

Thermodynamics & Statistical Physics

Chapter 3. Phase transition of single-component system

Yuan-Chuan Zou
zouyc@hust.edu.cn

School of Physics, Huazhong University of Science and Technology

December 30, 2013

Table of contents

① Phase Transition of Single-Component System

- 3.1 Criterion of thermal equilibrium
- 3.2 Basic equations of open system
- 3.3 Equilibrium of single-component multi-phase system
- 3.4 Properties of equilibrium of s-c multi-phase system
- 3.5 Critical point and phase change between gas and liquid
- 3.7 Classification of the phase transition
- 3.9 Landau's approximation for the continuous phase transition

§3.1 Criterion of thermal equilibrium

§3.1 Criterion of thermal equilibrium

- How to determine the thermal equilibrium **mathematically**?

§3.1 Criterion of thermal equilibrium

- How to determine the thermal equilibrium **mathematically**?
- 2nd law of thermodynamics

§3.1 Criterion of thermal equilibrium

- How to determine the thermal equilibrium **mathematically**?
- 2nd law of thermodynamics
 - Entropy never decreases for isolated system.

§3.1 Criterion of thermal equilibrium

- How to determine the thermal equilibrium **mathematically**?
- 2nd law of thermodynamics
 - Entropy never decreases for isolated system.
 - System is directed in increasing the entropy.

§3.1 Criterion of thermal equilibrium

- How to determine the thermal equilibrium **mathematically**?
- 2nd law of thermodynamics
 - Entropy never decreases for isolated system.
 - System is directed in increasing the entropy.
 - In the thermal equilibrium state, entropy is the maximum.

§3.1 Criterion of thermal equilibrium

- How to determine the thermal equilibrium **mathematically**?
- 2nd law of thermodynamics
 - Entropy never decreases for isolated system.
 - System is directed in increasing the entropy.
 - In the thermal equilibrium state, entropy is the maximum.
 - The criterion: virtual variation $\Delta S \leq 0$

§3.1 Criterion of thermal equilibrium

- How to determine the thermal equilibrium **mathematically**?
- 2nd law of thermodynamics
 - Entropy never decreases for isolated system.
 - System is directed in increasing the entropy.
 - In the thermal equilibrium state, entropy is the maximum.
 - The criterion: virtual variation $\Delta S \leq 0$
- $\Delta S = 0$ neutral equilibrium,

§3.1 Criterion of thermal equilibrium

- How to determine the thermal equilibrium **mathematically**?
- 2nd law of thermodynamics
 - Entropy never decreases for isolated system.
 - System is directed in increasing the entropy.
 - In the thermal equilibrium state, entropy is the maximum.
 - The criterion: virtual variation $\Delta S \leq 0$
- $\Delta S = 0$ neutral equilibrium,
 $\Delta S < 0$ stable equilibrium.

Criterion of thermal equilibrium

- 2nd Taylor expansion of ΔS :

Criterion of thermal equilibrium

- 2nd Taylor expansion of ΔS :

$$\Delta S = \sum_i \frac{\partial S}{\partial x_i} \delta x_i + \frac{1}{2} \sum_{i,j} \frac{\partial^2 S}{\partial x_i \partial x_j} \delta x_i \delta x_j$$

Criterion of thermal equilibrium

- 2nd Taylor expansion of ΔS :

$$\begin{aligned}\Delta S &= \sum_i \frac{\partial S}{\partial x_i} \delta x_i + \frac{1}{2} \sum_{i,j} \frac{\partial^2 S}{\partial x_i \partial x_j} \delta x_i \delta x_j \\ &\equiv \delta S + \frac{1}{2} \delta^2 S\end{aligned}$$

Criterion of thermal equilibrium

- 2nd Taylor expansion of ΔS :

$$\begin{aligned}\Delta S &= \sum_i \frac{\partial S}{\partial x_i} \delta x_i + \frac{1}{2} \sum_{i,j} \frac{\partial^2 S}{\partial x_i \partial x_j} \delta x_i \delta x_j \\ &\equiv \delta S + \frac{1}{2} \delta^2 S\end{aligned}$$

- $\delta S = 0$, exists extremum;

Criterion of thermal equilibrium

- 2nd Taylor expansion of ΔS :

$$\begin{aligned}\Delta S &= \sum_i \frac{\partial S}{\partial x_i} \delta x_i + \frac{1}{2} \sum_{i,j} \frac{\partial^2 S}{\partial x_i \partial x_j} \delta x_i \delta x_j \\ &\equiv \delta S + \frac{1}{2} \delta^2 S\end{aligned}$$

- $\delta S = 0$, exists extremum;
 $\delta^2 S < 0$, exists maximum,

Criterion of thermal equilibrium

- 2nd Taylor expansion of ΔS :

$$\begin{aligned}\Delta S &= \sum_i \frac{\partial S}{\partial x_i} \delta x_i + \frac{1}{2} \sum_{i,j} \frac{\partial^2 S}{\partial x_i \partial x_j} \delta x_i \delta x_j \\ &\equiv \delta S + \frac{1}{2} \delta^2 S\end{aligned}$$

- $\delta S = 0$, exists extremum;
 $\delta^2 S < 0$, exists maximum, and if there are several maxima, the biggest one is the stable equilibrium state, and the others are semi-stable state.

Criterion of thermal equilibrium

- 2nd Taylor expansion of ΔS :

$$\begin{aligned}\Delta S &= \sum_i \frac{\partial S}{\partial x_i} \delta x_i + \frac{1}{2} \sum_{i,j} \frac{\partial^2 S}{\partial x_i \partial x_j} \delta x_i \delta x_j \\ &\equiv \delta S + \frac{1}{2} \delta^2 S\end{aligned}$$

- $\delta S = 0$, exists extremum;
 $\delta^2 S < 0$, exists maximum, and if there are several maxima, the biggest one is the stable equilibrium state, and the others are semi-stable state.
If $\delta S = 0$ and $\delta^2 S = 0$, higher order is needed.

Criterion of thermal equilibrium

- 2nd Taylor expansion of ΔS :

$$\begin{aligned}\Delta S &= \sum_i \frac{\partial S}{\partial x_i} \delta x_i + \frac{1}{2} \sum_{i,j} \frac{\partial^2 S}{\partial x_i \partial x_j} \delta x_i \delta x_j \\ &\equiv \delta S + \frac{1}{2} \delta^2 S\end{aligned}$$

- $\delta S = 0$, exists extremum;
 $\delta^2 S < 0$, exists maximum, and if there are several maxima, the biggest one is the stable equilibrium state, and the others are semi-stable state.
If $\delta S = 0$ and $\delta^2 S = 0$, higher order is needed.
- Other criterion:

Criterion of thermal equilibrium

- 2nd Taylor expansion of ΔS :

$$\begin{aligned}\Delta S &= \sum_i \frac{\partial S}{\partial x_i} \delta x_i + \frac{1}{2} \sum_{i,j} \frac{\partial^2 S}{\partial x_i \partial x_j} \delta x_i \delta x_j \\ &\equiv \delta S + \frac{1}{2} \delta^2 S\end{aligned}$$

- $\delta S = 0$, exists extremum;
 $\delta^2 S < 0$, exists maximum, and if there are several maxima, the biggest one is the stable equilibrium state, and the others are semi-stable state.
If $\delta S = 0$ and $\delta^2 S = 0$, higher order is needed.
- Other criterion:
in isothermal and isochoric process $\Delta F > 0$,

Criterion of thermal equilibrium

- 2nd Taylor expansion of ΔS :

$$\begin{aligned}\Delta S &= \sum_i \frac{\partial S}{\partial x_i} \delta x_i + \frac{1}{2} \sum_{i,j} \frac{\partial^2 S}{\partial x_i \partial x_j} \delta x_i \delta x_j \\ &\equiv \delta S + \frac{1}{2} \delta^2 S\end{aligned}$$

- $\delta S = 0$, exists extremum;
 $\delta^2 S < 0$, exists maximum, and if there are several maxima, the biggest one is the stable equilibrium state, and the others are semi-stable state.
If $\delta S = 0$ and $\delta^2 S = 0$, higher order is needed.
- Other criterion:
in isothermal and isochoric process $\Delta F > 0$,
in isothermal and isobaric process: $\Delta G > 0$, ...

Example: Condition for the isolated uniform thermal equilibrium state and the stability criterion.

Example: Condition for the isolated uniform thermal equilibrium state and the stability criterion.

- The isolated system's small part (T, p) ,
and the other (almost whole) part (T_0, p_0) ,

Example: Condition for the isolated uniform thermal equilibrium state and the stability criterion.

- The isolated system's small part (T, p) , and the other (almost whole) part (T_0, p_0) , a virtual variation $\delta U, \delta V$ and $\delta U_0, \delta V_0$.

Example: Condition for the isolated uniform thermal equilibrium state and the stability criterion.

- The isolated system's small part (T, p) , and the other (almost whole) part (T_0, p_0) , a virtual variation $\delta U, \delta V$ and $\delta U_0, \delta V_0$.
- The whole system does not change (constraint):
$$\delta U + \delta U_0 = 0,$$
$$\delta V + \delta V_0 = 0.$$

Example: Condition for the isolated uniform thermal equilibrium state and the stability criterion.

- The isolated system's small part (T, p) , and the other (almost whole) part (T_0, p_0) , a virtual variation $\delta U, \delta V$ and $\delta U_0, \delta V_0$.
- The whole system does not change (constraint):
$$\delta U + \delta U_0 = 0,$$
$$\delta V + \delta V_0 = 0.$$
- Total change of the entropy: $\Delta \tilde{S} = \Delta S + \Delta S_0$.

Example: Condition for the isolated uniform thermal equilibrium state and the stability criterion.

- The isolated system's small part (T, p) , and the other (almost whole) part (T_0, p_0) , a virtual variation $\delta U, \delta V$ and $\delta U_0, \delta V_0$.
- The whole system does not change (constraint):
$$\delta U + \delta U_0 = 0,$$
$$\delta V + \delta V_0 = 0.$$
- Total change of the entropy: $\Delta \tilde{S} = \Delta S + \Delta S_0$.
- Condition for equilibrium: $\Delta \tilde{S} < 0$, i.e.,
$$\Delta S + \Delta S_0 < 0, \text{ or } \delta \tilde{S} = 0, \delta^2 \tilde{S} < 0.$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU + pdV}{T}$, then

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU + pdV}{T}$, then
$$\delta\tilde{S} = \frac{\delta U + p\delta V}{T} + \frac{\delta U_0 + p_0\delta V_0}{T_0}$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU + pdV}{T}$, then

$$\delta \tilde{S} = \frac{\delta U + p\delta V}{T} + \frac{\delta U_0 + p_0\delta V_0}{T_0} = \frac{\delta U + p\delta V}{T} + \frac{-\delta U - p_0\delta V}{T_0}$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU + pdV}{T}$, then
$$\delta \tilde{S} = \frac{\delta U + p\delta V}{T} + \frac{\delta U_0 + p_0\delta V_0}{T_0} = \frac{\delta U + p\delta V}{T} + \frac{-\delta U - p_0\delta V}{T_0}$$
$$= \delta U \left(\frac{1}{T} - \frac{1}{T_0} \right) + \delta V \left(\frac{p}{T} - \frac{p_0}{T_0} \right)$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU + pdV}{T}$, then

$$\begin{aligned}\delta\tilde{S} &= \frac{\delta U + p\delta V}{T} + \frac{\delta U_0 + p_0\delta V_0}{T_0} = \frac{\delta U + p\delta V}{T} + \frac{-\delta U - p_0\delta V}{T_0} \\ &= \delta U \left(\frac{1}{T} - \frac{1}{T_0} \right) + \delta V \left(\frac{p}{T} - \frac{p_0}{T_0} \right) = 0.\end{aligned}$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU + pdV}{T}$, then
$$\delta \tilde{S} = \frac{\delta U + p\delta V}{T} + \frac{\delta U_0 + p_0\delta V_0}{T_0} = \frac{\delta U + p\delta V}{T} + \frac{-\delta U - p_0\delta V}{T_0}$$
$$= \delta U \left(\frac{1}{T} - \frac{1}{T_0} \right) + \delta V \left(\frac{p}{T} - \frac{p_0}{T_0} \right) = 0.$$
- $\therefore \delta U$ and δV are independent, $\therefore T = T_0, p = p_0$.

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU+pdV}{T}$, then
$$\delta\tilde{S} = \frac{\delta U+p\delta V}{T} + \frac{\delta U_0+p_0\delta V_0}{T_0} = \frac{\delta U+p\delta V}{T} + \frac{-\delta U-p_0\delta V}{T_0}$$
$$= \delta U \left(\frac{1}{T} - \frac{1}{T_0} \right) + \delta V \left(\frac{p}{T} - \frac{p_0}{T_0} \right) = 0.$$
- $\therefore \delta U$ and δV are independent, $\therefore T = T_0, p = p_0$.
- Stability criterion: $\delta^2\tilde{S} < 0$.

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU + pdV}{T}$, then
$$\delta \tilde{S} = \frac{\delta U + p\delta V}{T} + \frac{\delta U_0 + p_0\delta V_0}{T_0} = \frac{\delta U + p\delta V}{T} + \frac{-\delta U - p_0\delta V}{T_0}$$
$$= \delta U \left(\frac{1}{T} - \frac{1}{T_0} \right) + \delta V \left(\frac{p}{T} - \frac{p_0}{T_0} \right) = 0.$$
- $\therefore \delta U$ and δV are independent, $\therefore T = T_0, p = p_0$.
- Stability criterion: $\delta^2 \tilde{S} < 0$.
- $\delta^2 S = \frac{\partial^2 S}{\partial U^2}(\delta U)^2 + 2\frac{\partial^2 S}{\partial U \partial V}\delta U\delta V + \frac{\partial^2 S}{\partial V^2}(\delta V)^2,$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU+pdV}{T}$, then

$$\delta\tilde{S} = \frac{\delta U+p\delta V}{T} + \frac{\delta U_0+p_0\delta V_0}{T_0} = \frac{\delta U+p\delta V}{T} + \frac{-\delta U-p_0\delta V}{T_0}$$

$$= \delta U \left(\frac{1}{T} - \frac{1}{T_0} \right) + \delta V \left(\frac{p}{T} - \frac{p_0}{T_0} \right) = 0.$$
- $\therefore \delta U$ and δV are independent, $\therefore T = T_0, p = p_0$.
- Stability criterion: $\delta^2\tilde{S} < 0$.
- $\delta^2 S = \frac{\partial^2 S}{\partial U^2}(\delta U)^2 + 2\frac{\partial^2 S}{\partial U\partial V}\delta U\delta V + \frac{\partial^2 S}{\partial V^2}(\delta V)^2,$
 $\delta^2 S_0 = \frac{\partial^2 S_0}{\partial U_0^2}(\delta U_0)^2 + 2\frac{\partial^2 S_0}{\partial U_0\partial V_0}\delta U_0\delta V_0 + \frac{\partial^2 S_0}{\partial V_0^2}(\delta V_0)^2.$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU+pdV}{T}$, then

$$\delta\tilde{S} = \frac{\delta U+p\delta V}{T} + \frac{\delta U_0+p_0\delta V_0}{T_0} = \frac{\delta U+p\delta V}{T} + \frac{-\delta U-p_0\delta V}{T_0}$$

$$= \delta U \left(\frac{1}{T} - \frac{1}{T_0} \right) + \delta V \left(\frac{p}{T} - \frac{p_0}{T_0} \right) = 0.$$
- $\therefore \delta U$ and δV are independent, $\therefore T = T_0, p = p_0$.
- Stability criterion: $\delta^2\tilde{S} < 0$.
- $$\delta^2 S = \frac{\partial^2 S}{\partial U^2}(\delta U)^2 + 2\frac{\partial^2 S}{\partial U\partial V}\delta U\delta V + \frac{\partial^2 S}{\partial V^2}(\delta V)^2,$$

$$\delta^2 S_0 = \frac{\partial^2 S_0}{\partial U_0^2}(\delta U_0)^2 + 2\frac{\partial^2 S_0}{\partial U_0\partial V_0}\delta U_0\delta V_0 + \frac{\partial^2 S_0}{\partial V_0^2}(\delta V_0)^2.$$
- The intensive quantities for the two system:
 $s \sim s_0, u \sim u_0, v \sim v_0$, where $S = ns$.

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU+pdV}{T}$, then

$$\delta\tilde{S} = \frac{\delta U+p\delta V}{T} + \frac{\delta U_0+p_0\delta V_0}{T_0} = \frac{\delta U+p\delta V}{T} + \frac{-\delta U-p_0\delta V}{T_0}$$

$$= \delta U \left(\frac{1}{T} - \frac{1}{T_0} \right) + \delta V \left(\frac{p}{T} - \frac{p_0}{T_0} \right) = 0.$$
- $\therefore \delta U$ and δV are independent, $\therefore T = T_0, p = p_0$.
- Stability criterion: $\delta^2\tilde{S} < 0$.
- $\delta^2 S = \frac{\partial^2 S}{\partial U^2}(\delta U)^2 + 2\frac{\partial^2 S}{\partial U\partial V}\delta U\delta V + \frac{\partial^2 S}{\partial V^2}(\delta V)^2$,
 $\delta^2 S_0 = \frac{\partial^2 S_0}{\partial U_0^2}(\delta U_0)^2 + 2\frac{\partial^2 S_0}{\partial U_0\partial V_0}\delta U_0\delta V_0 + \frac{\partial^2 S_0}{\partial V_0^2}(\delta V_0)^2$.
- The intensive quantities for the two system:
 $s \sim s_0, u \sim u_0, v \sim v_0$, where $S = ns$.
- $\frac{\partial^2 S}{\partial U^2} = \frac{1}{n} \frac{\partial^2 s}{\partial u^2}$, ($\delta U = -\delta U_0, \delta V = -\delta V_0$.)

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- Basic equation: $dS = \frac{dU+pdV}{T}$, then

$$\delta\tilde{S} = \frac{\delta U+p\delta V}{T} + \frac{\delta U_0+p_0\delta V_0}{T_0} = \frac{\delta U+p\delta V}{T} + \frac{-\delta U-p_0\delta V}{T_0}$$

$$= \delta U \left(\frac{1}{T} - \frac{1}{T_0} \right) + \delta V \left(\frac{p}{T} - \frac{p_0}{T_0} \right) = 0.$$
- $\because \delta U$ and δV are independent, $\therefore T = T_0, p = p_0$.
- Stability criterion: $\delta^2\tilde{S} < 0$.
- $$\delta^2 S = \frac{\partial^2 S}{\partial U^2}(\delta U)^2 + 2\frac{\partial^2 S}{\partial U\partial V}\delta U\delta V + \frac{\partial^2 S}{\partial V^2}(\delta V)^2,$$

$$\delta^2 S_0 = \frac{\partial^2 S_0}{\partial U_0^2}(\delta U_0)^2 + 2\frac{\partial^2 S_0}{\partial U_0\partial V_0}\delta U_0\delta V_0 + \frac{\partial^2 S_0}{\partial V_0^2}(\delta V_0)^2.$$
- The intensive quantities for the two system:
 $s \sim s_0, u \sim u_0, v \sim v_0$, where $S = ns$.
- $$\frac{\partial^2 S}{\partial U^2} = \frac{1}{n} \frac{\partial^2 s}{\partial u^2}, (\delta U = -\delta U_0, \delta V = -\delta V_0.)$$

$$\Rightarrow \delta^2 S_0 \ll \delta^2 S$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- $\delta^2 \tilde{S} \simeq \delta^2 S = \frac{\partial^2 S}{\partial U^2} (\delta U)^2 + 2 \frac{\partial^2 S}{\partial U \partial V} \delta U \delta V + \frac{\partial^2 S}{\partial V^2} (\delta V)^2.$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- $$\begin{aligned}\delta^2 \tilde{S} &\simeq \delta^2 S = \frac{\partial^2 S}{\partial U^2} (\delta U)^2 + 2 \frac{\partial^2 S}{\partial U \partial V} \delta U \delta V + \frac{\partial^2 S}{\partial V^2} (\delta V)^2. \\ &= \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial U} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial U} \delta V \right] \delta U \\ &\quad + \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial V} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial V} \delta V \right] \delta V,\end{aligned}$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

$$\begin{aligned}
 \bullet \quad \delta^2 \tilde{S} &\simeq \delta^2 S = \frac{\partial^2 S}{\partial U^2} (\delta U)^2 + 2 \frac{\partial^2 S}{\partial U \partial V} \delta U \delta V + \frac{\partial^2 S}{\partial V^2} (\delta V)^2. \\
 &= \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial U} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial U} \delta V \right] \delta U \\
 &\quad + \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial V} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial V} \delta V \right] \delta V, \\
 \left(dS = \frac{dU + p dV}{T} \Rightarrow \left(\frac{\partial S}{\partial U} \right)_V &= \frac{1}{T}, \left(\frac{\partial S}{\partial V} \right)_U = \frac{p}{T}. \right)
 \end{aligned}$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- $$\begin{aligned}\delta^2 \tilde{S} &\simeq \delta^2 S = \frac{\partial^2 S}{\partial U^2} (\delta U)^2 + 2 \frac{\partial^2 S}{\partial U \partial V} \delta U \delta V + \frac{\partial^2 S}{\partial V^2} (\delta V)^2. \\ &= \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial U} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial U} \delta V \right] \delta U \\ &\quad + \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial V} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial V} \delta V \right] \delta V, \\ &\left(dS = \frac{dU + p dV}{T} \Rightarrow \left(\frac{\partial S}{\partial U} \right)_V = \frac{1}{T}, \left(\frac{\partial S}{\partial V} \right)_U = \frac{p}{T} \right) \\ \delta^2 S &= \left[\frac{\partial}{\partial U} \left(\frac{1}{T} \right) \delta U + \frac{\partial}{\partial V} \left(\frac{1}{T} \right) \delta V \right] \delta U \\ &\quad + \left[\frac{\partial}{\partial U} \left(\frac{p}{T} \right) \delta U + \frac{\partial}{\partial V} \left(\frac{p}{T} \right) \delta V \right] \delta V\end{aligned}$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- $$\begin{aligned}\delta^2 \tilde{S} &\simeq \delta^2 S = \frac{\partial^2 S}{\partial U^2} (\delta U)^2 + 2 \frac{\partial^2 S}{\partial U \partial V} \delta U \delta V + \frac{\partial^2 S}{\partial V^2} (\delta V)^2. \\ &= \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial U} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial U} \delta V \right] \delta U \\ &\quad + \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial V} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial V} \delta V \right] \delta V, \\ \left(dS = \frac{dU + p dV}{T} \Rightarrow \left(\frac{\partial S}{\partial U} \right)_V &= \frac{1}{T}, \left(\frac{\partial S}{\partial V} \right)_U = \frac{p}{T}. \right) \\ \delta^2 S &= \left[\frac{\partial}{\partial U} \left(\frac{1}{T} \right) \delta U + \frac{\partial}{\partial V} \left(\frac{1}{T} \right) \delta V \right] \delta U \\ &\quad + \left[\frac{\partial}{\partial U} \left(\frac{p}{T} \right) \delta U + \frac{\partial}{\partial V} \left(\frac{p}{T} \right) \delta V \right] \delta V \\ &= \delta \left(\frac{1}{T} \right) \delta U + \delta \left(\frac{p}{T} \right) \delta V.\end{aligned}$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- $$\delta^2 \tilde{S} \simeq \delta^2 S = \frac{\partial^2 S}{\partial U^2} (\delta U)^2 + 2 \frac{\partial^2 S}{\partial U \partial V} \delta U \delta V + \frac{\partial^2 S}{\partial V^2} (\delta V)^2.$$

$$= \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial U} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial U} \delta V \right] \delta U$$

$$+ \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial V} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial V} \delta V \right] \delta V,$$

$$\left(dS = \frac{dU + p dV}{T} \Rightarrow \left(\frac{\partial S}{\partial U} \right)_V = \frac{1}{T}, \left(\frac{\partial S}{\partial V} \right)_U = \frac{p}{T} \right)$$
- $$\delta^2 S = \left[\frac{\partial}{\partial U} \left(\frac{1}{T} \right) \delta U + \frac{\partial}{\partial V} \left(\frac{1}{T} \right) \delta V \right] \delta U$$

$$+ \left[\frac{\partial}{\partial U} \left(\frac{p}{T} \right) \delta U + \frac{\partial}{\partial V} \left(\frac{p}{T} \right) \delta V \right] \delta V$$

$$= \delta \left(\frac{1}{T} \right) \delta U + \delta \left(\frac{p}{T} \right) \delta V.$$
- Convert to (T, V) :

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- $$\delta^2 \tilde{S} \simeq \delta^2 S = \frac{\partial^2 S}{\partial U^2} (\delta U)^2 + 2 \frac{\partial^2 S}{\partial U \partial V} \delta U \delta V + \frac{\partial^2 S}{\partial V^2} (\delta V)^2.$$

$$= \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial U} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial U} \delta V \right] \delta U$$

$$+ \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial V} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial V} \delta V \right] \delta V,$$

$$\left(dS = \frac{dU + p dV}{T} \Rightarrow \left(\frac{\partial S}{\partial U} \right)_V = \frac{1}{T}, \left(\frac{\partial S}{\partial V} \right)_U = \frac{p}{T} \right)$$
- $$\delta^2 S = \left[\frac{\partial}{\partial U} \left(\frac{1}{T} \right) \delta U + \frac{\partial}{\partial V} \left(\frac{1}{T} \right) \delta V \right] \delta U$$

$$+ \left[\frac{\partial}{\partial U} \left(\frac{p}{T} \right) \delta U + \frac{\partial}{\partial V} \left(\frac{p}{T} \right) \delta V \right] \delta V$$

$$= \delta \left(\frac{1}{T} \right) \delta U + \delta \left(\frac{p}{T} \right) \delta V.$$
- Convert to (T, V) :
- $$\delta U = \left(\frac{\partial U}{\partial T} \right)_V \delta T + \left(\frac{\partial U}{\partial V} \right)_T \delta V$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

- $$\delta^2 \tilde{S} \simeq \delta^2 S = \frac{\partial^2 S}{\partial U^2} (\delta U)^2 + 2 \frac{\partial^2 S}{\partial U \partial V} \delta U \delta V + \frac{\partial^2 S}{\partial V^2} (\delta V)^2.$$

$$= \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial U} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial U} \delta V \right] \delta U$$

$$+ \left[\frac{\partial}{\partial U} \frac{\partial S}{\partial V} \delta U + \frac{\partial}{\partial V} \frac{\partial S}{\partial V} \delta V \right] \delta V,$$

$$\left(dS = \frac{dU + pdV}{T} \Rightarrow \left(\frac{\partial S}{\partial U} \right)_V = \frac{1}{T}, \left(\frac{\partial S}{\partial V} \right)_U = \frac{p}{T} \right)$$
- $$\delta^2 S = \left[\frac{\partial}{\partial U} \left(\frac{1}{T} \right) \delta U + \frac{\partial}{\partial V} \left(\frac{1}{T} \right) \delta V \right] \delta U$$

$$+ \left[\frac{\partial}{\partial U} \left(\frac{p}{T} \right) \delta U + \frac{\partial}{\partial V} \left(\frac{p}{T} \right) \delta V \right] \delta V$$

$$= \delta \left(\frac{1}{T} \right) \delta U + \delta \left(\frac{p}{T} \right) \delta V.$$
- Convert to (T, V) :
- $$\delta U = \left(\frac{\partial U}{\partial T} \right)_V \delta T + \left(\frac{\partial U}{\partial V} \right)_T \delta V$$

$$= C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V.$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

$$\delta^2 S = \delta \left(\frac{1}{T} \right) \delta U + \delta \left(\frac{p}{T} \right) \delta V,$$

$$\delta U = C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V$$

- $$\delta \frac{1}{T} = \left(\frac{\partial}{\partial T} \frac{1}{T} \right)_V \delta T + \left(\frac{\partial}{\partial V} \frac{1}{T} \right)_T \delta V$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

$$\delta^2 S = \delta \left(\frac{1}{T} \right) \delta U + \delta \left(\frac{p}{T} \right) \delta V,$$

$$\delta U = C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V$$

- $$\delta \frac{1}{T} = \left(\frac{\partial}{\partial T} \frac{1}{T} \right)_V \delta T + \left(\frac{\partial}{\partial V} \frac{1}{T} \right)_T \delta V = -\frac{1}{T^2} \delta T.$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

$$\delta^2 S = \delta \left(\frac{1}{T} \right) \delta U + \delta \left(\frac{p}{T} \right) \delta V,$$

$$\delta U = C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V$$

- $\delta \frac{1}{T} = \left(\frac{\partial}{\partial T} \frac{1}{T} \right)_V \delta T + \left(\frac{\partial}{\partial V} \frac{1}{T} \right)_T \delta V = -\frac{1}{T^2} \delta T.$
- $\delta \frac{p}{T} = \left(\frac{\partial}{\partial T} \frac{p}{T} \right)_V \delta T + \left(\frac{\partial}{\partial V} \frac{p}{T} \right)_T \delta V$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

$$\delta^2 S = \delta \left(\frac{1}{T} \right) \delta U + \delta \left(\frac{p}{T} \right) \delta V,$$

$$\delta U = C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V$$

- $\delta \frac{1}{T} = \left(\frac{\partial}{\partial T} \frac{1}{T} \right)_V \delta T + \left(\frac{\partial}{\partial V} \frac{1}{T} \right)_T \delta V = -\frac{1}{T^2} \delta T.$
- $\delta \frac{p}{T} = \left(\frac{\partial}{\partial T} \frac{p}{T} \right)_V \delta T + \left(\frac{\partial}{\partial V} \frac{p}{T} \right)_T \delta V$
$$= \left[p \left(\frac{\partial}{\partial T} \frac{1}{T} \right)_V + \frac{1}{T} \left(\frac{\partial p}{\partial T} \right)_V \right] \delta T$$
$$+ \left[p \left(\frac{\partial}{\partial V} \frac{1}{T} \right)_T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \right] \delta V$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

$$\delta^2 S = \delta \left(\frac{1}{T} \right) \delta U + \delta \left(\frac{p}{T} \right) \delta V,$$

$$\delta U = C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V$$

- $\delta \frac{1}{T} = \left(\frac{\partial}{\partial T} \frac{1}{T} \right)_V \delta T + \left(\frac{\partial}{\partial V} \frac{1}{T} \right)_T \delta V = -\frac{1}{T^2} \delta T.$
- $\delta \frac{p}{T} = \left(\frac{\partial}{\partial T} \frac{p}{T} \right)_V \delta T + \left(\frac{\partial}{\partial V} \frac{p}{T} \right)_T \delta V$

$$= \left[p \left(\frac{\partial}{\partial T} \frac{1}{T} \right)_V + \frac{1}{T} \left(\frac{\partial p}{\partial T} \right)_V \right] \delta T$$

$$+ \left[p \left(\frac{\partial}{\partial V} \frac{1}{T} \right)_T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \right] \delta V$$

$$= \left[-\frac{p}{T^2} + \frac{1}{T} \left(\frac{\partial p}{\partial T} \right)_V \right] \delta T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \delta V$$

Condition for the isolated uniform thermal equilibrium state and the stability criterion. (Continuing...)

$$\delta^2 S = \delta \left(\frac{1}{T} \right) \delta U + \delta \left(\frac{p}{T} \right) \delta V,$$

$$\delta U = C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V$$

- $\delta \frac{1}{T} = \left(\frac{\partial}{\partial T} \frac{1}{T} \right)_V \delta T + \left(\frac{\partial}{\partial V} \frac{1}{T} \right)_T \delta V = -\frac{1}{T^2} \delta T.$
- $\delta \frac{p}{T} = \left(\frac{\partial}{\partial T} \frac{p}{T} \right)_V \delta T + \left(\frac{\partial}{\partial V} \frac{p}{T} \right)_T \delta V$

$$= \left[p \left(\frac{\partial}{\partial T} \frac{1}{T} \right)_V + \frac{1}{T} \left(\frac{\partial p}{\partial T} \right)_V \right] \delta T$$

$$+ \left[p \left(\frac{\partial}{\partial V} \frac{1}{T} \right)_T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \right] \delta V$$

$$= \left[-\frac{p}{T^2} + \frac{1}{T} \left(\frac{\partial p}{\partial T} \right)_V \right] \delta T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \delta V$$

$$= \frac{1}{T^2} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \delta V.$$

The stability criterion. (Continuing...)

$$\bullet \delta^2 \tilde{S} \simeq \delta^2 S = -\frac{1}{T^2} \delta T \left\{ C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V \right\} \\ + \left\{ \frac{1}{T^2} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \delta V \right\} \delta V$$

The stability criterion. (Continuing...)

$$\begin{aligned} \bullet \quad \delta^2 \tilde{S} &\simeq \delta^2 S = -\frac{1}{T^2} \delta T \left\{ C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V \right\} \\ &+ \left\{ \frac{1}{T^2} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \delta V \right\} \delta V \\ &= -\frac{C_V}{T^2} (\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T (\delta V)^2 \end{aligned}$$

The stability criterion. (Continuing...)

$$\begin{aligned} \bullet \quad \delta^2 \tilde{S} &\simeq \delta^2 S = -\frac{1}{T^2} \delta T \left\{ C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V \right\} \\ &+ \left\{ \frac{1}{T^2} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \delta V \right\} \delta V \\ &= -\frac{C_V}{T^2} (\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T (\delta V)^2 < 0. \end{aligned}$$

The stability criterion. (Continuing...)

- $\delta^2 \tilde{S} \simeq \delta^2 S = -\frac{1}{T^2} \delta T \left\{ C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V \right\}$
 $+ \left\{ \frac{1}{T^2} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \delta V \right\} \delta V$
 $= -\frac{C_V}{T^2} (\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T (\delta V)^2 < 0.$
- δT and δV is independent, and $(\delta T)^2 > 0$, $(\delta V)^2 > 0$.

The stability criterion. (Continuing...)

- $\delta^2 \tilde{S} \simeq \delta^2 S = -\frac{1}{T^2} \delta T \left\{ C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V \right\}$
 $+ \left\{ \frac{1}{T^2} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \delta V \right\} \delta V$
 $= -\frac{C_V}{T^2} (\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T (\delta V)^2 < 0.$
- δT and δV is independent, and $(\delta T)^2 > 0$, $(\delta V)^2 > 0$.
- $\therefore C_V > 0$, $\left(\frac{\partial p}{\partial V} \right)_T < 0$.

The stability criterion. (Continuing...)

- $$\delta^2 \tilde{S} \simeq \delta^2 S = -\frac{1}{T^2} \delta T \left\{ C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V \right\}$$

$$+ \left\{ \frac{1}{T^2} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \delta V \right\} \delta V$$

$$= -\frac{C_V}{T^2} (\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T (\delta V)^2 < 0.$$
- δT and δV is independent, and $(\delta T)^2 > 0$, $(\delta V)^2 > 0$.
- $\therefore C_V > 0$, $\left(\frac{\partial p}{\partial V} \right)_T < 0$.
- Meaning of $C_V > 0$: suppose $T \gtrsim T_0$, small part loses heat, as $C_V > 0$, temperature decreases, system goes back to the equilibrium.

The stability criterion. (Continuing...)

- $$\delta^2 \tilde{S} \simeq \delta^2 S = -\frac{1}{T^2} \delta T \left\{ C_V \delta T + \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta V \right\}$$

$$+ \left\{ \frac{1}{T^2} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] \delta T + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T \delta V \right\} \delta V$$

$$= -\frac{C_V}{T^2} (\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V} \right)_T (\delta V)^2 < 0.$$
- δT and δV is independent, and $(\delta T)^2 > 0$, $(\delta V)^2 > 0$.
- $\therefore C_V > 0$, $\left(\frac{\partial p}{\partial V} \right)_T < 0$.
- Meaning of $C_V > 0$: suppose $T \gtrsim T_0$, small part loses heat, as $C_V > 0$, temperature decreases, system goes back to the equilibrium.
- $\left(\frac{\partial p}{\partial V} \right)_T < 0$: imaging the small part shrinks, $\Delta V < 0$, as $\left(\frac{\partial p}{\partial V} \right)_T < 0$, $\Delta p > 0$, $p > p_0$, small part expands.

Table of contents

① Phase Transition of Single-Component System

- 3.1 Criterion of thermal equilibrium
- 3.2 Basic equations of open system
- 3.3 Equilibrium of single-component multi-phase system
- 3.4 Properties of equilibrium of s-c multi-phase system
- 3.5 Critical point and phase change between gas and liquid
- 3.7 Classification of the phase transition
- 3.9 Landau's approximation for the continuous phase transition

§3.2 Basic equations of open system

§3.2 Basic equations of open system

- Single-component system: one chemical component,

§3.2 Basic equations of open system

- Single-component system: one chemical component, with several phases
→ single-component multi-phase system.

§3.2 Basic equations of open system

- Single-component system: one chemical component, with several phases
→ single-component multi-phase system.
- For each phase, there is a set of parameters $(p, V, T \dots)$.

§3.2 Basic equations of open system

- Single-component system: one chemical component, with several phases
→ single-component multi-phase system.
- For each phase, there is a set of parameters $(p, V, T \dots)$.
- Major difference: amount of substance is not conserved for one phase.

§3.2 Basic equations of open system

- Single-component system: one chemical component, with several phases
→ single-component multi-phase system.
- For each phase, there is a set of parameters $(p, V, T \dots)$.
- Major difference: amount of substance is not conserved for one phase.
- Chemical parameter n should be introduced.

§3.2 Basic equations of open system

- Single-component system: one chemical component, with several phases
→ single-component multi-phase system.
- For each phase, there is a set of parameters $(p, V, T \dots)$.
- Major difference: amount of substance is not conserved for one phase.
- Chemical parameter n should be introduced.
- Define chemical potential:
$$\mu \equiv \frac{G}{n} = G_m = u - Ts + pv.$$

§3.2 Basic equations of open system

- Single-component system: one chemical component, with several phases
→ single-component multi-phase system.
- For each phase, there is a set of parameters $(p, V, T \dots)$.
- Major difference: amount of substance is not conserved for one phase.
- Chemical parameter n should be introduced.
- Define chemical potential:
$$\mu \equiv \frac{G}{n} = G_m = u - Ts + pv.$$
- $d\mu = V_m dp - S_m dT$ or $vd p - sd T$.

Basic equations of open system

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.
- $dU = d(nu) = n(Tds - pdv) + udn$

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.
- $dU = d(nu) = n(Tds - pdv) + udn$
 $= n(Tds - pdv) + (\mu + Ts - pv)dn$

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.
- $dU = d(nu) = n(Tds - pdv) + udn$
 $= n(Tds - pdv) + (\mu + Ts - pv)dn$
 $= T[d(ns) - sdn] - p[d(nv) - vdn] + (\mu + Ts - pv)dn$

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.
- $dU = d(nu) = n(Tds - pdv) + udn$
 $= n(Tds - pdv) + (\mu + Ts - pv)dn$
 $= T[d(ns) - sdn] - p[d(nv) - vdn] + (\mu + Ts - pv)dn$
 $= TdS - Tsdn - pdV + pvdn + \mu dn + Tsdn - pvdn$

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.
- $dU = d(nu) = n(Tds - pdv) + udn$
 $= n(Tds - pdv) + (\mu + Ts - pv)dn$
 $= T[d(ns) - sdn] - p[d(nv) - vdn] + (\mu + Ts - pv)dn$
 $= TdS - \cancel{Tsdn} - pdV + \cancel{pvdn} + \mu dn + \cancel{Tsdn} - \cancel{pvdn}$
 $= TdS - pdV + \mu dn.$

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.
- $dU = d(nu) = n(Tds - pdv) + udn$
 $= n(Tds - pdv) + (\mu + Ts - pv)dn$
 $= T[d(ns) - sdn] - p[d(nv) - vdn] + (\mu + Ts - pv)dn$
 $= TdS - Tsdn - pdV + pvdn + \mu dn + Tsdn - pvdn$
 $= TdS - pdV + \mu dn.$
- Similarly, $dH = TdS + Vdp + \mu dn,$

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.
- $dU = d(nu) = n(Tds - pdv) + udn$
 $= n(Tds - pdv) + (\mu + Ts - pv)dn$
 $= T[d(ns) - sdn] - p[d(nv) - vdn] + (\mu + Ts - pv)dn$
 $= TdS - Tsdn - pdV + pvdn + \mu dn + Tsdn - pvdn$
 $= TdS - pdV + \mu dn.$
- Similarly, $dH = TdS + Vdp + \mu dn$,
 $dF = -SdT - pdV + \mu dn,$

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.
- $dU = d(nu) = n(Tds - pdv) + udn$
 $= n(Tds - pdv) + (\mu + Ts - pv)dn$
 $= T[d(ns) - sdn] - p[d(nv) - vdn] + (\mu + Ts - pv)dn$
 $= TdS - Tsdn - pdV + pvdn + \mu dn + Tsdn - pvdn$
 $= TdS - pdV + \mu dn.$
- Similarly, $dH = TdS + Vdp + \mu dn,$
 $dF = -SdT - pdV + \mu dn,$
 $dG = -SdT + Vdp + \mu dn.$

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.
- $dU = d(nu) = n(Tds - pdv) + udn$
 $= n(Tds - pdv) + (\mu + Ts - pv)dn$
 $= T[d(ns) - sdn] - p[d(nv) - vdn] + (\mu + Ts - pv)dn$
 $= TdS - Tsdn - pdV + pvdn + \mu dn + Tsdn - pvdn$
 $= TdS - pdV + \mu dn.$
- Similarly, $dH = TdS + Vdp + \mu dn$,
 $dF = -SdT - pdV + \mu dn$,
 $dG = -SdT + Vdp + \mu dn.$
- Define a new function, grand thermodynamic potential $J \equiv F - \mu n$

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.
- $dU = d(nu) = n(Tds - pdv) + udn$
 $= n(Tds - pdv) + (\mu + Ts - pv)dn$
 $= T[d(ns) - sdn] - p[d(nv) - vdn] + (\mu + Ts - pv)dn$
 $= TdS - Tsdn - pdV + pvdn + \mu dn + Tsdn - pvdn$
 $= TdS - pdV + \mu dn.$
- Similarly, $dH = TdS + Vdp + \mu dn$,
 $dF = -SdT - pdV + \mu dn$,
 $dG = -SdT + Vdp + \mu dn.$
- Define a new function, grand thermodynamic potential $J \equiv F - \mu n = F - G = -pV$.

Basic equations of open system

- The internal energy $U = nu$, and $du = Tds - pdv$.
- $dU = d(nu) = n(Tds - pdv) + udn$
 $= n(Tds - pdv) + (\mu + Ts - pv)dn$
 $= T[d(ns) - sdn] - p[d(nv) - vdn] + (\mu + Ts - pv)dn$
 $= TdS - Tsdn - pdV + pvdn + \mu dn + Tsdn - pvdn$
 $= TdS - pdV + \mu dn.$
- Similarly, $dH = TdS + Vdp + \mu dn$,
 $dF = -SdT - pdV + \mu dn$,
 $dG = -SdT + Vdp + \mu dn.$
- Define a new function, grand thermodynamic potential $J \equiv F - \mu n = F - G = -pV$.
- $dJ = -SdT - pdV - nd\mu$.

Table of contents

① Phase Transition of Single-Component System

- 3.1 Criterion of thermal equilibrium
- 3.2 Basic equations of open system
- 3.3 Equilibrium of single-component multi-phase system
- 3.4 Properties of equilibrium of s-c multi-phase system
- 3.5 Critical point and phase change between gas and liquid
- 3.7 Classification of the phase transition
- 3.9 Landau's approximation for the continuous phase transition

§3.3 Equilibrium of single-component multi-phase system

§3.3 Equilibrium of single-component multi-phase system

- Considering isolated two-phase system, α, β .

§3.3 Equilibrium of single-component multi-phase system

- Considering isolated two-phase system, α, β .
- State parameter $(U^\alpha, V^\alpha, n^\alpha, T^\alpha), (U^\beta, V^\beta, n^\beta, T^\beta)$

§3.3 Equilibrium of single-component multi-phase system

- Considering isolated two-phase system, α, β .
- State parameter $(U^\alpha, V^\alpha, n^\alpha, T^\alpha), (U^\beta, V^\beta, n^\beta, T^\beta)$
- Constraints:
$$\begin{cases} U^\alpha + U^\beta = \text{Const.} \\ V^\alpha + V^\beta = \text{Const.} \\ n^\alpha + n^\beta = \text{Const.} \end{cases}$$

§3.3 Equilibrium of single-component multi-phase system

- Considering isolated two-phase system, α, β .
- State parameter $(U^\alpha, V^\alpha, n^\alpha, T^\alpha), (U^\beta, V^\beta, n^\beta, T^\beta)$
- Constraints:
$$\begin{cases} U^\alpha + U^\beta = \text{Const.} \\ V^\alpha + V^\beta = \text{Const.} \\ n^\alpha + n^\beta = \text{Const.} \end{cases} \Rightarrow \begin{cases} \delta U^\alpha + \delta U^\beta = 0 \\ \delta V^\alpha + \delta V^\beta = 0 \\ \delta n^\alpha + \delta n^\beta = 0 \end{cases} .$$

§3.3 Equilibrium of single-component multi-phase system

- Considering isolated two-phase system, α, β .
- State parameter $(U^\alpha, V^\alpha, n^\alpha, T^\alpha), (U^\beta, V^\beta, n^\beta, T^\beta)$
- Constraints:

$$\begin{cases} U^\alpha + U^\beta = \text{Const.} \\ V^\alpha + V^\beta = \text{Const.} \\ n^\alpha + n^\beta = \text{Const.} \end{cases} \Rightarrow \begin{cases} \delta U^\alpha + \delta U^\beta = 0 \\ \delta V^\alpha + \delta V^\beta = 0 \\ \delta n^\alpha + \delta n^\beta = 0 \end{cases} .$$
- $\delta S^\alpha = \frac{\delta U^\alpha + p^\alpha \delta V^\alpha - \mu^\alpha \delta n^\alpha}{T^\alpha}, \delta S^\beta = \frac{\delta U^\beta + p^\beta \delta V^\beta - \mu^\beta \delta n^\beta}{T^\beta}.$

§3.3 Equilibrium of single-component multi-phase system

- Considering isolated two-phase system, α, β .
- State parameter $(U^\alpha, V^\alpha, n^\alpha, T^\alpha), (U^\beta, V^\beta, n^\beta, T^\beta)$
- Constraints:

$$\begin{cases} U^\alpha + U^\beta = \text{Const.} \\ V^\alpha + V^\beta = \text{Const.} \\ n^\alpha + n^\beta = \text{Const.} \end{cases} \Rightarrow \begin{cases} \delta U^\alpha + \delta U^\beta = 0 \\ \delta V^\alpha + \delta V^\beta = 0 \\ \delta n^\alpha + \delta n^\beta = 0 \end{cases} .$$
- $\delta S^\alpha = \frac{\delta U^\alpha + p^\alpha \delta V^\alpha - \mu^\alpha \delta n^\alpha}{T^\alpha}, \delta S^\beta = \frac{\delta U^\beta + p^\beta \delta V^\beta - \mu^\beta \delta n^\beta}{T^\beta}.$
- $\delta S = \delta S^\alpha + \delta S^\beta =$

$$\delta U^\alpha \left(\frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) + \delta V^\alpha \left(\frac{p^\alpha}{T^\alpha} - \frac{p^\beta}{T^\beta} \right) - \delta n^\alpha \left(\frac{\mu^\alpha}{T^\alpha} - \frac{\mu^\beta}{T^\beta} \right)$$

§3.3 Equilibrium of single-component multi-phase system

- Considering isolated two-phase system, α, β .
- State parameter $(U^\alpha, V^\alpha, n^\alpha, T^\alpha), (U^\beta, V^\beta, n^\beta, T^\beta)$
- Constraints:

$$\begin{cases} U^\alpha + U^\beta = \text{Const.} \\ V^\alpha + V^\beta = \text{Const.} \\ n^\alpha + n^\beta = \text{Const.} \end{cases} \Rightarrow \begin{cases} \delta U^\alpha + \delta U^\beta = 0 \\ \delta V^\alpha + \delta V^\beta = 0 \\ \delta n^\alpha + \delta n^\beta = 0 \end{cases} .$$
- $\delta S^\alpha = \frac{\delta U^\alpha + p^\alpha \delta V^\alpha - \mu^\alpha \delta n^\alpha}{T^\alpha}, \delta S^\beta = \frac{\delta U^\beta + p^\beta \delta V^\beta - \mu^\beta \delta n^\beta}{T^\beta}.$
- $\delta S = \delta S^\alpha + \delta S^\beta =$

$$\delta U^\alpha \left(\frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) + \delta V^\alpha \left(\frac{p^\alpha}{T^\alpha} - \frac{p^\beta}{T^\beta} \right) - \delta n^\alpha \left(\frac{\mu^\alpha}{T^\alpha} - \frac{\mu^\beta}{T^\beta} \right) = 0.$$

§3.3 Equilibrium of single-component multi-phase system

- Considering isolated two-phase system, α, β .
- State parameter $(U^\alpha, V^\alpha, n^\alpha, T^\alpha), (U^\beta, V^\beta, n^\beta, T^\beta)$
- Constraints:

$$\begin{cases} U^\alpha + U^\beta = \text{Const.} \\ V^\alpha + V^\beta = \text{Const.} \\ n^\alpha + n^\beta = \text{Const.} \end{cases} \Rightarrow \begin{cases} \delta U^\alpha + \delta U^\beta = 0 \\ \delta V^\alpha + \delta V^\beta = 0 \\ \delta n^\alpha + \delta n^\beta = 0 \end{cases} .$$
- $\delta S^\alpha = \frac{\delta U^\alpha + p^\alpha \delta V^\alpha - \mu^\alpha \delta n^\alpha}{T^\alpha}, \delta S^\beta = \frac{\delta U^\beta + p^\beta \delta V^\beta - \mu^\beta \delta n^\beta}{T^\beta}.$
- $\delta S = \delta S^\alpha + \delta S^\beta =$
 $\delta U^\alpha \left(\frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) + \delta V^\alpha \left(\frac{p^\alpha}{T^\alpha} - \frac{p^\beta}{T^\beta} \right) - \delta n^\alpha \left(\frac{\mu^\alpha}{T^\alpha} - \frac{\mu^\beta}{T^\beta} \right) = 0.$
- Equilibrium condition: $T^\alpha = T^\beta$ (thermodynamics),
 $p^\alpha = p^\beta$ (mechanics), $\mu^\alpha = \mu^\beta$ (phase).

Equilibrium of single-component multi-phase system

$$\delta S = \delta U^\alpha \left(\frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) + \delta V^\alpha \left(\frac{p^\alpha}{T^\alpha} - \frac{p^\beta}{T^\beta} \right) - \delta n^\alpha \left(\frac{\mu^\alpha}{T^\alpha} - \frac{\mu^\beta}{T^\beta} \right) = 0.$$

- If the equilibrium condition is not satisfied, system proceeds in the direction of increasing the entropy.

Equilibrium of single-component multi-phase system

$$\delta S = \delta U^\alpha \left(\frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) + \delta V^\alpha \left(\frac{p^\alpha}{T^\alpha} - \frac{p^\beta}{T^\beta} \right) - \delta n^\alpha \left(\frac{\mu^\alpha}{T^\alpha} - \frac{\mu^\beta}{T^\beta} \right) = 0.$$

- If the equilibrium condition is not satisfied, system proceeds in the direction of increasing the entropy.
- 1° $T^\alpha \neq T^\beta$, direction: $\delta U^\alpha \left(\frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) > 0$. If $T^\alpha > T^\beta$, direction: $\delta U^\alpha < 0$. Energy $\alpha \rightarrow \beta$, from high T to low T .

Equilibrium of single-component multi-phase system

$$\delta S = \delta U^\alpha \left(\frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) + \delta V^\alpha \left(\frac{p^\alpha}{T^\alpha} - \frac{p^\beta}{T^\beta} \right) - \delta n^\alpha \left(\frac{\mu^\alpha}{T^\alpha} - \frac{\mu^\beta}{T^\beta} \right) = 0.$$

- If the equilibrium condition is not satisfied, system proceeds in the direction of increasing the entropy.
- 1° $T^\alpha \neq T^\beta$, direction: $\delta U^\alpha \left(\frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) > 0$. If $T^\alpha > T^\beta$, direction: $\delta U^\alpha < 0$. Energy $\alpha \rightarrow \beta$, from high T to low T .
- 2° If $p^\alpha > p^\beta$ ($T^\alpha = T^\beta$), direction: $\delta V^\alpha \frac{p^\alpha - p^\beta}{T} > 0$, i.e., $\delta V^\alpha > 0$, phase α (high pressure) expands.

Equilibrium of single-component multi-phase system

$$\delta S = \delta U^\alpha \left(\frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) + \delta V^\alpha \left(\frac{p^\alpha}{T^\alpha} - \frac{p^\beta}{T^\beta} \right) - \delta n^\alpha \left(\frac{\mu^\alpha}{T^\alpha} - \frac{\mu^\beta}{T^\beta} \right) = 0.$$

- If the equilibrium condition is not satisfied, system proceeds in the direction of increasing the entropy.
- 1° $T^\alpha \neq T^\beta$, direction: $\delta U^\alpha \left(\frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) > 0$. If $T^\alpha > T^\beta$, direction: $\delta U^\alpha < 0$. Energy $\alpha \rightarrow \beta$, from high T to low T .
- 2° If $p^\alpha > p^\beta$ ($T^\alpha = T^\beta$), direction: $\delta V^\alpha \frac{p^\alpha - p^\beta}{T} > 0$, i.e., $\delta V^\alpha > 0$, phase α (high pressure) expands.
- 3° If $\mu^\alpha > \mu^\beta$ ($T^\alpha = T^\beta$), direction: $-\delta n^\alpha \frac{\mu^\alpha - \mu^\beta}{T}$, i.e., $\delta n^\alpha < 0$, matter changes phase from α (high μ , chemical potential) to β .

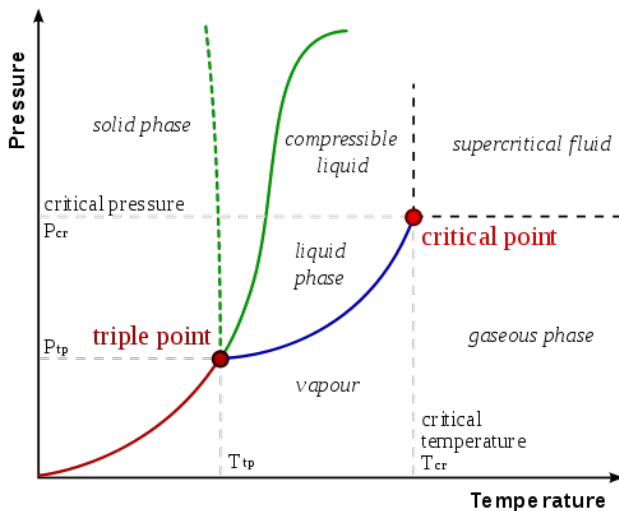
Table of contents

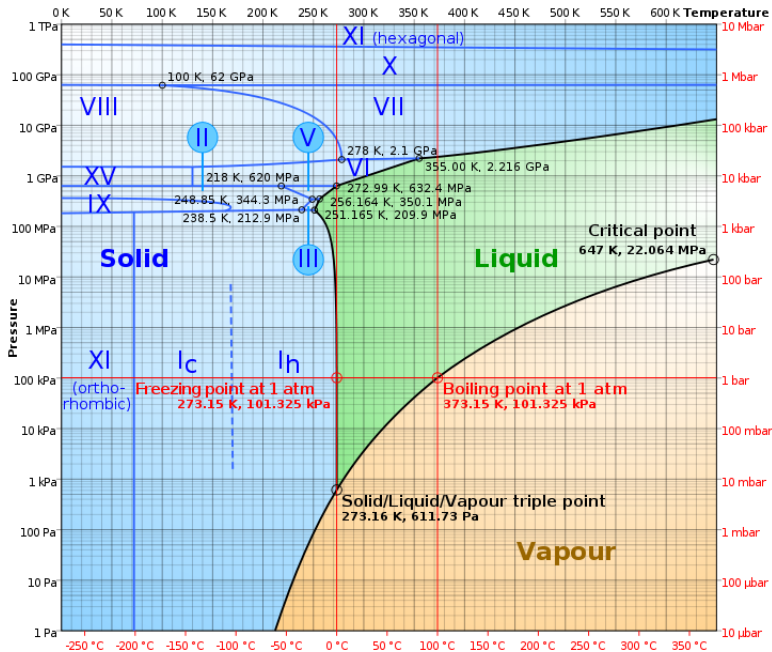
① Phase Transition of Single-Component System

- 3.1 Criterion of thermal equilibrium
- 3.2 Basic equations of open system
- 3.3 Equilibrium of single-component multi-phase system
- 3.4 Properties of equilibrium of s-c multi-phase system
- 3.5 Critical point and phase change between gas and liquid
- 3.7 Classification of the phase transition
- 3.9 Landau's approximation for the continuous phase transition

§3.4 Properties of equilibrium of single-component multi-phase system – Phase diagram of water

§3.4 Properties of equilibrium of single-component multi-phase system – Phase diagram of water





Properties of equilibrium of s-c multi-phase system

- For three phases coexistence:

Properties of equilibrium of s-c multi-phase system

- For three phases coexistence:

$$T^{\alpha} = T^{\beta} = T^{\gamma} = T,$$

Properties of equilibrium of s-c multi-phase system

- For three phases coexistence:

$$T^{\alpha} = T^{\beta} = T^{\gamma} = T,$$

$$p^{\alpha} = p^{\beta} = p^{\gamma} = p,$$

Properties of equilibrium of s-c multi-phase system

- For three phases coexistence:

$$T^{\alpha} = T^{\beta} = T^{\gamma} = T,$$

$$p^{\alpha} = p^{\beta} = p^{\gamma} = p,$$

$$\mu^{\alpha}(T, p) = \mu^{\beta}(T, p) = \mu^{\gamma}(T, p)$$

(Triple point).

Properties of equilibrium of s-c multi-phase system

- For three phases coexistence:

$$T^{\alpha} = T^{\beta} = T^{\gamma} = T,$$

$$p^{\alpha} = p^{\beta} = p^{\gamma} = p,$$

$$\mu^{\alpha}(T, p) = \mu^{\beta}(T, p) = \mu^{\gamma}(T, p)$$

(Triple point).

- For two phases coexistence curve:

Properties of equilibrium of s-c multi-phase system

- For three phases coexistence:

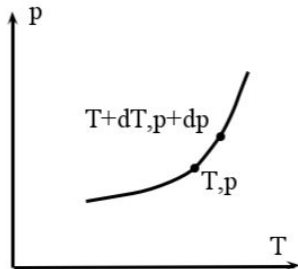
$$T^{\alpha} = T^{\beta} = T^{\gamma} = T,$$

$$p^{\alpha} = p^{\beta} = p^{\gamma} = p,$$

$$\mu^{\alpha}(T, p) = \mu^{\beta}(T, p) = \mu^{\gamma}(T, p)$$

(Triple point).

- For two phases coexistence curve:



Properties of equilibrium of s-c multi-phase system

- For three phases coexistence:

$$T^{\alpha} = T^{\beta} = T^{\gamma} = T,$$

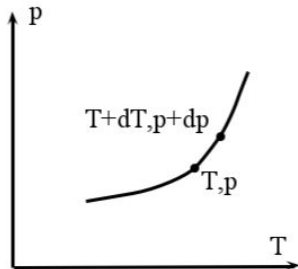
$$p^{\alpha} = p^{\beta} = p^{\gamma} = p,$$

$$\mu^{\alpha}(T, p) = \mu^{\beta}(T, p) = \mu^{\gamma}(T, p)$$

(Triple point).

- For two phases coexistence curve:

- $\mu^{\alpha}(T, p) = \mu^{\beta}(T, p),$



Properties of equilibrium of s-c multi-phase system

- For three phases coexistence:

$$T^{\alpha} = T^{\beta} = T^{\gamma} = T,$$

$$p^{\alpha} = p^{\beta} = p^{\gamma} = p,$$

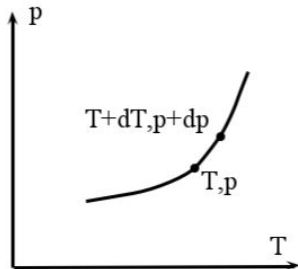
$$\mu^{\alpha}(T, p) = \mu^{\beta}(T, p) = \mu^{\gamma}(T, p)$$

(Triple point).

- For two phases coexistence curve:

- $\mu^{\alpha}(T, p) = \mu^{\beta}(T, p),$

$$\Rightarrow d\mu^{\alpha} = d\mu^{\beta}.$$



Properties of equilibrium of s-c multi-phase system

- For three phases coexistence:

$$T^{\alpha} = T^{\beta} = T^{\gamma} = T,$$

$$p^{\alpha} = p^{\beta} = p^{\gamma} = p,$$

$$\mu^{\alpha}(T, p) = \mu^{\beta}(T, p) = \mu^{\gamma}(T, p)$$

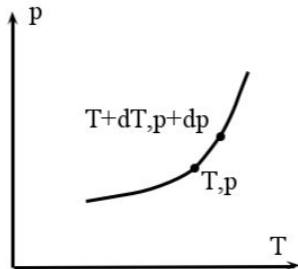
(Triple point).

- For two phases coexistence curve:

- $\mu^{\alpha}(T, p) = \mu^{\beta}(T, p),$

$$\Rightarrow d\mu^{\alpha} = d\mu^{\beta}.$$

- $\because d\mu = \frac{V}{n}dp - \frac{S}{n}dT = V_m dp - S_m dT,$



Properties of equilibrium of s-c multi-phase system

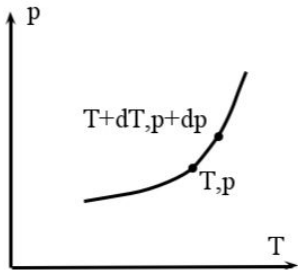
- For three phases coexistence:

$$T^{\alpha} = T^{\beta} = T^{\gamma} = T,$$

$$p^{\alpha} = p^{\beta} = p^{\gamma} = p,$$

$$\mu^{\alpha}(T, p) = \mu^{\beta}(T, p) = \mu^{\gamma}(T, p)$$

(Triple point).



- For two phases coexistence curve:

- $\mu^{\alpha}(T, p) = \mu^{\beta}(T, p),$

$$\Rightarrow d\mu^{\alpha} = d\mu^{\beta}.$$

- $\because d\mu = \frac{V}{n}dp - \frac{S}{n}dT = V_m dp - S_m dT,$

- $\therefore V_m^{\alpha} dp - S_m^{\alpha} dT = V_m^{\beta} dp - S_m^{\beta} dT$

Properties of equilibrium of s-c multi-phase system

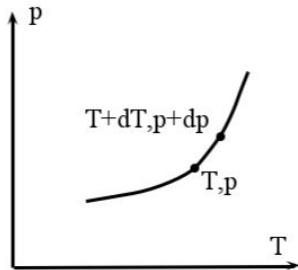
- For three phases coexistence:

$$T^{\alpha} = T^{\beta} = T^{\gamma} = T,$$

$$p^{\alpha} = p^{\beta} = p^{\gamma} = p,$$

$$\mu^{\alpha}(T, p) = \mu^{\beta}(T, p) = \mu^{\gamma}(T, p)$$

(Triple point).



- For two phases coexistence curve:

$$\mu^{\alpha}(T, p) = \mu^{\beta}(T, p),$$

$$\Rightarrow d\mu^{\alpha} = d\mu^{\beta}.$$

$$\because d\mu = \frac{V}{n}dp - \frac{S}{n}dT = V_m dp - S_m dT,$$

$$\therefore V_m^{\alpha} dp - S_m^{\alpha} dT = V_m^{\beta} dp - S_m^{\beta} dT \Rightarrow \frac{dp}{dT} = \frac{S_m^{\beta} - S_m^{\alpha}}{V_m^{\beta} - V_m^{\alpha}}.$$

Properties of equilibrium of s-c multi-phase system

$$\frac{dp}{dT} = \frac{S_m^\beta - S_m^\alpha}{V_m^\beta - V_m^\alpha}$$

Properties of equilibrium of s-c multi-phase system

$$\frac{dp}{dT} = \frac{S_m^\beta - S_m^\alpha}{V_m^\beta - V_m^\alpha}$$

- During the phase transition (notice, reversible),

Properties of equilibrium of s-c multi-phase system

$$\frac{dp}{dT} = \frac{S_m^\beta - S_m^\alpha}{V_m^\beta - V_m^\alpha}$$

- During the phase transition (notice, reversible),
 $dS = \delta Q/T \Rightarrow \Delta Q = T\Delta S.$

Properties of equilibrium of s-c multi-phase system

$$\frac{dp}{dT} = \frac{S_m^\beta - S_m^\alpha}{V_m^\beta - V_m^\alpha}$$

- During the phase transition (notice, reversible),
 $dS = \delta Q/T \Rightarrow \Delta Q = T\Delta S$.
- Define L: Latent heat of phase change of 1 mol matter, then $L = T(S_m^\beta - S_m^\alpha)$.

Properties of equilibrium of s-c multi-phase system

$$\frac{dp}{dT} = \frac{S_m^\beta - S_m^\alpha}{V_m^\beta - V_m^\alpha}$$

- During the phase transition (notice, reversible),
 $dS = \delta Q/T \Rightarrow \Delta Q = T\Delta S$.
- Define L: Latent heat of phase change of 1 mol matter, then $L = T(S_m^\beta - S_m^\alpha)$.
- $\therefore \frac{dp}{dT} = \frac{L}{T(V_m^\beta - V_m^\alpha)}$ called Clapeyron's equation.

Properties of equilibrium of s-c multi-phase system

$$\frac{dp}{dT} = \frac{S_m^\beta - S_m^\alpha}{V_m^\beta - V_m^\alpha}$$

- During the phase transition (notice, reversible),
 $dS = \delta Q/T \Rightarrow \Delta Q = T\Delta S$.
- Define L: Latent heat of phase change of 1 mol matter, then $L = T(S_m^\beta - S_m^\alpha)$.
- $\therefore \frac{dp}{dT} = \frac{L}{T(V_m^\beta - V_m^\alpha)}$ called Clapeyron's equation.
- The slope of phase-boundary curve is available in theory then (phase-boundary curve comes from experiment).

Table of contents

1 Phase Transition of Single-Component System

- 3.1 Criterion of thermal equilibrium
- 3.2 Basic equations of open system
- 3.3 Equilibrium of single-component multi-phase system
- 3.4 Properties of equilibrium of s-c multi-phase system
- 3.5 Critical point and phase change between gas and liquid
- 3.7 Classification of the phase transition
- 3.9 Landau's approximation for the continuous phase transition

§3.5 Critical point and phase change between gas and liquid

§3.5 Critical point and phase change between gas and liquid

- There exists a critical point between gas and liquid. Beyond that point, gas and liquid have no difference.

§3.5 Critical point and phase change between gas and liquid

- There exists a critical point between gas and liquid. Beyond that point, gas and liquid have no difference.
- A critical point for solid and liquid?

§3.5 Critical point and phase change between gas and liquid

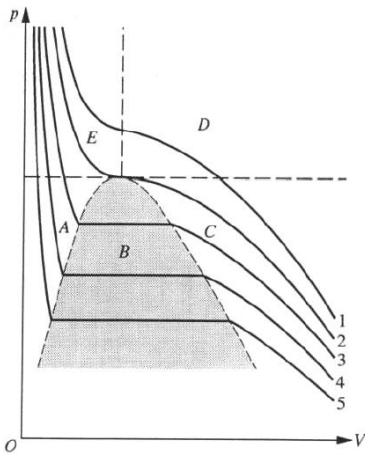
- There exists a critical point between gas and liquid. Beyond that point, gas and liquid have no difference.
- A critical point for solid and liquid?
- A critical point for solid and gas?

§3.5 Critical point and phase change between gas and liquid

- There exists a critical point between gas and liquid. Beyond that point, gas and liquid have no difference.
- A critical point for solid and liquid?
- A critical point for solid and gas?
- P-V diagram for liquid and gas:

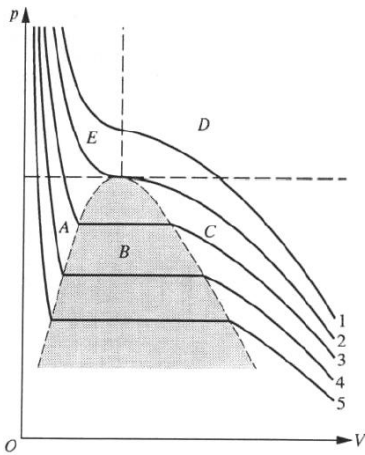
§3.5 Critical point and phase change between gas and liquid

- There exists a critical point between gas and liquid. Beyond that point, gas and liquid have no difference.
- A critical point for solid and liquid?
- A critical point for solid and gas?
- P-V diagram for liquid and gas:



§3.5 Critical point and phase change between gas and liquid

- There exists a critical point between gas and liquid. Beyond that point, gas and liquid have no difference.
- A critical point for solid and liquid?
- A critical point for solid and gas?
- P-V diagram for liquid and gas:



Stability condition: $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0.$

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- Consider two-phase coexistence, but very close to the critical point.

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- Consider two-phase coexistence, but very close to the critical point.
- Volume of liquid V_m , gas $V_m + \delta V_m$ (close to each other).

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- Consider two-phase coexistence, but very close to the critical point.
- Volume of liquid V_m , gas $V_m + \delta V_m$ (close to each other).
- Pressure: $p(V_m + \delta V_m, T) = p(V_m, T)$,

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- Consider two-phase coexistence, but very close to the critical point.
- Volume of liquid V_m , gas $V_m + \delta V_m$ (close to each other).
- Pressure: $p(V_m + \delta V_m, T) = p(V_m, T)$,
while $p(V_m + \delta V_m, T) =$
 $p(V_m, T) + \left(\frac{\partial p}{\partial V_m}\right)_T \delta V_m + \frac{1}{2} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^2,$

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- Consider two-phase coexistence, but very close to the critical point.
- Volume of liquid V_m , gas $V_m + \delta V_m$ (close to each other).
- Pressure: $p(V_m + \delta V_m, T) = p(V_m, T)$,
while $p(V_m + \delta V_m, T) =$
 $p(V_m, T) + \left(\frac{\partial p}{\partial V_m}\right)_T \delta V_m + \frac{1}{2} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^2,$
- $\Rightarrow \left(\frac{\partial p}{\partial V_m}\right)_T + \frac{1}{2} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T \delta V_m = 0$.

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- Consider two-phase coexistence, but very close to the critical point.
- Volume of liquid V_m , gas $V_m + \delta V_m$ (close to each other).
- Pressure: $p(V_m + \delta V_m, T) = p(V_m, T)$,
while $p(V_m + \delta V_m, T) =$
 $p(V_m, T) + \left(\frac{\partial p}{\partial V_m}\right)_T \delta V_m + \frac{1}{2} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^2,$
- $\Rightarrow \left(\frac{\partial p}{\partial V_m}\right)_T + \frac{1}{2} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T \delta V_m = 0$.
- When $T \rightarrow T_c$, $\delta V_m \rightarrow 0$,

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- Consider two-phase coexistence, but very close to the critical point.
- Volume of liquid V_m , gas $V_m + \delta V_m$ (close to each other).
- Pressure: $p(V_m + \delta V_m, T) = p(V_m, T)$,
while $p(V_m + \delta V_m, T) =$
 $p(V_m, T) + \left(\frac{\partial p}{\partial V_m}\right)_T \delta V_m + \frac{1}{2} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^2,$
- $\Rightarrow \left(\frac{\partial p}{\partial V_m}\right)_T + \frac{1}{2} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T \delta V_m = 0.$
- When $T \rightarrow T_c$, $\delta V_m \rightarrow 0$,
 \therefore in the critical point: $\left(\frac{\partial p}{\partial V_m}\right)_T = 0.$

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- $\Delta S_m = \delta S_m + \frac{1}{2}\delta^2 S_m + \frac{1}{3!}\delta^3 S_m + \frac{1}{4!}\delta^4 S_m + \dots,$

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- $\Delta S_m = \delta S_m + \frac{1}{2}\delta^2 S_m + \frac{1}{3!}\delta^3 S_m + \frac{1}{4!}\delta^4 S_m + \dots$,
where $\delta^2 S_m = -\frac{C_{V,m}}{T^2}(\delta T)^2 + \frac{1}{T}\left(\frac{\partial p}{\partial V_m}\right)_T(\delta V_m)^2$
(3.1.13);

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- $\Delta S_m = \delta S_m + \frac{1}{2}\delta^2 S_m + \frac{1}{3!}\delta^3 S_m + \frac{1}{4!}\delta^4 S_m + \dots,$

where $\delta^2 S_m = -\frac{C_{V,m}}{T^2}(\delta T)^2 + \frac{1}{T}\left(\frac{\partial p}{\partial V_m}\right)_T (\delta V_m)^2$
(3.1.13);

for phase coexistence, $\delta T = 0$, and also $\left(\frac{\partial p}{\partial V_m}\right)_T = 0$,

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- $\Delta S_m = \delta S_m + \frac{1}{2}\delta^2 S_m + \frac{1}{3!}\delta^3 S_m + \frac{1}{4!}\delta^4 S_m + \dots,$

where $\delta^2 S_m = -\frac{C_{V,m}}{T^2}(\delta T)^2 + \frac{1}{T}\left(\frac{\partial p}{\partial V_m}\right)_T(\delta V_m)^2$
(3.1.13);

for phase coexistence, $\delta T = 0$, and also $\left(\frac{\partial p}{\partial V_m}\right)_T = 0$,
 $\Rightarrow \delta^2 S_m = 0$.

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

- $\Delta S_m = \delta S_m + \frac{1}{2}\delta^2 S_m + \frac{1}{3!}\delta^3 S_m + \frac{1}{4!}\delta^4 S_m + \dots,$

where $\delta^2 S_m = -\frac{C_{V,m}}{T^2}(\delta T)^2 + \frac{1}{T}\left(\frac{\partial p}{\partial V_m}\right)_T (\delta V_m)^2$
(3.1.13);

for phase coexistence, $\delta T = 0$, and also $\left(\frac{\partial p}{\partial V_m}\right)_T = 0$,
 $\Rightarrow \delta^2 S_m = 0$.

- To get stability, one needs $\delta^3 S_m < 0$, or $\delta^4 S_m < 0$ if $\delta^3 S_m = 0$, or ...

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\bullet \Delta S_m = \delta S_m + \frac{1}{2}\delta^2 S_m + \frac{1}{3!}\delta^3 S_m + \frac{1}{4!}\delta^4 S_m + \dots,$$

$$\text{where } \delta^2 S_m = -\frac{C_{V,m}}{T^2}(\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V_m}\right)_T (\delta V_m)^2$$

$$(3.1.13);$$

for phase coexistence, $\delta T = 0$, and also $\left(\frac{\partial p}{\partial V_m}\right)_T = 0$,
 $\Rightarrow \delta^2 S_m = 0$.

- To get stability, one needs $\delta^3 S_m < 0$, or $\delta^4 S_m < 0$ if $\delta^3 S_m = 0$, or ...
- $\delta^3 S_m = \delta(\delta^2 S_m)$

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\bullet \Delta S_m = \delta S_m + \frac{1}{2}\delta^2 S_m + \frac{1}{3!}\delta^3 S_m + \frac{1}{4!}\delta^4 S_m + \dots,$$

$$\text{where } \delta^2 S_m = -\frac{C_{V,m}}{T^2}(\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V_m}\right)_T (\delta V_m)^2$$

(3.1.13);

for phase coexistence, $\delta T = 0$, and also $\left(\frac{\partial p}{\partial V_m}\right)_T = 0$,
 $\Rightarrow \delta^2 S_m = 0$.

• To get stability, one needs $\delta^3 S_m < 0$, or $\delta^4 S_m < 0$ if $\delta^3 S_m = 0$, or ...

$$\bullet \delta^3 S_m = \delta(\delta^2 S_m) = \frac{\partial}{\partial T}(\delta^2 S_m)\delta T + \frac{\partial}{\partial V_m}(\delta^2 S_m)\delta V_m$$

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\bullet \Delta S_m = \delta S_m + \frac{1}{2}\delta^2 S_m + \frac{1}{3!}\delta^3 S_m + \frac{1}{4!}\delta^4 S_m + \dots,$$

$$\text{where } \delta^2 S_m = -\frac{C_{V,m}}{T^2}(\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V_m}\right)_T (\delta V_m)^2$$

$$(3.1.13);$$

for phase coexistence, $\delta T = 0$, and also $\left(\frac{\partial p}{\partial V_m}\right)_T = 0$,
 $\Rightarrow \delta^2 S_m = 0$.

- To get stability, one needs $\delta^3 S_m < 0$, or $\delta^4 S_m < 0$ if $\delta^3 S_m = 0$, or ...

$$\bullet \delta^3 S_m = \delta(\delta^2 S_m) = \frac{\partial}{\partial T}(\delta^2 S_m)\delta T + \frac{\partial}{\partial V_m}(\delta^2 S_m)\delta V_m$$

$$= \frac{\partial}{\partial V_m}(\delta^2 S_m)\delta V_m$$

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\bullet \Delta S_m = \delta S_m + \frac{1}{2}\delta^2 S_m + \frac{1}{3!}\delta^3 S_m + \frac{1}{4!}\delta^4 S_m + \dots,$$

$$\text{where } \delta^2 S_m = -\frac{C_{V,m}}{T^2}(\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V_m}\right)_T (\delta V_m)^2$$

$$(3.1.13);$$

for phase coexistence, $\delta T = 0$, and also $\left(\frac{\partial p}{\partial V_m}\right)_T = 0$,
 $\Rightarrow \delta^2 S_m = 0$.

- To get stability, one needs $\delta^3 S_m < 0$, or $\delta^4 S_m < 0$ if $\delta^3 S_m = 0$, or ...

$$\bullet \delta^3 S_m = \delta(\delta^2 S_m) = \frac{\partial}{\partial T}(\delta^2 S_m)\delta T + \frac{\partial}{\partial V_m}(\delta^2 S_m)\delta V_m$$

$$= \frac{\partial}{\partial V_m}(\delta^2 S_m)\delta V_m = \frac{\partial}{\partial V_m}\left(\frac{1}{T} \left(\frac{\partial p}{\partial V_m}\right)_T (\delta V_m)^2\right)\delta V_m$$

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\bullet \Delta S_m = \delta S_m + \frac{1}{2}\delta^2 S_m + \frac{1}{3!}\delta^3 S_m + \frac{1}{4!}\delta^4 S_m + \dots,$$

$$\text{where } \delta^2 S_m = -\frac{C_{V,m}}{T^2}(\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V_m}\right)_T (\delta V_m)^2$$

$$(3.1.13);$$

for phase coexistence, $\delta T = 0$, and also $\left(\frac{\partial p}{\partial V_m}\right)_T = 0$,
 $\Rightarrow \delta^2 S_m = 0$.

- To get stability, one needs $\delta^3 S_m < 0$, or $\delta^4 S_m < 0$ if $\delta^3 S_m = 0$, or ...

$$\bullet \delta^3 S_m = \delta(\delta^2 S_m) = \frac{\partial}{\partial T}(\delta^2 S_m)\delta T + \frac{\partial}{\partial V_m}(\delta^2 S_m)\delta V_m$$

$$= \frac{\partial}{\partial V_m}(\delta^2 S_m)\delta V_m = \frac{\partial}{\partial V_m}\left(\frac{1}{T} \left(\frac{\partial p}{\partial V_m}\right)_T (\delta V_m)^2\right)\delta V_m$$

$$= \frac{1}{T} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^3.$$

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\delta^3 S_m = \frac{1}{T} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^3$$

- If $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T \neq 0$, it means the stability is determined by the virtual variation $\delta V_m > 0$ or < 0 , which is unreasonable.

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\delta^3 S_m = \frac{1}{T} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^3$$

- If $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T \neq 0$, it means the stability is determined by the virtual variation $\delta V_m > 0$ or < 0 , which is unreasonable.

Therefore, $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0$.

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\delta^3 S_m = \frac{1}{T} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^3$$

- If $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T \neq 0$, it means the stability is determined by the virtual variation $\delta V_m > 0$ or < 0 , which is unreasonable.

Therefore, $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0$.

- I.e., $\delta^3 S_m = 0$.

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\delta^3 S_m = \frac{1}{T} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^3$$

- If $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T \neq 0$, it means the stability is determined by the virtual variation $\delta V_m > 0$ or < 0 , which is unreasonable.

Therefore, $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0$.

- I.e., $\delta^3 S_m = 0$.
- $\delta^4 S_m = \delta(\delta^3 S_m)$

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\delta^3 S_m = \frac{1}{T} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^3$$

- If $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T \neq 0$, it means the stability is determined by the virtual variation $\delta V_m > 0$ or < 0 , which is unreasonable.

Therefore, $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0$.

- I.e., $\delta^3 S_m = 0$.
- $\delta^4 S_m = \delta(\delta^3 S_m) = \frac{1}{T} \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T (\delta V_m)^4$

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\delta^3 S_m = \frac{1}{T} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^3$$

- If $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T \neq 0$, it means the stability is determined by the virtual variation $\delta V_m > 0$ or < 0 , which is unreasonable.

Therefore, $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0$.

- I.e., $\delta^3 S_m = 0$.
- $\delta^4 S_m = \delta(\delta^3 S_m) = \frac{1}{T} \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T (\delta V_m)^4 < 0$ (required),

Provement of $\left(\frac{\partial p}{\partial V_m}\right)_T = 0, \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0, \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

$$\delta^3 S_m = \frac{1}{T} \left(\frac{\partial^2 p}{\partial V_m^2}\right)_T (\delta V_m)^3$$

- If $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T \neq 0$, it means the stability is determined by the virtual variation $\delta V_m > 0$ or < 0 , which is unreasonable.

Therefore, $\left(\frac{\partial^2 p}{\partial V_m^2}\right)_T = 0$.

- I.e., $\delta^3 S_m = 0$.
- $\delta^4 S_m = \delta(\delta^3 S_m) = \frac{1}{T} \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T (\delta V_m)^4 < 0$ (required),
- $\therefore \left(\frac{\partial^3 p}{\partial V_m^3}\right)_T < 0$.

Table of contents

1 Phase Transition of Single-Component System

- 3.1 Criterion of thermal equilibrium
- 3.2 Basic equations of open system
- 3.3 Equilibrium of single-component multi-phase system
- 3.4 Properties of equilibrium of s-c multi-phase system
- 3.5 Critical point and phase change between gas and liquid
- 3.7 Classification of the phase transition
- 3.9 Landau's approximation for the continuous phase transition

§3.7 Classification of the phase transition

§3.7 Classification of the phase transition

- Properties of normal phase transition: latent heat $L = T(S_m^\beta - S_m^\alpha)$ (entropy jump), volume jump.

§3.7 Classification of the phase transition

- Properties of normal phase transition: latent heat $L = T(S_m^\beta - S_m^\alpha)$ (entropy jump), volume jump.
- $d\mu = -S_m dT + V_m dp$

§3.7 Classification of the phase transition

- Properties of normal phase transition: latent heat $L = T(S_m^\beta - S_m^\alpha)$ (entropy jump), volume jump.
- $d\mu = -S_m dT + V_m dp$
 $\Rightarrow S_m = -\left(\frac{\partial \mu}{\partial T}\right)_p, V_m = \left(\frac{\partial \mu}{\partial p}\right)_T.$

§3.7 Classification of the phase transition

- Properties of normal phase transition: latent heat $L = T(S_m^\beta - S_m^\alpha)$ (entropy jump), volume jump.
- $d\mu = -S_m dT + V_m dp$
 $\Rightarrow S_m = -\left(\frac{\partial \mu}{\partial T}\right)_p, V_m = \left(\frac{\partial \mu}{\partial p}\right)_T$.
1st order of phase transition: 1st order of partial differential of μ not continuous.

§3.7 Classification of the phase transition

- Properties of normal phase transition: latent heat $L = T(S_m^\beta - S_m^\alpha)$ (entropy jump), volume jump.
- $d\mu = -S_m dT + V_m dp$
 $\Rightarrow S_m = -\left(\frac{\partial \mu}{\partial T}\right)_p, V_m = \left(\frac{\partial \mu}{\partial p}\right)_T$.
1st order of phase transition: 1st order of partial differential of μ not continuous.
- 2nd order: $c_p = T\left(\frac{\partial s}{\partial T}\right)_p = -T\frac{\partial^2 \mu}{\partial T^2},$

§3.7 Classification of the phase transition

- Properties of normal phase transition: latent heat $L = T(S_m^\beta - S_m^\alpha)$ (entropy jump), volume jump.
- $d\mu = -S_m dT + V_m dp$
 $\Rightarrow S_m = -\left(\frac{\partial\mu}{\partial T}\right)_p, V_m = \left(\frac{\partial\mu}{\partial p}\right)_T$.
1st order of phase transition: 1st order of partial differential of μ not continuous.
- 2nd order: $c_p = T\left(\frac{\partial s}{\partial T}\right)_p = -T\frac{\partial^2\mu}{\partial T^2}$,
 $\alpha = \frac{1}{v}\left(\frac{\partial v}{\partial T}\right)_p = \frac{1}{v}\frac{\partial^2\mu}{\partial T\partial p}$,

§3.7 Classification of the phase transition

- Properties of normal phase transition: latent heat $L = T(S_m^\beta - S_m^\alpha)$ (entropy jump), volume jump.
- $d\mu = -S_m dT + V_m dp$
 $\Rightarrow S_m = -\left(\frac{\partial\mu}{\partial T}\right)_p, V_m = \left(\frac{\partial\mu}{\partial p}\right)_T$.
 1st order of phase transition: 1st order of partial differential of μ not continuous.
- 2nd order: $c_p = T\left(\frac{\partial s}{\partial T}\right)_p = -T\frac{\partial^2\mu}{\partial T^2}$,
 $\alpha = \frac{1}{v}\left(\frac{\partial v}{\partial T}\right)_p = \frac{1}{v}\frac{\partial^2\mu}{\partial T\partial p}, \kappa_T = -\frac{1}{v}\left(\frac{\partial v}{\partial p}\right)_T = -\frac{1}{v}\frac{\partial^2\mu}{\partial p^2}$

§3.7 Classification of the phase transition

- Properties of normal phase transition: latent heat $L = T(S_m^\beta - S_m^\alpha)$ (entropy jump), volume jump.

- $d\mu = -S_m dT + V_m dp$
 $\Rightarrow S_m = -\left(\frac{\partial\mu}{\partial T}\right)_p, V_m = \left(\frac{\partial\mu}{\partial p}\right)_T.$

1st order of phase transition: 1st order of partial differential of μ not continuous.

- 2nd order: $c_p = T\left(\frac{\partial s}{\partial T}\right)_p = -T\frac{\partial^2\mu}{\partial T^2},$

$$\alpha = \frac{1}{v}\left(\frac{\partial v}{\partial T}\right)_p = \frac{1}{v}\frac{\partial^2\mu}{\partial T\partial p}, \quad \kappa_T = -\frac{1}{v}\left(\frac{\partial v}{\partial p}\right)_T = -\frac{1}{v}\frac{\partial^2\mu}{\partial p^2}$$

not continuous.

§3.7 Classification of the phase transition

- Properties of normal phase transition: latent heat $L = T(S_m^\beta - S_m^\alpha)$ (entropy jump), volume jump.
- $d\mu = -S_m dT + V_m dp$
 $\Rightarrow S_m = -\left(\frac{\partial\mu}{\partial T}\right)_p, V_m = \left(\frac{\partial\mu}{\partial p}\right)_T$.
1st order of phase transition: 1st order of partial differential of μ not continuous.
- 2nd order: $c_p = T\left(\frac{\partial s}{\partial T}\right)_p = -T\frac{\partial^2\mu}{\partial T^2}$,
 $\alpha = \frac{1}{v}\left(\frac{\partial v}{\partial T}\right)_p = \frac{1}{v}\frac{\partial^2\mu}{\partial T\partial p}, \kappa_T = -\frac{1}{v}\left(\frac{\partial v}{\partial p}\right)_T = -\frac{1}{v}\frac{\partial^2\mu}{\partial p^2}$
not continuous.
- 3rd order, 4th order ...

§3.7 Classification of the phase transition

- Properties of normal phase transition: latent heat $L = T(S_m^\beta - S_m^\alpha)$ (entropy jump), volume jump.
- $d\mu = -S_m dT + V_m dp$
 $\Rightarrow S_m = -\left(\frac{\partial\mu}{\partial T}\right)_p, V_m = \left(\frac{\partial\mu}{\partial p}\right)_T$.
1st order of phase transition: 1st order of partial differential of μ not continuous.
- 2nd order: $c_p = T\left(\frac{\partial s}{\partial T}\right)_p = -T\frac{\partial^2\mu}{\partial T^2}$,
 $\alpha = \frac{1}{v}\left(\frac{\partial v}{\partial T}\right)_p = \frac{1}{v}\frac{\partial^2\mu}{\partial T\partial p}, \kappa_T = -\frac{1}{v}\left(\frac{\partial v}{\partial p}\right)_T = -\frac{1}{v}\frac{\partial^2\mu}{\partial p^2}$
not continuous.
- 3rd order, 4th order ...
- Non-1st order – continuous phase transition.

Table of contents

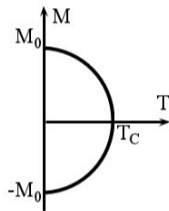
① Phase Transition of Single-Component System

- 3.1 Criterion of thermal equilibrium
- 3.2 Basic equations of open system
- 3.3 Equilibrium of single-component multi-phase system
- 3.4 Properties of equilibrium of s-c multi-phase system
- 3.5 Critical point and phase change between gas and liquid
- 3.7 Classification of the phase transition
- 3.9 Landau's approximation for the continuous phase transition

§3.9 Landau's approximation for the continuous phase transition

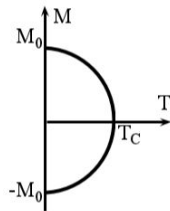
§3.9 Landau's approximation for the continuous phase transition

- Ferromagnetic \leftrightarrow paramagnetic



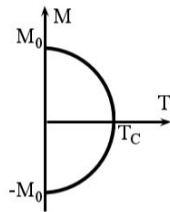
§3.9 Landau's approximation for the continuous phase transition

- Ferromagnetic \leftrightarrow paramagnetic
- Introduce order parameter: M for the spontaneous magnetization.



§3.9 Landau's approximation for the continuous phase transition

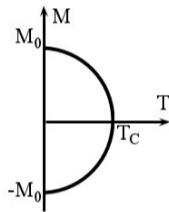
- Ferromagnetic \leftrightarrow paramagnetic
- Introduce order parameter: M for the spontaneous magnetization.
- Expand free energy F near T_c :



§3.9 Landau's approximation for the continuous phase transition

- Ferromagnetic \leftrightarrow paramagnetic
- Introduce order parameter: M for the spontaneous magnetization.
- Expand free energy F near T_c :

$$F(T, M) = F_0(T) + \frac{1}{2}a(T)M^2 + \frac{1}{4}b(T)M^4 + \dots$$

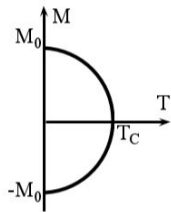


§3.9 Landau's approximation for the continuous phase transition

- Ferromagnetic \leftrightarrow paramagnetic
- Introduce order parameter: M for the spontaneous magnetization.
- Expand free energy F near T_c :

$$F(T, M) = F_0(T) + \frac{1}{2}a(T)M^2 + \frac{1}{4}b(T)M^4 + \dots$$

odd order of M does not exist as $F(M) = F(-M)$.



§3.9 Landau's approximation for the continuous phase transition

- Ferromagnetic \leftrightarrow paramagnetic
- Introduce order parameter: M for the spontaneous magnetization.
- Expand free energy F near T_c :

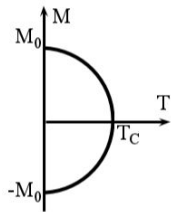
$$F(T, M) = F_0(T) + \frac{1}{2}a(T)M^2 + \frac{1}{4}b(T)M^4 + \dots$$

odd order of M does not exist as $F(M) = F(-M)$.

- For isothermal and isochoric process, criterion for stable equilibrium: $\delta F = 0$, $\delta^2 F > 0$:

$$\frac{\partial F}{\partial M} = M(a + bM^2) = 0,$$

$$\frac{\partial^2 F}{\partial M^2} = a + 3bM^2 > 0.$$



Landau's approximation for the continuous phase transition

- $\frac{\partial F}{\partial M} = M(a + bM^2) = 0$ (1), $\frac{\partial^2 F}{\partial M^2} = a + 3bM^2 > 0$ (2)

Landau's approximation for the continuous phase transition

- $\frac{\partial F}{\partial M} = M(a + bM^2) = 0$ (1), $\frac{\partial^2 F}{\partial M^2} = a + 3bM^2 > 0$ (2)
- (1) $\rightarrow M = 0$, or $M = \pm\sqrt{-a/b}$.

Landau's approximation for the continuous phase transition

- $\frac{\partial F}{\partial M} = M(a + bM^2) = 0$ (1), $\frac{\partial^2 F}{\partial M^2} = a + 3bM^2 > 0$ (2)
- (1) $\rightarrow M = 0$, or $M = \pm\sqrt{-a/b}$.
- From the figure, $T \rightarrow T_c^-$, $M \rightarrow 0$, i.e., $\sqrt{-a/b} \rightarrow 0$,
 $\therefore a(T \rightarrow T_c^-) \rightarrow 0$.

Landau's approximation for the continuous phase transition

- $\frac{\partial F}{\partial M} = M(a + bM^2) = 0$ (1), $\frac{\partial^2 F}{\partial M^2} = a + 3bM^2 > 0$ (2)
- (1) $\rightarrow M = 0$, or $M = \pm\sqrt{-a/b}$.
- From the figure, $T \rightarrow T_c^-$, $M \rightarrow 0$, i.e., $\sqrt{-a/b} \rightarrow 0$,
 $\therefore a(T \rightarrow T_c^-) \rightarrow 0$.
Simply define $a = a_0 \frac{T - T_c}{T_c}$, $b = \text{Const.}$

Landau's approximation for the continuous phase transition

- $\frac{\partial F}{\partial M} = M(a + bM^2) = 0$ (1), $\frac{\partial^2 F}{\partial M^2} = a + 3bM^2 > 0$ (2)
- (1) $\rightarrow M = 0$, or $M = \pm\sqrt{-a/b}$.
- From the figure, $T \rightarrow T_c^-$, $M \rightarrow 0$, i.e., $\sqrt{-a/b} \rightarrow 0$,
 $\therefore a(T \rightarrow T_c^-) \rightarrow 0$.
Simply define $a = a_0 \frac{T - T_c}{T_c}$, $b = \text{Const.}$
- Put $M = \pm\sqrt{-a/b}$ into (2), $\therefore a < 0$ ($T < T_c$),

Landau's approximation for the continuous phase transition

- $\frac{\partial F}{\partial M} = M(a + bM^2) = 0$ (1), $\frac{\partial^2 F}{\partial M^2} = a + 3bM^2 > 0$ (2)
- (1) $\rightarrow M = 0$, or $M = \pm\sqrt{-a/b}$.
- From the figure, $T \rightarrow T_c^-$, $M \rightarrow 0$, i.e., $\sqrt{-a/b} \rightarrow 0$,
 $\therefore a(T \rightarrow T_c^-) \rightarrow 0$.
Simply define $a = a_0 \frac{T - T_c}{T_c}$, $b = \text{Const.}$
- Put $M = \pm\sqrt{-a/b}$ into (2), $\therefore a < 0$ ($T < T_c$),
 $\Rightarrow b > 0$.

Landau's approximation for the continuous phase transition

- $\frac{\partial F}{\partial M} = M(a + bM^2) = 0$ (1), $\frac{\partial^2 F}{\partial M^2} = a + 3bM^2 > 0$ (2)
- (1) $\rightarrow M = 0$, or $M = \pm \sqrt{-a/b}$.
- From the figure, $T \rightarrow T_c^-$, $M \rightarrow 0$, i.e., $\sqrt{-a/b} \rightarrow 0$,
 $\therefore a(T \rightarrow T_c^-) \rightarrow 0$.
Simply define $a = a_0 \frac{T - T_c}{T_c}$, $b = \text{Const.}$
- Put $M = \pm \sqrt{-a/b}$ into (2), $\therefore a < 0$ ($T < T_c$),
 $\Rightarrow b > 0$.
For $T > T_c$, $a > 0$.

Landau's approximation for the continuous phase transition

- $\frac{\partial F}{\partial M} = M(a + bM^2) = 0$ (1), $\frac{\partial^2 F}{\partial M^2} = a + 3bM^2 > 0$ (2)
- (1) $\rightarrow M = 0$, or $M = \pm\sqrt{-a/b}$.
- From the figure, $T \rightarrow T_c^-$, $M \rightarrow 0$, i.e., $\sqrt{-a/b} \rightarrow 0$,
 $\therefore a(T \rightarrow T_c^-) \rightarrow 0$.

Simply define $a = a_0 \frac{T - T_c}{T_c}$, $b = \text{Const.}$

- Put $M = \pm\sqrt{-a/b}$ into (2), $\therefore a < 0$ ($T < T_c$),
 $\Rightarrow b > 0$.

For $T > T_c$, $a > 0$.

- $T < T_c$, $a < 0$, from (2), $M = \pm\sqrt{-a/b}$ is the stable solution.

Landau's approximation for the continuous phase transition

- $\frac{\partial F}{\partial M} = M(a + bM^2) = 0$ (1), $\frac{\partial^2 F}{\partial M^2} = a + 3bM^2 > 0$ (2)
- (1) $\rightarrow M = 0$, or $M = \pm\sqrt{-a/b}$.
- From the figure, $T \rightarrow T_c^-$, $M \rightarrow 0$, i.e., $\sqrt{-a/b} \rightarrow 0$,
 $\therefore a(T \rightarrow T_c^-) \rightarrow 0$.

Simply define $a = a_0 \frac{T - T_c}{T_c}$, $b = \text{Const.}$

- Put $M = \pm\sqrt{-a/b}$ into (2), $\therefore a < 0$ ($T < T_c$),
 $\Rightarrow b > 0$.

For $T > T_c$, $a > 0$.

- $T < T_c$, $a < 0$, from (2), $M = \pm\sqrt{-a/b}$ is the stable solution.

$T > T_c$, $a > 0$, from (2), $M = 0$ is the stable solution.

Table of contents

① Phase Transition of Single-Component System

- 3.1 Criterion of thermal equilibrium
- 3.2 Basic equations of open system
- 3.3 Equilibrium of single-component multi-phase system
- 3.4 Properties of equilibrium of s-c multi-phase system
- 3.5 Critical point and phase change between gas and liquid
- 3.7 Classification of the phase transition
- 3.9 Landau's approximation for the continuous phase transition