Thermodynamics & Statistical Physics Exercises

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- For a given distribution $(\varepsilon_l, \omega_l, a_l)$,
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 - $\Omega_{\text{B.S.}} = \prod \frac{(\omega_l + a_l 1)!}{a_l!(\omega_l 1)!} = \prod \frac{a_l!}{a_l!} = 1;$
- for Fermion gas, the number of micro-states (6.5.7): $\Omega_{\text{F.D.}} = \prod \frac{\omega_l!}{a_l!(\omega_l a_l)!} = \prod \frac{1}{a_l!(1 a_l)!} = 1;$
- There is no maximum for the number of micro-states.

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 According to the Maxwell's speed distribution, the number of molecules at (v, v + dv): $dN(v) = 4\pi N(\frac{m}{2kT})^{\frac{3}{2}} e^{-\frac{\dot{m}}{2kT}v^2} v^2 dv.$ The most probable speed: $v_m = \sqrt{\frac{2kT}{m}}$. $\frac{N_m}{N} = \int_0^{v_m} 4\pi (\frac{m}{2\pi kT})^{\frac{3}{2}} e^{-\frac{m}{2kT}v^2} v^2 dv$ $= \int_0^{(\sqrt{\frac{m}{2kT}})\sqrt{\frac{2kT}{m}}} \frac{4}{\sqrt{\pi}} e^{-\frac{m}{2kT}v^2} (\sqrt{\frac{m}{2kT}}v)^2 d(\sqrt{\frac{m}{2kT}}v)$ $=\frac{4}{\sqrt{\pi}}\int_0^1 x^2 e^{-x^2} dx$

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• According to the Maxwell's speed distribution, the number of molecules at $(v,v+\mathrm{d}v)$: $\mathrm{d}N(v)=4\pi N(\tfrac{m}{2kT})^{\frac{3}{2}}e^{-\frac{m}{2kT}v^2}v^2\mathrm{d}v.$

The most probable speed:
$$v_m = \sqrt{\frac{2kT}{m}}$$
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$$\frac{N_m}{N} = \int_0^{v_m} 4\pi \left(\frac{m}{2\pi kT}\right)^{\frac{3}{2}} e^{-\frac{m}{2kT}v^2} v^2 dv$$

$$= \int_0^{(\sqrt{\frac{m}{2kT}})\sqrt{\frac{2kT}{m}}} \frac{4}{\sqrt{\pi}} e^{-\frac{m}{2kT}v^2} \left(\sqrt{\frac{m}{2kT}}v\right)^2 d\left(\sqrt{\frac{m}{2kT}}v\right)$$

$$= \frac{4}{\sqrt{\pi}} \int_0^1 x^2 e^{-x^2} dx$$

$$\approx 0.4276:$$

Not depending on T.

A8.21 Prove: for monatomic classical ideal gas, entropy:

$$S = \frac{5}{2}Nk - Nk\ln(n\lambda^3).$$

$$S = Nk(\ln Z_1 - \beta \frac{\partial}{\partial \beta} \ln Z_1) - k \ln N!$$
 (7.1.13')

$$S = Nk(\ln Z_1 - \beta \frac{\partial}{\partial \beta} \ln Z_1) - k \ln N! \quad (7.1.13')$$

$$\simeq Nk\{\ln[V(\frac{2\pi m}{h^2 \beta})^{\frac{3}{2}}] - \beta \frac{\partial}{\partial \beta} \ln[V(\frac{2\pi m}{h^2 \beta})^{\frac{3}{2}}]\} - kN(\ln N - 1)$$

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$$= Nk\{\ln[V(\frac{2\pi mkT}{h^2})^{3/2}] - \beta \frac{\partial}{\partial \beta} \ln \beta^{-3/2}] - (\ln N - 1)\}$$

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$$= Nk\{-\ln(n\lambda^3) + \frac{5}{2}\} \quad \text{(before (7.2.7))}$$

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$$= \frac{5}{2}Nk - Nk\ln(n\lambda^3);$$

A8.21 Prove: for monatomic classical ideal gas, entropy:

$$S = \frac{5}{2}Nk - Nk\ln(n\lambda^3).$$

• Partition function: $Z_1 = V(\frac{2\pi m}{h^2 \beta})^{3/2}$. (7.5.21) Entropy:

$$S = Nk(\ln Z_1 - \beta \frac{\partial}{\partial \beta} \ln Z_1) - k \ln N! \quad (7.1.13')$$

$$\simeq Nk\{\ln[V(\frac{2\pi m}{h^2\beta})^{\frac{3}{2}}] - \beta \frac{\partial}{\partial \beta} \ln[V(\frac{2\pi m}{h^2\beta})^{\frac{3}{2}}]\} - kN(\ln N - 1)$$

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$$= Nk\{-\ln(n\lambda^3) + \frac{5}{2}\} \quad \text{(before (7.2.7))}$$

$$= \frac{5}{2}Nk - Nk\ln(n\lambda^3);$$

• $S = \frac{3}{2}Nk\ln T + Nk\ln\frac{V}{N} + \frac{3}{2}Nk\left[\frac{5}{3} + \ln\frac{2\pi mk}{h^2}\right]$. (7.6.2)

A8.26 Monatomic classical ideal gas, calculate the probability for atoms at $(\varepsilon, \varepsilon + d\varepsilon)$, $\rho(\varepsilon)$.

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• The Maxwell's speed distribution, the number of molecules at (v, v + dv):

$$\rho(v) dv = \frac{dN(v)}{N} = 4\pi (\frac{m}{2\pi kT})^{\frac{3}{2}} e^{-\frac{m}{2kT}v^2} v^2 dv.$$

• The Maxwell's speed distribution, the number of molecules at $(v,v+\mathrm{d}v)$: $\rho(v)\mathrm{d}v = \frac{\mathrm{d}N(v)}{N} = 4\pi(\frac{m}{2\pi kT})^{\frac{3}{2}}e^{-\frac{m}{2kT}v^2}v^2\mathrm{d}v.$ Convert to energy $(v=\sqrt{\frac{2\varepsilon}{m}})$:

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$$\rho(v)dv = \frac{dN(v)}{N} = 4\pi \left(\frac{m}{2\pi kT}\right)^{\frac{3}{2}} e^{-\frac{m}{2kT}v^2} v^2 dv.$$

Convert to energy
$$(v = \sqrt{\frac{2\varepsilon}{m}})$$
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$$\rho(\varepsilon)d\varepsilon = 4\pi \left(\frac{m}{2\pi kT}\right)^{\frac{3}{2}} e^{-\frac{m}{2kT}\frac{2\varepsilon}{m}} \frac{2\varepsilon}{m} d\sqrt{\frac{2\varepsilon}{m}}$$

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• The Maxwell's speed distribution, the number of molecules at $(v, v + \mathrm{d}v)$:

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Convert to energy $(v = \sqrt{\frac{2\varepsilon}{m}})$:

$$\rho(\varepsilon)d\varepsilon = 4\pi \left(\frac{m}{2\pi kT}\right)^{\frac{3}{2}} e^{-\frac{m}{2kT}\frac{2\varepsilon}{m}} \frac{2\varepsilon}{m} d\sqrt{\frac{2\varepsilon}{m}}$$
$$= \frac{2}{\sqrt{\pi}} (kT)^{-\frac{3}{2}} \varepsilon^{\frac{1}{2}} e^{-\frac{\varepsilon}{2kT}} d\varepsilon;$$

• Compare with the black-body:

$$\rho(\omega)d\omega = \frac{V}{\pi^2 c^3} \frac{\omega^2}{e^{\frac{\hbar \omega}{kT}} - 1} d\omega.$$
 (8.4.6)

•
$$n\lambda^3 = n(\frac{h^2}{2\pi mkT})^{\frac{3}{2}}$$
. $n = ?$

• $n\lambda^3=n(\frac{h^2}{2\pi mkT})^{\frac{3}{2}}.$ n=? Solar photosphere density $2\times 10^{-4}{\rm kg~m^{-3}}$ (wikipedia);

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• $n\lambda^3 = n(\frac{h^2}{2\pi mkT})^{\frac{3}{2}}$. n=? Solar photosphere density $2\times 10^{-4}{\rm kg~m^{-3}}$ (wikipedia); number density $n\simeq \frac{2\times 10^{-4}{\rm kg~m^{-3}}}{1.67\times 10^{-27}{\rm kg}}\simeq 10^{23}{\rm m^{-3}}$. $n\lambda^3\simeq 10^{23}{\rm m^{-3}}[\frac{(6.626\times 10^{-34}{\rm J\cdot s})^2}{2\pi\cdot 1.67\times 10^{-27}{\rm kg}\cdot k\cdot \frac{1{\rm eV}}{k}}]^{3/2}$

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 $\begin{array}{l} \bullet \ n\lambda^3 = n(\frac{h^2}{2\pi mkT})^{\frac{3}{2}}. \ n = ? \\ \text{Solar photosphere density } 2\times 10^{-4} \mathrm{kg \, m^{-3}} \ \text{(wikipedia);} \\ \text{number density } n \simeq \frac{2\times 10^{-4} \mathrm{kg \, m^{-3}}}{1.67\times 10^{-27} \mathrm{kg}} \simeq 10^{23} \mathrm{m^{-3}}. \\ n\lambda^3 \simeq 10^{23} \mathrm{m^{-3}} [\frac{(6.626\times 10^{-34}\mathrm{J \cdot s})^2}{2\pi \cdot 1.67\times 10^{-27} \mathrm{kg} \cdot k \cdot \frac{\mathrm{1eV}}{k}}]^{3/2} \\ \simeq 10^{23} \mathrm{m^{-3}} [\frac{(6.626\times 10^{-34}\mathrm{J \cdot s})^2}{2\pi \cdot 1.67\times 10^{-27} \mathrm{kg \cdot 1.6}\times 10^{-19}\mathrm{J}}]^{3/2} \simeq 4.2\times 10^{-10} \\ \end{array}$

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• 1. $\frac{3}{2}kT = \bar{\varepsilon} = 1eV$;

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• 1. $\frac{3}{2}kT = \bar{\varepsilon} = 1\text{eV}; \Rightarrow T \simeq 7.7 \times 10^3\text{K}.$

• 2. Energy level: $\varepsilon_n = -\frac{13.6}{n^2} \text{eV}$;

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• 2. Energy level: $\varepsilon_n = -\frac{13.6}{n^2} \text{eV}$; Boltzmann distribution: $a_l = \omega_l e^{-\alpha - \beta \varepsilon_l}$. so. $\frac{N_3}{N_1} = \frac{a_3}{a_1} = 3^2 \cdot e^{\frac{\varepsilon_1 - \varepsilon_3}{kT}}$

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- $N_{e^{\pm}} = \int_0^\infty \frac{1}{e^{\frac{\varepsilon}{kT}} + 1} 2 \cdot \frac{V4\pi p^2 \mathrm{d}p}{h^3}$.
- $kT \gg m_e c^2$, relativistic, $\varepsilon \simeq pc$, $N_{e^{\pm}} = \int_0^{\infty} \frac{1}{e^{\frac{1}{kT}} + 1} 2 \cdot \frac{V 4\pi p^2 dp}{h^3} = \frac{V (kT)^3}{\pi^2 (\hbar c)^3} \int_0^{\infty} \frac{x^2 dx}{e^x + 1}$.
- $kT \ll m_e c^2$, $\varepsilon = \sqrt{(pc)^2 + (m_e c^2)^2} \simeq m_e c^2 + \frac{p^2}{2m_e}$ and $e^{\frac{\varepsilon}{kT}} \gg 1$, $N_{e^{\pm}} = \int_0^{\infty} \frac{1}{e^{\frac{m_e c^2 + \frac{p^2}{2m_e}}{kT}}} 2 \cdot \frac{V4\pi p^2 \mathrm{d}p}{h^3}$ $= e^{-\frac{m_e c^2}{kT}} \int_0^{\infty} e^{-\frac{p^2}{2m_e kT}} 2 \cdot \frac{V4\pi p^2 \mathrm{d}p}{h^3} = 2V(\frac{2\pi m_e kT}{h^2})^{\frac{3}{2}} e^{-\frac{m_e c^2}{kT}}.$

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- ullet Two cases according to the velocity comparing with c.
- a. Non-relativistic case, $\varepsilon = \frac{p^2}{2m}$, number of state: $D(\varepsilon)\mathrm{d}\varepsilon = g \cdot \frac{V4\pi p^2\mathrm{d}p}{h^3} = C\varepsilon^{1/2}\mathrm{d}\varepsilon$. average energy: $\bar{\varepsilon} = \frac{\int_0^{\varepsilon_F} \varepsilon^{3/2}\mathrm{d}\varepsilon}{\int_0^{\varepsilon_F} \varepsilon^{1/2}\mathrm{d}\varepsilon}$

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A9.23 Fermi energy of Fermi gas is ε_F . Calculate the average energy at T=0K.

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- a. Non-relativistic case, $\varepsilon = \frac{p^2}{2m}$, number of state: $D(\varepsilon)\mathrm{d}\varepsilon = g \cdot \frac{V4\pi p^2\mathrm{d}p}{h^3} = C\varepsilon^{1/2}\mathrm{d}\varepsilon$. average energy: $\bar{\varepsilon} = \frac{\int_0^{\varepsilon_F} \varepsilon^{3/2}\mathrm{d}\varepsilon}{\int_0^{\varepsilon_F} \varepsilon^{1/2}\mathrm{d}\varepsilon} = \frac{3}{5}\varepsilon_F$.
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$$\mathcal{Z}(\alpha, \beta, V) = e^{a[\frac{4}{15}(-\alpha)^{5/2} + \frac{\pi^2}{6}(-\alpha)^{1/2} - \frac{7\pi^4}{1440}(-\alpha)^{-3/2}]\beta^{-3/2}}$$
, where $a = \frac{2\pi g V (2m)^{3/2}}{h^3}$, g is the inner degeneracy. Then calculate the grand potential $J(T, V, \mu)$. (non-relativistic)

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 $\ln \mathcal{Z} = \sum \omega_l \ln(1 + e^{-\alpha - \beta \varepsilon_l})$

Fermi statistics

A9.32 Prove \mathcal{Z} for strong degenerated Fermi gas:

$$\mathcal{Z}(\alpha,\beta,V) = e^{a[rac{4}{15}(-lpha)^{5/2} + rac{\pi^2}{6}(-lpha)^{1/2} - rac{7\pi^4}{1440}(-lpha)^{-3/2}]eta^{-3/2}}$$
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• $D(\varepsilon)d\varepsilon = g \cdot \frac{V4\pi p^2 dp}{h^3} = \frac{2\pi gV(2m)^{3/2}}{h^3} \varepsilon^{1/2} d\varepsilon = a\varepsilon^{1/2} d\varepsilon.$ $\ln \mathcal{Z} = \sum_{l} \omega_l \ln(1 + e^{-\alpha - \beta \varepsilon_l})$ $= \int_0^\infty D(\varepsilon) \ln(1 + e^{-\alpha - \beta \varepsilon}) d\varepsilon$

$$\mathcal{Z}(\alpha,\beta,V)=e^{a[\frac{4}{15}(-\alpha)^{5/2}+\frac{\pi^2}{6}(-\alpha)^{1/2}-\frac{7\pi^4}{1440}(-\alpha)^{-3/2}]\beta^{-3/2}}, \text{ where } a=\frac{2\pi gV(2m)^{3/2}}{h^3}, \ g \text{ is the inner degeneracy. Then calculate the grand potential } J(T,V,\mu). \text{ (non-relativistic)}$$

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$$D(\varepsilon)d\varepsilon = g \cdot \frac{V4\pi p^2 dp}{h^3} = \frac{2\pi g V(2m)^{3/2}}{h^3} \varepsilon^{1/2} d\varepsilon = a\varepsilon^{1/2} d\varepsilon.$$

 $\ln \mathcal{Z} = \sum_{l} \omega_l \ln(1 + e^{-\alpha - \beta \varepsilon_l})$
 $= \int_0^\infty D(\varepsilon) \ln(1 + e^{-\alpha - \beta \varepsilon}) d\varepsilon$
 $= a \int_0^\infty \varepsilon^{1/2} \ln(1 + e^{-\alpha - \beta \varepsilon}) d\varepsilon$

$$\mathcal{Z}(\alpha,\beta,V) = e^{a[\frac{4}{15}(-\alpha)^{5/2}+\frac{\pi^2}{6}(-\alpha)^{1/2}-\frac{7\pi^4}{1440}(-\alpha)^{-3/2}]\beta^{-3/2}}$$
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 $= a \{ \frac{2}{3} \varepsilon^{3/2} \ln(1 + e^{-\alpha - \beta \varepsilon}) |_0^\infty + \frac{2}{3} \int_0^\infty \varepsilon^{3/2} \frac{\beta e^{-\alpha - \beta \varepsilon}}{1 + e^{-\alpha - \beta \varepsilon}} d\varepsilon \}$

$$\mathcal{Z}(\alpha,\beta,V) = e^{a[\frac{4}{15}(-\alpha)^{5/2}+\frac{\pi^2}{6}(-\alpha)^{1/2}-\frac{7\pi^4}{1440}(-\alpha)^{-3/2}]\beta^{-3/2}}$$
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..(A grand dic of phys problems & solutions 5, P_{322} .)

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$$\begin{aligned} \bullet & D(\varepsilon) \mathrm{d}\varepsilon = g \cdot \frac{V4\pi p^2 \mathrm{d}p}{h^3} = \frac{2\pi g V(2m)^{3/2}}{h^3} \varepsilon^{1/2} \mathrm{d}\varepsilon = a\varepsilon^{1/2} \mathrm{d}\varepsilon. \\ & \ln \mathcal{Z} = \sum \omega_l \ln(1 + e^{-\alpha - \beta \varepsilon_l}) \\ & = \int_0^\infty D(\varepsilon) \ln(1 + e^{-\alpha - \beta \varepsilon}) \mathrm{d}\varepsilon \\ & = a \int_0^\infty \varepsilon^{1/2} \ln(1 + e^{-\alpha - \beta \varepsilon}) \mathrm{d}\varepsilon \\ & = a \{\frac{2}{3}\varepsilon^{3/2} \ln(1 + e^{-\alpha - \beta \varepsilon})|_0^\infty + \frac{2}{3} \int_0^\infty \varepsilon^{3/2} \frac{\beta e^{-\alpha - \beta \varepsilon}}{1 + e^{-\alpha - \beta \varepsilon}} \mathrm{d}\varepsilon \} \\ & .. (\text{A grand dic of phys problems \& solutions 5, P}_{322}.) \\ & = a [\frac{4}{15}(-\alpha)^{5/2} + \frac{\pi^2}{6}(-\alpha)^{1/2} - \frac{7\pi^4}{1440}(-\alpha)^{-3/2}]\beta^{-3/2}. \end{aligned}$$

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$$\begin{split} \bullet \ &D(\varepsilon) \mathrm{d}\varepsilon = g \cdot \frac{V4\pi p^2 \mathrm{d}p}{h^3} = \frac{2\pi g V(2m)^{3/2}}{h^3} \varepsilon^{1/2} \mathrm{d}\varepsilon = a\varepsilon^{1/2} \mathrm{d}\varepsilon. \\ &\ln \mathcal{Z} = \sum \omega_l \ln(1 + e^{-\alpha - \beta \varepsilon_l}) \\ &= \int_0^\infty D(\varepsilon) \ln(1 + e^{-\alpha - \beta \varepsilon}) \mathrm{d}\varepsilon \\ &= a \int_0^\infty \varepsilon^{1/2} \ln(1 + e^{-\alpha - \beta \varepsilon}) \mathrm{d}\varepsilon \\ &= a \{\frac{2}{3}\varepsilon^{3/2} \ln(1 + e^{-\alpha - \beta \varepsilon})|_0^\infty + \frac{2}{3} \int_0^\infty \varepsilon^{3/2} \frac{\beta e^{-\alpha - \beta \varepsilon}}{1 + e^{-\alpha - \beta \varepsilon}} \mathrm{d}\varepsilon \} \\ &..(\text{A grand dic of phys problems \& solutions 5, P}_{322}.) \\ &= a [\frac{4}{15}(-\alpha)^{5/2} + \frac{\pi^2}{6}(-\alpha)^{1/2} - \frac{7\pi^4}{1440}(-\alpha)^{-3/2}]\beta^{-3/2}. \\ \bullet \ &J = -kT \ln \mathcal{Z} = -a (\frac{4}{15}\mu^{5/2} + \frac{\pi^2}{6\beta^2}\mu^{1/2} - \frac{7\pi^2}{1440\beta^4}\mu^{-3/2}). \end{split}$$

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• We have $D(\varepsilon) = a\varepsilon^{\frac{1}{2}}$, where $a = \frac{2\pi gV(2m)^{3/2}}{h^3}$; The total number $N = \int_0^{\mu_0} D(\varepsilon) \mathrm{d}\varepsilon = \frac{2}{3}a\mu_0^{\frac{3}{2}}$. $\Rightarrow \mu_0 = \mu(T=0\mathrm{K}) = (\frac{3N}{2a})^{\frac{2}{3}}......(1)$ • $\ln \mathcal{Z} = a[\frac{4}{15}(-\alpha)^{5/2} + \frac{\pi^2}{6}(-\alpha)^{1/2} - \frac{7\pi^4}{1440}(-\alpha)^{-3/2}]\beta^{-3/2}$ $= \frac{2}{5}N\beta\mu_0^{-\frac{3}{2}}\mu^{\frac{5}{2}}[1 + \frac{5}{12}(\frac{\pi}{\beta\mu})^2 - \frac{7}{384}(\frac{\pi}{\beta\mu})^4]......(2)$

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 $N = -(\frac{\partial \ln \mathcal{Z}}{\partial \alpha})_{\beta,V}$

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- $\ln \mathcal{Z} = a \left[\frac{4}{15} (-\alpha)^{5/2} + \frac{\pi^2}{6} (-\alpha)^{1/2} \frac{7\pi^4}{1440} (-\alpha)^{-3/2} \right] \beta^{-3/2}$ $= \frac{2}{5} N \beta \mu_0^{-\frac{3}{2}} \mu^{\frac{5}{2}} [1 + \frac{5}{12} (\frac{\pi}{\beta \mu})^2 - \frac{7}{384} (\frac{\pi}{\beta \mu})^4] \dots (2)$ $N = -(\frac{\partial \ln \mathcal{Z}}{\partial \alpha})_{\beta, V}$ $= a\left[\frac{2}{3}(-\alpha)^{3/2} + \frac{\pi^2}{12}(-\alpha)^{-1/2} - \frac{7\pi^4}{960}(-\alpha)^{-5/2}\right]\beta^{-3/2}$ $= \frac{2}{3}a(\mu^{\frac{3}{2}} + \frac{\pi^2}{8\beta^2}\mu^{-\frac{1}{2}} + \frac{7\pi^2}{640\beta^4}\mu^{-\frac{5}{2}}).....(3)$
- Combine (1) and (3), $\mu = \mu_0 \left[1 + \frac{1}{8} \left(\frac{\pi}{\beta \mu}\right)^2 + \frac{7}{640} \left(\frac{\pi}{\beta \mu}\right)^4\right]^{-\frac{2}{3}}$

$$\ln \mathcal{Z} = \frac{2}{5} N \beta \mu_0^{-\frac{3}{2}} \mu^{\frac{5}{2}} \left[1 + \frac{5}{12} \left(\frac{\pi}{\beta \mu} \right)^2 - \frac{7}{384} \left(\frac{\pi}{\beta \mu} \right)^4 \right] \dots (2)$$

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• $\mu = \mu_0 [1 + \frac{1}{8} (\frac{\pi}{\beta \mu})^2 + \frac{7}{640} (\frac{\pi}{\beta \mu})^4]^{-\frac{2}{3}}$. For the strong degenerated Fermi gas, $kT \ll \mu_0$, $\mu \simeq \mu_0$, so, $\beta \mu \gg 1$; and define $\theta = \frac{\pi}{\beta \mu_0} = \frac{\pi kT}{\mu_0} \ll 1$.

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... For the strong degenerated Fermi gas, calculate μ , p, $U, F, S \text{ and } H. \mid \mu \simeq \mu_0 (1 - \frac{1}{12} \theta^2 - \frac{1}{20} \theta^4).....(4)$

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- Enthalpy: $H = U + pV = \frac{5}{2}pV \simeq N\mu_0(1 + \frac{5}{12}\theta^2 - \frac{1}{16}\theta^4).$

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 $\beta \equiv \frac{1}{p} \left(\frac{\partial p}{\partial T} \right)_V = \frac{5\pi^2 k^2}{6u^2} \left(1 - \frac{5}{12} \theta^2 \right) T.$

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 $\kappa_T \equiv -\frac{1}{V} (\frac{\partial V}{\partial p})_T = \frac{3}{2n\mu_0} (1 - \frac{1}{12}\theta^2).$

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• We have $S=\frac{\pi k}{2}N\theta(1-\frac{1}{10}\theta^2)\simeq \frac{\pi k}{2}N\theta=\frac{\pi^2k^2}{2b}NTV^{\frac{2}{3}},$ where $\theta=\frac{\pi}{\beta\mu_0}=\frac{\pi kT}{\mu_0}$, $\mu_0=(\frac{3N}{2a})^{\frac{2}{3}}$, $a=2\pi gV(\frac{2m}{h^2})^{\frac{3}{2}},$ $b=\frac{h^2}{2m}(\frac{3N}{4\pi a})^{\frac{2}{3}}.$

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Comparing with ideal gas: $pV^{\gamma} = \text{Const.}$ ($\gamma = 5/3$, monatomic molecular) (1.8.4)

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- Using these two, $p \propto T^{5/2}$.

December 30, 2013

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$$\begin{split} \bullet \ D(\varepsilon) \mathrm{d}\varepsilon &= g \cdot \frac{V4\pi p^2 \mathrm{d}p}{h^3} = \frac{4\pi g V}{(hc)^3} \varepsilon^2 \mathrm{d}\varepsilon = a \varepsilon^2 \mathrm{d}\varepsilon, \\ \text{where } a &= \frac{4\pi g V}{(hc)^3}. \\ \ln \mathcal{Z} &= \sum \omega_l \ln(1 + e^{-\alpha - \beta \varepsilon_l}) \\ &= \int_0^\infty D(\varepsilon) \ln(1 + e^{-\alpha - \beta \varepsilon}) \mathrm{d}\varepsilon \\ &= a \int_0^\infty \varepsilon^2 \ln(1 + e^{-\alpha - \beta \varepsilon}) \mathrm{d}\varepsilon \\ &= a \{\frac{1}{3}\varepsilon^3 \ln(1 + e^{-\alpha - \beta \varepsilon})|_0^\infty + \frac{1}{3} \int_0^\infty \varepsilon^3 \frac{\beta e^{-\alpha - \beta \varepsilon}}{1 + e^{-\alpha - \beta \varepsilon}} \mathrm{d}\varepsilon \} \\ ..(\text{A grand dic of phys problems \& solutions 5, P}_{332}.) \end{split}$$

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A9.55 He white dwarf in $T_{\rm WD} = 10^7 {\rm K}$ at center.

 $M_{\rm WD} \sim 1 M_{\odot} \simeq 2 \times 10^{30} {\rm kg}$, $R_{\rm WD} \sim 1 R_{\oplus} \simeq 6 \times 10^6 {\rm m}$.

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Fermion gas in astrophysics

A9.55 He white dwarf in $T_{
m WD}=10^7{
m K}$ at center.

$$M_{
m WD} \sim 1 M_{\odot} \simeq 2 imes 10^{30} {
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, $R_{
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- $kT \simeq 862.5 \text{eV}$, ionized.
- Number density: $n_e = 2 \cdot \frac{M_{
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$$\simeq \frac{2 \cdot 2 \times 10^{30} \text{kg}}{\frac{4}{3}\pi (6 \times 10^6 \text{m})^3 \cdot 4 \times 1.67 \times 10^{-27} \text{kg}}$$

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- i. Non-relativistic: $n\lambda^3 = n(\frac{h^2}{2\pi m kT})^{\frac{3}{2}}$

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- Strong degenerated.

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... He white dwarf.

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 $\simeq 8.5 \times 10^{-14} \, \mathrm{J} \simeq 0.53 \, \mathrm{MeV}$. (This is kinetic energy,

mildly relativistic.)

• Internal energy: $U=\int_0^{\mu_0} \varepsilon D(\varepsilon) \mathrm{d}\varepsilon$

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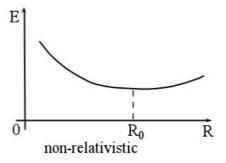
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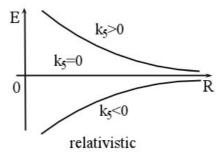
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References

 A: "A grand dictionary of physics problems and solutions 5", Science press, 2005

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