

1 Physical Simulation

1.1 Rigid Body Time Integrator

$$L(c^i, \theta^i, c^{i+1}, \theta^{i+1}) = \sum_{bodies} \left[\frac{\rho V}{2} \left\| \frac{c^{i+1} - c^i}{h} \right\|^2 + \frac{\rho}{2} \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \omega(\theta^i, \theta^{i+1}) \right] - V(c^{i+1}, \theta^{i+1})$$

$$\frac{\partial}{\partial c^i} L(c^i, \theta^i, c^{i+1}, \theta^{i+1}) = \sum_{bodies} \left[-\rho V \frac{c^{i+1} - c^i}{h^2} \right]^T$$

$$\begin{aligned} \frac{\partial}{\partial \theta^i} L(c^i, \theta^i, c^{i+1}, \theta^{i+1}) &= \sum_{bodies} \frac{\partial}{\partial \theta^i} \left[\frac{\rho}{2} \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \omega(\theta^i, \theta^{i+1}) \right] \\ &= \sum_{bodies} \left[\rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I \left(\frac{\partial}{\partial \theta^i} [R_{-\theta^i} \omega(\theta^i, \theta^{i+1})] \right) \right] \\ &= \sum_{bodies} \left[\rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I \left(\frac{\partial}{\partial \theta^i} R_{-\theta^i} \omega(\theta^i, \theta^{i+1}) + R_{-\theta^i} \frac{\partial}{\partial \theta^i} \omega(\theta^i, \theta^{i+1}) \right) \right] \\ &= \sum_{bodies} \left[\rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I \left(-R_{-\theta^i} [\omega(\theta^i, \theta^{i+1})]_{\times} T(-\theta_i) + R_{-\theta^i} \frac{-1}{h} T(h\omega(\theta^i, \theta^{i+1}))^{-1} T(-\theta^i) \right) \right] \\ &= \sum_{bodies} \left[\rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left(-[\omega(\theta^i, \theta^{i+1})]_{\times} T(-\theta_i) + \frac{-1}{h} T(h\omega(\theta^i, \theta^{i+1}))^{-1} T(-\theta^i) \right) \right] \end{aligned}$$

$$\frac{\partial}{\partial c^{i+1}} L(c^i, \theta^i, c^{i+1}, \theta^{i+1}) = \sum_{bodies} \left[\rho V \frac{c^{i+1} - c^i}{h^2} \right]^T - \frac{\partial}{\partial c^{i+1}} V(c^{i+1}, \theta^{i+1})$$

$$\begin{aligned} \frac{\partial}{\partial \theta^{i+1}} L(c^i, \theta^i, c^{i+1}, \theta^{i+1}) &= \sum_{bodies} \frac{\partial}{\partial \theta^{i+1}} \left[\frac{\rho}{2} \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \omega(\theta^i, \theta^{i+1}) \right] - \frac{\partial}{\partial \theta^{i+1}} V(\theta^i, \theta^{i+1}) \\ &= \sum_{bodies} \left[\rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left(\frac{\partial}{\partial \theta^{i+1}} \omega(\theta^i, \theta^{i+1}) \right) \right] - \frac{\partial}{\partial \theta^{i+1}} V(\theta^i, \theta^{i+1}) \\ &= \sum_{bodies} \left[\rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left(\frac{1}{h} T(-h\omega(\theta^i, \theta^{i+1}))^{-1} T(-\theta^{i+1}) \right) \right] - \frac{\partial}{\partial \theta^{i+1}} V(\theta^i, \theta^{i+1}) \end{aligned}$$

1.2 Euler-Lagrange Equations

$$d_2 L(q^{i-1}, q^i) + d_1 L(q^i, q^{i+1}) = 0$$

For the tangential velocity, this is simply particle motion, and follows the force = mass velocity form.

$$\begin{aligned} \sum_{bodies} \left[\rho V \frac{c^i - c^{i-1}}{h^2} \right]^T - \frac{\partial}{\partial c^i} V(c^i, \theta^i) + \sum_{bodies} \left[-\rho V \frac{c^{i+1} - c^i}{h^2} \right]^T &= 0 \\ \sum_{bodies} \left[\rho V \frac{c^{i+1} - c^i}{h^2} \right]^T &= \sum_{bodies} \left[\rho V \frac{c^i - c^{i-1}}{h^2} \right]^T - \frac{\partial}{\partial c^i} V(c^i, \theta^i) \\ \sum_{bodies} [\rho V \dot{c}_{i+1}] &= \sum_{bodies} [\rho V \dot{c}_i] - h \frac{\partial}{\partial c^i} V(c^i, \theta^i)^T \\ c_{i+1} &= c^i + h \dot{c}^i \\ \dot{c}_{i+1} &= \dot{c}_i - h \frac{1}{\rho V} dV(c^i, \theta^i)^T \quad \forall i \in bodies \end{aligned}$$

Solving for angular velocity, we can see each rigid body is an independent system, and since the potential $V(c, \theta)$ in this case only includes gravity, $d_2 V(c, \theta) = 0$.

$$\begin{aligned} \left[\rho \omega(\theta^{i-1}, \theta^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} \left(\frac{1}{h} T(-h\omega(\theta^{i-1}, \theta^i))^{-1} T(-\theta^i) \right) \right] + \left[\rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left([\omega(\theta^i, \theta^{i+1})]_{\times} T(-\theta_i) + \frac{-1}{h} T(h\omega(\theta^i, \theta^{i+1}))^{-1} \right) T(-\theta^i) \right] &= 0 \\ \left[\rho \omega(\theta^{i-1}, \theta^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} \frac{1}{h} T(-h\omega(\theta^{i-1}, \theta^i))^{-1} \right] + \left[\rho \omega(\theta^i, \theta^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left([\omega(\theta^i, \theta^{i+1})]_{\times} T(-\theta_i) + \frac{-1}{h} T(h\omega(\theta^i, \theta^{i+1}))^{-1} \right) \right] &= 0 \end{aligned}$$

The goal is to compute the orientation at the next time step, θ^{i+1} . This can be done through solving for $\omega(\theta^i, \theta^{i+1})$, since $R_{h\omega(\theta^i, \theta^{i+1})} = R_{\theta^{i+1}} R_{-\theta^i}$.

Using properties of rotation matrices, we can see $R_{-\theta^i}^T R_{h\omega(\theta^i, \theta^{i+1})} = R_{\theta^{i+1}}$.

Let $\omega^i = \omega(\theta^{i-1}, \theta^i)$

$$\begin{aligned} \left[\rho (\omega^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} \frac{1}{h} T(-h\omega^i)^{-1} \right] + \left[\rho (\omega^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left([\omega^{i+1}]_{\times} T(-\theta_i) + \frac{-1}{h} T(h\omega^{i+1})^{-1} \right) \right] &= 0 \\ \left[(\omega^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} T(-h\omega^i)^{-1} \right] + \left[(\omega^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left(h[\omega^{i+1}]_{\times} T(-\theta_i) - T(h\omega^{i+1})^{-1} \right) \right] &= 0 \end{aligned}$$

We can use Newton's method to solve for ω^{i+1} , where

$$\begin{aligned} f(\omega^{i+1}) &= \left[(\omega^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} T(-h\omega^i)^{-1} \right] + \left[(\omega^{i+1})^T R_{-\theta^i}^T M_I R_{-\theta^i} \left(h[\omega^{i+1}]_{\times} T(-\theta_i) - T(h\omega^{i+1})^{-1} \right) \right] \\ &= 0 \\ df(\omega^{i+1}) &= R_{-\theta^i}^T M_I R_{-\theta^i} \left(h[\omega^{i+1}]_{\times} T(-\theta_i) - T(h\omega^{i+1})^{-1} \right) + (\omega^i)^T R_{-\theta^{i-1}}^T M_I R_{-\theta^{i-1}} \left(h d[\omega^{i+1}]_{\times} T(-\theta_i) - h dT(h\omega^{i+1})^{-1} \right) \\ &\approx R_{-\theta^i}^T M_I R_{-\theta^i} \left(-T(h\omega^{i+1})^{-1} \right) \end{aligned}$$

1.3 Integration over a Triangle

Given - arbitrary points of a triangle u, v, w , and a function f that you want to integrate over the face of the triangle.

Let's consider the basic function $f(u, v, w) = 1$ for now, so this integration should give us the area of the triangle - $\|(v - u) \times (w - u)\|$.

First step is to parameterize this into a function of two parameters, s, t , using the basis vectors $v - u$, and $w - u$.

This can be thought of sliding across the plane that is covered by the triangle.

This leads us to the double integral -

$$\begin{aligned} \int_0^1 \int_0^{1-s} f(s, t) \, dt \, ds &= \int_0^1 t|_0^{1-s} \, ds \\ &= \int_0^1 1 - s \, ds \\ &= s - \frac{1}{2}s^2 \Big|_0^1 \\ &= \frac{1}{2} \end{aligned}$$

The true area of this triangle should be $\frac{\|(v-u) \times (w-u)\|}{2}$, so we can see we're off by a factor of $\|(v - u) \times (w - u)\|$, which is caused by fact that these basis vectors are no longer unit length.