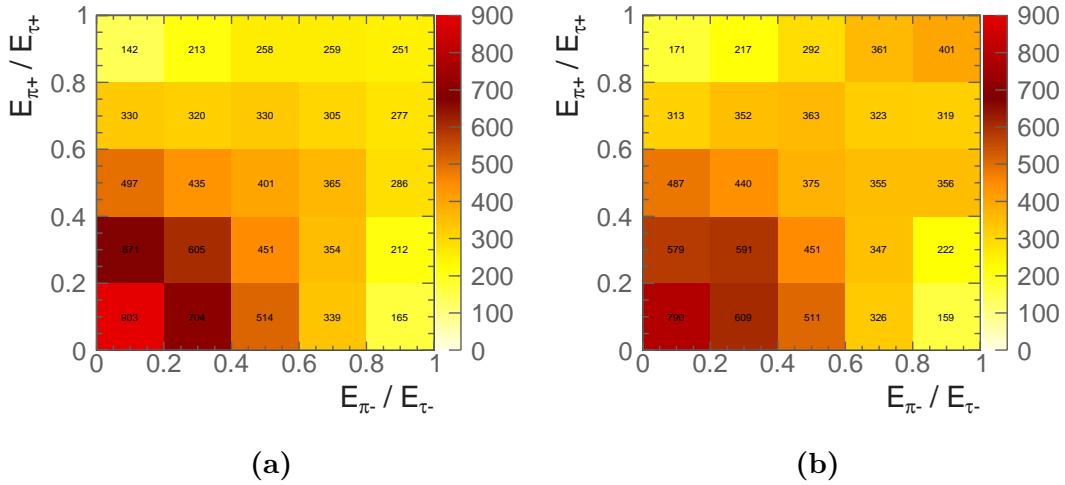
**Table 7.2:** Numbers of events with unmatched and matched taus identified in each of four steps to select best tau candidates in  $107 \text{ e}^+\text{e}^- \rightarrow \text{ZZ} \rightarrow \tau^+\tau^- \text{qq}$  events.

In the lab frame, the energies of the taus from  $\text{Z} \rightarrow \tau^+\tau^-$  can not be determined easily. However, in the Z rest frame, the energies of the taus can be determined, which are the half of the energy of the Z, i.e.  $E'_{\tau^+} = E'_{\tau^-} = \frac{1}{2}E'_Z$ .

Because the energies of the taus can be obtained in the Z rest frame, kinematic variables used in tau decay mode classification, for example  $\tilde{E}_C$ , are calculated in the Z rest frame as well. Therefore, tau decay products need to be boosted into the Z rest frame for the calculation of the kinematic variables.

To boost tau decay products into the Z rest frame, the four-momentum of the Z is needed. Since there are two Z in the  $\text{e}^+\text{e}^- \rightarrow \text{ZZ} \rightarrow \tau^+\tau^- \text{qq}$  event,  $Z_{\tau\tau}$  refers to the Z decaying to a tau pair and  $Z_{\text{qq}}$  refers to the Z decaying to a quark pair. The four-momentum of the  $Z_{\tau\tau}$  can be obtained from the recoil four-momentum of  $Z_{\text{qq}}$  against the centre-of-mass energy, where the four-momentum of  $Z_{\text{qq}}$  is measured directly from the particles that are not the tau candidates. A beam crossing angle of 14 mrad is considered.

The estimation of the four-momentum of  $Z_{\tau\tau}$  is improved by correcting the energy of the  $Z_{\tau\tau}$  to be half of the centre-of-mass energy, which improves the estimation of the tau energies. Figure 7.3 shows the two-dimensional distributions of  $E_{\pi^+}/E_{\tau^+}$  as a function of  $E_{\pi^-}/E_{\tau^-}$  from  $\text{Z} \rightarrow \tau^+\tau^-$  decays where both  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ .  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$  decay mode is determined using the truth information. True Monte Carlo tau decay particles are used to generate distributions shown in the figures. Figure 7.3a shows the distribution without the improved four-momentum of  $Z_{\tau\tau}$ . Figure 7.3b shows the distribution with the improved four-momentum of  $Z_{\tau\tau}$ . A better match with the distributions obtained in the generator level study in figure 2.5a is achieved with the improved four-momentum of  $Z_{\tau\tau}$ , which motivates the correction of the  $Z_{\tau\tau}$  energy to improve the measurement of the distributions of  $E_{\pi^+}/E_{\tau^+}$  and  $E_{\pi^-}/E_{\tau^-}$ .



**Figure 7.3:** Two-dimensional distributions of  $E_{\pi^+}/E_{\tau^+}$  as a function of  $E_{\pi^-}/E_{\tau^-}$  from  $Z \rightarrow \tau^+\tau^-$  decay where both  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ , with: a) the unimproved four-momentum of  $Z_{\tau\tau}$ ; and b) the improved four-momentum of  $Z_{\tau\tau}$ .  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$  decay mode is determined using the truth information. True Monte Carlo tau decay particles are used to generate distributions in figures.

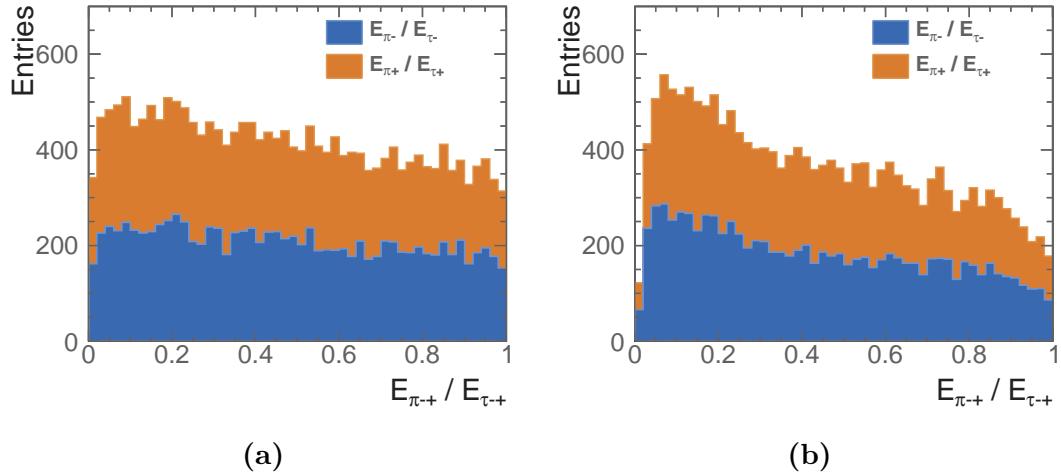
## 7.6 Tau decay mode classification

The tau decay mode classifier developed in chapter 6 can now be used to select the  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$  decay mode. In the classifier, variables regarding EM shower profiles, calorimeter hit information, and track information are not used (the last three rows in table 6.3) as the information was not available in the outputs of the standard version of PandoraPFA.

## 7.7 Tau pair polarisation correlations

Figure 7.4 shows the one-dimensional distributions of  $E_{\pi^+}/E_{\tau^+}$  and  $E_{\pi^-}/E_{\tau^-}$ , using generated  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  events. Both taus decay by  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ .  $E_{\tau^\pm}$  is half of the energy of  $Z_{\tau\tau}$  in  $Z_{\tau\tau}$  rest frame. At the generator level, the shape of the distribution decreases gradually with the increasing  $E_{\pi^\pm}/E_{\tau^\pm}$ . The decreasing shape is largely preserved in the full detector simulation. In the full detector simulation, events with  $E_{\pi^\pm}/E_{\tau^\pm}$  close to 0, which fall in the first bin in the figures, are not reconstructed correctly. When the energies of the tau are mostly carried away by the neutrinos, it is difficult to identify low-energy charged pions as tau decay products.

Figure 7.5 shows the two-dimensional distributions of  $E_{\pi^+}/E_{\tau^+}$  versus  $E_{\pi^-}/E_{\tau^-}$ , using generated  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  events. Both taus decay by  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ . Figure 7.5a shows the two-dimensional tau decay product energy fraction distribution obtained with

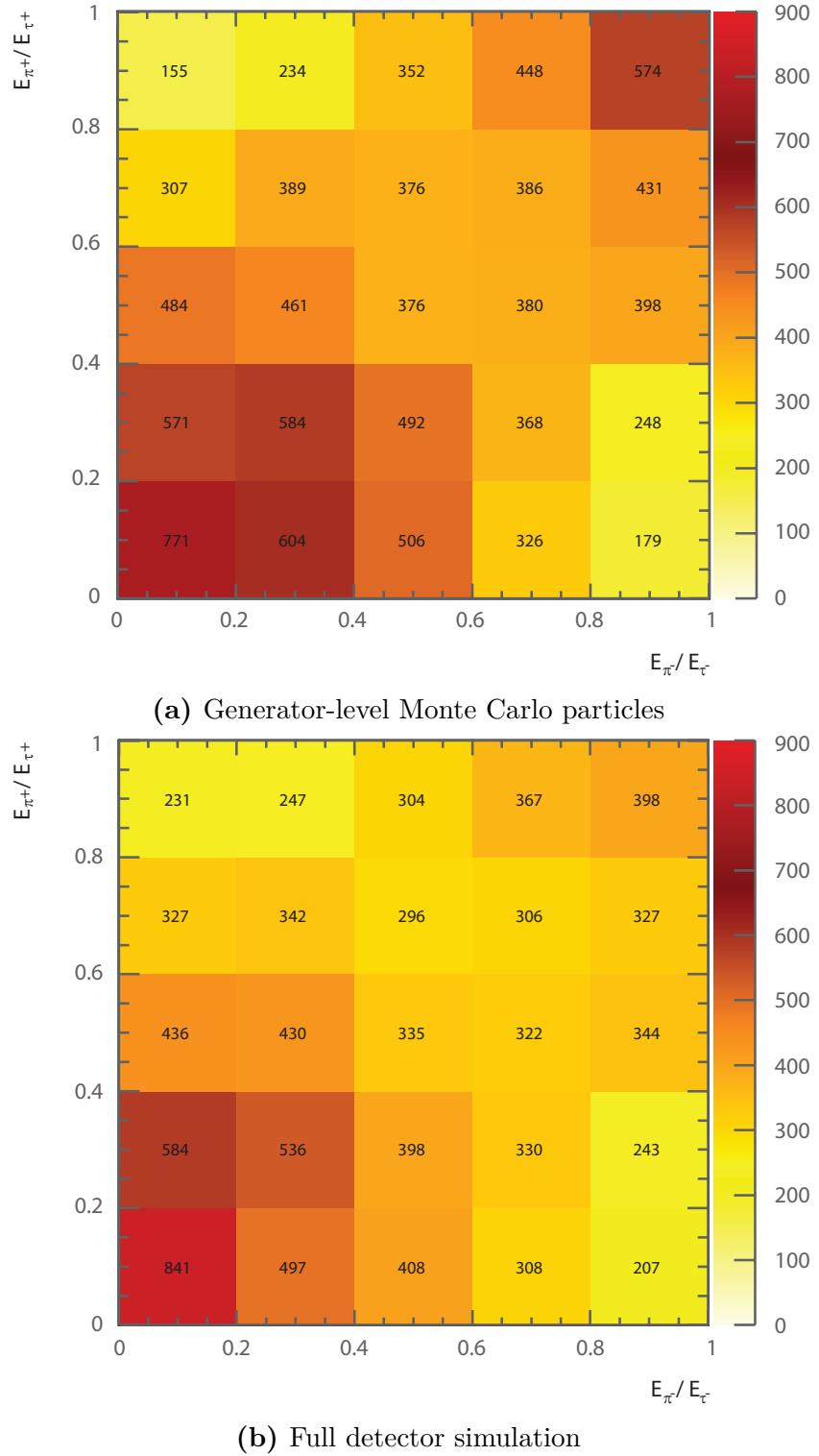


**Figure 7.4:** Distributions of  $E_{\pi^\pm}/E_{\tau^\pm}$  from  $Z \rightarrow \tau^+\tau^-$  decays where both  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ , in  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  events for: a) generator-level Monte Carlo particles, and b) the full detector simulation. The  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$  decay mode is selected using: a) the truth information; and b) the tau decay mode classifier.

the generator-level Monte Carlo particles. Figure 7.5b shows the distribution using the full detector simulation. A good match between the distributions obtained with the generator-level MC particles and the full detector simulation is achieved. Dark regions along the diagonal can be seen in both the distribution for the Monte Carlo particles and the distribution for the full detector simulation. In the  $Z \rightarrow \tau^+ \tau^-$  decays, where both  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ , an energetic  $\pi^\pm$  is likely to be associated with an energetic  $\pi^\mp$  and a low-energy  $\pi^\pm$  is likely to be associated with a low-energy  $\pi^\mp$ . Comparing the two figures, some events in the top right quadrant, corresponding to both  $\pi^\pm$  being energetic, are not reconstructed correctly in the full detector simulation, due to the failure of identifying the correct tau decay products.

## 7.8 Summary

This is a proof-of-principle demonstration that the generator-level pion energy fraction correlation can be reconstructed at the analysis level. The analysis contains several important steps: tau identification; kinematic reconstruction of the energies of the taus; the classification of the  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$  decay mode; and the reconstruction of the tau pair polarisation correlations discussed in section 2.9.



**Figure 7.5:** Two-dimensional distributions of  $E_{\pi^+}/E_{\tau^+}$  as a function of  $E_{\pi^-}/E_{\tau^-}$  from  $Z \rightarrow \tau^+\tau^-$  decays where both  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$ , in  $e^+e^- \rightarrow ZZ \rightarrow \tau^+\tau^-qq$  events for: a) generator-level Monte Carlo particles; and b) the full detector simulation. The  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$  decay mode is selected using: a) the truth information; and b) the tau decay mode classifier.

# Chapter 8

## Double Higgs Boson Production Analysis

*‘Life is really simple, but we insist on making it complicated.’*

— Confucius, 551 BC – 479 BC

Having discovered a Higgs-like particle the LHC in 2012 [10, 11], it became crucial to understand the interaction between the Higgs and other particles, and to determine whether it is the Standard Model Higgs. A number of Higgs theories beyond the Standard Model may be tested via the double Higgs production in an electron-positron collider [12, 13]. The study of double Higgs production would prevail the measurement of the Higgs trilinear self coupling,  $g_{\text{HHH}}$ , and the quartic coupling,  $g_{\text{WWHH}}$ . The precision for the measurement of  $g_{\text{HHH}}$  achievable by the Compact Linear Collider (CLIC), is superior to that at the LHC and the HL-LHC [21].

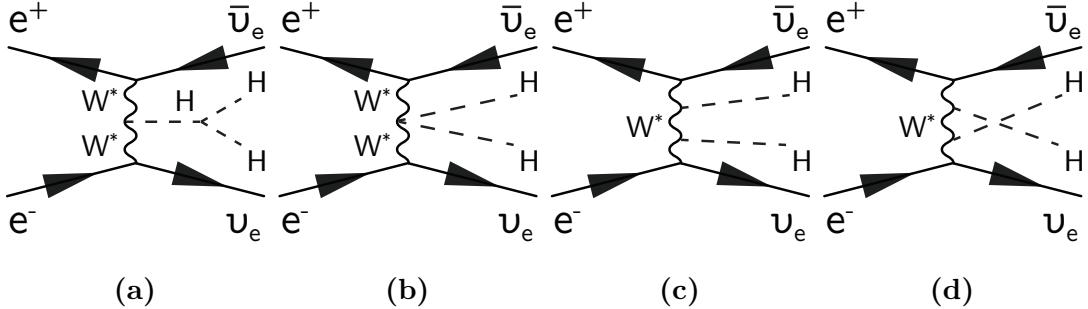
In  $e^+e^-$  collisions, there are two main challenges with the study of the double Higgs production,  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$ . Firstly, the process has a small cross section: 0.149 fb at  $\sqrt{s} = 1.4$  TeV and 0.588 fb at  $\sqrt{s} = 3$  TeV. The other challenge is that at high centre-of-mass energies, events are often boosted. Consequently, many final-state particles are in the forward region of the detector, where the reconstruction performance is inferior to the barrel region. In addition, particles can escape detection, causing a degradation in the event reconstruction performance.

In this chapter, a full CLIC\\_ILD detector simulation study has been performed for the double Higgs production process,  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$ , via  $W^+W^-$  fusion. Event generation

and simulation will be discussed first. An overview of the analysis, including lepton finding and jet reconstruction, is presented, followed by an optimised multivariate analysis to distinguish signal from background processes. The optimised event selection is used to derive an estimate of the uncertainty on  $g_{\text{HHH}}$  and  $g_{\text{WWHH}}$  measurements at CLIC. Part of this analysis has been published in [23].

## 8.1 Analysis strategy overview

The study of double Higgs production via  $W^+W^-$  fusion can probe the Higgs trilinear self coupling,  $g_{\text{HHH}}$ , and quartic coupling,  $g_{\text{WWHH}}$ . Leading-order Feynman diagrams for double Higgs production via  $W^+W^-$  fusion are shown in figure 8.1. The diagram shown in figure 8.1a contains the triple Higgs vertex, which is sensitive to the Higgs trilinear self coupling  $g_{\text{HHH}}$ . The diagram in the figure 8.1b is sensitive to the quartic coupling  $g_{\text{WWHH}}$ . Figures 8.1c and 8.1d show the Feynman diagrams for irreducible background processes in the study of  $g_{\text{HHH}}$  and  $g_{\text{WWHH}}$ .



**Figure 8.1:** The main Feynman diagrams for the leading-order  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  processes at CLIC.

Double Higgs production can also be produced via  $e^+e^- \rightarrow ZHH$ , where the  $Z$  decays to  $\nu\bar{\nu}$ . The  $ZHH$  process has been studied in  $e^+e^-$  collisions at  $\sqrt{s} = 500$  GeV [98]. However, for the CLIC energies of  $\sqrt{s} = 1.4$  TeV and 3 TeV, its contribution to the  $HH\nu\bar{\nu}$  final state is small compared to that of the  $W^+W^-$  fusion, and it can be neglected.

The two Higgs in the  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  decay to a range of particles. Hence double Higgs production has several distinct final-state topologies. The sub-channel with the largest cross section,  $HH \rightarrow b\bar{b}b\bar{b}$ , has been studied by collaborators at CERN. In this chapter, the  $HH \rightarrow b\bar{b}W^+W^-$  sub-channel is investigated. Firstly, the  $HH \rightarrow b\bar{b}W^+W^-$  sub-channel is studied for fully hadronic decays of the  $W^+W^-$ ; fully hadronic  $W^+W^-$  decays have the largest branching fraction and the lack of neutrinos in the final states allows each  $W$  to be reconstructed. The semi-leptonic final state of the  $W^+W^-$  system

in  $\text{HH} \rightarrow b\bar{b}W^+W^-$  is also studied. Here the presence of the neutrino in the final state makes it difficult to reconstruct the two Higgs bosons.

The process,  $\text{HH}\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e \rightarrow b\bar{b}qqqq\nu_e\bar{\nu}_e$ , results in a six quark final state with missing momentum. The high number of quarks requires an efficient jet reconstruction and a jet pairing algorithm to select the signal events. The two b quarks in the final state can be identified statistically with b jet tagging.

The chapter is organised as follows. Firstly, suitable signal and background processes are identified. Events with isolated high-energy leptons are discarded. Vertex information is used to identify b quark jets, in return to help to select signal events. The particles are clustered into jets and the jets are used as inputs for pre-selection and multivariate analysis.

The event reconstruction was first performed for  $\sqrt{s} = 1.4 \text{ TeV}$  and then  $\sqrt{s} = 3 \text{ TeV}$ , using the Marlin framework and reconstruction package in iLCSoft v01-16. More details on the reconstruction software can be found in chapter 4.

## 8.2 Monte Carlo sample generation

A full list of generated samples with their cross sections can be found in table 8.1. All samples were generated with the CLIC\_ILD detector model.

At high centre-of-mass energies, in addition to considering electron-electron interactions, electron-photon and photon-photon interactions are important as their interactions become significant. These photons are produced due to the high electric field generated by the colliding beams. Processes involving real photons from beamsstrahlung (BS) and “quasi-real” photons are generated separately. For the “quasi-real” photon initiated processes, the Equivalent Photon Approximation (EPA) has been used [99].

Background processes with multiple quarks and missing momentum in the final states are challenging to reject, as the topologies are similar to that of the signal events. Two example background processes are  $e^+e^- \rightarrow qqqq\nu\bar{\nu}$  and  $e^\pm\gamma \rightarrow \nu qqqq$ . For the same reason, single Higgs boson production, such as  $e^+e^- \rightarrow qqH\nu\bar{\nu}$ , has a similar final state to the signal events and is also difficult to reject.

Some processes are not considered in this analysis because they either have very different event topologies to the signal, or they have very small cross sections. For example,  $e^\pm\gamma \rightarrow qqH\ell$  is neglected as the cross section is very small, even at  $\sqrt{s} = 3 \text{ TeV}$ .

The background processes are generated according to the final states fermions and usually correspond to the contributions from multiple Feynman diagrams. These diagrams are already accounted for in the generated samples for explicit Higgs production. Therefore, to separate Higgs production from other processes, all background processes are generated with a Higgs boson mass of 14 TeV to ensure a negligible Higgs contribution. Processes involving Higgs production are simulated with a Higgs boson mass of 126 GeV.

The cross section of the signal,  $\text{HH} \rightarrow b\bar{b}W^+W^-$ , is scaled according to values listed in [100], as the values are more updated than the default Higgs branching ratios in the generator software.

The simulation and reconstruction chain is described in chapter 4. For some background processes, events are generated requiring that the invariant mass of the total momenta of all quarks is above 50 GeV or 120 GeV. This restricts the event generation to the region of phase space that could be populated by the signal processes.

Finally, the beam induced background,  $\gamma\gamma \rightarrow \text{hadrons}$ , is simulated and overlayed on all events. Details can be found in section 4.2.1.

## 8.3 Lepton identification

For the signal process,  $\text{HH} \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}qqqq$ , there is no primary lepton in the final state, whilst many background processes, such as  $qqqq\ell\nu$ , contain primary leptons in final states. Hence, efficiently rejecting events with primary leptons is an important step in the event selection. Primary leptons deposit energies in the tracking detector. The impact parameter to the interaction point of the fitted track of the primary lepton is typically small. At the same time, the primary leptons often have energies above 10 GeV and are isolated from other particles. High-energy electrons and muons are stable enough to deposit energies in the calorimeters. However, tau leptons are short lived with a typical decay lifetime of 290 fs [25]. They decay before reaching the vertex detector. Therefore, only the decay products of the tau leptons can be reconstructed.

### 8.3.1 Electron and muon identification

Two approaches to electron and muon identification were utilised, which are described below. The performance is summarised in table 8.6.

Process $\sqrt{s} = 1.4 \text{ TeV}$	$\sigma / \text{fb}$
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$	0.149
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-$ , hadronic	0.018
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}b\bar{b}\nu_e\bar{\nu}_e$	0.047
$e^+e^- \rightarrow HH \rightarrow \text{others}$	0.085
$e^+e^- \rightarrow q\bar{q}H\nu\bar{\nu}$	0.86
$e^+e^- \rightarrow c\bar{c}H\nu\bar{\nu}$	0.36
$e^+e^- \rightarrow b\bar{b}H\nu\bar{\nu}$	0.31
$e^+e^- \rightarrow qqqq$	1245.1
$e^+e^- \rightarrow qqqq\ell\ell$	62.1*
$e^+e^- \rightarrow qqqq\ell\nu$	110.4*
$e^+e^- \rightarrow qqqq\nu\bar{\nu}$	23.2*
$e^+e^- \rightarrow qq$	4009.5
$e^+e^- \rightarrow qq\ell\nu$	4309.7
$e^+e^- \rightarrow qq\ell\ell$	2725.8
$e^+e^- \rightarrow qq\nu\nu$	787.7
$e^\pm\gamma(\text{BS}) \rightarrow e^\pm qqqq$	2317
$e^\pm\gamma(\text{EPA}) \rightarrow e^\pm qqqq$	574
$e^\pm\gamma(\text{BS}) \rightarrow \nu qqqq$	159.1†
$e^\pm\gamma(\text{EPA}) \rightarrow \nu qqqq$	34.7†
$e^\pm\gamma(\text{BS}) \rightarrow qqH\nu$	31.5*
$e^\pm\gamma(\text{EPA}) \rightarrow qqH\nu$	6.78*
$\gamma(\text{BS})\gamma(\text{BS}) \rightarrow qqqq$	21406.2*
$\gamma(\text{BS})\gamma(\text{EPA}) \rightarrow qqqq$	4018.7*
$\gamma(\text{EPA})\gamma(\text{BS}) \rightarrow qqqq$	4034.8*
$\gamma(\text{EPA})\gamma(\text{EPA}) \rightarrow qqqq$	753.0*

**Table 8.1:** List of signal and background samples used in the double Higgs analysis with the corresponding cross sections at  $\sqrt{s} = 1.4 \text{ TeV}$ .  $q$  can be  $u$ ,  $d$ ,  $s$ ,  $b$  or  $t$ . Unless specified,  $q$ ,  $\ell$  and  $\nu$  represent either particles or the corresponding anti-particles.  $\gamma$  (BS) represents a real photon from beamstrahlung (BS).  $\gamma$  (EPA) represents a “quasi-real” photon, simulated with the Equivalent Photon Approximation. For processes labeled with \* and †, events are generated with the invariant mass of the total momenta of all quarks above 50 and 120 GeV, respectively.

### 8.3.1.1 IsolatedLeptonFinder

An optimised version of the existing ISOLATEDLEPTONFINDER reconstruction package is used. This algorithm identifies high energy electrons and muons that are isolated from other particles. The algorithm parameters were optimised by the collaborator using the  $\text{HH} \rightarrow b\bar{b}b\bar{b}$  as the signal process and the  $e^+e^- \rightarrow qqqq\ell\nu$  as the background process, as the background processes are the same for this analysis with  $\text{HH} \rightarrow b\bar{b}W^+W^-$  process.

Optimal values of the parameters of the ISOLATEDLEPTONFINDER are listed in table 8.2:  $E$  is the energy of the lepton;  $E_{ECAL}$  is the energy of the lepton deposited in the ECAL;  $E_{cone}$  is the total energy within a cone of an opening angle of  $\cos^{-1}(0.995)$  around the lepton; and the impact parameters,  $d_0$ ,  $z_0$ , and  $r_0$  are the closest Euclidean distance of the fitted track of the primary lepton to the interaction point in  $x$ - $y$  plane, in  $z$  direction, and in  $x$ - $y$ - $z$  three dimensional space, respectively.

ISOLATEDLEPTONFINDER	Selection
High Energy	$E > 15 \text{ GeV}$
$e^\pm$ ID	$\frac{E_{ECAL}}{E} > 0.9$
$\mu^\pm$ ID	$0.25 > \frac{E_{ECAL}}{E} > 0.05$
Primary Track	$d_0 < 0.02 \text{ mm}; z_0 < 0.03 \text{ mm}; r_0 < 0.04 \text{ mm}$
Isolation	$E_{cone}^2 \leqslant 5.7 \text{ GeV} \times E - 50 \text{ GeV}^2$

**Table 8.2:** Optimised parameters of the ISOLATEDLEPTONFINDER processor.

### 8.3.1.2 IsolatedLeptonIdentifier

A complimentary electron finder, ISOLATEDLEPTONIDENTIFIER, was developed to further identify isolated electrons and muons. Compared to the ISOLATEDLEPTONFINDER, the main difference is that the ISOLATEDLEPTONIDENTIFIER utilises particle ID information provided by the PandoraPFA reconstruction to identify leptons.

Table 8.3 lists the selection cuts for ISOLATEDLEPTONIDENTIFIER. The variables in the ISOLATEDLEPTONFINDER and the ISOLATEDLEPTONIDENTIFIER are defined in the same way. In addition:  $p_T$  is the transverse momentum;  $E_{cone1}$  and  $E_{cone2}$  are the total energy of PFOs within a cone around the lepton of an opening angle of  $\cos^{-1}(0.995)$  and  $\cos^{-1}(0.99)$  respectively.

The algorithm uses two sets of cuts to identify isolated leptons. If a PFO passes either set of cuts, it will be identified by the processor. The first set of cuts uses

the particle ID information from PandoraPFA, demanding a PandoraPFA electron or muon with high energy above 10 GeV and  $r_0 < 0.015$  mm. Afterwards, the lepton should either have  $p_T > 40$  GeV, or  $E \geq 23 \text{ GeV}^{\frac{1}{2}} \times \sqrt{E_{\text{cone}1}} + 5$  GeV. Figure 8.2a and 8.2b show the distributions of the  $p_T$  of identified electrons after  $E$  and  $r_0$  cuts, for  $\text{HH} \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}\text{qqqq}$  signal process and  $e^+e^- \rightarrow \text{qqqq}\ell\nu$  background process respectively. A cut of  $p_T > 40$  GeV preserves most signal events. Figure 8.2c and 8.2d show the distributions of  $23 \text{ GeV}^{\frac{1}{2}} \times \sqrt{E_{\text{cone}1}} + 5$  GeV as a function of  $E$  of identified electrons after  $E$  and  $r_0$  cuts, for  $\text{HH} \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}\text{qqqq}$  signal process and  $e^+e^- \rightarrow \text{qqqq}\ell\nu$  background process respectively. A cut along the two-dimensional histogram would discard background events and leave signal events intact.

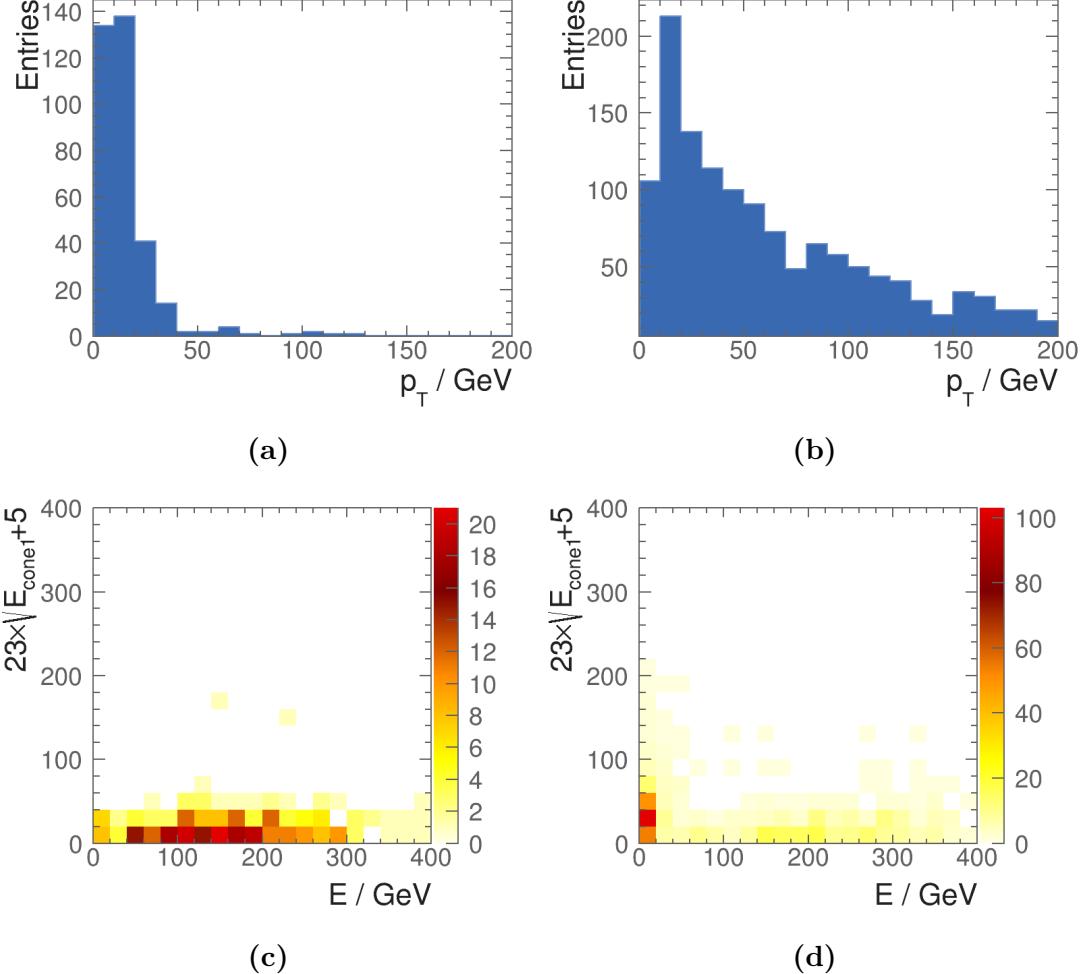
The second set of cuts is similar to the first set of cuts. Apart of the differences in the values of the cuts, lepton ID in the second set of cuts is determined using the fraction of the energy deposited in the ECAL as a function of the total energy,  $\frac{E_{\text{ECAL}}}{E}$ : if  $\frac{E_{\text{ECAL}}}{E} > 0.95$  then the PFO is an electron; and if  $0.2 > \frac{E_{\text{ECAL}}}{E} > 0.05$  then the PFO is a muon.

ISOLATEDLEPTONIDENTIFIER	Selection
High Energy	$E > 10$ GeV
$e^\pm$ ID	PandoraPFA reconstructed; $\frac{E_{\text{ECAL}}}{E} > 0.95$
$\mu^\pm$ ID	PandoraPFA reconstructed
Primary Track	$r_0 < 0.015$ mm
a) High Transverse Momentum, or	$p_T > 40$ GeV
b) Isolation	$E \geq 23 \text{ GeV}^{\frac{1}{2}} \times \sqrt{E_{\text{cone}1}} + 5$ GeV
High Energy	$E > 10$ GeV
$e^\pm$ ID	$\frac{E_{\text{ECAL}}}{E} > 0.95$
$\mu^\pm$ ID	$0.2 > \frac{E_{\text{ECAL}}}{E} > 0.05$
Primary Track	$r_0 < 0.5$ mm
a) High Transverse Momentum, or	$p_T > 40$ GeV
b) Isolation	$E \geq 28 \text{ GeV}^{\frac{1}{2}} \times \sqrt{E_{\text{cone}2}} + 30$ GeV

**Table 8.3:** Optimised parameters of the ISOLATEDLEPTONIDENTIFIER processor. A PFO needs to pass either set of cuts to be identified as a isolated electron or muon. Within a set of cuts, the PFO needs to satisfy either condition a) or b).

### 8.3.2 Tau lepton identification

The tau lepton has a short lifetime and decays before reaching the vertex detector and can only be identified through the reconstruction of its decay products. The leptonic



**Figure 8.2:** Distributions shown for  $p_T$  of identified electrons after  $E$  and  $r_0$  cuts, for: a)  $\text{HH} \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}\text{qqqq}$  signal process; and b)  $e^+e^- \rightarrow \text{qqqq}\ell\nu$  background process. Distributions shown for  $23\text{ GeV}^{\frac{1}{2}} \times \sqrt{E_{cone1}} + 5$  GeV as a function of  $E$  after  $E$  and  $r_0$  cuts, for: c)  $\text{HH} \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}\text{qqqq}$  signal process; and d)  $e^+e^- \rightarrow \text{qqqq}\ell\nu$  background process.

decay of tau lepton can be identified using the isolated lepton finder processors described above. Therefore in this section, tau identification will focus on the hadronic decay modes.

The existing TAU<sup>FINDER</sup> [101] reconstruction package has been optimised. In addition, a package, ISOLATED TAU IDENTIFIER, was developed to provide additional tau lepton identification.

### 8.3.2.1 TauFinder

The TAUFINDER works by identifying tau lepton decay products, and requiring the decay products to be isolated from other PFOs. To find the decay products, the algorithm starts with the highest energy track as a seed for the cone clustering algorithm. A cone with opening angle 0.03 rad with respect to the seed is formed. The PFOs within the cone are required to be consistent with the signature of a tau hadronic decay: no more than 3 charged particles in the cone; invariant mass of all PFOs in the cone less than 2 GeV; and few than 10 PFOs in the cone. The cone is also required to be isolated from other particles. To reduce fake rate, PFOs with low momentum (less than 1 GeV) are not used in tau finding, as they more likely come from  $\gamma\gamma \rightarrow$  hadrons background. The identified tau lepton and associated decay products are then not used in further tau finding. This tau lepton finding procedure iterates with other high-energy tracks as seeds.

The optimised parameters are listed in table 8.4. The optimisation is performed by the collaborator using  $\text{HH} \rightarrow b\bar{b}b\bar{b}$  signal process and the  $e^+e^- \rightarrow qqqq\ell\nu$  background process, by scanning the parameters to obtain a good background rejection rate with lowest signal rejection rate. Variables are defined in the same way as in previous sections. In addition:  $\theta_Z$  is the polar angle with respect to the beam axis;  $N_{X^+}$  and  $N_\tau$  are the number of charged particles and the number of PFOs in the tau cone respectively;  $m_\tau$  is the invariant mass of the sum of the PFOs in the tau candidate; and  $E_{cone}$  is the total energy of PFOs within a cone of an opening angle between 0.03 and 0.33 rad around tau seed track.

TAUFINDER	Selection
Veto $\gamma\gamma \rightarrow$ hadrons	$p_T < 1 \text{ GeV}$
Seed particle	$p_T > 10 \text{ GeV}$
Tau candidate cone opening angle	0.03 rad
Tau candidate rejection	$N_{X^+} > 3; N_\tau > 10; m_\tau > 2 \text{ GeV}$
Isolation	$E_{cone} < 3 \text{ GeV}$

**Table 8.4:** Optimised parameters of the TAUFINDER processor.

### 8.3.2.2 IsolatedTauIdentifier

The ISOLATEDTAUIDENTIFIER works in a similar way to the TAUFINDER. It identifies high momentum particles as tau seeds. Particles are iteratively added to a cone in the order of the ascending opening angle to the seed. The cone is called search cone,

which contains potential tau decay products. After each particle addition, the temporary search cone is then considered as a temporary tau candidate and tested for isolation and consistency with a tau hadronic decay signature. The temporary tau candidate only needs to pass one of the isolation conditions to be identified as a tau candidate. There are multiple isolation conditions for tau 1-prong decay and 3-prong decay, reflecting different topologies of tau decay final states. The isolation criteria typically demand few particles around the search cone and the total  $p_T$  in the search cone to be greater than a threshold.

The iterative particle addition procedure stops when the cone opening angle is larger than a threshold. If multiple temporary tau candidates of the same tau seed pass the selection, the one with smallest opening angle is chosen to form the final tau candidate. To reduce the fake rate from  $\gamma\gamma \rightarrow \text{hadrons}$  background, particles with energies less than 1 GeV are not considered.

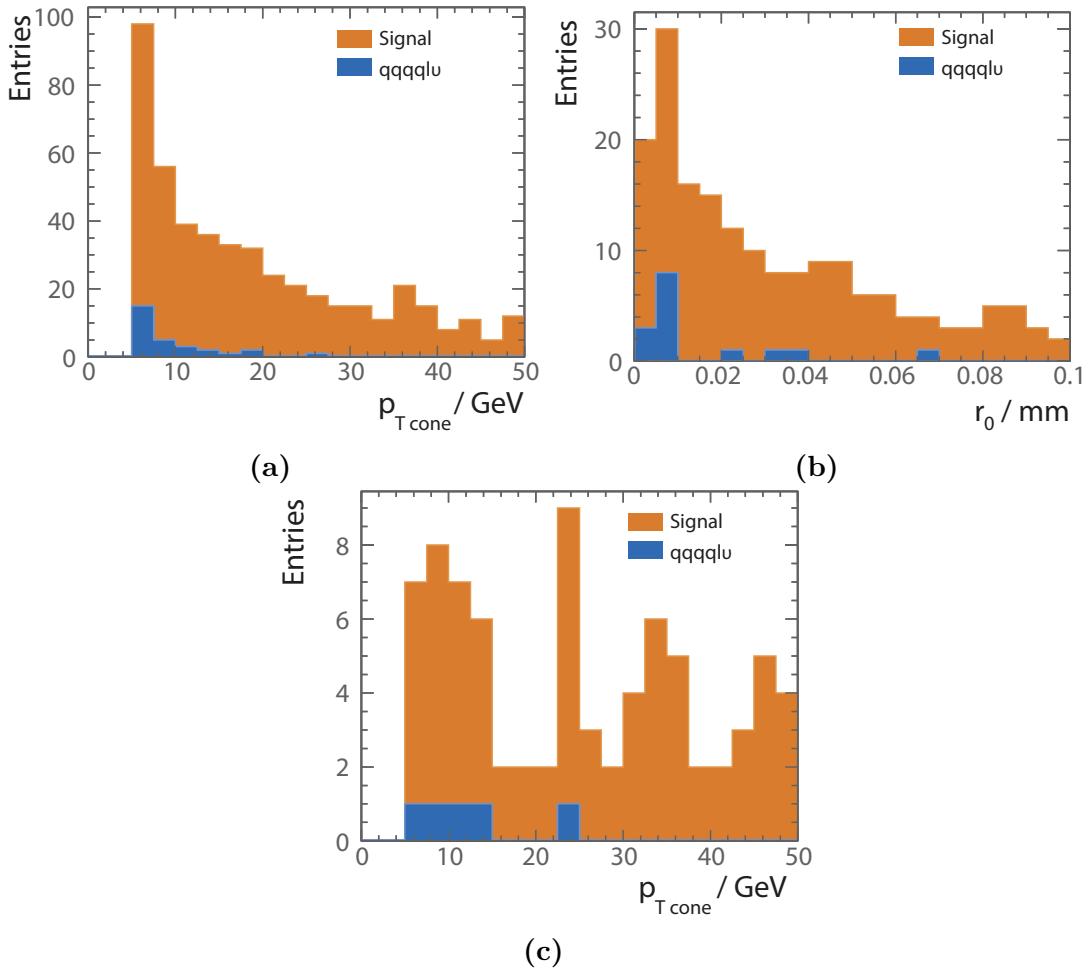
Table 8.4 lists the optimised parameters for ISOLATEDTAUIDENTIFIER. Variables are defined in the same way as those in previous sections. In addition,  $\theta_S$  is the opening angle of the search cone in rad;  $cone1$  and  $cone2$  are defined as a cone around the tau seed of an opening angle of  $\cos^{-1}(0.95)$ , and  $\cos^{-1}(0.99)$  respectively.

ISOLATEDTAUIDENTIFIER	Selection
Veto $\gamma\gamma \rightarrow \text{hadrons}$	$E < 1 \text{ GeV}$
Seed particle	$p_T > 5 \text{ GeV}$
Maximum search cone opening angle	$\theta_S \leq \cos^{-1}(0.999) \text{ GeV}$
Tau candidate rejection	$N_{X^+} \neq 1, 3; m_{PFO} > 3 \text{ GeV}$
Isolation 1 or	$N_{cone1} = 0; p_{Tcone} \geq 10 \text{ GeV}$
Isolation 2 or	$N_{X^+} = 1; N_{cone1} = 1; r_0 > 0.01 \text{ mm}$
Isolation 3 or	$N_{X^+} = 3; N_{cone1} = 1; p_{Tcone} \geq 10 \text{ GeV}; \theta_S < \cos^{-1}(0.9995)$
Isolation 4 or	$N_{X^+} = 1; N_{cone2} = 0; r_0 > 0.01 \text{ mm}; p_{Tcone} \geq 10 \text{ GeV}$
Isolation 5	$N_{X^+} = 3; N_{cone2} = 0; p_{Tcone} \geq 10 \text{ GeV}; \theta_S < \cos^{-1}(0.9995)$

**Table 8.5:** Optimised parameters of ISOLATEDTAUIDENTIFIER processor

Figure 8.3 shows the distributions of variables used in isolation criterion of tau candidate for  $\text{HH} \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}\text{qqqq}$  signal process (blue) and  $e^+e^- \rightarrow \text{qqqq}\ell\nu$  background process (orange). Figure 8.3a shows the distribution of the transverse momentum of the particles in the search cone, after selecting  $N_{cone1} = 0$ , used in the isolation criterion

1. The cut at  $p_{Tcone} \geq 10$  GeV selects more tau candidates in background events than in the signal events, where there should be no true high-energy isolated tau leptons in signal events. Shown in figure 8.3b, after selecting  $N_{X+} = 1$  and  $N_{cone1} = 1$ , the cut at  $r_0 > 0.01$  mm used in the isolation criterion 2 selects more true tau candidates in background events. Figure 8.3a shows the distribution of the transverse momentum of the particles in the search cone, after requiring  $N_{X+} = 3$ ,  $N_{cone1} = 1$  and  $\theta_S < \cos^{-1}(0.9995)$  for isolation criteria 3. The cut at  $p_{Tcone} \geq 10$  GeV selects tau candidates in background events.



**Figure 8.3:** Distributions to show isolation criteria: a)  $p_{Tcone}$  for isolation criterion 1 after selecting  $N_{cone1} = 0$ ; b)  $r_0$  for isolation criterion 2 after selecting  $N_{X+} = 1$ ,  $N_{cone1} = 1$ ; and c)  $p_{Tcone}$  for isolation criteria 3 after selecting  $N_{X+} = 3$ ,  $N_{cone1} = 1$ ,  $\theta_S < \cos^{-1}(0.9995)$ . Particle ID is from the truth information. Distributions are shown for tau candidates in  $\text{HH} \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}\text{qqqq}$  signal process (blue) and  $e^+e^- \rightarrow \text{qqql}\ell\nu$  background process (orange).

Relative to the TAUFINDER algorithm, the main difference is that the ISOLATEDTAUIDENTIFER adopts an iterative approach to build up a tau candidate, which allows a dynamic tau search cone size.

### 8.3.3 Very forward electron identification

At the high centre-of-mass energy of CLIC, particles produced are often highly boosted. Because of this, it is important to identify leptons in the forward calorimeters to aid the signal selection. In particular, photon–electron interactions can have energetic primary electrons in the forward calorimeters, the LumiCAL and/or the BeamCAL.

Because of the large background in the forward region, it is challenging to identify primary leptons. In the Monte Carlo production, particles including the primary leptons and beam induced background in the forward calorimeters are not simulated, due to the high demand on the computational resources. Instead, studies have been performed with particles simulated in the forward calorimeters to understand the primary lepton identification efficiencies [48, 102, 103]. The studied primary lepton identification efficiencies are then parameterised as a function of lepton energies. The parametrisation approach is adopted in this analysis.

Figure 8.4a shows the primary electron identification efficiencies in the BeamCAL as a function of polar angle for a 500 GeV electron. An external processor [102] has been developed to parameterise the primary electron identification efficiencies in the BeamCAL at  $\sqrt{s} = 3 \text{ TeV}$  as a function of electron energy and the polar angle. The full simulation study to obtain the primary electron identification efficiencies in the BeamCAL assumes a background integrated over 40 bunch crossings. The same primary electron identification efficiency is assumed for  $\sqrt{s} = 1.4 \text{ TeV}$  and  $\sqrt{s} = 3 \text{ TeV}$ . In the analysis for  $\sqrt{s} = 1.4 \text{ TeV}$ , the momenta of the electron is scaled down by a ratio of the centre-of-mass energies to use the external processor.

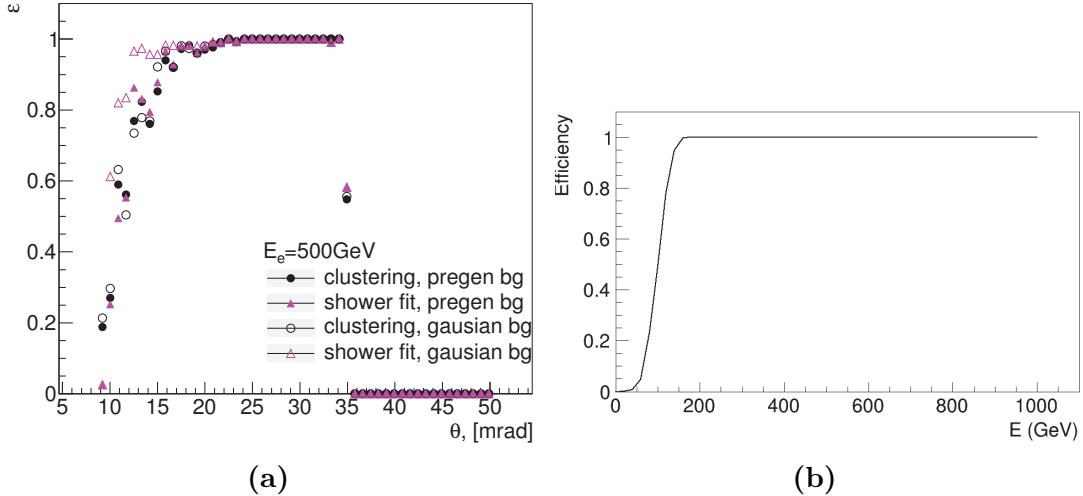
Figure 8.4b shows the primary electron identification efficiencies in the LumiCAL as a function of electron energy for a polar angle  $\theta = 50 \text{ mrad}$ . The efficiency is obtained from a full simulation study [103], assuming a background integrated over 100 bunch crossings. In this analysis, the primary electron identification efficiency as a function of electron energy is assumed to be parameterised by the curve in figure 8.4b. The polar angle dependency of the efficiency is not considered, due to the lack of study. The primary

electron identification efficiency curve in figure 8.4b takes the functional form of:

$$\varepsilon = \begin{cases} 0, & \text{if } E < 50 \text{ GeV}, \\ 0.99 \times \frac{\text{erf}(E/\text{GeV}-100)+1}{2}, & \text{otherwise,} \end{cases} \quad (8.1)$$

where  $E$  is the energy of the electron and  $\text{erf}$  is the error function.

Due to a lack of tracking detector in the very forward region, electrons and photons can not be differentiated. Therefore, both photons and electrons are identified in the forward calorimeters. Events with identified high-energy electrons and/or photons in the BeamCAL and/or LumiCAL are rejected.



**Figure 8.4:** a) 500 GeV electron identification efficiency in the BeamCAL as a function of polar angles, with different methods to model backgrounds: pregenerated and Gaussian, and two methods to identify electrons: clustering algorithm and shower fitting algorithm, obtained from a full simulation study in [102]. b) electron tagging efficiency in the LumiCAL as a function of the electron energy, for a polar angle  $\theta = 50$  mrad, obtained from a full simulation study in [103].

### 8.3.4 Lepton identification performance

The performances of the different lepton finding processors for signal events and the selected background processes are shown in table 8.6 for  $\sqrt{s} = 1.4$  TeV. Numbers in the table represent the fractions of events where no leptons are identified by the individual lepton finder. ISOLATEDLEPTONIDENTIFIER and ISOLATEDTAUIDENTIFIER reject more background events than the ISOLATEDLEPTONFINDER and TAUFINDER. By combining the processors, 86.6% of the signal events remain and 16.8% of the  $e^+e^- \rightarrow qqqq\ell\nu$  events survive after rejecting events where leptons are identified.

The forward lepton finders are most effective at rejecting background events with primary leptons in the forward region. Only 1% of signal events are rejected, but 47.4% of the  $e^-\gamma(BS) \rightarrow e^-qqqq$  background events are rejected.

Efficiency (1.4 TeV)	Signal	$e^+e^- \rightarrow qqqq\ell\nu$	$e^-\gamma(BS) \rightarrow e^-qqqq$
ISOLATEDLEPTONFINDER	99.3%	50.3%	87.3%
ISOLATEDLEPTONIDENTIFER	99.1%	39.9%	83.7%
TAUFINDER	97.5%	52.3%	90.4%
ISOLATEDTAUIDENTIFER	89.7%	38.5%	78.5%
Forward Finder Processors	98.9%	95.1%	53.6%
Combined	86.6%	16.8%	30.8%

**Table 8.6:** The performances of the lepton finding algorithms for the signal events and selected background events at  $\sqrt{s} = 1.4$  TeV.  $\gamma$  (BS) represents a real photon from beamstrahlung (BS). Numbers represent the fractions of events where no leptons are identified by the individual lepton finder.

The lepton finding processors were optimised with events at  $\sqrt{s} = 1.4$  TeV. It was found that the same set of parameters is also effective for  $\sqrt{s} = 3$  TeV. The performances of the lepton finders at  $\sqrt{s} = 3$  TeV are shown in table 8.7.

When comparing the lepton finding performances at  $\sqrt{s} = 1.4$  TeV and  $\sqrt{s} = 3$  TeV, the performance for  $\sqrt{s} = 1.4$  TeV is better. This is because at  $\sqrt{s} = 3$  TeV, particles tend to be boosted more and the spatial separation between particles is smaller due to the higher multiplicities. Consequently particles are less isolated from each other. The higher centre-of-mass energy also affects the performance of the forward lepton finder. Whilst at  $\sqrt{s} = 1.4$  TeV, the forward finder only rejects 5% of the  $e^+e^- \rightarrow qqqq\ell\nu$  background events and 1% of the signal events, at  $\sqrt{s} = 3$  TeV it rejects 19% of events from the same background process and 4% of the signal events, as more leptons are boosted into the forward region.

## 8.4 Jet reconstruction

The signal process,  $HH \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}qqqq$ , is a six-quark final state, which will result in multiple reconstructed jets. The pairing of jets to form the H,  $W^+$  and  $W^-$  in the event is an essential part of the event reconstruction. In this section, the optimisation of the jet reconstruction is discussed.

Efficiency (3 TeV)	Signal	$e^+e^- \rightarrow qqqq\ell\nu$	$e^-\gamma(\text{BS}) \rightarrow e^-qqqq$
ISOLATEDLEPTONFINDER	99.5%	66.8%	88.8%
ISOLATEDLEPTONIDENTIFER	99.0%	52.5%	82.2%
TAUFINDER	97.7%	79.5%	76.7%
ISOLATEDTAUIDENTIFER	86.3%	60.3%	92.6%
Forward Finder Processors	95.9%	80.7%	55.4%
Combined	81.0%	23.3%	33.4%

**Table 8.7:** The performances of the lepton finding algorithms for the signal events and selected background events at  $\sqrt{s} = 3$  TeV.  $\gamma$  (BS) represents a real photon from beamstrahlung (BS). Numbers represent the fractions of events where no leptons are identified by the individual lepton finder.

#### 8.4.1 Jet reconstruction optimisation

Jet reconstruction algorithms cluster particles into jets. For this analysis, longitudinal invariant  $k_t$  jet algorithm is chosen for the jet clustering, as discussed in section 4.4.2. The free parameter for  $k_t$  algorithm is the  $R$  parameter, which controls the radius of the jet. The optimal jet clustering will also depend on the centre-of-mass energy of the event. This is particularly important at CLIC because of the large beam induced background from relative low  $p_T$  particles. Hence a suitable level of background suppression needs to be chosen, which is incorporated in the choice of the PFO collection.

The use of the  $k_t$  jet algorithm in exclusive modes allows some particles to be clustered into beam jet, which is not used in the subsequent event reconstruction.

The value of the  $R$  parameter and the PFO collection are chosen to optimise the invariant mass and mass resolution of H and W. To choose the optimal parameters,  $HH \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}qqqq$  events are processed through  $k_t$  jet algorithm in the six-jet exclusive mode. The six jets are paired using the MC truth information by examining the decay chain of MC particles. Four invariant mass distributions are obtained: two Higgs masses ( $m_{H_{bb}}$  and  $m_{H_{WW^*}}$ ) and two W masses ( $m_W$  and  $m_{W^*}$ ). Here  $W^*$  indicates the off-mass-shell W boson. The MC paring is used to optimise the choice of parameters. It is not used in the subsequent analysis.

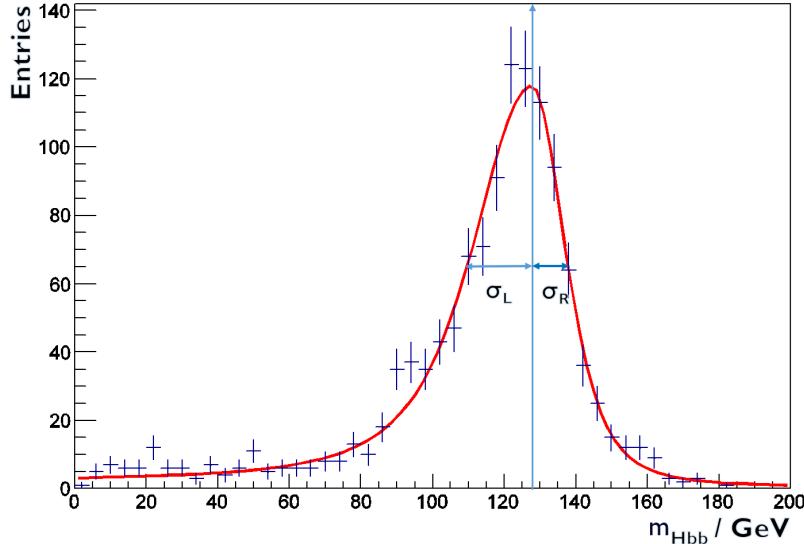
Three invariant mass distributions are considered:  $m_{H_{bb}}$ ,  $m_{H_{WW^*}}$ , and  $m_W$ . The optimal jet reconstruction should produce sharp mass peaks around the simulated particle masses. For example, figure 8.5 shows the  $m_{H_{bb}}$  invariant mass distribution for  $R = 1.3$  using the loose PFO collection for samples at  $\sqrt{s} = 3$  TeV. An analytical functional form is fitted to describe the shape. The fitting function is a Gaussian-like

function. Additional parameters are used in the fitting function to describe the tails of the distribution. The fitting function takes the form of

$$f(m) = A \exp \left\{ -\frac{(m - \mu)^2}{g} \right\}, \quad (8.2)$$

$$g = \begin{cases} 2\sigma_L + \alpha_L(m - \mu), & \text{if } m < \mu, \\ 2\sigma_R + \alpha_R(m - \mu), & \text{if } m \geq \mu, \end{cases} \quad (8.3)$$

where:  $\mu$  is the fitted mass peak position;  $\sigma_L$  and  $\sigma_R$  allow for an asymmetrical width of the distribution;  $\alpha_L$  and  $\alpha_R$  account for a constant tail of the distribution; and  $A$  is a normalisation factor.



**Figure 8.5:** A typical example of the reconstructed  $m_{H_{bb}}$  mass distribution for  $R = 1.3$  using loose PFO collection for  $HH \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}qqqq$  samples at  $\sqrt{s} = 3$  TeV. The fitting function is superimposed in red. The arrow shows the fitted peak position.

To parameterise the performance of different jet algorithm settings, the overall relative width is used, defined as  $(\sigma_L + \sigma_R)/M$ . A smaller width indicates a better mass resolution. The fitted  $H_{bb}$ ,  $H_{WW^*}$ , and  $W$  masses are studied for  $R$  values between 0.5 and 1.3, and with the three possible PFO collections: loose, normal, and tight.

Figure 8.6 shows the variation of the mass peak position and its relative width as a function of  $R$  and PFO collections, for  $m_{H_{bb}}$ ,  $m_{H_{WW^*}}$ , and  $m_W$ . The mass peak position,  $\mu$ , increases as  $R$  increases. This is because more particles are included in jets with

increasing jet radii. For the relative width, the values for  $H_{bb}$  increase with increasing jet radii, but the values for  $H_{WW^*}$  decrease with increasing jet radii. This is due to a compensating effect; the invariant mass for  $H_{WW^*}$  is formed from four jets, which prefers a large jet radius, whereas the invariant mass for  $H_{bb}$  is obtained from two jets, which favours a small jet radius.

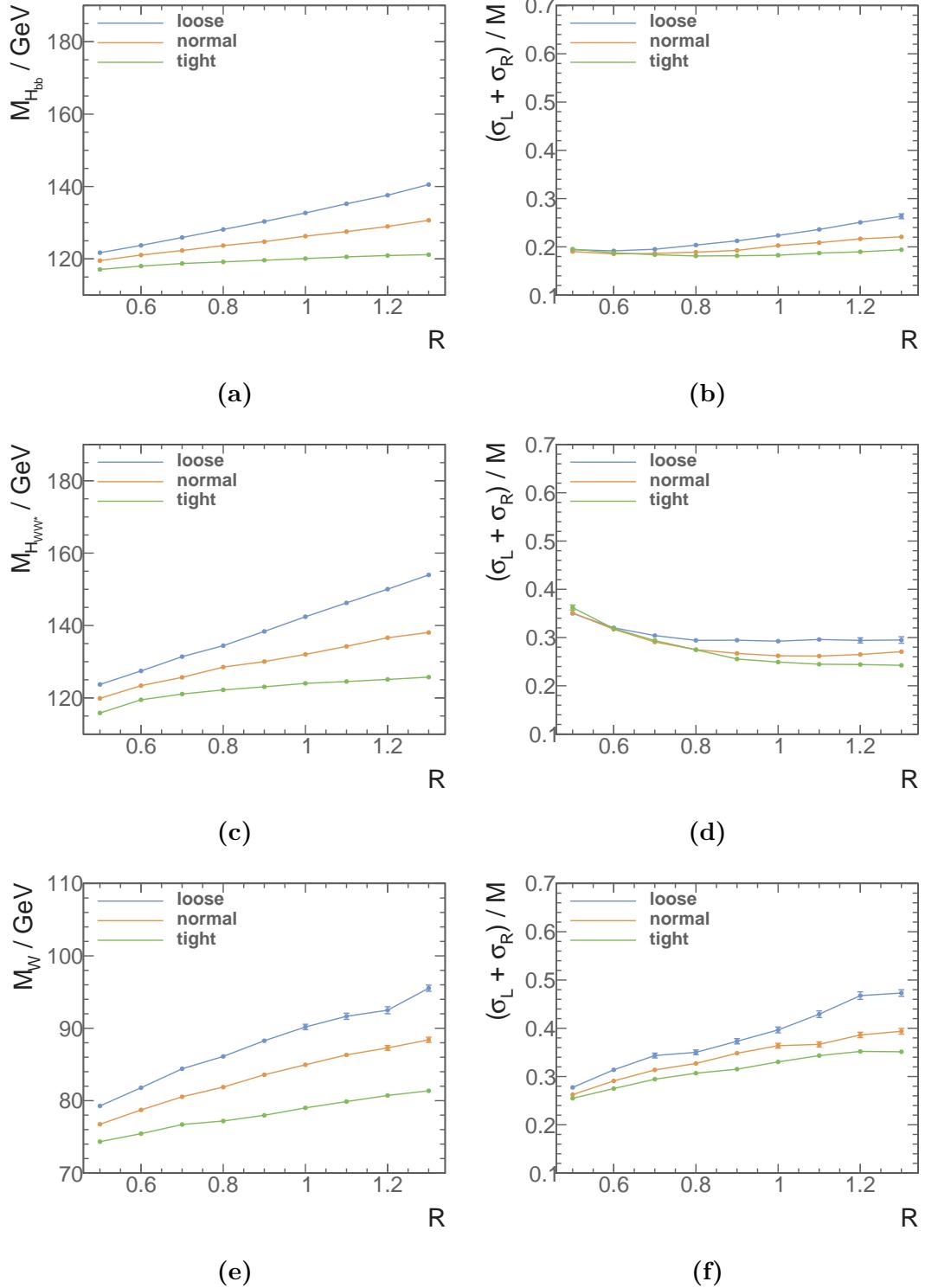
The choice of PFO collection impacts number of PFOs in the event. The loose PFO selection has the most PFOs in the event and, therefore, the largest invariant mass and worst mass resolution.

Based on the results summarised in figure 8.6 for this analysis, it was decided to use  $R = 0.7$  with the selected PFO collection. This choice gives good fitted mass peak positions for  $H_{bb}$ ,  $H_{WW^*}$  and  $W$ . The extracted fitted parameters of optimal jet reconstructions are summarised in table 8.8.

Fitted jet parameter	$\sqrt{s} = 1.4 \text{ TeV}$	$\sqrt{s} = 3 \text{ TeV}$
$\mu_{H_{bb}}$	$122.3 \pm 0.2$	$119.1 \pm 0.3$
$\sigma_{L,H_{bb}}$	$15.2 \pm 0.2$	$15.0 \pm 0.3$
$\sigma_{R,H_{bb}}$	$7.6 \pm 0.2$	$8.4 \pm 0.2$
$\mu_{H_{WW^*}}$	$125.7 \pm 0.2$	$123.0 \pm 0.3$
$\sigma_{L,H_{WW^*}}$	$29.4 \pm 0.3$	$36.6 \pm 0.6$
$\sigma_{R,H_{WW^*}}$	$7.2 \pm 0.2$	$7.4 \pm 0.2$
$\mu_W$	$80.5 \pm 0.2$	$78.1 \pm 0.3$
$\sigma_{L,W}$	$16.2 \pm 0.3$	$13.1 \pm 0.4$
$\sigma_{R,W}$	$9.0 \pm 0.2$	$9.5 \pm 0.2$

**Table 8.8:** The fitted mass parameters for  $\sqrt{s} = 1.4 \text{ TeV}$  analysis:  $R = 0.7$  using the selected PFO collection, and for  $\sqrt{s} = 3 \text{ TeV}$  analysis:  $R = 0.7$  using the tight selected PFO collection.

A separate jet reconstruction optimisation is performed for  $\sqrt{s} = 3 \text{ TeV}$  analysis. Figure 8.7 shows the variation of fitted mass peak positions and the relative mass resolutions for  $H_{bb}$ ,  $H_{WW^*}$ , and  $W$  as function of  $R$  and PFO collections. The relative mass resolution of  $W$  boson quickly degrades with an increasing  $R$ . The fitted mass peak positions also increases more rapidly with the increase of  $R$ , compared with the fitted mass peak positions at  $\sqrt{s} = 1.4 \text{ TeV}$ . This is because at a higher centre-of-mass energy, more beam induced background particles are produced. The background particles, if included in the jets, will increase the invariant masses of the fitted physical bosons. Based



**Figure 8.6:** Distributions of: a) fitted mass peak positions for  $H_{bb}$ ; b) relative mass peak widths for  $H_{bb}$ ; c) fitted mass peak positions for  $H_{WW^*}$ ; d) relative mass peak widths for  $H_{WW^*}$ ; e) fitted mass peak positions for  $W$ ; and f) relative mass peak widths for  $W$ . All plots show the variation of the fitted masses and mass resolutions as a function of  $R$  for loose, normal, and tight selected PFO collections at  $\sqrt{s} = 1.4$  TeV, using  $HH \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}qqqq$  samples.

on this study,  $R = 0.7$  with the tight selected PFO collection was chosen for the  $\sqrt{s} = 3\text{ TeV}$  analysis. With chosen parameters, the better relative mass resolutions compensate for the invariant masses being slightly smaller than simulated values. The extracted fitted parameters of optimal jet reconstructions at  $\sqrt{s} = 3\text{ TeV}$  are summarised in table 8.8.

## 8.5 Jet flavour tagging

As the signal process,  $\text{HH} \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}\text{qqqq}$ , contains two b quarks in the final state, identifying jets originated from b quarks is an important part of the event selection.

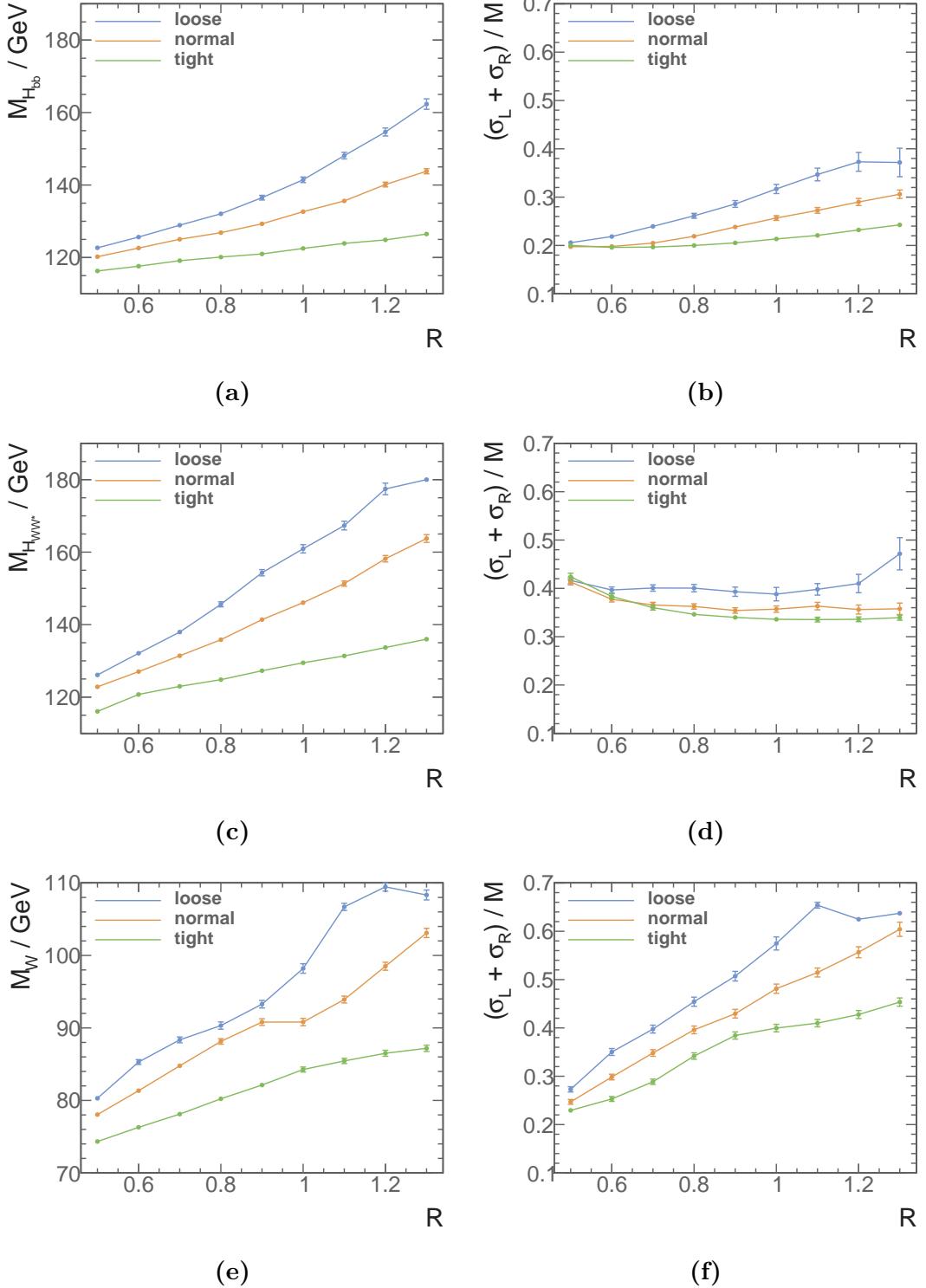
The flavour tagging processor, LCFIPLUS [104] is used. The processor is based on the LCFIVERTEX package [105], which was used in the simulation studies for the ILC Letter of Intent [28, 106] and the CLIC Concept Design Report [2].

After the previous jet clustering step, particles which are not in the beam jet are used as inputs for the flavour tagging. The flavour tagging algorithm identifies vertices, then re-cluster particles into jets. Lastly, the algorithm decides if a jet is a b-quark jet or a c-quark jet.

The vertex finding algorithms perform vertex fitting and identify primary and secondary vertices. There are two vertex refining algorithm. First algorithm rejects the topology of a neutral particle that decays into pairs of charged particles, which can be mistaken as the decay of b or c quarks. The second algorithm is performed after the re-clustering step to reconstruct more secondary vertices, with additional information from the jet clustering.

The jet re-clustering algorithm is a modified Durham algorithm, with additional constraint of the secondary vertices and the muons, which are identified from semi-leptonic decay of the quarks, falling into the same jet as the quarks. This ensures the topology of the jet remains consistent with the hadronic decays of heavy quarks.

Having obtained re-clustered jets, the next step is the flavour tagging, which uses a multivariate classifier to determine if a jet is from b quark or c quark. LCFIPLUS uses the Boosted Decision Tree MVA MULTICLASS classifier as implemented in the TMVA software package [78]. There are four categories for classification: jets with zero, one, two properly reconstructed vertices, or a single-track pseudo-vertex. A jet can be classified into one of three classes: b jet, c jet, or a light flavour quark (u, d or s) jet.

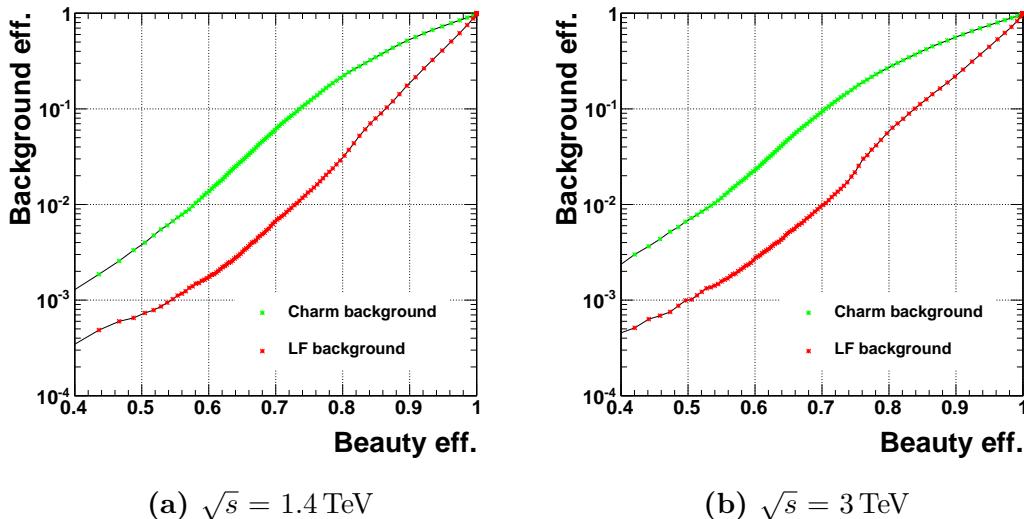


**Figure 8.7:** Distributions of: a) fitted mass peak positions of  $H_{bb}$ ; b) relative mass peak widths of  $H_{bb}$ ; c) fitted mass peak positions of  $H_{WW^*}$ ; d) relative mass peak widths of  $H_{WW^*}$ ; e) fitted mass peak positions of  $W$ ; and f) relative mass peak widths of  $W$ . All plots show the variation of fitted masses and mass resolutions as a function of  $R$  for loose, normal, and tight selected PFO collections at  $\sqrt{s} = 3 \text{ TeV}$ , using  $\text{HH} \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}qqqq$  samples.

The MVA MULTICLASS classifier was trained with  $e^+e^- \rightarrow Z\bar{\nu}\nu$  event at  $\sqrt{s} = 1.4$  TeV, where Z decays to  $b\bar{b}$ ,  $c\bar{c}$ , or  $u\bar{u}/d\bar{d}/s\bar{s}$ . The training sample contains missing momentum, which is similar to the signal sample. The training sample only has two quarks in the final states, which reduces the error in jet clustering and provides a good ground truth for training. The MVA classicisation efficiency with the training samples is shown in figure 8.8a.

Having trained the MVA classifier, the MVA classifier is applied to the samples. Under the signal hypothesis, the re-cluster algorithm is set to find six jets. The normalised distribution of the highest b-jet tag value for the  $HH \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}qqqq$  sample is shown in figure 8.9.

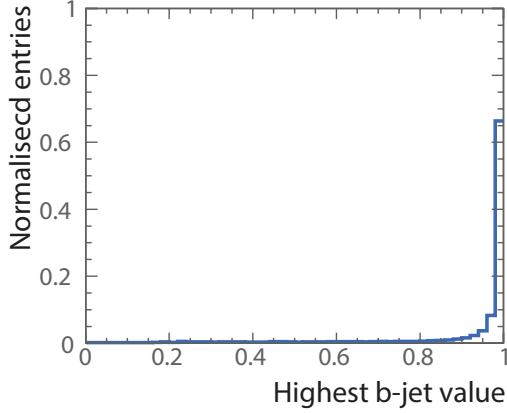
For the  $\sqrt{s} = 3$  TeV analysis, the MVA classifier is re-trained with  $e^+e^- \rightarrow Z\bar{\nu}\nu$  event at  $\sqrt{s} = 3$  TeV. The performance of the flavour tagging with training samples is shown in figure 8.8b. The performance at  $\sqrt{s} = 3$  TeV is slightly worse than at  $\sqrt{s} = 1.4$  TeV, because at the higher centre-of-mass energy, jets are more collimated and more difficult to separate.



**Figure 8.8:** Performance of b-jet tagging with  $e^+e^- \rightarrow Z\bar{\nu}\nu$  samples, where Z decays to  $b\bar{b}$ ,  $c\bar{c}$ , or  $u\bar{u}/d\bar{d}/s\bar{s}$  at: a)  $\sqrt{s} = 1.4$  TeV; and b)  $\sqrt{s} = 3$  TeV.

### 8.5.1 Mutually exclusive cuts for $HH \rightarrow b\bar{b}W^+W^-$ and $HH \rightarrow b\bar{b}b\bar{b}$

The two  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  final states with the largest branching fractions are  $HH \rightarrow b\bar{b}b\bar{b}$  (31.5%) and  $HH \rightarrow b\bar{b}W^+W^-$  (25.9%). These two final states have different topologies



**Figure 8.9:** The distribution of the highest b-jet value for the  $\text{HH} \rightarrow b\bar{b}W^+W^- \rightarrow b\bar{b}\text{qqqq}$  events at  $\sqrt{s} = 1.4$  TeV. The area under the curve is normalised to unity.

and are subjects of two analysis strategies. The  $\text{HH} \rightarrow b\bar{b}W^+W^-$  final state is the subject of this thesis. The study of the  $\text{HH} \rightarrow b\bar{b}b\bar{b}$  final state is the subject of an independent analysis. Because the results of the two studies are subsequently combined, a set of cuts are designed to separate samples, for both signal and background events, into two mutually exclusive sets for two independent analyses. This ensures there are no correlations between two analyses.

The most distinctive difference between the  $\text{HH} \rightarrow b\bar{b}W^+W^-$  and  $\text{HH} \rightarrow b\bar{b}b\bar{b}$  sub-channels is the different jet multiplicity and the different number of b-jets in the final state. Consequently variables relating to the number of b-jets and total number of jets are suitable for separating the two sub-channels.

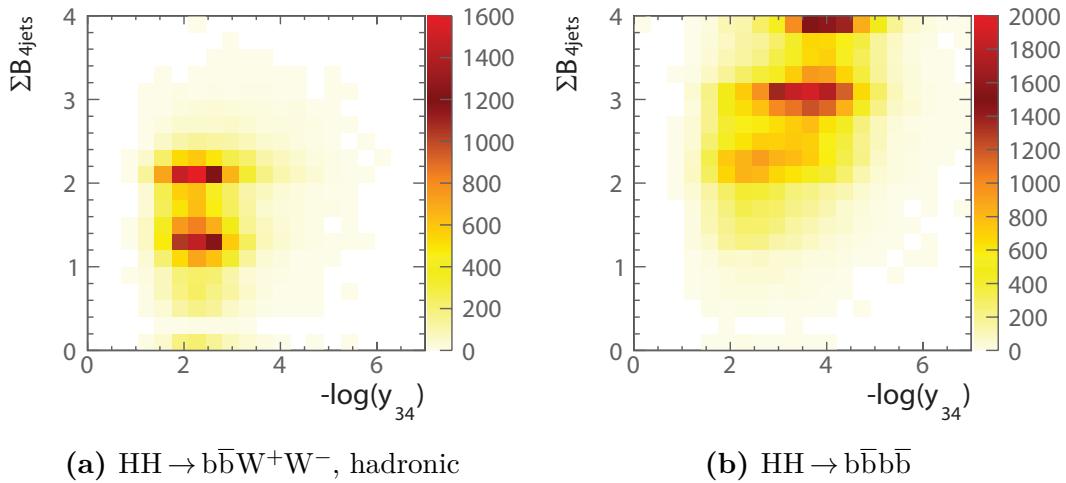
Figure 8.10 shows the sum of b-jet tag values, when the event is clustered into four jets, as a function of  $-\log(y_{34})$  for the hadronic  $W^+W^-$  decay in  $\text{HH} \rightarrow b\bar{b}W^+W^-$  and  $\text{HH} \rightarrow b\bar{b}b\bar{b}$ . As expected, the two sub-channels can be clearly separated in this two dimensional phase space. A rectangular cut can be used to separate the phase space into two spaces, denoted as  $S$  and  $\neg S$ . The hadronic  $W^+W^-$  decay in  $\text{HH} \rightarrow b\bar{b}W^+W^-$  events should be contained in phase space  $S$ , and the  $\text{HH} \rightarrow b\bar{b}b\bar{b}$  events should be contained in phase space  $\neg S$ .

The optimal cuts are chosen such that they maximise:

$$\varepsilon = \frac{N_{\text{HH} \rightarrow b\bar{b}W^+W^-, \text{hadronic}} \in S}{N_{\text{HH} \rightarrow b\bar{b}W^+W^-, \text{hadronic}}} \times \frac{N_{\text{HH} \rightarrow b\bar{b}b\bar{b}} \in \neg S}{N_{\text{HH} \rightarrow b\bar{b}b\bar{b}}}, \quad (8.4)$$

where  $N \in S$  indicates number of events in the phase space  $S$ .

Several combinations of pairs of variables were considered. In each case, the product of the fraction of the sub-channel events in each space,  $\varepsilon$  was maximised. This procedure identified  $\sum B_{4\text{jets}} < 2.3$ ,  $-\log(y_{34}) < 3.7$  as the best choice with 86% of the hadronic  $W^+W^-$  decay in  $\text{HH} \rightarrow b\bar{b}W^+W^-$  events are in  $S$  and 78% of the  $\text{HH} \rightarrow b\bar{b}b\bar{b}$  events are in  $\neg S$ . The full list of fraction of events after passing mutually exclusive cuts for individual background processes are listed in table 8.9.



**Figure 8.10:** The two-dimensional distribution of sum of b-jet tag values against  $-\log(y_{34})$  for: a) hadronic  $W^+W^-$  decay of  $\text{HH} \rightarrow b\bar{b}W^+W^-$ ; and b)  $\text{HH} \rightarrow b\bar{b}b\bar{b}$  events at  $\sqrt{s} = 1.4$  TeV. The sum of b-jet tag values is calculated for the cases where events are clustered into four jets.

## 8.6 Jet pairing

Events are reconstructed assuming the  $\text{HH} \rightarrow b\bar{b}W^+W^-$  signal topology. The six jets are obtained from the jet re-clustering step in the LCFIPLUS processor. The next step is to group jets according to signal event topology. Jets are paired up such that there are two jets for  $H \rightarrow b\bar{b}$ , two jets for hadronic decay of a  $W$ , and two jets for hadronic decay of a  $W^*$ . In addition, the two  $W$ s should be from the  $H$  boson decay.

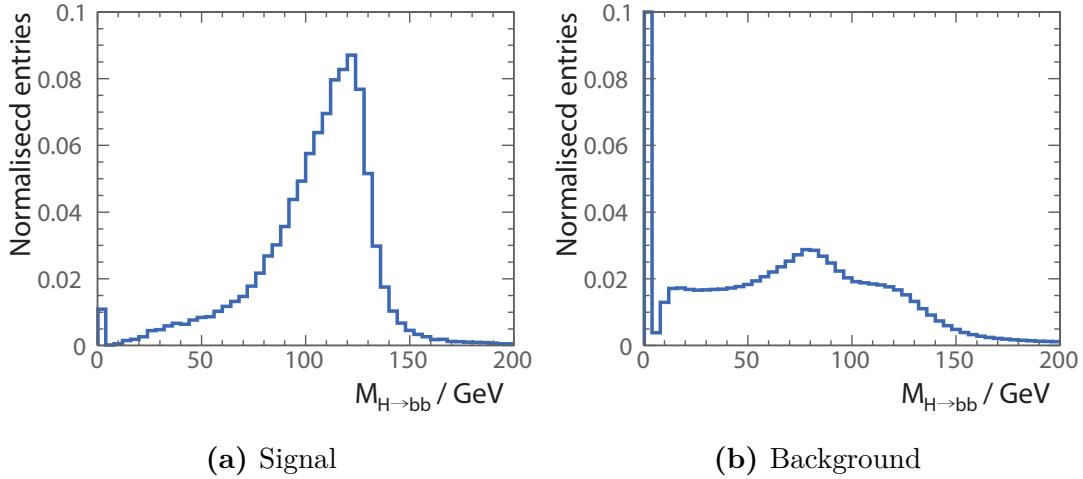
Six jets are associated to  $H_{bb}$ ,  $W$  and  $W^*$ . There are 90 possible permutations for associating six jets to  $H_{bb}$ ,  $W$  and  $W^*$ . The best permutation is obtained by choosing the permutation that gives the smallest  $\chi^2$  quantity, representing the consistency of the hypothesis with the signal topology:

$$\chi^2 = \left( \frac{m_{ij} - \mu_{H_{bb}}}{\sigma'_{H_{bb}}} \right)^2 + \left( \frac{m_{klmn} - \mu_{H_{WW^*}}}{\sigma'_{H_{WW^*}}} \right)^2 + \left( \frac{m_{kl} - \mu_W}{\sigma'_W} \right)^2, \quad (8.5)$$

where the indices represent the six jets. The parameter  $\mu$  and  $\sigma'$  are the expected peak and (asymmetric) width of the reconstructed mass distributions given in table 8.8, defined as:

$$\sigma'_{H_{bb}} = \begin{cases} \sigma_{L,H_{bb}}, & \text{if } m_{ij} < \mu_{H_{bb}}, \\ \sigma_{R,H_{bb}}, & \text{otherwise,} \end{cases} \text{ etc.} \quad (8.6)$$

A jet pairing is only considered when at least one of the jets associated to the  $H_{bb}$  decay has a b-jet tag  $> 0.2$ . Of these combinations of jets, the jet pairing giving smallest  $\chi^2$  is selected. Figure 8.11 shows the normalised distribution of  $m_{H_{bb}}$  after jet pairing, for the signal process,  $HH \rightarrow b\bar{b}W^+W^-$  and the sum of all background processes. For the signal process, the distribution peaks around the expected mass of  $m_{H_{bb}}$ . Around 1% of signal events have no solutions for the jet pairing, i.e. no jet has a b-jet tag  $> 0.2$ . These events are no longer considered in the analysis. The full list of fraction of events surviving after this jet pairing selection are listed in table 8.9 for signal  $HH \rightarrow b\bar{b}W^+W^-$  process and all background processes.



**Figure 8.11:** The distribution of  $m_{H_{bb}}$  for: a) the signal process, hadronic  $W^+W^-$  decay of  $HH \rightarrow b\bar{b}W^+W^-$ ; and b) the sum of all background processes normalised to the respective cross sections. The area under the curve is normalised to unity. All plots are shown for  $\sqrt{s} = 1.4 \text{ TeV}$ .

## 8.7 Pre-selection

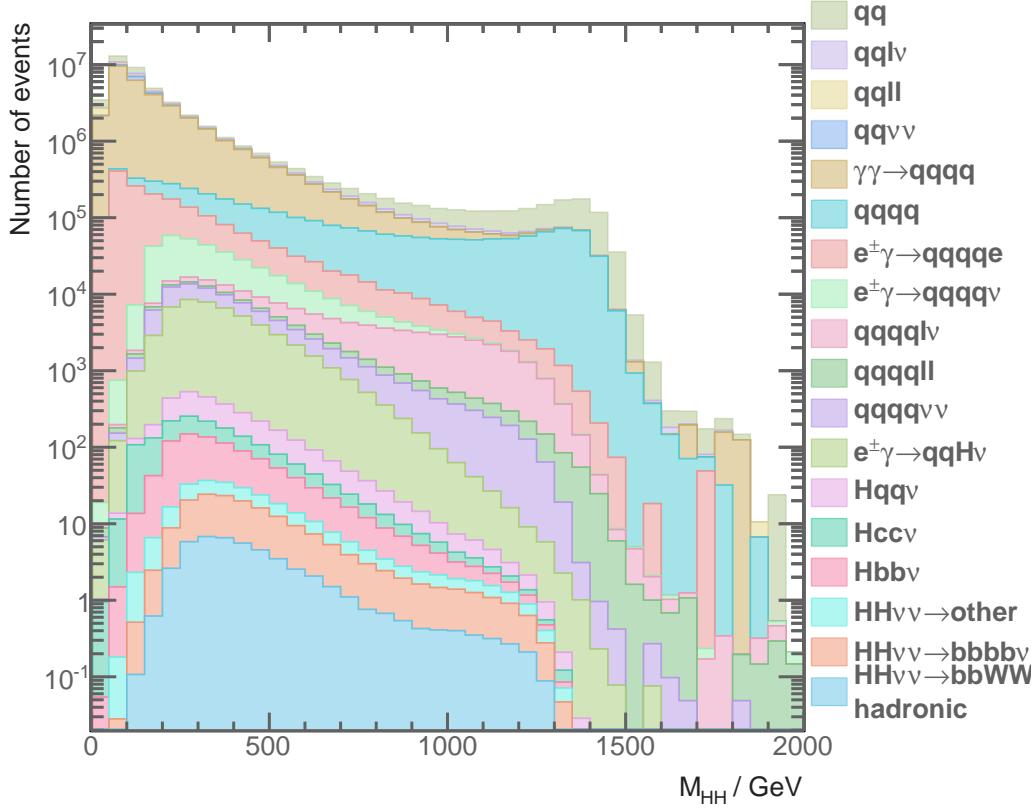
After the association of jets to candidate bosons are made under hypothesis that an event is a signal event, kinematic and topological variables can be calculated. A set of

$\sqrt{s} = 1.4 \text{ TeV}$	N	Lepton evto	$b\bar{b}W^+W^- / b\bar{b}W^+W^-$ separation	Valid jet Pairing
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e$ , hadronic	27.9	89.7%	79.1%	78.3%
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}b\bar{b}\nu_e\bar{\nu}_e$	67.6	90.8%	18.0%	18.0%
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow \text{other}$	128.0	40.8%	35.8%	31.2%
$e^+e^- \rightarrow q\bar{q}H\nu\bar{\nu}$	1290	72.8%	69.7%	57.7%
$e^+e^- \rightarrow c\bar{c}H\nu\bar{\nu}$	540	74.7%	59.8%	52.7%
$e^+e^- \rightarrow b\bar{b}H\nu\bar{\nu}$	465	74.3%	32.2%	31.8%
$e^+e^- \rightarrow qqqq$	1867650	79.9%	64.0%	38.6%
$e^+e^- \rightarrow qqqq\ell\ell$	93150	8.9%	8.2%	4.7%
$e^+e^- \rightarrow qqqq\ell\nu$	165600	16.5%	14.6%	13.3%
$e^+e^- \rightarrow qqqq\nu\bar{\nu}$	34800	87.6%	82.0%	46.8%
$e^+e^- \rightarrow qq$	6014250	81.0%	57.8%	39.0%
$e^+e^- \rightarrow qq\ell\nu$	6464550	22.5%	17.0%	10.5%
$e^+e^- \rightarrow qq\ell\ell$	4088700	19.4%	18.6%	12.4%
$e^+e^- \rightarrow qq\nu\nu$	1181550	91.8%	74.0%	47.3%
$e^\pm\gamma(\text{BS}) \rightarrow e^\pm qqqq$	2606625	34.2%	33.5%	22.9%
$e^\pm\gamma(\text{EPA}) \rightarrow e^\pm qqqq$	861000.0	16.4%	15.8%	10.7%
$e^\pm\gamma(\text{BS}) \rightarrow \nu qqqq$	178987.5	85.6%	81.3%	54.4%
$e^\pm\gamma(\text{EPA}) \rightarrow \nu qqqq$	52050	44.5%	42.0%	27.4%
$e^\pm\gamma(\text{BS}) \rightarrow qqH\nu$	35437.5	70.7%	65.0%	55.4%
$e^\pm\gamma(\text{EPA}) \rightarrow qqH\nu$	10170	37.0%	33.8%	28.8%
$\gamma(\text{BS})\gamma(\text{BS}) \rightarrow qqqq$	2054951.5	85.6%	81.3%	54.0%
$\gamma(\text{BS})\gamma(\text{EPA}) \rightarrow qqqq$	4521037.5	49.6%	48.5%	32.9%
$\gamma(\text{EPA})\gamma(\text{BS}) \rightarrow qqqq$	4539150	49.6%	48.5%	32.9%
$\gamma(\text{EPA})\gamma(\text{EPA}) \rightarrow qqqq$	1129500	31.0%	30.1%	20.5%

**Table 8.9:** The table shows the expected number of events, before cuts and after successive cuts: the lepton veto,  $HH \rightarrow b\bar{b}W^+W^- / HH \rightarrow b\bar{b}b\bar{b}$  separation, and valid jet pairing. Table shows the signal and background events at  $\sqrt{s} = 1.4 \text{ TeV}$ , assuming an integrated luminosity of  $1500 \text{ fb}^{-1}$ . q can be u, d, s, b or t. Unless specified, q,  $\ell$  and  $\nu$  represent either particles or the corresponding anti-particles.

pre-selection cuts are placed to discard the phase space dominated by background events. Cuts on  $p_T$ , b-jet tag, and invariant mass of the double Higgs system are used.

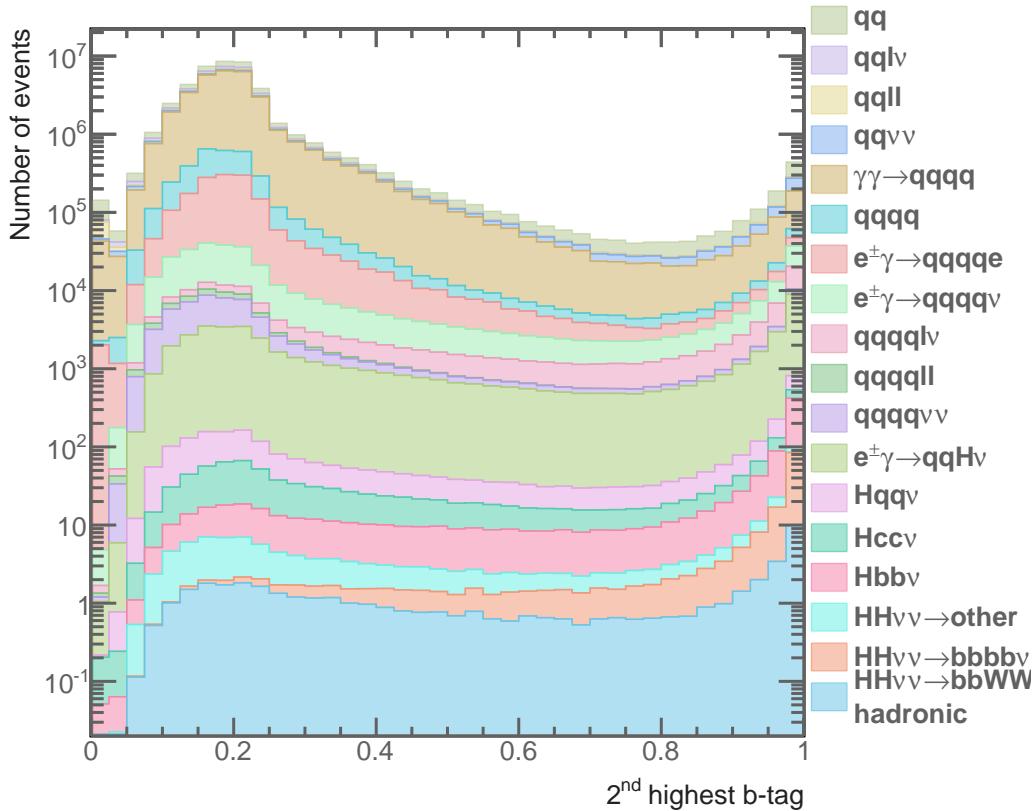
Since both Higgs bosons are on mass shell, the invariant mass of the double Higgs system is large. Consequently, a cut on  $m_{HH} > 150 \text{ GeV}$ , as shown in Figure 8.12, removes a small amount of signal events, but discard lots of background events, especially  $\gamma\gamma \rightarrow \text{qqqq}$  events.



**Figure 8.12:** Distributions of the invariant mass of the two Higgs system for  $\sqrt{s} = 1.4 \text{ TeV}$ , assuming an intergraded luminosity of  $1500 \text{ fb}^{-1}$ .

Many background events do not have b-quark jets in the final state. Therefore, by requiring the second highest b-jet tag value greater than 0.2, as shown in Figure 8.13, background events with no b quarks in final states are removed.

The signal final states have neutrinos and hence missing momentum in the events. Therefore, the transverse momentum of the two Higgs system is non zero. A cut of



**Figure 8.13:** Distributions of the second highest b-jet tag value for  $\sqrt{s} = 1.4$  TeV, assuming an intergraded luminosity of  $1500 \text{ fb}^{-1}$ .

$p_T > 30$  GeV, as shown in figure 8.14, is extremely effective against background processes with no neutrinos in the final state.

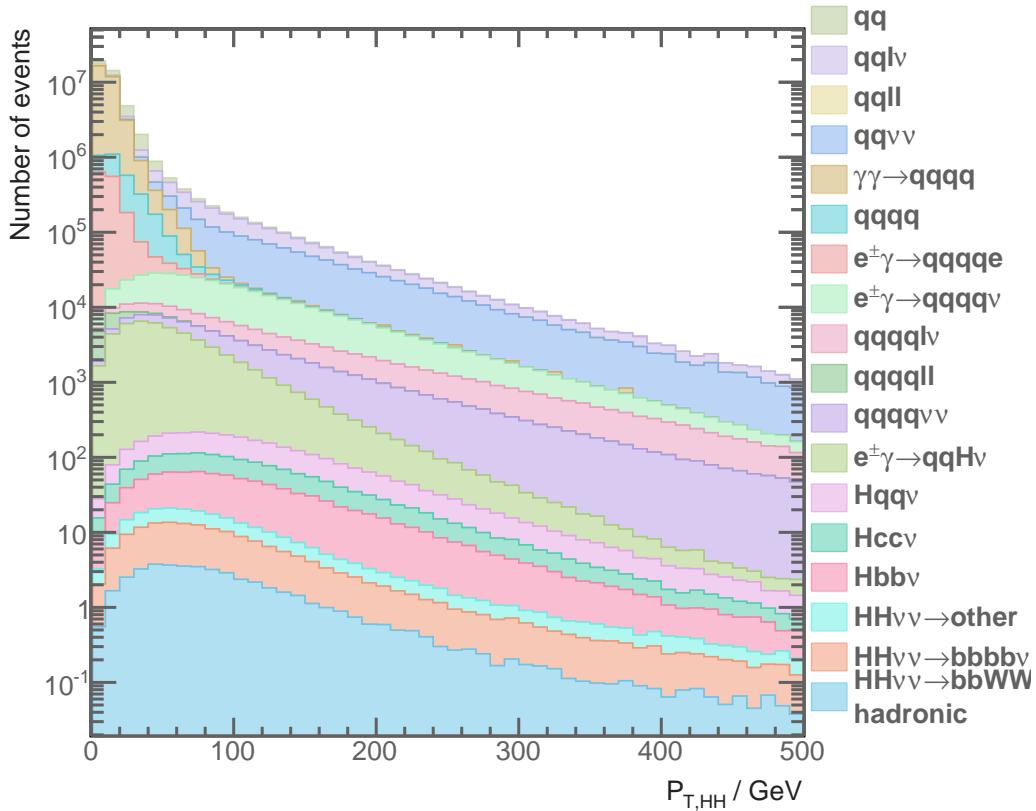
The full list of fraction of events after each pre-selection cut can be found in table 8.10.

## 8.8 MVA variables

Having extracted information about leptons, b-jets, and jet pairing, a number of variables are used to differentiate the signal to the background events. These variables are the basis of the subsequent MVA event selection. The variables used are listed in table 8.11. The distributions of the four most powerful discriminators are shown in figure 8.15.

Process	$m_{HH} > 150 \text{ GeV}$	$B_2 > 0.2$	$p_T > 30 \text{ GeV}$
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e$ , hadronic	78.1%	66.3%	59.7%
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}b\bar{b}\nu_e\bar{\nu}_e$	17.8%	17.4%	15.4%
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow$ other	30.5%	23.0%	20.5%
$e^+e^- \rightarrow q\bar{q}H\nu\bar{\nu}$	56.8%	42.3%	39.5%
$e^+e^- \rightarrow c\bar{c}H\nu\bar{\nu}$	44.8%	34.1%	31.7%
$e^+e^- \rightarrow b\bar{b}H\nu\bar{\nu}$	30.7%	27.0%	25.2%
$e^+e^- \rightarrow qqqq$	36.1%	13.2%	3.4%
$e^+e^- \rightarrow qqqq\ell\ell$	4.7%	1.5%	0.3%
$e^+e^- \rightarrow qqqq\ell\nu$	13.2%	10.7%	9.8%
$e^+e^- \rightarrow qqqq\nu\bar{\nu}$	46.1%	17.7%	16.6%
$e^+e^- \rightarrow qq$	8.1%	3.7%	0.8%
$e^+e^- \rightarrow q\bar{q}\ell\nu$	3.1%	1.2%	0.9%
$e^+e^- \rightarrow q\bar{q}\ell\ell$	0.7%	0.4%	0.1%
$e^+e^- \rightarrow qq\nu\nu$	9%	4.3%	4.0%
$e^\pm\gamma(\text{BS}) \rightarrow e^\pm qqqq$	10.1%	4.1%	0.4%
$e^\pm\gamma(\text{EPA}) \rightarrow e^\pm qqqq$	5.1%	2.0%	0.3%
$e^\pm\gamma(\text{BS}) \rightarrow \nu qqqq$	53.0%	28.0%	25.1%
$e^-\gamma(\text{EPA}) \rightarrow \nu qqqq$	26.7%	13.8%	12.5%
$e^\pm\gamma(\text{BS}) \rightarrow qqH\nu$	54.3%	40.3%	30.6%
$e^-\gamma(\text{EPA}) \rightarrow qqH\nu$	28.2%	20.9%	16.1%
$\gamma(\text{BS})\gamma(\text{BS}) \rightarrow qqqq$	23.1%	9.2%	0.3%
$\gamma(\text{BS})\gamma(\text{EPA}) \rightarrow qqqq$	13.6%	5.4%	0.4%
$\gamma(\text{EPA})\gamma(\text{BS}) \rightarrow qqqq$	13.6%	5.4%	0.3%
$\gamma(\text{EPA})\gamma(\text{EPA}) \rightarrow qqqq$	8.6%	3.5%	0.3%

**Table 8.10:** The table shows the expected number of events after successive cuts: invariant mass of the two Higgs system  $> 150 \text{ GeV}$ , the second highest b-jet tag value  $> 0.2$ , and the transverse momentum of the two Higgs system  $> 30 \text{ GeV}$ . All cuts include the lepton veto,  $HH \rightarrow b\bar{b}W^+W^-/HH \rightarrow b\bar{b}b\bar{b}$  separation, and valid jet pairing. Table shows the signal and background events at  $\sqrt{s} = 1.4 \text{ TeV}$ , assuming an integrated luminosity of  $1500 \text{ fb}^{-1}$ . q can be u, d, s, b or t. Unless specified, q,  $\ell$  and  $\nu$  represent either particles or the corresponding anti-particles.



**Figure 8.14:** Distributions of the transverse momentum of the two Higgs system for  $\sqrt{s} = 1.4$  TeV, assuming an intergraded luminosity of  $1500 \text{ fb}^{-1}$ .

### 8.8.1 Invariant mass variables

Four invariant masses are used in the MVA event selection: the invariant mass of  $H_{bb}$  ( $m_{H_{bb}}$ ), the invariant mass of  $H_{WW^*}$  ( $m_{H_{WW^*}}$ ), the invariant mass of  $W$  ( $m_W$ ), and the invariant mass of the double Higgs system ( $m_{HH}$ ).

After the jet pairing under the hypothesis of the signal events, the distributions of the invariant mass of the physical bosons of the signal events have peaks around the expect masses, where the distributions of the background events do not have such resonance structure. Shown in figure 8.15a, the distributions of the invariant mass of the  $H_{bb}$  is different to the distributions of the background events. Similarly, the distributions of the invariant mass of the  $H_{WW^*}$ , shown in figure 8.15b have a different peak position to the distributions of the background events. The invariant mass of the double Higgs system

in the signal events is large due to the presence of two on-mass-shell Higgs bosons, which is also different to the distribution of the background events

### 8.8.2 Energy and momentum variables

Six energy and momentum variables participate in the MVA event selection: the energy of the off-mass-shell W ( $E_{W^*}$ ), the energy of the missing momenta ( $E_{mis}$ ), the transverse momentum of  $H_{bb}$  ( $p_{TH_{bb}}$ ), the transverse momentum of  $H_{WW^*}$  ( $p_{TH_{WW^*}}$ ), the transverse momentum of W ( $p_{TW}$ ), and the transverse momentum of the double Higgs system ( $p_{THH}$ ).

For the off-mass-shell W, the energy is used instead of the invariant mass, as invariant mass distribution of  $W^*$  does not have a resonance structure. The energy of the missing momenta is powerful against background events with no neutrinos in the final states. The missing momenta is calculated by assuming the collision at  $\sqrt{s}$  and a beam crossing angle of 20 mrad. Other momentum variables correspond to the same physical bosons or the double Higgs system used in the invariant mass variables, for the same reason that the distributions of these momentum variables are different for the signal events and the background events.

### 8.8.3 Lab-frame angle variables

Four lab-frame angle variables are in the MVA event selection: the pseudorapidity of the missing momenta ( $\eta_{mis}$ ), the acollinearity of the two jets associated with  $H_{bb}$  ( $A_{H_{bb}}$ ), the acollinearity of the two jets associated with  $H_{WW^*}$  ( $A_{H_{WW^*}}$ ), and the acollinearity of the two Higgs bosons ( $A_{HH}$ ).

The pseudorapidity of the missing momenta is used, instead of the polar angle, because the forward polar angles are transformed to a larger range in the pseudorapidity. The pseudorapidity of the missing momenta is defined as

$$\eta_{mis} \equiv -\ln \left[ \tan \left( \frac{\theta_{mis}}{2} \right) \right], \quad (8.7)$$

where  $\theta_{mis}$  is the polar angle of the missing momenta measured in a spherical polar coordinate system.

Acollinearity is a measures the angle between the two momenta. The definition for the acollinearity for momenta  $i$  and momenta  $j$  is

$$A_{ij} = \pi - \cos^{-1} (\hat{\mathbf{p}}_i \cdot \hat{\mathbf{p}}_j), \quad (8.8)$$

where  $\hat{\mathbf{p}}_i$  is the unit momentum three-vector of momenta  $i$ . The distribution of the  $A_{H_{bb}}$ , shown in figure 8.15c, peaks at the value of 0 or  $\pi$  for many background events, which are not the same as the signal events. For the same reason, the distributions of  $A_{H_{WW^*}}$  and  $A_{HH}$  are different for the signal and the background events.

### 8.8.4 Boosted-frame angle variables

The MVA event selection also uses five boosted-frame angle variables: the angle between two jets associated with  $H_{bb}$  in the  $H_{bb}$  decay rest frame ( $\cos(\theta_{H_{bb}}^*)$ ), the angle between two W associated with  $H_{WW^*}$  in the  $H_{WW^*}$  decay rest frame ( $\cos(\theta_{H_{WW^*}}^*)$ ), the angle between two jets associated with W in the W decay rest frame ( $\cos(\theta_W^*)$ ), the angle between two jets associated with  $W^*$  in the  $W^*$  decay rest frame ( $\cos(\theta_{W^*}^*)$ ), and the angle between two Higgs bosons in two Higgs bosons decay rest frame ( $\cos(\theta_{HH}^*)$ ).

These variables are some of the most powerful variables. For example,  $\cos(\theta_{H_{bb}}^*)$  for the signal events has a uniform distribution, shown in figure 8.15d, as it is equally likely for two quarks to decay in any open angle in the  $H_{bb}$  decay rest frame. For the background events, by pairing jets under the signal event hypothesis,  $\cos(\theta_{H_{bb}}^*)$  does not have a flat distribution. The  $\cos(\theta_{H_{bb}}^*)$  distribution for the background events peaks at 1.

### 8.8.5 Event shape variables

Five event shapes variables are used in the MVA event selection: the absolute value of the sphericity ( $|\mathbf{S}|$ ), the negative logarithm of  $y_{23}$  ( $-\ln(y_{23})$ ), the negative logarithm of  $y_{34}$  ( $-\ln(y_{34})$ ), the negative logarithm of  $y_{45}$  ( $-\ln(y_{45})$ ), the negative logarithm of  $y_{56}$  ( $-\ln(y_{56})$ )).

The sphericity,  $\mathbf{S}$ , is a measurement of the spherically symmetry of the event, which will be different for the signal and background events. The sphericity is derived from the

sphericity tensor [107]. The sphericity tensor is defined as

$$\mathbf{S}^{\alpha\beta} = \frac{\sum_i p_i^\alpha p_i^\beta}{\sum_i |\vec{p}_i|^2}, \quad (8.9)$$

where  $\vec{p}_i$  is the momentum vector of the particle  $i$ ; index  $i$  is summed over all particles in the event; and  $\alpha$  and  $\beta$  refer to the x, y, z coordinate axis. Eigenvalues of  $\mathbf{S}$  tensor, denoted with  $\lambda_1$ ,  $\lambda_2$ ,  $\lambda_3$ , can be found via diagonalisation of the matrix  $\mathbf{S}$ . The normalisation condition requires  $\lambda_1 \geq \lambda_2 \geq \lambda_3$  and  $\lambda_1 + \lambda_2 + \lambda_3 = 1$ . Sphericity,  $S$ , is defined in terms of  $\lambda$ ,

$$S = \frac{3}{2}(\lambda_1 + \lambda_2). \quad (8.10)$$

$S$ , is 0 for a perfect pencil-like back-to-back two-jet event, and 1 for a perfect spherically symmetric event.

### 8.8.6 b and c tag variables

Six b-jet and c-jet tag variables are used in the MVA event selection: the highest b-jet tag value of the two jets associated with  $H_{bb}$  ( $B_{1,H_{bb}}$ ), the lowest b-jet tag value of the two jets associated with  $H_{bb}$  ( $B_{2,H_{bb}}$ ), the highest b-jet tag value of the two jets associated with  $W$  ( $B_{1,W}$ ), the highest b-jet tag value of the two jets associated with  $W^*$  ( $B_{1,W^*}$ ), the highest c-jet tag value of the two jets associated with  $H_{bb}$  ( $C_{1,H_{bb}}$ ), and the highest c-jet tag value of the two jets associated with  $W$  ( $C_{1,W}$ ).

As mentioned in the flavour tagging section, these b-jet and c-jet tag variables are useful to separate the signal events from the background events which do not have b quarks in the final states.

### 8.8.7 Particle number variables

The last four variables used in the MVA event selection are the particle number variables: the number of particles associated with  $H_{bb}$  ( $H_{bb}$ ), the number of particles associated with  $H_{WW^*}$  ( $H_{WW^*}$ ), the number of particles associated with  $W$  ( $W$ ), the number of particles associated with  $W^*$  ( $W^*$ ). These variables are effective to differentiate the signal events from the background events with fewer than six quarks in final states.

An optimal set of 32 variables are chosen for the best MVA performance, whilst no strong ( $> 80\%$ ) pair-wise correlation exists between any two variables.

Category	Variable
Invariant mass	$m_{H_{bb}}, m_{H_{WW^*}}, m_W, m_{HH}$
Energy and momentum	$E_{W^*}, E_{mis}, p_{TH_{bb}}, p_{TH_{WW^*}}, p_{TW}, p_{THH}$
Lab-frame angles	$\eta_{mis}, A_{H_{bb}}, A_W, A_{HH}$
Boosted-frame angles	$\cos(\theta_{H_{bb}}^*), \cos(\theta_{H_{WW^*}}^*), \cos(\theta_W^*), \cos(\theta_{W^*}^*), \cos(\theta_{HH}^*)$
Event shape	$ \mathbf{S} , -\ln(y_{23}), -\ln(y_{34}), -\ln(y_{45}), -\ln(y_{56})$
b and c tag	$B_{1,H_{bb}}, B_{2,H_{bb}}, B_{1,W}, B_{1,W^*}, C_{1,H_{bb}}, C_{1,W}$
Particle number	$N_{H_{bb}}, N_{H_{WW^*}}, N_W, N_{W^*}$

**Table 8.11:** Variables used in the MVA event selection for  $\sqrt{s} = 1.4$  TeV

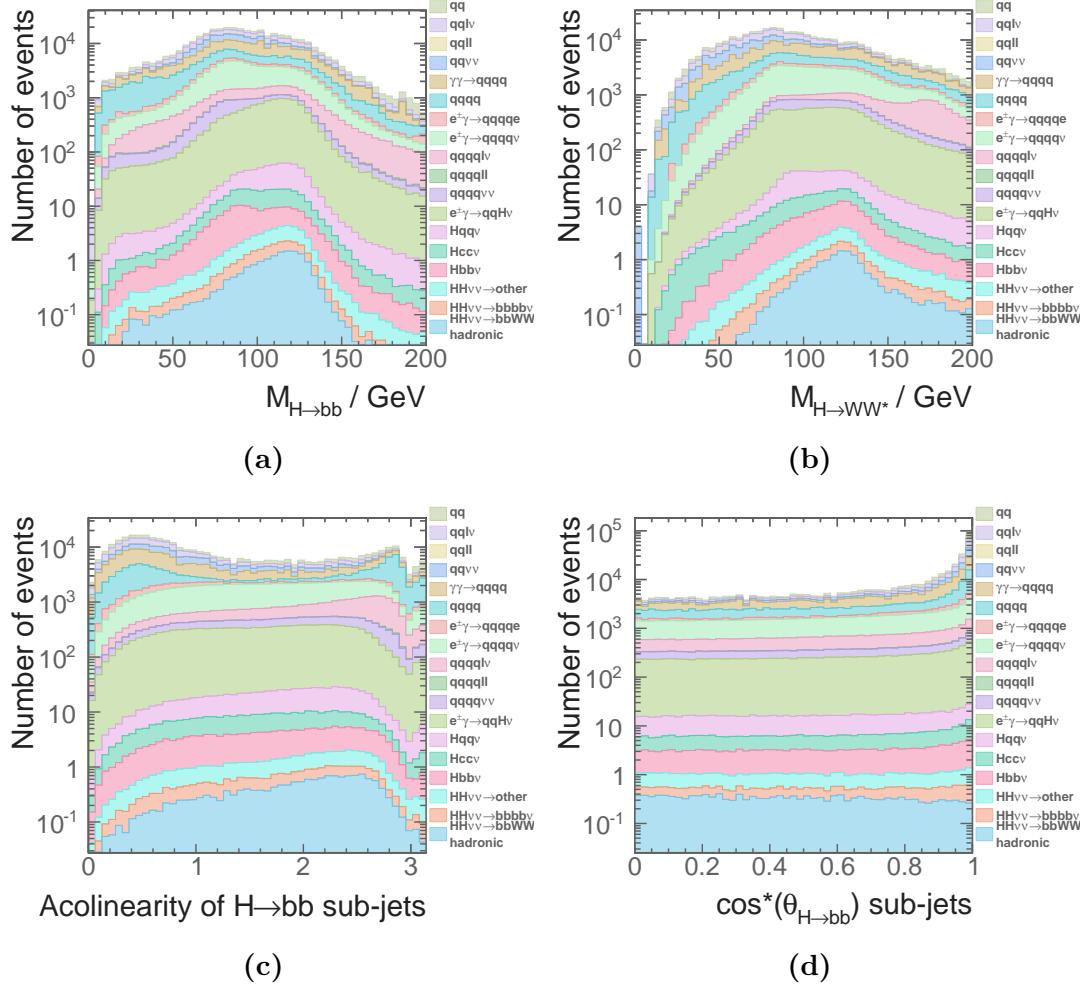
### 8.8.8 Cuts to aid the MVA

A set of cuts reduce the range of invariant masses variables in order to increase the effectiveness of the MVA event selection. Occasionally, extreme values of the invariant masses variables skew the distributions. Therefore by limiting the range of the variables, the MVA classifier could focus on the phase spaces with high event densities. The cuts require the invariant mass of the  $H_{bb} < 500$  GeV, the invariant mass of the  $H_{WW^*} < 800$  GeV, the invariant mass of the  $W < 200$  GeV, and the invariant mass of the double Higgs system  $< 1400$  GeV.

## 8.9 Multivariate analysis

After gathering information and applying pre-selection cuts, signal events are selected using the multivariate analysis (MVA) with Boosted Decision Tree classifier (BDT), as implemented in the TMVA [78]. The parameters of the boosted decision tree classifier were optimised and checked for overtraining, following the strategy outlined in section 4.5. Half of the events were used for training, and the other half used for testing and classifier optimisation. The optimised parameters are listed in table 8.12.

After dividing all events into a training set and a testing set, in the training stage of the MVA classifier, the training signal events are the hadronic  $W^+W^-$  decay of the  $HH \rightarrow b\bar{b}W^+W^-$  events in the training set. The training background events are all events without double higgs production in the training set. However, for the extraction of the  $g_{HHH}$  and  $g_{WWHH}$ , all events with double higgs production are sensitive to the couplings



**Figure 8.15:** Distributions of the four variables with highest discriminating power: a) the invariant mass of  $H_{bb}$ ; b) the invariant mas of  $H_{WW^*}$ ; c) the acolinearity of the two jets associated with  $H_{bb}$ ; and d) the opening angle of the two jets associated with  $H_{bb}$  in the decay rest frame of the  $H_{bb}$ . All plots assume an intergraded luminosity of  $1500 \text{ fb}^{-1}$  at  $\sqrt{s} = 1.4 \text{ TeV}$  after all pre-selection cuts applied before the MVA event selection.

Therefore, at the applying stage of the MVA classifier, all events in the testing set are used.

## 8.10 Signal selection results

Number of events passed the MVA event selection at  $\sqrt{s} = 1.4 \text{ TeV}$ , assuming a luminosity of  $1500 \text{ fb}^{-1}$ , are listed in table 8.13 for individual processes. A few background processes have non-zero events after the MVA event selection.  $e^+e^- \rightarrow q\bar{q}H\nu\bar{\nu}$  events are difficult to discard because its topology, one Higgs and neutrinos, is very similar to the signal

Parameter	Value
Depth of tree	4
Number of trees	4000
The minimum number of events in a node	0.25% of the total events
Boosting	adaptive boost
Learning rate of the adaptive boost	0.5
Metric for the optimal cuts	Gini Index
Bagging fraction	0.5
Number of bins per variables	40
End node output	$x \in [0, 1]$
DO <sub>P</sub> RESELECTION	yes

**Table 8.12:** Optimised parameters of the boosted decision tree classifier used in the MVA event selection. See section 4.5.6.1 for detailed explanations of variables.

event topology. Similarly,  $e^+e^- \rightarrow qqqq\ell\nu$  events can be confused with the signal events when the lepton is undetected in the forward region, or the energy of the lepton is too low to be tagged.  $e^+e^- \rightarrow qqqq\nu\bar{\nu}$  events can also have a similar topology to the signal events. Other background processes that are not discarded after the MVA are the electron–photon and photon–photon interactions with the same final states as the processes above.

Before interpreting the result for analysis at  $\sqrt{s} = 1.4$  TeV, the analyses at  $\sqrt{s} = 3$  TeV and the semi-leptonic channel of  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e$  are presented.

## 8.11 $\sqrt{s} = 3$ TeV analysis

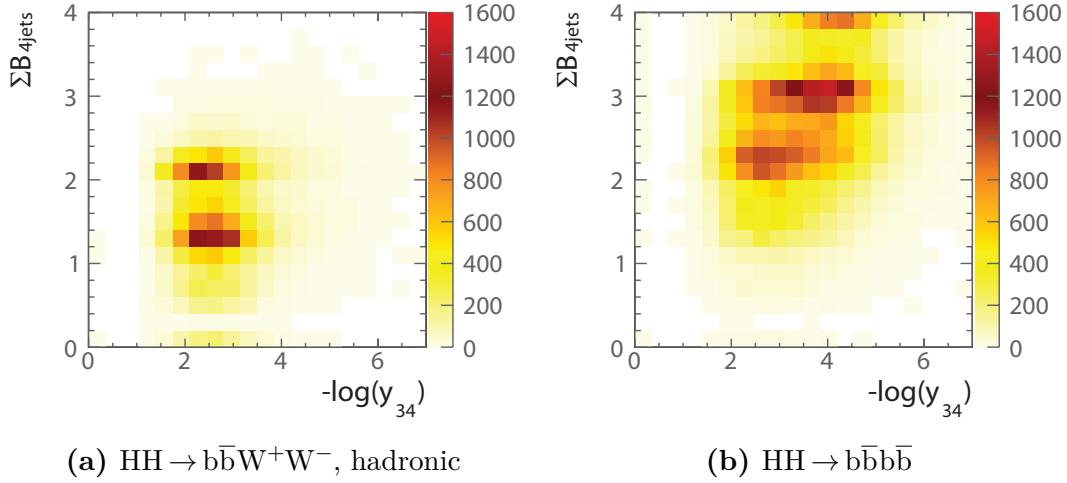
The hadronic  $W^+W^-$  decay of the  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu\bar{\nu}$  at  $\sqrt{s} = 3$  TeV analysis follows the same strategy as the analysis at  $\sqrt{s} = 1.4$  TeV. Lepton finding, jet pairing and flavouring tagging have been discussed in previous sections. The differences, which have not been mentioned, will be highlighted in this section.

Cross sections of used samples are listed in table 8.14. The mutually exclusive cuts to separate events into two independent sets are almost identical to the cuts used in the  $\sqrt{s} = 1.4$  TeV analysis. Figure 8.16 shows the sum of b-jet tag values, when the event is clustered into four jets, as a function of  $-\log(y_{34})$  for the hadronic  $W^+W^-$  decay in  $HH \rightarrow b\bar{b}W^+W^-$  and  $HH \rightarrow b\bar{b}b\bar{b}$  sub-channels. The optimised cuts are

$\sqrt{s} = 1.4 \text{ TeV}$	N	$\varepsilon_{\text{presel}}$	$\varepsilon_{\text{MVA}}$	$N_{\text{MVA}}$
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e$ , hadronic	27.9	59.8%	8.2%	1.29
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}b\bar{b}\nu_e\bar{\nu}_e$	67.6	15.4%	0.5%	0.05
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow \text{other}$	128.0	20.4%	1.7%	0.45
$e^+e^- \rightarrow q\bar{q}H\nu\bar{\nu}$	1290	39.5%	0.05%	0.29
$e^+e^- \rightarrow c\bar{c}H\nu\bar{\nu}$	540	31.6%	0.1%	0.16
$e^+e^- \rightarrow b\bar{b}H\nu\bar{\nu}$	465	24.7%	0.3%	0.37
$e^+e^- \rightarrow qq\bar{q}\bar{q}$	1867650	3.3%	-	-
$e^+e^- \rightarrow qq\bar{q}\bar{q}\ell\bar{\ell}$	93150	0.3%	-	-
$e^+e^- \rightarrow qq\bar{q}\bar{q}\ell\nu$	165600	9.8%	0.01%	2.06
$e^+e^- \rightarrow qq\bar{q}\bar{q}\nu\bar{\nu}$	34800	16.5%	0.002%	0.10
$e^+e^- \rightarrow qq$	6014250	0.8%	-	-
$e^+e^- \rightarrow qq\ell\nu$	6464550	0.9%	-	-
$e^+e^- \rightarrow qq\ell\bar{\ell}$	4088700	0.08%	-	-
$e^+e^- \rightarrow qq\nu\bar{\nu}$	1181550	4.0%	-	-
$e^\pm\gamma(BS) \rightarrow e^\pm qq\bar{q}\bar{q}$	2606625	0.3%	-	-
$e^\pm\gamma(EPA) \rightarrow e^\pm qq\bar{q}\bar{q}$	861000	0.3%	-	-
$e^\pm\gamma(BS) \rightarrow \nu qq\bar{q}\bar{q}$	178987.5	25.7%	0.005%	2.05
$e^\pm\gamma(EPA) \rightarrow \nu qq\bar{q}\bar{q}$	52050	12.5%	0.004%	0.27
$e^\pm\gamma(BS) \rightarrow qqH\nu$	35437.5	30.7%	0.02%	2.16
$e^\pm\gamma(EPA) \rightarrow qqH\nu$	10170.0	16.1%	0.06%	0.95
$\gamma(BS)\gamma(BS) \rightarrow qq\bar{q}\bar{q}$	2054951.5	0.2%	-	-
$\gamma(BS)\gamma(EPA) \rightarrow qq\bar{q}\bar{q}$	4521037.5	0.4%	-	-
$\gamma(EPA)\gamma(BS) \rightarrow qq\bar{q}\bar{q}$	4539150.0	0.3%	-	-
$\gamma(EPA)\gamma(EPA) \rightarrow qq\bar{q}\bar{q}$	1129500.0	0.3%	-	-

**Table 8.13:** List of signal and background events with selection efficiency and number of events at  $\sqrt{s} = 1.4 \text{ TeV}$ , assuming an integrated of  $1500 \text{ fb}^{-1}$ . The number of events (N), the selection efficiencies of pre-selection cuts ( $\varepsilon_{\text{presel}}$ ), the selection efficiencies of the MVA event selection after pre-selection cuts ( $\varepsilon_{\text{MVA}}$ ), and the number of events after the MVA event selection ( $N_{\text{MVA}}$ ) are shown. The entries marked with “-” represent numbers less than 0.01. q can be u, d, s, b or t.

$\Sigma B_{4\text{jets}} < 2.3$ ,  $-\log(y_{34}) < 3.6$ . The selection efficiencies of evens after lepton veto, the mutually exclusive cuts and the jet pairing for individual processes are shown in table C.1.



**Figure 8.16:** The two-dimensional distribution of sum of b-jet tag values against  $-\log(y_{34})$  for: a) hadronic  $W^+W^-$  decay of  $\text{HH} \rightarrow b\bar{b}W^+W^-$ ; and b)  $\text{HH} \rightarrow b\bar{b}b\bar{b}$  events at  $\sqrt{s} = 3$  TeV. The sum of b-jet tag values is calculated for the cases where events are clustered into four jets.

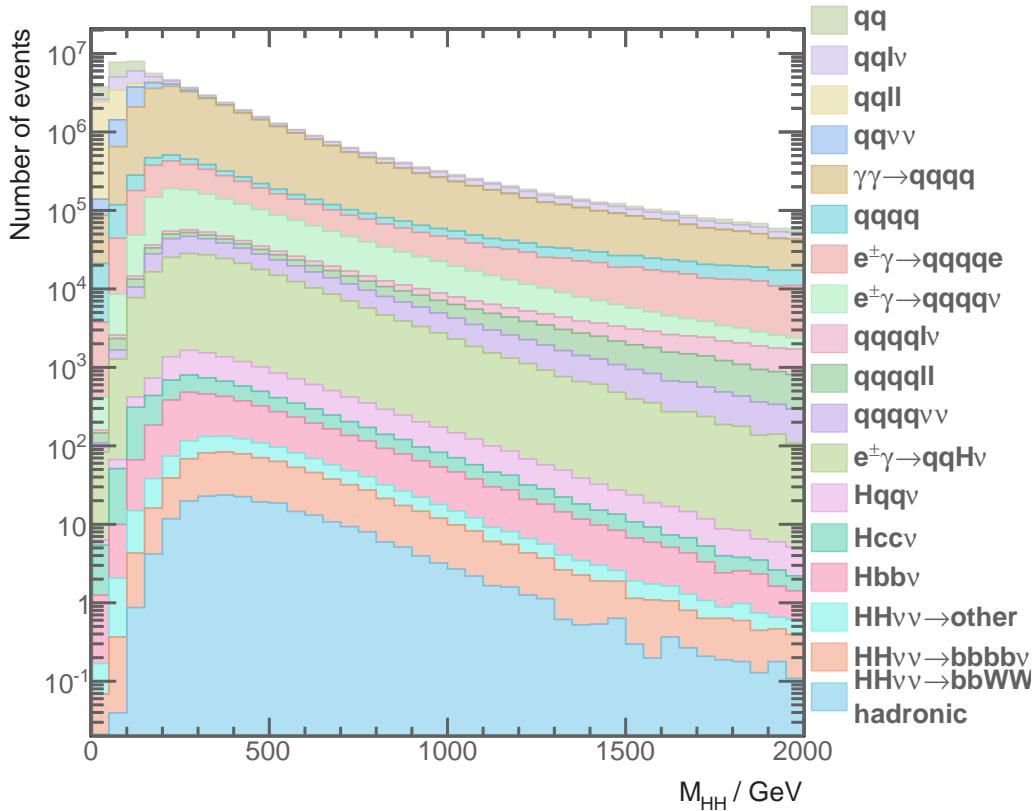
The pre-selection cuts at  $\sqrt{s} = 3$  TeV use the same cut on  $m_{\text{HH}}$ . The cut on b-jet tag is different because the performance of flavour tagging is worse at  $\sqrt{s} = 3$  TeV in comparison to the performance at  $\sqrt{s} = 1.4$  TeV. Figure 8.18 shows the distribution of the highest b-jet tag value, where the cut above 0.7 helps to reduce background events with no b quarks in final states. Figure 8.17 shows the distribution of the invariant mass of the two Higgs system, where the cut above 150 GeV is effective against samples with two-quark final states. The fractions of events passing each pre-selection cut for individual processes are listed in table C.2.

The cuts to aid the MVA at  $\sqrt{s} = 3$  TeV are largely the same as the ones at  $\sqrt{s} = 1.4$  TeV, apart from the difference on the cut of the invariant mass of HH due to a higher centre-of-mass energy. The cuts are the invariant mass of the  $H_{bb} < 500$  GeV, the invariant mass of the  $H_{WW^*} < 800$  GeV, the invariant mass of the  $W < 200$  GeV, and the invariant mass of the double Higgs system  $< 3000$  GeV.

The same set of variables are used in the MVA as in the analysis at  $\sqrt{s} = 1.4$  TeV. The optimised parameters of the Boosted Decision Tree classifier are the same. The efficiencies of the MVA event selection and the numbers of events after the MVA event selection are listed in table 8.15. Background processes that are dominant after the MVA

Process	$\sigma(\sqrt{s} = 3 \text{ TeV}) / \text{fb}$
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$	0.588
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-$ , hadronic	0.07
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}b\bar{b}\nu_e\bar{\nu}_e$	0.19
$e^+e^- \rightarrow HH \rightarrow \text{others}$	0.34
$e^+e^- \rightarrow q\bar{q}H\nu\bar{\nu}$	3.06
$e^+e^- \rightarrow c\bar{c}H\nu\bar{\nu}$	1.15
$e^+e^- \rightarrow b\bar{b}H\nu\bar{\nu}$	1.78
$e^+e^- \rightarrow qqqq$	546.5*
$e^+e^- \rightarrow qqqq\ell\ell$	169.3*
$e^+e^- \rightarrow qqqq\ell\nu$	106.6*
$e^+e^- \rightarrow qqqq\nu\bar{\nu}$	71.5*
$e^+e^- \rightarrow qq$	2948.9
$e^+e^- \rightarrow qq\ell\nu$	5561.1
$e^+e^- \rightarrow qq\ell\ell$	3319.6
$e^+e^- \rightarrow qq\nu\nu$	1317.5
$e^\pm\gamma(\text{BS}) \rightarrow e^\pm qqqq$	2536.3*
$e^\pm\gamma(\text{EPA}) \rightarrow e^\pm qqqq$	575.7*
$e^\pm\gamma(\text{BS}) \rightarrow \nu qqqq$	524.8*
$e^\pm\gamma(\text{EPA}) \rightarrow \nu qqqq$	108.4*
$e^\pm\gamma(\text{BS}) \rightarrow qqH\nu$	117.1*
$e^\pm\gamma(\text{EPA}) \rightarrow qqH\nu$	22.4*
$\gamma(\text{BS})\gamma(\text{BS}) \rightarrow qqqq$	13050.3*
$\gamma(\text{BS})\gamma(\text{EPA}) \rightarrow qqqq$	2420.6*
$\gamma(\text{EPA})\gamma(\text{BS}) \rightarrow qqqq$	2423.1*
$\gamma(\text{EPA})\gamma(\text{EPA}) \rightarrow qqqq$	402.7*

**Table 8.14:** List of signal and background samples used in the double Higgs analysis with the corresponding cross sections at  $\sqrt{s} = 3 \text{ TeV}$ .  $q$  can be  $u$ ,  $d$ ,  $s$ ,  $b$  or  $t$ . Unless specified,  $q$ ,  $\ell$  and  $\nu$  represent either particles or the corresponding anti-particles.  $\gamma$  (BS) represents a real photon from beamstrahlung (BS).  $\gamma$  (EPA) represents a “quasi-real” photon, simulated with the Equivalent Photon Approximation. For processes labelled with \*, events are generated with the invariant mass of the total momenta of all quarks above 50 GeV.



**Figure 8.17:** Distributions of the invariant mass of the two Higgs system for  $\sqrt{s} = 3 \text{ TeV}$ , assuming an intergraded luminosity of  $2000 \text{ fb}^{-1}$ .

event selection are almost identical to those at  $\sqrt{s} = 1.4 \text{ TeV}$ . Hence see section 8.10 for discussion.

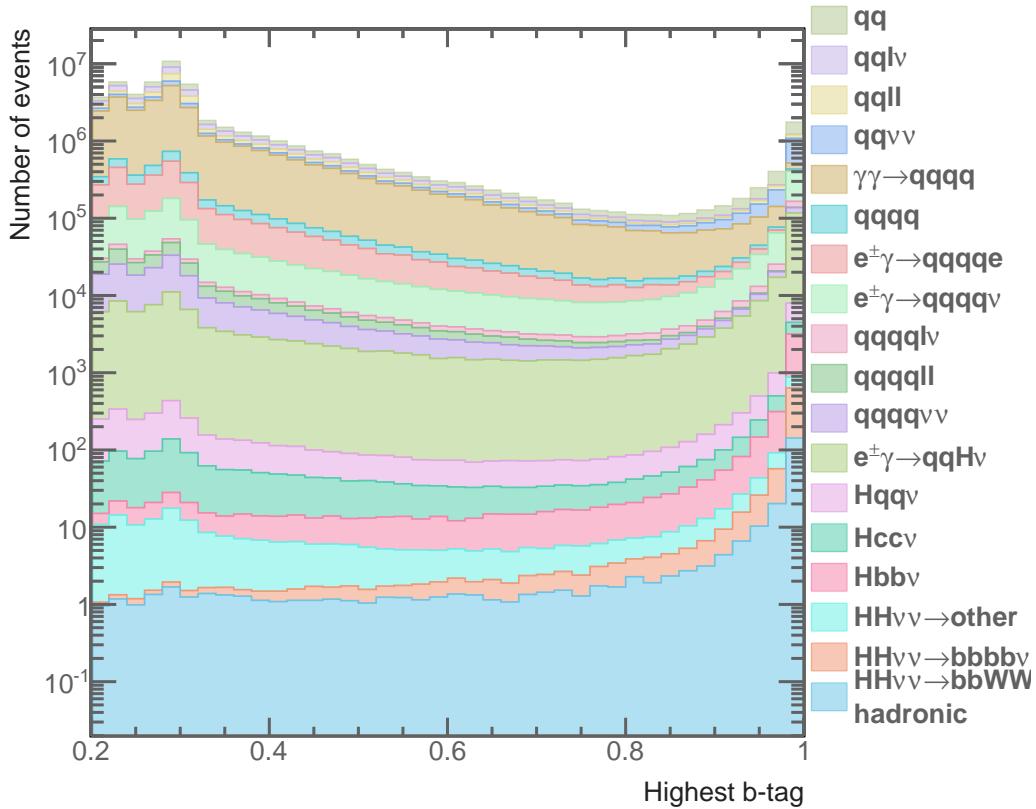
## 8.12 Semi-leptonic decay at $\sqrt{s} = 3 \text{ TeV}$ analysis

The final analysis is the semi-leptonic  $W^+W^-$  decay of  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu\bar{\nu}$  at  $\sqrt{s} = 3 \text{ TeV}$ . The semi-leptonic decay analysis at  $\sqrt{s} = 1.4 \text{ TeV}$  was also performed. However there are not enough signal events to have a meaningful discussion for the analysis at  $\sqrt{s} = 1.4 \text{ TeV}$ . Hence, only the semi-leptonic decay analysis at  $\sqrt{s} = 3 \text{ TeV}$  is presented.

The strategy of the semi-leptonic decay analysis is very similar to the hadronic decay analysis. The main difference are that there is one lepton in the final state and the

$\sqrt{s} = 3 \text{ TeV}$	N	$\varepsilon_{\text{presel}}$	$\varepsilon_{\text{MVA}}$	$N_{\text{MVA}}$
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e$ , hadronic	146.0	61.7%	11.6%	9.89
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}b\bar{b}\nu_e\bar{\nu}_e$	355.0	18.8%	1.5%	1.05
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow \text{other}$	675.0	20.0%	3.6%	4.51
$e^+e^- \rightarrow q\bar{q}H\nu\bar{\nu}$	6120	36.0%	0.4%	9.42
$e^+e^- \rightarrow c\bar{c}H\nu\bar{\nu}$	2300	26.3%	0.5%	3.13
$e^+e^- \rightarrow b\bar{b}H\nu\bar{\nu}$	3560	25.8%	1.2%	6.82
$e^+e^- \rightarrow qq\bar{q}\bar{q}$	1093000	1.4%	0.01%	1.43
$e^+e^- \rightarrow qq\bar{q}\bar{q}\ell\ell$	338600	0.6%	-	-
$e^+e^- \rightarrow qq\bar{q}\bar{q}\ell\nu$	213200	7.3%	0.05%	8.35
$e^+e^- \rightarrow qq\bar{q}\bar{q}\nu\bar{\nu}$	143000	9.0%	0.05%	6.35
$e^+e^- \rightarrow qq$	5897800	1.4%	-	-
$e^+e^- \rightarrow qq\ell\nu$	11121800	0.1%	-	-
$e^+e^- \rightarrow qq\ell\ell$	6639200	0.4%	-	-
$e^+e^- \rightarrow qq\nu\nu$	2635000	3.1%	-	-
$e^\pm\gamma(\text{BS}) \rightarrow e^\pm qq\bar{q}\bar{q}$	4007354	0.7%	-	-
$e^\pm\gamma(\text{EPA}) \rightarrow e^\pm qq\bar{q}\bar{q}$	1151200	0.4%	-	-
$e^\pm\gamma(\text{BS}) \rightarrow \nu qq\bar{q}\bar{q}$	829184	16.4%	0.04%	61.0
$e^\pm\gamma(\text{EPA}) \rightarrow \nu qq\bar{q}\bar{q}$	216800	7.6%	0.04%	6.0
$e^\pm\gamma(\text{BS}) \rightarrow qqH\nu$	185018	30.2%	0.2%	121.7
$e^\pm\gamma(\text{EPA}) \rightarrow qqH\nu$	46800.0	15.3%	0.2%	18.1
$\gamma(\text{BS})\gamma(\text{BS}) \rightarrow qq\bar{q}\bar{q}$	18009414	1.6%	-	-
$\gamma(\text{BS})\gamma(\text{EPA}) \rightarrow qq\bar{q}\bar{q}$	3824548	1.0%	-	-
$\gamma(\text{EPA})\gamma(\text{BS}) \rightarrow qq\bar{q}\bar{q}$	3828498	1.0%	-	-
$\gamma(\text{EPA})\gamma(\text{EPA}) \rightarrow qq\bar{q}\bar{q}$	805400	0.6%	-	-

**Table 8.15:** List of signal and background events with selection efficiency and number of events at  $\sqrt{s} = 3 \text{ TeV}$ , assuming an integrated luminosity of  $2000 \text{ fb}^{-1}$ . The number of events (N), the selection efficiencies of pre-selection cuts ( $\varepsilon_{\text{presel}}$ ), the selection efficiencies of the MVA event selection after pre-selection cuts ( $\varepsilon_{\text{MVA}}$ ), and the number of events after the MVA event selection ( $N_{\text{MVA}}$ ) are shown. The entries marked with “-” represent numbers less than 0.01. q can be u, d, s, b or t. Unless specified, q,  $\ell$  and  $\nu$  represent either particles or the corresponding anti-particles.



**Figure 8.18:** Distributions of the highest b-jet tag value for  $\sqrt{s} = 3 \text{ TeV}$ , assuming an integrated luminosity of  $2000 \text{ fb}^{-1}$ .

final state has four quarks instead of six.  $H_{bb}$  and  $W$  can not be reconstructed due to the leptonic decay of one of the  $W$ . Hence, the signal events are selected when there is one identified lepton using the same lepton finding processors. The jet reconstruction parameters are the same as hadronic decay analysis at the  $\sqrt{s} = 3 \text{ TeV}$ . There are no mutually exclusive cuts since there is no semi-leptonic analysis in the  $HH \rightarrow b\bar{b}b\bar{b}$  analysis.

The pre-selection cuts are similar to the cuts in the hadronic analysis. The invariant mass of the double Higgs system is required to be above  $150 \text{ GeV}$ . The highest b-jet tag value is higher than  $0.2$ . The transverse momentum of the double Higgs system is higher than  $30 \text{ GeV}$ .

Variables used in the MVA classifier, listed in table 8.16, belong to a reduced set of the variables used in the hadronic decay analysis, as  $H_{bb}$  and  $W$  can not be reconstructed in the semi-hadronic decay analysis. For the same reason, the cuts to aid the MVA are

reduced to the invariant mass of  $H_{bb} < 500 \text{ GeV}$  and the invariant mass of the double Higgs system  $< 3000 \text{ GeV}$ .

Figure 8.17 lists the selection efficiency and number of events after the MVA event selection for individual processes at  $\sqrt{s} = 3 \text{ TeV}$ , assuming a luminosity of  $2000 \text{ fb}^{-1}$ . Almost all background processes are non-zero after the MVA event selection. Nevertheless, dominant background processes are almost identical to the hadronic decay analysis at  $\sqrt{s} = 3 \text{ TeV}$ . Hence discussion of the MVA event selection is provided in section 8.10.

Category	Variable
Invariant mass	$m_{H_{bb}}, m_W, m_{HH}$
Energy and momentum	$E_{mis}, p_{TH_{bb}}, p_{TW}, p_{THH}$
Lab-frame angles	$\theta_{mis}, A_{H_{bb}}, A_W, A_{HH}$
Boosted-frame frames	$\cos(\theta_{H_{bb}}^*), \cos(\theta_{HH}^*)$
Event shape	$ S , -\ln(y_{23}), -\ln(y_{34}), -\ln(y_{45}), -\ln(y_{56})$
b and c tag	$B_{1,H_{bb}}, B_{2,H_{bb}}, B_{1,W}, C_{1,H_{bb}}, C_{1,W}$
Particle number	$N_{H_{bb}}, N_W$

**Table 8.16:** Variables used in the MVA event selection for the semi-leptonic  $W^+W^-$  decay of  $HH \rightarrow b\bar{b}W^+W^-$  analysis at  $\sqrt{s} = 3 \text{ TeV}$ .

## 8.13 Result interpretation

The numbers of signal events and background events and significance after the MVA event selection for analyses at the  $\sqrt{s} = 1.4 \text{ TeV}$  and  $\sqrt{s} = 3 \text{ TeV}$  are listed in table 8.18. The significance is defined as the number of signal events divided by the square root of the sum of the signal and background events. The  $e^+e^-$  collision at  $\sqrt{s} = 1.4 \text{ TeV}$  assumes an integrated luminosity of  $1500 \text{ fb}^{-1}$ , whilst that of the  $\sqrt{s} = 3 \text{ TeV}$  assumes an integrated luminosity of  $2000 \text{ fb}^{-1}$ . For the hadronic  $W^+W^-$  decay of  $HH \rightarrow b\bar{b}W^+W^-$  analysis at  $\sqrt{s} = 1.4 \text{ TeV}$ , the numbers of signal and background events after the MVA selection are 1.79 and 8.41 respectively. For the hadronic analysis at  $\sqrt{s} = 3 \text{ TeV}$ , the respective numbers of the signal and background events after the MVA selection are 15.45 and 242.28. For the semi-leptonic  $W^+W^-$  decay of  $HH \rightarrow b\bar{b}W^+W^-$  analysis at  $\sqrt{s} = 3 \text{ TeV}$ , the numbers of the signal and background events after the MVA selection are 31.24 and 3612.39 respectively.

$\sqrt{s} = 3 \text{ TeV}$	N	$\varepsilon_{\text{presel}}$	$\varepsilon_{\text{MVA}}$	$N_{\text{MVA}}$
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e$ , semi-leptonic	96.8	44.6%	21.9%	13.11
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}b\bar{b}\nu_e\bar{\nu}_e$	355.0	13.3%	10.9%	5.38
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow \text{other}$	724.2	13.1%	13.6%	12.75
$e^+e^- \rightarrow q\bar{q}H\nu\bar{\nu}$	6120	7.4%	13.7%	62.63
$e^+e^- \rightarrow c\bar{c}H\nu\bar{\nu}$	2300	6.3%	12.1%	17.10
$e^+e^- \rightarrow b\bar{b}H\nu\bar{\nu}$	3560	15.9%	5.1%	18.03
$e^+e^- \rightarrow qqqq$	1093000	0.6%	0.2%	15.04
$e^+e^- \rightarrow qqqq\ell\ell$	338600	1.0%	0.06%	1.85
$e^+e^- \rightarrow qqqq\ell\nu$	213200	27.6%	0.5%	270.33
$e^+e^- \rightarrow qqqq\nu\bar{\nu}$	143000	1.9%	1.6%	43.78
$e^+e^- \rightarrow qq$	5897800	0.4%	0.3%	60.82
$e^+e^- \rightarrow qq\ell\nu$	11121800	0.3%	0.08%	21.24
$e^+e^- \rightarrow qq\ell\ell$	6639200	0.6%	0.2%	84.14
$e^+e^- \rightarrow qq\nu\nu$	2635000	0.4%	0.9%	92.55
$e^\pm\gamma(\text{BS}) \rightarrow e^\pm qqqq$	4007354	1.2%	-	-
$e^\pm\gamma(\text{EPA}) \rightarrow e^\pm qqqq$	1151200	1.1%	-	-
$e^\pm\gamma(\text{BS}) \rightarrow \nu qqqq$	829184	3.6%	1.5%	452.45
$e^\pm\gamma(\text{EPA}) \rightarrow \nu qqqq$	216800	11.0%	0.9%	200.65
$e^\pm\gamma(\text{BS}) \rightarrow qqH\nu$	185018	7.9%	10.4%	1521.93
$e^-\gamma(\text{EPA}) \rightarrow qqH\nu$	46800	22.8%	7.1%	750.85
$\gamma(\text{BS})\gamma(\text{BS}) \rightarrow qqqq$	18009414	0.4%	-	-
$\gamma(\text{BS})\gamma(\text{EPA}) \rightarrow qqqq$	3824548	1.0%	-	-
$\gamma(\text{EPA})\gamma(\text{BS}) \rightarrow qqqq$	3828498	1.0%	0.08%	28.85
$\gamma(\text{EPA})\gamma(\text{EPA}) \rightarrow qqqq$	805400	1.1%	-	-

**Table 8.17:** List of signal and background events with selection efficiency and number of events at  $\sqrt{s} = 3 \text{ TeV}$  for semi-leptonic  $W^+W^-$  decay of  $HH \rightarrow b\bar{b}W^+W^-$  analysis, assuming an integrated luminosity of  $2000 \text{ fb}^{-1}$ . The number of events (N), the selection efficiencies of pre-selection cuts ( $\varepsilon_{\text{presel}}$ ), the selection efficiencies of MVA after pre-selection cuts ( $\varepsilon_{\text{MVA}}$ ), and the number of events after MVA ( $N_{\text{MVA}}$ ) are shown. The entries marked with “-” represent numbers less than 0.01. q can be u, d, s, b or t. Unless specified, q,  $\ell$  and  $\nu$  represent either particles or the corresponding anti-particles.

Process	$N_S$	$N_B$	$\frac{N_S}{\sqrt{N_S + N_B}}$
$HH \rightarrow b\bar{b}W^+W^-$ , hadronic, $\sqrt{s} = 1.4$ TeV	1.79	8.41	0.56
$HH \rightarrow b\bar{b}W^+W^-$ , hadronic, $\sqrt{s} = 3$ TeV	15.45	242.28	0.96
$HH \rightarrow b\bar{b}W^+W^-$ , semi-leptonic, $\sqrt{s} = 3$ TeV	31.24	3612.39	0.52

**Table 8.18:** Number of signal ( $N_S$ ) and background ( $N_B$ ) events, and significance ( $N_S/\sqrt{N_S + N_B}$ ) after MVA event selections for  $HH \rightarrow b\bar{b}W^+W^-$  analyses at  $\sqrt{s} = 1.4$  TeV and 3 TeV. The  $e^+e^-$  collision at  $\sqrt{s} = 1.4$  TeV assumes an integrated luminosity of  $1500 \text{ fb}^{-1}$ , whilst that of  $\sqrt{s} = 3$  TeV assumes an integrated luminosity of  $2000 \text{ fb}^{-1}$ .

The expected uncertainties on the measurement of the double Higgs production cross sections are estimated to be the inverse of the significance [3]:

$$\frac{\Delta [\sigma (HH\nu_e\bar{\nu}_e)]}{\sigma (HH\nu_e\bar{\nu}_e)} = \begin{cases} 179\%, & \text{at } \sqrt{s} = 1.4 \text{ TeV}, \\ 92\%, & \text{at } \sqrt{s} = 3 \text{ TeV}, \end{cases} \quad (8.11)$$

The expected uncertainty at  $\sqrt{s} = 3$  TeV combines the results from analyses for the hadronic and semi-leptonic decay sub-channels.

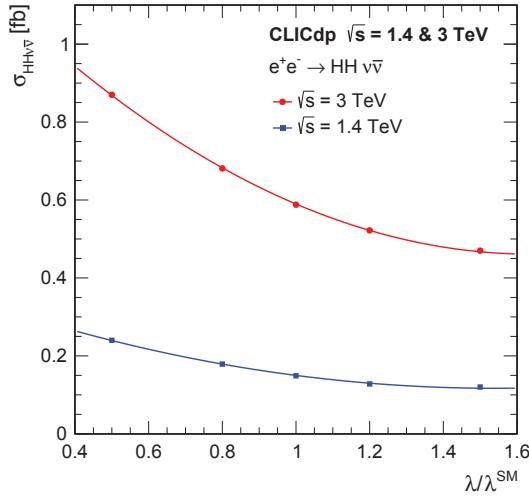
The Higgs trilinear self coupling  $g_{HHH}$  is related to the double Higgs production cross section via [23]:

$$\frac{\Delta g_{HHH}}{g_{HHH}} \approx \kappa \cdot \frac{\Delta [\sigma (HH\nu_e\bar{\nu}_e)]}{\sigma (HH\nu_e\bar{\nu}_e)}, \quad (8.12)$$

The coefficient  $\kappa$  can be determined by parameterising the  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  cross section as a function of the coupling  $g_{HHH}$ . Figure 8.19 shows the  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  cross sections as a function of the coupling  $g_{HHH}$  for  $\sqrt{s} = 1.4$  TeV and  $\sqrt{s} = 3$  TeV. At the SM  $g_{HHH}$  value, the coefficient  $\kappa$  is 1.22 at  $\sqrt{s} = 1.4$  TeV, and 1.47 at  $\sqrt{s} = 3$  TeV.

The uncertainties on measurement of the Higgs trilinear self coupling,  $g_{HHH}$ , from  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e$  analyses, are hence obtained via equation 8.12:

$$\frac{\Delta g_{HHH}}{g_{HHH}} \approx \begin{cases} 218\%, & \text{at } \sqrt{s} = 1.4 \text{ TeV}, \\ 135\%, & \text{at } \sqrt{s} = 3 \text{ TeV}. \end{cases} \quad (8.13)$$



**Figure 8.19:** Cross section for the  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  process as a function of the ratio  $\lambda/\lambda_{SM}$  at  $\sqrt{s} = 1.4 \text{ TeV}$  and  $3 \text{ TeV}$ , taken from [23]. Here  $\lambda$  is the Higgs trilinear self coupling,  $g_{HHH}$ .

Since the leading-order Feynman diagrams for the double Higgs boson production include a t-channel WW-fusion process, the cross section of the double Higgs production can be enhanced by using a polarised electron beam. For a electron beam polarisation of  $P(e^-) = 80\%$ , the uncertainties of the coupling  $g_{HHH}$  become:

$$\frac{\Delta g_{HHH}}{g_{HHH}} \approx \begin{cases} 163\%, & \text{at } \sqrt{s} = 1.4 \text{ TeV}, \\ 97\%, & \text{at } \sqrt{s} = 3 \text{ TeV}. \end{cases} \quad (8.14)$$

When the analyses at both  $\sqrt{s} = 1.4 \text{ TeV}$  and  $\sqrt{s} = 3 \text{ TeV}$  are combined, the uncertainty of the coupling  $g_{HHH}$  improves to 99% with the unpolarised beam, and to 87% with the polarised beam of  $P(e^-) = 80\%$ .

When the analyses for  $HH \rightarrow b\bar{b}W^+W^-$  and  $HH \rightarrow b\bar{b}b\bar{b}$  sub-channels are combined, the expected uncertainties on the double Higgs production cross section measurements are:

$$\frac{\Delta [\sigma(HH\nu_e\bar{\nu}_e)]}{\sigma(HH\nu_e\bar{\nu}_e)} = \begin{cases} 44\%, & \text{at } \sqrt{s} = 1.4 \text{ TeV}, \\ 20\%, & \text{at } \sqrt{s} = 3 \text{ TeV}, \end{cases} \quad (8.15)$$

This translates to uncertainties on the measurement of the Higgs trilinear self coupling  $g_{\text{HHH}}$ , via equation 8.12, with unpolarised beams:

$$\frac{\Delta g_{\text{HHH}}}{g_{\text{HHH}}} \approx \begin{cases} 54\%, & \text{at } \sqrt{s} = 1.4 \text{ TeV}, \\ 29\%, & \text{at } \sqrt{s} = 3 \text{ TeV}. \end{cases} \quad (8.16)$$

## 8.14 Simultaneous couplings extraction

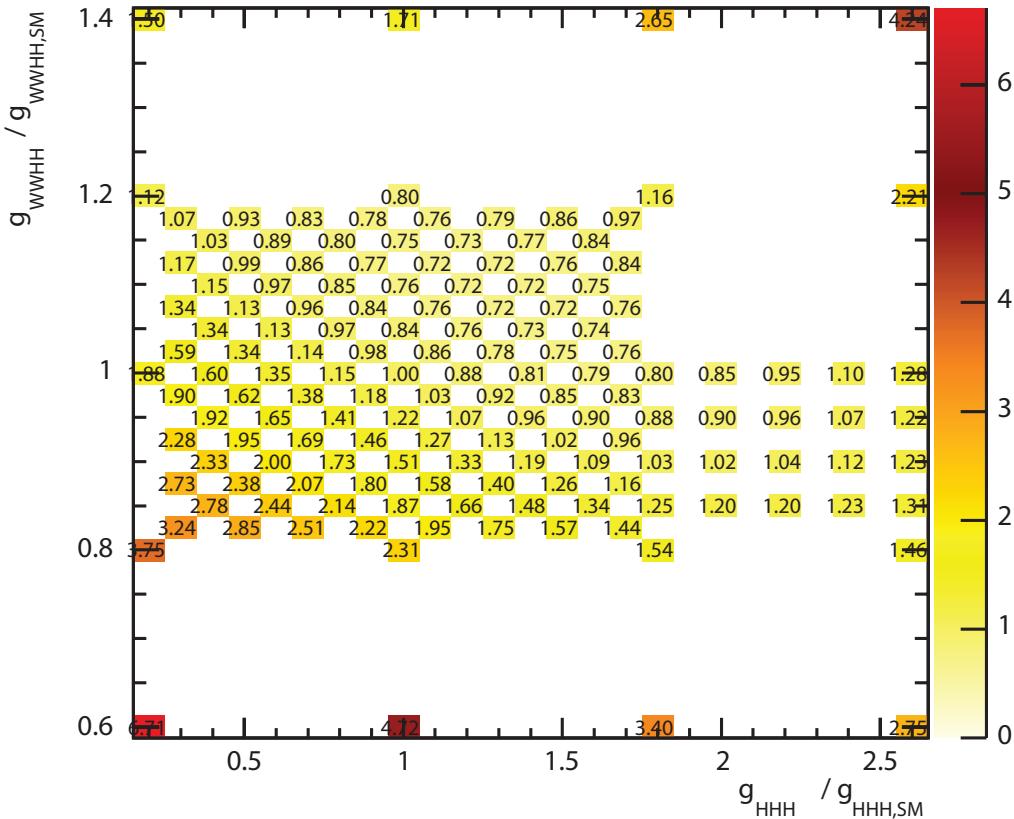
The study of the double Higgs production via  $W^+W^-$  fusion can probe the Higgs trilinear self coupling,  $g_{\text{HHH}}$ , and quartic coupling,  $g_{\text{WWHH}}$ . A two-dimensional  $g_{\text{HHH}}$  and  $g_{\text{WWHH}}$  couplings extraction is performed using the results of the hadronic analysis at  $\sqrt{s} = 3 \text{ TeV}$ . The integrated luminosity in this section is assumed to be  $3000 \text{ fb}^{-1}$  at  $\sqrt{s} = 3 \text{ TeV}$  to reflect the updated CLIC running scenario [108]. A simple scaling is applied to the number of events after the MVA event selection for the analysis at  $\sqrt{s} = 3 \text{ TeV}$  to adapt to the change in the luminosity.

The  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  events with non-SM  $g_{\text{HHH}}$  and  $g_{\text{WWHH}}$  couplings were generated and reconstructed. The normalised cross sections of the  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  as a function of  $g_{\text{HHH}}$  and  $g_{\text{WWHH}}$  are shown in figure 8.20. Around the SM coupling values, the cross section increases with the decrease of  $g_{\text{HHH}}$  and with the increase of  $g_{\text{WWHH}}$ .

These generated non-SM coupling  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  events went through the analysis chain described in this chapter with the same cuts and the same MVA classifier applied. The same of background events from the  $\sqrt{s} = 3 \text{ TeV}$  analysis were added to the  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  events with non-SM couplings. Figure 8.21 shows the signal significance after the MVA event selection of the double Higgs events with hadronic  $W^+W^-$  decay of  $HH \rightarrow b\bar{b}W^+W^-$  sub-channel as a function of  $g_{\text{HHH}}$  and  $g_{\text{WWHH}}$ .

Two kinematic variables that are sensitive to the change of the couplings, motivated in see section 2.8, are used for the extraction of the couplings: the invariant mass of the two Higgs system,  $m_{\text{HH}}$ , and the scalar sum of the two Higgs transverse momentum,  $H_T$ .

Events are subsequently binned using the kinematic variables. Two bins in  $H_T$  are obtained by dividing the  $H_T$  distribution at 200 GeV. Four bins in  $m_{\text{HH}}$  are obtained by dividing the  $m_{\text{HH}}$  distribution at 400, 560, and 720 GeV. This results in events being divided into eight kinematic bins.



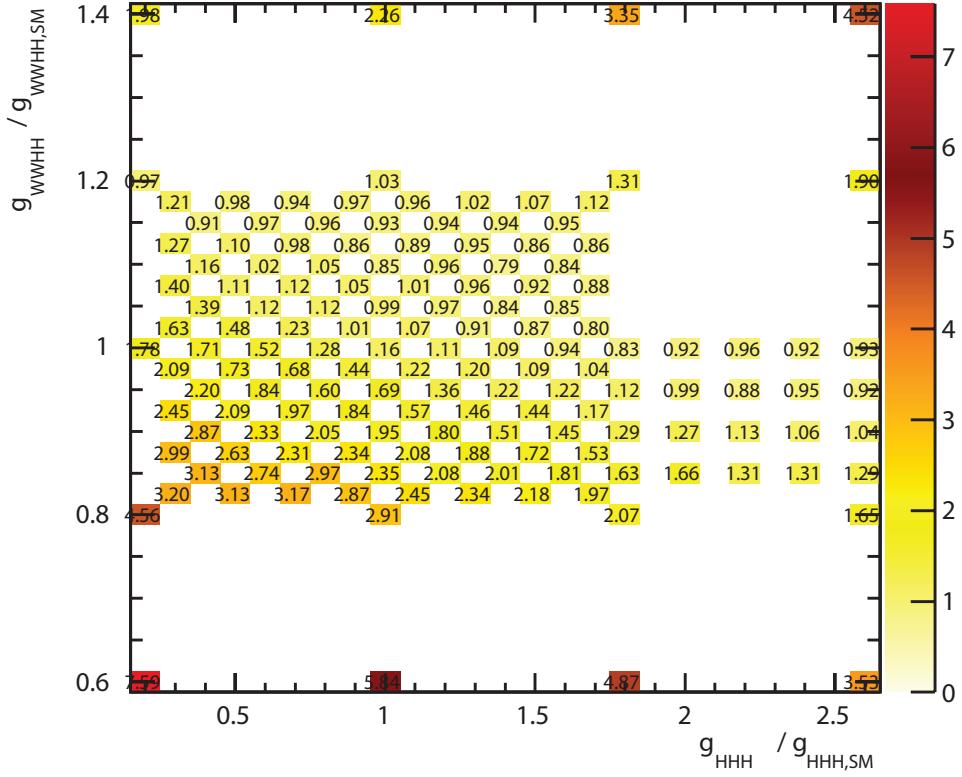
**Figure 8.20:** Normalised cross section for the  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  process as a function of the  $g_{HHH}/g_{HHH}^{SM}$  and  $g_{WWHH}/g_{WWHH}^{SM}$  at  $\sqrt{s} = 3$  TeV. All cross sections are normalised to the cross section at the SM coupling value.

To access the change in the  $m_{HH}$  and  $H_T$  distributions for non-SM coupling samples, the  $\chi^2$  function is used:

$$\chi^2 = \sum_i^8 \frac{(N_i - N_{i,observed})^2}{N_i}, \quad (8.17)$$

where  $N_i$  is the number of expected events in the kinematic bin  $i$  in a non-SM coupling sample; and  $N_{i,observed}$  is the number of observed event in the kinematic bin  $i$ . Here the observed sample is assumed to be the SM coupling sample. The  $\chi^2$  is summed over all kinematic bins. By construction, the SM coupling sample has a  $\chi^2$  of 0. Figure 8.22 shows the  $\chi^2$  as a function of  $g_{HHH}$  and  $g_{WWHH}$  using samples from two sub-channels: hadronic  $W^+W^-$  decay of  $HH \rightarrow b\bar{b}W^+W^-$  and  $HH \rightarrow b\bar{b}b\bar{b}$ . The  $\chi^2$  values for the  $HH \rightarrow b\bar{b}b\bar{b}$  sub-channel are larger because the signal significance of  $HH \rightarrow b\bar{b}b\bar{b}$  sub-channel is much higher than that of the hadronic  $W^+W^-$  decay of  $HH \rightarrow b\bar{b}W^+W^-$  channel.

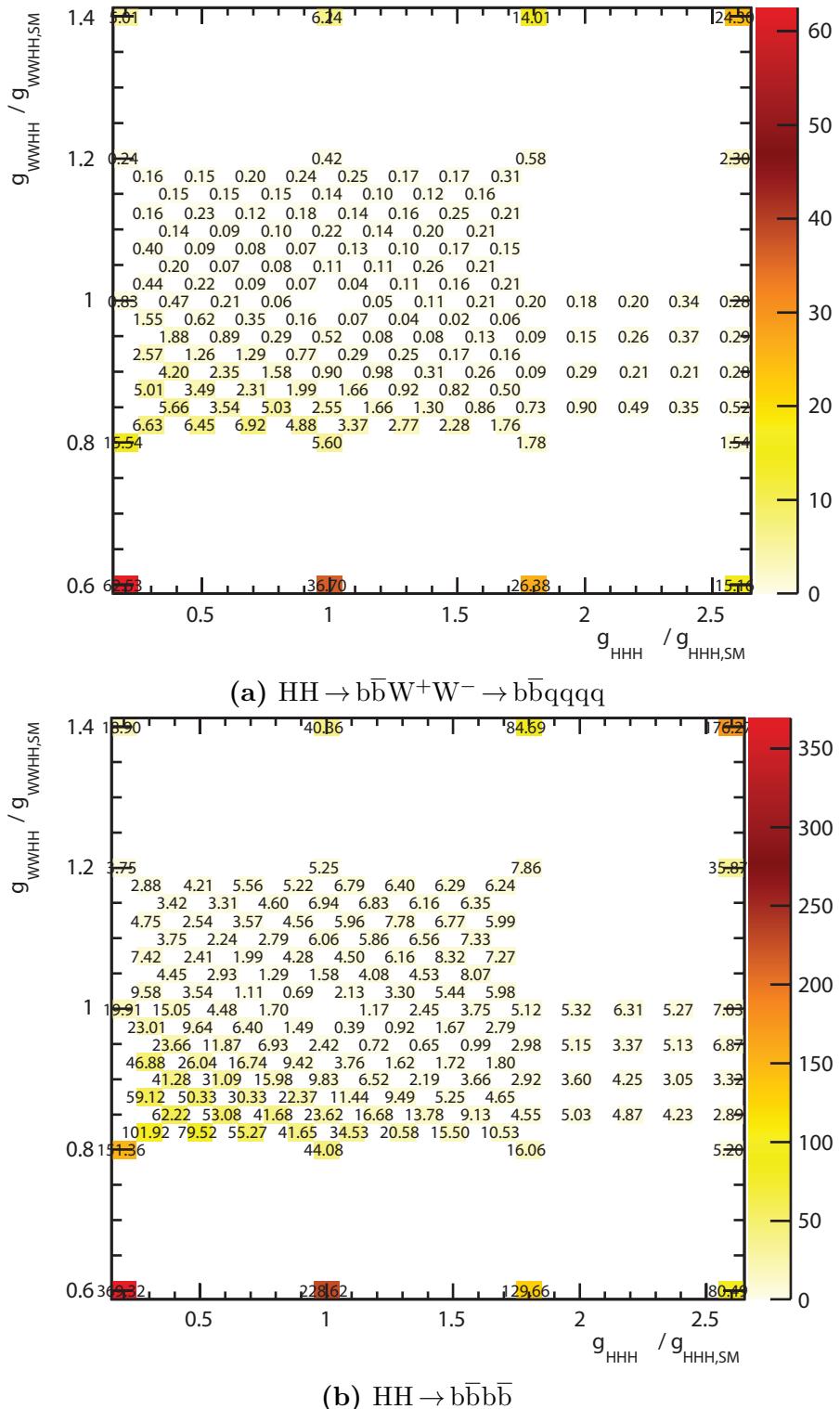
Two sub-channels, hadronic  $W^+W^-$  decay of  $HH \rightarrow b\bar{b}W^+W^-$  and  $HH \rightarrow b\bar{b}b\bar{b}$ , are combined to increase the statistical precision on the coupling measurements. To minimise



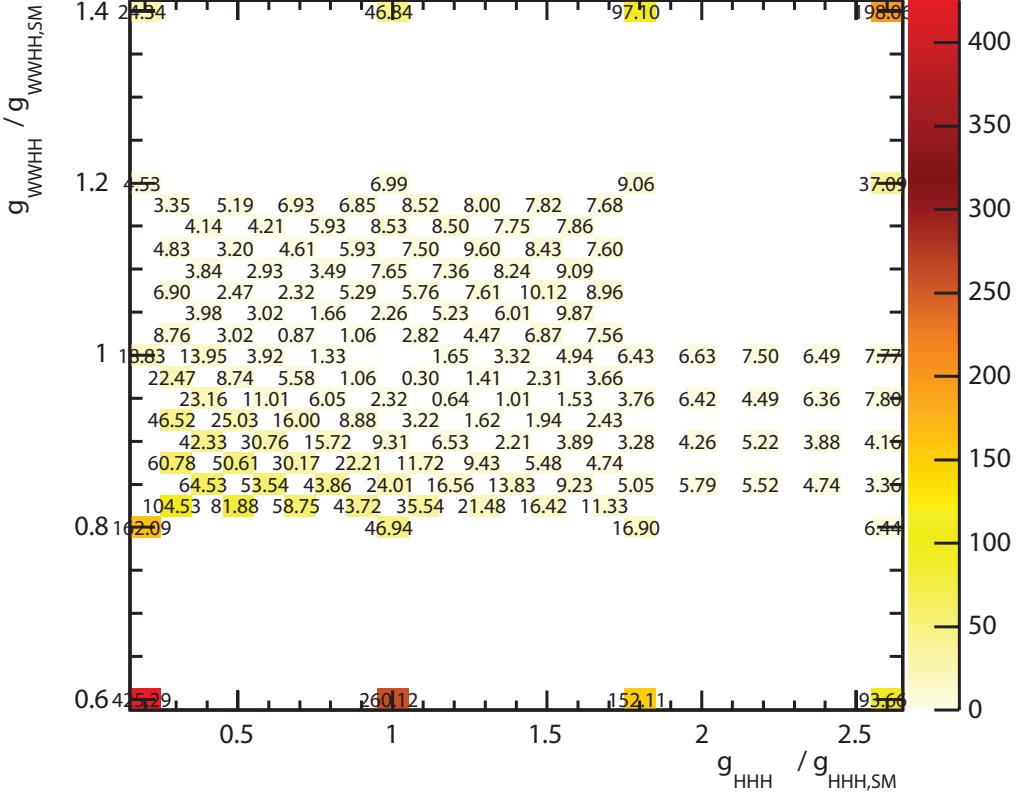
**Figure 8.21:** The significance for the  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  process as a function of the  $g_{\text{HHH}}/g_{\text{HHH,SM}}$  and  $g_{\text{WWHH}}/g_{\text{WWHH,SM}}$  at  $\sqrt{s} = 3 \text{ TeV}$ , using hadronic  $W^+W^-$  decay of  $HH \rightarrow b\bar{b}W^+W^-$  sub-channel, assuming an integrated luminosity of  $3000 \text{ fb}^{-1}$ .

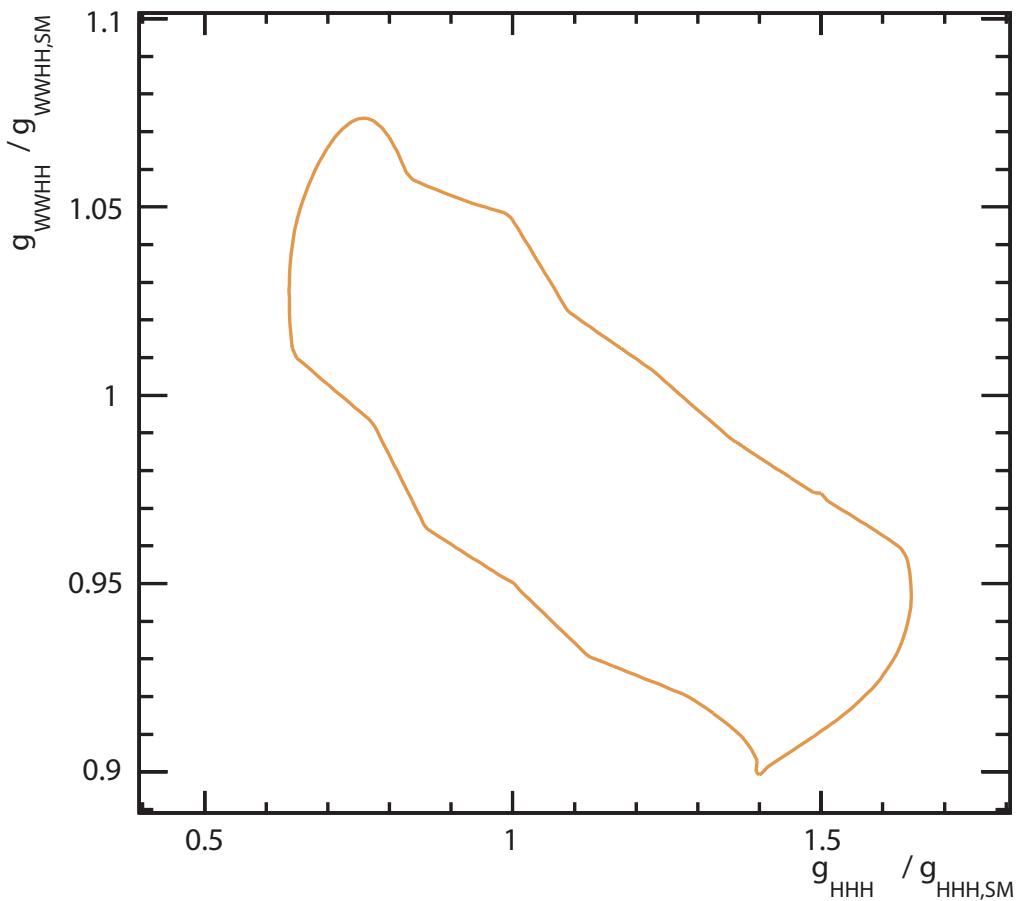
the statistical fluctuations when generating samples with different non-SM couplings, a toy MC experiment is performed. The SM coupling sample is treated as a data template set. 100000 data sets are generated by fluctuating the event number in each kinematic bin in the data template according to the Poisson distribution with a mean that is equal to the event number in the bin. The  $\chi^2$  is calculated using these generated data sets as the observed data ( $N_{i,\text{observed}}$  in equation 8.17). The  $\chi^2$  is then averaged over the number of data sets (100000) and normalised such that the  $\chi^2$  at the SM couplings is 0. Figure 8.23 shows the normalised  $\chi^2$  after averaging over 100000 toy MC experiments as a function of  $g_{\text{HHH}}/g_{\text{HHH,SM}}$  and  $g_{\text{WWHH}}/g_{\text{WWHH,SM}}$ .

Since there are two couplings in this  $\chi^2$  surface, the degree of freedom is 2. A contour of 68% confidence ( $\chi^2 = 2.3$ ) can be drawn by interpolating between points on the  $\chi^2$  surface in figure 8.23. Figure 8.24 shows the contour of 68% confidence of the measurements of the  $g_{\text{HHH}}$  and  $g_{\text{WWHH}}$ . The counter can be sliced one dimensionally to extract the uncertainty of the measurements of one coupling for a given value of the



**Figure 8.22:** The  $\chi^2$  for the  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e$  process as a function of the  $g_{\text{HHH}}/g_{\text{HHH,SM}}$  and  $g_{\text{WWHH}}/g_{\text{WWHH,SM}}$  at  $\sqrt{s} = 3 \text{ TeV}$ , using: a) hadronic  $W^+W^-$  decay of  $\text{HH} \rightarrow b\bar{b}W^+W^-$ ; b) and  $\text{HH} \rightarrow b\bar{b}b\bar{b}$  sub-channels, assuming an integrated luminosity of  $3000 \text{ fb}^{-1}$ .





**Figure 8.24:** Contour plot of 68% confidence ( $\chi^2 = 2.3$ ) , after averaging the toy MC experiments, as a function of  $g_{\text{HHH}} / g_{\text{HHH,SM}}$  and  $g_{\text{WWHH}} / g_{\text{WWHH,SM}}$ , combining hadronic  $W^+W^-$  decay of  $\text{HH} \rightarrow \text{bb}W^+W^-$  and  $\text{HH} \rightarrow \text{bbbb}$  sub-channels, assuming an integrated luminosity of  $3000 \text{ fb}^{-1}$ .



# Chapter 9

## Summary

*‘To know what you know and what you do not know, that is true knowledge.’*

— Confucius, 551 BC – 479 BC

This chapter summarises key results presented in analyses in this thesis.

In chapter 5, a set of photon reconstruction algorithms developed in PandoraPFA are discussed. Using the ILD detector model, the single photon reconstruction efficiency is above 98% for photons with energies above 2GeV, and above 99.5% for photons with energies above 100 GeV. The photon fragments produced during the event reconstruction have been greatly reduced. The ability to separate spatially close photons and the jet energy resolution have improved, as a result of a better photon reconstruction. Using a two-photon-per-event sample with photon energies of 500 GeV and 50 GeV, the average numbers of photons and particles beyond a distance separation of 20 mm are both less than 2.05, where the true value is 2. The distance of the minimal resolved photon pairs is reduced to 6 mm for two photons with same energies, and 10 mm for two photons with different energies. The jet energy resolution has been improved for  $\sqrt{s} = 360$  GeV and 500 GeV. The photon confusion terms, except at  $\sqrt{s} = 91$  GeV, have been reduced to 0.9%.

In chapter 6, a high classification rate of the tau lepton seven major decay modes is achieved using the  $e^+e^- \rightarrow \tau^+\tau^-$  events in the ILD detector. The tau decay mode classification is used for the ECAL optimisation study with different ECAL cell sizes and different centre-of-mass energies. The efficiency of the tau decay mode classification

decreases with an increase of the centre-of-mass energy and with the increasing ECAL cell sizes. The sensitivity of  $\varepsilon_{had}$  to different cell sizes is stronger at high centre-of-mass energies. With decay products being spatially close at high centre-of-mass energies, it is more beneficial to have a small ECAL cell size to reconstruct individual particles.

Chapter 7 presents a proof-of-principle demonstration that the generator-level pion energy fraction correlation can be reconstructed at the analysis level. The analysis contains several important steps: tau identification; kinematic reconstruction of the energies of the taus; the classification of the  $\tau^\pm \rightarrow \pi^\pm \nu_\tau$  decay mode; and the reconstruction of the tau pair polarisation correlations discussed in section 2.9.

In chapter 8, the analyses of the  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e$  channel for CLIC at  $\sqrt{s} = 1.4$  TeV and  $\sqrt{s} = 3$  TeV are performed. The significance of the signal events are 0.56 and 1.09, assuming an integrated luminosity of  $1500\text{ fb}^{-1}$  and  $2000\text{ fb}^{-1}$ , for  $\sqrt{s} = 1.4$  TeV and  $\sqrt{s} = 3$  TeV respectively. The uncertainty on measurement of the Higgs trilinear self coupling,  $g_{HHH}$ , from  $e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e$  analysis is obtained:

$$\frac{\Delta g_{HHH}}{g_{HHH}} = \begin{cases} 218\%, & \text{at } \sqrt{s} = 1.4 \text{ TeV}, \\ 135\%, & \text{at } \sqrt{s} = 3 \text{ TeV}. \end{cases} \quad (9.1)$$

When analysis at both  $HH \rightarrow b\bar{b}W^+W^-$  and  $HH \rightarrow b\bar{b}b\bar{b}$  sub-channels are combined at  $\sqrt{s} = 3$  TeV, assuming an integrated luminosity of  $3000\text{ fb}^{-1}$ , the simultaneous extraction of the uncertainty on the measurement of the  $g_{HHH}$  and  $g_{WWHH}$  yields:

$$\frac{\Delta g_{WWHH}}{g_{WWHH}} \simeq 4.9\%, \text{ for } g_{HHH} = g_{HHH}^{\text{SM}}, \quad (9.2)$$

$$\frac{\Delta g_{HHH}}{g_{HHH}} \simeq 29\%, \text{ for } g_{WWHH} = g_{WWHH}^{\text{SM}}. \quad (9.3)$$

The statistical precisions on the measurements of  $g_{WWHH}$  and  $g_{HHH}$  are much better at CLIC than at current LHC or high luminosity upgraded LHC [20].

# Appendix A

## Generation Parameters

*‘Opportunities multiply as they are seized.’*

— Du Mu, 803 – 852

Particle masses and widths used to generate SM samples for studies with CLIC detectors, used in chapter 8, are listed in table A.1. The Higgs boson mass is specified for individual samples.

Particle	Mass (GeV/c <sup>2</sup> )	Width (GeV/c <sup>2</sup> )
u, d, s quarks	0	0
c quark	0.54	0
b quark	2.9	0
t quark	174	1.37
W boson	80.45	2.071
Z boson	91.188	2.478

**Table A.1:** Particle masses and widths used for the generation of SM samples for CLIC detectors, taken from [2].



## Appendix B

# Photon Reconstruction in PandoraPFA

*'Thank you Mario! But our Princess is in another castle!'*

— Toad, Super Mario Bros [109], 1985

Here are extra tables for the chapter 5.

$E_f \leq 1 \text{ GeV}$	Photon–photon	Photon–neutral-hadron
Transverse shower comparison, or	$d < 30 \text{ mm}; \frac{E_{p1}}{E_m+E_f} > 0.9;$ $\frac{E_{p2}}{E_f} < 0.5; E_{p1} > E_m$	-
Low energy fragment, or	$d < 20 \text{ mm}; E_f < 0.4 \text{ GeV}$	$d < 20 \text{ mm}; d_c < 40 \text{ mm}$
Small fragment 1, or	$d < 30 \text{ mm}; N_{\text{calo}} < 40;$ $d_c < 50 \text{ mm}$	$d < 50 \text{ mm}; N_{\text{calo}} < 10;$ $d_h < 50 \text{ mm}$
Small fragment 2, or	$d < 50 \text{ mm}; N_{\text{calo}} < 20$	-
Small fragment forward region, or	$N_{\text{calo}} < 40; d_c < 60 \text{ mm};$ $E_f < 0.6 \text{ GeV};$ $ \cos(\theta_Z)  > 0.7$	-
Relative low energy fragment	$d < 40 \text{ mm}; d_h < 20 \text{ mm};$ $\frac{E_f}{E_m} < 0.01$	$d < 40 \text{ mm}; d_h < 15 \text{ mm};$ $\frac{E_f}{E_m} < 0.01$
$E_f > 1 \text{ GeV}$	Photon–photon	Photon–neutral-hadron
Transverse shower comparison, or	$\frac{E_{p1}}{E_m+E_f} > 0.9; E_{p2} = 0 \text{ or}$ $(\frac{E_{p2}}{E_f} < 0.5, E_{p1} > E_m)$	$\frac{E_{p1}}{E_m+E_f} > 0.9; E_{p2} = 0 \text{ or}$ $(\frac{E_{p2}}{E_f} < 0.5, E_{p1} > E_m)$
Relative low energy fragment 1, or	$d < 40 \text{ mm}; d_h < 20 \text{ mm};$ $\frac{E_f}{E_m} < 0.02$	$d < 40 \text{ mm}; d_h < 20 \text{ mm};$ $\frac{E_f}{E_m} < 0.02$
Relative low energy fragment 2, or	-	$d < 40 \text{ mm}; d_h < 20 \text{ mm};$ $\frac{E_f}{E_m} < 0.1; E_f > 10 \text{ GeV}$
Relative low energy fragment 3	-	$d < 20 \text{ mm}; d_h < 20 \text{ mm};$ $\frac{E_f}{E_m} < 0.2; E_f > 10 \text{ GeV}$

**Table B.1:** The cuts for merging photon–photon-fragment pairs and photon–neutral-hadron-fragment pairs for both low energy and high energy fragments, after charged hadron reconstruction. Variables  $d$ ,  $d_c$  and  $d_h$  are the mean energy weighted intra-layer distance of the pair, the distance between centroids, the minimum distance between calorimeter hits of the pair, respectively. Variables  $E_m$  and  $E_f$  are the main photon energy and the fragment energy, respectively. Variables  $E_{p1}$  and  $E_{p2}$  are the energies the two largest peaks, found by 2D PEAK FINDING algorithm, ordered by descending energy, respectively.  $N_{\text{calo}}$  is the number of the calorimeter hits in the fragment.  $|\cos(\theta_Z)|$  is the absolute cosine of the polar angle, where beam direction is the z-axis.

$E_f \leq 1 \text{ GeV}$	Photon–photon	Photon–neutral-hadron
Transverse shower comparison, or	$d < 20 \text{ mm}; \frac{E_{p1}}{E_m+E_f} > 0.9;$ $E_{p2} = 0 \text{ or } (\frac{E_{p2}}{E_f} < 0.5,$ $E_{p1} > E_m)$	$d < 20 \text{ mm}; \frac{E_{p1}}{E_m+E_f} > 0.9;$ $E_{p2} = 0 \text{ or } (\frac{E_{p2}}{E_f} < 0.5,$ $E_{p1} > E_m)$
Low energy fragment, or	$d < 20 \text{ mm}; E_f < 0.2 \text{ GeV}$	-
Small fragment 1, or	$d < 30 \text{ mm}; N_{\text{calo}} < 20;$ $d_h < 13 \text{ mm}$	$d < 50 \text{ mm}; N_{\text{calo}} < 10;$ $d_h < 50 \text{ mm}$
Small fragment 2, or	$d_c < 30 \text{ mm}; N_{\text{calo}} < 10;$ $d_h < 13 \text{ mm}$	-
Relative low energy fragment	-	$d < 40 \text{ mm}; d_h < 15 \text{ mm};$ $\frac{E_f}{E_m} < 0.01$
$E_f > 1 \text{ GeV}$	Photon–photon	Photon–neutral-hadron
Transverse shower comparison, or	$d < 20 \text{ mm}; \frac{E_{p1}}{E_m+E_f} > 0.9;$ $E_{p2} = 0 \text{ or } (\frac{E_{p2}}{E_f} < 0.5,$ $E_{p1} > E_m)$	$d < 20 \text{ mm}; \frac{E_{p1}}{E_m+E_f} > 0.9;$ $E_{p2} = 0 \text{ or } (\frac{E_{p2}}{E_f} < 0.5,$ $E_{p1} > E_m)$
Relative low energy fragment	-	$d < 40 \text{ mm}; d_h < 20 \text{ mm};$ $\frac{E_f}{E_m} < 0.02$

**Table B.2:** The cuts for merging photon–photon-fragment pairs and photon–neutral-hadron-fragment pairs for both low energy and high energy fragments, immediately after photon reconstruction. Variables  $d$ ,  $d_c$  and  $d_h$  are the mean energy weighted intra-layer distance of the pair, the distance between centroids, the minimum distance between calorimeter hits of the pair, respectively. Variables  $E_m$  and  $E_f$  are the main photon energy and the fragment energy, respectively. Variables  $E_{p1}$  and  $E_{p2}$  are the energies the two largest peaks, found by 2D PEAK FINDING algorithm, ordered by descending energy, respectively.  $N_{\text{calo}}$  is the number of the calorimeter hits in the fragment.



## Appendix C

# Double Higgs Boson Production Analysis

*'I was an adventurer like you, then I took an arrow in the knee.'*

— The town guard, Skyrim [110], 2011

Here are extra tables and plots for chapter 8.

### C.1 Hadronic decay at $\sqrt{s} = 3 \text{ TeV}$ analysis

$\sqrt{s} = 3 \text{ TeV}$	N	Lepton evto	$b\bar{b}W^+W^- / b\bar{b}W^+W^-$ separation	Valid jet Pairing
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e$ , hadronic	146.0	80.9%	72.8%	72.1%
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}b\bar{b}\nu_e\bar{\nu}_e$	355.0	83.5%	20.5%	20.5%
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow \text{other}$	675.0	40.1%	34.3%	20.5%
$e^+e^- \rightarrow q\bar{q}H\nu\bar{\nu}$	6120	67.7%	61.9%	61.9%
$e^+e^- \rightarrow c\bar{c}H\nu\bar{\nu}$	2300	69.1%	53.0%	48.8%
$e^+e^- \rightarrow b\bar{b}H\nu\bar{\nu}$	3560	70.1%	30.9%	30.6%
$e^+e^- \rightarrow qqqq$	1093000	62.4%	44.9%	34.9%
$e^+e^- \rightarrow qqqq\ell\ell$	338600	21.4%	19.6%	13.3%
$e^+e^- \rightarrow qqqq\ell\nu$	213200	23.3%	19.5%	16.3%
$e^+e^- \rightarrow qqqq\nu\bar{\nu}$	143000	80.7%	71.4%	50.7%
$e^+e^- \rightarrow qq$	5897800	72.9%	63.9%	55.4%
$e^+e^- \rightarrow qq\ell\nu$	11121800	34.0%	24.7%	20.5%
$e^+e^- \rightarrow qq\ell\ell$	6639200	43.1%	41.7%	37.0%
$e^+e^- \rightarrow qq\nu\nu$	2635000	84.6%	63.8%	53.2%
$e^\pm\gamma(\text{BS}) \rightarrow e^\pm qqqq$	4007354	31.0%	28.2%	21.1%
$e^\pm\gamma(\text{EPA}) \rightarrow e^\pm qqqq$	1151200	15.9%	14.5%	10.9%
$e^\pm\gamma(\text{BS}) \rightarrow \nu qqqq$	829184	78.3%	68.8%	53.3%
$e^\pm\gamma(\text{EPA}) \rightarrow \nu qqqq$	216800	39.6%	35.0%	26.9%
$e^\pm\gamma(\text{BS}) \rightarrow qqH\nu$	185018.0	64.0%	55.4%	49.8%
$e^\pm\gamma(\text{EPA}) \rightarrow qqH\nu$	46800	32.9%	28.8%	25.9%
$\gamma(\text{BS})\gamma(\text{BS}) \rightarrow qqqq$	18009414	71.6%	65.5%	49.4%
$\gamma(\text{BS})\gamma(\text{EPA}) \rightarrow qqqq$	3824548	44.3%	40.6%	30.6%
$\gamma(\text{EPA})\gamma(\text{BS}) \rightarrow qqqq$	3828498	44.3%	40.7%	30.7%
$\gamma(\text{EPA})\gamma(\text{EPA}) \rightarrow qqqq$	805400	29.0%	26.7%	20.1%

**Table C.1:** The table shows the expected number of events, before cuts and after successive cuts: the lepton veto,  $HH \rightarrow b\bar{b}W^+W^- / HH \rightarrow b\bar{b}b\bar{b}$  separation, and valid jet pairing, for the signal and background events at  $\sqrt{s} = 3 \text{ TeV}$ , assuming an integrated luminosity of  $2000 \text{ fb}^{-1}$ .  $q$  can be  $u, d, s, b$  or  $t$ . Unless specified,  $q, \ell$  and  $\nu$  represent either particles or the corresponding anti-particles.

Process	$m_{\text{HH}} > 150 \text{ GeV}$	$B_1 > 0.7$
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}W^+W^-\nu_e\bar{\nu}_e$ , hadronic	71.7%	61.8%
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow b\bar{b}b\bar{b}\nu_e\bar{\nu}_e$	20.2%	18.8%
$e^+e^- \rightarrow HH\nu_e\bar{\nu}_e \rightarrow \text{other}$	30.2%	20.0%
$e^+e^- \rightarrow q\bar{q}H\nu\bar{\nu}$	53.1%	36.0%
$e^+e^- \rightarrow c\bar{c}H\nu\bar{\nu}$	43.8%	26.3%
$e^+e^- \rightarrow b\bar{b}H\nu\bar{\nu}$	29.6%	25.9%
$e^+e^- \rightarrow qqqq$	26.5%	1.7%
$e^+e^- \rightarrow qqqq\ell\ell$	12.8%	0.7%
$e^+e^- \rightarrow qqqq\ell\nu$	16.0%	7.9%
$e^+e^- \rightarrow qqqq\nu\bar{\nu}$	49.7%	9.0%
$e^+e^- \rightarrow qq$	8.3%	1.4%
$e^+e^- \rightarrow qq\ell\nu$	6.0%	0.1%
$e^+e^- \rightarrow qq\ell\ell$	1.9%	0.4%
$e^+e^- \rightarrow qq\nu\nu$	16.6%	3.1%
$e^\pm\gamma(\text{BS}) \rightarrow e^\pm qqqq$	19.4%	0.7%
$e^\pm\gamma(\text{EPA}) \rightarrow e^\pm qqqq$	9.9%	0.4%
$e^\pm\gamma(\text{BS}) \rightarrow \nu qqqq$	51.3%	16.4%
$e^\pm\gamma(\text{EPA}) \rightarrow \nu qqqq$	26.0%	7.7%
$e^\pm\gamma(\text{BS}) \rightarrow qqH\nu$	47.9%	30.3%
$e^-\gamma(\text{EPA}) \rightarrow qqH\nu$	25.0%	15.8%
$\gamma(\text{BS})\gamma(\text{BS}) \rightarrow qqqq$	44.5%	1.7%
$\gamma(\text{BS})\gamma(\text{EPA}) \rightarrow qqqq$	27.4%	1.0%
$\gamma(\text{EPA})\gamma(\text{BS}) \rightarrow qqqq$	27.5%	1.0%
$\gamma(\text{EPA})\gamma(\text{EPA}) \rightarrow qqqq$	18.0%	0.7%

**Table C.2:** The table shows the expected number of events after successive cuts: invariant mass of the two Higgs system  $> 150 \text{ GeV}$ , and the highest b-jet tag value  $> 0.7$ . All cuts include the lepton veto,  $HH \rightarrow b\bar{b}W^+W^- / HH \rightarrow b\bar{b}b\bar{b}$  separation, and valid jet pairing. Table shows the signal and background events at  $\sqrt{s} = 3 \text{ TeV}$ , assuming an integrated luminosity of  $2000 \text{ fb}^{-1}$ .  $q$  can be  $u, d, s, b$  or  $t$ . Unless specified,  $q, \ell$  and  $\nu$  represent either particles or the corresponding anti-particles.



# Colophon

This thesis was made in L<sup>A</sup>T<sub>E</sub>X 2 <sub>$\varepsilon$</sub>  using the “heptesis” class [111].



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