
QHD- ∞ Heisenberg Equations of Motion and the Cross Terms Approximation

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Introduction

The evolution of the average of an observable A in the Heisenberg representation is given by the equation of motion (EOM):

$$i \hbar \frac{d}{dt} \langle A \rangle = \langle [A, H] \rangle$$

Consider the averages of momentum, position and their products $\langle p \rangle$, $\langle q \rangle$, $\langle p^2 \rangle$, $\langle q^2 \rangle$, $\langle pq \rangle$, $\langle p^3 \rangle$, $\langle p^2 q \rangle$... The EOMs for the average values are coupled and, in general, form an infinite hierarchy of equations. This document shows how to get some of the equations of that infinite hierarchy. Subscripts will be used in this document to represent different particles (different modes). Products are symmetrized, for example $\langle p_1 \cdot q_1 \rangle_s = \left\langle \frac{p_1 \cdot q_1 + q_1 \cdot p_1}{2} \right\rangle$, and cross terms are approximated, for instance $\langle p_1 \cdot p_2 \rangle \approx \langle p_1 \rangle \langle p_2 \rangle$. Here we obtain equations from the appendix of Pahl and Prezhdo in J. Chem Phys. Vol 116 No. 20, May 2002, Pages 8704-8712

<http://homepage.cem.itesm.mx/lgoomez/quantum/QHDHigherOrders.pdf> .

Calculations are performed using QUANTUM, a free *Mathematica* add-on that can be downloaded from:

<http://homepage.cem.itesm.mx/lgoomez/quantum/>

Load the Package

First load the Quantum`QHD` package. Write:

Needs["Quantum`QHD`"]

then press at the same time the keys `SHIFT-ENTER` to evaluate. *Mathematica* will load the package and print a welcome message:

Needs ["Quantum`QHD`"]

Quantum`QHD`

A Mathematica package for Quantized Hamilton

Dynamics approximation to Heisenberg Equations of Motion

by José Luis Gómez-Muñoz

based on the original idea of Kirill Igumenshchev

This add-on does NOT work properly with the debugger turned on. Therefore the debugger must NOT be checked in the Evaluation menu of Mathematica.

Execute SetQHDAliases[] in order to use the keyboard to enter QHD objects

SetQHDAliases[] must be executed again in each new notebook that is created

In order to use the keyboard to enter quantum objects write:

SetQHDAliases[];

then press at the same time the keys **SHIFT-ENTER** to evaluate. Remember that SetQuantumAliases[] must be evaluated again in each new notebook:

SetQHDAliases[]

ALIASES:

[ESC]on[ESC]	• Quantum concatenation symbol
[ESC]time[ESC]	t Time symbol
[ESC]hb[ESC]	\hbar Reduced Planck's constant (h bar)
[ESC]ii[ESC]	i Imaginary I symbol
[ESC]inf[ESC]	∞ Infinity symbol
[ESC]->[ESC]	\rightarrow Option (Rule) symbol
[ESC]ave[ESC]	$\langle \square \rangle$ Quantum average template
[ESC]expec[ESC]	$\langle \square \rangle$ Quantum average template
[ESC]symm[ESC]	$(\square \cdot \square)_s$ Symmetrized quantum product template
[ESC]comm[ESC]	$[\square, \square]_-$ Commutator template
[ESC]po[ESC]	$(\square)^\square$ Power template
[ESC]su[ESC]	\square_\square Subscripted variable template
[ESC]posu[ESC]	\square_\square^\square Power of a subscripted variable template
[ESC]fra[ESC]	\square — Fraction template
[ESC]eva[ESC]	$\square/.{\square \rightarrow \square, \square \rightarrow \square}$ Evaluation (ReplaceAll) template

SetQHDAliases[] must be executed again in each

new notebook that is created, only one time per notebook.

This Hamiltonian and commutation relationships are those used by E. Pahl and O V. Prezhdo in their paper.

In order to enter the templates and symbols \square , \square_a , $[[\square, \square]]$ and i you can either use the QHD palette (toolbar) or press the keys [ESC]su[ESC], [ESC]supo[ESC], [ESC]comm[ESC] and [ESC]ii[ESC]. The Hamiltonian that will be used in this document is stored in the variable *inh* below:

```
SetQuantumObject[q, p];
[[q1, p1]]_ = i * h;
[[q2, p2]]_ = i * h;
[[q1, p2]]_ = 0;
[[q2, p1]]_ = 0;
[[q1, q2]]_ = 0;
[[p1, p2]]_ = 0;
inh =  $\frac{k_1}{2} p_1^2 + \frac{a_2}{2} q_1^2 + \frac{k_2}{2} p_2^2 + b_1 * q_2 + \frac{1}{2} b_2 * q_2^2 + \frac{b_3}{3} q_2^3 + \lambda * q_1^2 \cdot q_2$ 
```

$$\frac{1}{2} k_1 p_1^2 + \frac{1}{2} k_2 p_2^2 + \frac{1}{2} a_2 q_1^2 + b_1 q_2 + \frac{1}{2} b_2 q_2^2 + \frac{1}{3} b_3 q_2^3 + \lambda q_1^2 \cdot q_2$$

Heisenberg Equations of Motion (EOM)

The Heisenberg EOM:

$$i \hbar \frac{d}{dt} \langle p_1 \rangle = \langle [p_1, H] \rangle$$

can be calculated using the QHD-*Mathematica* command QHDEOM, as shown below. The first argument specifies the QHD-order, which is set below to ∞ (type [ESC]inf[ESC]) so that no QHD closure is done. The second argument is the observable, p_1 in this example. The third argument is the Hamiltonian, remember that it was stored in variable *inh* above in this document. The output of QHDEOM can be stored in a variable; it is stored in *infeom*, then it is shown in a nice format with the command QHDForm below:

```
infeom = QHDEOM[ $\infty$ , p1, inh];
QHDForm[infeom]
```

No closure was applied
 QHDAproximantFunction->QHDCrossTermsApproximant.

$$\frac{d \langle p_1 \rangle}{dt} = -a_2 \langle q_1 \rangle - 2 \lambda \langle q_1 \rangle \langle q_2 \rangle$$

The command QHDDifferentialEquations formats the output of QHDEOM as a standard *Mathematica* equation, as those that can be part of the input of standard *Mathematica* commands like DSolve and NDSolve. The output *infeom* that was stored above is formatted that way below:

```
QHDDifferentialEquations[infeom]
```

$$\{ \langle p_1 \rangle' [t] == -a_2 \langle q_1 \rangle [t] - 2 \lambda \langle q_1 \rangle [t] \langle q_2 \rangle [t] \}$$

A different symbol for time can be specified. The operator \rightarrow can be entered pressing the keys [ESC][MINUS][GREATERTHAN][ESC], see the equations below:

```
QHDDifferentialEquations[infeom, QHDSymbolForTime  $\rightarrow$  var]
```

$$\{\langle p_1 \rangle'[\text{var}] = -a_2 \langle q_1 \rangle[\text{var}] - 2 \lambda \langle q_1 \rangle[\text{var}] \langle q_2 \rangle[\text{var}]\}$$

Hierarchy of Heisenberg EOMs

The EOM calculated above for $\langle p_1 \rangle$ depends on $\langle q_1 \rangle$ and $\langle q_2 \rangle$. The EOM of those two average values will depend on other new ones, like $\langle p_2 \rangle$ and $\langle p_1 \cdot q_1 \rangle$, and so on, thus an infinite hierarchy of equations is obtained. The command QHDHierarchy can be used in order to obtain the first few equations of that hierarchy. It has the same syntax as the command QHDEOM that was used above; the first argument specifies the QHD-order, which is set below to ∞ so that no QHD closure is done. The second argument is the initial observable, p_1 in this example. The third argument is the Hamiltonian, remember that it was stored in variable *inhf* above in this document. This hierarchy of equations is infinite, calculations stop at second order in the products of the observables. Those products are symmetrized, for example $\langle p_1 \cdot q_1 \rangle_s = \left\langle \frac{p_1 \cdot q_1 + q_1 \cdot p_1}{2} \right\rangle$, and cross terms are approximated, for instance $\langle p_1 \cdot p_2 \rangle \approx \langle p_1 \rangle \langle p_2 \rangle$. The output of QHDHierarchy is stored in the variable *inhfier*, and it is shown using the QHDForm command below:

```
inhfier = QHDHierarchy[ $\infty$ , p1, inhf];
QHDForm[ inhfier ]
```

Calculations were stopped at order QHDMaxOrder \rightarrow 2
QHDApproximantFunction \rightarrow QHDCrossTermsApproximant.

$$\frac{d \langle p_1 \rangle}{dt} = -a_2 \langle q_1 \rangle - 2 \lambda \langle q_1 \rangle \langle q_2 \rangle$$

$$\frac{d \langle q_1 \rangle}{dt} = k_1 \langle p_1 \rangle$$

$$\frac{d \langle q_2 \rangle}{dt} = k_2 \langle p_2 \rangle$$

$$\frac{d \langle p_2 \rangle}{dt} = -b_1 - \lambda \langle q_1^2 \rangle - b_2 \langle q_2 \rangle - b_3 \langle q_2^2 \rangle$$

$$\frac{d \langle q_1^2 \rangle}{dt} = 2 k_1 \langle p_1 q_1 \rangle_s$$

$$\frac{d \langle q_2^2 \rangle}{dt} = 2 k_2 \langle p_2 q_2 \rangle_s$$

$$\frac{d \langle p_1 q_1 \rangle_s}{dt} = k_1 \langle p_1^2 \rangle - a_2 \langle q_1^2 \rangle - 2 \lambda \langle q_1^2 \rangle \langle q_2 \rangle$$

$$\frac{d \langle p_2 q_2 \rangle_s}{dt} = k_2 \langle p_2^2 \rangle - b_1 \langle q_2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2 \rangle - b_2 \langle q_2^2 \rangle - b_3 \langle q_2^3 \rangle$$

$$\frac{d \langle p_1^2 \rangle}{dt} = -2 a_2 \langle p_1 q_1 \rangle_s - 4 \lambda \langle q_2 \rangle \langle p_1 q_1 \rangle_s$$

$$\frac{d \langle p_2^2 \rangle}{dt} = -2 b_1 \langle p_2 \rangle - 2 \lambda \langle p_2 \rangle \langle q_1^2 \rangle - 2 b_2 \langle p_2 q_2 \rangle_s - 2 b_3 \langle p_2 q_2^2 \rangle_s$$

Not all the variables have their EOM included in the hierarchy, because infinite order ∞ means that no closure procedure was applied. The command QHDForm accepts the option QHDNeedClosureStyle, which allows to specify a different style for those variables whose EOM is not included in the hierarchy, see the bottom of the table below:

```
QHDForm[ infhier, QHDNeedClosureStyle -> {Bold, Larger, Blue} ]
```

Calculations were stopped at order QHDMaxOrder→2
 QHDApproximantFunction→QHDCrossTermsApproximant.

$$\frac{d \langle p_1 \rangle}{dt} = -a_2 \langle q_1 \rangle - 2 \lambda \langle q_1 \rangle \langle q_2 \rangle$$

$$\frac{d \langle q_1 \rangle}{dt} = k_1 \langle p_1 \rangle$$

$$\frac{d \langle q_2 \rangle}{dt} = k_2 \langle p_2 \rangle$$

$$\frac{d \langle p_2 \rangle}{dt} = -b_1 - \lambda \langle q_1^2 \rangle - b_2 \langle q_2 \rangle - b_3 \langle q_2^2 \rangle$$

$$\frac{d \langle q_1^2 \rangle}{dt} = 2 k_1 \langle p_1 q_1 \rangle_s$$

$$\frac{d \langle q_2^2 \rangle}{dt} = 2 k_2 \langle p_2 q_2 \rangle_s$$

$$\frac{d \langle p_1 q_1 \rangle_s}{dt} = k_1 \langle p_1^2 \rangle - a_2 \langle q_1^2 \rangle - 2 \lambda \langle q_1^2 \rangle \langle q_2 \rangle$$

$$\frac{d \langle p_2 q_2 \rangle_s}{dt} = - \langle \mathbf{q}_2^3 \rangle b_3 + k_2 \langle p_2^2 \rangle - b_1 \langle q_2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2 \rangle - b_2 \langle q_2^2 \rangle$$

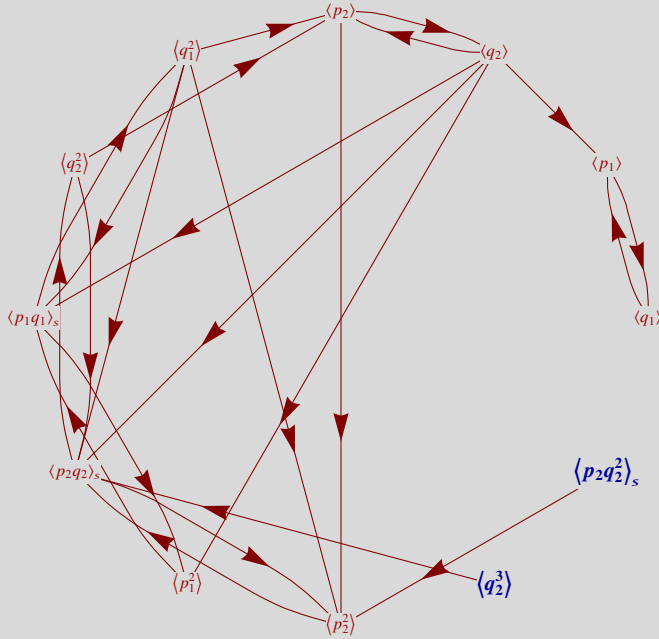
$$\frac{d \langle p_1^2 \rangle}{dt} = -2 a_2 \langle p_1 q_1 \rangle_s - 4 \lambda \langle q_2 \rangle \langle p_1 q_1 \rangle_s$$

$$\frac{d \langle p_2^2 \rangle}{dt} = -2 \langle \mathbf{p}_2 \mathbf{q}_2^2 \rangle_s b_3 - 2 b_1 \langle p_2 \rangle - 2 \lambda \langle p_2 \rangle \langle q_1^2 \rangle - 2 b_2 \langle p_2 q_2 \rangle_s$$

A hierarchy can be shown in a graph using the command QHDGraphPlot on the output of QHDHierarchy. Each arrow points from a first dynamical variable to a second dynamical variable that includes the first one in its EOM; please compare the table above with the graph below:

QHDGraphPlot[infhier]

Calculations were stopped at order QHDMaxOrder→2
 QHDAproximantFunction→QHDCrossTermsApproximant.



The command `QHDDifferentialEquations` formats the output of `QHDHierarchy` as standard *Mathematica* equations, as those that can be part of the input of standard *Mathematica* commands like `DSolve` and `NDSolve`. The output *infhier* that was stored above is formatted that way below:

QHDDifferentialEquations[infhier]

$$\left\{ \begin{aligned} \langle p_1 \rangle' [t] &= -a_2 \langle q_1 \rangle [t] - 2 \lambda \langle q_1 \rangle [t] \langle q_2 \rangle [t], \quad \langle q_1 \rangle' [t] = k_1 \langle p_1 \rangle [t], \\ \langle q_2 \rangle' [t] &= k_2 \langle p_2 \rangle [t], \quad \langle p_2 \rangle' [t] = -b_1 - \lambda \langle q_1^2 \rangle [t] - b_2 \langle q_2 \rangle [t] - b_3 \langle q_2^2 \rangle [t], \\ \langle q_1^2 \rangle' [t] &= 2 k_1 \langle (p_1 \cdot q_1)_s \rangle [t], \quad \langle q_2^2 \rangle' [t] = 2 k_2 \langle (p_2 \cdot q_2)_s \rangle [t], \\ \langle (p_1 \cdot q_1)_s \rangle' [t] &= k_1 \langle p_1^2 \rangle [t] - a_2 \langle q_1^2 \rangle [t] - 2 \lambda \langle q_1^2 \rangle [t] \langle q_2 \rangle [t], \\ \langle (p_2 \cdot q_2)_s \rangle' [t] &= k_2 \langle p_2^2 \rangle [t] - b_1 \langle q_2 \rangle [t] - \lambda \langle q_1^2 \rangle [t] \langle q_2 \rangle [t] - b_2 \langle q_2^2 \rangle [t] - b_3 \langle q_2^3 \rangle [t], \\ \langle p_1^2 \rangle' [t] &= -2 a_2 \langle (p_1 \cdot q_1)_s \rangle [t] - 4 \lambda \langle q_2 \rangle [t] \langle (p_1 \cdot q_1)_s \rangle [t], \\ \langle p_2^2 \rangle' [t] &= -2 b_1 \langle p_2 \rangle [t] - 2 \lambda \langle p_2 \rangle [t] \langle q_1^2 \rangle [t] - 2 b_2 \langle (p_2 \cdot q_2)_s \rangle [t] - 2 b_3 \langle (p_2 \cdot q_2)_s \rangle [t] \end{aligned} \right\}$$

A different symbol for time can be specified. The operator \rightarrow can be entered pressing the keys `[ESC][MINUS][GREATERTHAN][ESC]`, see the equations below:

```
QHDDifferentialEquations[infhier, QHDSymbolForTime → z]
```

$$\begin{aligned} & \{ \langle p_1 \rangle' [z] = -a_2 \langle q_1 \rangle [z] - 2 \lambda \langle q_1 \rangle [z] \langle q_2 \rangle [z], \langle q_1 \rangle' [z] = k_1 \langle p_1 \rangle [z], \\ & \langle q_2 \rangle' [z] = k_2 \langle p_2 \rangle [z], \langle p_2 \rangle' [z] = -b_1 - \lambda \langle q_1^2 \rangle [z] - b_2 \langle q_2 \rangle [z] - b_3 \langle q_2^2 \rangle [z], \\ & \langle q_1^2 \rangle' [z] = 2 k_1 \langle (p_1 \cdot q_1)_s \rangle [z], \langle q_2^2 \rangle' [z] = 2 k_2 \langle (p_2 \cdot q_2)_s \rangle [z], \\ & \langle (p_1 \cdot q_1)_s \rangle' [z] = k_1 \langle p_1^2 \rangle [z] - a_2 \langle q_1^2 \rangle [z] - 2 \lambda \langle q_1^2 \rangle [z] \langle q_2 \rangle [z], \\ & \langle (p_2 \cdot q_2)_s \rangle' [z] = k_2 \langle p_2^2 \rangle [z] - b_1 \langle q_2 \rangle [z] - \lambda \langle q_1^2 \rangle [z] \langle q_2 \rangle [z] - b_2 \langle q_2^2 \rangle [z] - b_3 \langle q_2^3 \rangle [z], \\ & \langle p_1^2 \rangle' [z] = -2 a_2 \langle (p_1 \cdot q_1)_s \rangle [z] - 4 \lambda \langle q_2 \rangle [z] \langle (p_1 \cdot q_1)_s \rangle [z], \langle p_2^2 \rangle' [z] = \\ & -2 b_1 \langle p_2 \rangle [z] - 2 \lambda \langle p_2 \rangle [z] \langle q_1^2 \rangle [z] - 2 b_2 \langle (p_2 \cdot q_2)_s \rangle [z] - 2 b_3 \langle (p_2 \cdot q_2^2)_s \rangle [z] \} \end{aligned}$$

Cross Terms Approximations

The command QHDSymbolForTime accepts several options, among them QHDAproximantFunction, see the evaluation below:

```
Options[QHDHierarchy]
```

```
{QHDSymbolForTime → 2, QHDSymbolForTime → ħ, QHDLabel → Automatic,
  QHDAproximantFunction → QHDCrossTermsApproximant}
```

The default approximant is QHDCrossTermsApproximant, we can see what is the effect of this approximant below:

```
? QHDCrossTermsApproximant
```

QHDCrossTermsApproximant[expr] approximates expected values of crossterms $\langle p_1^n \cdot p_2^m \rangle \approx \langle p_1^n \rangle \langle p_2^m \rangle$ in expr

EOM Hierarchy without Cross Terms Approximation

We can get another hierarchy for the same problem without the cross terms approximation, adding as a last argument of QHDSymbolForTime the option QHDAproximantFunction→Identity, as shown below:

```

inhierNoApp = QHDHierarchy[∞, p1, inh, QHDAproximantFunction → Identity];
QHDForm[ inhierNoApp ]

```

Calculations were stopped at order QHDMaXOrder→2

$$\frac{d \langle p_1 \rangle}{dt} = -a_2 \langle q_1 \rangle - 2 \lambda \langle q_1 q_2 \rangle_s$$

$$\frac{d \langle q_1 \rangle}{dt} = k_1 \langle p_1 \rangle$$

$$\frac{d \langle q_1 q_2 \rangle_s}{dt} = k_2 \langle p_2 q_1 \rangle_s + k_1 \langle p_1 q_2 \rangle_s$$

$$\frac{d \langle p_2 q_1 \rangle_s}{dt} = -b_1 \langle q_1 \rangle - \lambda \langle q_1^3 \rangle + k_1 \langle p_1 p_2 \rangle_s - b_2 \langle q_1 q_2 \rangle_s - b_3 \langle q_1 q_2^2 \rangle_s$$

$$\frac{d \langle p_1 q_2 \rangle_s}{dt} = k_2 \langle p_1 p_2 \rangle_s - a_2 \langle q_1 q_2 \rangle_s - 2 \lambda \langle q_1 q_2^2 \rangle_s$$

$$\frac{d \langle p_1 p_2 \rangle_s}{dt} = -b_1 \langle p_1 \rangle - a_2 \langle p_2 q_1 \rangle_s - \lambda \langle p_1 q_1^2 \rangle_s - b_2 \langle p_1 q_2 \rangle_s - b_3 \langle p_1 q_2^2 \rangle_s - 2 \lambda \langle p_2 q_1 q_2 \rangle_s$$

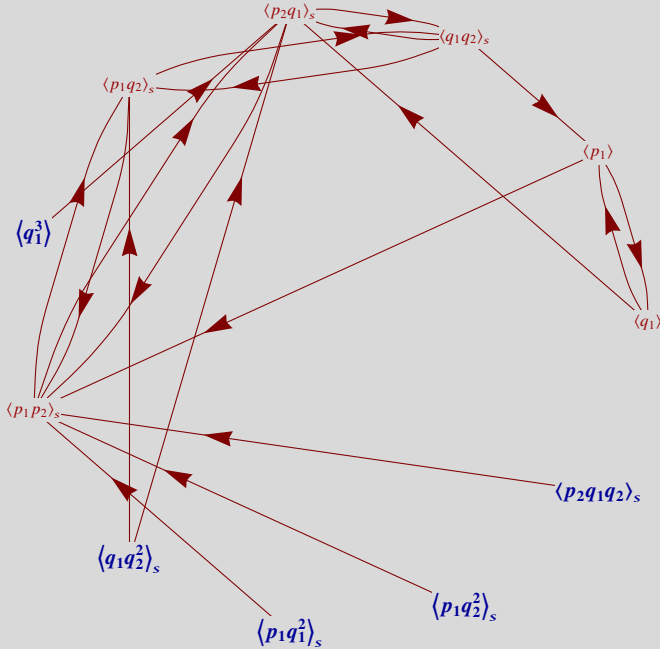
A hierarchy can be shown in a graph using the command QHDGraphPlot on the output of QHDHierarchy. Each arrow points from a first dynamical variable to a second dynamical variable that includes the first one in its EOM; please compare the table above with the graph below:

```

QHDGraphPlot[ inhierNoApp ]

```

Calculations were stopped at order QHDMaXOrder→2



Products of operators were symmetrized, for example: $\langle p_2 \cdot q_1 \cdot q_2 \rangle_s = \frac{1}{2} \langle p_2 \cdot q_1 \cdot q_2 \rangle + \frac{1}{2} \langle q_2 \cdot q_1 \cdot p_2 \rangle$

Hierarchy Up To Third Order

QHDHierarchy stops calculating EOM up to the order specified by its option QHDMaxOrder, which has a default value of 2. Here we obtain the hierarchy up to order 3, this calculation takes several minutes in a laptop computer, see the table below:

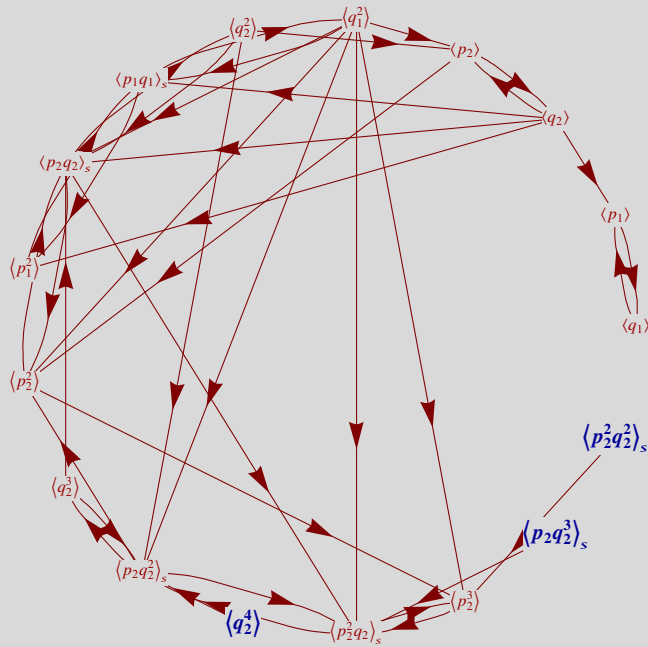
```
infhier3 = QHDHierarchy[∞, p1, infh, QHDMaxOrder → 3];
QHDForm[ infhier3 ]
```

Calculations were stopped at order QHDMaxOrder→3 QHDApproximantFunction→QHDCrossTermsApproximant.
$\frac{d \langle p_1 \rangle}{dt} = -a_2 \langle q_1 \rangle - 2 \lambda \langle q_1 \rangle \langle q_2 \rangle$
$\frac{d \langle q_1 \rangle}{dt} = k_1 \langle p_1 \rangle$
$\frac{d \langle q_2 \rangle}{dt} = k_2 \langle p_2 \rangle$
$\frac{d \langle p_2 \rangle}{dt} = -b_1 - \lambda \langle q_1^2 \rangle - b_2 \langle q_2 \rangle - b_3 \langle q_2^2 \rangle$
$\frac{d \langle q_1^2 \rangle}{dt} = 2 k_1 \langle p_1 q_1 \rangle_s$
$\frac{d \langle q_2^2 \rangle}{dt} = 2 k_2 \langle p_2 q_2 \rangle_s$
$\frac{d \langle p_1 q_1 \rangle_s}{dt} = k_1 \langle p_1^2 \rangle - a_2 \langle q_1^2 \rangle - 2 \lambda \langle q_1^2 \rangle \langle q_2 \rangle$
$\frac{d \langle p_2 q_2 \rangle_s}{dt} = k_2 \langle p_2^2 \rangle - b_1 \langle q_2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2 \rangle - b_2 \langle q_2^2 \rangle - b_3 \langle q_2^3 \rangle$
$\frac{d \langle p_1^2 \rangle}{dt} = -2 a_2 \langle p_1 q_1 \rangle_s - 4 \lambda \langle q_2 \rangle \langle p_1 q_1 \rangle_s$
$\frac{d \langle p_2^2 \rangle}{dt} = -2 b_1 \langle p_2 \rangle - 2 \lambda \langle p_2 \rangle \langle q_1^2 \rangle - 2 b_2 \langle p_2 q_2 \rangle_s - 2 b_3 \langle p_2 q_2^2 \rangle_s$
$\frac{d \langle q_2^3 \rangle}{dt} = 3 k_2 \langle p_2 q_2^2 \rangle_s$
$\frac{d \langle p_2 q_2^2 \rangle_s}{dt} = -b_1 \langle q_2^2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2^2 \rangle - b_2 \langle q_2^3 \rangle - b_3 \langle q_2^4 \rangle + 2 k_2 \langle p_2^2 q_2 \rangle_s$
$\frac{d \langle p_2^2 q_2 \rangle_s}{dt} = k_2 \langle p_2^3 \rangle - 2 b_1 \langle p_2 q_2 \rangle_s - 2 \lambda \langle q_1^2 \rangle \langle p_2 q_2 \rangle_s - 2 b_2 \langle p_2 q_2^2 \rangle_s - 2 b_3 \langle p_2 q_2^3 \rangle_s$
$\frac{d \langle p_2^3 \rangle}{dt} = -\hbar^2 b_3 - 3 b_1 \langle p_2^2 \rangle - 3 \lambda \langle p_2^2 \rangle \langle q_1^2 \rangle - 3 b_2 \langle p_2^2 q_2 \rangle_s - 3 b_3 \langle p_2^2 q_2^2 \rangle_s$

A hierarchy can be shown in a graph using the command QHDGraphPlot on the output of QHDHierarchy. Each arrow points from a first dynamical variable to a second dynamical variable that includes the first one in its EOM; please compare the table above with the graph below:

QHDGraphPlot [infhier3]

Calculations were stopped at order QHDMaxOrder→3
 QHDAproximantFunction→QHDCrossTermsApproximant.



EOM Hierarchy Up To Fourth Order

Here we obtain the hierarchy up to order 4, this calculation takes several minutes in a laptop computer, see the table below:

```

inhier4 = QHDSierarchy[∞, p1, inh, QHDSMaxOrder → 4];
QHDSForm[ inhier4 ]

```

Calculations were stopped at order QHDSMaxOrder→4
 QHDSApproximantFunction→QHDSCrossTermsApproximant.

$$\frac{d \langle p_1 \rangle}{dt} = -a_2 \langle q_1 \rangle - 2 \lambda \langle q_1 \rangle \langle q_2 \rangle$$

$$\frac{d \langle q_1 \rangle}{dt} = k_1 \langle p_1 \rangle$$

$$\frac{d \langle q_2 \rangle}{dt} = k_2 \langle p_2 \rangle$$

$$\frac{d \langle p_2 \rangle}{dt} = -b_1 - \lambda \langle q_1^2 \rangle - b_2 \langle q_2 \rangle - b_3 \langle q_2^2 \rangle$$

$$\frac{d \langle q_1^2 \rangle}{dt} = 2 k_1 \langle p_1 q_1 \rangle_s$$

$$\frac{d \langle q_2^2 \rangle}{dt} = 2 k_2 \langle p_2 q_2 \rangle_s$$

$$\frac{d \langle p_1 q_1 \rangle_s}{dt} = k_1 \langle p_1^2 \rangle - a_2 \langle q_1^2 \rangle - 2 \lambda \langle q_1^2 \rangle \langle q_2 \rangle$$

$$\frac{d \langle p_2 q_2 \rangle_s}{dt} = k_2 \langle p_2^2 \rangle - b_1 \langle q_2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2 \rangle - b_2 \langle q_2^2 \rangle - b_3 \langle q_2^3 \rangle$$

$$\frac{d \langle p_1^2 \rangle}{dt} = -2 a_2 \langle p_1 q_1 \rangle_s - 4 \lambda \langle q_2 \rangle \langle p_1 q_1 \rangle_s$$

$$\frac{d \langle p_2^2 \rangle}{dt} = -2 b_1 \langle p_2 \rangle - 2 \lambda \langle p_2 \rangle \langle q_1^2 \rangle - 2 b_2 \langle p_2 q_2 \rangle_s - 2 b_3 \langle p_2 q_2^2 \rangle_s$$

$$\frac{d \langle q_2^3 \rangle}{dt} = 3 k_2 \langle p_2 q_2^2 \rangle_s$$

$$\frac{d \langle p_2 q_2^2 \rangle_s}{dt} = -b_1 \langle q_2^2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2^2 \rangle - b_2 \langle q_2^3 \rangle - b_3 \langle q_2^4 \rangle + 2 k_2 \langle p_2^2 q_2 \rangle_s$$

$$\frac{d \langle q_2^4 \rangle}{dt} = 4 k_2 \langle p_2 q_2^3 \rangle_s$$

$$\frac{d \langle p_2^2 q_2 \rangle_s}{dt} = k_2 \langle p_2^3 \rangle - 2 b_1 \langle p_2 q_2 \rangle_s - 2 \lambda \langle q_1^2 \rangle \langle p_2 q_2 \rangle_s - 2 b_2 \langle p_2 q_2^2 \rangle_s - 2 b_3 \langle p_2 q_2^3 \rangle_s$$

$$\frac{d \langle p_2 q_2^3 \rangle_s}{dt} = \frac{3 \hbar^2 k_2}{2} - b_1 \langle q_2^3 \rangle - \lambda \langle q_1^2 \rangle \langle q_2^3 \rangle - b_2 \langle q_2^4 \rangle - b_3 \langle q_2^5 \rangle + 3 k_2 \langle p_2^2 q_2^2 \rangle_s$$

$$\frac{d \langle p_2^3 \rangle}{dt} = -\hbar^2 b_3 - 3 b_1 \langle p_2^2 \rangle - 3 \lambda \langle p_2^2 \rangle \langle q_1^2 \rangle - 3 b_2 \langle p_2^2 q_2 \rangle_s - 3 b_3 \langle p_2^2 q_2^2 \rangle_s$$

$$\frac{d \langle p_2^2 q_2^2 \rangle_s}{dt} = 2 k_2 \langle p_2^3 q_2 \rangle_s - 2 b_1 \langle p_2 q_2^2 \rangle_s - 2 \lambda \langle q_1^2 \rangle \langle p_2 q_2^2 \rangle_s - 2 b_2 \langle p_2 q_2^3 \rangle_s - 2 b_3 \langle p_2 q_2^4 \rangle_s$$

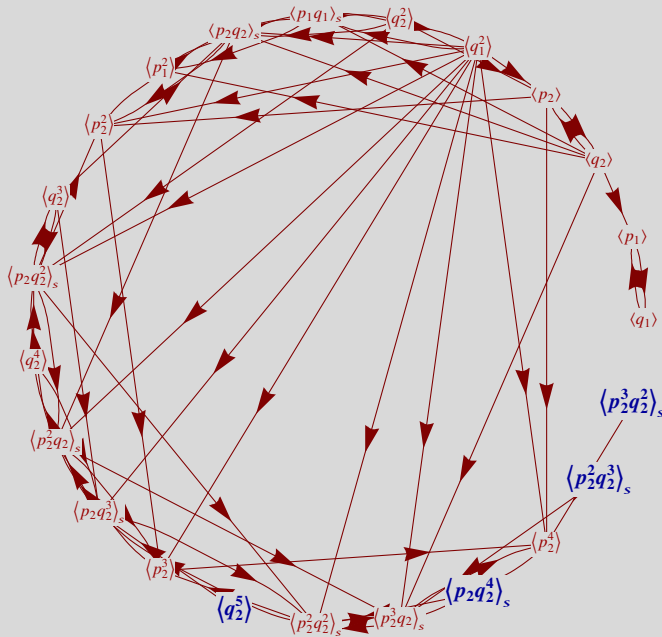
$$\frac{d \langle p_2^3 q_2 \rangle_s}{dt} = -\frac{3}{2} \hbar^2 b_2 + k_2 \langle p_2^4 \rangle - 4 \hbar^2 b_3 \langle q_2 \rangle - 3 b_1 \langle p_2^2 q_2 \rangle_s - 3 \lambda \langle q_1^2 \rangle \langle p_2^2 q_2 \rangle_s - 3 b_2 \langle p_2^2 q_2^2 \rangle_s - 3 b_3 \langle p_2^2 q_2^3 \rangle_s$$

$$\frac{d \langle p_2^4 \rangle}{dt} = -4 \hbar^2 b_3 \langle p_2 \rangle - 4 b_1 \langle p_2^3 \rangle - 4 \lambda \langle p_2^3 \rangle \langle q_1^2 \rangle - 4 b_2 \langle p_2^3 q_2 \rangle_s - 4 b_3 \langle p_2^3 q_2^2 \rangle_s$$

A hierarchy can be shown in a graph using the command QHDSGraphPlot on the output of QHDSHierarchy. Each arrow points from a first dynamical variable to a second dynamical variable that includes the first one in its EOM; please compare the table above with the graph below:

```
QHDGraphPlot [infhier4]
```

Calculations were stopped at order QHDMaxOrder→4
 QHDAproximantFunction→QHDCrossTermsApproximant.



Saving the QHD hierarchies for future use

The calculation QHD hierarchies is time consuming, therefore it is a good idea to save them for future use. The simplest way to do it is using the standard *Mathematica* command Put. The QHD hierarchy is stored in the file hier4.m, see the command below:

```
Put [infhier4, "hier4.m"]
```

The command FileNames["*.m"] can be used to see the files that were created by the command Put (together with any other preexisting file with .m extension), see the command below:

```
FileNames ["*.m"]
```

```
{hier4.m, qhd2.m, qhd3.m, qhd4.m}
```

Loading and Formating the Hierarchy

Next commands have been written as they would be used in a different *Mathematica* session, using the command Get to load the

hierarchy that was saved above (instead of calculating it again). Notice that the commutation relation and the Hamiltonian are defined again, and its important to use those that correspond to the hierarchy that is loaded. The hierarchy is shown below:

```
Needs["Quantum`QHD`"];
SetQHDAliases[];
SetQuantumObject[q, p];
[[q1, p1]]_ = i * h;
[[q2, p2]]_ = i * h;
[[q1, p2]]_ = 0;
[[q2, p1]]_ = 0;
[[q1, q2]]_ = 0;
[[p1, p2]]_ = 0;
inhf =  $\frac{k_1}{2} p_1^2 + \frac{a_2}{2} q_1^2 + \frac{k_2}{2} p_2^2 + b_1 * q_2 + \frac{1}{2} b_2 * q_2^2 + \frac{b_3}{3} q_2^3 + \lambda * q_1^2 \cdot q_2$ ;
myhierarchy = Get["hier4.m"];
QHDForm[myhierarchy]
```

Calculations were stopped at order QHDMaXOrder→4 QHDApProximantFunction→QHDCrossTermsApproximant.
$\frac{d \langle p_1 \rangle}{dt} = -a_2 \langle q_1 \rangle - 2 \lambda \langle q_1 \rangle \langle q_2 \rangle$
$\frac{d \langle q_1 \rangle}{dt} = k_1 \langle p_1 \rangle$
$\frac{d \langle q_2 \rangle}{dt} = k_2 \langle p_2 \rangle$
$\frac{d \langle p_2 \rangle}{dt} = -b_1 - \lambda \langle q_1^2 \rangle - b_2 \langle q_2 \rangle - b_3 \langle q_2^2 \rangle$
$\frac{d \langle q_1^2 \rangle}{dt} = 2 k_1 \langle p_1 q_1 \rangle_s$
$\frac{d \langle q_2^2 \rangle}{dt} = 2 k_2 \langle p_2 q_2 \rangle_s$
$\frac{d \langle p_1 q_1 \rangle_s}{dt} = k_1 \langle p_1^2 \rangle - a_2 \langle q_1^2 \rangle - 2 \lambda \langle q_1^2 \rangle \langle q_2 \rangle$
$\frac{d \langle p_2 q_2 \rangle_s}{dt} = k_2 \langle p_2^2 \rangle - b_1 \langle q_2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2 \rangle - b_2 \langle q_2^2 \rangle - b_3 \langle q_2^3 \rangle$
$\frac{d \langle p_1^2 \rangle}{dt} = -2 a_2 \langle p_1 q_1 \rangle_s - 4 \lambda \langle q_2 \rangle \langle p_1 q_1 \rangle_s$
$\frac{d \langle p_2^2 \rangle}{dt} = -2 b_1 \langle p_2 \rangle - 2 \lambda \langle p_2 \rangle \langle q_1^2 \rangle - 2 b_2 \langle p_2 q_2 \rangle_s - 2 b_3 \langle p_2 q_2^2 \rangle_s$
$\frac{d \langle q_2^3 \rangle}{dt} = 3 k_2 \langle p_2 q_2^2 \rangle_s$
$\frac{d \langle p_2 q_2^2 \rangle_s}{dt} = -b_1 \langle q_2^2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2^2 \rangle - b_2 \langle q_2^3 \rangle - b_3 \langle q_2^4 \rangle + 2 k_2 \langle p_2^2 q_2 \rangle_s$
$\frac{d \langle q_2^4 \rangle}{dt} = 4 k_2 \langle p_2 q_2^3 \rangle_s$
$\frac{d \langle p_2^2 q_2 \rangle_s}{dt} = k_2 \langle p_2^3 \rangle - 2 b_1 \langle p_2 q_2 \rangle_s - 2 \lambda \langle q_1^2 \rangle \langle p_2 q_2 \rangle_s - 2 b_2 \langle p_2 q_2^2 \rangle_s - 2 b_3 \langle p_2 q_2^3 \rangle_s$
$\frac{d \langle p_2 q_2^3 \rangle_s}{dt} = \frac{3 \hbar^2 k_2}{2} - b_1 \langle q_2^3 \rangle - \lambda \langle q_1^2 \rangle \langle q_2^3 \rangle - b_2 \langle q_2^4 \rangle - b_3 \langle q_2^5 \rangle + 3 k_2 \langle p_2^2 q_2^2 \rangle_s$
$\frac{d \langle p_2^3 \rangle}{dt} = -\hbar^2 b_3 - 3 b_1 \langle p_2^2 \rangle - 3 \lambda \langle p_2^2 \rangle \langle q_1^2 \rangle - 3 b_2 \langle p_2^2 q_2 \rangle_s - 3 b_3 \langle p_2^2 q_2^2 \rangle_s$
$\frac{d \langle p_2^2 q_2^2 \rangle_s}{dt} = 2 k_2 \langle p_2^3 q_2 \rangle_s - 2 b_1 \langle p_2 q_2^2 \rangle_s - 2 \lambda \langle q_1^2 \rangle \langle p_2 q_2^2 \rangle_s - 2 b_2 \langle p_2 q_2^3 \rangle_s - 2 b_3 \langle p_2 q_2^4 \rangle_s$
$\frac{d \langle p_2^3 q_2 \rangle_s}{dt} = -\frac{3}{2} \hbar^2 b_2 + k_2 \langle p_2^4 \rangle - 4 \hbar^2 b_3 \langle q_2 \rangle - 3 b_1 \langle p_2^2 q_2 \rangle_s - 3 \lambda \langle q_1^2 \rangle \langle p_2^2 q_2 \rangle_s - 3 b_2 \langle p_2^2 q_2^2 \rangle_s - 3 b_3 \langle p_2^2 q_2^3 \rangle_s$
$\frac{d \langle p_2^4 \rangle}{dt} = -4 \hbar^2 b_3 \langle p_2 \rangle - 4 b_1 \langle p_2^3 \rangle - 4 \lambda \langle p_2^3 \rangle \langle q_1^2 \rangle - 4 b_2 \langle p_2^3 q_2 \rangle_s - 4 b_3 \langle p_2^3 q_2^2 \rangle_s$

The hierarchy can be shown with a specific label in the first row, see the table below:

QHDForm[myhierarchy,

QHDLLabel → "Compare with Pahl and Prezhdo\nJ.Chem.Phys., Vol 116, No.
20, May 2002\nPages 8704-8712\nEquations A2 of the Appendix"]

Compare with Pahl and Prezhdo

J.Chem.Phys., Vol 116, No. 20, May 2002

Pages 8704-8712

Equations A2 of the Appendix

$$\frac{d \langle p_1 \rangle}{dt} = -a_2 \langle q_1 \rangle - 2 \lambda \langle q_1 \rangle \langle q_2 \rangle$$

$$\frac{d \langle q_1 \rangle}{dt} = k_1 \langle p_1 \rangle$$

$$\frac{d \langle q_2 \rangle}{dt} = k_2 \langle p_2 \rangle$$

$$\frac{d \langle p_2 \rangle}{dt} = -b_1 - \lambda \langle q_1^2 \rangle - b_2 \langle q_2 \rangle - b_3 \langle q_2^2 \rangle$$

$$\frac{d \langle q_1^2 \rangle}{dt} = 2 k_1 \langle p_1 q_1 \rangle_s$$

$$\frac{d \langle q_2^2 \rangle}{dt} = 2 k_2 \langle p_2 q_2 \rangle_s$$

$$\frac{d \langle p_1 q_1 \rangle_s}{dt} = k_1 \langle p_1^2 \rangle - a_2 \langle q_1^2 \rangle - 2 \lambda \langle q_1^2 \rangle \langle q_2 \rangle$$

$$\frac{d \langle p_2 q_2 \rangle_s}{dt} = k_2 \langle p_2^2 \rangle - b_1 \langle q_2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2 \rangle - b_2 \langle q_2^2 \rangle - b_3 \langle q_2^3 \rangle$$

$$\frac{d \langle p_1^2 \rangle}{dt} = -2 a_2 \langle p_1 q_1 \rangle_s - 4 \lambda \langle q_2 \rangle \langle p_1 q_1 \rangle_s$$

$$\frac{d \langle p_2^2 \rangle}{dt} = -2 b_1 \langle p_2 \rangle - 2 \lambda \langle p_2 \rangle \langle q_1^2 \rangle - 2 b_2 \langle p_2 q_2 \rangle_s - 2 b_3 \langle p_2 q_2^2 \rangle_s$$

$$\frac{d \langle q_2^3 \rangle}{dt} = 3 k_2 \langle p_2 q_2^2 \rangle_s$$

$$\frac{d \langle p_2 q_2^2 \rangle_s}{dt} = -b_1 \langle q_2^2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2^2 \rangle - b_2 \langle q_2^3 \rangle - b_3 \langle q_2^4 \rangle + 2 k_2 \langle p_2^2 q_2 \rangle_s$$

$$\frac{d \langle q_2^4 \rangle}{dt} = 4 k_2 \langle p_2 q_2^3 \rangle_s$$

$$\frac{d \langle p_2^2 q_2 \rangle_s}{dt} = k_2 \langle p_2^3 \rangle - 2 b_1 \langle p_2 q_2 \rangle_s - 2 \lambda \langle q_1^2 \rangle \langle p_2 q_2 \rangle_s - 2 b_2 \langle p_2 q_2^2 \rangle_s - 2 b_3 \langle p_2 q_2^3 \rangle_s$$

$$\frac{d \langle p_2 q_2^3 \rangle_s}{dt} = \frac{3 \hbar^2 k_2}{2} - b_1 \langle q_2^3 \rangle - \lambda \langle q_1^2 \rangle \langle q_2^3 \rangle - b_2 \langle q_2^4 \rangle - b_3 \langle q_2^5 \rangle + 3 k_2 \langle p_2^2 q_2^2 \rangle_s$$

$$\frac{d \langle p_2^3 \rangle}{dt} = -\hbar^2 b_3 - 3 b_1 \langle p_2^2 \rangle - 3 \lambda \langle p_2^2 \rangle \langle q_1^2 \rangle - 3 b_2 \langle p_2^2 q_2 \rangle_s - 3 b_3 \langle p_2^2 q_2^2 \rangle_s$$

$$\frac{d \langle p_2^2 q_2^2 \rangle_s}{dt} = 2 k_2 \langle p_2^3 q_2 \rangle_s - 2 b_1 \langle p_2 q_2^2 \rangle_s - 2 \lambda \langle q_1^2 \rangle \langle p_2 q_2^2 \rangle_s - 2 b_2 \langle p_2 q_2^3 \rangle_s - 2 b_3 \langle p_2 q_2^4 \rangle_s$$

$$\frac{d \langle p_2^3 q_2 \rangle_s}{dt} = -\frac{3}{2} \hbar^2 b_2 + k_2 \langle p_2^4 \rangle - 4 \hbar^2 b_3 \langle q_2 \rangle - 3 b_1 \langle p_2^2 q_2 \rangle_s - 3 \lambda \langle q_1^2 \rangle \langle p_2^2 q_2 \rangle_s - 3 b_2 \langle p_2^2 q_2^2 \rangle_s - 3 b_3 \langle p_2^2 q_2^3 \rangle_s$$

$$\frac{d \langle p_2^3 \rangle}{dt} = -4 \hbar^2 b_3 \langle p_2 \rangle - 4 b_1 \langle p_2^3 \rangle - 4 \lambda \langle p_2^3 \rangle \langle q_1^2 \rangle - 4 b_2 \langle p_2^3 q_2 \rangle_s - 4 b_3 \langle p_2^3 q_2^2 \rangle_s$$

The hierarchy can be shown without any label, see the table below:

QHDForm[myhierarchy, QHDLLabel → None]

$\frac{d \langle p_1 \rangle}{dt} = -a_2 \langle q_1 \rangle - 2 \lambda \langle q_1 \rangle \langle q_2 \rangle$
$\frac{d \langle q_1 \rangle}{dt} = k_1 \langle p_1 \rangle$
$\frac{d \langle q_2 \rangle}{dt} = k_2 \langle p_2 \rangle$
$\frac{d \langle p_2 \rangle}{dt} = -b_1 - \lambda \langle q_1^2 \rangle - b_2 \langle q_2 \rangle - b_3 \langle q_2^2 \rangle$
$\frac{d \langle q_1^2 \rangle}{dt} = 2 k_1 \langle p_1 q_1 \rangle_s$
$\frac{d \langle q_2^2 \rangle}{dt} = 2 k_2 \langle p_2 q_2 \rangle_s$
$\frac{d \langle p_1 q_1 \rangle_s}{dt} = k_1 \langle p_1^2 \rangle - a_2 \langle q_1^2 \rangle - 2 \lambda \langle q_1^2 \rangle \langle q_2 \rangle$
$\frac{d \langle p_2 q_2 \rangle_s}{dt} = k_2 \langle p_2^2 \rangle - b_1 \langle q_2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2 \rangle - b_2 \langle q_2^2 \rangle - b_3 \langle q_2^3 \rangle$
$\frac{d \langle p_1^2 \rangle}{dt} = -2 a_2 \langle p_1 q_1 \rangle_s - 4 \lambda \langle q_2 \rangle \langle p_1 q_1 \rangle_s$
$\frac{d \langle p_2^2 \rangle}{dt} = -2 b_1 \langle p_2 \rangle - 2 \lambda \langle p_2 \rangle \langle q_1^2 \rangle - 2 b_2 \langle p_2 q_2 \rangle_s - 2 b_3 \langle p_2 q_2^2 \rangle_s$
$\frac{d \langle q_2^3 \rangle}{dt} = 3 k_2 \langle p_2 q_2^2 \rangle_s$
$\frac{d \langle p_2 q_2^2 \rangle_s}{dt} = -b_1 \langle q_2^2 \rangle - \lambda \langle q_1^2 \rangle \langle q_2^2 \rangle - b_2 \langle q_2^3 \rangle - b_3 \langle q_2^4 \rangle + 2 k_2 \langle p_2^2 q_2 \rangle_s$
$\frac{d \langle q_2^4 \rangle}{dt} = 4 k_2 \langle p_2 q_2^3 \rangle_s$
$\frac{d \langle p_2^2 q_2 \rangle_s}{dt} = k_2 \langle p_2^3 \rangle - 2 b_1 \langle p_2 q_2 \rangle_s - 2 \lambda \langle q_1^2 \rangle \langle p_2 q_2 \rangle_s - 2 b_2 \langle p_2 q_2^2 \rangle_s - 2 b_3 \langle p_2 q_2^3 \rangle_s$
$\frac{d \langle p_2 q_2^3 \rangle_s}{dt} = \frac{3 \hbar^2 k_2}{2} - b_1 \langle q_2^3 \rangle - \lambda \langle q_1^2 \rangle \langle q_2^3 \rangle - b_2 \langle q_2^4 \rangle - b_3 \langle q_2^5 \rangle + 3 k_2 \langle p_2^2 q_2^2 \rangle_s$
$\frac{d \langle p_2^3 \rangle}{dt} = -\hbar^2 b_3 - 3 b_1 \langle p_2^2 \rangle - 3 \lambda \langle p_2^2 \rangle \langle q_1^2 \rangle - 3 b_2 \langle p_2^2 q_2 \rangle_s - 3 b_3 \langle p_2^2 q_2^2 \rangle_s$
$\frac{d \langle p_2^2 q_2^2 \rangle_s}{dt} = 2 k_2 \langle p_2^3 q_2 \rangle_s - 2 b_1 \langle p_2 q_2^2 \rangle_s - 2 \lambda \langle q_1^2 \rangle \langle p_2 q_2^2 \rangle_s - 2 b_2 \langle p_2 q_2^3 \rangle_s - 2 b_3 \langle p_2 q_2^4 \rangle_s$
$\frac{d \langle p_2^3 q_2 \rangle_s}{dt} = -\frac{3}{2} \hbar^2 b_2 + k_2 \langle p_2^4 \rangle - 4 \hbar^2 b_3 \langle q_2 \rangle - 3 b_1 \langle p_2^2 q_2 \rangle_s - 3 \lambda \langle q_1^2 \rangle \langle p_2^2 q_2 \rangle_s - 3 b_2 \langle p_2^2 q_2^2 \rangle_s - 3 b_3 \langle p_2^2 q_2^3 \rangle_s$
$\frac{d \langle p_2^4 \rangle}{dt} = -4 \hbar^2 b_3 \langle p_2 \rangle - 4 b_1 \langle p_2^3 \rangle - 4 \lambda \langle p_2^3 \rangle \langle q_1^2 \rangle - 4 b_2 \langle p_2^3 q_2 \rangle_s - 4 b_3 \langle p_2^3 q_2^2 \rangle_s$

Several options can be specified for QHDForm. Averages that need closure can be shown in a different style, see the bottom of the table below:


```

QHDForm[ myhierarchy,
  QHDLLabel → None,
  QHDSymbolForTime → T,
  QHDClosureStyle → {Black, Bold},
  QHDNeedClosureStyle → {Blue, Bold, Larger}]

```

$\frac{d \langle p_1 \rangle}{d T} = -2 \lambda \langle q_1 \rangle \langle q_2 \rangle - \langle q_1 \rangle a_2$
$\frac{d \langle q_1 \rangle}{d T} = \langle p_1 \rangle k_1$
$\frac{d \langle q_2 \rangle}{d T} = \langle p_2 \rangle k_2$
$\frac{d \langle p_2 \rangle}{d T} = -\lambda \langle q_1^2 \rangle - b_1 - \langle q_2 \rangle b_2 - \langle q_2^2 \rangle b_3$
$\frac{d \langle q_1^2 \rangle}{d T} = 2 \langle p_1 q_1 \rangle_s k_1$
$\frac{d \langle q_2^2 \rangle}{d T} = 2 \langle p_2 q_2 \rangle_s k_2$
$\frac{d \langle p_1 q_1 \rangle_s}{d T} = -2 \lambda \langle q_1^2 \rangle \langle q_2 \rangle - \langle q_1^2 \rangle a_2 + \langle p_1^2 \rangle k_1$
$\frac{d \langle p_2 q_2 \rangle_s}{d T} = -\lambda \langle q_1^2 \rangle \langle q_2 \rangle - \langle q_2 \rangle b_1 - \langle q_2^2 \rangle b_2 - \langle q_2^3 \rangle b_3 + \langle p_2^2 \rangle k_2$
$\frac{d \langle p_1^2 \rangle}{d T} = -4 \lambda \langle q_2 \rangle \langle p_1 q_1 \rangle_s - 2 \langle p_1 q_1 \rangle_s a_2$
$\frac{d \langle p_2^2 \rangle}{d T} = -2 \lambda \langle p_2 \rangle \langle q_1^2 \rangle - 2 \langle p_2 \rangle b_1 - 2 \langle p_2 q_2 \rangle_s b_2 - 2 \langle p_2 q_2^2 \rangle_s b_3$
$\frac{d \langle q_2^3 \rangle}{d T} = 3 \langle p_2 q_2^2 \rangle_s k_2$
$\frac{d \langle p_2 q_2^2 \rangle_s}{d T} = -\lambda \langle q_1^2 \rangle \langle q_2^2 \rangle - \langle q_2^2 \rangle b_1 - \langle q_2^3 \rangle b_2 - \langle q_2^4 \rangle b_3 + 2 \langle p_2^2 q_2 \rangle_s k_2$
$\frac{d \langle q_2^4 \rangle}{d T} = 4 \langle p_2 q_2^3 \rangle_s k_2$
$\frac{d \langle p_2^2 q_2 \rangle_s}{d T} = -2 \lambda \langle q_1^2 \rangle \langle p_2 q_2 \rangle_s - 2 \langle p_2 q_2 \rangle_s b_1 - 2 \langle p_2 q_2^2 \rangle_s b_2 - 2 \langle p_2 q_2^3 \rangle_s b_3 + \langle p_2^3 \rangle k_2$
$\frac{d \langle p_2 q_2^3 \rangle_s}{d T} = -\lambda \langle q_1^2 \rangle \langle q_2^3 \rangle - \langle q_2^3 \rangle b_1 - \langle q_2^4 \rangle b_2 - \langle q_2^5 \rangle b_3 + \frac{3 \hbar^2 k_2}{2} + 3 \langle p_2^2 q_2^2 \rangle_s k_2$
$\frac{d \langle p_2^3 \rangle}{d T} = -3 \lambda \langle p_2^2 \rangle \langle q_1^2 \rangle - 3 \langle p_2^2 \rangle b_1 - 3 \langle p_2^2 q_2 \rangle_s b_2 - \hbar^2 b_3 - 3 \langle p_2^2 q_2^2 \rangle_s b_3$
$\frac{d \langle p_2^2 q_2^2 \rangle_s}{d T} = -2 \lambda \langle q_1^2 \rangle \langle p_2 q_2^2 \rangle_s - 2 \langle p_2 q_2^2 \rangle_s b_1 - 2 \langle p_2 q_2^3 \rangle_s b_2 - 2 \langle p_2 q_2^4 \rangle_s b_3 + 2 \langle p_2^3 q_2 \rangle_s k_2$
$\frac{d \langle p_2^3 q_2 \rangle_s}{d T} = -3 \lambda \langle q_1^2 \rangle \langle p_2^2 q_2 \rangle_s - 3 \langle p_2^2 q_2 \rangle_s b_1 - \frac{3 \hbar^2 b_2}{2} - 3 \langle p_2^2 q_2^2 \rangle_s b_2 - 4 \hbar^2 \langle q_2 \rangle b_3 - 3 \langle p_2^2 q_2^3 \rangle_s b_3 + \langle p_2^4 \rangle k_2$
$\frac{d \langle p_2^4 \rangle}{d T} = -4 \lambda \langle p_2^3 \rangle \langle q_1^2 \rangle - 4 \langle p_2^3 \rangle b_1 - 4 \langle p_2^3 q_2 \rangle_s b_2 - 4 \hbar^2 \langle p_2 \rangle b_3 - 4 \langle p_2^3 q_2^2 \rangle_s b_3$

Export can be used to generate a file for L^AT_EX editors, to be included in a paper, see the command below:

```

Export["myh4.tex", QHDForm[ myhierarchy,
  QHDLLabel → None] ]

```

myh4.tex

The content of the exported L^AT_EX file can be seen below:

FilePrint["myh4.tex"]

```
%% AMS-LaTeX Created by Wolfram Mathematica 8.0 : www.wolfram.com
```

```
\documentclass{article}
\usepackage{amsmath, amssymb, graphics, setspace}

\newcommand{\mathsym}[1]{{{}}
\newcommand{\unicode}[1]{{{}}

\newcounter{mathematicapage}
\begin{document}

\[\begin{array}{|l|}
\hline
\frac{d\left\langle p_1\right\rangle}{dt}=-a_2\left\langle q_1\right\rangle-2\lambda
\\
\hline
\frac{d\left\langle q_1\right\rangle}{dt}=k_1\left\langle p_1\right\rangle
\\
\hline
\frac{d\left\langle q_2\right\rangle}{dt}=k_2\left\langle p_2\right\rangle
\\
\hline
\frac{d\left\langle p_2\right\rangle}{dt}=-b_1-\lambda\left\langle q_1^2\right\rangle-2\lambda\left\langle q_2^2\right\rangle
\\
\hline
\frac{d\left\langle q_1^2\right\rangle}{dt}=2k_1\left\langle p_1q_1\right\rangle-2\lambda\left\langle q_1^2\right\rangle-2\lambda\left\langle q_2^2\right\rangle
\\
\hline
\frac{d\left\langle q_2^2\right\rangle}{dt}=2k_2\left\langle p_2q_2\right\rangle-2\lambda\left\langle q_1^2\right\rangle-2\lambda\left\langle q_2^2\right\rangle
\\
\hline
\frac{d\left\langle p_1q_1\right\rangle}{dt}=k_1\left\langle p_1^2\right\rangle-2\lambda\left\langle q_1^2\right\rangle-2\lambda\left\langle q_2^2\right\rangle
\\
\hline
\frac{d\left\langle p_2q_2\right\rangle}{dt}=k_2\left\langle p_2^2\right\rangle-2\lambda\left\langle q_1^2\right\rangle-2\lambda\left\langle q_2^2\right\rangle-2b_2\left\langle q_2^2\right\rangle
\\
\hline
\frac{d\left\langle p_1^2\right\rangle}{dt}=-2a_2\left\langle p_1q_1\right\rangle-2\lambda\left\langle p_1q_1\right\rangle
\\
\hline
\frac{d\left\langle p_2^2\right\rangle}{dt}=-2b_1\left\langle p_2\right\rangle-2\lambda\left\langle q_1^2\right\rangle-2b_2\left\langle p_2q_2\right\rangle-2b_3\left\langle p_2q_2\right\rangle
\\
\hline
\frac{d\left\langle q_2^3\right\rangle}{dt}=3k_2\left\langle p_2q_2^2\right\rangle
\\
\hline
\frac{d\left\langle p_2q_2^2\right\rangle}{dt}=-b_1\left\langle q_2^2\right\rangle-2b_2\left\langle q_2^3\right\rangle-2b_3\left\langle q_2^4\right\rangle
\\
\hline
\frac{d\left\langle q_2^4\right\rangle}{dt}=4k_2\left\langle p_2q_2^3\right\rangle
\\
\hline
\frac{d\left\langle p_2^2q_2\right\rangle}{dt}=k_2\left\langle p_2^3\right\rangle-2\lambda\left\langle q_1^2\right\rangle-2\lambda\left\langle p_2q_2\right\rangle-2b_3\left\langle p_2q_2^3\right\rangle
\\
\hline
\frac{d\left\langle p_2q_2^3\right\rangle}{dt}=\frac{3}{2}k_2\left\langle q_1^2\right\rangle-2b_2\left\langle q_2^4\right\rangle-2b_3\left\langle p_2^2q_2^2\right\rangle
\\
\hline\end{array}
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\frac{d\left\langle p_2^3\right\rangle}{dt}=-\hbar^2 b_3^{-3} b_1 \left\langle p_2^2\right\rangle
\left\langle q_1^2\right\rangle^{-3} b_2 \left\langle p_2^2 q_2\right\rangle_{\mathit{s}}^{-3} b_3
\\
\hline
\frac{d\left\langle p_2^2 q_2^2\right\rangle}{dt}=2 k_2 \left\langle p_2^2\right\rangle
\left\langle q_2^2\right\rangle_{\mathit{s}}^{-2} \lambda \left\langle q_1^2\right\rangle \left\langle p_2^2 q_2^3\right\rangle_{\mathit{s}}^{-2} b_3 \left\langle p_2 q_2^4\right\rangle_{\mathit{s}} \backslash
\hline
\frac{d\left\langle p_2^3 q_2\right\rangle}{dt}=-\frac{3}{2} \hbar^2 b_3
b_1 \left\langle q_2\right\rangle^{-3} b_1 \left\langle p_2^2 q_2\right\rangle_{\mathit{s}}^{-3} \lambda
p_2^2 q_2\right\rangle_{\mathit{s}}^{-3} b_2 \left\langle p_2^2 q_2^2\right\rangle_{\mathit{s}}
\\
\hline
\frac{d\left\langle p_2^4\right\rangle}{dt}=-4 \hbar^2 b_3 \left\langle p_2\right\rangle
\left\langle p_2^3\right\rangle \left\langle q_1^2\right\rangle^{-4} b_2 \left\langle p_2^3 q_2\right\rangle_{\mathit{s}}
\\
\hline
\end{array}
\\
\end{document}

```

Compare equations in this document with equations from the appendix of Pahl and Prezhdo in J. Chem Phys. Vol 116 No. 20, May 2002, Pages 8704-8712

<http://homepage.cem.itesm.mx/lgozquez/quantum/QHDHigherOrders.pdf> .

Calculations are performed using QUANTUM, a free *Mathematica* add-on that can be downloaded from

<http://homepage.cem.itesm.mx/lgozquez/quantum/>

by José Luis Gómez-Muñoz

<http://homepage.cem.itesm.mx/lgozquez/quantum/>

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