

# 1 Physical Simulation

## 1.1 Rigid Body Time Integrator

$$L(c^i, \theta^i, c^{i+1}, \theta^{i+1}) = \sum_{bodies} \left[ \frac{\rho V}{2} \left\| \frac{c^{i+1} - c^i}{h} \right\|^2 + \frac{\rho}{2} \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \omega(\theta^i, \theta^{i+1}) \right] - V(c^{i+1}, \theta^{i+1})$$

$$\frac{\partial}{\partial c^i} L(c^i, \theta^i, c^{i+1}, \theta^{i+1}) = \sum_{bodies} \left[ -\rho V \frac{c^{i+1} - c^i}{h^2} \right]^T$$

$$\begin{aligned} \frac{\partial}{\partial \theta^i} L(c^i, \theta^i, c^{i+1}, \theta^{i+1}) &= \sum_{bodies} \frac{\partial}{\partial \theta^i} \left[ \frac{\rho}{2} \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \omega(\theta^i, \theta^{i+1}) \right] \\ &= \sum_{bodies} \left[ \rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I \left( \frac{\partial}{\partial \theta^i} [R_{-\theta^i} \omega(\theta^i, \theta^{i+1})] \right) \right] \\ &= \sum_{bodies} \left[ \rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I \left( \frac{\partial}{\partial \theta^i} R_{-\theta^i} \omega(\theta^i, \theta^{i+1}) + R_{-\theta^i} \frac{\partial}{\partial \theta^i} \omega(\theta^i, \theta^{i+1}) \right) \right] \\ &= \sum_{bodies} \left[ \rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I \left( -R_{-\theta^i} [\omega(\theta^i, \theta^{i+1})]_{\times} T(-\theta_i) + R_{-\theta^i} \frac{-1}{h} T(h\omega(\theta^i, \theta^{i+1}))^{-1} T(-\theta^i) \right) \right] \\ &= \sum_{bodies} \left[ \rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left( -[\omega(\theta^i, \theta^{i+1})]_{\times} T(-\theta_i) + \frac{-1}{h} T(h\omega(\theta^i, \theta^{i+1}))^{-1} T(-\theta^i) \right) \right] \end{aligned}$$

$$\frac{\partial}{\partial c^{i+1}} L(c^i, \theta^i, c^{i+1}, \theta^{i+1}) = \sum_{bodies} \left[ \rho V \frac{c^{i+1} - c^i}{h^2} \right]^T - \frac{\partial}{\partial c^{i+1}} V(c^{i+1}, \theta^{i+1})$$

$$\begin{aligned} \frac{\partial}{\partial \theta^{i+1}} L(c^i, \theta^i, c^{i+1}, \theta^{i+1}) &= \sum_{bodies} \frac{\partial}{\partial \theta^{i+1}} \left[ \frac{\rho}{2} \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \omega(\theta^i, \theta^{i+1}) \right] - \frac{\partial}{\partial \theta^{i+1}} V(\theta^i, \theta^{i+1}) \\ &= \sum_{bodies} \left[ \rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left( \frac{\partial}{\partial \theta^{i+1}} \omega(\theta^i, \theta^{i+1}) \right) \right] - \frac{\partial}{\partial \theta^{i+1}} V(\theta^i, \theta^{i+1}) \\ &= \sum_{bodies} \left[ \rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left( \frac{1}{h} T(-h\omega(\theta^i, \theta^{i+1}))^{-1} T(-\theta^{i+1}) \right) \right] - \frac{\partial}{\partial \theta^{i+1}} V(\theta^i, \theta^{i+1}) \end{aligned}$$

## 1.2 Euler-Lagrange Equations

$$d_2 L(q^{i-1}, q^i) + d_1 L(q^i, q^{i+1}) = 0$$

For the tangential velocity, this is simply particle motion, and follows the force = mass velocity form.

$$\begin{aligned} \sum_{bodies} \left[ \rho V \frac{c^i - c^{i-1}}{h^2} \right]^T - \frac{\partial}{\partial c^i} V(c^i, \theta^i) + \sum_{bodies} \left[ -\rho V \frac{c^{i+1} - c^i}{h^2} \right]^T &= 0 \\ \sum_{bodies} \left[ \rho V \frac{c^{i+1} - c^i}{h^2} \right]^T &= \sum_{bodies} \left[ \rho V \frac{c^i - c^{i-1}}{h^2} \right]^T - \frac{\partial}{\partial c^i} V(c^i, \theta^i) \\ \sum_{bodies} \left[ \rho V \dot{c}_{i+1} \right] &= \sum_{bodies} \left[ \rho V \dot{c}_i \right] - h \frac{\partial}{\partial c^i} V(c^i, \theta^i)^T \\ c_{i+1} &= c^i + h \dot{c}^i \\ \dot{c}_{i+1} &= \dot{c}_i - h \frac{1}{\rho V} dV(c^i, \theta^i)^T \quad \forall i \in bodies \end{aligned}$$

Solving for angular velocity, we can see each rigid body is an independent system, and since the potential  $V(c, \theta)$  in this case only includes gravity,  $d_2 V(c, \theta) = 0$ .

$$\begin{aligned} \left[ \rho \omega(\theta^{i-1}, \theta^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} \left( \frac{1}{h} T(-h\omega(\theta^{i-1}, \theta^i))^{-1} T(-\theta^i) \right) \right] &+ \left[ \rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left( -[\omega(\theta^i, \theta^{i+1})]_{\times} T(-\theta_i) + \frac{-1}{h} T(h\omega(\theta^i, \theta^{i+1}))^{-1} \right) T(-\theta^i) \right] = 0 \\ \left[ \rho \omega(\theta^{i-1}, \theta^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} \frac{1}{h} T(-h\omega(\theta^{i-1}, \theta^i))^{-1} \right] &+ \left[ \rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left( -[\omega(\theta^i, \theta^{i+1})]_{\times} T(-\theta_i) + \frac{-1}{h} T(h\omega(\theta^i, \theta^{i+1}))^{-1} \right) \right] = 0 \end{aligned}$$

The goal is to compute the orientation at the next time step,  $\theta^{i+1}$ . This can be done through solving for  $\omega(\theta^i, \theta^{i+1})$ , since  $R_{h\omega(\theta^i, \theta^{i+1})} = R_{\theta^{i+1}} R_{-\theta^i}$ .

Using properties of rotation matrices, we can see  $R_{-\theta^i}^T R_{h\omega(\theta^i, \theta^{i+1})} = R_{\theta^{i+1}}$ .

Let  $\omega^i = \omega(\theta^{i-1}, \theta^i)$

$$\begin{aligned} \left[ \rho(\omega^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} \frac{1}{h} T(-h\omega^i)^{-1} \right] &+ \left[ \rho(\omega^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left( -[\omega^{i+1}]_{\times} T(-\theta_i) + \frac{-1}{h} T(h\omega^{i+1})^{-1} \right) \right] = 0 \\ \left[ (\omega^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} T(-h\omega^i)^{-1} \right] &+ \left[ (\omega^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left( -h[\omega^{i+1}]_{\times} T(-\theta_i) + T(h\omega^{i+1})^{-1} \right) \right] = 0 \end{aligned}$$

We can use Newton's method to solve for  $\omega^{i+1}$ , where

$$\begin{aligned} f(\omega^{i+1}) &= \left[ (\omega^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} T(-h\omega^i)^{-1} \right] + \left[ (\omega^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left( -h[\omega^{i+1}]_{\times} T(-\theta_i) + T(h\omega^{i+1})^{-1} \right) \right] \\ &= 0 \\ df(\omega^{i+1}) &= R_{-\theta^i}^T M_I R_{-\theta^i} \left( -h[\omega^{i+1}]_{\times} T(-\theta_i) + T(h\omega^{i+1})^{-1} \right) + (\omega^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} \left( -h d[\omega^{i+1}]_{\times} T(-\theta_i) + h dT(h\omega^{i+1})^{-1} \right) \\ &\approx R_{-\theta^i}^T M_I R_{-\theta^i} \left( T(h\omega^{i+1})^{-1} \right) \end{aligned}$$

### 1.3 Integration over a Triangle

Given - arbitrary points of a triangle  $u, v, w$ , and a function  $f$  that you want to integrate over the face of the triangle.

Let's consider the basic function  $f(u, v, w) = 1$  for now, so this integration should give us the area of the triangle -  $\|(v - u) \times (w - u)\|$ .

First step is to parameterize this into a function of two parameters,  $s, t$ , using the basis vectors  $v - u$ , and  $w - u$ .

This can be thought of sliding across the plane that is covered by the triangle.

This leads us to the double integral -

$$\begin{aligned} \int_0^1 \int_0^{1-s} f(s, t) \, dt \, ds &= \int_0^1 t|_0^{1-s} \, ds \\ &= \int_0^1 1 - s \, ds \\ &= s - \frac{1}{2}s^2 \Big|_0^1 \\ &= \frac{1}{2} \end{aligned}$$

The true area of this triangle should be  $\frac{\|(v-u) \times (w-u)\|}{2}$ , so we can see we're off by a factor of  $\|(v - u) \times (w - u)\|$ , which is caused by fact that these basis vectors are no longer unit length.