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# What controls the variability of CO<sub>2</sub> fluxes in eastern boundary upwelling systems?

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#### **Abstract**

Eastern Boundary Upwelling Systems (EBUS) are governed by alongshore, equatorward winds that force cold, corrosive, and nutrient-enriched waters to the surface. These regions are biologically productive, compensating for the effect of upwelling of carbon-rich waters on pCO2. This leads to a variable mosaic of air-sea CO<sub>2</sub> fluxes with some of the highest flux density globally. This variability also dictates when we might expect the anthropogenic signal of CO<sub>2</sub> flux to become emergent in each system. In this study, we diagnose the physical and biological mechanisms that control historical (1920-2015) CO<sub>2</sub> fluxes in the four major EBUS: the California (CalCS), Humboldt (HumCS), Canary (CanCS), and Benguela Currents (BenCS). We utilize biogeochemical output from the CESM Large Ensemble, a global coupled climate model ensemble that is forced under historical and RCP8.5 radiative forcing. Differences between simulations can be attributed entirely to internal climate variability, as simulations are generated by introducing round-off perturbations to the initial atmospheric temperature. This experimental setup provides us with 34 independent and unique representations of the natural climate system, allowing us to robustly assess variability in CO<sub>2</sub> fluxes. We find that anomalous CO<sub>2</sub> flux in the CalCS and CanCS is most associated with oscillations in subtropical gyres: the North Pacific Gyre Oscillation (NPGO) and the North Atlantic Oscillation (NAO), respectively. The CalCS (CanCS) has anomalous uptake (outgassing) of carbon during the positive phase of the NPGO (NAO). The HumCS responds mainly to El Niño Southern Oscillation (ENSO), with anomalous uptake of CO2 during an El Niño event. Variations in dissolved inorganic carbon (DIC) and sea surface temperatures (SST) are the major contributors to these anomalous fluxes, and are generally driven by changes to upwelling, the mixed layer depth, and biology. A better understanding of the sensitivity of CO<sub>2</sub> fluxes in EBUS to internal climate variability might lead to some short-term predictive skill in the ocean-atmosphere carbon cycle. Skillful prediction would be particularly useful in forecasting and managing the onset of ocean acidification in systems that have a naturally low pH and carbonate ion concentration.

# 1 Introduction

The four major Eastern Boundary Upwelling Systems (EBUS) occur at the eastern edges of subtropical gyres in the Atlantic and Pacific oceans – the California Current System (CalCS), Humboldt Current System (HumCS), Canary Current System (CanCS), and Benguela Current System (BenCS). These regions are characterized by semi-permanent or permanent equatorward winds that drive both coastal Ekman upwelling and curl-driven Ekman pumping within the first 200km of the coastline (Chavez and Messié, 2009). Upwelling delivers deep waters with respired nutrients to the surface, fueling fisheries that are highly productive with respect to the small surface area they cover (Ryther, 1969). This process also supplies waters with an elevated dissolved inorganic carbon (DIC) content, which enhances the partial pressure of carbon dioxide ( $pCO_2$ ) and reduces the pH and carbonate ion concentration. In turn, these systems are naturally sensitive to the impacts of ocean acidification (Hauri et al., 2009; Bakun et al., 2015; Chan et al., 2017).

The carbonate chemistry of EBUS is controlled by a complex interplay of physical and biological processes: entrainment of subsurface waters, horizontal advection, upwelling and vertical mixing, temperature changes, photosynthesis, respiration, and calcium carbonate formation and dissolution (DeGrandpre et al., 1998; King et al., 2007). These terms combine to dictate oceanic  $pCO_2$ , which drives the  $pCO_2$  gradient between the ocean and atmosphere ( $\Delta pCO_2$ ), thus contributing to the magnitude and determining the direction of air-sea  $CO_2$  fluxes.

Despite coastal oceans around the world contributing modestly to the global carbon cycle, they are characterized by a high  $CO_2$  flux density, or the magnitude of air-sea carbon exchange per unit area (Laruelle et al., 2010, 2014; Gruber, 2015). Low-latitude upwelling systems, such as the HumCS and CanCS, tend to be net outgassing systems, due to their relatively warm waters and persistent upwelling. Because of their colder temperatures and active biology, mid-latitude systems, such as the CalCS and BenCS, act as weak  $CO_2$  sinks that can become  $CO_2$  sources during certain seasons (Borges and Frankignoulle, 2002; Hales et al., 2005; Cai et al., 2006; Gregor and Monteiro, 2013). Surface ocean  $pCO_2$  and thus air-sea  $CO_2$  flux in EBUS exhibits high temporal variability with statistical power at seasonal, sub-seasonal, and interannual time



scales (Friederich et al., 2002; González-Dávila et al., 2009; Leinweber et al., 2009; Evans et al., 2011; Turi et al., 2014). Although the pronounced temporal variability of CO<sub>2</sub> fluxes in EBUS has been documented by a number of studies, little work has been done to associate it directly with internal climate variability.

It has long been recognized that EBUS have intense variations in biology that are coupled to large-scale physical variability (*e.g.*, Chelton et al., 1982; Barber and Chavez, 1983; Barber and Chávez, 1986). Studies have associated this variability with major climate indices, such as the El Niño Southern Oscillation (ENSO; Barber and Chavez, 1983; Barber and Chávez, 1986; Lynn and Bograd, 2002; Chavez et al., 2002; Escribano et al., 2004; Frischknecht et al., 2015), Pacific Decadal Oscillation (PDO; Mantua et al., 1997; Chhak and Di Lorenzo, 2007; Chenillat et al., 2012), North Pacific Gyre Oscillation (NPGO; Di Lorenzo et al., 2008, 2009; Chenillat et al., 2012), and North Atlantic Oscillation (NAO; Borges et al., 2003; Cropper et al., 2014). A similar analysis for EBUS CO<sub>2</sub> fluxes is necessary for the community. Identifying robust relationships between modes of climate variability and CO<sub>2</sub> fluxes would (1) advance our understanding of the complex carbonate system in EBUS and (2) provide some level of predictability in CO<sub>2</sub> flux anomalies. EBUS are naturally sensitive to ocean acidification and have internal variability that rivals the magnitude of the seasonal cycle. Thus, forecasting anomalous CO<sub>2</sub> fluxes could aid in detecting the early onset of seasonal ocean acidification (Landschützer et al., 2018; Kwiatkowski and Orr, 2018; Hauck, 2018).

Previous studies have utilized observations (*e.g.*, Boyd et al., 1987; Friederich et al., 2002; Chavez et al., 2002; Santana-Casiano et al., 2007; Di Lorenzo et al., 2008, 2009) and output from high-resolution hindcast simulations (Jacox et al., 2015; Frischknecht et al., 2015, 2017; Turi et al., 2017; Mogollón and Calil, 2017) to explore the relationship between climate variability and EBUS biogeochemistry. However, single realizations, whether modeled or observed, provide a limited sample size of internal variability. At most, hindcast model runs and observational time series capture 10 El Niño and 10 La Niña events, and only 1 phase of the PDO. This is problematic, because locally generated variability in EBUS can cause a diversity of responses to climate modes such as ENSO (Frischknecht et al., 2017; Fiedler and Mantua, 2017). Thus, one requires many more case studies to robustly assess the response of EBUS biogeochem-

istry to internal variability. A solution to this problem is to use a single-model ensemble that is derived by introducing round-off level perturbations to the initial state of the climate system. This gives rise to a set of realizations with unique representations of internal climate variability and gives one access to many hundred ENSO events, rather than just a handful. By performing historical experiments with increasing atmospheric  $CO_2$  rather than a long control run, we can account for variability in anthropogenic  $CO_2$  in the ocean as well as potential modifications to the frequency and amplitude of internal variability with climate change (e.g., Timmermann et al., 1999).

In this study, we utilize output from the single-model Community Earth System Model "Large Ensemble" (CESM-LENS; Kay et al., 2015; Lovenduski et al., 2016) to identify major modes of climate variability that are associated with anomalous CO<sub>2</sub> fluxes in the major EBUS. We expand on this by investigating the physical and biological drivers that underpin these anomalies. The single-model ensemble is necessary for such an analysis, since the forced signal can be removed to generate residual simulations that solely represent CO<sub>2</sub> flux responses to internal climate variability (Thompson et al., 2015). Since the simulations are forced with historical CO<sub>2</sub> emissions, each member accounts for the combined response of natural and anthropogenic CO<sub>2</sub> to climate variability. Furthermore, the availability of 34 simulations allows us to find statistically robust relationships between anomalous fluxes and internal variability. This experimental setup addresses the data limitation issues of an observational study as well as the single realization problem of a model hindcast or a multi-model ensemble.

## 2 Methods and Model Evaluation

# 2.1 Model Configuration and Upwelling Regions

We utilize monthly output from 34 members of the CESM-LENS, which is derived from a fully coupled Atmosphere-Ocean General Circulation Model (AOGCM) with ocean biogeochemistry (Kay et al., 2015; Lovenduski et al., 2016). Round-off level perturbations are made to the atmospheric temperature in 1920, leading to an ensemble of simulations that diverge solely due to

the influence of internally generated variability. This provides us with a set of 34 independent representations of climate variability, with which we can robustly assess the controls on air-sea CO<sub>2</sub> flux variability in EBUS. The ensemble is forced with historical radiative forcing from 1920–2005 and RCP8.5 radiative forcing from 2006–2100. The ocean model has nominal 1° horizontal resolution with vertical resolution of 10 m through the upper 250 m, thus resolving the Ekman layer. Due to the coarse horizontal resolution, neither curl-driven nor coastal upwelling is directly resolved, but both are represented in the model. A more detailed discussion of coastal upwelling in the CESM-LENS for the CalCS in particular can be found in Brady et al. (2017).

Upwelling regions were confined to approximately the  $10^{\circ}$  latitude of most active upwelling as defined by Chavez and Messié (2009), although the CanCS domain was shifted north by  $9^{\circ}$  to capture the more intense upwelling off the Western Sahara in CESM1 (Table 1). They were then limited to the first 800km in the offshore direction. The black outlines in Figure 1e-h display these regions. The ensemble mean – which represents both the seasonality and anthropogenic trend for  $CO_2$  flux – was removed from each simulation for each upwelling system to create internally generated residuals. All output was analyzed at monthly resolution.

## 2.2 Statistical Analysis and Model Equations

Air-sea CO<sub>2</sub> fluxes in CESM are computed following the parameterization of Wanninkhof (2014):

$$F = k \cdot K_0 \cdot (pCO_2^o - pCO_2^a), \tag{1}$$

where k represents the gas transfer velocity (dependent on the wind speed squared),  $K_0$  the solubility of  $CO_2$  in seawater, and  $pCO_2^o$  and  $pCO_2^o$  the partial pressures of  $CO_2$  in the surface ocean and atmosphere, respectively.

We use a linear Taylor expansion to quantify the relative contribution of each variable to the overall CO<sub>2</sub> flux anomaly in response to internally generated variability following Lovenduski

et al. (2007) and Turi et al. (2014),

$$\Delta F = \frac{\partial F}{\partial U} \Delta U + \frac{\partial F}{\partial p C O_2^{oc}} \Delta p C O_2^{oc}, \tag{2}$$

where  $\frac{\partial F}{\partial U}$  and  $\frac{\partial F}{\partial pCO_2^{oc}}$  are determined from the model equations and mean values in each EBUS, and the influence of sea ice was assumed to be negligible.  $\Delta$ 's represent the linear regression of the given variable's residuals onto a climate index. The contributions from  $\Delta pCO_2^{oc}$  is further decomposed into DIC, Alk, SST, and salinity terms:

$$\Delta pCO_2^{oc} = \frac{\partial pCO_2^{oc}}{\partial DIC} \Delta DIC + \frac{\partial pCO_2^{oc}}{\partial Alk} \Delta Alk + \frac{\partial pCO_2^{oc}}{\partial T} \Delta T + \frac{\partial pCO_2^{oc}}{\partial S} \Delta S.$$
 (3)

Because DIC and Alk can be diluted by freshwater fluxes, we introduce salinity-normalized DIC (sDIC) and Alk (sAlk),

$$\Delta F = \frac{\partial F}{\partial U} \Delta U + \frac{S}{S_0} \frac{\partial F}{\partial DIC} \Delta sDIC + \frac{S}{S_0} \frac{\partial F}{\partial Alk} \Delta sAlk + \frac{\partial F}{\partial fw} \Delta fw + \frac{\partial F}{\partial T} \Delta T + \frac{\partial F}{\partial S} \Delta S. \tag{4}$$

To compensate for autocorrelation that is characteristic of climate indices and is also introduced from smoothing, we replace the *t*-statistic sample size N with an effective sample size,  $N_{eff}$ :

$$N_{eff} = N\left(\frac{1 - r_1 r_2}{1 + r_1 r_2}\right) \tag{5}$$

where  $r_1$  and  $r_2$  are the lag-1 autocorrelation coefficients of the two time series being correlated (Bretherton et al., 1999; Lovenduski and Gruber, 2005).  $N_{eff}$  represents the number of statistically independent measurements.

## 2.3 Model Evaluation

CESM-LENS air-sea CO<sub>2</sub> fluxes were compared to the observationally-based SOM-FFN (Self-Organizing Map-Feed Forward Network) product from Landschützer et al. (2017) along the four

major EBUS outlined by Chavez and Messié (2009). The SOM-FFN was generated by a two step process. First, the global oceans were grouped into 16 biogeochemical provinces based on common relationships between SSTs, sea surface salinity, mixed layer depth, and  $pCO_2$  climatology from Takahashi et al. (2009). Secondly, nonlinear relationships were determined between an expanded set of predictor variables and the Surface Ocean Carbon Atlas version 4 (Bakker et al., 2016) database of surface ocean  $CO_2$  measurements to interpolate  $pCO_2$  to monthly resolution spanning 1982-2015 at  $1^ox1^o$  global resolution. Extensive details on and validation of the procedure can be found in Landschützer et al. (2013) and Landschützer et al. (2016).

Figure 1 compares the historical climatology (1982–2015) between the SOM-FFN (a-d) and the CESM-LENS (e-h). The Pacific systems are particularly well-modeled. The CESM-LENS simulates the meridional gradient of poleward uptake and equatorward outgassing of CO<sub>2</sub> in the CalCS (Figure 1e). For this system and all other EBUS, we do not expect the model to resolve nearshore outgassing, as coastal upwelling here occurs at the sub-grid scale. In the HumCS, the model depicts the strong outgassing that is characteristic of a tropical upwelling system (Figure 1f). The CO<sub>2</sub> flux climatology in the Atlantic systems is more biased in the CESM-LENS. While the SOM-FFN portrays a meridional gradient of relatively weak CO<sub>2</sub> fluxes in the CanCS, the CESM-LENS simulates strong outgassing along the Western Sahara (Figure 1c and g). An important caveat is that the data density of pCO<sub>2</sub> in EBUS informing the SOM-FFN is on the order of the Southern Ocean, a notably undersampled region (Bakker et al., 2016; Laruelle et al., 2017). In turn, the EBUS CO<sub>2</sub> fluxes are being informed by remote biogeochemical provinces more often than other regions of the ocean. The BenCS has the most biased CO<sub>2</sub> flux climatology of the major EBUS in CESM-LENS. Although it simulates the meridional gradient portrayed in the SOM-FFN, the ougassing cell is nearly 10° too far south and is significantly stronger than in the SOM-FFN (Figure 1d and h).

The BenCS has the largest physical biases in CESM-LENS than all other EBUS. Its SST bias is in excess of  $7^{\circ}$ C with the nominal  $2^{\circ}$  atmospheric resolution. Further, it only improves to a  $5^{\circ}$ C bias at  $0.5^{\circ}$  atmospheric resolution (Gent et al., 2010). This bias is likely driven by the fact that the Angola-Benguela Front is simulated too far south, in addition to deficiencies in up-

welling and meridional transport that are driven by unrealistic alongshore wind stress structure (Small et al., 2015). Because of these deficiencies that are specific to the BenCS, we will only discuss its internal variability in CO<sub>2</sub> fluxes in Section 3.1, but will not perform a full analysis on its connections to larger-scale climate variability.

#### 3 Results

## 3.1 Internal Variability in Upwelling Systems

We emphasize the magnitude of internal variability in EBUS CO<sub>2</sub> fluxes in Figure 2 by showing the ensemble mean standard deviation of air-sea CO<sub>2</sub> flux residuals (ensemble mean subtracted) at each location across the global ocean. Save for the Southern Ocean and subpolar Arctic, the EBUS emerge as significant regions influenced by internal variability on a global scale. The HumCS, CanCS, and BenCS in particular have some of the highest internally driven CO<sub>2</sub> fluxes globally. The CalCS has comparatively low internal variability in CO<sub>2</sub> fluxes. Coastally, the EBUS are distinct from the major western boundary currents, which appear to be influenced very little by internal variability (Figure 2).

Figure 3 (a–d) displays time series of historical  $CO_2$  flux from 1920–2015 and the mean seasonal cycle (e–h) for each of the four EBUS. The black line depicts the seasonal cycle, the red line the anthropogenically forced trend, and the gray shading the component due to internal variability. Values for each of these components and the linear intercept are reported in Table 1. The largest absolute internal variability is found in the BenCS and HumCS with values of 0.98 mol m<sup>-2</sup> yr<sup>-1</sup> and 1.20 mol m<sup>-2</sup> yr<sup>-1</sup>, respectively (Table 1). The BenCS is uniquely exposed to variability from the Southern Ocean and Agulhas Current (Reason et al., 2006). The HumCS likely has intense variability due to its direct connection to the tropical Pacific Ocean and thus rapid communication with ENSO (e.g., Colas et al., 2008; Montes et al., 2011).

All four systems have statistically significant trends toward a greater  $CO_2$  sink due to the invasion of anthropogenic carbon (Figure 3; Table 1). This forces the HumCS, CanCS, and BenCS to act as intermittent sinks by 2015 in some realizations due to the combination of the

long-term trend and internal variability. The HumCS and BenCS have the largest uptake of anthropogenic CO<sub>2</sub> over the historical period, which is on the order of the magnitude of their seasonal cycles over the course of 96 years. The CanCS is a unique case, where the anthropogenic trend is more than double the magnitude of its seasonal cycle (Table 1). The seasonal cycle of all EBUS excluding the CanCS is approximately sinusoidal with an outgassing peak in the late boreal summer (Figure 3e–h). The CanCS has a relatively weak bi-modal outgassing peak in late winter and early summer, but is otherwise relatively flat (Figure 3g). It has by far the smallest seasonal cycle of the four systems (Table 1).

The magnitude of internal variability is greater than that of the seasonal cycle for the majority of systems. The non-seasonal component of variability is 59% for the HumCS, 73% for the CanCS, and 56% for the BenCS (Table 1). Only the CalCS has a stronger seasonal cycle of CO<sub>2</sub> flux than internal variability, but the non-seasonal component still accounts for 33% of the variability in this system (Table 1). Perhaps for the CalCS, more significant internal variability would be captured at a higher resolution that resolves coastal upwelling, such as in Turi et al. (2014, Figure 8c). Lastly, internal variability in CO<sub>2</sub> fluxes tends to be phase-locked with its seasonal cycle for all of the EBUS, as the peak magnitudes of internal variability track the ridges and troughs of the seasonal component (Figure 3, a–d). This is most clear in the CalCS, HumCS, and BenCS. The CanCS has more complex behavior in the residuals due to the bimodal peaks of its seasonal cycle (Figure 3g).

## 3.2 California Current

Our primary goal for each EBUS was to identify the dominating mode of climate variability associated with its CO<sub>2</sub> flux residuals. We followed a top-down approach, correlating area-weighted residuals from the black boxes in Figure 1e-h for each simulation with every grid cell globally for a set of predictor variables' residuals: SST, sea level pressure (SLP), 10m wind speed, and wind stress curl. We then assessed the ensemble mean of the regressions to determine the mode of climate variability associated with the given global spatial pattern. Figure 4 displays one ensemble mean regression case for the CalCS (a), HumCS (b), and CanCS (c).

The global regression between CalCS  $CO_2$  flux residuals and SSTa yields a map clearly indicative of Pacific Decadal Variability, due to the zonal dipole of correlations in the North Pacific (Figure 4a; Mantua and Hare, 2002; Di Lorenzo et al., 2008). Although similar in structure to the PDO, this map most closely resembles the NPGO (Di Lorenzo et al., 2008). In fact, correlations between the NPGO with annual smoothing and CalCS  $CO_2$  flux yields an r-value of -0.49  $\pm$  0.04. In comparison, linear correlations with the PDO result in an r-value of 0.24  $\pm$  0.05. Thus, we highlight the NPGO as the major mode of climate variability associated with anomalous  $CO_2$  flux in the CalCS. We computed the NPGO index in CESM-LENS following Di Lorenzo and Mantua (2016).

Figure 5a depicts the results of a linear Taylor expansion for CalCS  $CO_2$  flux residuals regressed onto the  $1\sigma$  positive phase of the NPGO (Eq. 4). Since the area-weighted analysis might be sensitive to the box chosen to represent the CalCS, we also include correlations between individual grid cells and the NPGO in Figure 5b. We find that the CalCS responds uniformly to the NPGO with anomalous uptake of  $CO_2$ , intensifying the mean state of the system as an uptake site (Figure 5). The direct regression of  $\Delta F$  onto the NPGO results in an anomalous uptake of 0.10 mol m<sup>-2</sup> yr<sup>-1</sup> (Table 2), which is roughly 24% of the long-term historical mean of -0.42 mol m<sup>-2</sup> yr<sup>-1</sup>. The primary contributions to this uptake anomaly come from SST and sDIC, which act in opposition to one another. This is coherent with our definition of the NPGO as the oceanic expression of the atmospheric North Pacific Oscillation (NPO; Di Lorenzo et al., 2008). A positive NPO (and thus NPGO) increases upwelling-favorable conditions in the CalCS through intensification of the North Pacific High (Di Lorenzo et al., 2008). This leads to enhanced upwelling of cold subsurface waters, which increases the  $CO_2$  solubility of the system, contributing toward the uptake anomaly.

Nearly in balance with this term is the influence of sDIC. The enhanced upwelling also delivers remineralized carbon from depth, which increases surface sDIC, contributing to an opposing outgassing anomaly. In fact it is the minor uptake contributions from the remaining terms – wind, salinity, sAlk, and freshwater flux – that pushes the system in favor of anomalous uptake. The CalCS has the largest relative ensemble spread in sDIC and sAlk (Figure 5; Table 2). This is potentially due to inter-simulation variability in the upwelling response to NPGO or in

the source waters feeding the upwelling. Although the linear Taylor expansions approximates a  $CO_2$  flux anomaly nearly half that of the direct regression of  $\Delta F$  onto the NPGO, it is still of the same sign. This discrepancy is due to the influence of higher-order and cross-derivative terms that we did not account for in our approximation.

We also performed this analysis for the CalCS response to a  $1\sigma$  positive (warm) phase of the PDO (Figure 6a and b). Every simulation displayed a dipole response to the PDO, with anomalous uptake in the nearshore region south of Cape Mendocino, and anomalous outgassing elsewhere in the domain (Figure 6c). This was the only case in which we found a non-uniform response across all simulations to any mode of climate variability investigated. Both the nearshore and offshore regions have modest correlations with the PDO, with r-values of -0.16  $\pm$  0.03 and 0.28  $\pm$  0.05, respectively. The positive phase of the PDO results in anomalously warm SSTs along the CalCS and causes shallower upwelling cells with higher retention of nutrient- and carbon-depleted surface waters (Chhak and Di Lorenzo, 2007). This aligns with the inverted contributions of SST and sDIC in Figure 6a and b relative to the contributions of these terms in response to the NPGO (Figure 5a).

The warming of CalCS SSTs during a positive phase of PDO causes a reduction of  $CO_2$  solubility and thus a tendency toward outgassing (Figure 6). The shallow upwelling cells with less remineralized carbon contribute toward anomalous uptake of  $CO_2$  throughout the system. Note that the nearshore decomposition in Figure 6a has a y-axis range four times smaller than that of the offshore decomposition. This slight uptake anomaly is the result of a delicate balance of minor terms, where the sDIC reduction slightly outweighs the warming effect. On the other hand, the offshore region has contributions from SST and sDIC that are as much as triple the magnitude as that for the NPGO (Table 2). Despite the sDIC reduction being larger than the SST term, the reduced sAlk is substantial enough to cause a slight outgassing anomaly offshore (Figure 6b).

Interestingly, the direct response of winds to the NPGO and PDO plays a negligible role in influencing anomalous  $CO_2$  flux in the CalCS (Table 2). Although  $\Delta U$  in response to the NPGO and PDO is on the order of the HumCS and CanCS,  $\frac{\partial F}{\partial U}$  is 3–10 times smaller than the other systems.  $\frac{\partial F}{\partial U}$  is based on the climatological mean U,  $\Delta pCO_2$ , and Schmidt number. The CalCS

has the smallest mean  $\Delta pCO_2$  of the EBUS – just  $0.2\mu$ atm. This causes  $CO_2$  flux in the system to be relatively insensitive to fluctuations in the wind.

#### 3.3 Humboldt Current

Figure 4b shows the ensemble mean global correlation between the HumCS and SSTa. This clearly displays ENSO as the major influence on  $CO_2$  flux anomalies in the HumCS, with regions of high correlation focused around the equatorial Pacific. Correlations between HumCS  $CO_2$  flux anomalies and the Nino3 index resulted in an r-value of -0.40  $\pm$  0.04. Similar results were found for the Nino3.4 index (-0.38  $\pm$  0.04) and the Nino4 index (-0.36  $\pm$  0.05). We chose the Nino3 index as our primary predictor of HumCS  $CO_2$  flux anomalies, since it is more eastern-focused and thus captures the stronger spatial correlations closest to the HumCS (Figure 4b).

We present the results of a linear Taylor expansion for HumCS  $CO_2$  flux residuals regressed onto a  $1^o$  El Niño in Figure 7 (Eq. 4). We find that the HumCS responds with a near-uniform  $CO_2$  uptake anomaly, resulting in a weakening of the climatological outgassing site (Figure 7b). Although there is a small region in the northern HumCS that responds with an outgassing anomaly, it is nowhere near as coherent across the ensemble as was the spatial dipole response of the CalCS to the PDO (Figure 6c). The direct regression of  $\Delta F$  onto the Nino3 index results in an anomalous uptake of 0.49 mol m $^{-2}$  yr $^{-1}$ , which is approximately 18% of the long-term historical mean of 2.8 mol m $^{-2}$  yr $^{-1}$ . As in the case of the CalCS, the two major terms contributing to the uptake anomaly are sDIC and SST, which are in opposition to one another. We would anticipate this to be the case, as an El Niño event induces warming along the HumCS as well as reduces the efficacy of upwelling due to the presence of an anomalously deep thermocline (Strub et al., 1998).

In CESM-LENS, the HumCS experiences a warming of  $0.7^{\circ}$ C for a  $1^{\circ}$  El Niño, which results in an outgassing pressure of 0.3 mol m<sup>-2</sup> yr<sup>-1</sup> (Table 2). However, sDIC in the system is reduced by 13.2 mmol m<sup>-3</sup> for the same event, which translates to a large uptake contribution of 0.8 mol m<sup>-2</sup> yr<sup>-1</sup> (Table 2). This is an enormous change in sDIC, which is partially driven by the high subsurface DIC bias in the east equatorial Pacific in CESM (see Lovenduski et al.,

2015, their Figure 2). The large sDIC reduction is due to weakened upwelling and a deepening of the thermocline by advected midequatorial waters.

Lastly, there is a minor outgassing anomaly of 0.06 mol m<sup>-2</sup> yr<sup>-1</sup> in response to a slight intensification of winds during El Niño (Table 2). While upwelling-favorable winds tend to decrease along Chile during an El Niño, they generally persist or strengthen along Peru (Wyrtki, 1975; Enfield, 1981; Huyer et al., 1987). Despite the significant contributions of wind speed, SST, and sAlk toward outgassing, the large reduction in sDIC drives an uptake anomaly that weakens the HumCS outgassing during an El Niño event.

# 3.4 Canary Current

The global correlation between CanCS  $CO_2$  flux anomalies and SLPa are displayed in Figure 4c. Here a region of high positive correlation emerges just northwest of Africa. This encircles the climatological position of the Azores High, the atmospheric subtropical gyre which forces the CanCS. The climate index that most directly captures variability in the Azores High is the NAO, and will thus be considered the main mode of climate variability that modulates anomalous  $CO_2$  flux in the CanCS. We find modest correlations of  $0.28 \pm 0.03$  between annually smoothed CanCS  $CO_2$  flux anomalies and the NAO.

Grid cell correlations between CanCS  $CO_2$  flux anomalies and the NAO are displayed in Figure 8b. The CanCS has a nearly uniform response of increased outgassing during the positive phase of the NAO. The direct regression of  $\Delta F$  onto a  $1\sigma$  NAO results in an outgassing anomaly of 0.2 mol m<sup>-2</sup> yr<sup>-1</sup> (Table 2), which is 21% of the historical mean of 0.95 mol m<sup>-2</sup> yr<sup>-1</sup>. Also note that the linear Taylor approximation aligns exactly with the direct regression. As with the other EBUS, the major contributors toward this anomaly are sDIC and SST (Figure 8a). Their directions align with that of the CalCS response to the NPGO (Figure 5a). This is expected, as the NAO is the atmospheric analogue to the NPGO in the Atlantic – it describes modifications to the intensity of atmospheric gyre circulation between the Azores High and Icelandic Low.

During the positive phase of the NAO, a stronger Azores High leads to intensified alongshore winds and thus more vigorous upwelling. This brings up additional deep cold water which in turn increases the CO<sub>2</sub> solubility of the system, tending toward an uptake anomaly of 0.15 mol m<sup>-2</sup> yr<sup>-1</sup> (Table 2). On the other hand, the increased sDIC from intensified upwelling is double the magnitude of the SST contribution, leading to an outgassing anomaly of 0.33 mol m<sup>-2</sup> yr<sup>-1</sup>. This high sDIC response is driven both by a high  $\Delta sDIC$  on 3.9 mmol m<sup>-3</sup> per  $1\sigma$  NAO as well as the fact that the CanCS has the highest  $\frac{\partial F}{\partial sDIC}$  of the major EBUS. Increased winds of 0.3 m s<sup>-1</sup> per  $1\sigma$  NAO lead to a significant outgassing pressure of 0.05 mol m<sup>-2</sup> yr This is due both to a high system sensitivity,  $\frac{\partial F}{\partial U}$ , to changes in wind and a high wind anomaly in response to the NAO. Ultimately, intensified winds and an anomalous increase in sDIC due to enhanced upwelling counteracts the solubility effects of colder SSTs. This leads to the highest relative CO<sub>2</sub> flux anomaly of any system, with a 21% increase in outgassing per  $1\sigma$  NAO event.

## 4 Summary and Conclusions

We utilize a 34-member single-model ensemble to investigate the relationship between internal climate variability and anomalous  $CO_2$  fluxes in the major EBUS over the historical period (1920–2015). We find that the magnitude of internal variability in EBUS  $CO_2$  fluxes is large and is only rivaled globally by the Southern Ocean and subpolar Arctic. For all EBUS but the CalCS, internal variability in  $CO_2$  fluxes is larger than the seasonal cycle. The highest absolute magnitude of internal variability is in the BenCS and HumCS, with values of 0.98 mol m<sup>-2</sup> yr<sup>-1</sup> and 1.20 mol m<sup>-2</sup> yr<sup>-1</sup>, respectively (Table 1). We identify the major mode of climate variability associated with  $CO_2$  flux residuals for three of the four systems, and then perform a linear Taylor expansion to explore how wind speed, SST, salinity, sDIC, sAlk, and freshwater fluxes individually contribute to the total anomaly. The BenCS was not analyzed in this way, due to significant model biases in alongshore winds, upwelling, and SSTs in CESM (Small et al., 2015).

We find that oscillations in the subtropical anticyclonic gyres exert the most influence on CO<sub>2</sub> fluxes in the CalCS and CanCS. CanCS CO<sub>2</sub> flux anomalies are associated mainly with the NAO, which in its positive phase reflects an intensification of the Azores High. The CalCS is modulated mainly by the NPGO, the oceanic expression of the atmospheric NPO, which in its positive phase describes an enhanced North Pacific High (Di Lorenzo et al., 2008). Anoma-

lously cold waters from upwelling increase  $CO_2$  solubility and contribute toward  $CO_2$  uptake. The increased sDIC delivery from upwelling acts in opposition to the SST effect, contributing toward outgassing. The sDIC and SST contributions nearly exactly balance each other in the CalCS, so it is the minor contributions from winds, salinity, sAlk, and freshwater fluxes that tip the system toward a slight uptake anomaly. In contrast, the sDIC effect in the CanCS nearly doubles the SST effect. The outgassing pressure from increased sDIC is further reinforced by the relatively large increase in wind speed in response to the NAO and the system's high sensitivity to changes in winds, which is approximately 3 times greater than the CalCS. Interestingly, the CalCS experiences an uptake anomaly in response to enhanced upwelling-favorable gyre circulation, while the CanCS has increased outgassing. However, the mean state of both systems is intensified during a positive NPGO and NAO.  $\frac{\partial F}{\partial U}$  is negative in the CalCS, but positive in the CanCS, i.e., increased winds drive uptake (outgassing) anomalies in the CalCS (CanCS). The sensitivity term is directly dependent on  $\Delta pCO_2$ , which favors uptake (outgassing) in the CalCS (CanCS).

We also investigated the CalCS response to the PDO in two sub-regions of the system that captured its unique dipole response of anomalous uptake nearshore and outgassing offshore (Figure 6). The nearshore region experiences an uptake anomaly which is driven by a delicate balance between reduced sDIC and warmer SSTs. The offshore region experiences the same sign changes in sDIC and SST, but with contributions an order of magnitude larger in size. It is actually the reduction in sAlk offshore that contributes toward outgassing and causes the slight outgassing anomaly overall. The dipole is thus not due to an inverted response of sDIC and SST to the PDO, but potentially to variability in the source waters feeding the two regions of the CalCS.

We show that CO<sub>2</sub> flux anomalies in the HumCS are mostly driven by ENSO, due to its direct connection with the equatorial Pacific. This is largely due to our definition of the HumCS, which in fact only captures the northern Humboldt Current System. We might anticipate that the Chilean portion of the system would be more closely related to the CalCS and CanCS, and thus responsive to anticyclonic gyre oscillations due to its closer proximity to the South Pacific High. We find that the HumCS has weakened outgassing during El Niño due to a large

anomalous reduction in sDIC. The sDIC response is large enough to counteract the outgassing pressure from warmer SSTs, increased winds, and reduced sAlk.

In summary, we find that variations in sDIC and SST exert the most influence on anomalous CO<sub>2</sub> fluxes in the CalCS, CanCS, and HumCS. Further, these terms always act in opposition to one another. Secondary to these terms are wind speed and sAlk. Although their contributions do not rival those of SST and sDIC in magnitude, they act to further reinforce anomalies or to tip the balance toward outgassing or uptake when sDIC and SST are of equal magnitude. In all systems, salinity and freshwater fluxes have negligible contributions toward the total CO<sub>2</sub> flux anomaly. Future work should investigate the sDIC term more closely, which is modified both by changes to physical circulation (e.g., upwelling, advection, variability in the mixed layer depth) and biology (photosynthesis and calcium carbonate formation and dissolution). However, biology should generally act in opposition to circulation effects in the EBUS: intensified upwelling increases sDIC, but also delivers nutrients which enhance productivity that reduces sDIC through photosynthesis. However, partitioning the contributions toward sDIC anomalies is still a useful exercise to compare the relative influence of biology between different EBUS.

Our study serves as a starting point toward better understanding how internal climate variability modulates CO<sub>2</sub> fluxes in the major EBUS. It is limited by our use of a single model ensemble and by the coarseness of our climate model. The community would benefit from future studies involving multiple single-model ensembles, which would reduce uncertainty due to structural biases, such as in the dynamics of the BenCS and the elevated sub-surface DIC concentration in the east equatorial Pacific. Due to model resolution, we do not directly resolve the coastal upwelling process which induces vigorous outgassing within the first O(10km) of the coastline. This problem could be mitigated by nesting high-resolution EBUS ROMS simulations within a coarser global ensemble or by using regional mesh refinement techniques. This would allow the remote propagation of climate variability into the EBUS, while avoiding the high computational cost of running multiple high-resolution global simulations. In particular, the BenCS requires significant attention. We find pronounced internal variability in CO<sub>2</sub> fluxes in the BenCS in CESM-LENS that warrants investigation in a high-resolution model specific to the BenCS. We anticipate that these results and further investigation of the relationship be-

tween internal climate variability and anomalous  $CO_2$  fluxes in EBUS will be useful for the rapidly developing subseasonal to decadal prediction community. Skillful prediction of climate variability, such as ENSO, the NPGO, and NAO, could be linked directly to anomalous fluxes of  $CO_2$  in the major EBUS. As these systems are naturally sensitive to the undersaturation of calcium carbonate, these predictions could aid in detecting and managing the onset of seasonal ocean acidification.

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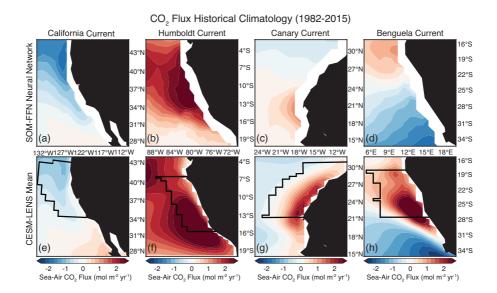
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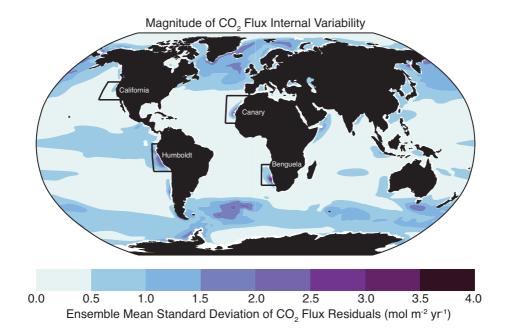
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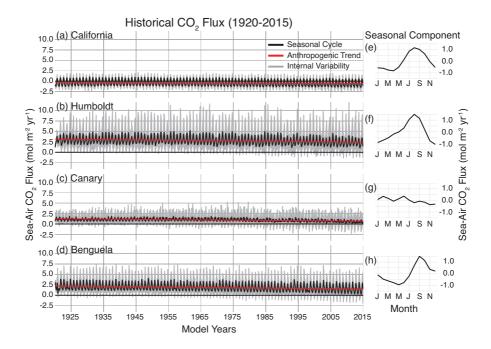
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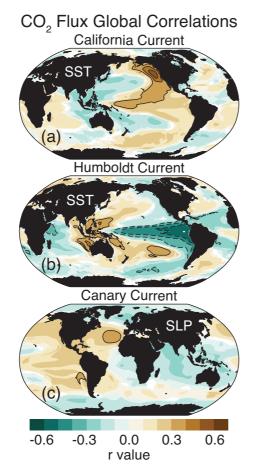
**Fig. 1.** Comparison of  $CO_2$  flux climatology from 1982–2015 between the SOM-FFN (a–d) and the CESM-LENS (e–h). Red denotes outgassing of  $CO_2$  from the ocean to the atmosphere, while blue represents uptake of  $CO_2$  by the ocean. Black lines in e–h follow the model grid and show the region used in each EBUS for statistical analysis, which is based on the  $10^{\circ}$  latitude of most active upwelling from Chavez and Messié (2009) and confined to 800km offshore.



**Fig. 2.** Magnitude of internal variability in  $CO_2$  flux from 1920–2015 in the CESM-LENS. Residuals were generated by removing the ensemble mean – which represents the seasonality and forced signal – from each realization. Internal variability was then quantified by taking the ensemble mean standard deviation of the residuals from 1920–2015.



**Fig. 3.** Time series of historical  $CO_2$  flux (1920–2015) in the CESM-LENS for each of the four studied upwelling systems (a–d). The ensemble mean yields both the seasonal cycle (black) and the anthropogenic trend (red). Gray shading shows the bounds of the maximum and minimum realizations due to internal variability. Table 1 displays the intercept, seasonality, internal variability, and anthropogenic trend for each system. Plots e–h show the mean seasonal cycle for 1920–2015 for each system.



**Fig. 4.** Correlations between area—weighted  $CO_2$  flux residuals in the statistical study regions outlined in black in Figure 1 (e–g) and SSTa (a–b; California, Humboldt) and SLPa (c; Canary) grid cells globally. Brown colors indicate that positive values of SSTa/SLPa correlate with outgassing, and blue with uptake. Contour lines begin at  $\pm$  0.3 and progress in intervals of 0.1. Correlations were performed for each realization individually and the ensemble mean of those global correlations are presented here.

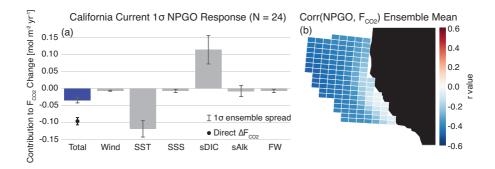
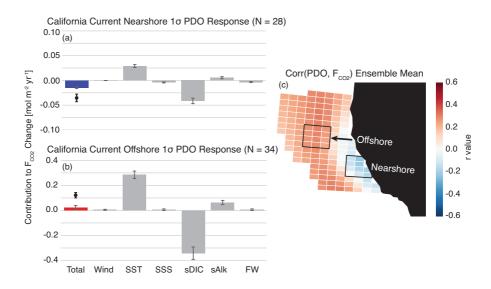


Fig. 5. Linear Taylor expansion from Eq. (4) for CalCS  $CO_2$  flux residuals regressed onto the NPGO (a). Gray bars represent the ensemble mean contributions of each variable to the  $CO_2$  flux anomaly. Error bars represent the one standard deviation spread of the full ensemble. The individual bars sum to the "total" bar to approximate the direct regression of  $\Delta F_{CO_2}$  onto the NPGO, which is denoted as the black dot with its associated ensemble spread. The ensemble mean grid cell correlations between  $CO_2$  flux residuals and the NPGO in the CalCS study region are displayed in (b). Positive correlations are associated with outgassing, negative with uptake. Values and ensemble spread for each bar are presented in Table 2.



**Fig. 6.** As in Figure 5, but in response to the PDO for a nearshore region (a) and offshore region (b). Note that the offshore decomposition (b) has a y-axis range four times that of the nearshore decomposition (a).

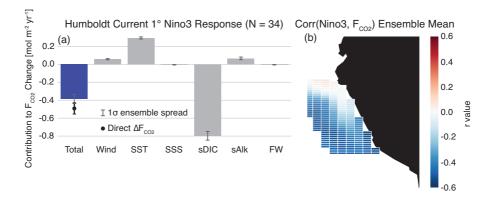


Fig. 7. As in Figure 5, but for the Humboldt Current response to the Nino3 index.

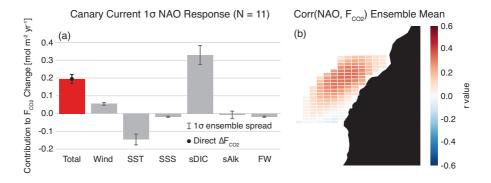


Fig. 8. As in Figure 5, but for the Canary Current response to the NAO.

**Table 1.** Statistics for  $CO_2$  fluxes in the CalCS, HumCS, CanCS, and BenCS from 1920–2015. The seasonal component is computed as the standard deviation of the ensemble mean after removing a fourth-order polynomial fit to remove the anthropogenic trend. The internal component is computed as the ensemble mean standard deviation of the residuals. The trend is computed as the first-order ordinary least squares regression coefficient. The intercept is derived from this linear regression.

Upwelling System	Intercept <sup>1</sup>	Seasonal <sup>1</sup>	Internal <sup>1</sup>	Trend <sup>23</sup>	Non-Seasonal Component (%)
California (34° N–44° N)	-0.27	0.71	0.33	-0.31	31
Humboldt (16° S–6° S)	3.16	0.83	1.20	-0.71	59
Canary (21° N–31° N)	1.23	0.23	0.62	-0.56	73
Benguela (28° S–18° S)	2.25	0.77	0.98	-0.76	56

 $<sup>^{1}</sup>$  mol m $^{-2}$  yr $^{-1}$ 

<sup>2</sup>mol m<sup>-2</sup> yr<sup>-1</sup> 96yr<sup>-1</sup>

 $<sup>^3</sup>$  All trends are significant to  $\alpha = 0.05$  for a one-sided Mann-Kendall Test

**Table 2.** Estimated contributions to  $CO_2$  flux anomalies,  $\Delta F$  using Eq. (4).

Quantity	CalCS – NPGO	CalCS – PDOo	CalCS – PDOn	HumCS – Nino3	CanCS – NAO				
Individual Terms									
$\frac{\partial F}{\partial U}\Delta U$	$-0.01 \pm 0.0$	$0.01 \pm 0.0$	$0.0 \pm 0.0$	$0.06\pm0.01$	$0.05\pm0.01$				
$\frac{\partial F}{\partial T}\Delta T$	$-0.12 \pm 0.02$	$0.28 \pm 0.03$	$0.03 \pm 0.0$	$0.29 \pm 0.01$	$-0.15 \pm 0.03$				
$\frac{\partial F}{\partial S}\Delta S$	$-0.01 \pm 0.0$	$0.01\pm0.01$	$0.0 \pm 0.0$	$-0.0 \pm 0.0$	$-0.02 \pm 0.0$				
$\frac{S}{S_0} \frac{\partial F}{\partial DIC} \Delta sDIC$	$0.11 \pm 0.04$	$-0.34 \pm 0.05$	$-0.04 \pm 0.01$	$-0.8 \pm 0.05$	$0.33 \pm 0.05$				
$\frac{S}{S_0} \frac{\partial F}{\partial Alk} \Delta s Alk$	$-0.01 \pm 0.02$	$0.06\pm0.02$	$0.01 \pm 0.0$	$0.07\pm0.01$	$-0.01 \pm 0.02$				
$\begin{array}{l} \frac{\partial F}{\partial U} \Delta U \\ \frac{\partial F}{\partial T} \Delta T \\ \frac{\partial F}{\partial S} \Delta S \\ \frac{S}{S_0} \frac{\partial F}{\partial DIC} \Delta s DIC \\ \frac{S}{S_0} \frac{\partial F}{\partial Alk} \Delta s Alk \\ \frac{\partial F}{\partial fw} \Delta fw \end{array}$	$-0.01\pm0.0$	$0.01\pm0.01$	$0.0 \pm 0.0$	$0.0 \pm 0.0$	$-0.02 \pm 0.0$				
Sum of Terms Versus Modeled									
$\Sigma$	$-0.03 \pm 0.01$	$0.02\pm0.02$	$-0.01 \pm 0.0$	$-0.38 \pm 0.05$	$0.21\pm0.03$				
$\Delta F_{mod}$	$-0.10 \pm 0.01$	$0.12\pm0.02$	$-0.04 \pm 0.01$	$-0.49 \pm 0.06$	$0.2 \pm 0.02$				