Thermodynamics & Statistical Physics Chapter 8. Bose statistics and Fermi statistics

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• For Boson system, $a_l = \frac{\omega_l}{e^{\alpha+\beta\varepsilon_l}-1}$.

$$\mathcal{Z} = \prod \mathcal{Z}_l = \prod (1 - e^{-\alpha - \beta \varepsilon_l})^{-\omega_l}$$

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• For Boson system, $a_l = \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} - 1}$. Define Grand partition function:

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 $\Rightarrow N = -\frac{\partial}{\partial \alpha} \ln \mathcal{Z}$.

$$\frac{\mathcal{Z} = \prod \mathcal{Z}_l = \prod (1 - e^{-\alpha - \beta \varepsilon_l})^{-\omega_l}}{\ln \mathcal{Z} = -\sum \omega_l \ln (1 - e^{-\alpha - \beta \varepsilon_l})}.$$

- $\frac{\partial}{\partial \alpha} \ln \mathcal{Z} = -\sum \frac{\omega_l}{1 e^{-\alpha \beta \varepsilon_l}} \cdot e^{-\alpha \beta \varepsilon_l} = -\sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} 1}$ $=-\sum a_l$. $\Rightarrow N = -\frac{\partial}{\partial \alpha} \ln \mathcal{Z}$.
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- $\frac{\partial}{\partial \alpha} \ln \mathcal{Z} = -\sum_{l=\alpha-\beta \varepsilon_l} \frac{\omega_l}{1 e^{-\alpha-\beta \varepsilon_l}} \cdot e^{-\alpha-\beta \varepsilon_l} = -\sum_{l=\alpha+\beta \varepsilon_l} \frac{\omega_l}{1 e^{-\alpha-\beta \varepsilon_l}}$ $=-\sum a_l$. $\Rightarrow |N = -\frac{\partial}{\partial \alpha} \ln \mathcal{Z}|.$
- $\frac{\partial}{\partial \beta} \ln \mathcal{Z} = -\sum_{l=e^{-\alpha-\beta\varepsilon_l}} \frac{\omega_l}{1-e^{-\alpha-\beta\varepsilon_l}} \cdot e^{-\alpha-\beta\varepsilon_l} \cdot \varepsilon_l$

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• For Boson system, $a_l = \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} - 1}$. Define Grand partition function:

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- $\frac{\partial}{\partial \beta} \ln \mathcal{Z} = -\sum \frac{\omega_l}{1 e^{-\alpha \beta \varepsilon_l}} \cdot e^{-\alpha \beta \varepsilon_l} \cdot \varepsilon_l = -\sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} 1} \varepsilon_l$ $=-\sum a_l \varepsilon_l$. $\Rightarrow |U = -\frac{\partial}{\partial \beta} \ln \mathcal{Z}|.$

$$\ln \mathcal{Z} = -\sum \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

• Notice $Y = \sum \frac{\partial \varepsilon_l}{\partial u} a_l = \sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} - 1} \frac{\partial \varepsilon_l}{\partial u}$.

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$$Y=\sum rac{\partial arepsilon_l}{\partial y}a_l=\sum rac{\omega_l}{e^{lpha+eta arepsilon_l-1}}rac{\partial arepsilon_l}{\partial y}.$$

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$$\frac{\partial}{\partial y} \ln \mathcal{Z} = -\sum_{l} \frac{\omega_l}{1 - e^{-\alpha - \beta \varepsilon_l}} e^{-\alpha - \beta \varepsilon_l} \cdot \beta \frac{\partial \varepsilon_l}{\partial y}$$

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$$\frac{\partial}{\partial y} \ln \mathcal{Z} = -\sum_{l=e^{-\alpha-\beta\varepsilon_l}} \frac{\omega_l}{1-e^{-\alpha-\beta\varepsilon_l}} e^{-\alpha-\beta\varepsilon_l} \cdot \beta \frac{\partial \varepsilon_l}{\partial y} = -\beta \sum_{l=e^{\alpha+\beta\varepsilon_l}} \frac{\omega_l}{\partial y} \frac{\partial \varepsilon_l}{\partial y}.$$

$$\Rightarrow Y = -\frac{1}{\beta} \frac{\partial}{\partial y} \ln \mathcal{Z}.$$

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$$\frac{\partial}{\partial y} \ln \mathcal{Z} = -\sum \frac{\omega_l}{1 - e^{-\alpha - \beta \varepsilon_l}} e^{-\alpha - \beta \varepsilon_l} \cdot \beta \frac{\partial \varepsilon_l}{\partial y} = -\beta \sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} - 1} \frac{\partial \varepsilon_l}{\partial y}.$$

$$\Rightarrow V = -\frac{1}{2} \frac{\partial}{\partial z} \ln \mathcal{Z}. \text{ Pressure } p = \frac{1}{2} \frac{\partial}{\partial z} \ln \mathcal{Z}.$$

$$\Rightarrow Y = -\frac{1}{\beta} \frac{\partial}{\partial y} \ln \mathcal{Z}$$
. Pressure $p = \frac{1}{\beta} \frac{\partial}{\partial V} \ln \mathcal{Z}$.

$$\ln \mathcal{Z} = -\sum \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

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$$\begin{split} &\frac{\partial}{\partial y} \ln \mathcal{Z} = -\sum \frac{\omega_l}{1 - e^{-\alpha - \beta \varepsilon_l}} e^{-\alpha - \beta \varepsilon_l} \cdot \beta \frac{\partial \varepsilon_l}{\partial y} = -\beta \sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} - 1} \frac{\partial \varepsilon_l}{\partial y}. \\ &\Rightarrow \boxed{Y = -\frac{1}{\beta} \frac{\partial}{\partial y} \ln \mathcal{Z}}. \quad \text{Pressure } p = \frac{1}{\beta} \frac{\partial}{\partial V} \ln \mathcal{Z}. \end{split}$$

• Together with expresses of N, U, Y,

$$\ln \mathcal{Z} = -\sum \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

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$$\beta(\mathrm{d}U - Y\mathrm{d}y + \frac{\alpha}{\beta}\mathrm{d}N)$$

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 Pressure $p = \frac{1}{\beta} \frac{\partial}{\partial V} \ln \mathcal{Z}.$

• Together with expresses of N, U, Y,

$$\beta(dU - Ydy + \frac{\alpha}{\beta}dN) = -\beta d\frac{\partial \ln \mathcal{Z}}{\partial \beta} + \frac{\partial \ln \mathcal{Z}}{\partial y}dy - \alpha d\frac{\partial \ln \mathcal{Z}}{\partial \alpha}$$
$$= -d(\beta \frac{\partial \ln \mathcal{Z}}{\partial \beta}) + \frac{\partial \ln \mathcal{Z}}{\partial \beta}d\beta + \frac{\partial \ln \mathcal{Z}}{\partial y}dy - d(\alpha \frac{\partial \ln \mathcal{Z}}{\partial \alpha}) + \frac{\partial \ln \mathcal{Z}}{\partial \alpha}d\alpha$$

$$\ln \mathcal{Z} = -\sum \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

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$$\beta(\mathrm{d}U - Y\mathrm{d}y + \frac{\alpha}{\beta}\mathrm{d}N) = -\beta\mathrm{d}\frac{\partial\ln\mathcal{Z}}{\partial\beta} + \frac{\partial\ln\mathcal{Z}}{\partial y}\mathrm{d}y - \alpha\mathrm{d}\frac{\partial\ln\mathcal{Z}}{\partial\alpha}$$

$$= -\mathrm{d}(\beta\frac{\partial\ln\mathcal{Z}}{\partial\beta}) + \frac{\partial\ln\mathcal{Z}}{\partial\beta}\mathrm{d}\beta + \frac{\partial\ln\mathcal{Z}}{\partial y}\mathrm{d}y - \mathrm{d}(\alpha\frac{\partial\ln\mathcal{Z}}{\partial\alpha}) + \frac{\partial\ln\mathcal{Z}}{\partial\alpha}\mathrm{d}\alpha$$

$$= -\mathrm{d}(\beta\frac{\partial\ln\mathcal{Z}}{\partial\beta}) - \mathrm{d}(\alpha\frac{\partial\ln\mathcal{Z}}{\partial\alpha}) + (\frac{\partial\ln\mathcal{Z}}{\partial\beta}\mathrm{d}\beta + \frac{\partial\ln\mathcal{Z}}{\partial y}\mathrm{d}y + \frac{\partial\ln\mathcal{Z}}{\partial\alpha}\mathrm{d}\alpha)$$

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$$\Rightarrow Y = -\frac{1}{\beta} \frac{\partial}{\partial y} \ln \mathcal{Z}. \text{ Pressure } p = \frac{1}{\beta} \frac{\partial}{\partial V} \ln \mathcal{Z}.$$

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$$\beta(\mathrm{d}U - Y\mathrm{d}y + \frac{\alpha}{\beta}\mathrm{d}N) = -\beta\mathrm{d}\frac{\partial\ln\mathcal{Z}}{\partial\beta} + \frac{\partial\ln\mathcal{Z}}{\partial y}\mathrm{d}y - \alpha\mathrm{d}\frac{\partial\ln\mathcal{Z}}{\partial\alpha}$$

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$$= -\mathrm{d}(\beta\frac{\partial\ln\mathcal{Z}}{\partial\beta}) - \mathrm{d}(\alpha\frac{\partial\ln\mathcal{Z}}{\partial\alpha}) + (\frac{\partial\ln\mathcal{Z}}{\partial\beta}\mathrm{d}\beta + \frac{\partial\ln\mathcal{Z}}{\partial y}\mathrm{d}y + \frac{\partial\ln\mathcal{Z}}{\partial\alpha}\mathrm{d}\alpha)$$

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$$\ln \mathcal{Z} = -\sum \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

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$$\Rightarrow Y = -\frac{1}{\beta} \frac{\partial}{\partial y} \ln \mathcal{Z}. \text{ Pressure } p = \frac{1}{\beta} \frac{\partial}{\partial V} \ln \mathcal{Z}.$$

• Together with expresses of N, U, Y,

$$\beta(\mathrm{d}U - Y\mathrm{d}y + \frac{\alpha}{\beta}\mathrm{d}N) = -\beta\mathrm{d}\frac{\partial\ln\mathcal{Z}}{\partial\beta} + \frac{\partial\ln\mathcal{Z}}{\partial\gamma}\mathrm{d}y - \alpha\mathrm{d}\frac{\partial\ln\mathcal{Z}}{\partial\alpha}$$

$$= -\mathrm{d}(\beta\frac{\partial\ln\mathcal{Z}}{\partial\beta}) + \frac{\partial\ln\mathcal{Z}}{\partial\beta}\mathrm{d}\beta + \frac{\partial\ln\mathcal{Z}}{\partial\gamma}\mathrm{d}y - \mathrm{d}(\alpha\frac{\partial\ln\mathcal{Z}}{\partial\alpha}) + \frac{\partial\ln\mathcal{Z}}{\partial\alpha}\mathrm{d}\alpha$$

$$= -\mathrm{d}(\beta\frac{\partial\ln\mathcal{Z}}{\partial\beta}) - \mathrm{d}(\alpha\frac{\partial\ln\mathcal{Z}}{\partial\alpha}) + (\frac{\partial\ln\mathcal{Z}}{\partial\beta}\mathrm{d}\beta + \frac{\partial\ln\mathcal{Z}}{\partial\gamma}\mathrm{d}y + \frac{\partial\ln\mathcal{Z}}{\partial\alpha}\mathrm{d}\alpha)$$

$$= -\mathrm{d}(\beta\frac{\partial\ln\mathcal{Z}}{\partial\beta}) - \mathrm{d}(\alpha\frac{\partial\ln\mathcal{Z}}{\partial\alpha}) + \mathrm{d}\ln\mathcal{Z}$$

$$= \mathrm{d}(\ln\mathcal{Z} - \alpha\frac{\partial\ln\mathcal{Z}}{\partial\alpha} - \beta\frac{\partial\ln\mathcal{Z}}{\partial\beta}).$$

§8. Bose statistics and Fermi statistics 8.1 Thermal parameters in statistics

$$\ln \mathcal{Z} = -\sum \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

• Eq. (4.1.12), $dU = TdS - pdV + \sum \mu_i dn_i$.

$$\ln \mathcal{Z} = -\sum \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

• Eq. (4.1.12), $dU = TdS - pdV + \sum \mu_i dn_i$. $-p \to Y$, $dV \to dy$, $n_i \to N$,

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• Eq. (4.1.12), $dU = TdS - pdV + \sum \mu_i dn_i$. $-n \to Y$. $dV \to dy$, $n_i \to N$, $\Rightarrow dS = \frac{1}{T}(dU - Ydy - \mu dN);$

$$\ln \mathcal{Z} = -\sum \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

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§8. Bose statistics and Fermi statistics

$$\ln \mathcal{Z} = -\sum \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

• Eq. (4.1.12), $dU = TdS - pdV + \sum \mu_i dn_i$. $-p \to Y$. $dV \to du$. $n_i \to N$. $\Rightarrow dS = \frac{1}{T}(dU - Ydy - \mu dN);$ $d(\ln \mathcal{Z} - \alpha \frac{\partial \ln \mathcal{Z}}{\partial \alpha} - \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta}) = \beta (dU - Y dy + \frac{\alpha}{\beta} dN).$ $\Rightarrow \beta = \frac{1}{kT}, \ \alpha = -\beta \mu = -\frac{\mu}{kT}.$ β —temperature, α —chemical potential.

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• Eq. (4.1.12), $dU = TdS - pdV + \sum \mu_i dn_i$. $-p \to Y$. $dV \to du$. $n_i \to N$. $\Rightarrow dS = \frac{1}{T}(dU - Ydy - \mu dN);$ $d(\ln \mathcal{Z} - \alpha \frac{\partial \ln \mathcal{Z}}{\partial \alpha} - \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta}) = \beta (dU - Y dy + \frac{\alpha}{\beta} dN).$ $\Rightarrow \beta = \frac{1}{kT}, \ \alpha = -\beta \mu = -\frac{\mu}{kT}.$ β —temperature, α —chemical potential.

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• $S = k(\ln \mathcal{Z} - \alpha \frac{\partial \ln \mathcal{Z}}{\partial \alpha} - \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta})$

$$\ln \mathcal{Z} = -\sum \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

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- $S = k(\ln \mathcal{Z} \alpha \frac{\partial \ln \mathcal{Z}}{\partial \alpha} \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta})$ $=k[-\sum \omega_l \ln(1-e^{-\alpha-\beta\varepsilon_l})+\alpha\sum \frac{\omega_l}{e^{\alpha+\beta\varepsilon_l-1}}$ $+\beta\sum_{\alpha}\frac{\omega_{l}\varepsilon_{l}}{\alpha+\beta\varepsilon_{l}-1}$];

$$\ln \mathcal{Z} = -\sum \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

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$$\begin{split} \bullet \ S &= k (\ln \mathcal{Z} - \alpha \frac{\partial \ln \mathcal{Z}}{\partial \alpha} - \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta}) \\ &= k [-\sum_{e} \omega_l \ln (1 - e^{-\alpha - \beta \varepsilon_l}) + \alpha \sum_{e} \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} - 1} \\ &+ \beta \sum_{e} \frac{\omega_l \varepsilon_l}{e^{\alpha + \beta \varepsilon_l} - 1}]; \\ \text{Notice:} \ \ a_l &= \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} - 1} \Rightarrow e^{\alpha + \beta \varepsilon_l} - 1 = \frac{\omega_l}{a_l}, \ \text{then} \end{split}$$

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• Eq. (4.1.12), $dU = TdS - pdV + \sum \mu_i dn_i$. $-p \to Y$. $dV \to du$. $n_i \to N$. $\Rightarrow dS = \frac{1}{T}(dU - Ydy - \mu dN);$ $d(\ln \mathcal{Z} - \alpha \frac{\partial \ln \mathcal{Z}}{\partial \alpha} - \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta}) = \beta (dU - Y dy + \frac{\alpha}{\beta} dN).$ $\Rightarrow \beta = \frac{1}{kT}, \ \alpha = -\beta \mu = -\frac{\mu}{kT}.$ β —temperature, α —chemical potential.

•
$$S = k(\ln \mathcal{Z} - \alpha \frac{\partial \ln \mathcal{Z}}{\partial \alpha} - \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta})$$

 $= k[-\sum_{\alpha l} \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l}) + \alpha \sum_{\alpha l} \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} - 1}];$

Notice: $a_l = \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} - 1} \Rightarrow e^{\alpha + \beta \varepsilon_l} - 1 = \frac{\omega_l}{\alpha}$, then $1 - e^{-\alpha - \beta \varepsilon_l} = 1 - (\frac{\omega_l}{a_l} + 1)^{-1} = 1 - \frac{a_l}{\omega_l + a_l} = \frac{\omega_l}{\omega_l + a_l}$ §8. Bose statistics and Fermi statistics 8.1 Thermal parameters in statistics

$$e^{\alpha+\beta\varepsilon_l}-1=\frac{\omega_l}{a_l},\ 1-e^{-\alpha-\beta\varepsilon_l}=\frac{\omega_l}{\omega_l+a_l}$$

•
$$S = k\left[-\sum_{\substack{\omega_l \varepsilon_l \\ e^{\alpha + \beta \varepsilon_l} - 1}} \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l}) + \alpha \sum_{\substack{\alpha \leq l \\ e^{\alpha + \beta \varepsilon_l} - 1}} \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l}}\right]$$

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§8. Bose statistics and Fermi statistics 8.1 Thermal parameters in statistics

$$e^{\alpha+\beta\varepsilon_l}-1=\frac{\omega_l}{a_l}$$
, $1-e^{-\alpha-\beta\varepsilon_l}=\frac{\omega_l}{\omega_l+a_l}$

•
$$S = k\left[-\sum_{\substack{\omega_l \varepsilon_l \\ e^{\alpha + \beta \varepsilon_l - 1}}} \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l}) + \alpha \sum_{\substack{e^{\alpha + \beta \varepsilon_l - 1} \\ e^{\alpha + \beta \varepsilon_l - 1}}}\right]$$

= $k\left[-\sum_{\substack{\omega_l \ln \frac{\omega_l}{a_l + \omega_l}}} + \sum_{\substack{\omega_l \frac{\alpha + \beta \varepsilon_l}{e^{\alpha + \beta \varepsilon_l - 1}}}}\right]$

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= $k \left[-\sum_{\substack{\omega_l \ln \frac{\omega_l}{a_l + \omega_l}}} + \sum_{\substack{\omega_l \ln \frac{\omega_l/a_l + 1}{\omega_l/a_l}}} \right]$

$$e^{\alpha+\beta\varepsilon_l}-1=\frac{\omega_l}{a_l}$$
, $1-e^{-\alpha-\beta\varepsilon_l}=\frac{\omega_l}{\omega_l+a_l}$

•
$$S = k \left[-\sum_{\substack{\omega_l \varepsilon_l \\ e^{\alpha + \beta \varepsilon_l - 1}}} \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l}) + \alpha \sum_{\substack{e^{\alpha_l + \beta \varepsilon_l - 1} \\ e^{\alpha + \beta \varepsilon_l - 1}}} \right]$$

= $k \left[-\sum_{\substack{\omega_l \ln \frac{\omega_l}{a_l + \omega_l}}} + \sum_{\substack{\omega_l \frac{\alpha + \beta \varepsilon_l}{e^{\alpha + \beta \varepsilon_l - 1}}}} \right]$
= $k \left[-\sum_{\substack{\omega_l \ln \frac{\omega_l}{a_l + \omega_l}}} + \sum_{\substack{\omega_l \ln \frac{\omega_l / a_l + 1}{\omega_l / a_l}}} \right]$
= $k \left[-\sum_{\substack{\omega_l \ln \frac{\omega_l}{a_l + \omega_l}}} + \sum_{\substack{\alpha_l \ln \frac{\omega_l + a_l}{a_l}}} \right]$

$$e^{\alpha+\beta\varepsilon_l}-1=\frac{\omega_l}{a_l}$$
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= $k \left[\sum_{\alpha_l \log \frac{\omega_l}{a_l + \omega_l}} + \sum_{\alpha_l \log \frac{\omega_l}{a_l}} + \sum_{\alpha_l \log \frac{\omega_l}{a_l}} \right]$

$$e^{\alpha+\beta\varepsilon_l}-1=\frac{\omega_l}{a_l}$$
, $1-e^{-\alpha-\beta\varepsilon_l}=\frac{\omega_l}{\omega_l+a_l}$

•
$$S = k\left[-\sum_{\substack{\omega_l \in l \\ e^{\alpha + \beta \varepsilon_l - 1}}} \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l}) + \alpha \sum_{\substack{\alpha \in \alpha + \beta \varepsilon_l - 1 \\ e^{\alpha + \beta \varepsilon_l - 1}}}\right]$$

= $k\left[-\sum_{\alpha l} \omega_l \ln \frac{\omega_l}{a_l + \omega_l} + \sum_{\alpha l} \omega_l \frac{\alpha + \beta \varepsilon_l}{e^{\alpha + \beta \varepsilon_l - 1}}\right]$
= $k\left[-\sum_{\alpha l} \omega_l \ln \frac{\omega_l}{a_l + \omega_l} + \sum_{\alpha l} \omega_l \frac{\ln(\omega_l/a_l + 1)}{\omega_l/a_l}\right]$
= $k\left[-\sum_{\alpha l} \omega_l \ln \frac{\omega_l}{a_l + \omega_l} + \sum_{\alpha l} a_l \ln \frac{\omega_l + a_l}{a_l}\right]$
= $k\left[\sum_{\alpha l} \omega_l \ln(a_l + \omega_l) - \sum_{\alpha l} \omega_l \ln(\omega_l + a_l)\right]$
= $k\left[\sum_{\alpha l} (a_l + \omega_l) \ln(a_l + \omega_l) - \sum_{\alpha l} \omega_l \ln(\omega_l - \sum_{\alpha l} a_l \ln a_l)\right]$

$$e^{\alpha+\beta\varepsilon_l}-1=\frac{\omega_l}{a_l}$$
, $1-e^{-\alpha-\beta\varepsilon_l}=\frac{\omega_l}{\omega_l+a_l}$

•
$$S = k\left[-\sum_{\substack{\omega_l \varepsilon_l \\ e^{\alpha+\beta\varepsilon_l}-1}} \omega_l \ln(1-e^{-\alpha-\beta\varepsilon_l}) + \alpha \sum_{\substack{e^{\alpha+\beta\varepsilon_l - 1} \\ e^{\alpha+\beta\varepsilon_l}-1}} \right]$$

$$= k\left[-\sum_{\substack{\omega_l \ln \frac{\omega_l }{a_l+\omega_l}}} + \sum_{\substack{\omega_l \frac{\alpha+\beta\varepsilon_l }{e^{\alpha+\beta\varepsilon_l - 1}}}} \right]$$

$$= k\left[-\sum_{\substack{\omega_l \ln \frac{\omega_l }{a_l+\omega_l}}} + \sum_{\substack{\omega_l \ln \frac{\omega_l/a_l}{\omega_l/a_l}}} \right]$$

$$= k\left[-\sum_{\substack{\omega_l \ln \frac{\omega_l }{a_l+\omega_l}}} + \sum_{\substack{\alpha_l \ln \frac{\omega_l+a_l}{a_l}}} \right]$$

$$= k\left[\sum_{\substack{\omega_l \ln a_l}} \omega_l \ln(a_l+\omega_l) - \sum_{\substack{\omega_l \ln \omega_l }} \omega_l \ln(\omega_l+a_l) - \sum_{\substack{\omega_l \ln a_l}} \omega_l \ln(\omega_l+a_l) \right]$$

$$= k\left[\sum_{\substack{\alpha_l \ln a_l}} (a_l+\omega_l) \ln(a_l+\omega_l) - \sum_{\substack{\omega_l \ln a_l}} \omega_l \ln(a_l+\omega_l) - \sum_{\substack{\alpha_l \ln a_l}} \omega_l \ln(a_l+\omega_l) \right].$$
Comparing with eq.(6.7.4),

88. Bose statistics and Fermi statistics 8.1 Thermal parameters in statistics

$$e^{\alpha+\beta\varepsilon_l}-1=\frac{\omega_l}{a_l}$$
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Thermal parameters in statistics for Fermion system

Thermal parameters in statistics for Fermion system

Grand partition function:

$$\mathcal{Z} = \prod \mathcal{Z}_l = \prod (1 + e^{-\alpha - \beta \varepsilon_l})^{\omega_l}.$$

$$\ln \mathcal{Z} = \sum \omega_l \ln(1 + e^{-\alpha - \beta \varepsilon_l}).$$

- $U = -\frac{\partial}{\partial \beta} \ln \mathcal{Z}$.
- $\bullet \mid Y = -\frac{1}{\beta} \frac{\partial}{\partial y} \ln \mathcal{Z} \mid.$
- $\bullet | S = k \ln \Omega |$

Method for Bose (Fermi) statistics:

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$$\{\varepsilon_l\}, \{\omega_l\} \to \ln \mathcal{Z}(\alpha, \beta, y)$$

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$$J = U - TS - \mu N$$
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- Take (8.1.10) in, $J = U \frac{1}{k\beta} \cdot k(\ln \mathcal{Z} + \alpha N + \beta U) \mu N$

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- Take (8.1.10) in, $J = U - \frac{1}{k\beta} \cdot k(\ln \mathcal{Z} + \alpha N + \beta U) - \mu N$ $=U-\frac{1}{\beta}(\ln \mathcal{Z}-\mu\beta N+\beta U)-\mu N=-\frac{1}{\beta}\ln \mathcal{Z}$ $=-kT \ln \mathcal{Z}$

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 - 8.3 Bose-Einstein Condensate (BEC)
 - 8.4 Photon gas
 - 8.5 Free electron gas in metal

• Non-degeneration condition: $e^{-\alpha} \ll 1$ or $n\lambda^3 \ll 1$. (7.2.6, 7.2.7).

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$$D(\varepsilon)d\varepsilon = g\frac{2\pi V}{h^3}(2m)^{\frac{3}{2}}\varepsilon^{\frac{1}{2}}d\varepsilon$$

Number of particles:

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Number of particles:

$$N = \sum a_l$$

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$$N = \sum a_l = \sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} \pm 1}$$



$$D(\varepsilon)d\varepsilon = g^{\frac{2\pi V}{h^3}}(2m)^{\frac{3}{2}}\varepsilon^{\frac{1}{2}}d\varepsilon$$

$$N = \sum a_l = \sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l \pm 1}} = g^{\frac{2\pi V}{h^3}} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\alpha + \beta \varepsilon} \pm 1}.$$



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• Internal energy:

$$D(\varepsilon)d\varepsilon = g\frac{2\pi V}{h^3}(2m)^{\frac{3}{2}}\varepsilon^{\frac{1}{2}}d\varepsilon$$

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• Internal energy:

$$U = \sum a_l \varepsilon_l$$

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Internal energy:

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$$N = \sum a_l = \sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} \pm 1} = g \frac{2\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\alpha + \beta \varepsilon} \pm 1}.$$

• Internal energy:

$$U = \sum a_l \varepsilon_l = \sum \frac{\omega_l \varepsilon_l}{e^{\alpha + \beta \varepsilon_l \pm 1}} = g \frac{2\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{3}{2}} d\varepsilon}{e^{\alpha + \beta \varepsilon} \pm 1}.$$

$$D(\varepsilon)d\varepsilon = g\frac{2\pi V}{h^3}(2m)^{\frac{3}{2}}\varepsilon^{\frac{1}{2}}d\varepsilon$$

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• Set $x = \beta \varepsilon$,

$$D(\varepsilon)d\varepsilon = g\frac{2\pi V}{h^3}(2m)^{\frac{3}{2}}\varepsilon^{\frac{1}{2}}d\varepsilon$$

$$N = \sum a_l = \sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} \pm 1} = g \frac{2\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\alpha + \beta \varepsilon} \pm 1}.$$

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• Set $x = \beta \varepsilon$,

$$N = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{\alpha + x} \pm 1}$$

$$D(\varepsilon)d\varepsilon = g\frac{2\pi V}{h^3}(2m)^{\frac{3}{2}}\varepsilon^{\frac{1}{2}}d\varepsilon$$

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$$U = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty \frac{x^{\frac{3}{2}} dx}{e^{\alpha + x} \pm 1}.$$

$$D(\varepsilon)d\varepsilon = g\frac{2\pi V}{h^3}(2m)^{\frac{3}{2}}\varepsilon^{\frac{1}{2}}d\varepsilon$$

$$N = \sum a_l = \sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l \pm 1}} = g \frac{2\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\alpha + \beta \varepsilon} \pm 1}.$$

Internal energy:

$$U = \sum a_l \varepsilon_l = \sum_{\substack{e^{\alpha + \beta \varepsilon_l \pm 1}}} = g^{\frac{2\pi V}{h^3}} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{3}{2}} d\varepsilon}{e^{\alpha + \beta \varepsilon \pm 1}}.$$

• Set $x = \beta \varepsilon$,

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• As $e^{-\alpha} < 1$, x > 0, $\Rightarrow e^{-\alpha - x} < 1$.

$$D(\varepsilon)d\varepsilon = g^{\frac{2\pi V}{h^3}}(2m)^{\frac{3}{2}}\varepsilon^{\frac{1}{2}}d\varepsilon$$

$$N = \sum a_l = \sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l \pm 1}} = g \frac{2\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\alpha + \beta \varepsilon} \pm 1}.$$

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$$U = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty \frac{x^{\frac{3}{2}} dx}{e^{\alpha + x} \pm 1}.$$

• As $e^{-\alpha}<1$, x>0, $\Rightarrow e^{-\alpha-x}<1$. $\frac{1}{1+e^{-\alpha-x}}\simeq 1\mp e^{-\alpha-x};$

$$D(\varepsilon)d\varepsilon = g\frac{2\pi V}{h^3}(2m)^{\frac{3}{2}}\varepsilon^{\frac{1}{2}}d\varepsilon$$

$$N = \sum a_l = \sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l \pm 1}} = g \frac{2\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\alpha + \beta \varepsilon} \pm 1}.$$

Internal energy:

$$U = \sum a_l \varepsilon_l = \sum_{\substack{e^{\alpha + \beta \varepsilon_l \pm 1}}} = g^{\frac{2\pi V}{h^3}} (2m)^{\frac{3}{2}} \int_0^{\infty} \frac{\varepsilon^{\frac{3}{2}} d\varepsilon}{e^{\alpha + \beta \varepsilon} \pm 1}.$$

• Set $x = \beta \varepsilon$,

$$N = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{\alpha + x} \pm 1},$$

$$U = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty \frac{x^{\frac{3}{2}} dx}{e^{\alpha + x} \pm 1}.$$

 $\begin{array}{l} \bullet \text{ As } e^{-\alpha} < 1, \ x > 0, \Rightarrow e^{-\alpha - x} < 1. \\ \frac{1}{1 \pm e^{-\alpha - x}} \simeq 1 \mp e^{-\alpha - x}; \\ \frac{1}{e^{\alpha + x} \pm 1} = \frac{1}{e^{\alpha + x}(1 \pm e^{-\alpha - x})} \end{array}$

$$D(\varepsilon)d\varepsilon = g\frac{2\pi V}{h^3}(2m)^{\frac{3}{2}}\varepsilon^{\frac{1}{2}}d\varepsilon$$

$$N = \sum a_l = \sum \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l \pm 1}} = g \frac{2\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\alpha + \beta \varepsilon} \pm 1}.$$

Internal energy:

$$U = \sum a_l \varepsilon_l = \sum_{\substack{e^{\alpha + \beta \varepsilon_l \pm 1} \\ e^{\alpha + \beta \varepsilon_l \pm 1}}} = g^{\frac{2\pi V}{h^3}} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{3}{2}} d\varepsilon}{e^{\alpha + \beta \varepsilon} \pm 1}.$$

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$$N = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{\alpha + x} \pm 1},$$

$$U = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty \frac{x^{\frac{3}{2}} dx}{e^{\alpha + x} \pm 1}.$$

• As $e^{-\alpha} < 1$, x > 0, $\Rightarrow e^{-\alpha - x} < 1$. $\frac{1}{1 \pm e^{-\alpha - x}} \simeq 1 \mp e^{-\alpha - x};$ $\frac{1}{e^{\alpha + x} \pm 1} = \frac{1}{e^{\alpha + x}(1 \pm e^{-\alpha - x})} \simeq e^{-\alpha - x}(1 \mp e^{-\alpha - x}).$

$$\frac{1}{e^{\alpha+x}\pm 1} \simeq e^{-\alpha-x} (1 \mp e^{-\alpha-x})$$

• Math:
$$\Gamma(n) \equiv \int_0^\infty e^{-x} x^{n-1} dx$$
; $\Gamma(n+\frac{1}{2}) = (n-\frac{1}{2}) \cdot (n-\frac{3}{2})...\frac{1}{2}\sqrt{\pi}$.

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•
$$N = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{\alpha + x} \pm 1}$$

$$\frac{1}{e^{\alpha+x}\pm 1} \simeq e^{-\alpha-x} (1 \mp e^{-\alpha-x})$$

- Math: $\Gamma(n) \equiv \int_0^\infty e^{-x} x^{n-1} dx$; $\Gamma(n+\frac{1}{2}) = (n-\frac{1}{2}) \cdot (n-\frac{3}{2}) \dots \frac{1}{2} \sqrt{\pi}$.
- $N = g^{\frac{2\pi V}{h^3}} (2mkT)^{\frac{3}{2}} \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{\alpha + x + 1}}$ $\simeq q^{\frac{2\pi V}{1.3}} (2mkT)^{\frac{3}{2}} \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{1}{2}} dx$

$$\frac{1}{e^{\alpha+x}\pm 1} \simeq e^{-\alpha-x} (1 \mp e^{-\alpha-x})$$

- Math: $\Gamma(n) \equiv \int_0^\infty e^{-x} x^{n-1} dx$; $\Gamma(n+\frac{1}{2}) = (n-\frac{1}{2}) \cdot (n-\frac{3}{2})...\frac{1}{2}\sqrt{\pi}.$
- $N = g^{\frac{2\pi V}{b^3}} (2mkT)^{\frac{3}{2}} \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{\alpha + x + 1}}$ $\simeq q^{\frac{2\pi V}{L^3}} (2mkT)^{\frac{3}{2}} \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{1}{2}} dx$ $=q^{\frac{2\pi V}{h^3}}(2mkT)^{\frac{3}{2}}e^{-\alpha}$ $\cdot \left| \int_0^\infty e^{-x} x^{\frac{1}{2}} dx \mp \int_0^\infty e^{-\alpha} e^{-2x} x^{\frac{1}{2}} dx \right|$

$$\frac{1}{e^{\alpha+x}+1} \simeq e^{-\alpha-x} (1 \mp e^{-\alpha-x})$$

- Math: $\Gamma(n) \equiv \int_0^\infty e^{-x} x^{n-1} dx$; $\Gamma(n + \frac{1}{2}) = (n \frac{1}{2}) \cdot (n \frac{3}{2}) \dots \frac{1}{2} \sqrt{\pi}$.
- $N = g^{\frac{2\pi V}{h^3}} (2mkT)^{\frac{3}{2}} \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{\alpha + x} \pm 1}$ $\simeq g^{\frac{2\pi V}{h^3}} (2mkT)^{\frac{3}{2}} \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{1}{2}} dx$ $= g^{\frac{2\pi V}{h^3}} (2mkT)^{\frac{3}{2}} e^{-\alpha}$ $\cdot \left[\int_0^\infty e^{-x} x^{\frac{1}{2}} dx \mp \int_0^\infty e^{-\alpha} e^{-2x} x^{\frac{1}{2}} dx \right]$ $= g^{\frac{2\pi V}{h^3}} (2mkT)^{\frac{3}{2}} e^{-\alpha}$ $\cdot \left[\int_0^\infty e^{-x} x^{\frac{1}{2}} dx \mp e^{-\alpha} 2^{-\frac{3}{2}} \int_0^\infty e^{-2x} (2x)^{\frac{1}{2}} d(2x) \right]$

$$\frac{1}{e^{\alpha+x}\pm 1} \simeq e^{-\alpha-x} (1 \mp e^{-\alpha-x})$$

- Math: $\Gamma(n) \equiv \int_0^\infty e^{-x} x^{n-1} dx$; $\Gamma(n + \frac{1}{2}) = (n \frac{1}{2}) \cdot (n \frac{3}{2}) \dots \frac{1}{2} \sqrt{\pi}$.
- $N = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{\alpha + x + 1}}$ $\simeq q^{\frac{2\pi V}{L^3}} (2mkT)^{\frac{3}{2}} \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{1}{2}} dx$ $=q^{\frac{2\pi V}{h^3}}(2mkT)^{\frac{3}{2}}e^{-\alpha}$ $\cdot \left| \int_0^\infty e^{-x} x^{\frac{1}{2}} dx \mp \int_0^\infty e^{-\alpha} e^{-2x} x^{\frac{1}{2}} dx \right|$ $=g^{\frac{2\pi V}{h^3}}(2mkT)^{\frac{3}{2}}e^{-\alpha}$ $\cdot \left| \int_0^\infty e^{-x} x^{\frac{1}{2}} dx \mp e^{-\alpha} 2^{-\frac{3}{2}} \int_0^\infty e^{-2x} (2x)^{\frac{1}{2}} d(2x) \right|$ $= g^{\frac{2\pi V}{h^3}} (2mkT)^{\frac{3}{2}} e^{-\alpha} [\Gamma(\frac{3}{2}) \mp e^{-\alpha} 2^{-\frac{3}{2}} \Gamma(\frac{3}{2})]$

$$\frac{1}{e^{\alpha+x}\pm 1} \simeq e^{-\alpha-x} (1 \mp e^{-\alpha-x})$$

• Math: $\Gamma(n) \equiv \int_0^\infty e^{-x} x^{n-1} dx$; $\Gamma(n+\frac{1}{2}) = (n-\frac{1}{2}) \cdot (n-\frac{3}{2})...\frac{1}{2}\sqrt{\pi}.$

•
$$N = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{\alpha + x} \pm 1}$$

 $\simeq g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{1}{2}} dx$
 $= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} e^{-\alpha}$
 $\cdot \left[\int_0^\infty e^{-x} x^{\frac{1}{2}} dx \mp \int_0^\infty e^{-\alpha} e^{-2x} x^{\frac{1}{2}} dx \right]$
 $= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} e^{-\alpha}$
 $\cdot \left[\int_0^\infty e^{-x} x^{\frac{1}{2}} dx \mp e^{-\alpha} 2^{-\frac{3}{2}} \int_0^\infty e^{-2x} (2x)^{\frac{1}{2}} d(2x) \right]$
 $= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} e^{-\alpha} \left[\Gamma(\frac{3}{2}) \mp e^{-\alpha} 2^{-\frac{3}{2}} \Gamma(\frac{3}{2}) \right]$
 $= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} e^{-\alpha} \frac{\sqrt{\pi}}{2} (1 \mp e^{-\alpha} 2^{-\frac{3}{2}}) \dots$

$$\frac{1}{e^{\alpha+x}\pm 1} = e^{-\alpha-x} (1 \mp e^{-\alpha-x}), \ \Gamma(n+\frac{1}{2}) = (n-\frac{1}{2}) \cdots \frac{1}{2} \sqrt{\pi}$$

•
$$N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha}).$$

$$\frac{1}{e^{\alpha+x}\pm 1} = e^{-\alpha-x} (1 \mp e^{-\alpha-x}), \ \Gamma(n+\frac{1}{2}) = (n-\frac{1}{2}) \cdots \frac{1}{2} \sqrt{\pi}$$

•
$$N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha}).$$

•
$$U = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty \frac{x^{\frac{3}{2}} dx}{e^{\alpha + x} \pm 1}$$

$$\frac{1}{e^{\alpha+x}\pm 1} = e^{-\alpha-x} (1 \mp e^{-\alpha-x}), \ \Gamma(n+\frac{1}{2}) = (n-\frac{1}{2}) \cdots \frac{1}{2} \sqrt{\pi}$$

•
$$N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha}).$$

•
$$U = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty \frac{x^{\frac{3}{2}} dx}{e^{\alpha + x} \pm 1}$$

 $\simeq g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{3}{2}} dx$

$$\frac{1}{e^{\alpha+x}\pm 1}=e^{-\alpha-x}(1\mp e^{-\alpha-x}), \ \Gamma(n+\frac{1}{2})=(n-\frac{1}{2})\cdots\frac{1}{2}\sqrt{\pi}$$

•
$$N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha}).$$

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$$U = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty \frac{x^{\frac{3}{2}} dx}{e^{\alpha + x} \pm 1}$$

 $\simeq g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{3}{2}} dx$
 $= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT e^{-\alpha} [\Gamma(\frac{5}{2}) \mp e^{-\alpha} 2^{-\frac{5}{2}} \Gamma(\frac{5}{2})]$

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$$\frac{1}{e^{\alpha+x}\pm 1}=e^{-\alpha-x}(1\mp e^{-\alpha-x}), \ \Gamma(n+\frac{1}{2})=(n-\frac{1}{2})\cdots\frac{1}{2}\sqrt{\pi}$$

•
$$N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}Ve^{-\alpha}(1 \mp 2^{-\frac{3}{2}}e^{-\alpha}).$$

•
$$U = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty \frac{x^{\frac{3}{2}} dx}{e^{\alpha + x} \pm 1}$$

 $\simeq g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{3}{2}} dx$
 $= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT e^{-\alpha} [\Gamma(\frac{5}{2}) \mp e^{-\alpha} 2^{-\frac{5}{2}} \Gamma(\frac{5}{2})]$
 $= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT e^{-\alpha} \frac{3}{2} \frac{\sqrt{\pi}}{2} (1 \mp e^{-\alpha} 2^{-\frac{5}{2}})$

$$\frac{1}{e^{\alpha+x}\pm 1} = e^{-\alpha-x} (1 \mp e^{-\alpha-x}), \ \Gamma(n+\frac{1}{2}) = (n-\frac{1}{2}) \cdots \frac{1}{2} \sqrt{\pi}$$

•
$$N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha}).$$

•
$$U = g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty \frac{x^{\frac{3}{2}} dx}{e^{\alpha + x} \pm 1}$$

 $\simeq g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{3}{2}} dx$
 $= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT e^{-\alpha} \left[\Gamma(\frac{5}{2}) \mp e^{-\alpha} 2^{-\frac{5}{2}} \Gamma(\frac{5}{2})\right]$
 $= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT e^{-\alpha} \frac{3}{2} \frac{\sqrt{\pi}}{2} (1 \mp e^{-\alpha} 2^{-\frac{5}{2}})$
 $= g (\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V \frac{3}{2} kT e^{-\alpha} (1 \mp 2^{-\frac{5}{2}} e^{-\alpha}).$

§8. Bose statistics and Fermi statistics 8.2
$$U$$
 of weak degenerated ideal Fermion(Boson) gas
$$\frac{1}{e^{\alpha+x}+1}=e^{-\alpha-x}(1\mp e^{-\alpha-x}), \ \Gamma(n+\frac{1}{2})=(n-\frac{1}{2})\cdots\frac{1}{2}\sqrt{\pi}$$

•
$$N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}Ve^{-\alpha}(1 \mp 2^{-\frac{3}{2}}e^{-\alpha}).$$

• $U = g\frac{2\pi V}{h^3}(2mkT)^{\frac{3}{2}}kT\int_0^\infty \frac{x^{\frac{3}{2}}dx}{e^{\alpha+x}\pm 1}$

$$\simeq g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{3}{2}} dx$$

$$= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT e^{-\alpha} \left[\Gamma(\frac{5}{2}) \mp e^{-\alpha} 2^{-\frac{5}{2}} \Gamma(\frac{5}{2})\right]$$

$$= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT e^{-\alpha} \frac{3}{2} \frac{\sqrt{\pi}}{2} (1 \mp e^{-\alpha} 2^{-\frac{5}{2}})$$

$$= g (\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V^{\frac{3}{2}} kT e^{-\alpha} (1 \mp 2^{-\frac{5}{2}} e^{-\alpha}).$$

$$U = g (\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V^{\frac{3}{2}} kT e^{-\alpha} (1 \mp 2^{-\frac{5}{2}} e^{-\alpha})$$

•
$$\frac{U}{N} = \frac{g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}V^{\frac{3}{2}}kTe^{-\alpha}(1\mp 2^{-\frac{5}{2}}e^{-\alpha})}{g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}Ve^{-\alpha}(1\mp 2^{-\frac{3}{2}}e^{-\alpha})}$$

$$\frac{1}{e^{\alpha+x}\pm 1} = e^{-\alpha-x} (1 \mp e^{-\alpha-x}), \ \Gamma(n+\frac{1}{2}) = (n-\frac{1}{2}) \cdots \frac{1}{2} \sqrt{\pi}$$

$$\begin{split} \bullet \ N &= g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}Ve^{-\alpha}(1\mp 2^{-\frac{3}{2}}e^{-\alpha}). \\ \bullet \ U &= g\frac{2\pi V}{h^3}(2mkT)^{\frac{3}{2}}kT\int_0^\infty \frac{x^{\frac{3}{2}}\mathrm{d}x}{e^{\alpha+x}\pm 1} \\ &\simeq g\frac{2\pi V}{h^3}(2mkT)^{\frac{3}{2}}kT\int_0^\infty e^{-\alpha-x}(1\mp e^{-\alpha-x})x^{\frac{3}{2}}\mathrm{d}x \\ &= g\frac{2\pi V}{h^3}(2mkT)^{\frac{3}{2}}kTe^{-\alpha}[\Gamma(\frac{5}{2})\mp e^{-\alpha}2^{-\frac{5}{2}}\Gamma(\frac{5}{2})] \\ &= g\frac{2\pi V}{h^3}(2mkT)^{\frac{3}{2}}kTe^{-\alpha}\frac{3}{2}\frac{\sqrt{\pi}}{2}(1\mp e^{-\alpha}2^{-\frac{5}{2}}) \\ &= g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}V^{\frac{3}{2}}kTe^{-\alpha}(1\mp 2^{-\frac{5}{2}}e^{-\alpha}). \\ \bullet \ \frac{U}{N} &= \frac{g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}V^{\frac{3}{2}}kTe^{-\alpha}(1\mp 2^{-\frac{5}{2}}e^{-\alpha})}{g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}V^{\frac{3}{2}}kTe^{-\alpha}(1\mp 2^{-\frac{3}{2}}e^{-\alpha})} = \frac{3}{2}kT\frac{1\mp 2^{-\frac{5}{2}}e^{-\alpha}}{1\mp 2^{-\frac{3}{2}}e^{-\alpha}} \end{split}$$

$$\frac{1}{e^{\alpha+x}\pm 1} = e^{-\alpha-x} (1 \mp e^{-\alpha-x}), \ \Gamma(n+\frac{1}{2}) = (n-\frac{1}{2}) \cdots \frac{1}{2} \sqrt{\pi}$$

•
$$N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha}).$$

$$\begin{split} \bullet \ U &= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty \frac{x^{\frac{3}{2}} dx}{e^{\alpha + x} \pm 1} \\ &\simeq g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{3}{2}} dx \\ &= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT e^{-\alpha} \left[\Gamma(\frac{5}{2}) \mp e^{-\alpha} 2^{-\frac{5}{2}} \Gamma(\frac{5}{2})\right] \\ &= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT e^{-\alpha} \frac{3}{2} \frac{\sqrt{\pi}}{2} (1 \mp e^{-\alpha} 2^{-\frac{5}{2}}) \\ &= g (\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V^{\frac{3}{2}} kT e^{-\alpha} (1 \mp 2^{-\frac{5}{2}} e^{-\alpha}). \\ &U = g (\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V^{\frac{3}{2}} kT e^{-\alpha} (1 \mp 2^{-\frac{5}{2}} e^{-\alpha}) - 3kT^{\frac{1}{2}} T^{\frac{5}{2}} e^{-\alpha} \end{split}$$

•
$$\frac{U}{N} = \frac{g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}V^{\frac{3}{2}}kTe^{-\alpha}(1\mp 2^{-\frac{5}{2}}e^{-\alpha})}{g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}Ve^{-\alpha}(1\mp 2^{-\frac{3}{2}}e^{-\alpha})} = \frac{3}{2}kT\frac{1\mp 2^{-\frac{5}{2}}e^{-\alpha}}{1\mp 2^{-\frac{3}{2}}e^{-\alpha}}$$

$$\simeq \frac{3}{2}kT(1\mp 2^{-\frac{5}{2}}e^{-\alpha})(1\pm 2^{-\frac{3}{2}}e^{-\alpha})$$

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$$\frac{1}{e^{\alpha+x}\pm 1}=e^{-\alpha-x}(1\mp e^{-\alpha-x}), \ \Gamma(n+\frac{1}{2})=(n-\frac{1}{2})\cdots\frac{1}{2}\sqrt{\pi}$$

•
$$N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha}).$$

$$\begin{split} \bullet \ U &= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty \frac{x^{\frac{7}{2}} dx}{e^{\alpha + x} \pm 1} \\ &\simeq g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT \int_0^\infty e^{-\alpha - x} (1 \mp e^{-\alpha - x}) x^{\frac{3}{2}} dx \\ &= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT e^{-\alpha} \big[\Gamma(\frac{5}{2}) \mp e^{-\alpha} 2^{-\frac{5}{2}} \Gamma(\frac{5}{2}) \big] \\ &= g \frac{2\pi V}{h^3} (2mkT)^{\frac{3}{2}} kT e^{-\alpha} \frac{3}{2} \frac{\sqrt{\pi}}{2} (1 \mp e^{-\alpha} 2^{-\frac{5}{2}}) \\ &= g (\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V^{\frac{3}{2}} kT e^{-\alpha} (1 \mp 2^{-\frac{5}{2}} e^{-\alpha}). \\ \bullet \ \frac{U}{N} &= \frac{g (\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V^{\frac{3}{2}} kT e^{-\alpha} (1 \mp 2^{-\frac{5}{2}} e^{-\alpha})}{g (\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha})} = \frac{3}{2} kT \frac{1 \mp 2^{-\frac{5}{2}} e^{-\alpha}}{1 \mp 2^{-\frac{3}{2}} e^{-\alpha}} \end{split}$$

$$\frac{1}{N} = \frac{1}{g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1\mp 2^{-\frac{3}{2}} e^{-\alpha})} = \frac{1}{2} kT \frac{1}{1\mp 2^{-\frac{3}{2}} e^{-\alpha}}
\simeq \frac{3}{2} kT (1\mp 2^{-\frac{5}{2}} e^{-\alpha}) (1\pm 2^{-\frac{3}{2}} e^{-\alpha})
= \frac{3}{2} kT [1+(\pm 2^{-\frac{3}{2}}\mp 2^{-\frac{5}{2}}) e^{-\alpha} - 2^{-4} e^{-2\alpha}]$$

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§8. Bose statistics and Fermi statistics 8.2
$$U$$
 of weak degenerated ideal Fermion(Boson) gas $\frac{1}{e^{\alpha+x}+1}=e^{-\alpha-x}(1\mp e^{-\alpha-x})$, $\Gamma(n+\frac{1}{2})=(n-\frac{1}{2})\cdots\frac{1}{2}\sqrt{\pi}$

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$$N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha}).$$

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$$g(\frac{2\pi mkT}{h^2})^{\frac{1}{2}V}e^{-\alpha}(1\mp 2^{-\frac{1}{2}}e^{-\alpha}) \qquad 1\mp 2^{-\frac{1}{2}}e^{-\alpha}$$

$$\simeq \frac{3}{2}kT(1\mp 2^{-\frac{5}{2}}e^{-\alpha})(1\pm 2^{-\frac{3}{2}}e^{-\alpha})$$

$$= \frac{3}{2}kT[1+(\pm 2^{-\frac{3}{2}}\mp 2^{-\frac{5}{2}})e^{-\alpha}-2^{-4}e^{-2\alpha}]$$

$$\simeq \frac{3}{2}kT(1\pm 2^{-\frac{5}{2}}e^{-\alpha}).$$

•
$$\frac{U}{N} \simeq \frac{3}{2}kT(1 \pm 2^{-\frac{5}{2}}e^{-\alpha}).$$

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- $\frac{U}{N} \simeq \frac{3}{2}kT(1 \pm 2^{-\frac{5}{2}}e^{-\alpha}).$
- For small $e^{-\alpha}$, $N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}Ve^{-\alpha}(1 \mp 2^{-\frac{3}{2}}e^{-\alpha})$

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- $\frac{U}{N} \simeq \frac{3}{2}kT(1 \pm 2^{-\frac{5}{2}}e^{-\alpha}).$
- $\begin{array}{l} \bullet \text{ For small } e^{-\alpha}, \ N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha}) \\ \simeq g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha}, \\ \Rightarrow e^{-\alpha} \simeq \frac{N}{V} (\frac{h^2}{2\pi mkT})^{\frac{3}{2}} g^{-1}. \end{array}$

- $\frac{U}{N} \simeq \frac{3}{2}kT(1 \pm 2^{-\frac{5}{2}}e^{-\alpha}).$
- For small $e^{-\alpha}$, $N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}Ve^{-\alpha}(1 \mp 2^{-\frac{3}{2}}e^{-\alpha})$ $\simeq q(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}Ve^{-\alpha}$ $\Rightarrow e^{-\alpha} \simeq \frac{N}{V} \left(\frac{h^2}{2\pi m kT}\right)^{\frac{3}{2}} g^{-1}.$
- $U \simeq N_{\frac{3}{2}}kT(1 \pm 2^{-\frac{5}{2}}e^{-\alpha})$

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U of weak degenerated ideal Fermion(Boson) gas

- $\frac{U}{N} \simeq \frac{3}{2}kT(1 \pm 2^{-\frac{5}{2}}e^{-\alpha}).$
- $$\begin{split} \text{For small } e^{-\alpha}, \ N &= g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha}) \\ &\simeq g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha}, \\ &\Rightarrow e^{-\alpha} \simeq \frac{N}{V} (\frac{h^2}{2\pi mkT})^{\frac{3}{2}} g^{-1}. \end{split}$$
- $U \simeq N \frac{3}{2} kT (1 \pm 2^{-\frac{5}{2}} e^{-\alpha})$ $\simeq \frac{3}{2} N kT [1 \pm 2^{-\frac{5}{2}} \frac{N}{V} (\frac{h^2}{2\pi m kT})^{\frac{3}{2}} g^{-1}]$

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U of weak degenerated ideal Fermion(Boson) gas

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- For small $e^{-\alpha}$, $N = g(\frac{2\pi mkT}{h^2})^{\frac{3}{2}} V e^{-\alpha} (1 \mp 2^{-\frac{3}{2}} e^{-\alpha})$ $\simeq q(\frac{2\pi mkT}{h^2})^{\frac{3}{2}}Ve^{-\alpha}$ $\Rightarrow e^{-\alpha} \simeq \frac{N}{V} \left(\frac{h^2}{2\pi m^k T}\right)^{\frac{3}{2}} g^{-1}.$
- $U \simeq N\frac{3}{2}kT(1 \pm 2^{-\frac{5}{2}}e^{-\alpha})$ $\simeq \frac{3}{2}NkT[1\pm 2^{-\frac{5}{2}}\frac{N}{V}(\frac{h^2}{2\pi mkT})^{\frac{3}{2}}g^{-1}]$ $=\frac{3}{2}NkT(1\pm 2^{-\frac{5}{2}}g^{-1}n\lambda^3).$

1st term, same as the Boltzmann distribution; 2nd term, modification by the identical principle of particles (Fermions, repel; Bosons, attract).

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 - ullet 8.2 U of weak degenerated ideal Fermion(Boson) gas
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 - 8.4 Photon gas
 - 8.5 Free electron gas in metal

• If even $n\lambda^3 < 1$ does not satisfy, for Bosons, Bose-Einstein Condensate.

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- $a_l = \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} 1} = \frac{\omega_l}{e^{\frac{\varepsilon_l \mu}{kT}} 1}$.

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$$a_l = \frac{\omega_l}{e^{\alpha + \beta \varepsilon_l} - 1} = \frac{\omega_l}{e^{\frac{\varepsilon_l - \mu}{kT}} - 1}$$
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- $n = \frac{1}{V} \sum a_l = \frac{1}{V} \sum \frac{\omega_l}{e^{\frac{\varepsilon_l \mu}{kT} 1}}$.

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- $a_l > 0$ requires: $e^{\frac{\varepsilon_l \mu}{kT}} 1 > 0 \Rightarrow \varepsilon_l > \mu$, so $\varepsilon_0 > \mu$.
- For ideal gas, $\varepsilon_0 = 0$, $\Rightarrow \mu < 0$.
- $n = \frac{1}{V} \sum a_l = \frac{1}{V} \sum_{\substack{\frac{\varepsilon_l \mu}{kT} 1}} \frac{\omega_l}{1}$. $T \downarrow$, $\mu \uparrow$, until 0. Corresponds to a critical temperature: T_c .

December 30, 2013

$$\bullet \ n = \frac{1}{V} \sum_{\substack{\frac{\varepsilon_l - 0}{e^{\frac{\varepsilon_l - 0}{kT_c}} - 1}}$$

•
$$n = \frac{1}{V} \sum_{\substack{\frac{\varepsilon_l - 0}{e^{\frac{1}{kT_c}} - 1}}} = \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\frac{\varepsilon}{kT_c}} - 1}$$

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•
$$n = \frac{1}{V} \sum_{\substack{\frac{\omega_l}{e^{\frac{1}{kT_c}} - 1}}} = \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\frac{\varepsilon}{kT_c}} - 1}$$

= $\frac{2\pi}{h^3} (2m)^{\frac{3}{2}} (kT_c)^{\frac{3}{2}} \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^{x} - 1}$

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 $= \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} (kT_c)^{\frac{3}{2}} \int_0^\infty \frac{x^{\frac{1}{2}} dx}{e^x - 1}$
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$$n = \frac{1}{V} \sum_{\substack{\frac{\omega_l}{e^{\frac{1}{kT_c}} - 1}}} = \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\frac{\varepsilon}{kT_c}} - 1}$$

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• Notice $n\lambda^3 = n(\frac{h^2}{2\pi mkT})^{\frac{3}{2}}$,

•
$$n = \frac{1}{V} \sum_{\substack{\frac{\omega_l}{e^{\frac{1}{2}-0}} = \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} \int_0^{\infty} \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\frac{\varepsilon}{kT_c}} - 1}}$$

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• Notice $n\lambda^3=n(\frac{h^2}{2\pi mkT})^{\frac{3}{2}}$, $n\lambda_c^3=n(\frac{h^2}{2\pi mkT_c})^{\frac{3}{2}}$

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$$n\lambda_c^3 = n(\frac{h^2}{2\pi mkT_c})^{\frac{3}{2}} = n(\frac{h^2}{2\pi mk})^{\frac{3}{2}} \frac{2.612(mk)^{\frac{3}{2}}}{(2\pi)^{\frac{3}{2}}h^3n}$$

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$$n = \frac{1}{V} \sum_{\substack{\frac{\omega_l}{e^{\frac{1}{kT_c}} - 1}}} = \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} \int_0^{\infty} \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\frac{\varepsilon}{kT_c}} - 1}$$

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•
$$n = \frac{1}{V} \sum_{\substack{\frac{\varepsilon_l - 0}{e^{\frac{1}{kT_c}} - 1}}} = \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\frac{\varepsilon}{kT_c}} - 1}$$

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- What's happening if $T < T_c$?

•
$$n = \frac{1}{V} \sum_{\substack{\frac{\varepsilon_l - 0}{e^{\frac{1}{kT_c}} - 1}}} = \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\frac{\varepsilon}{kT_c}} - 1}$$

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- What's happening if $T < T_c$?
- More and more particles stay at ground state.

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•
$$n = \frac{1}{V} \sum_{\substack{\frac{\omega_l}{e^{\frac{1}{k}T_c} - 1}}} = \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} \int_0^{\infty} \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{e^{\frac{\varepsilon}{k}T_c} - 1}$$

 $= \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} (kT_c)^{\frac{3}{2}} \int_0^{\infty} \frac{x^{\frac{1}{2}} dx}{e^{x} - 1}$
 $= \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} (kT_c)^{\frac{3}{2}} \frac{\sqrt{\pi}}{2} \cdot 2.612$
 $\Rightarrow T_c = \frac{2\pi}{2.612^{\frac{2}{3}}} \frac{\hbar^2}{mk} n^{\frac{2}{3}}.$

- Notice $n\lambda^3 = n(\frac{h^2}{2\pi m^{kT}})^{\frac{3}{2}}$, $n\lambda_c^3 = n(\frac{h^2}{2\pi mkT_c})^{\frac{3}{2}} = n(\frac{h^2}{2\pi mk})^{\frac{3}{2}} \frac{2.612(mk)^{\frac{2}{2}}}{(2\pi)^{\frac{3}{2}+2}} = 2.612.$
- What's happening if $T < T_c$?
- More and more particles stay at ground state. $\mu < 0$ still obeys, $\mu \to 0$;

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- What's happening if $T < T_c$?
- More and more particles stay at ground state. $\mu < 0$ still obeys, $\mu \to 0$; $\frac{2\pi}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} d\varepsilon}{\sqrt{\varepsilon}} < n$.

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, where $n_{\varepsilon > 0} = \frac{2\pi}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}} \mathrm{d}\varepsilon}{e^{\frac{\varepsilon}{kT}} - 1}$

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There is a break for C_V at $T = T_c$. (Jump on derivative of C_V , continuous phase transition.)

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- ⁴He at 3.13K, transfer to superfluid.
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- Magneto optical trap 3-D harmonic oscillator.

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$$\varepsilon_{n_x,n_y,n_z} = \hbar\omega_x(n_x + \frac{1}{2}) + \hbar\omega_y(n_y + \frac{1}{2}) + \hbar\omega_z(n_z + \frac{1}{2}).$$

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In case of $\frac{\hbar\omega}{kT}\ll 1$, summation to integration:

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In case of $\frac{\hbar\omega}{kT}\ll 1$, summation to integration:

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 $= (\frac{kT_c}{\hbar})^3(\omega_x\omega_y\omega_z)^{-1} \iiint \frac{1}{e^{x+y+z}-1} dx dy dz$

$$\begin{split} \bullet \ N &= \iiint \frac{1}{e^{\frac{\hbar}{kT_c}(n_x\omega_x + n_y\omega_y + n_z\omega_z)} - 1} \mathrm{d}n_x \mathrm{d}n_y \mathrm{d}n_z \\ &= \iiint \frac{[(\frac{\hbar}{kT_c})^3\omega_x\omega_y\omega_z]^{-1}}{e^{\frac{\hbar}{kT}(n_x\omega_x + n_y\omega_y + n_z\omega_z)} - 1} \mathrm{d}\frac{\hbar\omega_x}{kT_c} n_x \mathrm{d}\frac{\hbar\omega_y}{kT_c} n_y \mathrm{d}\frac{\hbar\omega_z}{kT_c} n_z \\ &= (\frac{kT_c}{\hbar})^3(\omega_x\omega_y\omega_z)^{-1} \iiint \frac{1}{e^{x+y+z} - 1} \mathrm{d}x \mathrm{d}y \mathrm{d}z \\ &\quad \left(\mathsf{define} \ \bar{\omega} = (\omega_x\omega_y\omega_z)^{\frac{1}{3}} \right) \end{split}$$

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$$N = \iiint \frac{1}{e^{\frac{\hbar}{kT_c}(n_x\omega_x + n_y\omega_y + n_z\omega_z)} - 1} dn_x dn_y dn_z$$

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•
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- Convert to frequency $\hbar\omega=\varepsilon=pc$, $\Rightarrow p=\frac{\hbar\omega}{c}$, $D(\omega)d\omega = \frac{8\pi V(\frac{\hbar\omega}{c})^2 d(\frac{\hbar\omega}{c})}{h^3} = \frac{V}{\pi^2 c^3} \omega^2 d\omega.$
- Number of photons at $(\omega, \omega + d\omega)$: $a(\omega)d\omega = \frac{V}{\pi^2 c^3} \frac{\omega^2 d\omega}{c^{\frac{\hbar\omega}{LT}}}.$
- Internal energy: $U(\omega,T)d\omega = \frac{V}{\pi^2c^3}\frac{\hbar\omega^3d\omega}{\hbar\omega}$.
- Total internal energy: $U = \int_0^\infty U(\omega, T) d\omega$ $= \int_0^\infty \frac{V}{\pi^2 c^3} \frac{\hbar \omega^3 d\omega}{e^{\frac{\hbar \omega}{kT}} - 1} = \int_0^\infty \frac{V \hbar}{\pi^2 c^3} \left(\frac{\hbar}{kT}\right)^{-4} \frac{\left(\frac{\hbar \omega}{kT}\right)^3 d\left(\frac{\hbar \omega}{kT}\right)}{e^{\frac{\hbar \omega}{kT}} - 1}$ $=\frac{V(kT)^4}{\pi^2c^3\hbar^3}\int_0^\infty \frac{x^3\mathrm{d}x}{c^x-1} = \frac{\pi^2k^4}{15c^3\hbar^3}T^4V.$ $(a=\frac{\pi^2k^4}{15c^3\hbar^3}).$

• Maximum for $U(\omega,T)d\omega = \frac{V}{\pi^2c^3} \frac{\hbar\omega^3d\omega}{e^{\frac{\hbar\omega}{kT}}-1}$:



• Maximum for $U(\omega,T)d\omega = \frac{V}{\pi^2 c^3} \frac{\hbar \omega^3 d\omega}{e^{\frac{\hbar \omega}{kT}} - 1}$: $\frac{\mathrm{d}}{\mathrm{d}x}\left(\frac{x^3}{e^x-1}\right) = 0$

• Maximum for $U(\omega,T)\mathrm{d}\omega = \frac{V}{\pi^2c^3}\frac{\hbar\omega^3\mathrm{d}\omega}{e^{\frac{\hbar\omega}{kT}}-1}$: $\frac{\mathrm{d}}{\mathrm{d}x}(\frac{x^3}{e^x-1}) = 0 \Rightarrow x_m = \frac{\hbar\omega_m}{kT} \simeq 2.822$, Wien's displacement law.

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- Grand partition function for the photons:

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- Grand partition function for the photons:

$$\ln \mathcal{Z} = -\sum_{l} \omega_{l} \ln(1 - e^{-\alpha - \beta \varepsilon_{l}})$$
$$= -\frac{V}{\pi^{2} c^{3}} \int_{0}^{\infty} \omega^{2} \ln(1 - e^{-\beta \hbar \omega}) d\omega$$

Photon gas $D(\omega)\mathrm{d}\omega = \frac{V}{\pi^2c^3}\omega^2\mathrm{d}\omega$

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- Maximum for $U(\omega,T)d\omega = \frac{V}{\pi^2 c^3} \frac{\hbar \omega^3 d\omega}{e^{\frac{\hbar \omega}{4T}} 1}$: $\frac{\mathrm{d}}{\mathrm{d}x}(\frac{x^3}{e^x-1})=0 \Rightarrow x_m=\frac{\hbar\omega_m}{kT}\simeq 2.822$, Wien's displacement law.
- Grand partition function for the photons:

$$\ln \mathcal{Z} = -\sum_{\tau} \omega_{l} \ln(1 - e^{-\alpha - \beta \varepsilon_{l}})
= -\frac{V}{\pi^{2} c^{3}} \int_{0}^{\infty} \omega^{2} \ln(1 - e^{-\beta \hbar \omega}) d\omega
= -\frac{V}{\pi^{2} c^{3}} (\beta \hbar)^{-3} \int_{0}^{\infty} (\beta \hbar \omega)^{2} \ln(1 - e^{-\beta \hbar \omega}) d(\beta \hbar \omega)
= -\frac{V}{\pi^{2} c^{3}} (\beta \hbar)^{-3} \int_{0}^{\infty} x^{2} \ln(1 - e^{-x}) dx
= -\frac{V}{\pi^{2} c^{3}} \frac{1}{(\beta \hbar)^{3}} \int_{0}^{\infty} \ln(1 - e^{-x}) d\frac{x^{3}}{3}$$

Photon gas $D(\omega)\mathrm{d}\omega = \frac{V}{\pi^2c^3}\omega^2\mathrm{d}\omega$

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- Grand partition function for the photons:

$$\ln \mathcal{Z} = -\sum_{\eta} \omega_{l} \ln(1 - e^{-\alpha - \beta \varepsilon_{l}})$$

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Photon gas
$$D(\omega)d\omega = \frac{V}{\pi^2 c^3} \omega^2 d\omega$$

- Maximum for $U(\omega,T)d\omega = \frac{V}{\pi^2 c^3} \frac{\hbar \omega^3 d\omega}{e^{\frac{\hbar \omega}{4T}} 1}$: $\frac{\mathrm{d}}{\mathrm{d}x}(\frac{x^3}{e^x-1})=0 \Rightarrow x_m=\frac{\hbar\omega_m}{kT}\simeq 2.822$, Wien's displacement law.
- Grand partition function for the photons:

$$\ln \mathcal{Z} = -\sum_{\eta} \omega_l \ln(1 - e^{-\alpha - \beta \varepsilon_l})$$

$$= -\frac{V}{\pi^2 c^3} \int_0^{\infty} \omega^2 \ln(1 - e^{-\beta \hbar \omega}) d\omega$$

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$$\ln \mathcal{Z} = -\sum_{n} \omega_{l} \ln(1 - e^{-\alpha - \beta \varepsilon_{l}})$$

$$= -\frac{V}{\pi^{2}c^{3}} \int_{0}^{\infty} \omega^{2} \ln(1 - e^{-\beta \hbar \omega}) d\omega$$

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$$= -\frac{V}{\pi^{2}c^{3}} \frac{1}{(\beta \hbar)^{3}} \left[\frac{x^{3}}{3} \ln(1 - e^{-x})|_{0}^{\infty} - \int_{0}^{\infty} \frac{x^{3}}{3} \cdot \frac{e^{-x}}{1 - e^{-x}} dx\right]$$

$$= \frac{V}{\pi^{2}c^{3}} \frac{1}{(\beta \hbar)^{3}} \frac{1}{3} \int_{0}^{\infty} \frac{x^{3}}{e^{x} - 1} dx = \frac{\pi^{2}V}{45c^{3}} \frac{1}{(\beta \hbar)^{3}}.$$

$$\ln \mathcal{Z} = \frac{\pi^2 V}{45c^3} \frac{1}{(\beta\hbar)^3}$$
.

$$U = -\frac{\partial}{\partial \beta} \ln \mathcal{Z}$$

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$$U = -\frac{\partial}{\partial \beta} \ln \mathcal{Z} = -\frac{\pi^2 V}{45c^3} \frac{1}{\hbar^3} \frac{-3}{\beta^4}$$

§8. Bose statistics and Fermi statistics

Photon gas
$$U=rac{\pi^2k^4}{15c^3\hbar^3}T^4V$$
, $\ln\mathcal{Z}=rac{\pi^2V}{45c^3}rac{1}{(\beta\hbar)^3}$.

8.4 Photon gas

• Internal energy:

$$U = -\frac{\partial}{\partial \beta} \ln \mathcal{Z} = -\frac{\pi^2 V}{45c^3} \frac{1}{\hbar^3} \frac{-3}{\beta^4} = \frac{\pi^2 k^4}{15c^3 \hbar^3} T^4 V.$$

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$$U=rac{\pi^2 k^4}{15c^3\hbar^3}T^4V$$
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• Pressure:

$$p = \frac{1}{\beta} \frac{\partial}{\partial V} \ln \mathcal{Z}$$

$$U = -\frac{\partial}{\partial \beta} \ln \mathcal{Z} = -\frac{\pi^2 V}{45c^3} \frac{1}{\hbar^3} \frac{-3}{\beta^4} = \frac{\pi^2 k^4}{15c^3 \hbar^3} T^4 V.$$

• Pressure:

$$p = \frac{1}{\beta} \frac{\partial}{\partial V} \ln \mathcal{Z} = \frac{1}{\beta} \frac{\pi^2}{45c^3} \frac{1}{(\beta \hbar)^3}$$

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$$U=rac{\pi^2k^4}{15c^3\hbar^3}T^4V$$
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$$S = k(\ln \mathcal{Z} - \alpha \frac{\partial \ln \mathcal{Z}}{\partial \alpha} - \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta})$$

$$U = -\frac{\partial}{\partial \beta} \ln \mathcal{Z} = -\frac{\pi^2 V}{45c^3} \frac{1}{\hbar^3} \frac{-3}{\beta^4} = \frac{\pi^2 k^4}{15c^3 \hbar^3} T^4 V.$$

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$$S = k(\ln \mathcal{Z} - \alpha \frac{\partial \ln \mathcal{Z}}{\partial \alpha} - \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta}) = k(\ln \mathcal{Z} - \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta})$$

$$U = -\frac{\partial}{\partial \beta} \ln \mathcal{Z} = -\frac{\pi^2 V}{45c^3} \frac{1}{\hbar^3} \frac{-3}{\beta^4} = \frac{\pi^2 k^4}{15c^3 \hbar^3} T^4 V.$$

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$$= k\left[\frac{\pi^2 V}{45c^3} \frac{1}{(\beta \hbar)^3} + \beta \frac{\pi^2 k^4}{15c^3 \hbar^3} T^4 V\right]$$

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• Entropy:

$$S = k(\ln \mathcal{Z} - \alpha \frac{\partial \ln \mathcal{Z}}{\partial \alpha} - \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta}) = k(\ln \mathcal{Z} - \beta \frac{\partial \ln \mathcal{Z}}{\partial \beta})$$

$$= k \left[\frac{\pi^2 V}{45c^3} \frac{1}{(\beta \hbar)^3} + \beta \frac{\pi^2 k^4}{15c^3 \hbar^3} T^4 V \right]$$

$$= k \left(\frac{\pi^2 V}{45c^3} \frac{k^3 T^3}{\hbar^3} + \frac{\pi^2 V}{15c^3} \frac{k^3 T^3}{\hbar^3} \right)$$

$$= \frac{4}{45} \frac{\pi^2 k^4 V}{c^3 \hbar^3} T^3. \text{ (same as (2.6.4))}$$

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 - 8.4 Photon gas
 - 8.5 Free electron gas in metal

• Degenerated Fermi gas $(e^{\alpha} \gg 1 \text{ or } n\lambda^3 \gg 1)$.

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- Degenerated Fermi gas $(e^{\alpha} \gg 1 \text{ or } n\lambda^3 \gg 1)$.
- Problem: in the metal, electron's contribution on C_V is negligible.

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- Example: Copper (Cu) $ho = 8.9 imes 10^3 \mathrm{kg} \cdot \mathrm{m}^{-3}$, atomic weight: 63, number density $n_e \sim n_{\rm Cu} \simeq \frac{8.9 \times 10^3 {\rm kg \cdot m^{-3}}}{63 \times 1.67 \times 10^{-27} {\rm kg}} \simeq 8.5 \times 10^{28} {\rm m^{-3}}.$ $n\lambda^3 = n(\frac{h^2}{2\pi m kT})^{\frac{3}{2}}$ $\simeq 8.5 \times 10^{28} \mathrm{m}^{-3} \left(\frac{(6.626 \times 10^{-34} \mathrm{J \cdot s})^2}{2\pi \times 9.1 \times 10^{-31} \mathrm{kg} \cdot 1.38 \times 10^{-23} \mathrm{J \cdot K}^{-1} \cdot 300 \mathrm{K} \frac{T}{200 \mathrm{M}}} \right)^{\frac{3}{2}}$ $\simeq 3400(\frac{T}{300K})^{-\frac{3}{2}}$

December 30, 2013

- Degenerated Fermi gas $(e^{\alpha} \gg 1 \text{ or } n\lambda^3 \gg 1)$.
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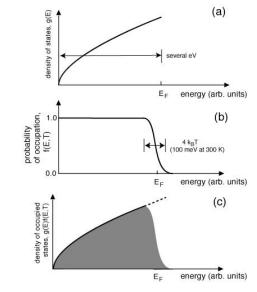
$$\bullet \ f \left\{ \begin{array}{l} > \frac{1}{2}, \ \varepsilon < \mu; \\ = \frac{1}{2}, \ \varepsilon = \mu; \\ < \frac{1}{2}, \ \varepsilon > \mu. \end{array} \right.$$

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$$R = \frac{4\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{3}{2}}}{e^{\frac{\varepsilon - \mu}{kT} + 1}} d\varepsilon,$$

$$E = \frac{4\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{3}{2}}}{e^{\frac{\varepsilon - \mu}{kT}} + 1} d\varepsilon.$$

• Estimate: only the electrons at $(\mu - kT, \mu + kT)$ contributes to the C_V .

$$N_{\rm eff} \simeq \frac{kT}{\mu} N$$
.

Acts like classical free particle, so $C_V \simeq \frac{3}{2} N_{\rm eff} k$

$$= \frac{3}{2} \frac{kT}{\mu} Nk \simeq \frac{3}{2} Nk \frac{T}{T_{\rm F}}.$$

Quantitative:

$$N = \frac{4\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}}}{e^{\frac{\varepsilon - \mu}{kT}} + 1} d\varepsilon,$$

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• Calculable in principle:

• Estimate: only the electrons at $(\mu - kT, \mu + kT)$ contributes to the C_V .

$$N_{\rm eff} \simeq \frac{kT}{\mu} N$$
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Acts like classical free particle, so $C_V \simeq \frac{3}{2} N_{\rm eff} k$ = $\frac{3}{2} \frac{kT}{\mu} N k \simeq \frac{3}{2} N k \frac{T}{T_{\rm E}}$.

• Quantitative:

$$N = \frac{4\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{1}{2}}}{e^{\frac{\varepsilon - \mu}{kT}} + 1} d\varepsilon,$$

$$E = \frac{4\pi V}{h^3} (2m)^{\frac{3}{2}} \int_0^\infty \frac{\varepsilon^{\frac{3}{2}}}{e^{\frac{\varepsilon - \mu}{kT}} + 1} d\varepsilon.$$

• Calculable in principle:

$$N = \frac{2}{3}C\mu^{\frac{3}{2}} \left[1 + \frac{\pi^{2}}{8} \left(\frac{kT}{\mu}\right)^{2}\right],$$

$$E = \frac{2}{5}C\mu^{\frac{5}{2}} \left[1 + \frac{5\pi^{2}}{8} \left(\frac{kT}{\mu}\right)^{2}\right], \text{ where } C = \frac{4\pi V}{h^{3}} (2m)^{\frac{3}{2}}.$$

88. Bose statistics and Fermi statistics 8.5 Free electron gas in metal

$$N = \frac{2}{3}C\mu^{\frac{3}{2}}\left[1 + \frac{\pi^2}{8}(\frac{kT}{\mu})^2\right], E = \frac{2}{5}C\mu^{\frac{5}{2}}\left[1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2\right]$$

• As $kT \ll \mu$, from N,

88. Bose statistics and Fermi statistics 8.5 Free electron gas in metal

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§8. Bose statistics and Fermi statistics 8.5 Free electron gas in metal

$$N = \frac{2}{3}C\mu^{\frac{3}{2}}\left[1 + \frac{\pi^2}{8}(\frac{kT}{\mu})^2\right], E = \frac{2}{5}C\mu^{\frac{5}{2}}\left[1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2\right]$$

• As $kT \ll \mu$, from N, $\mu(0)^{\frac{3}{2}} \simeq \frac{3N}{2C} = \mu^{\frac{3}{2}} [1 + \frac{\pi^2}{8} (\frac{kT}{\mu})^2]$

88. Bose statistics and Fermi statistics

$$N = \frac{2}{3}C\mu^{\frac{3}{2}}\left[1 + \frac{\pi^2}{8}(\frac{kT}{\mu})^2\right], E = \frac{2}{5}C\mu^{\frac{5}{2}}\left[1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2\right]$$

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88. Bose statistics and Fermi statistics

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§8. Bose statistics and Fermi statistics 8.5 Free electron gas in meta

$$N = \frac{2}{3}C\mu^{\frac{3}{2}} \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 \right], E = \frac{2}{5}C\mu^{\frac{5}{2}} \left[1 + \frac{5\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 \right]$$

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 $\simeq \mu^{\frac{3}{2}} [1 + \frac{\pi^2}{8} (\frac{kT}{\mu(0)})^2] \Rightarrow \mu(0) \simeq \mu [1 + \frac{\pi^2}{8} (\frac{kT}{\mu(0)})^2]^{\frac{2}{3}}$
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§8. Bose statistics and Fermi statistics 8.5 Free electron gas in metal

$$N = \frac{2}{3}C\mu^{\frac{3}{2}} \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 \right], E = \frac{2}{5}C\mu^{\frac{5}{2}} \left[1 + \frac{5\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 \right]$$

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$$N = \frac{2}{3}C\mu^{\frac{3}{2}} \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 \right], E = \frac{2}{5}C\mu^{\frac{5}{2}} \left[1 + \frac{5\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 \right]$$

$$\begin{array}{l} \bullet \text{ As } kT \ll \mu \text{, from } N \text{, } \mu(0)^{\frac{3}{2}} \simeq \frac{3N}{2C} = \mu^{\frac{3}{2}} \big[1 + \frac{\pi^2}{8} (\frac{kT}{\mu})^2 \big] \\ \simeq \mu^{\frac{3}{2}} \big[1 + \frac{\pi^2}{8} (\frac{kT}{\mu(0)})^2 \big] \Rightarrow \mu(0) \simeq \mu \big[1 + \frac{\pi^2}{8} (\frac{kT}{\mu(0)})^2 \big]^{\frac{2}{3}} \\ \Rightarrow \mu \simeq \mu(0) \big[1 + \frac{\pi^2}{8} (\frac{kT}{\mu(0)})^2 \big]^{-\frac{2}{3}} \simeq \mu(0) \big[1 - \frac{\pi^2}{12} (\frac{kT}{\mu(0)})^2 \big]. \end{array}$$

•
$$U = \frac{3}{5}N\mu \frac{1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2}{1 + \frac{\pi^2}{8}(\frac{kT}{\mu})^2}$$

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§8. Bose statistics and Fermi statistics 8.5 Free electron gas in metal

$$N = \frac{2}{3}C\mu^{\frac{3}{2}} \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 \right], E = \frac{2}{5}C\mu^{\frac{5}{2}} \left[1 + \frac{5\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 \right]$$

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$$U = \frac{3}{5}N\mu \frac{1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2}{1 + \frac{\pi^2}{2}(\frac{kT}{\mu})^2} \simeq \frac{3}{5}N\mu \left[1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2\right] \left[1 - \frac{\pi^2}{8}(\frac{kT}{\mu})^2\right]$$

8. Bose statistics and Fermi statistics 8.5 Free electron gas in metal

$$N = \frac{2}{3}C\mu^{\frac{3}{2}} \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu}\right)^2\right], E = \frac{2}{5}C\mu^{\frac{5}{2}} \left[1 + \frac{5\pi^2}{8} \left(\frac{kT}{\mu}\right)^2\right]$$

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$$U = \frac{3}{5}N\mu \frac{1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2}{1 + \frac{\pi^2}{8}(\frac{kT}{\mu})^2} \simeq \frac{3}{5}N\mu \left[1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2\right] \left[1 - \frac{\pi^2}{8}(\frac{kT}{\mu})^2\right]$$

 $\simeq \frac{3}{5}N\mu \left[1 + \frac{\pi^2}{2}(\frac{kT}{\mu})^2\right]$

8. Bose statistics and Fermi statistics 8.5 Free electron gas in metal

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$$U = \frac{3}{5}N\mu \frac{1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2}{1 + \frac{\pi^2}{8}(\frac{kT}{\mu})^2} \simeq \frac{3}{5}N\mu \left[1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2\right] \left[1 - \frac{\pi^2}{8}(\frac{kT}{\mu})^2\right]$$

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8. Bose statistics and Fermi statistics 8.5 Free electron gas in metal

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$$U = \frac{3}{5}N\mu \frac{1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2}{1 + \frac{\pi^2}{8}(\frac{kT}{\mu})^2} \simeq \frac{3}{5}N\mu \left[1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2\right] \left[1 - \frac{\pi^2}{8}(\frac{kT}{\mu})^2\right]$$

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 $\simeq \frac{3}{5}N\mu (0) \left[1 + \frac{5\pi^2}{12}(\frac{kT}{\mu(0)})^2\right].$

$$N = \frac{2}{3}C\mu^{\frac{3}{2}}\left[1 + \frac{\pi^2}{8}(\frac{kT}{\mu})^2\right], E = \frac{2}{5}C\mu^{\frac{5}{2}}\left[1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2\right]$$

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$$U = \frac{3}{5}N\mu \frac{1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2}{1 + \frac{\pi^2}{8}(\frac{kT}{\mu})^2} \simeq \frac{3}{5}N\mu \left[1 + \frac{5\pi^2}{8}(\frac{kT}{\mu})^2\right] \left[1 - \frac{\pi^2}{8}(\frac{kT}{\mu})^2\right]$$

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$$\simeq \frac{3}{5}N\mu (0) \left[1 - \frac{\pi^2}{12}(\frac{kT}{\mu(0)})^2\right] \left[1 + \frac{\pi^2}{2}(\frac{kT}{\mu(0)})^2\right]$$

$$\simeq \frac{3}{5}N\mu (0) \left[1 + \frac{5\pi^2}{12}(\frac{kT}{\mu(0)})^2\right].$$

• $C_V = (\frac{\partial U}{\partial T})_V = Nk \frac{\pi^2 k}{2\mu(0)} T = \frac{\pi^2}{2} Nk \frac{T}{T_E} \propto T$.

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