

# Dark Matter Particle Astronomy

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GRAPPA Institute - 10th October 2016



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## Observation of High-Energy Astrophysical Neutrinos in Three Years of IceCube Data

M. G. Aartsen,<sup>2</sup> M. Ackermann,<sup>45</sup> J. Adams,<sup>15</sup> J. A. Aguilar,<sup>23</sup> M. Ahlers,<sup>28</sup> M. Ahrens,<sup>36</sup> D. Altmann,<sup>22</sup> T. Anderson,<sup>42</sup>

## The Era of Neutrino Astronomy Has Begun

NOVEMBER 21, 2013

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Contacts: Heather Dewar 301-405-9267

Space.com &gt; Science &amp; Astronomy

## Neutrino Telescopes Launch New Era of Astronomy

By Tanya Lewis, Staff Writer | January 20, 2014 07:01am ET

  
150  
56

has been named the "2013 Breakthrough of the Year" by  
[Read more here.](#)

ng a telescope embedded in Antarctic ice have succeeded in a

## Viewpoint: The Beginning of Extra-Galactic Neutrino Astronomy

Eli Waxman, Particle Physics & Astrophysics Department, Weizmann Institute of Science, Israel

September 2, 2014 • *Physics* 7, 88

What can high-energy neutrinos tell us about astrophysical objects beyond our galaxy?



## Observation of Gravitational Waves from a Binary Black Hole Merger

B. P. Abbott *et al.*\*

(LIGO Scientific Collaboration and Virgo Collaboration)

(Received 21 January 2016; published 11 February 2016)

THE SCIENCES

# The Future of Gravitational Wave Astronomy

Fully opening this new window on the universe will take decades—even centuries

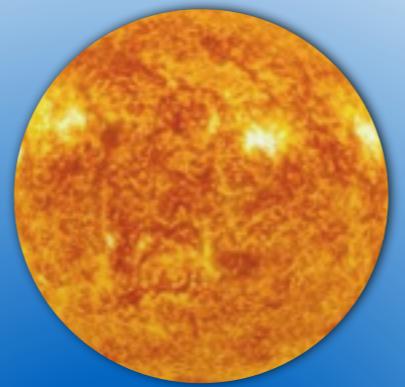
## Why gravitational wave astronomy has physicists so damn excited

Updated by Brian Resnick · @B\_resnick · brian@vox.com · Jun 21, 2016, 10:55a

## Second detection heralds the era of gravitational wave astronomy

June 17, 2016 by Paul Lasky, The Conversation





**NOT TO SCALE**

# Overview

Dark Matter (DM)

Direct detection of DM

Overcoming halo uncertainties in direct detection

BJK, Green [1207.2039, 1303.6868, 1312.1852]

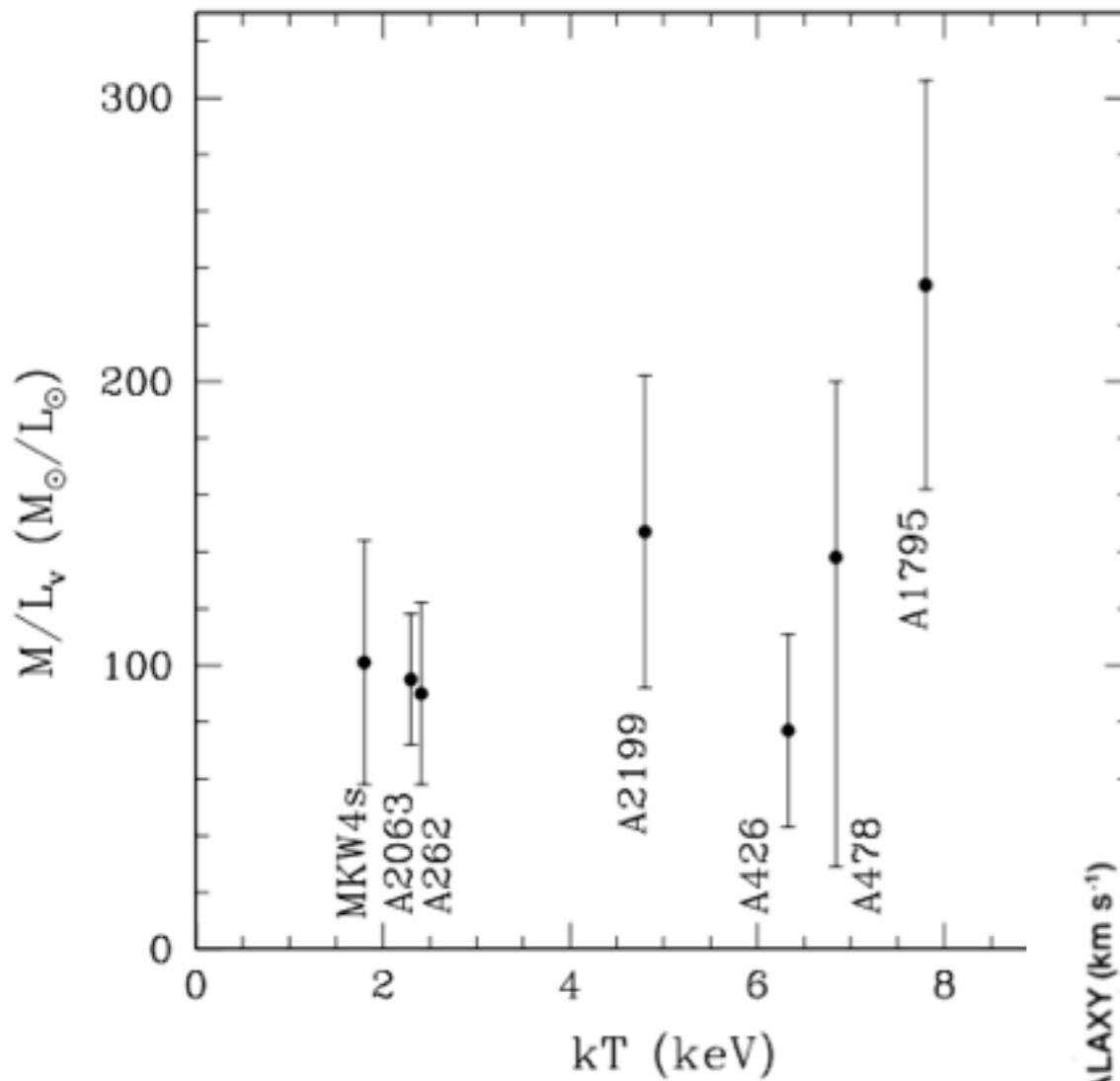
Probing low speed DM with neutrino telescopes

BJK, Fornasa, Green [1410.8051]

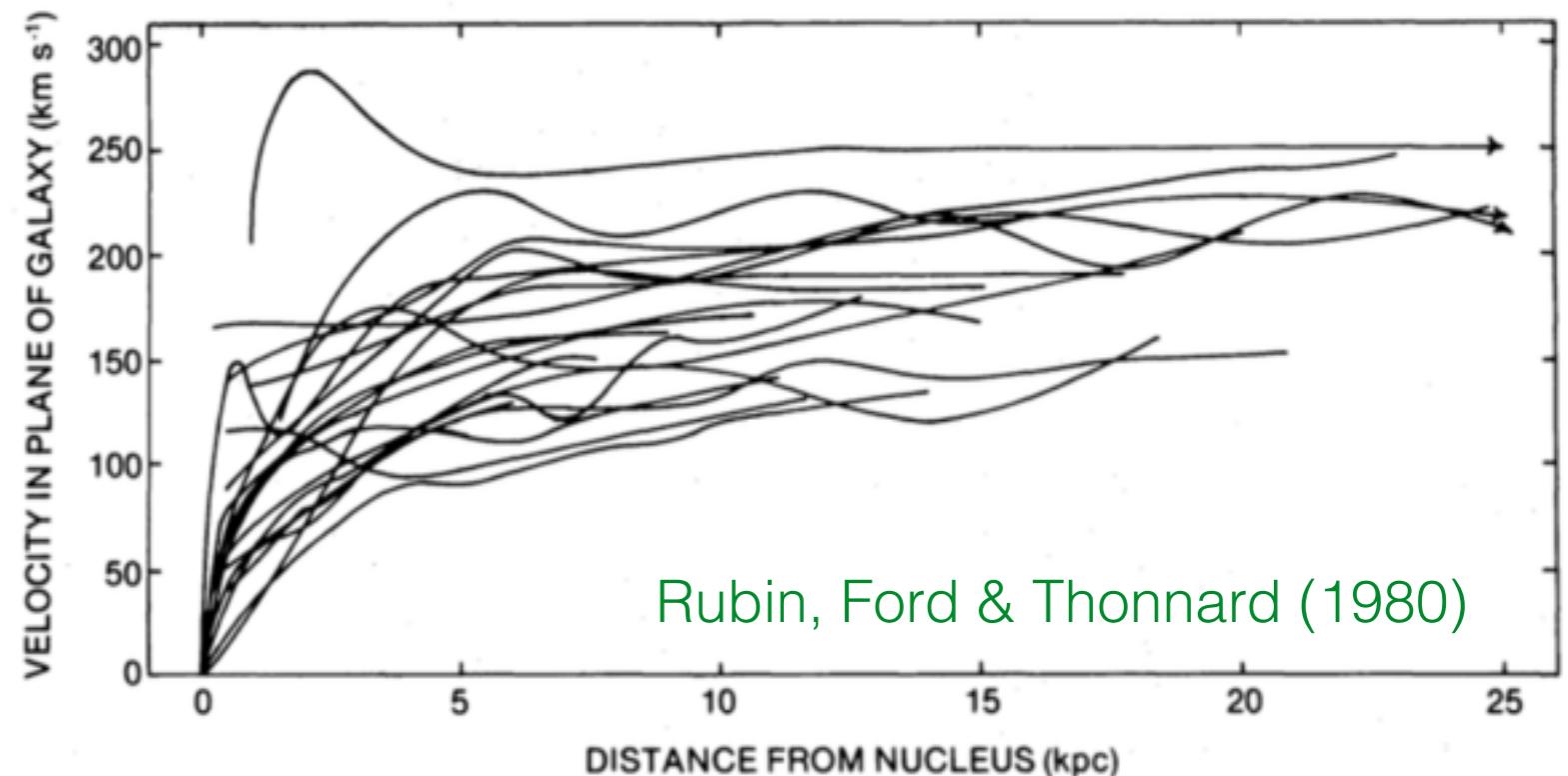
Measuring the DM velocity distribution  
with directional experiments

BJK [1502.04224]; BJK, O'Hare [1609.08630]

# Dark Matter



Hradecky et al. [astro-ph/0006397]



Rubin, Ford & Thonnard (1980)

# Dark Matter at the Sun's Radius

## Global

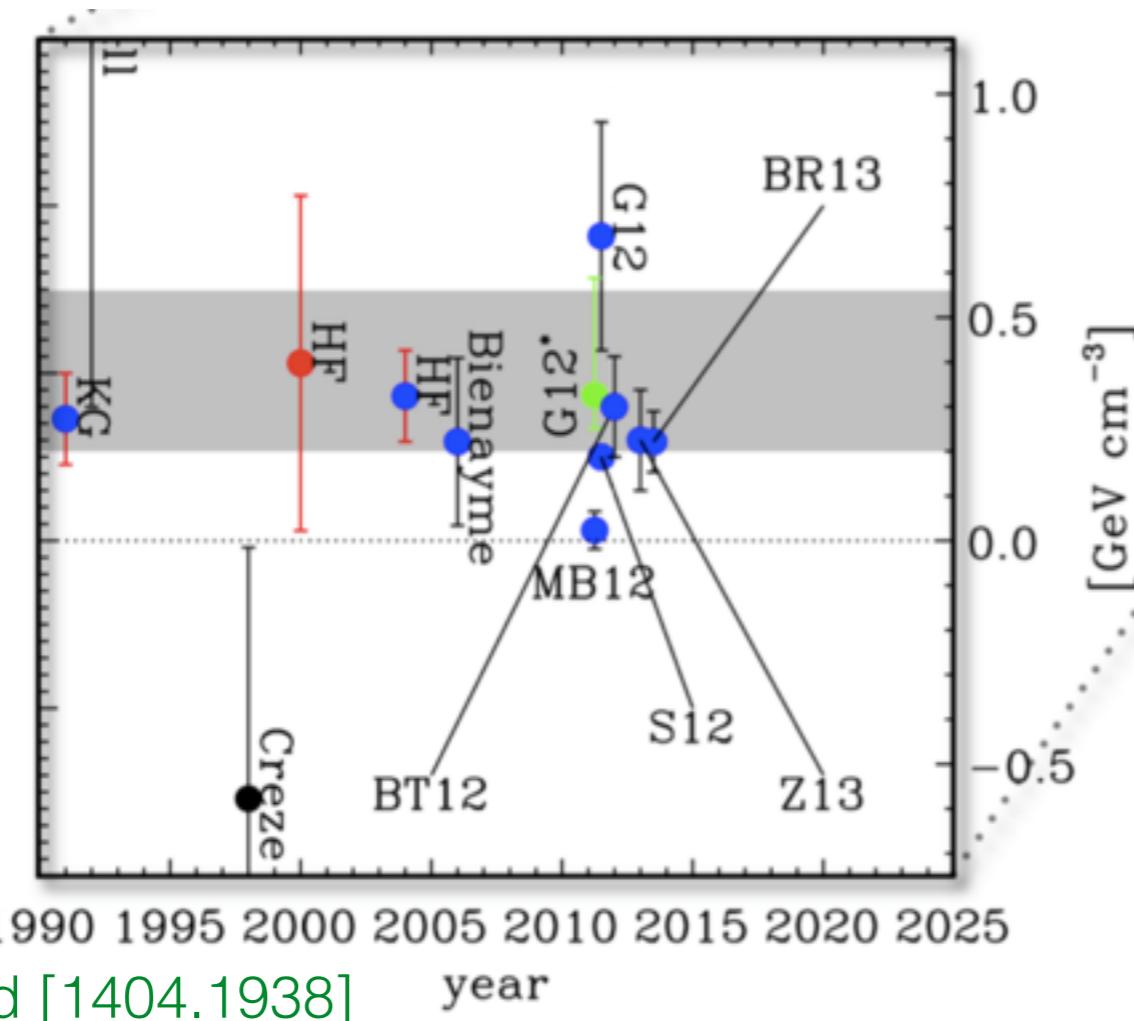
Model total mass distribution in Milky Way and extract DM density at Solar Radius ( $\sim 8$  kpc)

E.g. locco et al. [1502.03821]

## Local

Estimate local DM density from kinematics of local stars (assuming local disk equilibrium)

E.g. Garbari et al. [1206.0015]

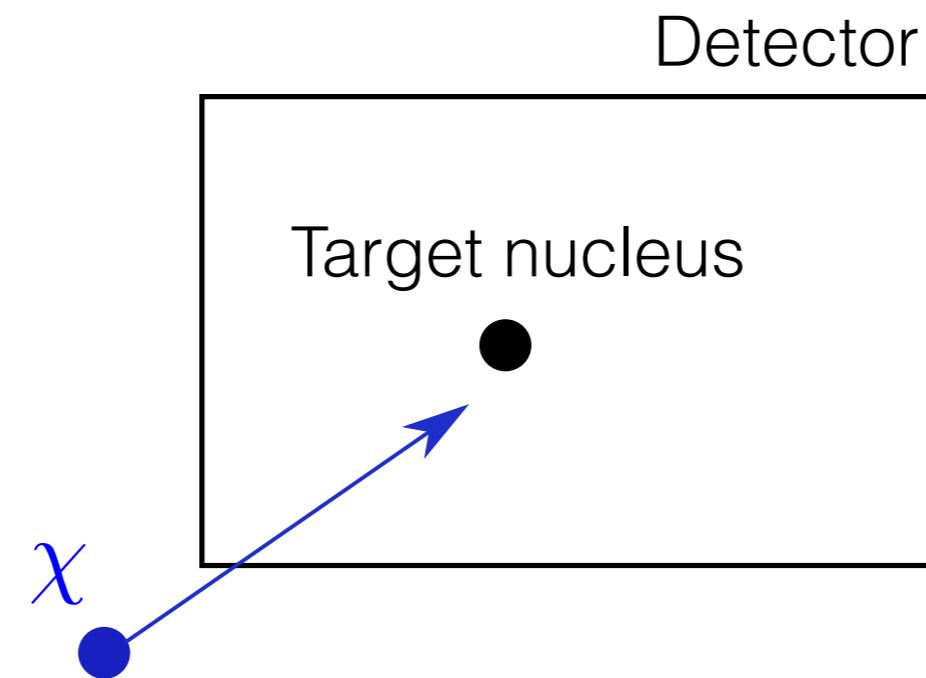


Values in the range:  
 $\rho_\chi \sim 0.2\text{--}0.8 \text{ GeV cm}^{-3}$

But **not** zero!  
c.f. Garbari et al. [1204.3924]

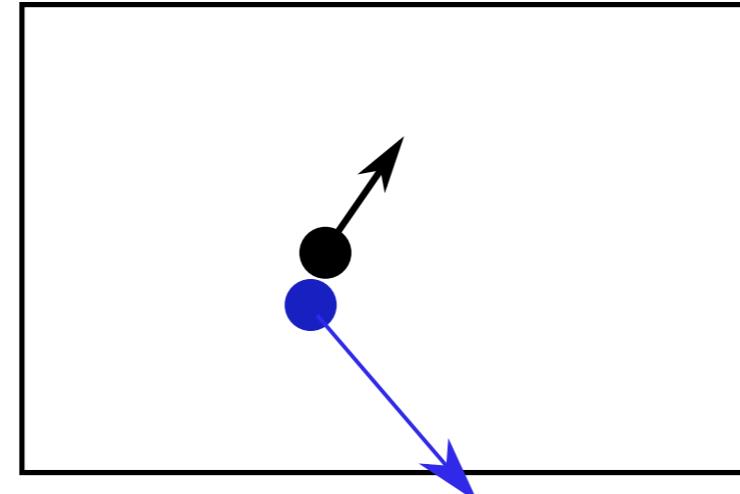
# Direct detection

$$m_\chi \gtrsim 1 \text{ GeV}$$
$$v \sim 10^{-3}$$



# Direct detection

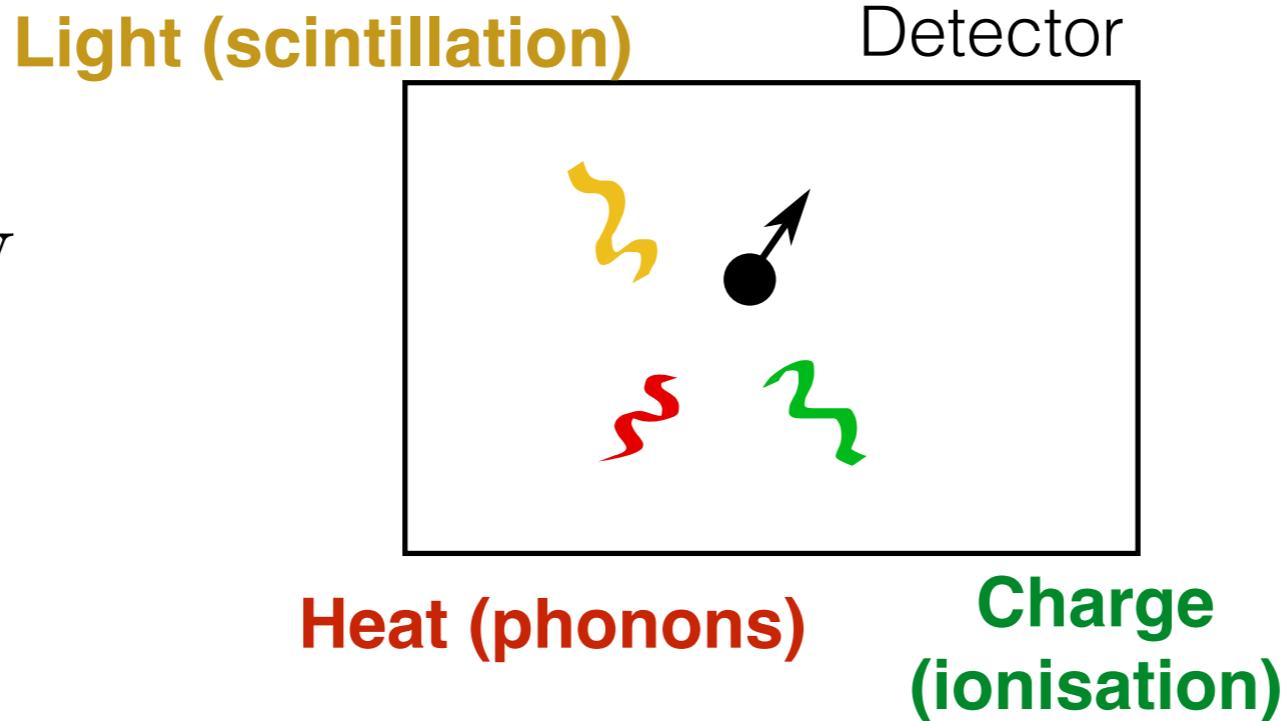
Detector



$$m_\chi \gtrsim 1 \text{ GeV}$$
$$v \sim 10^{-3}$$

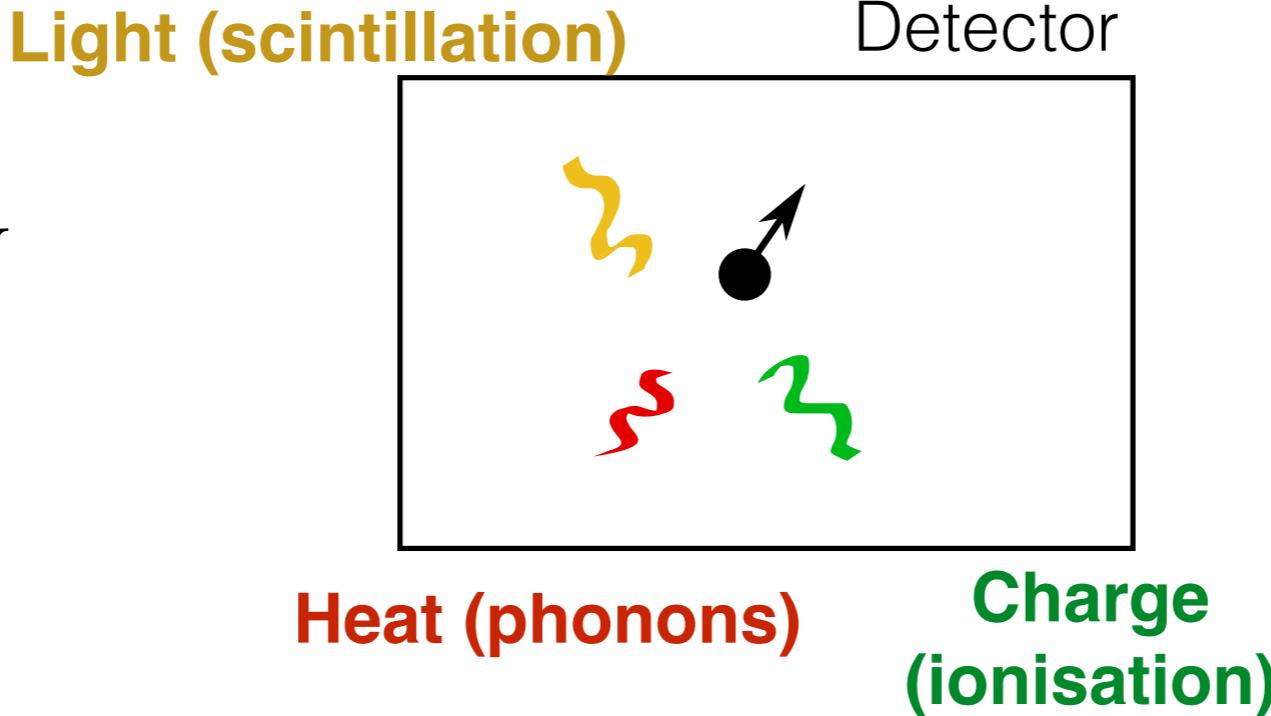
# Direct detection

$m_\chi \gtrsim 1 \text{ GeV}$   
 $v \sim 10^{-3}$



# Direct detection

$$m_\chi \gtrsim 1 \text{ GeV}$$
$$v \sim 10^{-3}$$



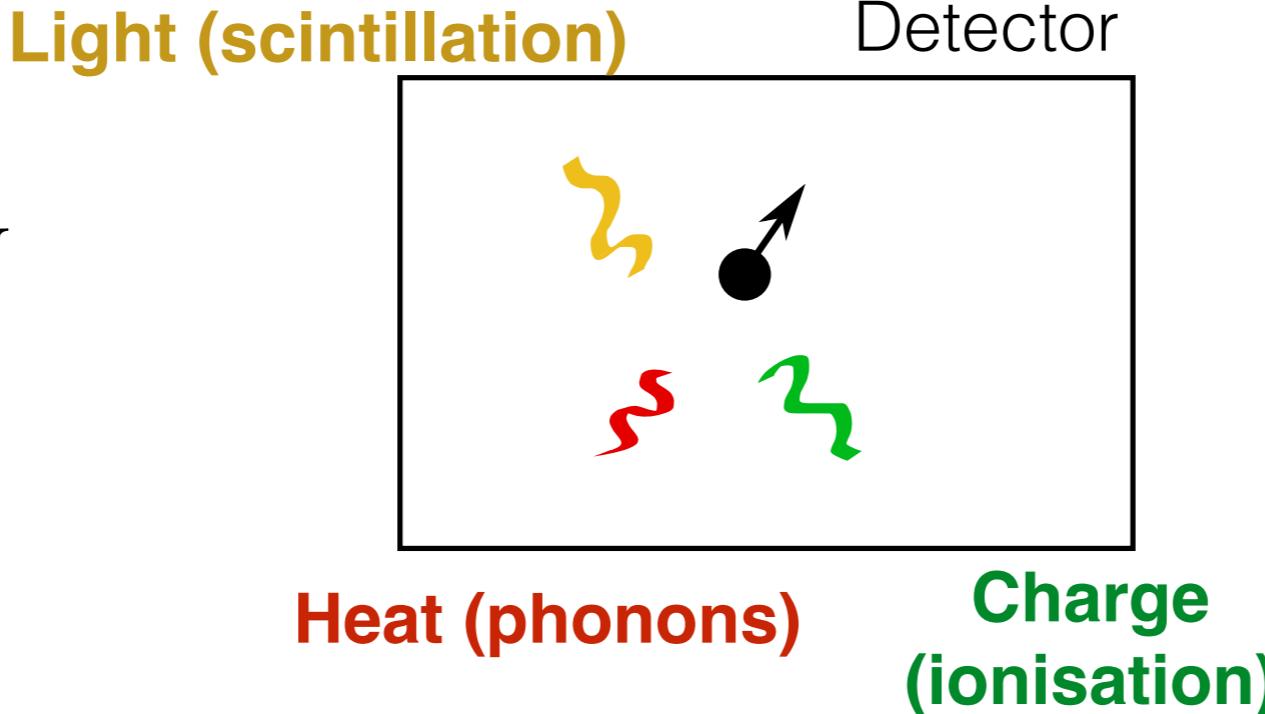
$$\frac{dR}{dE_R} = \frac{\rho_\chi}{m_\chi m_A} \int_{v_{\min}}^{\infty} v f(\mathbf{v}) \frac{d\sigma}{dE_R} d^3\mathbf{v}$$

Include all particles with enough speed to excite recoil of energy  $E_R$ :

$$v_{\min} = \sqrt{\frac{m_N E_R}{2\mu_{\chi N}^2}}$$

# Direct detection

$$m_\chi \gtrsim 1 \text{ GeV}$$
$$v \sim 10^{-3}$$



$$\frac{dR}{dE_R} = \frac{\rho_\chi}{m_\chi m_A} \int_{v_{\min}}^{\infty} v f(v) \frac{d\sigma}{dE_R} d^3v$$

Astro-physics

Particle and nuclear physics

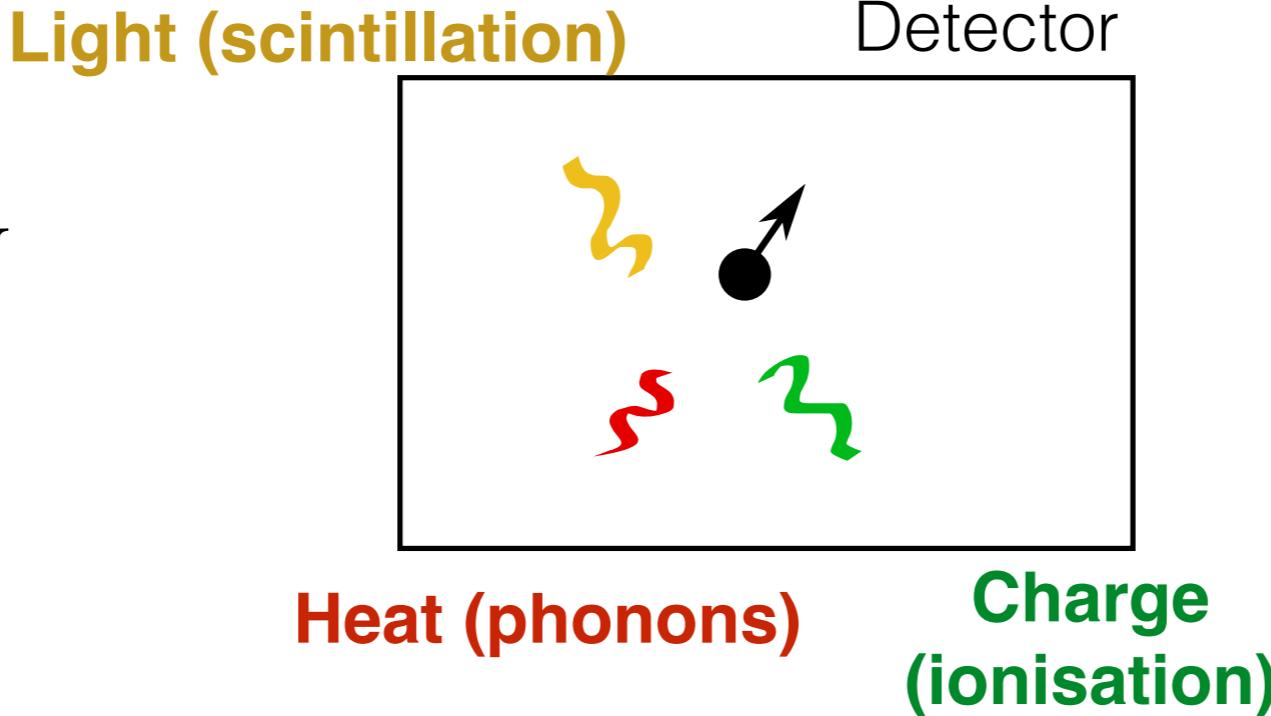
Diagram illustrating the differential rate of signal generation. The equation shows the rate  $dR/dE_R$  as a function of recoil energy  $E_R$ . The integrand consists of the density of dark matter  $\rho_\chi$ , the mass of the dark matter particle  $m_\chi$ , the mass of the target nucleus  $m_A$ , the velocity distribution  $v f(v)$ , the differential cross-section  $d\sigma/dE_R$ , and the volume element  $d^3v$ . Arrows point from the terms  $\rho_\chi/m_\chi m_A$  and  $v f(v) d\sigma/dE_R$  to the text "Astrophysics" and "Particle and nuclear physics" respectively, indicating their physical origins.

Include all particles with enough speed to excite recoil of energy  $E_R$ :

$$v_{\min} = \sqrt{\frac{m_N E_R}{2 \mu_{\chi N}^2}}$$

# Direct detection

$$m_\chi \gtrsim 1 \text{ GeV}$$
$$v \sim 10^{-3}$$



$$\frac{dR}{dE_R} = \frac{\rho_\chi}{m_\chi m_A} \int_{v_{\min}}^{\infty} v f(v) \frac{d\sigma}{dE_R} d^3v$$

Astro-physics

Particle and nuclear physics

Diagram illustrating the differential rate of signal detection. The equation shows the rate  $dR/dE_R$  as a function of recoil energy  $E_R$ . The terms include the dark matter density  $\rho_\chi$ , the mass of the dark matter particle  $m_\chi$ , the mass of the target nucleus  $m_A$ , the velocity distribution  $v f(v)$ , the differential cross-section  $d\sigma/dE_R$ , and the volume element  $d^3v$ .

Include all particles with enough speed to excite recoil of energy  $E_R$ :

$$v_{\min} = \sqrt{\frac{m_N E_R}{2 \mu_{\chi N}^2}}$$

But plenty of alternative ideas:  
DM-electron recoils [1108.5383]  
Superconducting detectors [1504.07237]  
Axion DM searches [1404.1455]

# Astrophysics of DM (the simple picture)

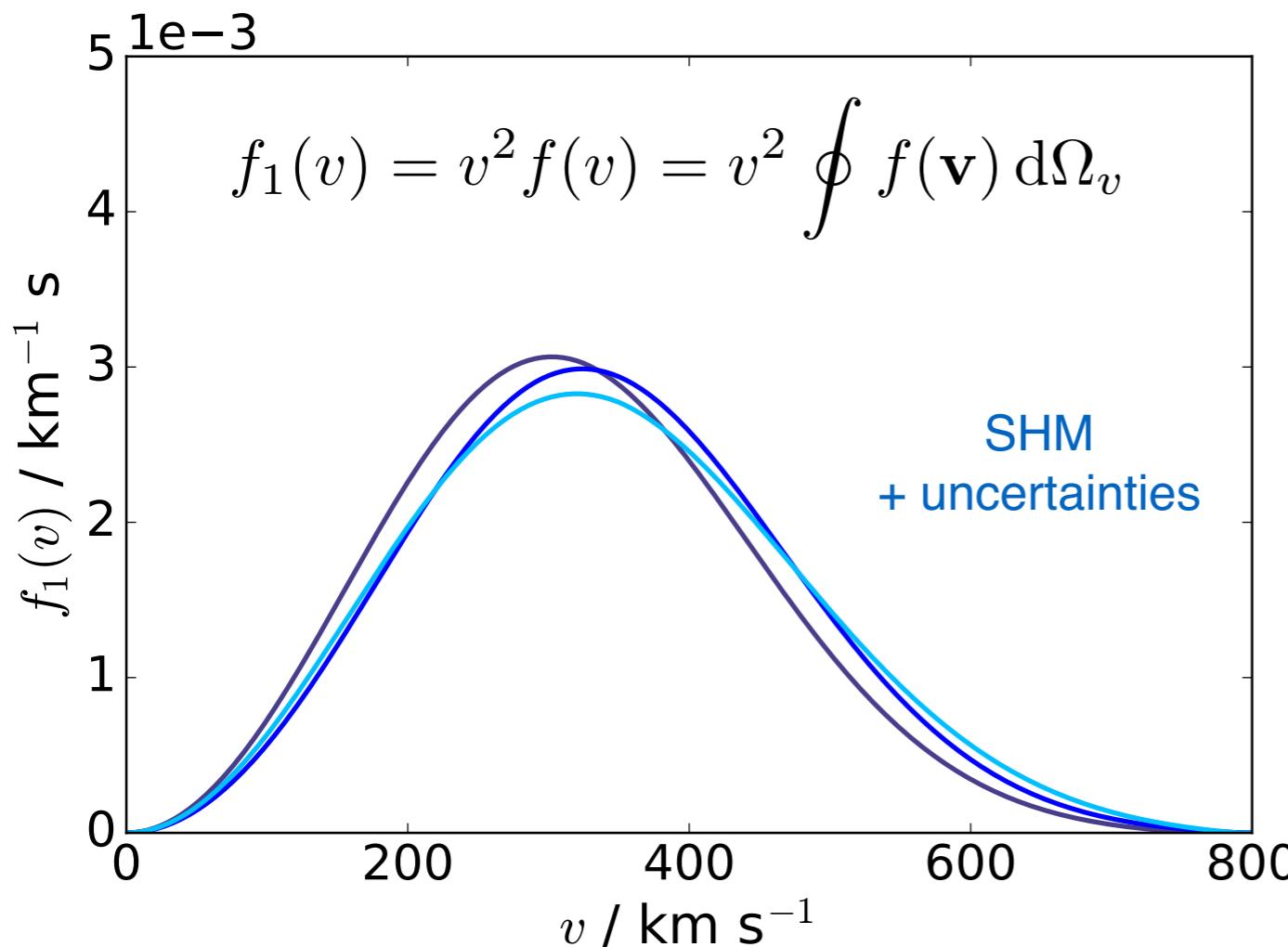
Standard Halo Model (**SHM**) is typically assumed: isotropic, spherically symmetric distribution of particles with  $\rho(r) \propto r^{-2}$ .

Leads to a Maxwell-Boltzmann (MB) distribution,

$$f_{\text{Lab}}(\mathbf{v}) = (2\pi\sigma_v^2)^{-3/2} \exp\left[-\frac{(\mathbf{v} - \mathbf{v}_e)^2}{2\sigma_v^2}\right] \Theta(|\mathbf{v} - \mathbf{v}_e| - v_{\text{esc}})$$

which is well matched in some hydro simulations.

[1601.04707, 1601.04725, 1601.05402]



$\mathbf{v}_e$  - Earth's Velocity

$$v_e \sim 220 - 250 \text{ km s}^{-1}$$

$$\sigma_v \sim 155 - 175 \text{ km s}^{-1}$$

Feast et al. [astro-ph/9706293],  
Bovy et al. [1209.0759]

$$v_{\text{esc}} = 533^{+54}_{-41} \text{ km s}^{-1}$$

Piffl et al. (RAVE) [1309.4293]

# Particle Physics of DM (the simple picture)

Typically assume contact interactions (heavy mediators).  
In the non-relativistic limit, obtain two main contributions.  
Write in terms of DM-proton cross section  $\sigma^p$ :

$$\frac{d\sigma^A}{dE_R} \propto \frac{\sigma^p}{\mu_{\chi p}^2 v^2} \mathcal{C}_A F^2(E_R)$$

Form factor accounts for  
loss of coherence at high  
energy

Enhancement factor different for:

*spin-independent (SI) interactions* -  $\mathcal{C}_A^{\text{SI}} \sim A^2$

*spin-dependent (SD) interactions* -  $\mathcal{C}_A^{\text{SD}} \sim (J+1)/J$

Interactions which are higher order in  $v$   
are possible. See the non-relativistic EFT  
of Fitzpatrick et al. [1203.3542]

# The final event rate

$$\frac{dR}{dE_R} \sim \int_{v_{\min}}^{\infty} v f(\mathbf{v}) \frac{d\sigma}{dE_R} d^3\mathbf{v} + \frac{d\sigma}{dE_R} \propto \frac{1}{v^2} \rightarrow \frac{dR}{dE_R} \sim \frac{\rho_\chi}{m_\chi} \mathcal{C}_A \eta(v_{\min})$$

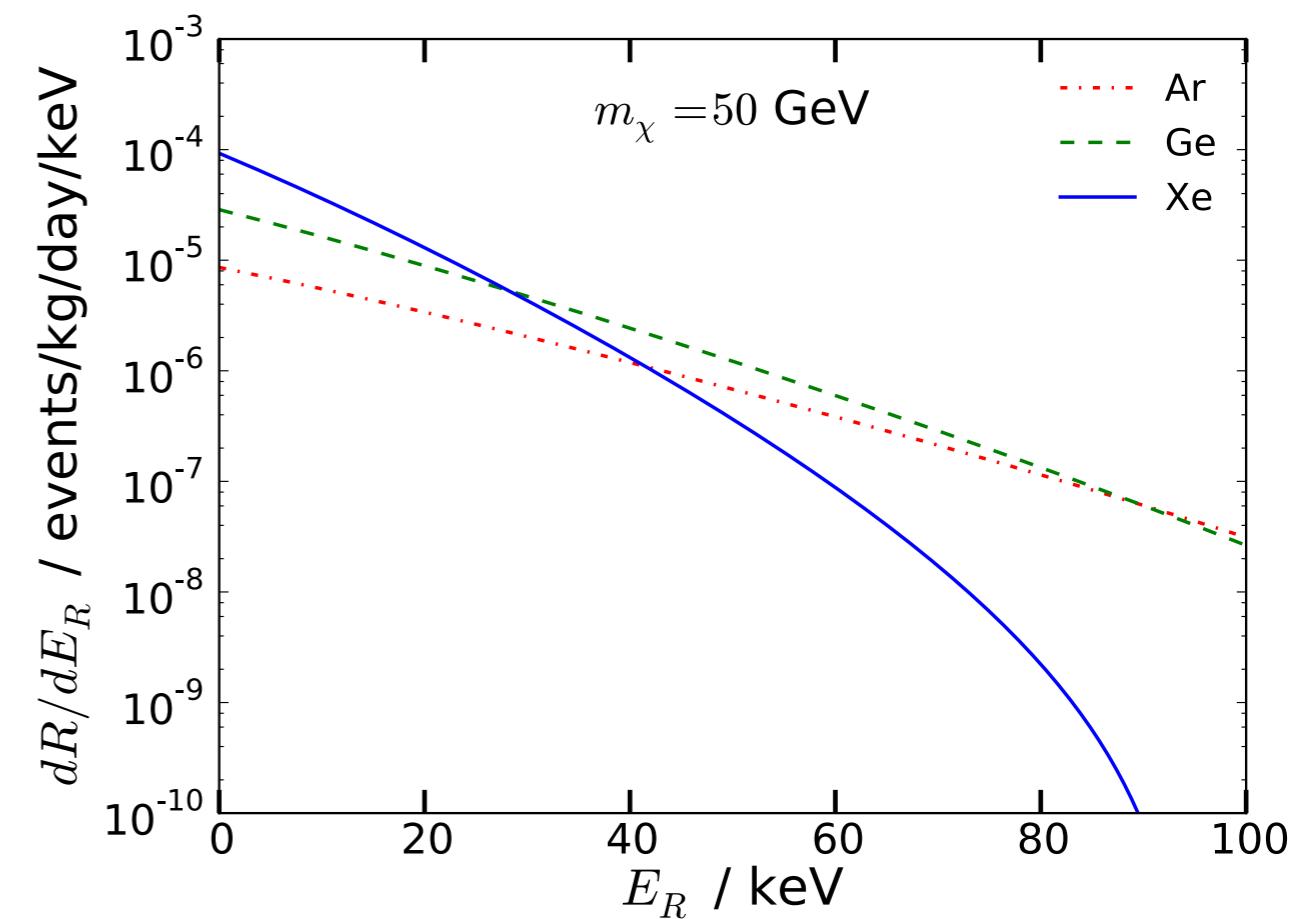
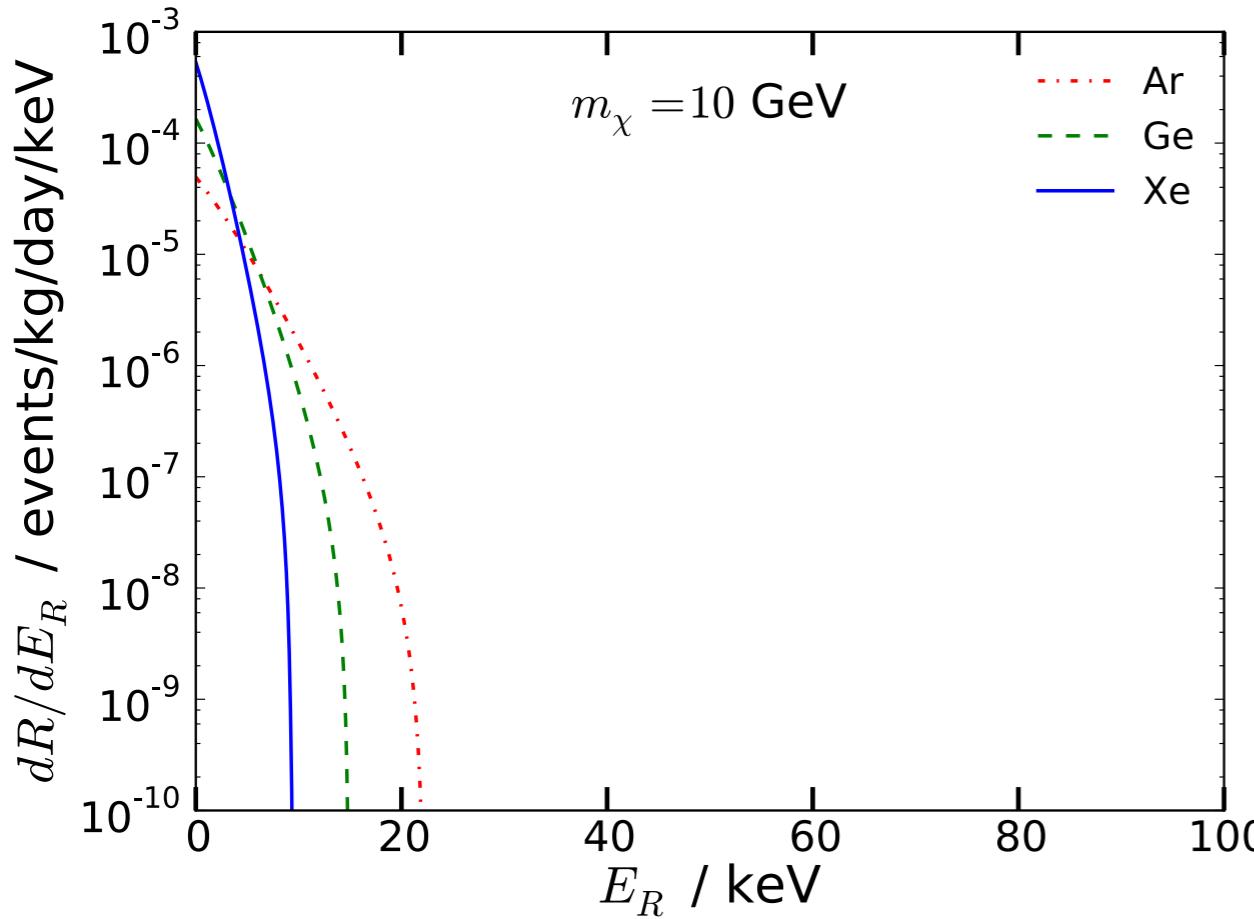
The ‘velocity integral’:

$$\eta(v_{\min}) \equiv \int_{v_{\min}}^{v_{\text{esc}}} \frac{f_1(v)}{v} dv$$

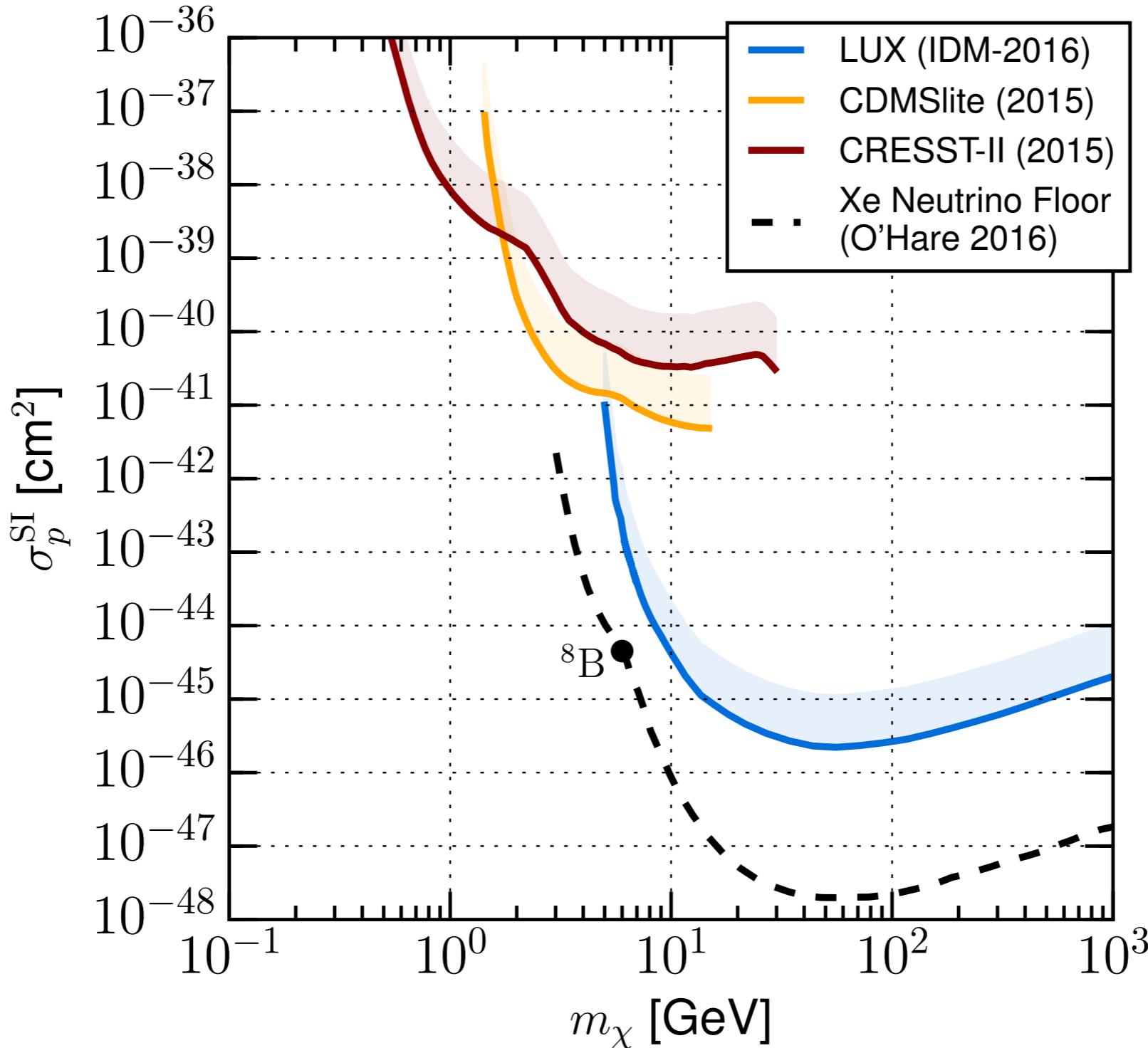
where

$$f_1(v) = v^2 \oint f(\mathbf{v}) d\Omega_v$$

SI interactions, SHM distribution



# The current landscape

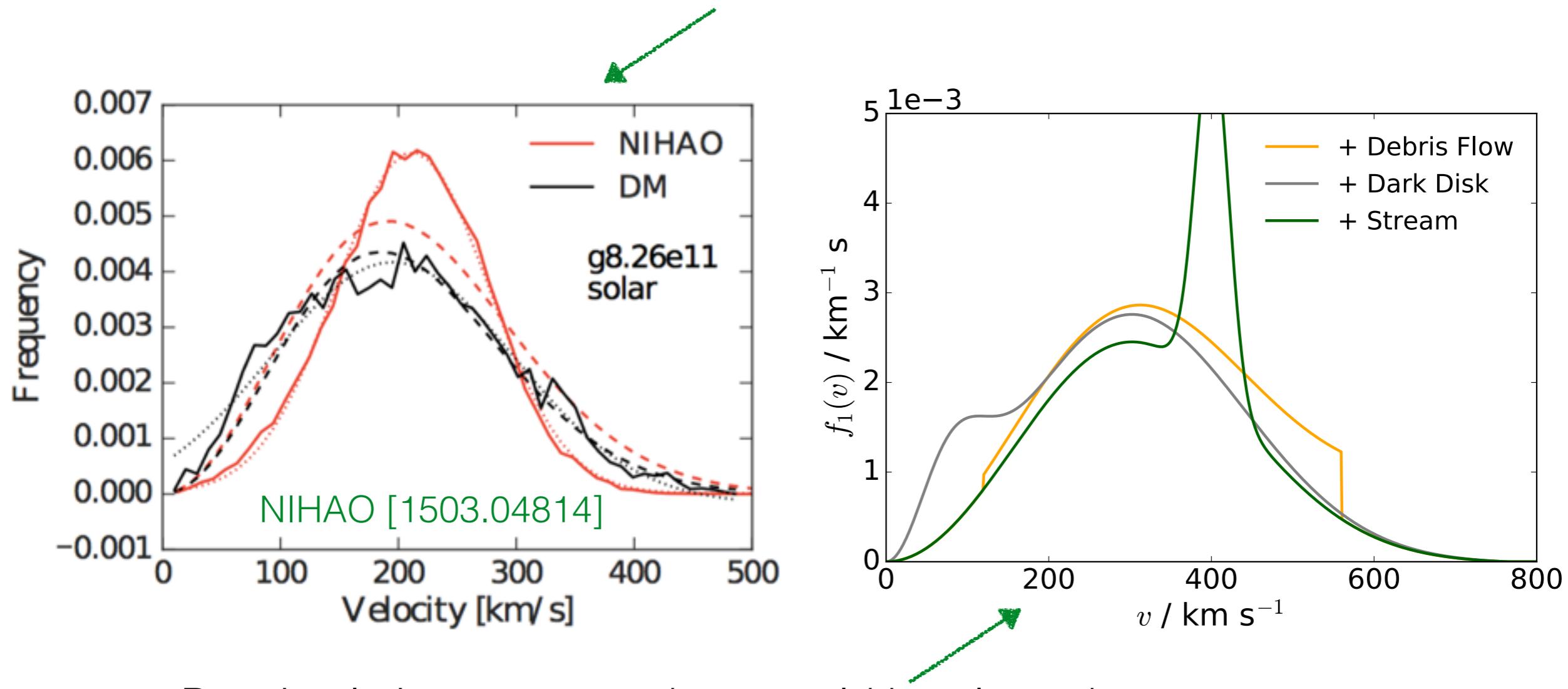


*Assuming the Standard Halo Model...*

# Overcoming halo uncertainties in direct detection

# Astrophysical uncertainties

The Standard Halo Model (SHM) has some inherent uncertainties.  
But there could also be deviations from MB form:



But simulations suggest there could be also substructure:

Debris flows      Kuhlen et al. [1202.0007]

Dark disk      Pillepich et al. [1308.1703], Schaller et al. [1605.02770]

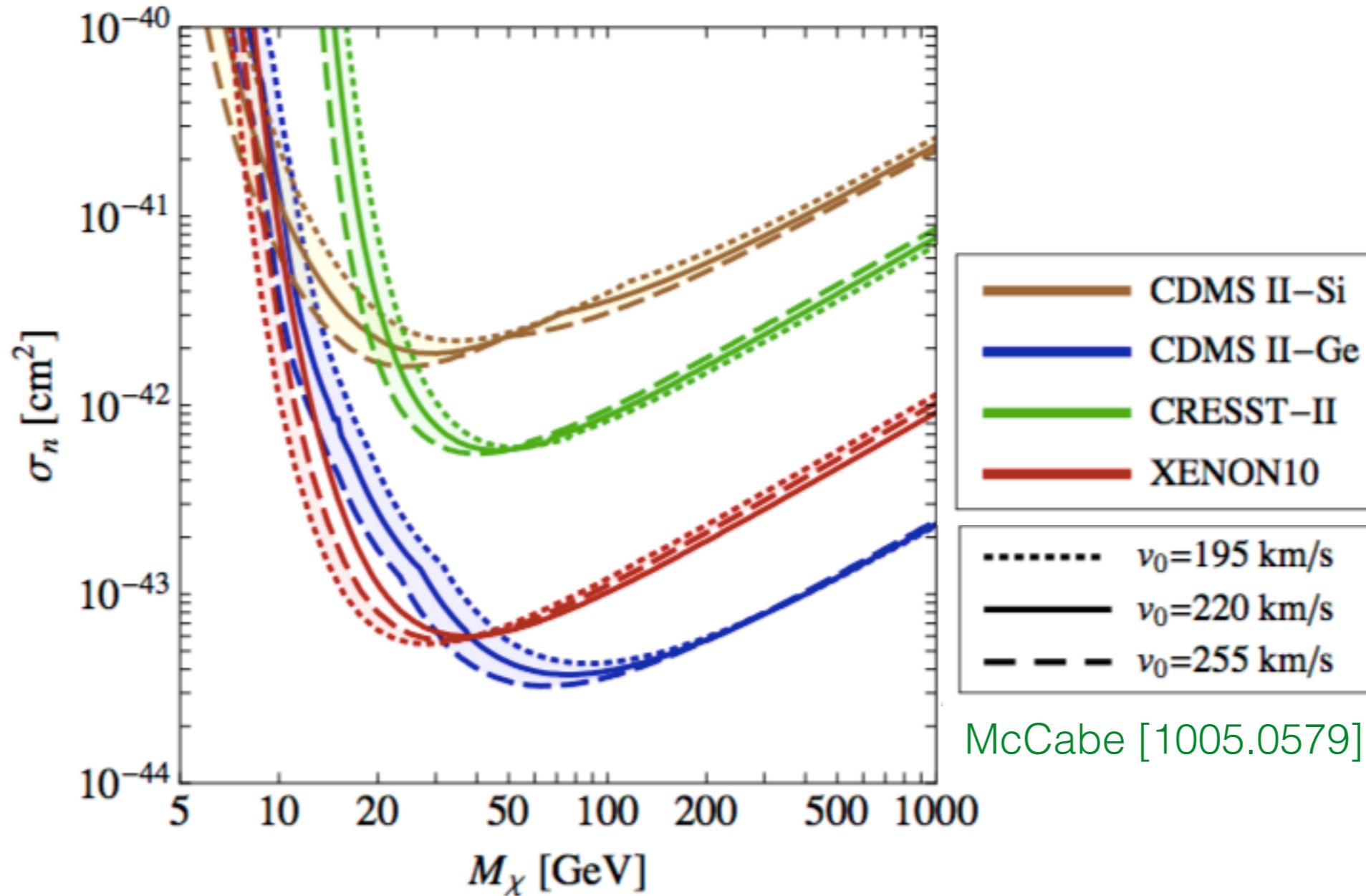
Tidal stream      Freese et al. [astro-ph/0309279, astro-ph/0310334]

# What could go wrong? (1)

Compare direct detection limits, incorporating SHM uncertainties

↳ may affect proper comparison/compatibility of results

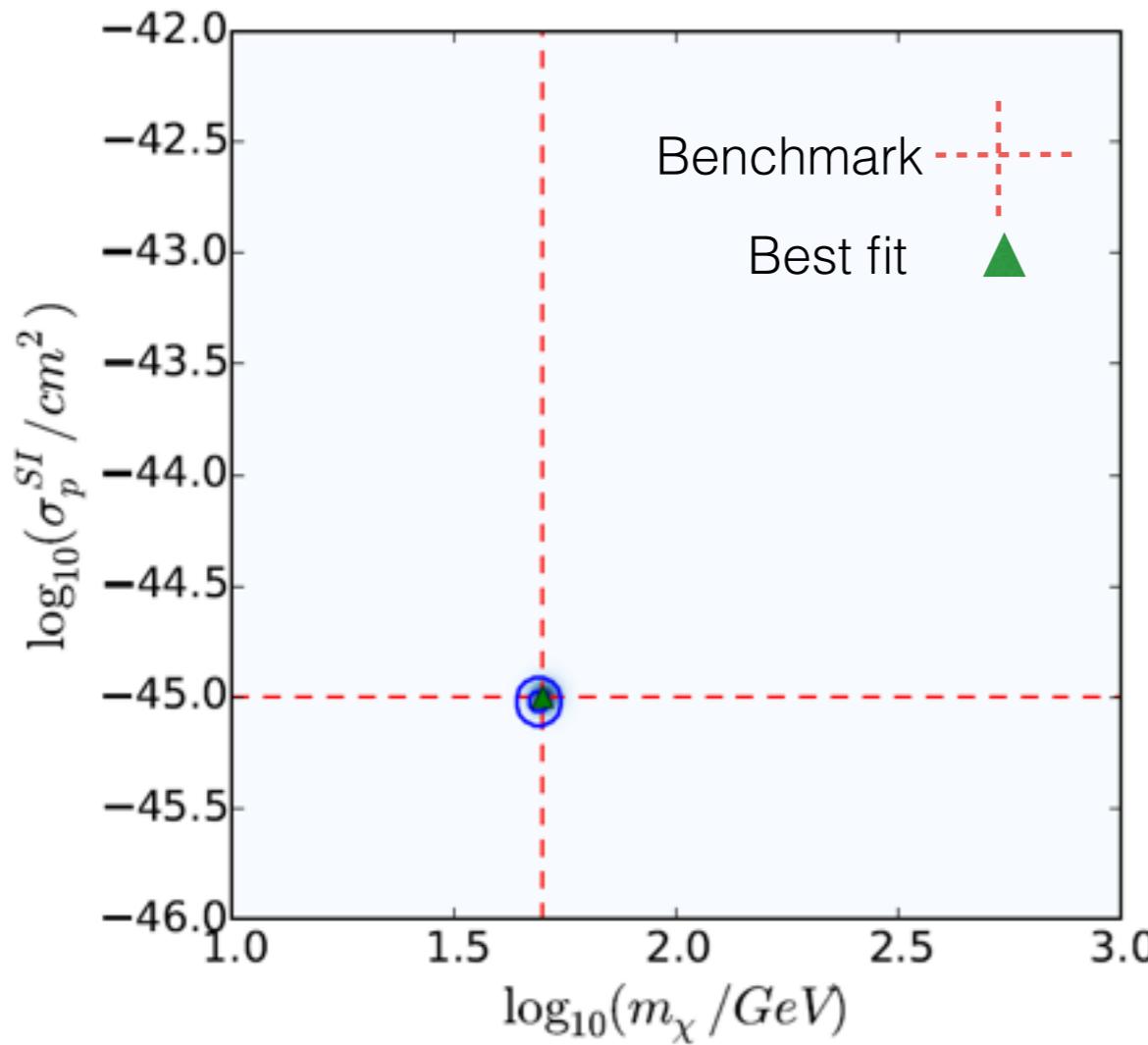
e.g. March-Russell at al. [0812.1931]



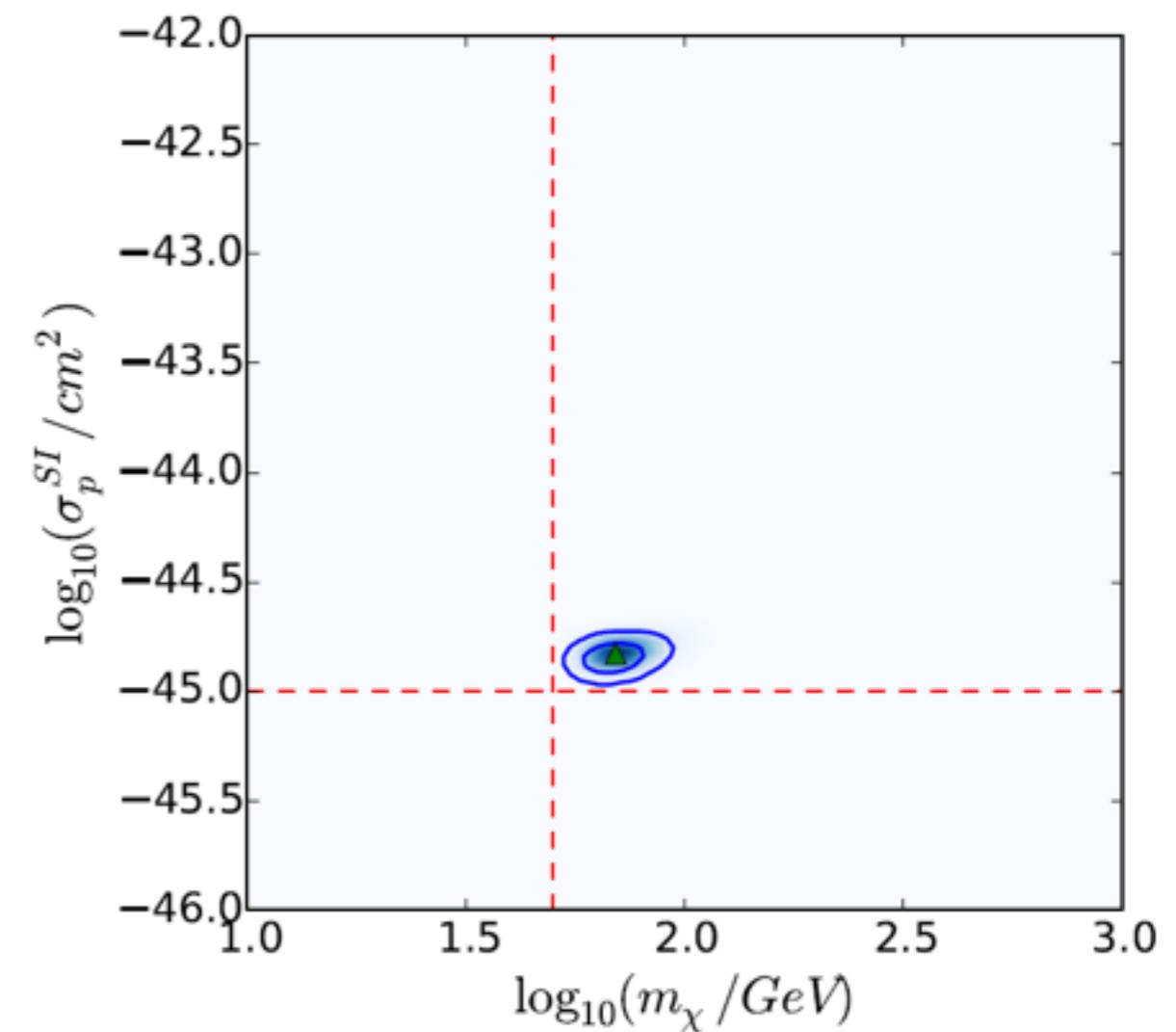
# What could go wrong? (2)

Generate mock data for several experiments, assuming a **stream** distribution, then try to reconstruct the mass and cross section assuming:

(correct) **stream** distribution



(incorrect) **SHM** distribution



# Halo-independent methods

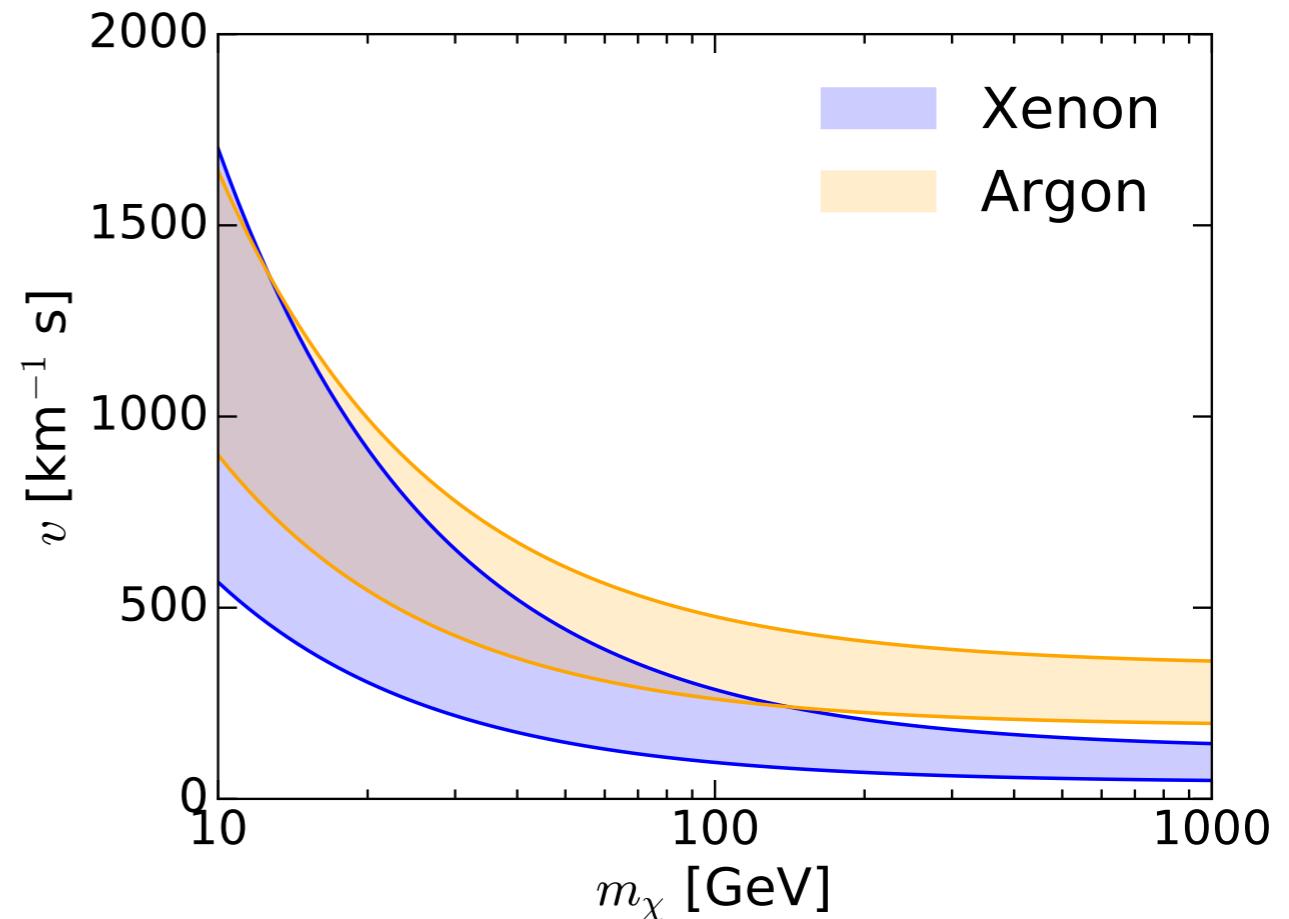
Experiments sensitive to a fixed range of recoil energies and therefore (through  $v_{\min}(E_R)$ ) a fixed range of speeds



Ask whether results are consistent over the range of speeds where two experiments overlap



Compare  $\eta(v_{\min}) \equiv \int_{v_{\min}}^{\infty} \frac{f_1(v)}{v} dv$   
(inferred from rate) over this limited range



Fox et al. [1011.1915, 1011.1910], but see also [1111.0292, 1107.0741, 1202.6359, 1304.6183, 1403.4606, 1403.6830, 1504.03333, 1607.02445, 1607.04418 and more...]

*But ideally we want to fit  $f_1(v)$ , the speed distribution.*

# Reconstructing the speed distribution

Write a *general parametrisation* for the speed distribution:

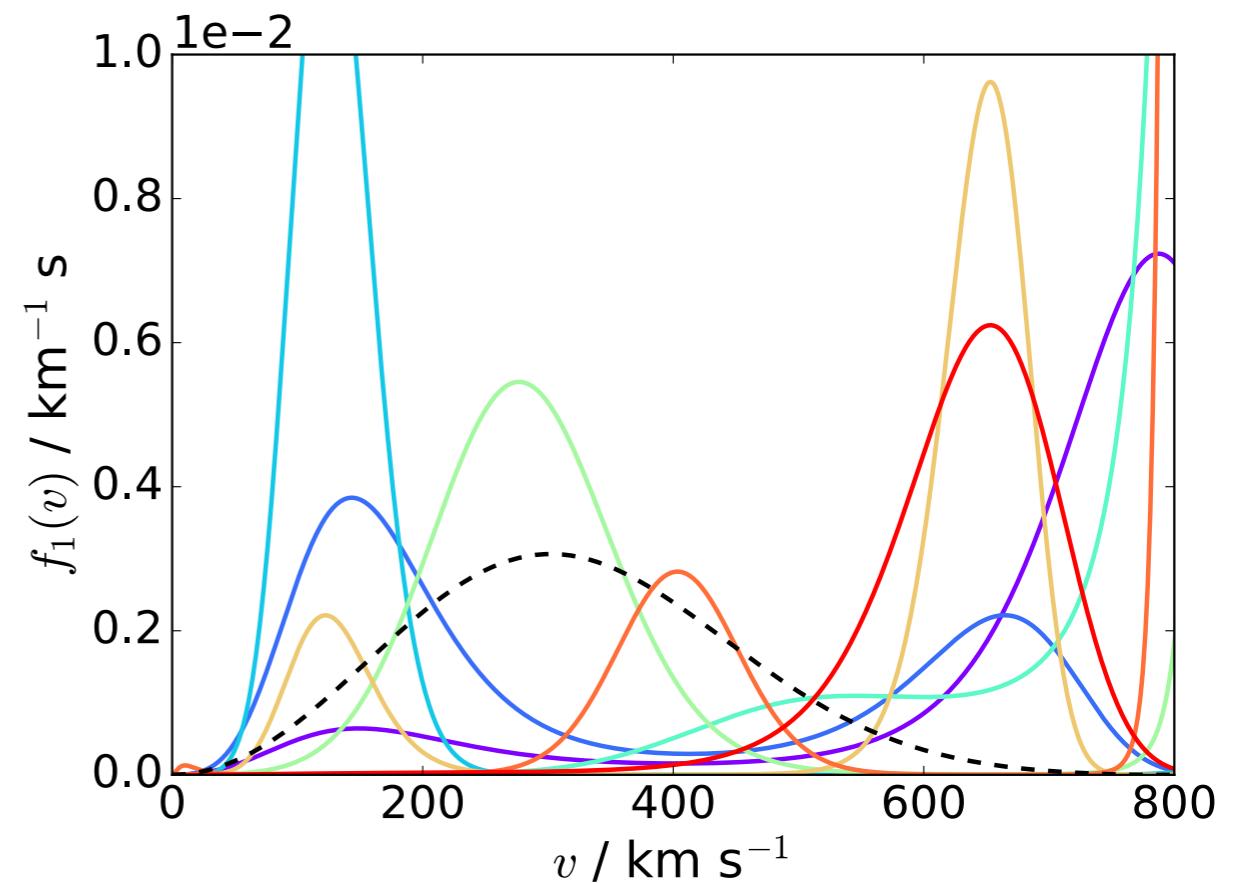
Peter [1103.5145]

$$f_1(v) = v^2 \exp \left( - \sum_{m=0}^{N-1} a_m v^m \right)$$

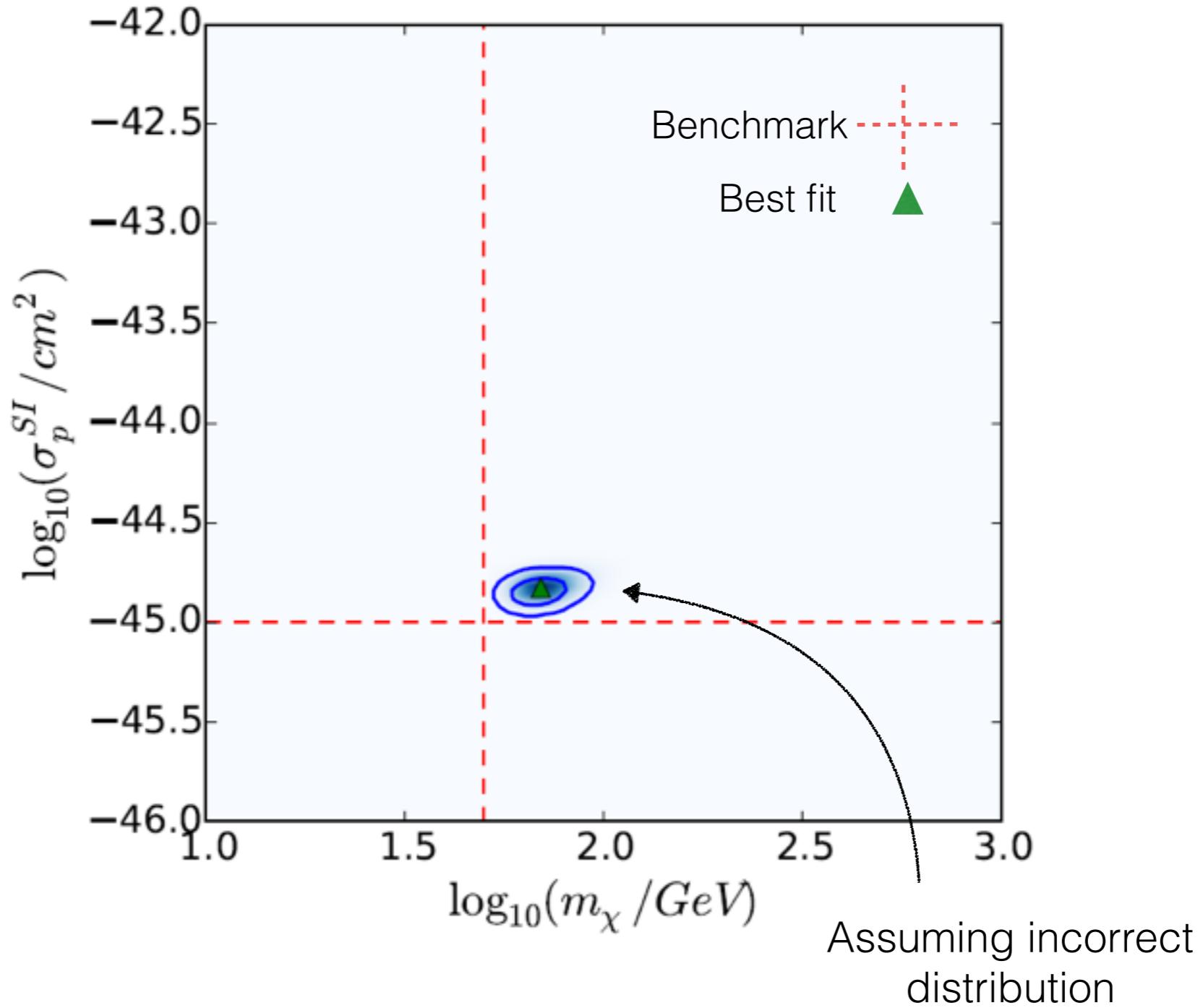
BJK & Green [1303.6868]

This form guarantees a distribution function which is *everywhere positive*.

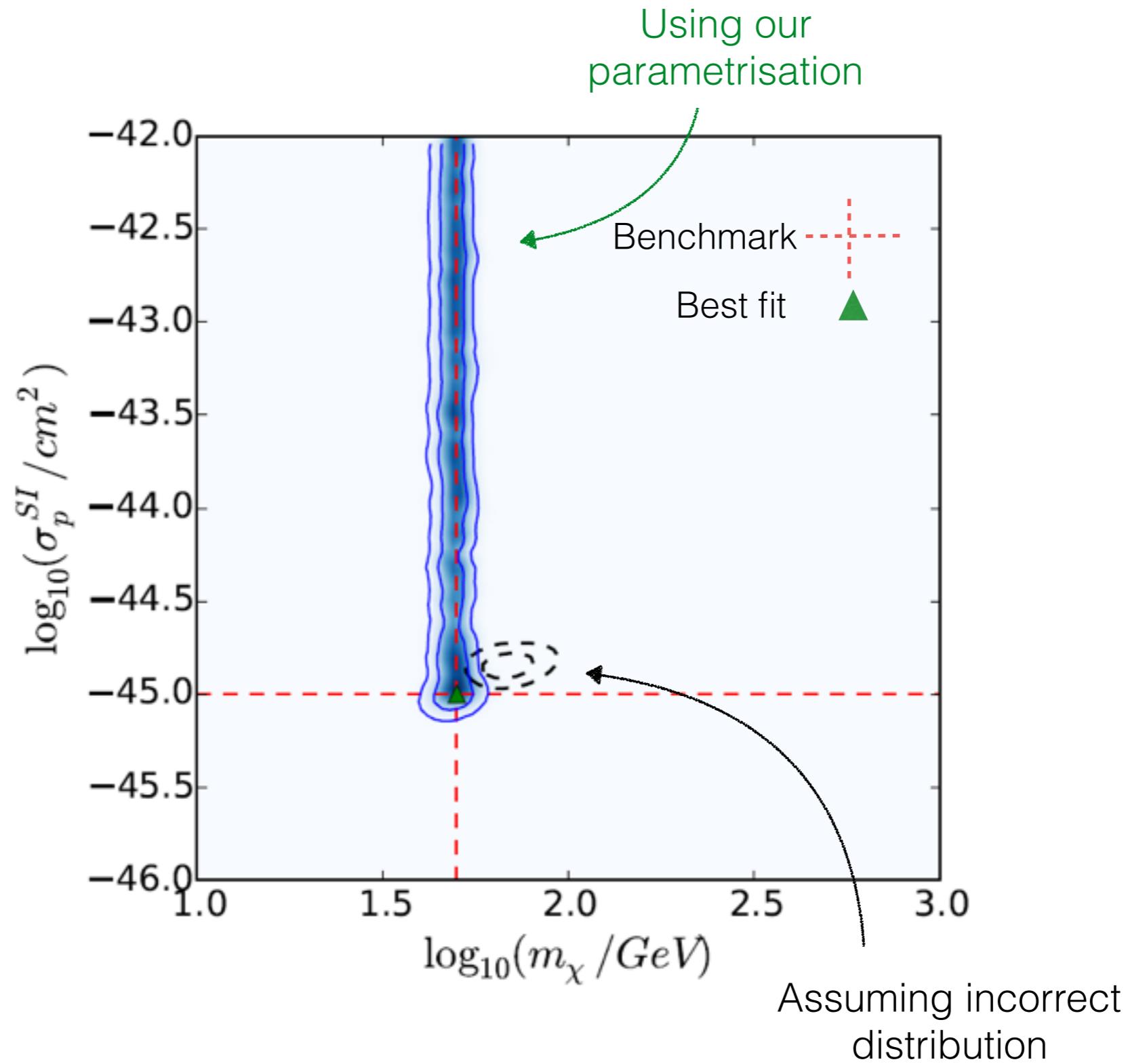
Now we attempt to fit the particle physics parameters  $(m_\chi, \sigma^p)$ , as well as the astrophysics parameters  $\{a_m\}$ .



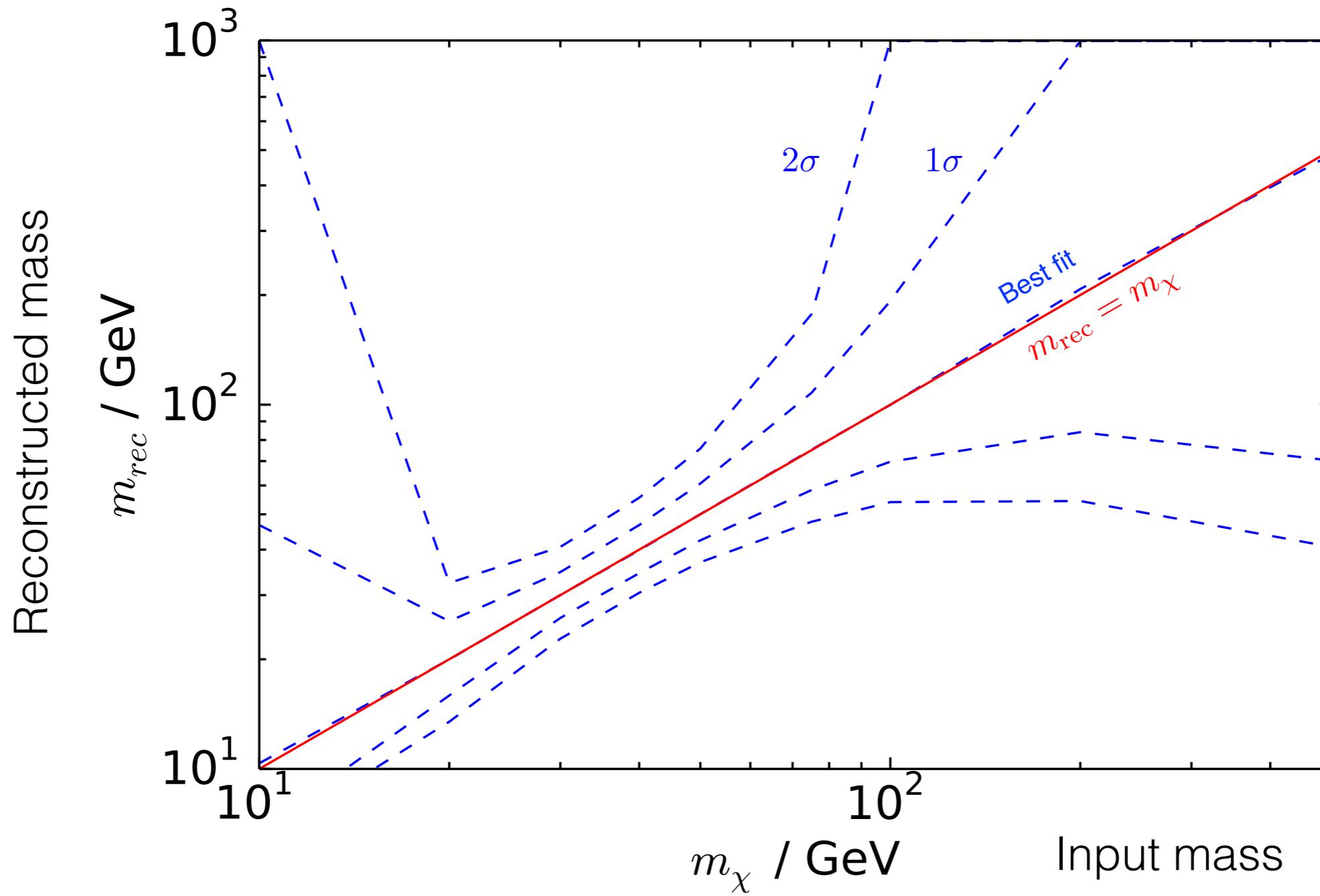
# Testing the parametrisation



# Testing the parametrisation

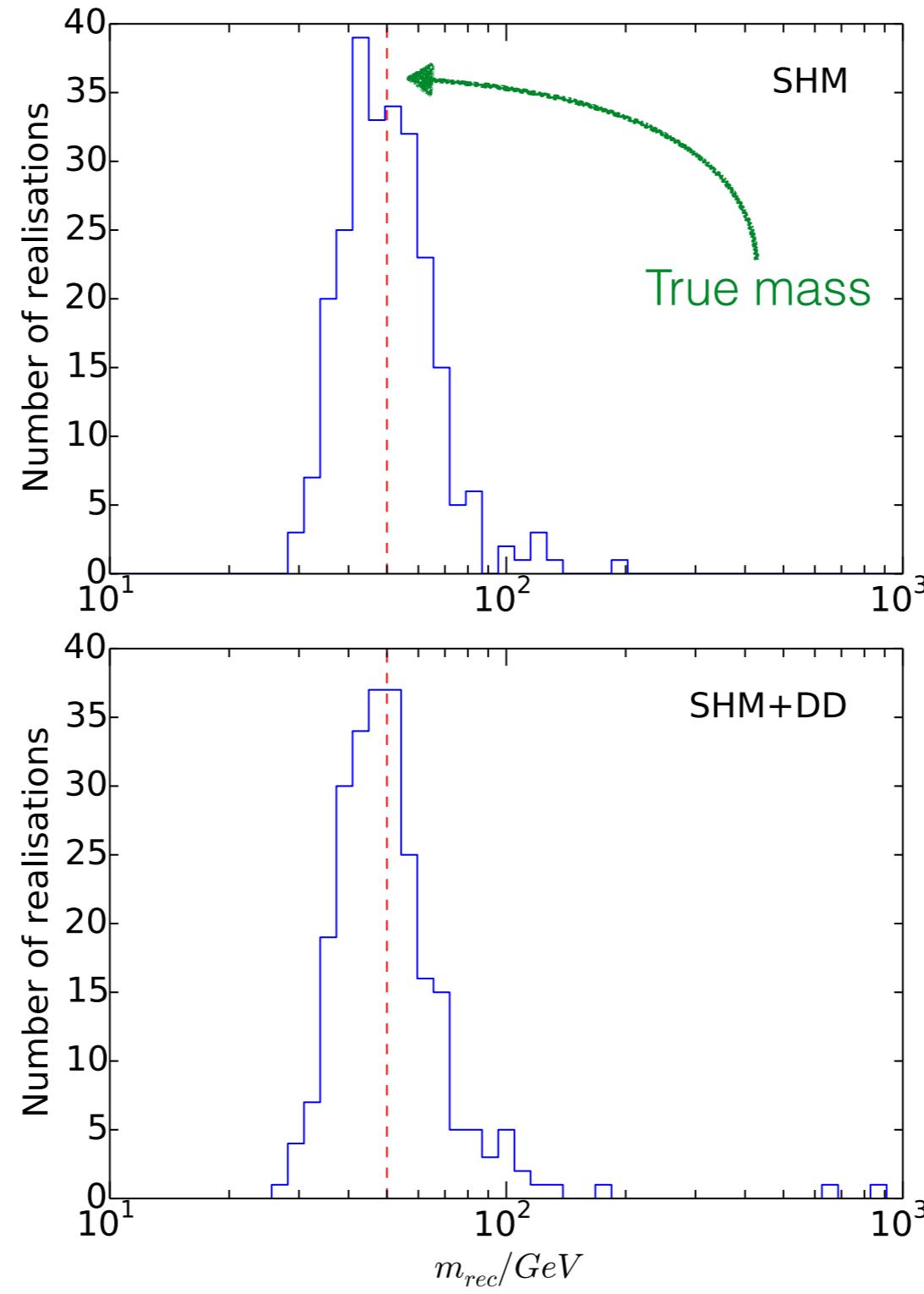


# Testing the parametrisation



BJK [1312.1852]

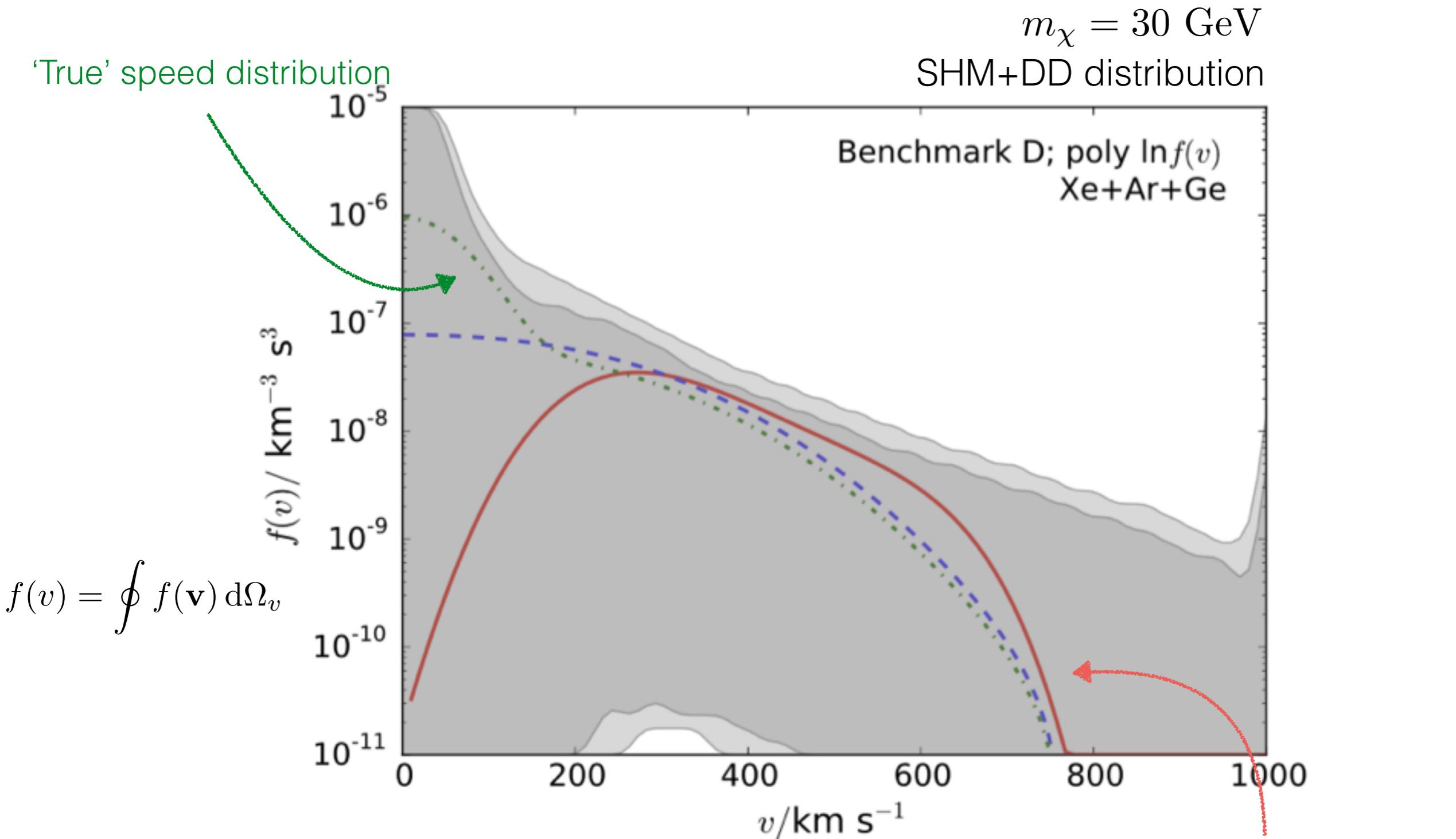
# Testing the parametrisation



Reconstructed mass

BJK [1312.1852]

# Reconstructing the speed distribution



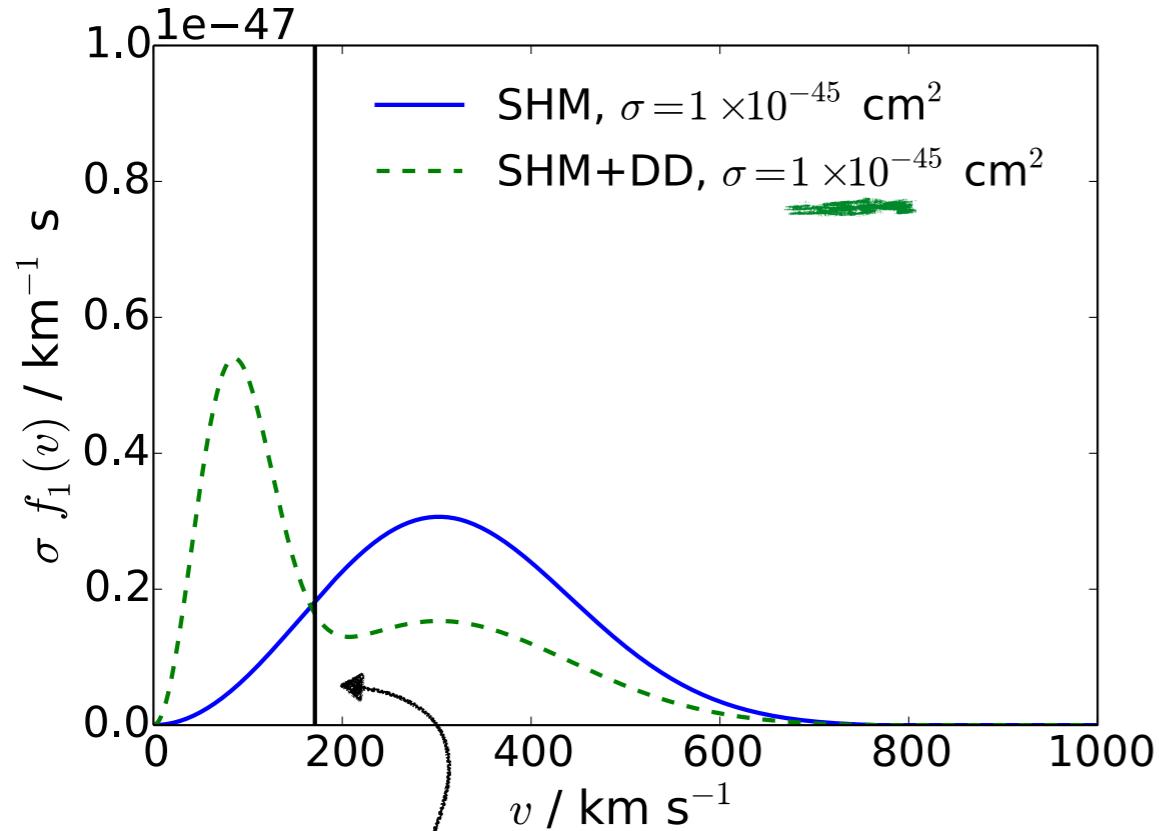
BJK, Fornasa, Green [1410.8051]

Bradley J Kavanagh (LPTHE, Paris)

• DM Particle Astronomy

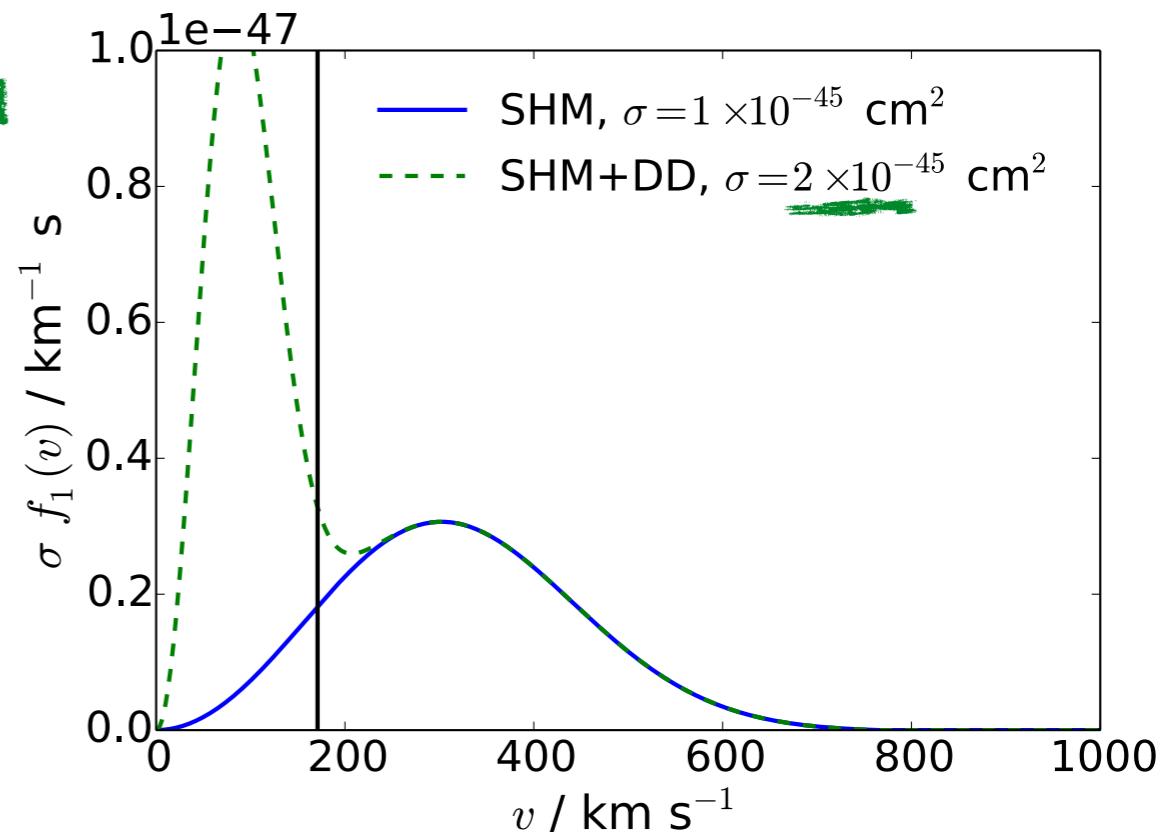
• GRAPPA Institute - 10th October 2016

# Cross section degeneracy



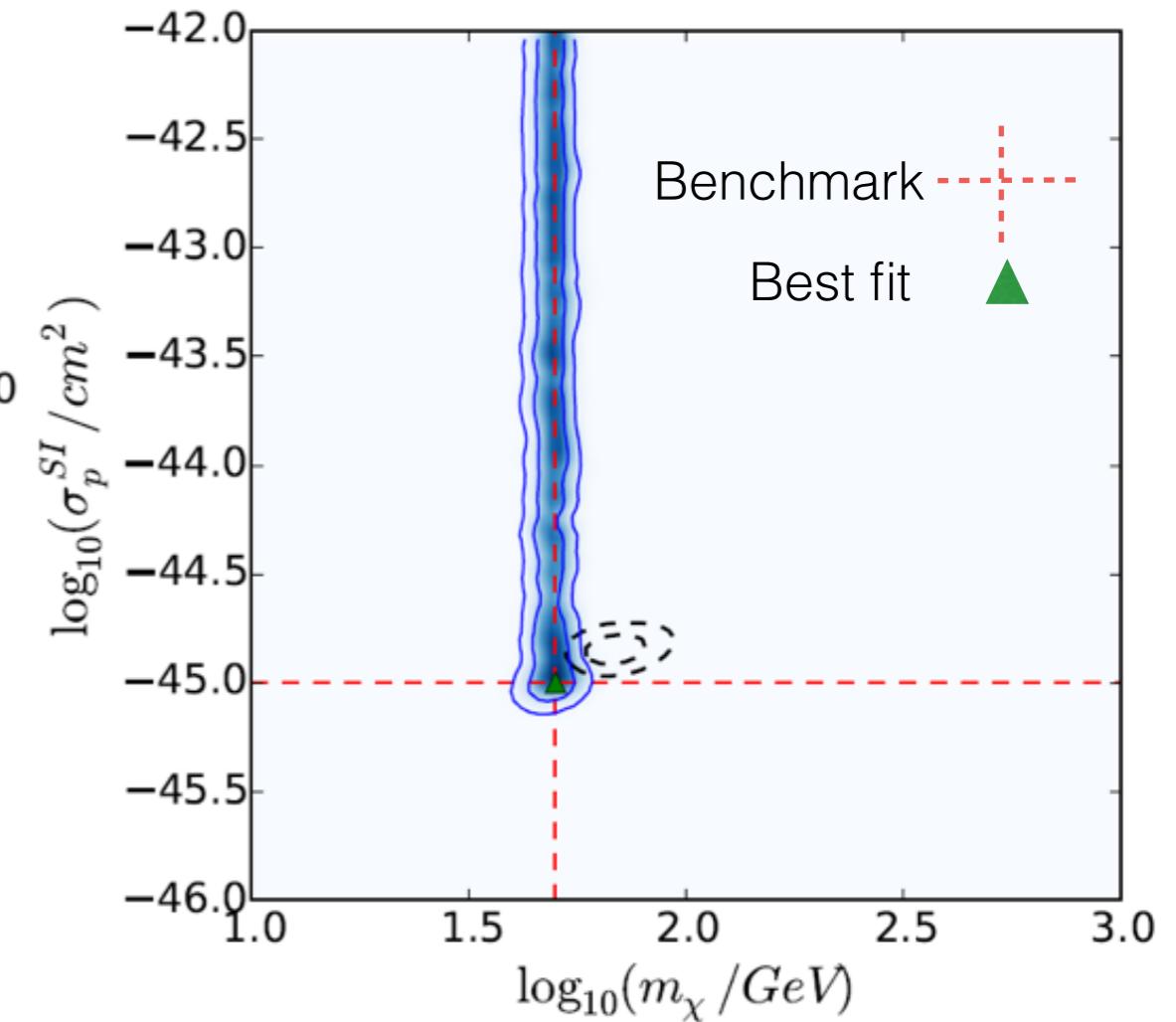
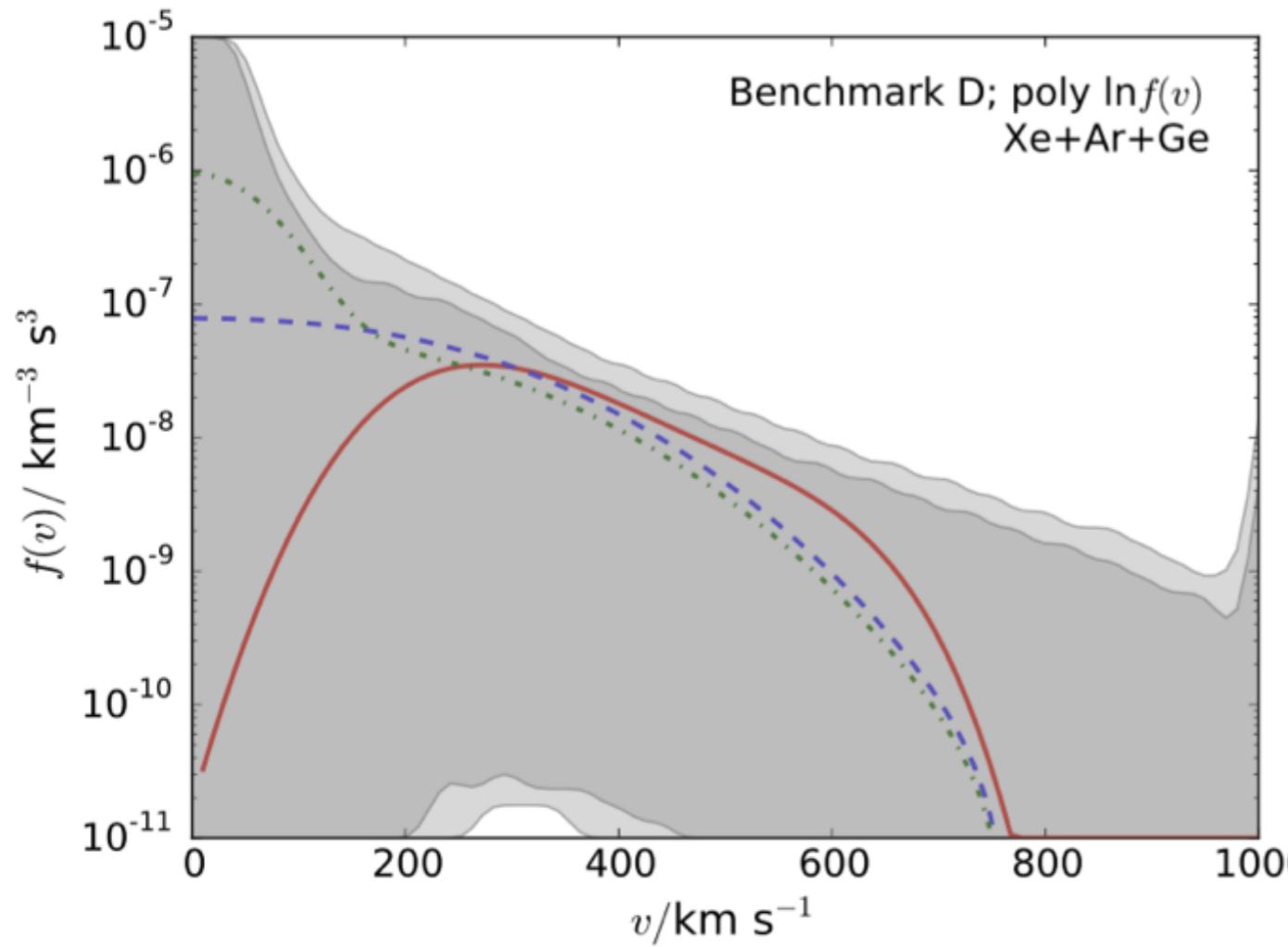
Minimum DM speed probed by  
a typical Xe experiment

$$\frac{dR}{dE_R} \propto \sigma \int_{v_{\min}}^{\infty} \frac{f_1(v)}{v} dv$$



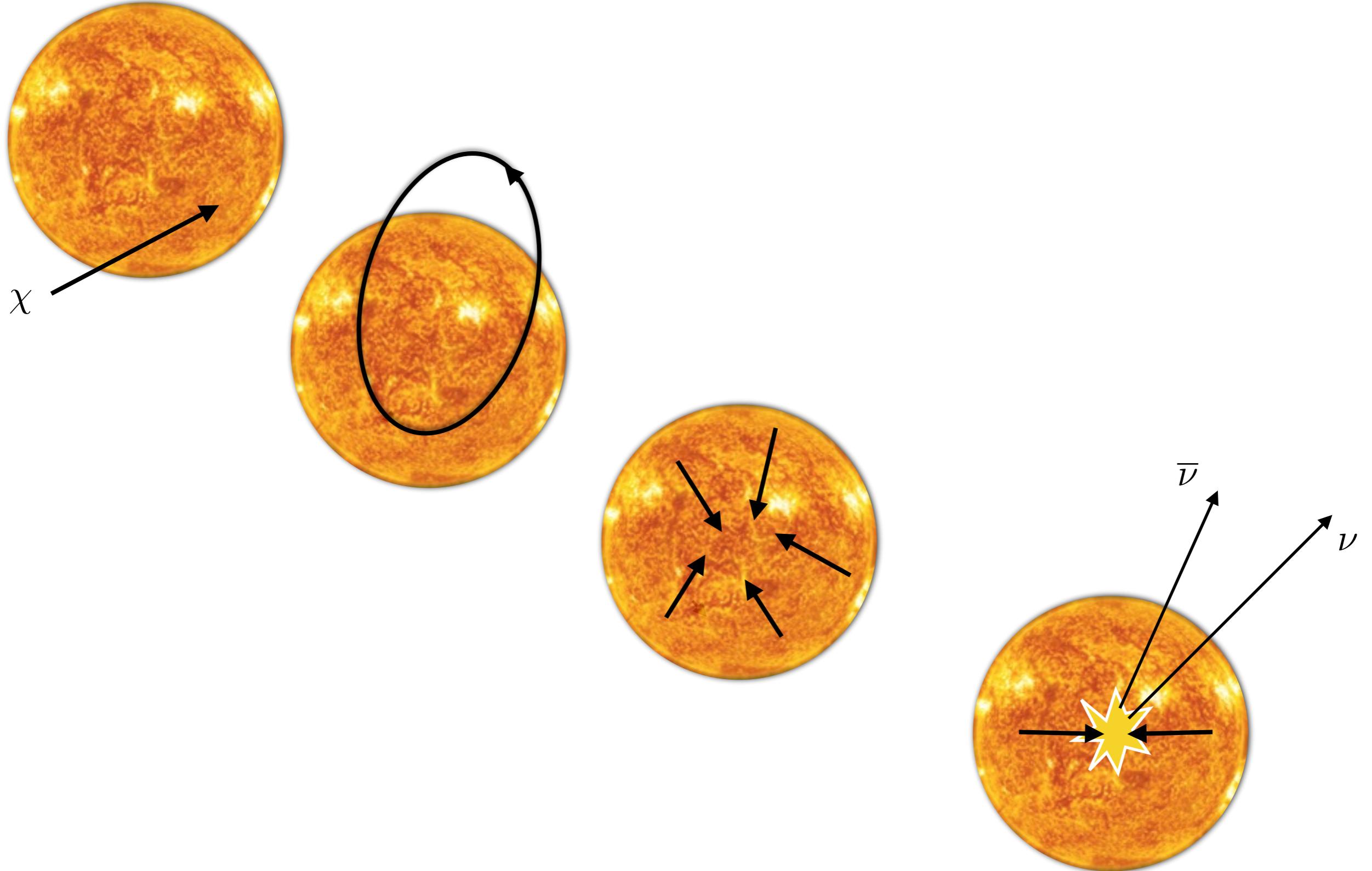
This is a problem for *any*  
astrophysics-independent method!

# Cross section degeneracy



# Neutrino telescopes

# DM capture in the Sun



# Incorporating IceCube

IceCube can detect the neutrinos from DM annihilation in the Sun

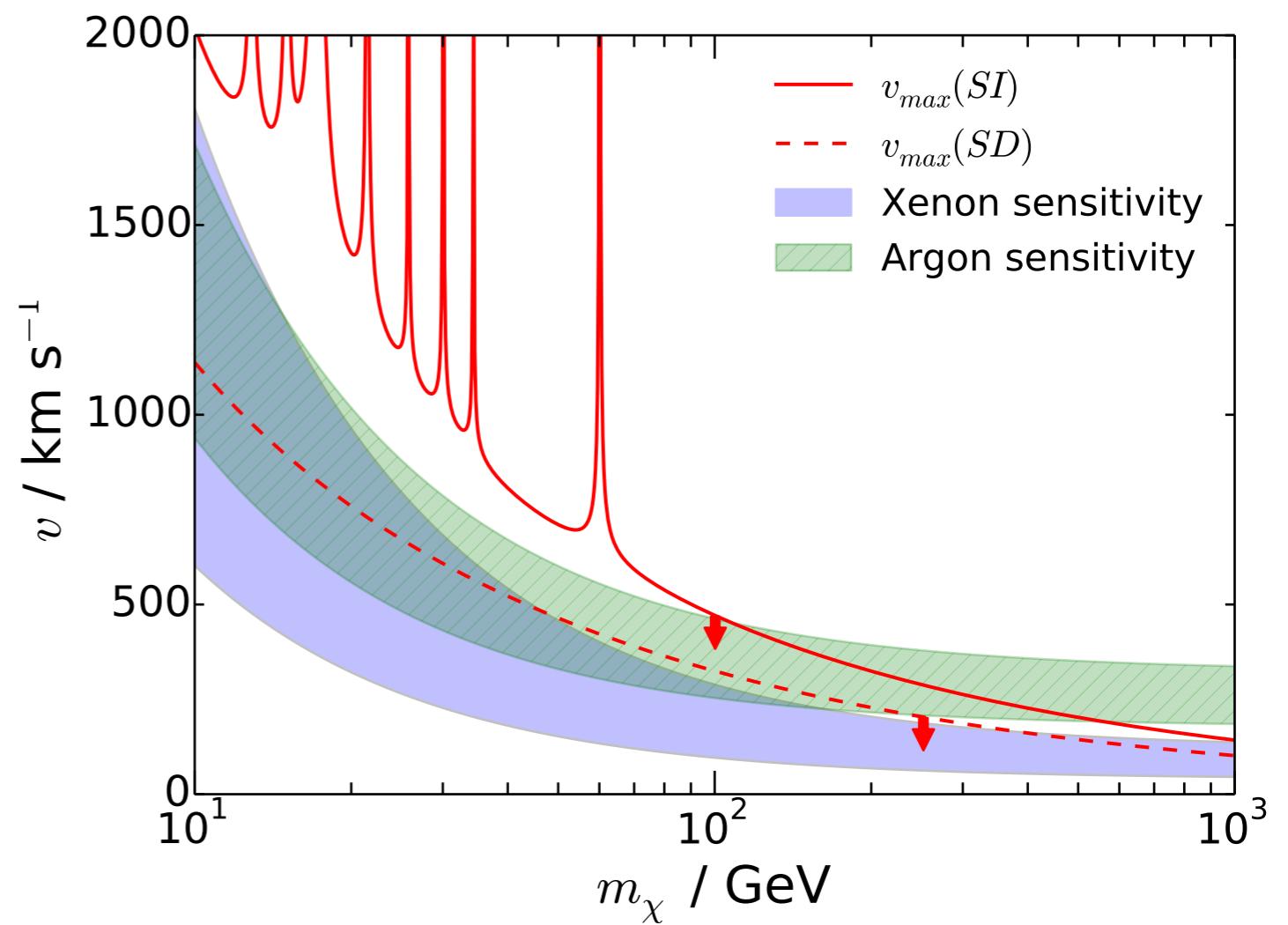
Assuming equilibrium in the Sun, rate is driven by solar capture of DM, which depends on the DM-nucleus scattering cross section

Crucially, only low energy DM particles are captured:

$$\frac{dC}{dV} \sim \sigma \int_0^{v_{\max}} \frac{f_1(v)}{v} dv$$

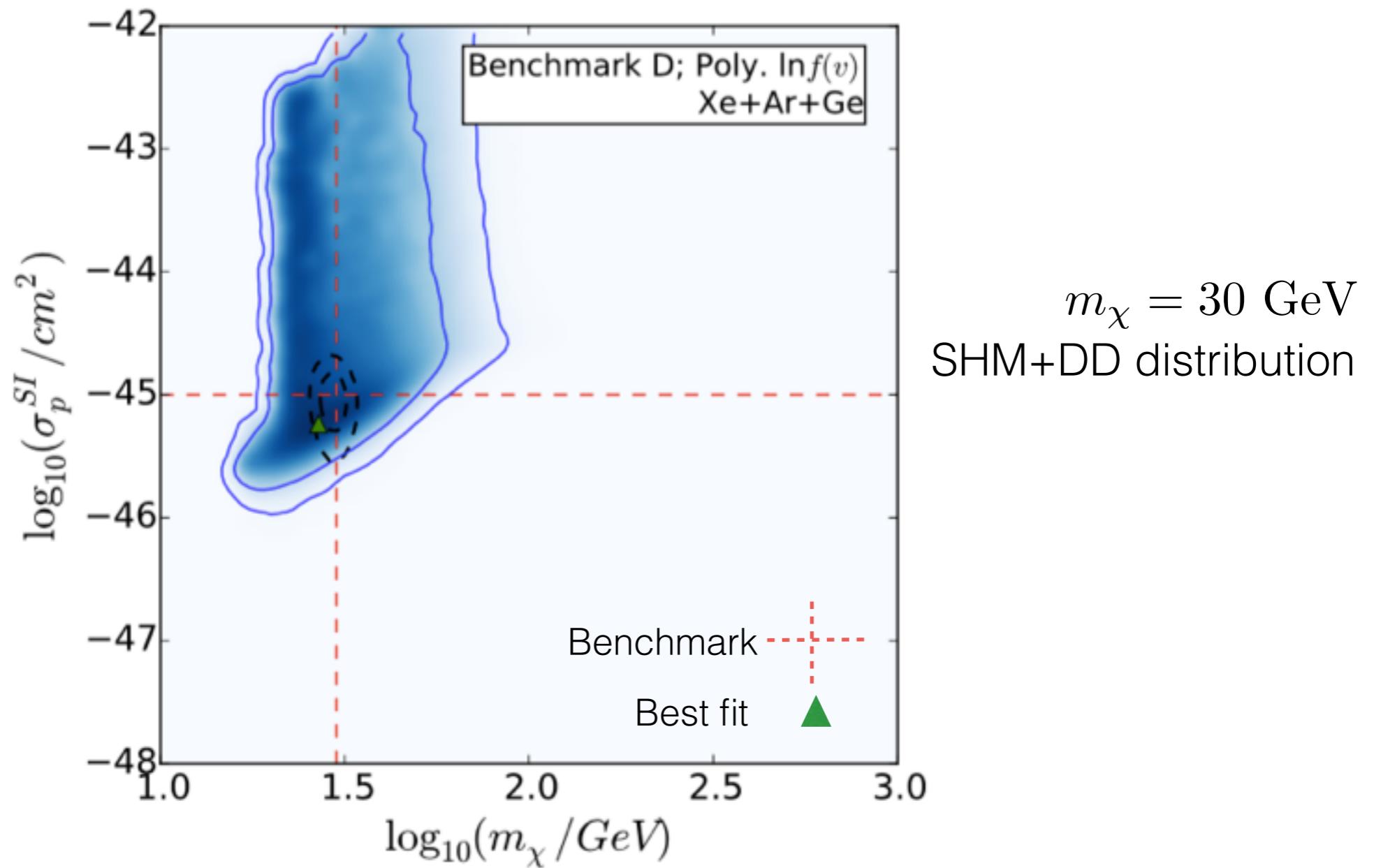
Gould (1991)

*If we also had a signal in IceCube, what could we do then?*



# Reconstructions without IceCube

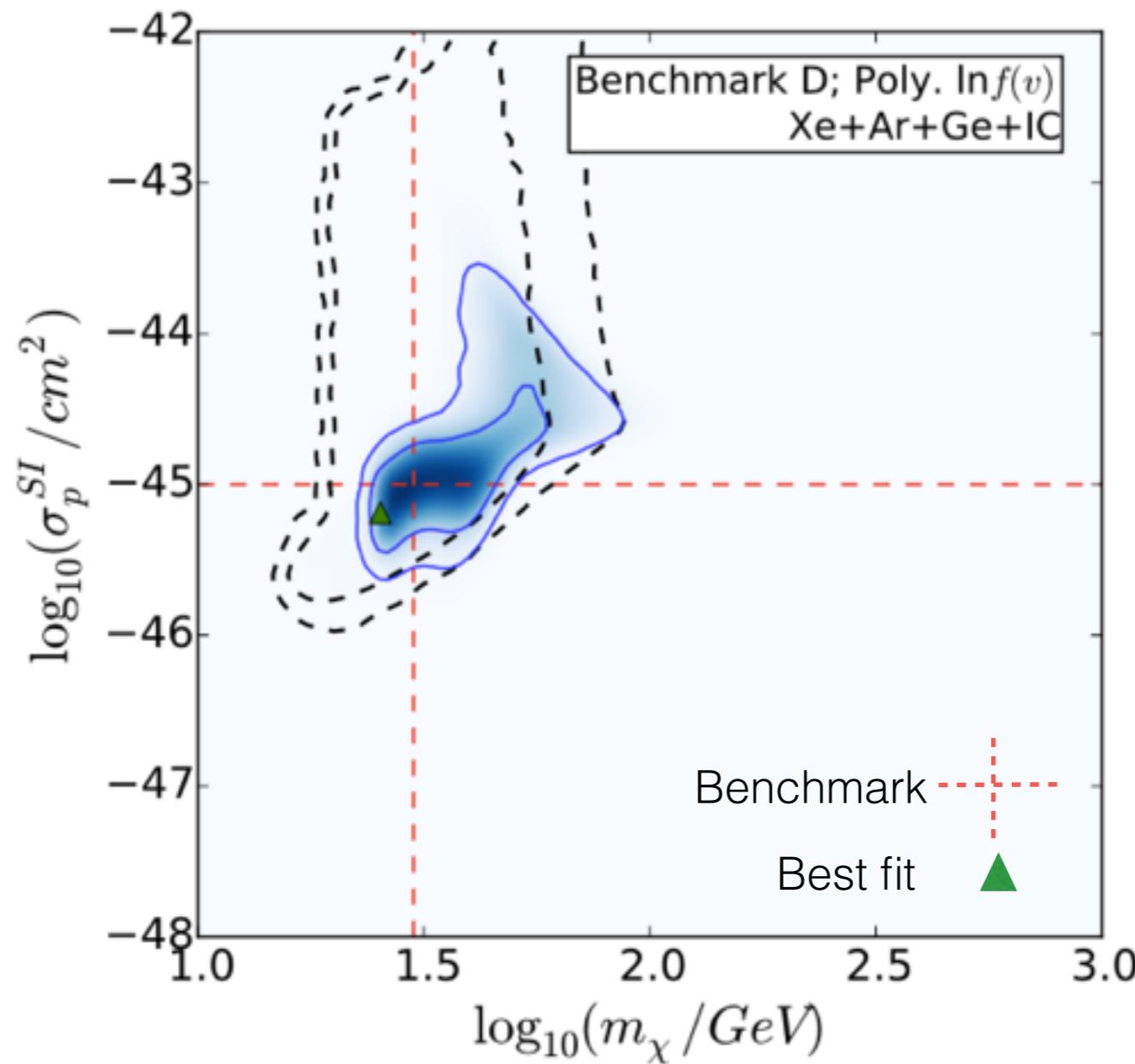
Mass and cross section reconstruction using three different direct detection experiments



BJK, Fornasa, Green [1410.8051]

# Reconstructions with IceCube

Mass and cross section reconstruction using three different direct detection experiments and an IceCube signal



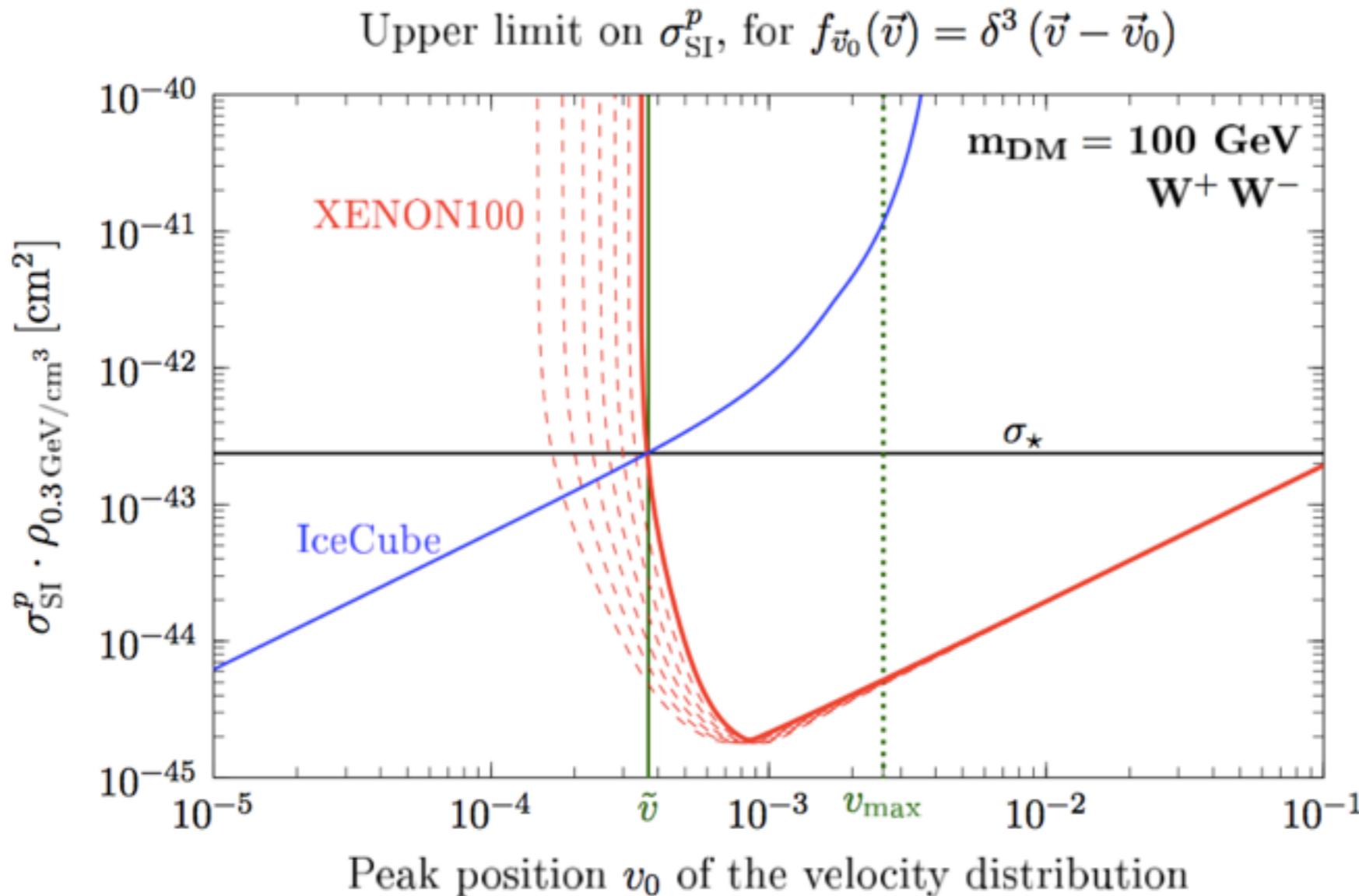
$m_\chi = 30 \text{ GeV}$   
SHM+DD distribution  
Annihilation to  $\nu_\mu \bar{\nu}_\mu$

Also works for other channels...almost everything produces neutrinos

BJK, Fornasa, Green [1410.8051]

# Halo-independent constraints

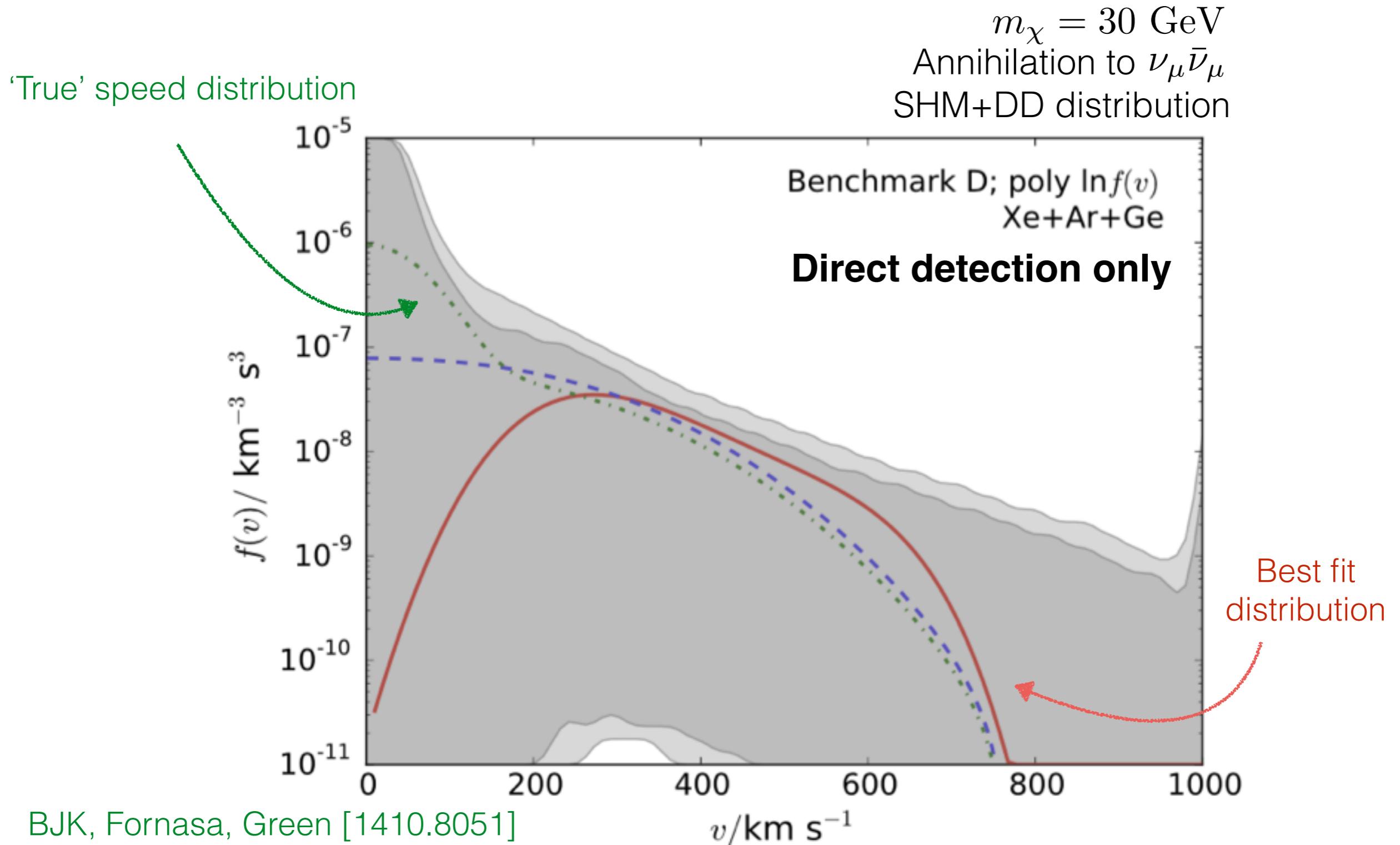
Combining limits from DD and IceCube also allows you to place halo-independent constraints on the DM-nucleon cross section



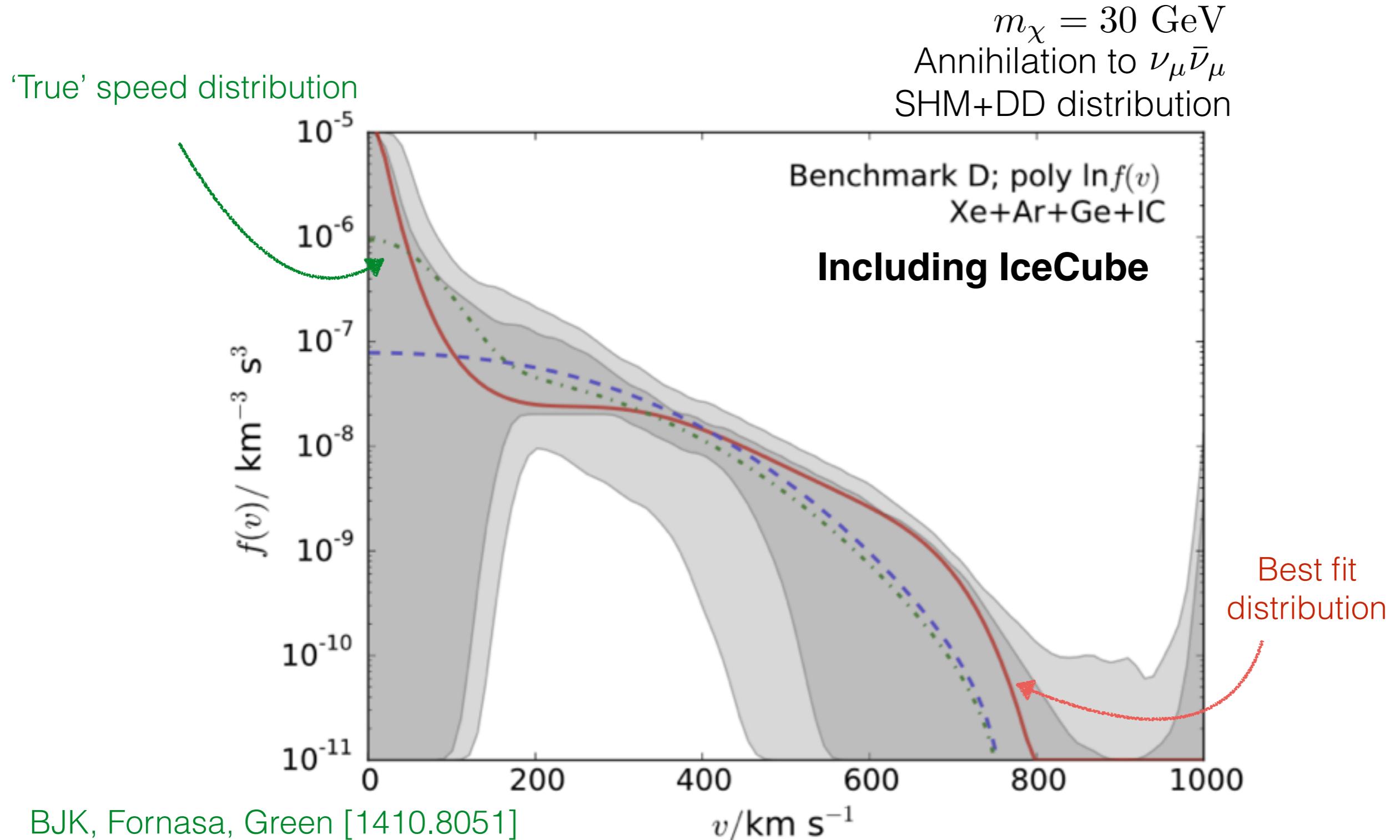
Ferrer et al. [1506.03386]

But see also Blennow et al. [1502.03342]

# Reconstructing the speed distribution



# Reconstructing the speed distribution



Constraints improved, but still difficult to distinguish underlying distributions...

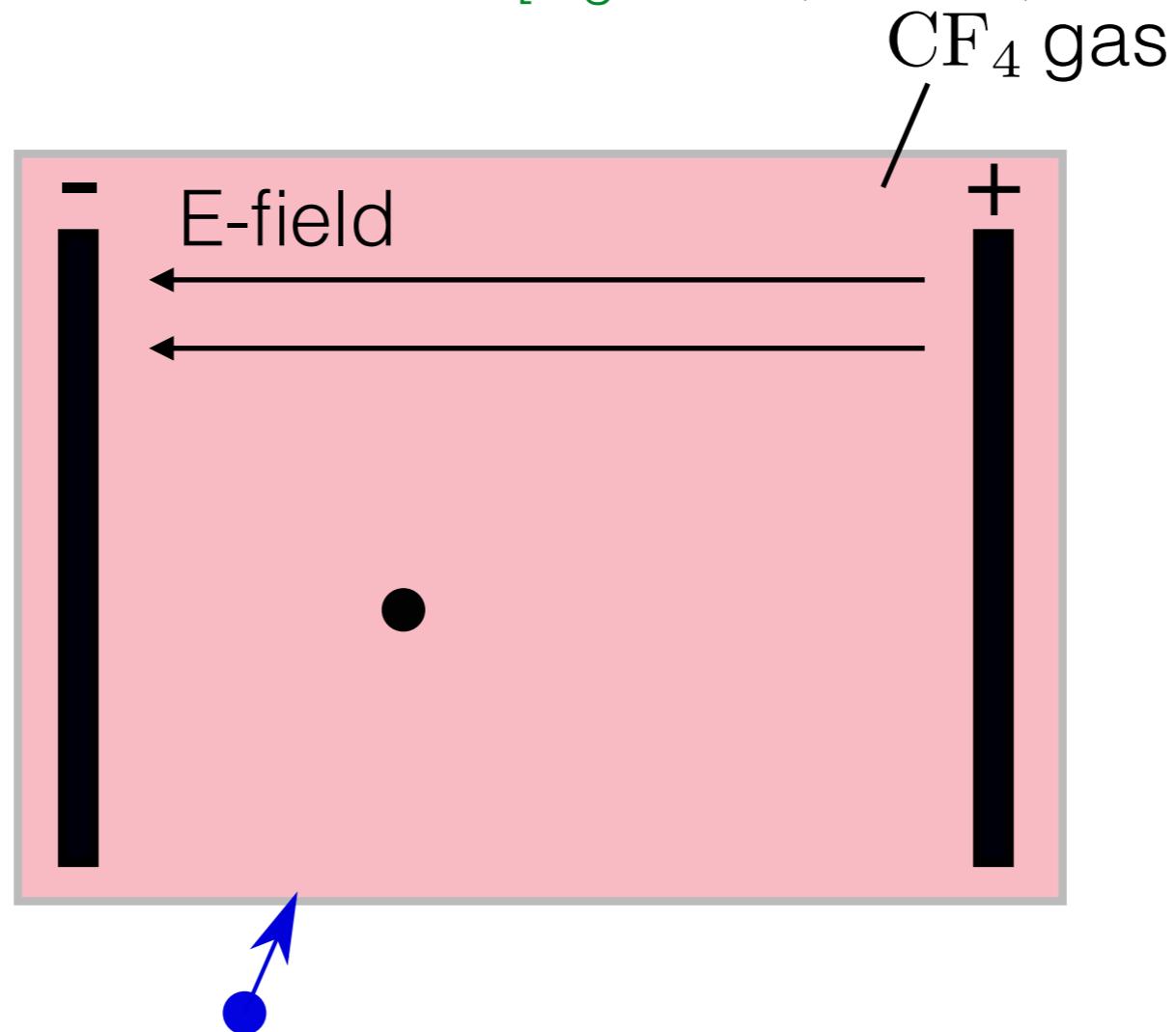
# Directional Detection

# Directional Detection

Try to measure both the energy *and the direction* of the recoil

Most mature technology is the gaseous Time Projection Chamber (TPC)

[e.g. DRIFT, MIMAC, DMTPC, NEWAGE, D3]

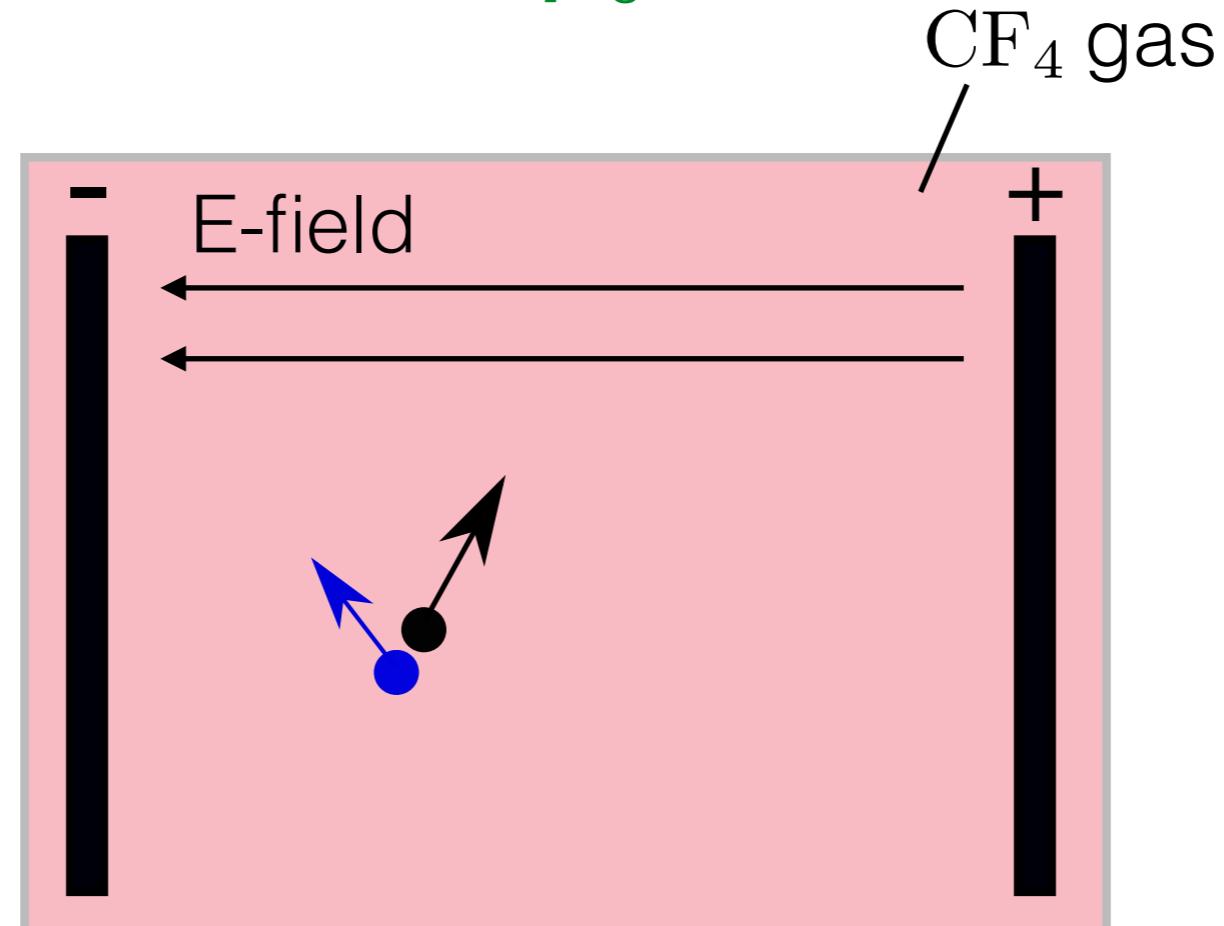


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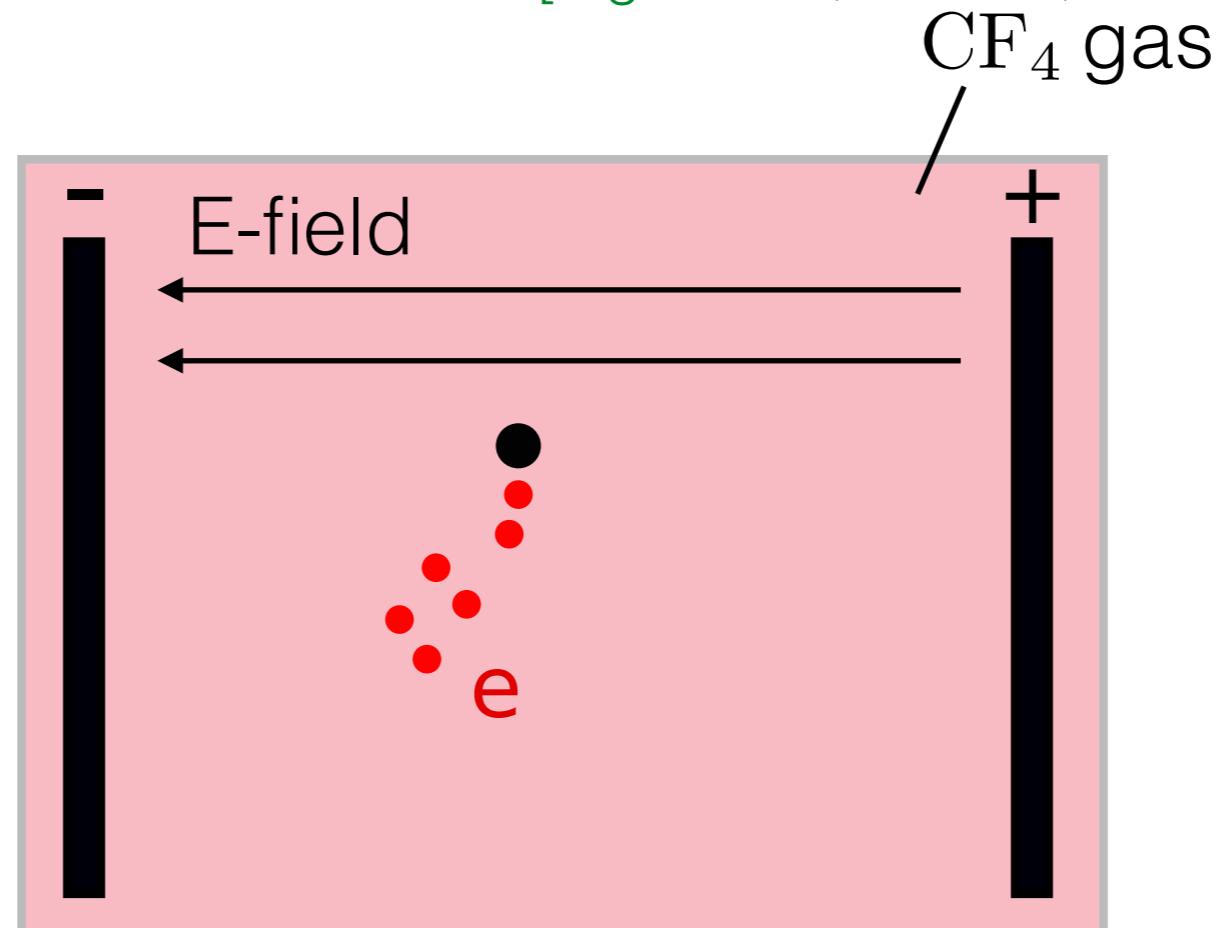


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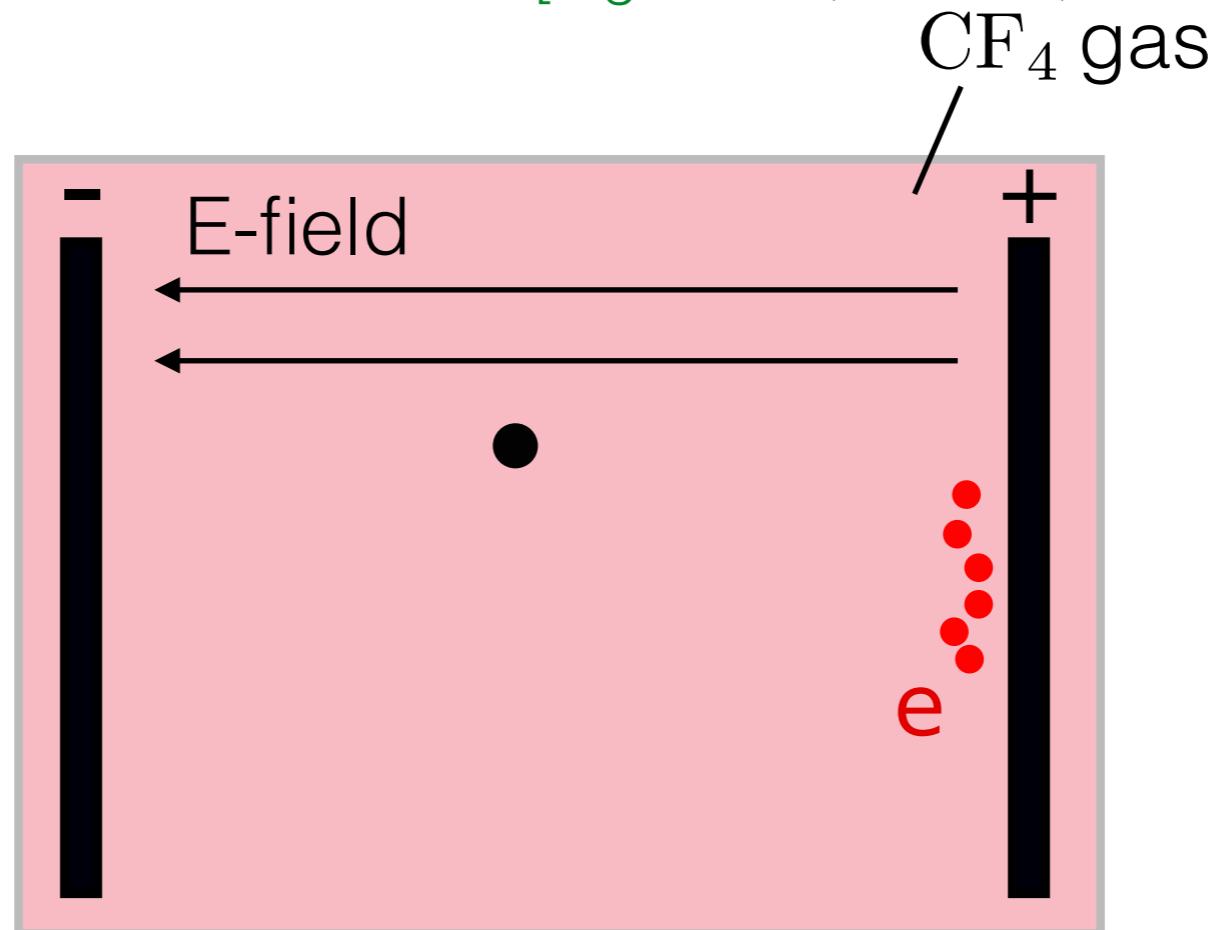


# Directional Detection

Try to measure both the energy *and the direction* of the recoil

Most mature technology is the gaseous Time Projection Chamber (TPC)

[e.g. DRIFT, MIMAC, DMTPC, NEWAGE, D3]



Get x,y of track from distribution of electrons hitting anode

Get z of track from timing of electrons hitting anode

# Directional recoil spectrum

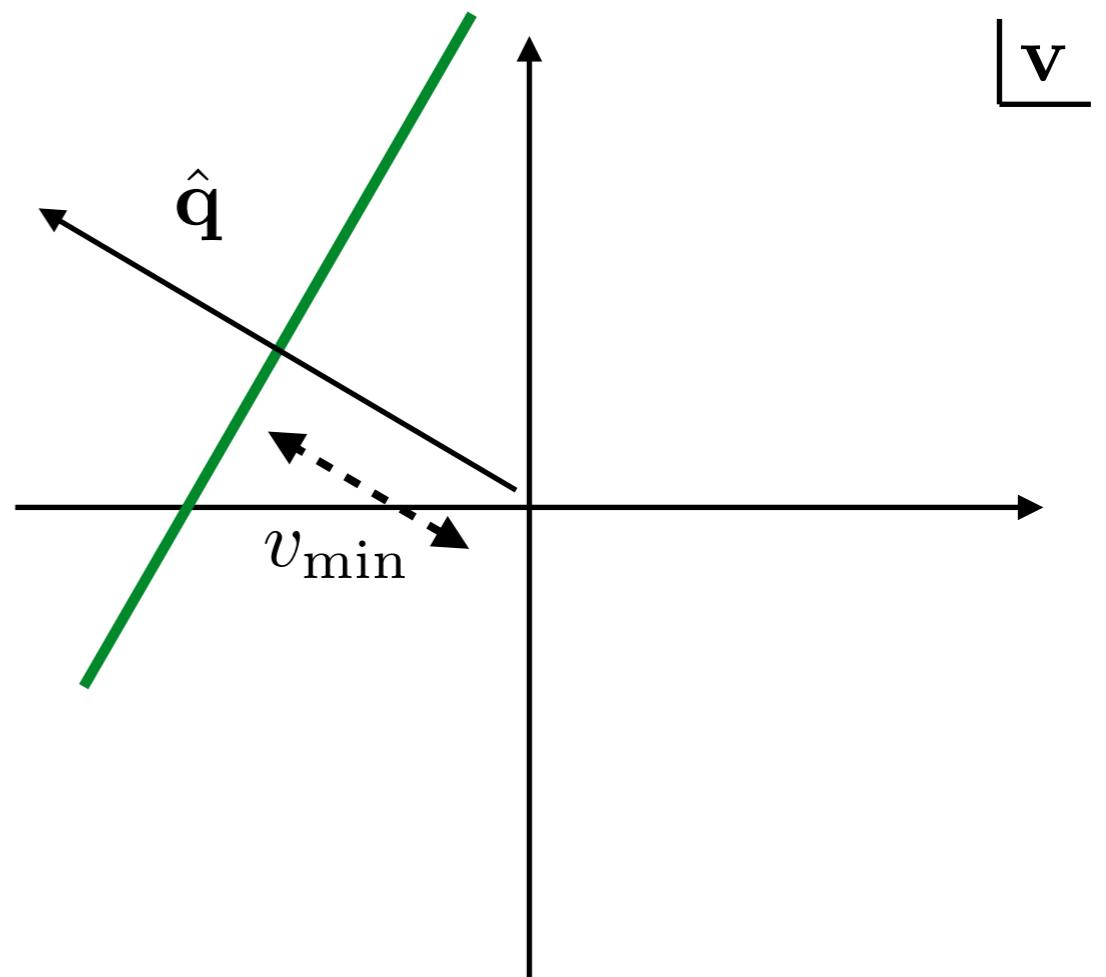
Rate of recoils in direction  $\hat{\mathbf{q}}$ :

$$\frac{dR}{dE_R d\Omega_q} = \frac{\rho_0}{4\pi\mu_{\chi p}^2 m_\chi} \sigma^p C_N F^2(E_R) \hat{f}(v_{\min}, \hat{\mathbf{q}})$$

$$v_{\min} = \sqrt{\frac{m_N E_R}{2\mu_{\chi N}^2}}$$

Radon Transform (RT):

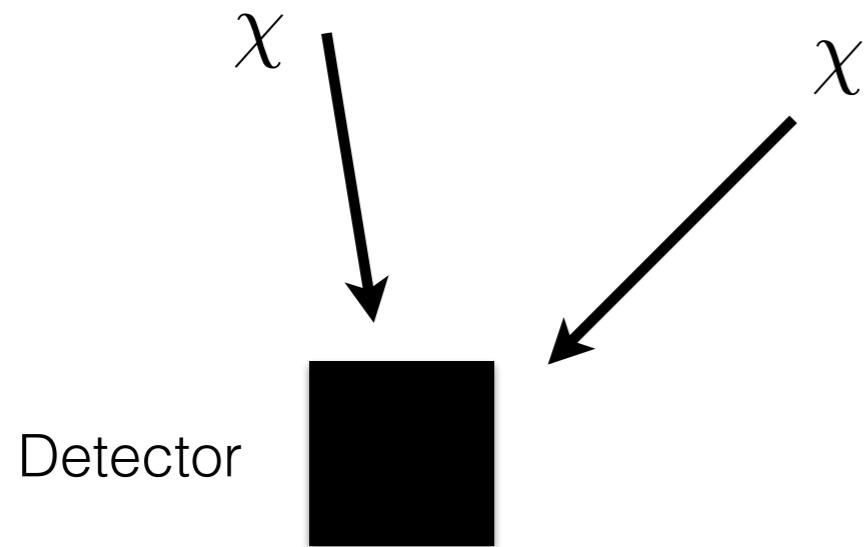
$$\hat{f}(v_{\min}, \hat{\mathbf{q}}) = \int_{\mathbb{R}^3} f(\mathbf{v}) \delta(\mathbf{v} \cdot \hat{\mathbf{q}} - v_{\min}) d^3\mathbf{v}$$



# DM velocity distribution

Experiments which are sensitive to the *direction* of the nuclear recoil can give us information about the full 3-D distribution of the *velocity vector*  $\mathbf{v} = (v_x, v_y, v_z)$ , not just the speed  $v = |\mathbf{v}|$

Mayet et al. [1602.03781]



But, we now have an *infinite* number of functions to parametrise (one for each incoming direction  $(\theta, \phi)$ )!

If we want to parametrise  $f(\mathbf{v})$ , we need some *basis functions* to make things more tractable:

$$f(\mathbf{v}) = f^1(v)A^1(\hat{\mathbf{v}}) + f^2(v)A^2(\hat{\mathbf{v}}) + f^3(v)A^3(\hat{\mathbf{v}}) + \dots$$

# Basis functions

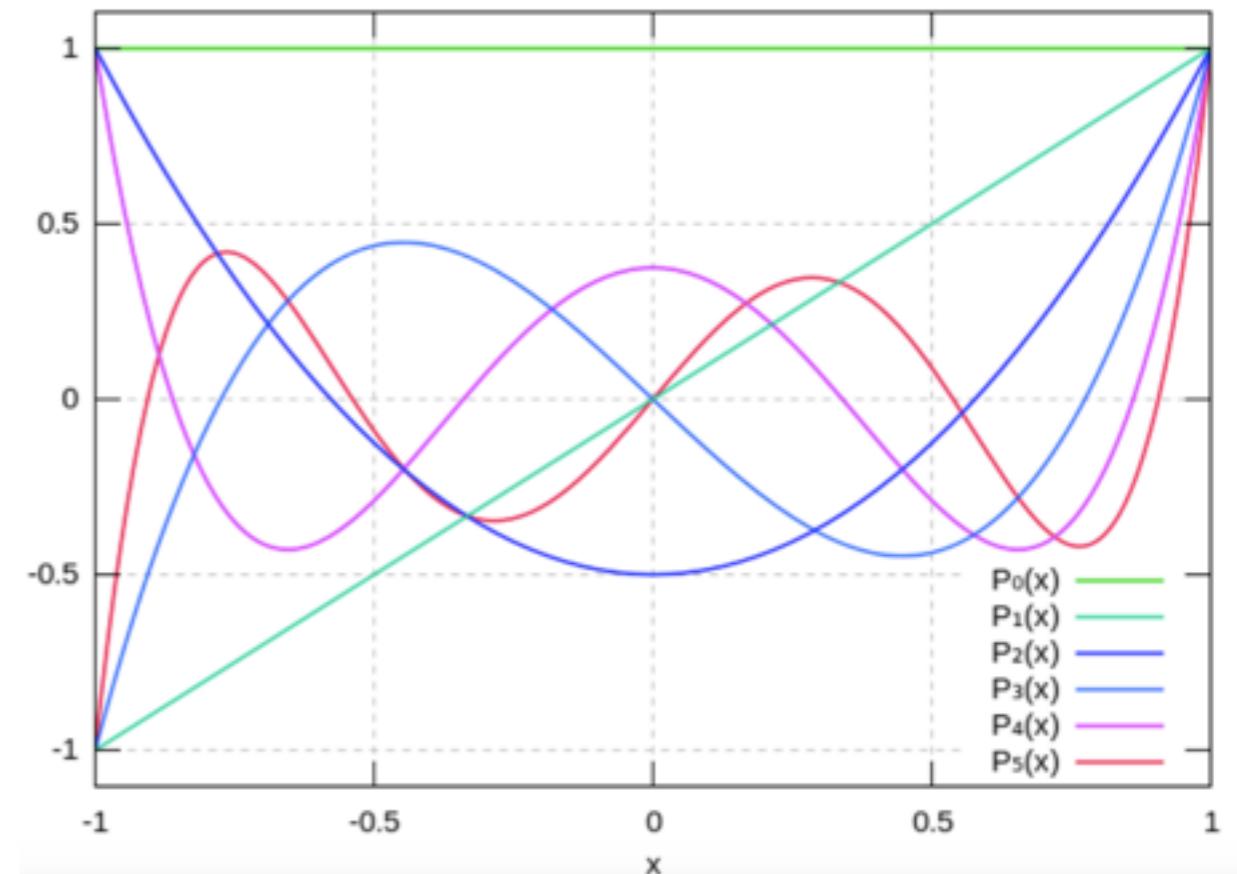
$$f(\mathbf{v}) = f^1(v)A^1(\hat{\mathbf{v}}) + f^2(v)A^2(\hat{\mathbf{v}}) + f^3(v)A^3(\hat{\mathbf{v}}) + \dots$$

One possible basis is spherical harmonics:

Alves et al. [1204.5487], Lee [1401.6179]

$$f(\mathbf{v}) = \sum_{lm} f_{lm}(v)Y_{lm}(\hat{\mathbf{v}})$$

$$Y_{l0}(\cos \theta)$$



However, they are not strictly positive definite.



Physical distribution functions must be positive!

# A discretised velocity distribution

Divide the velocity distribution into  $N = 3$  angular bins...

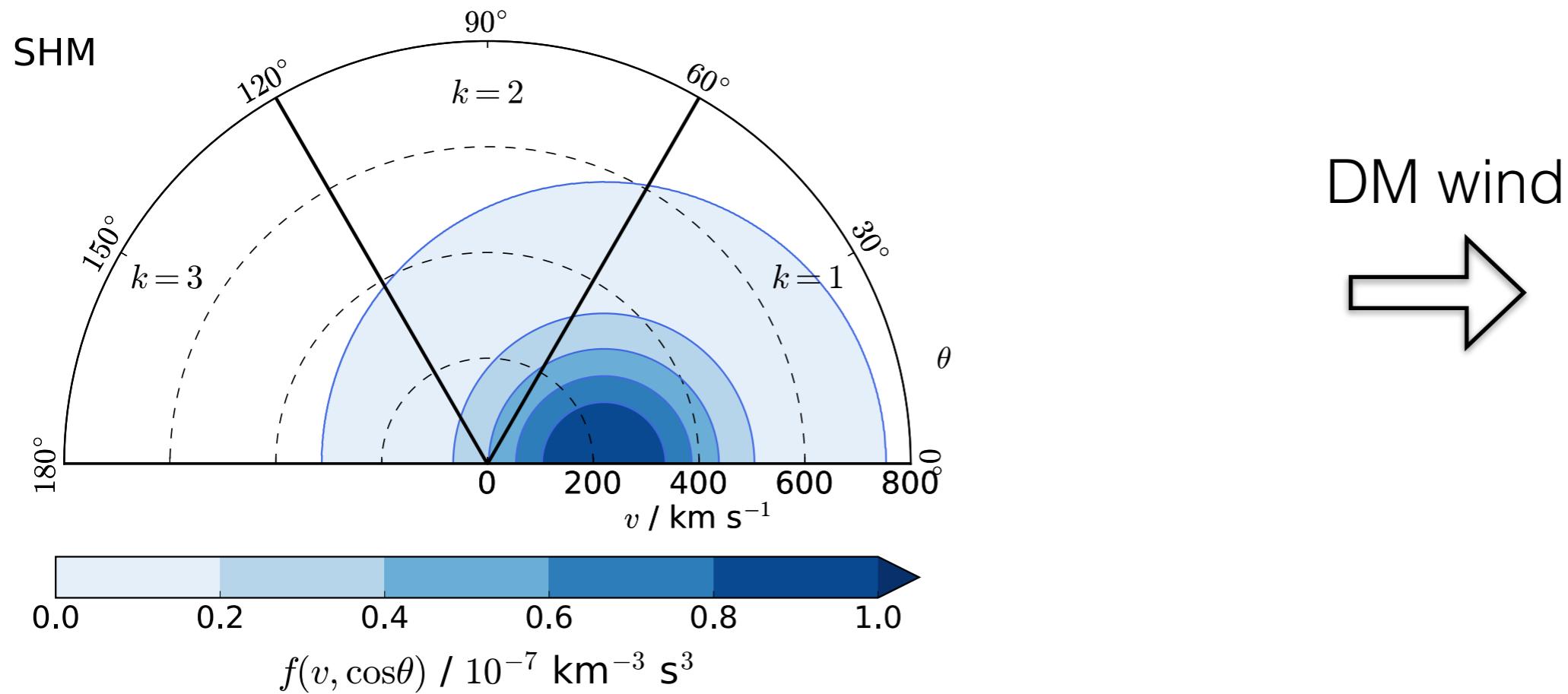
$$f(\mathbf{v}) = f(v, \cos\theta, \phi) = \begin{cases} f^1(v) & \text{for } \theta \in [0^\circ, 60^\circ] \\ f^2(v) & \text{for } \theta \in [60^\circ, 120^\circ] \\ f^3(v) & \text{for } \theta \in [120^\circ, 180^\circ] \end{cases}$$

BJK [1502.04224]

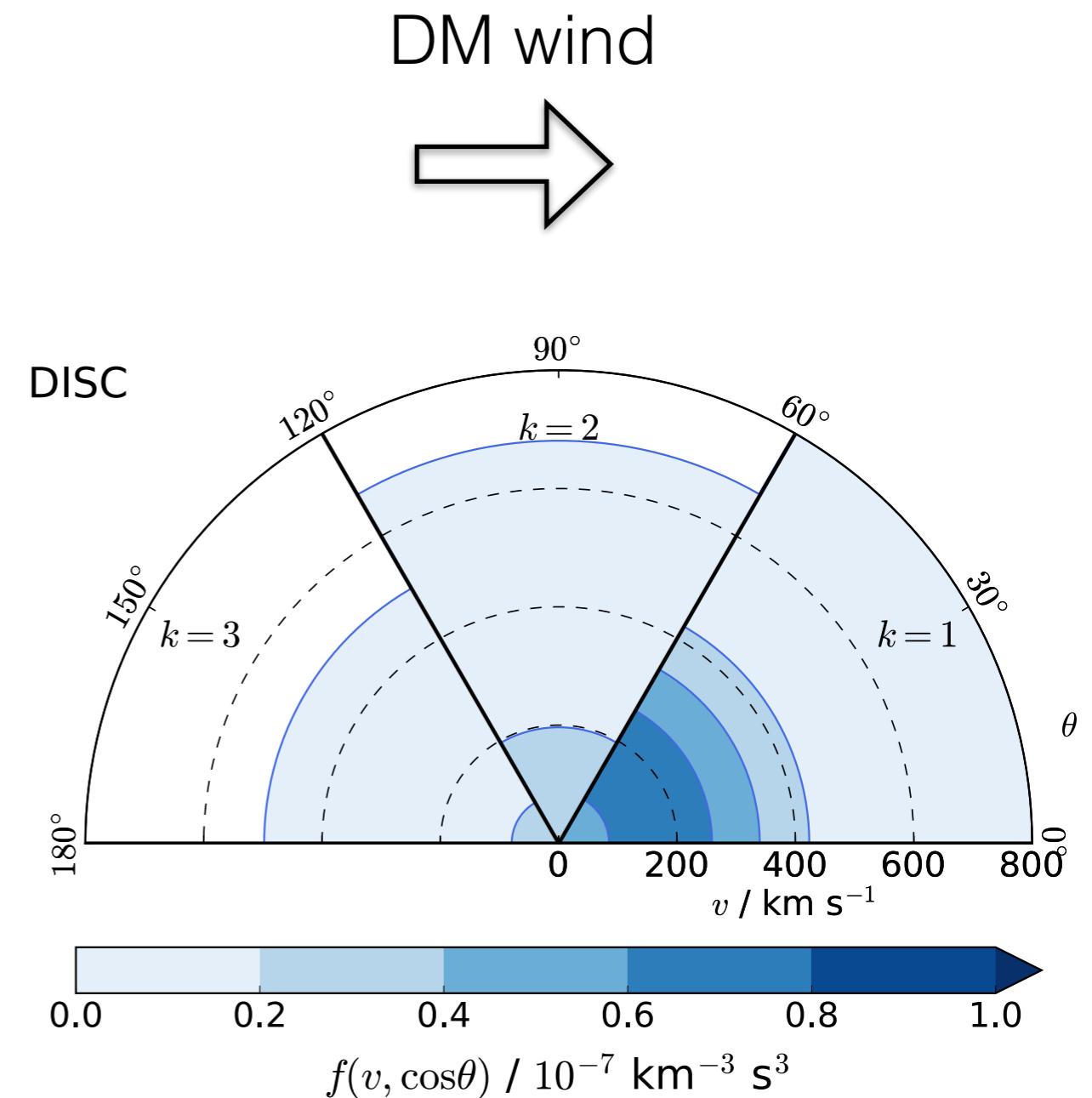
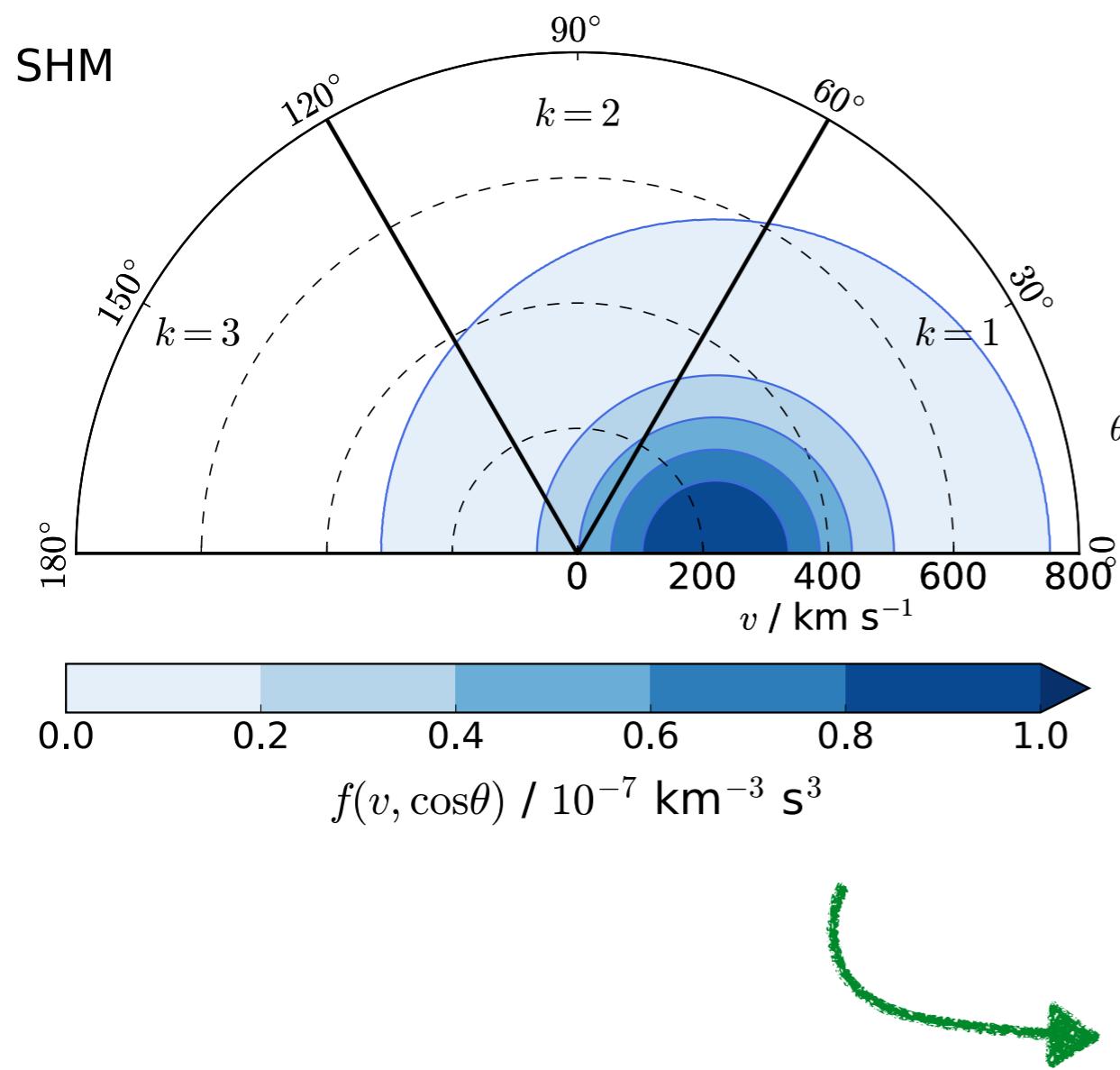
...and then parametrise  $f^k(v)$  within each angular bin  
(using the parametrisation we've already discussed)...

Calculating the event rate from such a distribution (especially for arbitrary  $N$ ) is non-trivial. But not impossible.

# An example: the SHM

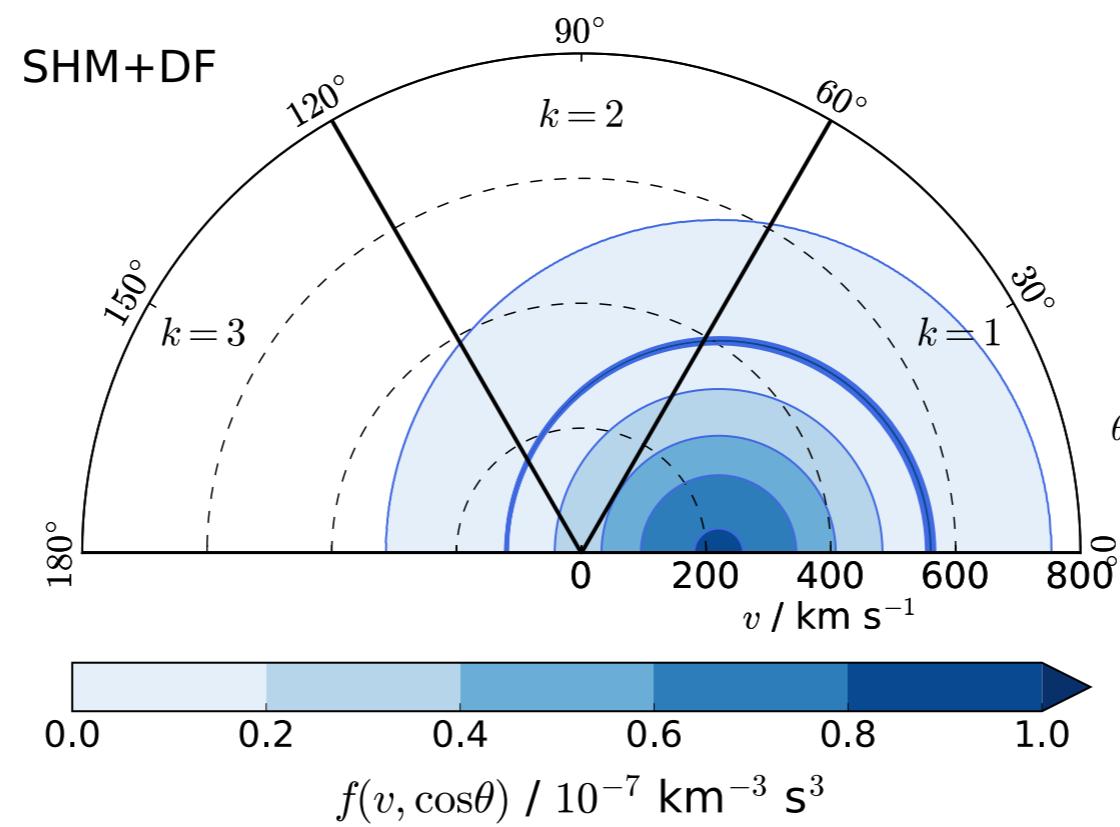
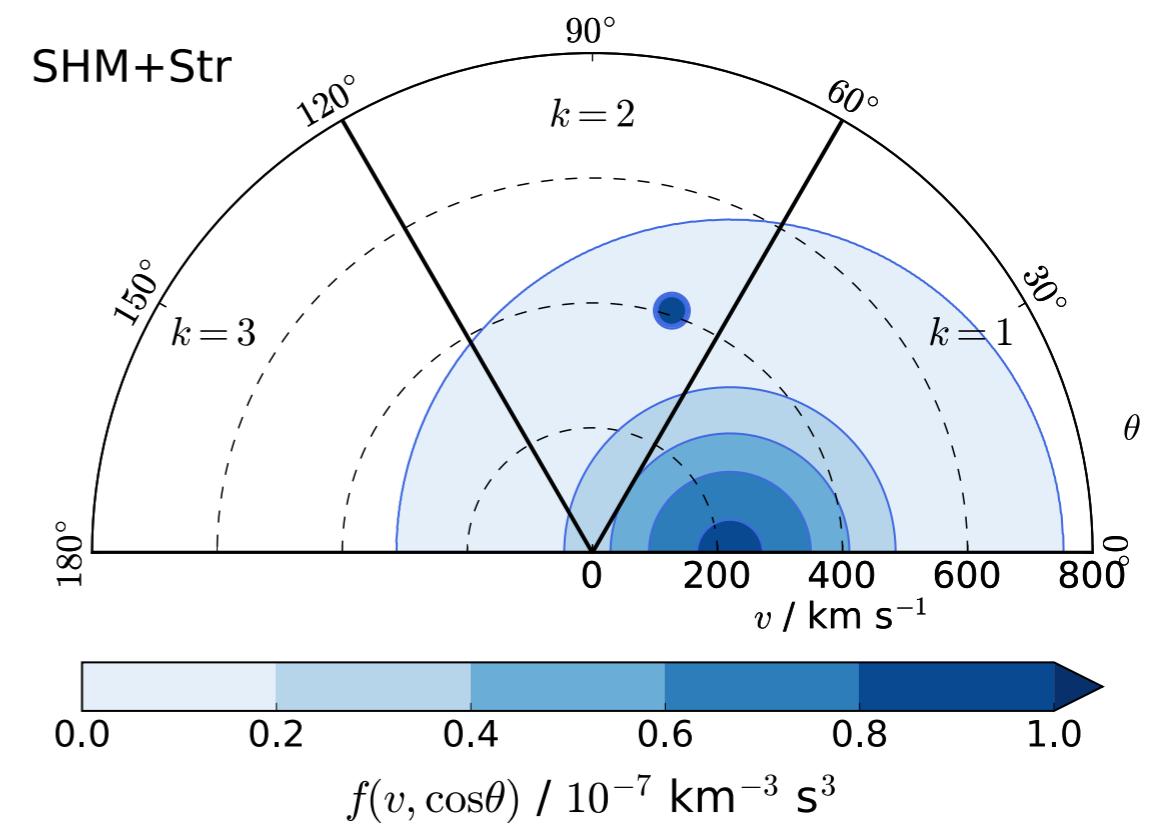
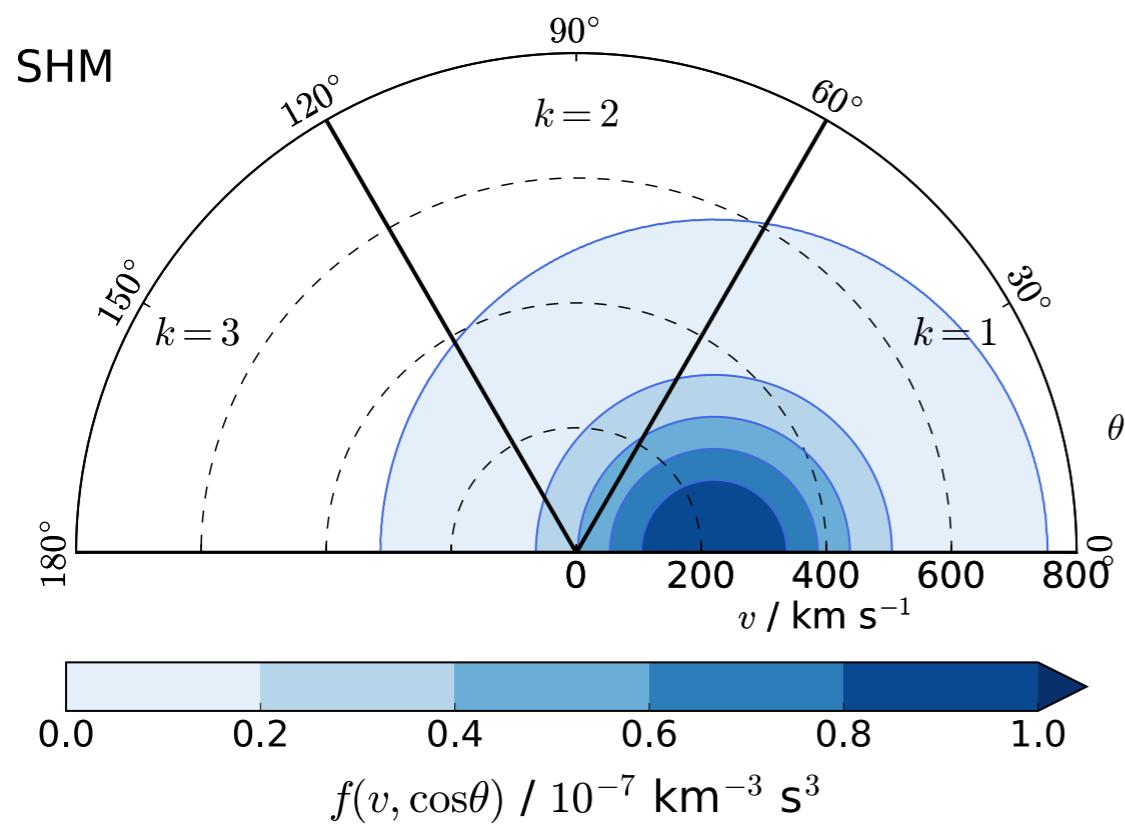


# An example: the SHM



But how well will this work?

# Benchmarks



# Reconstructions

BJK, CAJ O'Hare [1609.08630]

For a single particle physics benchmark  $(m_\chi, \sigma^p)$ , generate mock data in two *ideal* future directional detectors: Xenon-based [1503.03937] and Fluorine-based [1410.7821]

Then fit to the data ( $\sim 1000$  events) using 3 methods:

*Method A:  
Best Case*

Assume underlying velocity distribution is known exactly.

Fit  $m_\chi, \sigma_p$

*Method B:  
Reasonable Case*

Assume functional form of underlying velocity distribution is known.

Fit  $m_\chi, \sigma_p$  and theoretical parameters

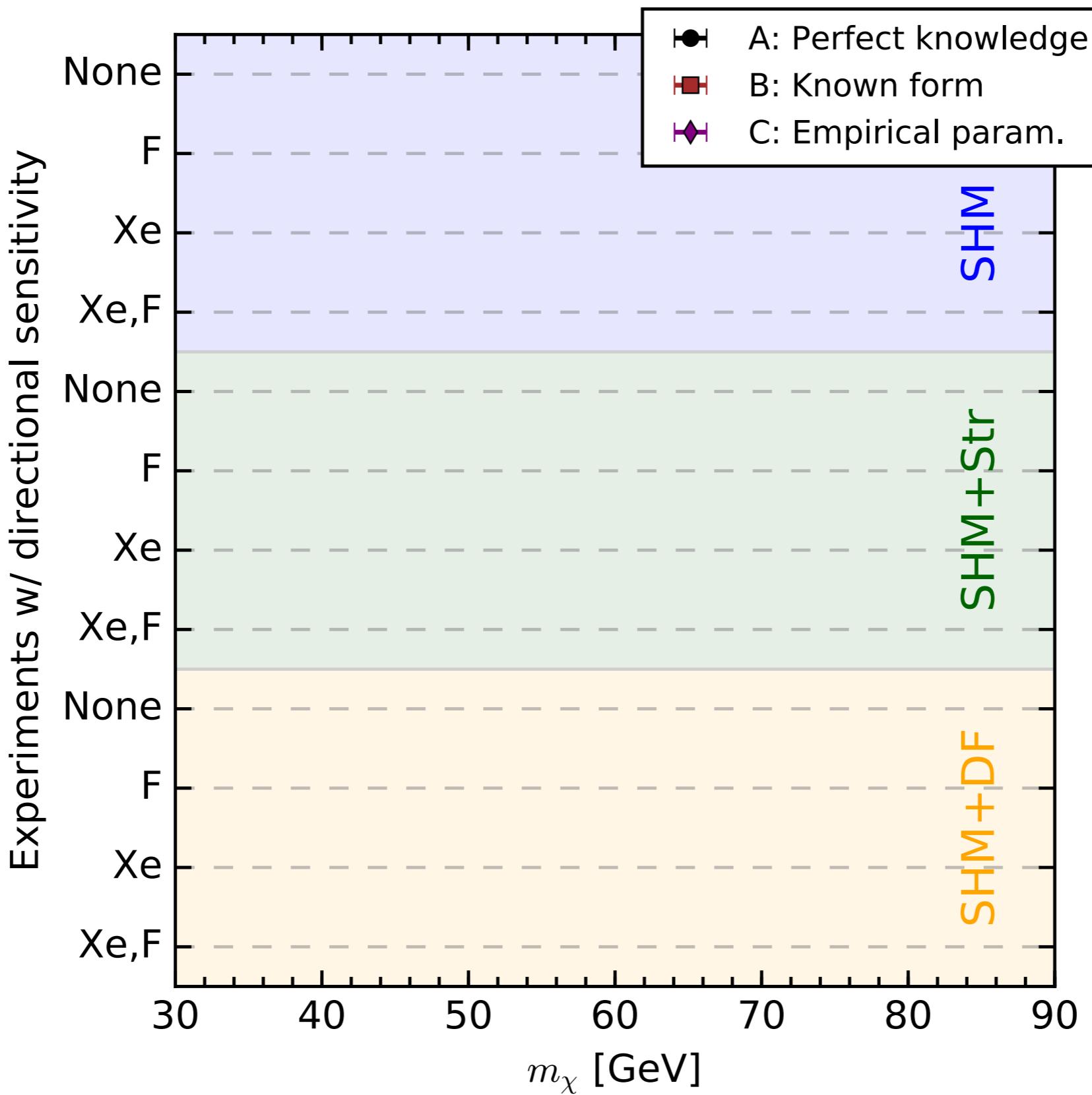
*Method C:  
Worst Case*

Assume nothing about the underlying velocity distribution.

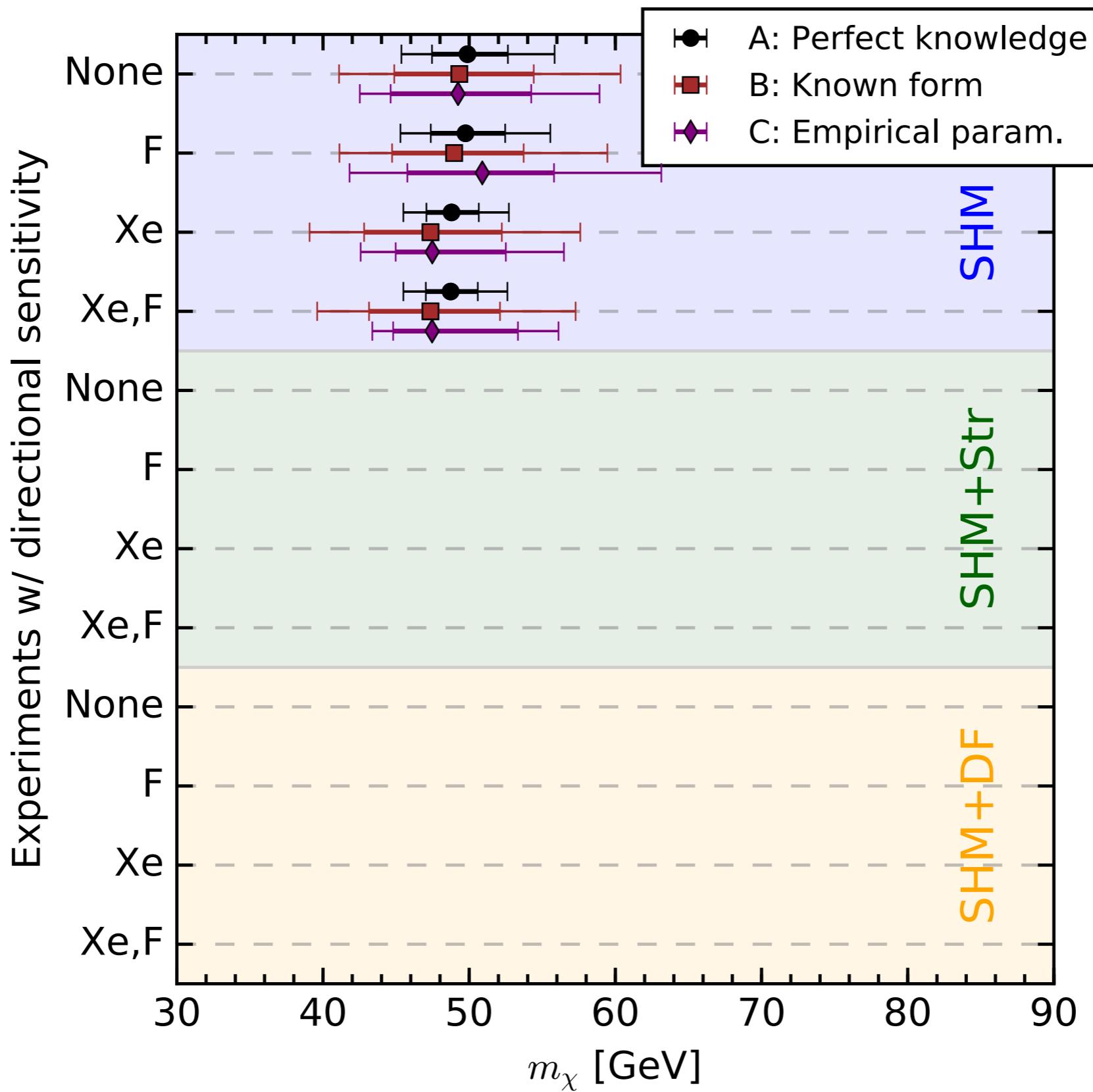
Fit  $m_\chi, \sigma_p$  and empirical parameters

Lee at al. [1202.5035]  
Billard et al. [1207.1050]

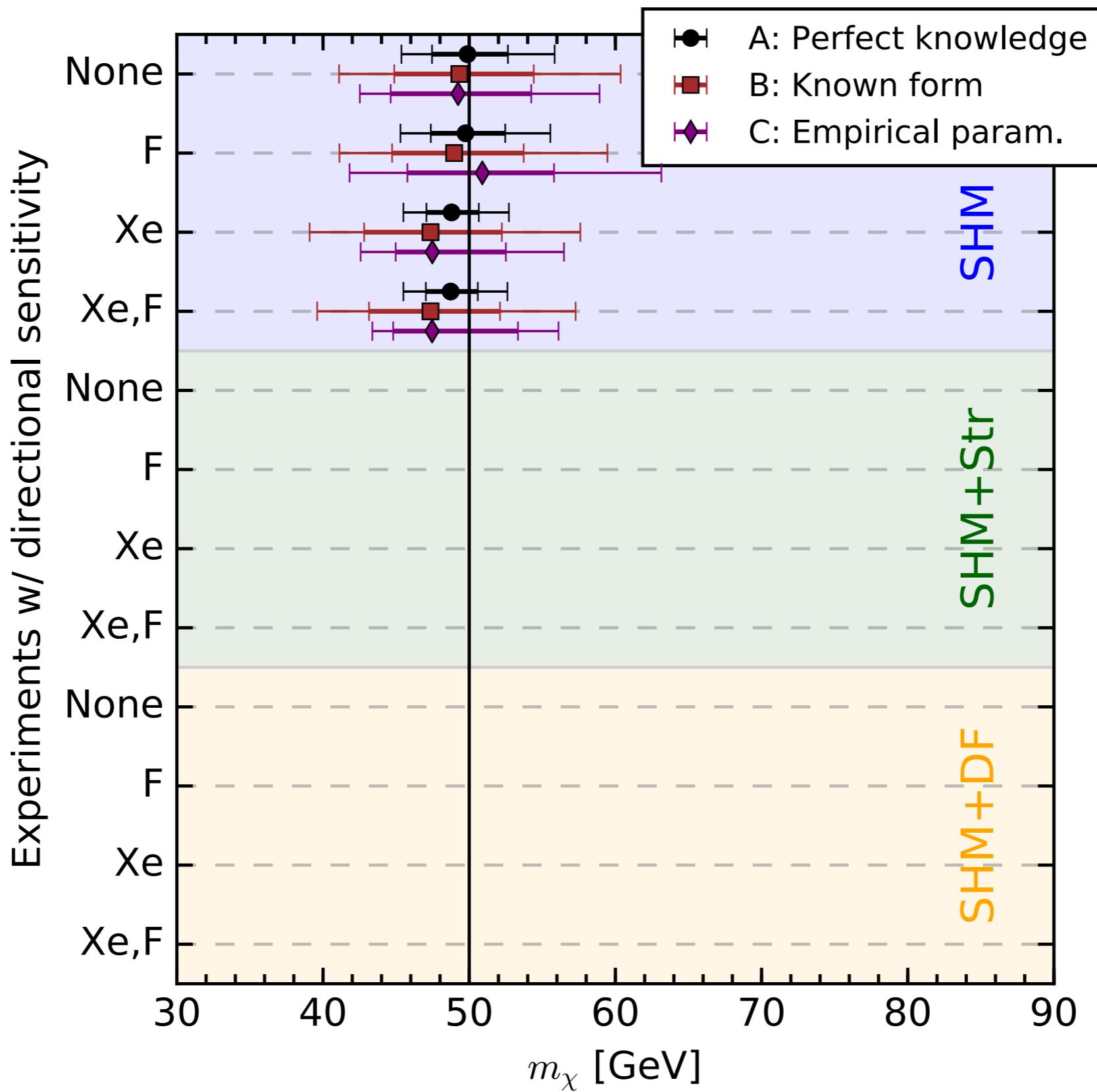
# Reconstructing the DM mass



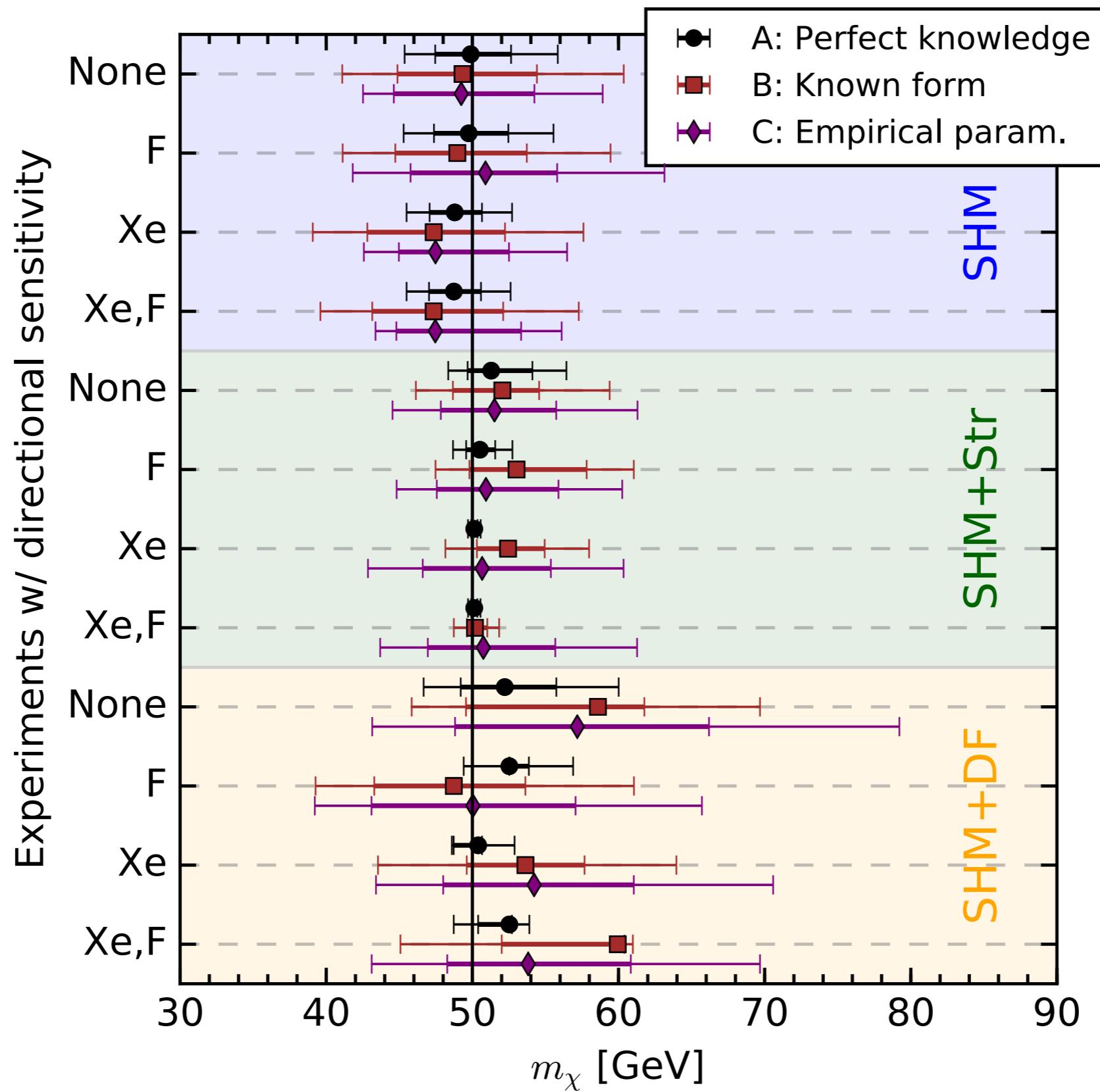
# Reconstructing the DM mass



# Reconstructing the DM mass



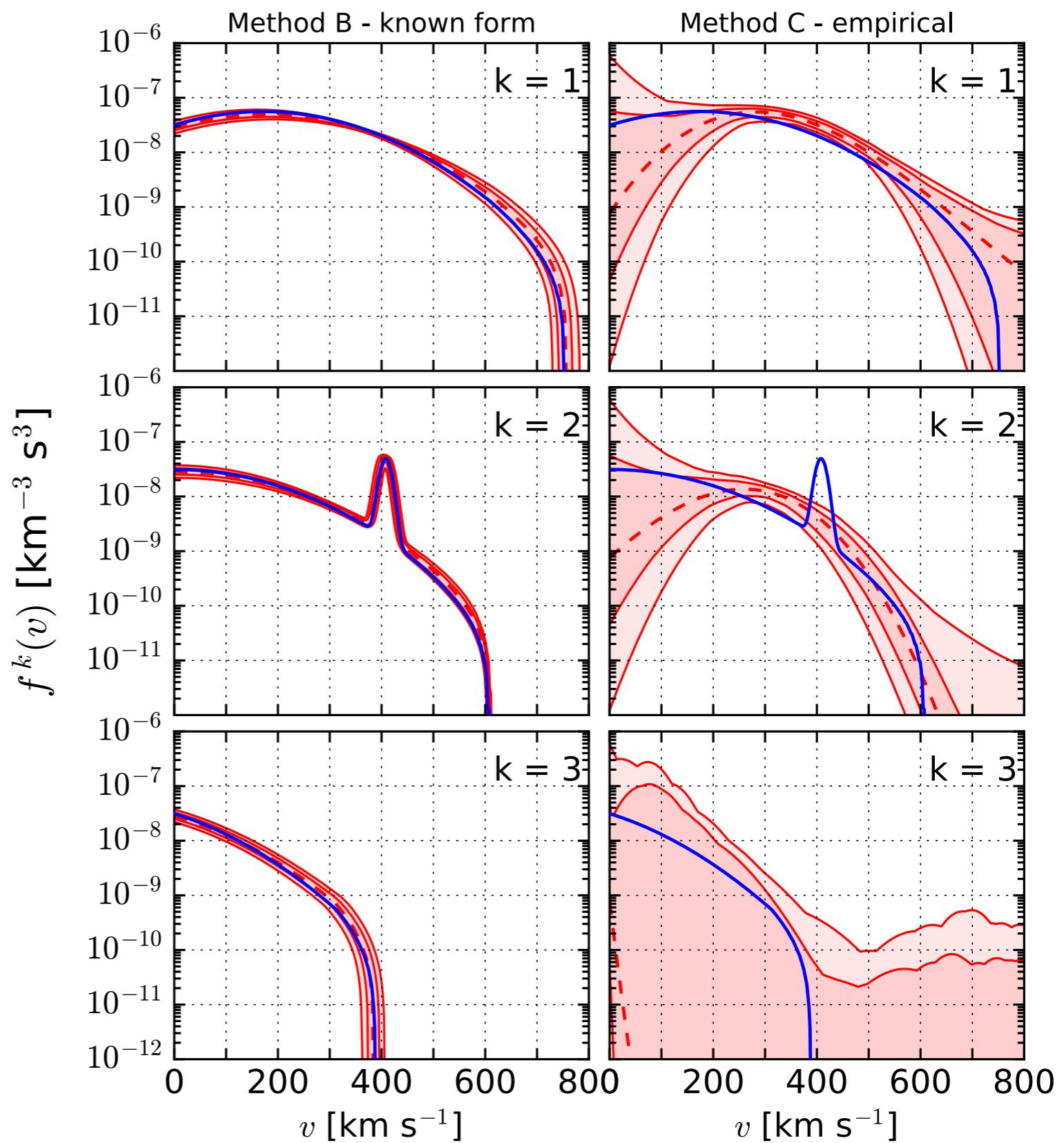
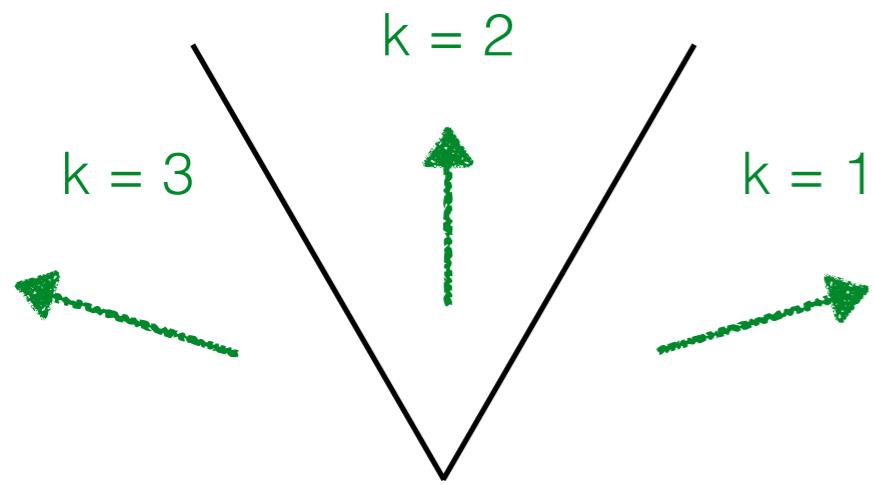
# Reconstructing the DM mass



# Shape of the velocity distribution

SHM+Stream distribution  
with directional sensitivity in  
Xe and F

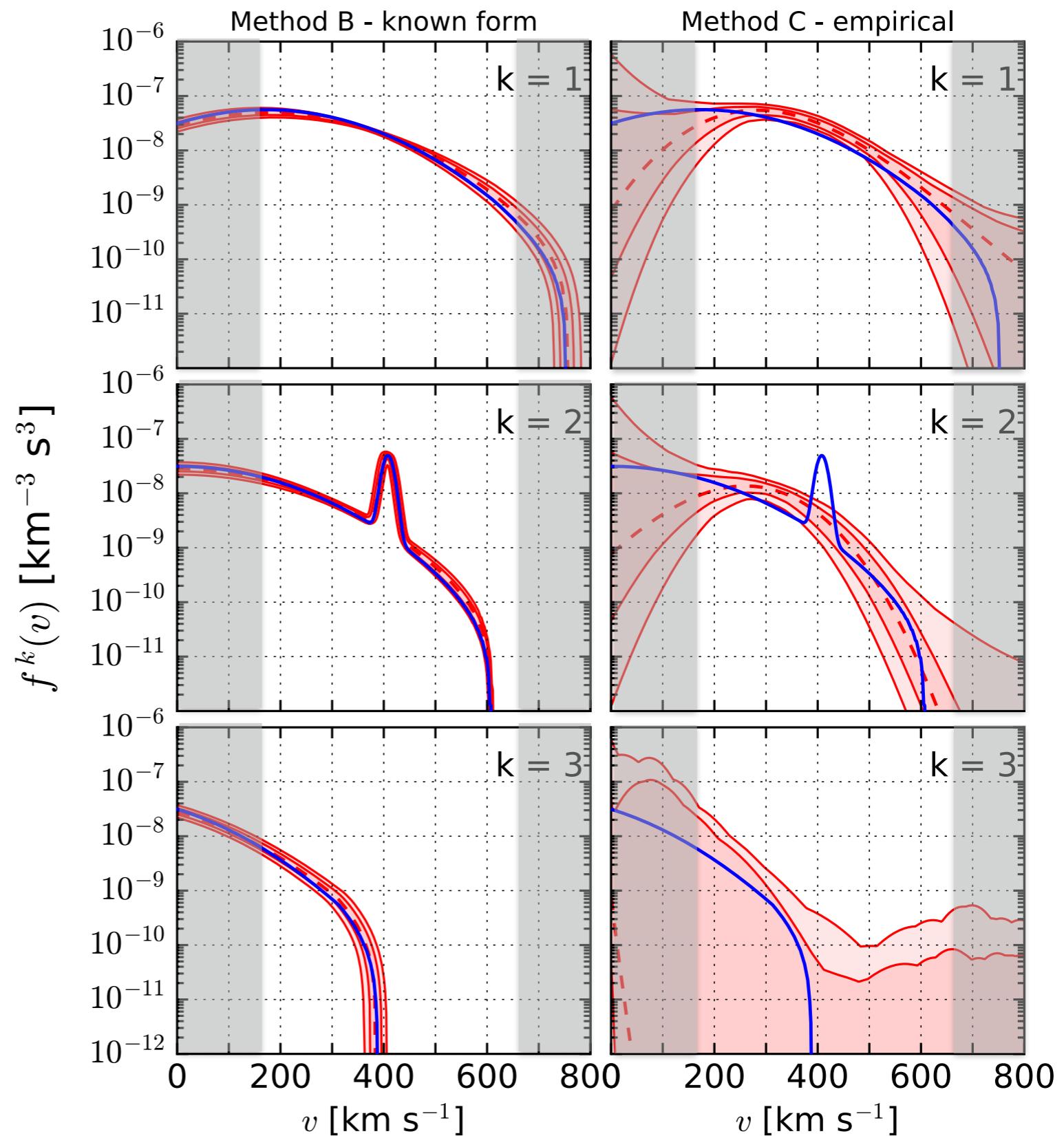
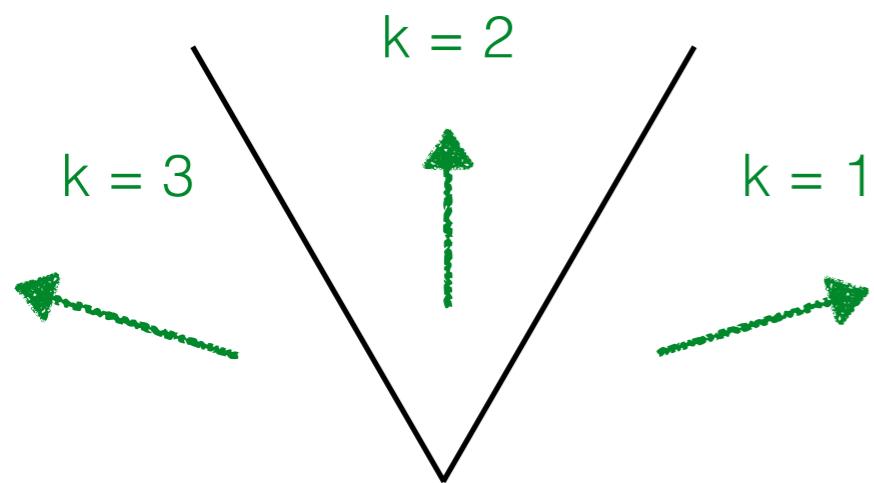
'True' velocity distribution —————  
Best fit distribution - - -  
(+68% and 95% intervals)



# Shape of the velocity distribution

SHM+Stream distribution  
with directional sensitivity in  
Xe and F

'True' velocity distribution —————  
Best fit distribution - - - - -  
(+68% and 95% intervals)



# Velocity parameters

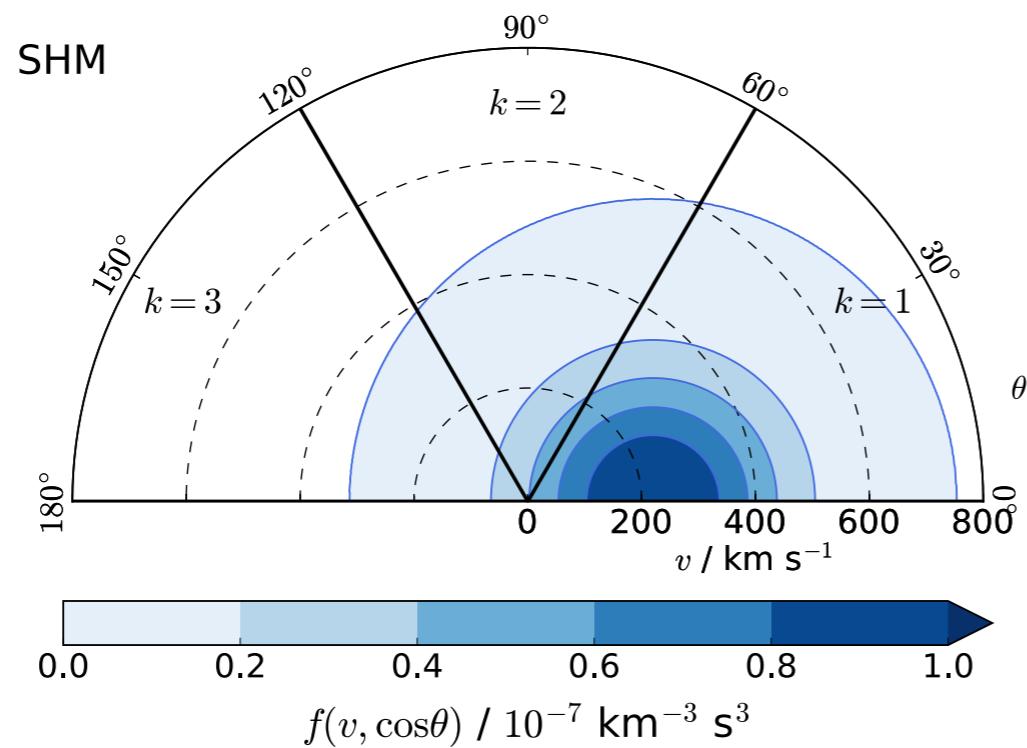
In order to compare distributions, calculate some derived parameters:

Average DM velocity  
parallel to Earth's motion

$$\rightarrow \langle v_y \rangle = \int dv \int_0^{2\pi} d\phi \int_{-1}^1 d \cos \theta (v \cos \theta) v^2 f(\mathbf{v})$$

Average DM velocity  
transverse to Earth's motion

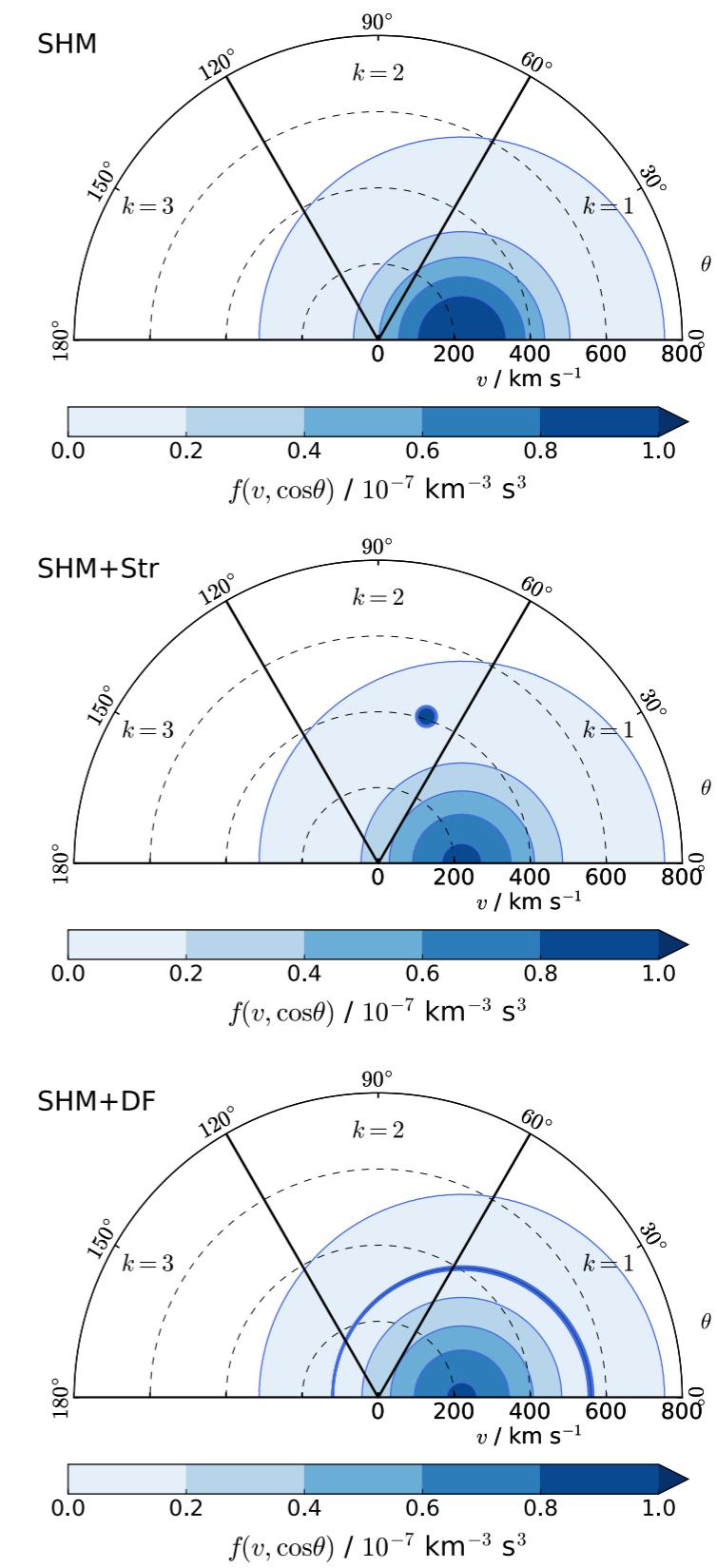
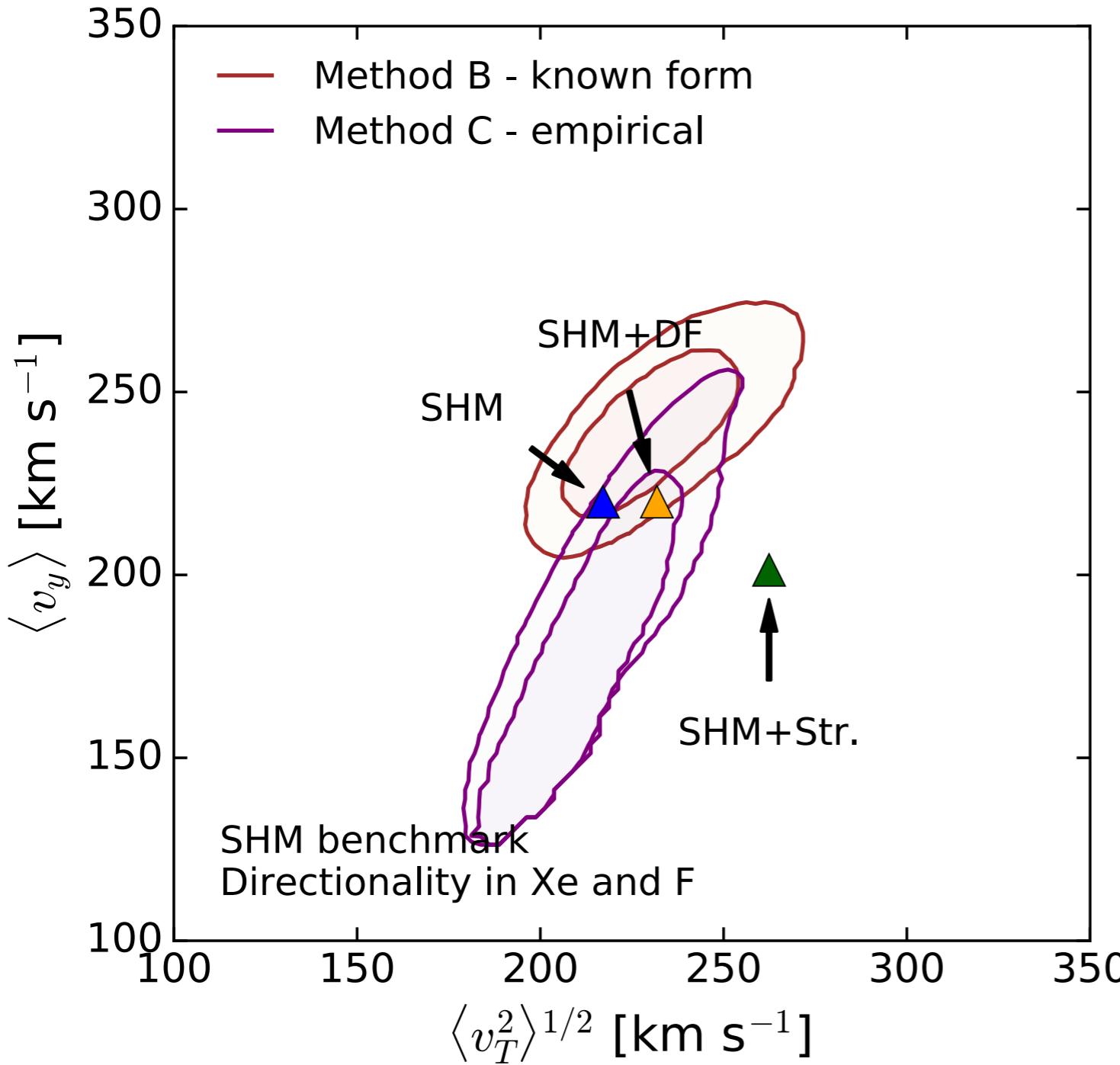
$$\rightarrow \langle v_T^2 \rangle = \int dv \int_0^{2\pi} d\phi \int_{-1}^1 d \cos \theta (v^2 \sin^2 \theta) v^2 f(\mathbf{v})$$



$$\begin{array}{c} \uparrow \\ \langle v_T^2 \rangle^{1/2} \\ \longrightarrow \langle v_y \rangle \end{array}$$

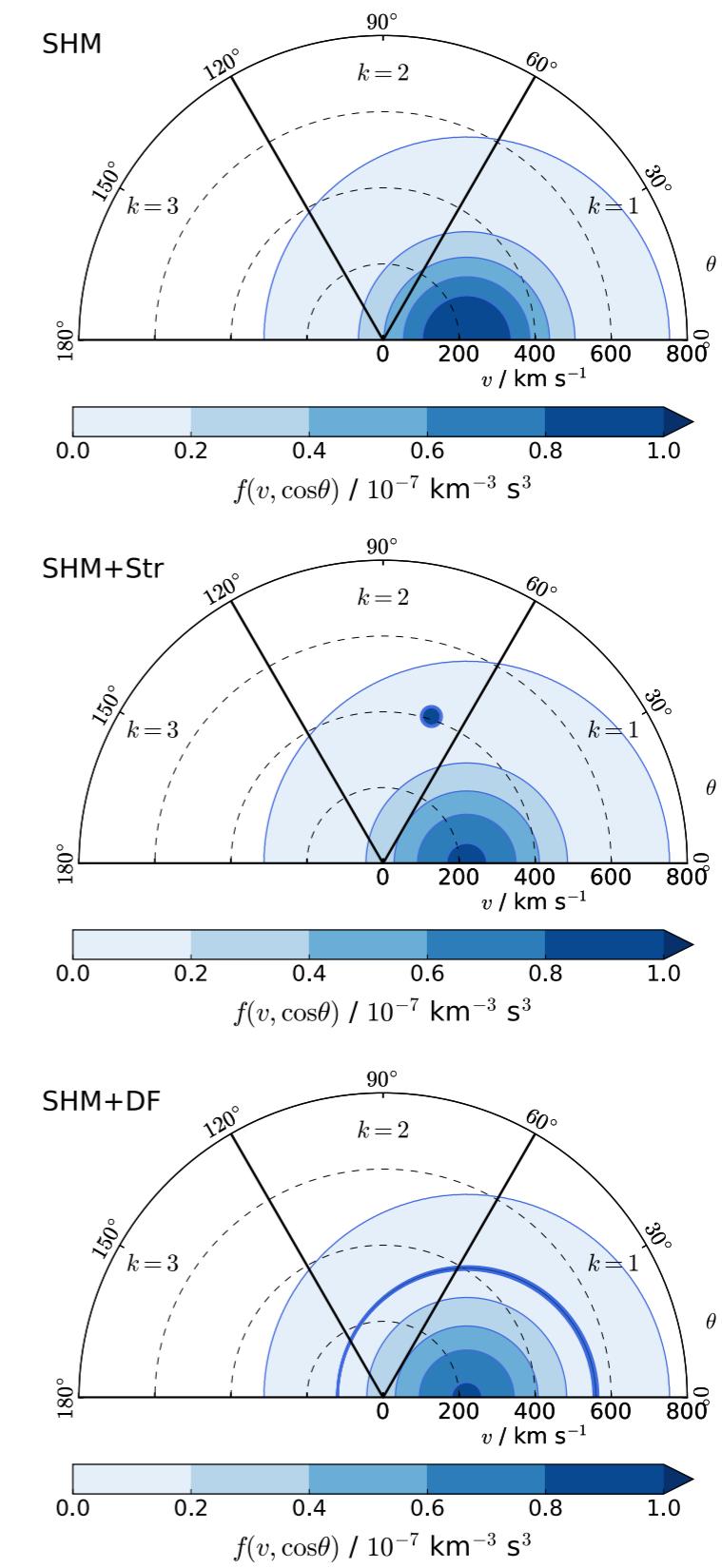
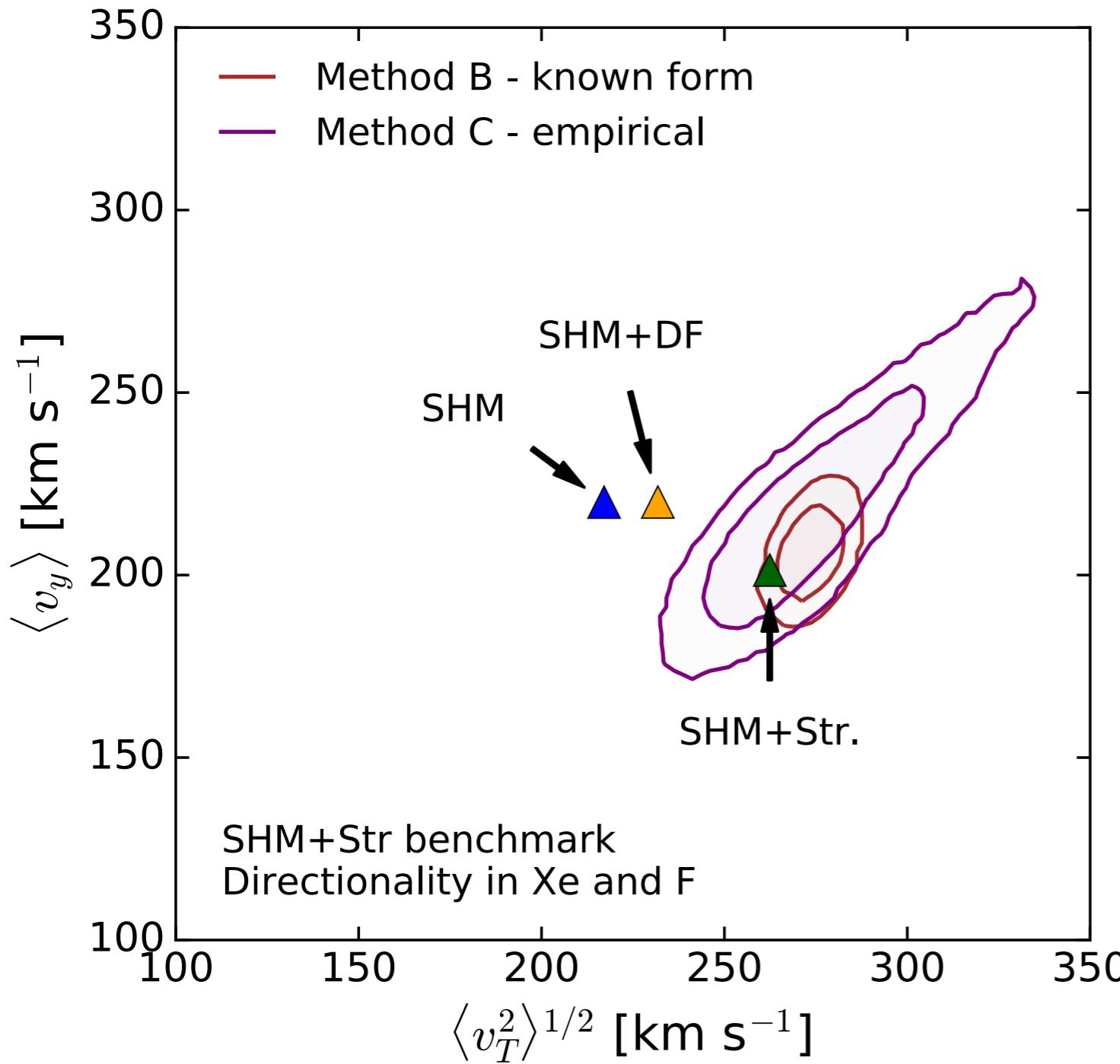
# Comparing distributions

Input distribution: SHM



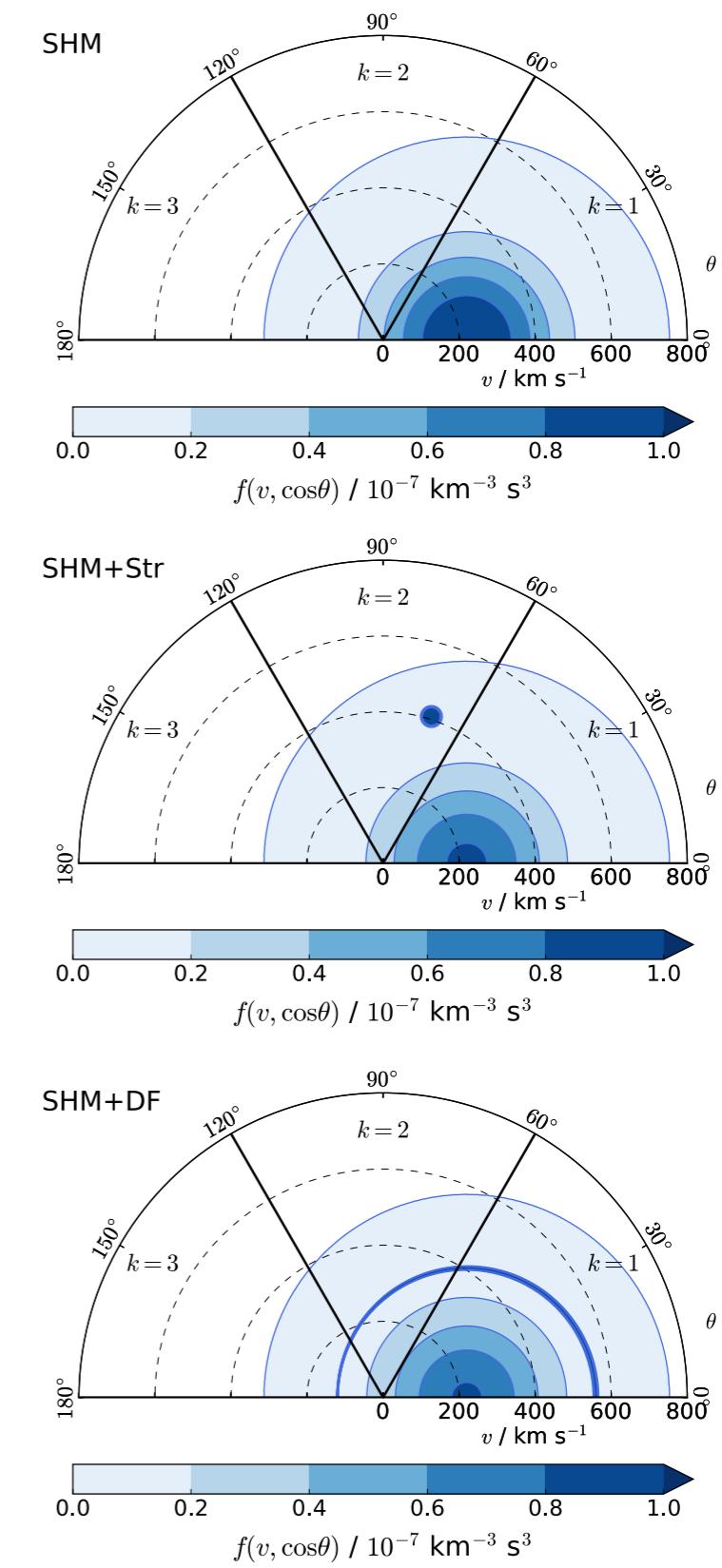
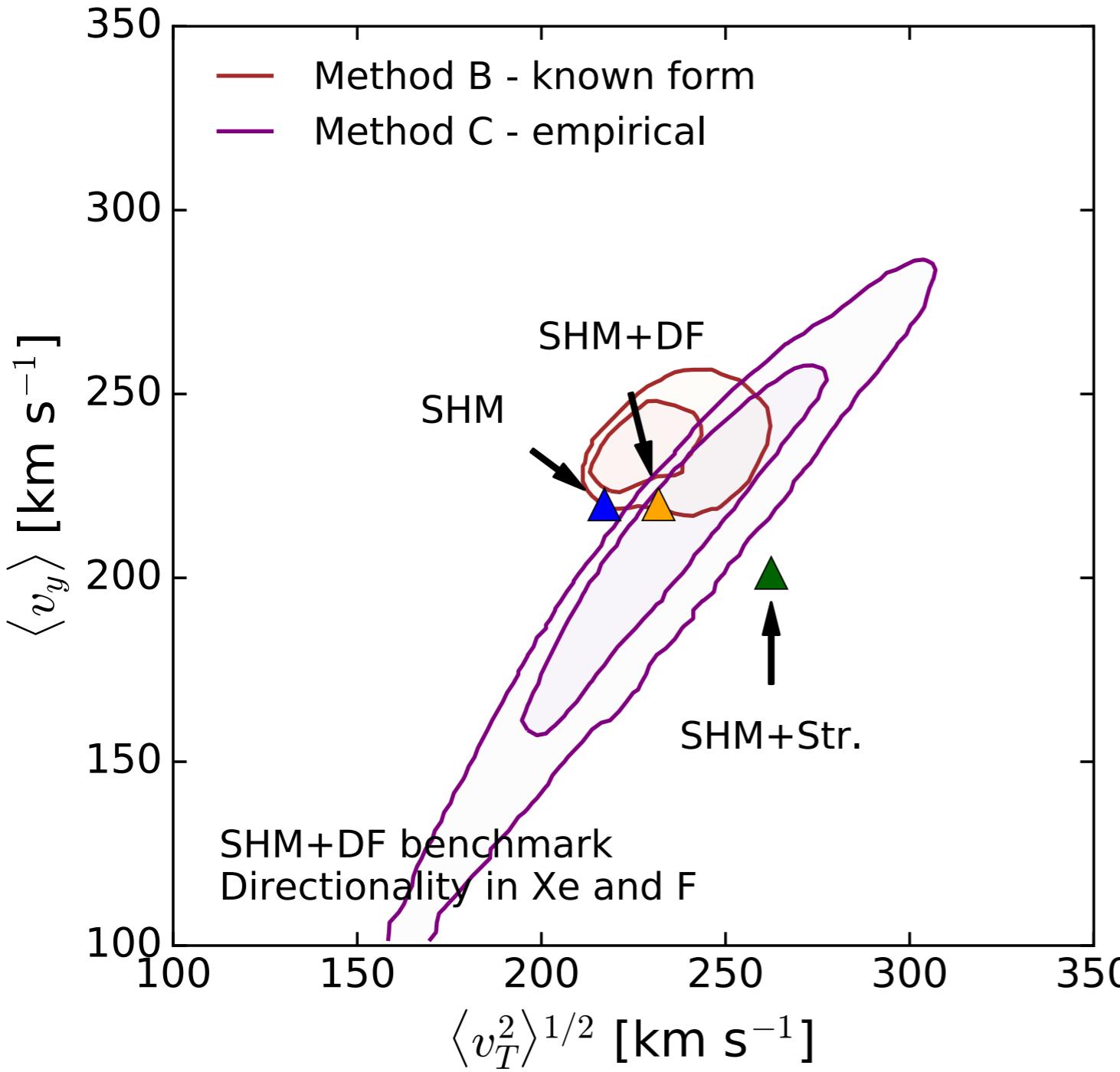
# Comparing distributions

Input distribution: SHM + Stream



# Comparing distributions

Input distribution: SHM + Debris Flow



# The strategy

## In case of signal break glass

Perform parameter estimation using two methods:  
'known' functional form vs. empirical parametrisation

→ Compare reconstructed particle parameters

Calculate derived parameters (such as  $\langle v_y \rangle$  and  $\langle v_T^2 \rangle^{1/2}$ )

→ Check for consistency with SHM

In case of inconsistency, look at reconstructed shape of  $f(v)$

→ Hint towards unexpected structure?

Fantin et al. [1108.4411], Fan et al. [1303.1521]

# Summary

With multiple direct detection experiments, astrophysical uncertainties can be overcome



Reconstruct DM mass *and* shape of speed distribution using a general empirical parametrisation

Information from solar capture and neutrino telescopes tells us about low speed DM particles



Recover full speed distribution & DM-nucleon cross section

Methods can be extended to directional detection without spoiling nice properties



Towards reconstructing full velocity distribution and helping discriminate different halo models

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## THE ERA OF DARK MATTER PARTICLE ASTRONOMY HAS BEGUN

By Jonathan Fakename, Staff Writer

| January 20, 2021 07:01am ET

