
Operator Diagonalization, Partial Trace and Partial Transpose

by José Luis Gómez-Muñoz

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This document is based on suggestions by Anirban Das

Introduction

This is a tutorial on the use of Quantum *Mathematica* add-on to calculate partial traces and partial transposes of operators in Dirac Notation

Load the Package

First load the Quantum`Computing` package. Write:

`Needs["Quantum`Computing`"]`

then press at the same time the keys `SHIFT-ENTER` to evaluate. *Mathematica* will load the package

```
Needs["Quantum`Computing`"]
```

```
Quantum`Computing` Version 2.2.0. (July 2010)
A Mathematica package for Quantum Computing
  in Dirac bra-ket notation and plotting of quantum circuits
by José Luis Gómez-Muñoz
```

```
Execute SetComputingAliases[] in order to use
  the keyboard to enter quantum objects in Dirac's notation
SetComputingAliases[] must be executed again in each new notebook that is created
```

In order to use the keyboard to enter quantum objects write:

`SetComputingAliases[];`

then press at the same time the keys `SHIFT-ENTER` to evaluate. The semicolon prevents *Mathematica* from printing the help message. Remember that `SetComputingAliases[]` must be evaluated again in each new notebook:

```
SetComputingAliases[];
```

Operator Diagonalization

`QuantumEigensystem` gives a list of eigenvalues and eigenvectors of an operator, in the same format as the standard *Mathematica* command `Eigensystem`:

```
QuantumEigensystem[C(1)[NOT2]]
```

```
{{-1, 1, 1, 1}, {- (|11, 02⟩ + |11, 12⟩)/√2, (|11, 02⟩ - |11, 12⟩)/√2, |01, 12⟩, |01, 02⟩}}
```

QuantumEigensystemForm gives a nicer formatting:

```
QuantumEigensystemForm[C(1)[NOT2]]
```

Eigenvalue	Eigenvector
-1	$-\frac{ 1_1, 0_2\rangle}{\sqrt{2}} + \frac{ 1_1, 1_2\rangle}{\sqrt{2}}$
1	$\frac{ 1_1, 0_2\rangle}{\sqrt{2}} + \frac{ 1_1, 1_2\rangle}{\sqrt{2}}$
1	$ 0_1, 1_2\rangle$
1	$ 0_1, 0_2\rangle$

TraditionalForm gives a format closer to the format used in papers:

```
TraditionalForm[QuantumEigensystemForm[C(1)[NOT2]]]
```

Eigenvalue	Eigenvector
-1	$\frac{ 11\rangle}{\sqrt{2}} - \frac{ 10\rangle}{\sqrt{2}}$
1	$\frac{ 10\rangle}{\sqrt{2}} + \frac{ 11\rangle}{\sqrt{2}}$
1	$ 01\rangle$
1	$ 00\rangle$

One simple way to force numerical evaluation is to multiply the operator times 1.0 (with a decimal point!)

```
TraditionalForm[QuantumEigensystemForm[1.0 * C(1)[NOT2]]]
```

Eigenvalue	Eigenvector
1.	1. 01⟩
1.	1. 00⟩
1.	0.707107 10⟩ + 0.707107 11⟩
-1.	0.707107 10⟩ - 0.707107 11⟩

A more interesting example:

$$\text{TraditionalForm}\left[\text{QuantumEigensystemForm}\left[C^{(\hat{1})}\left[NOT_2\right] \cdot \mathcal{H}_1\right]\right]$$

Eigenvalue	Eigenvector
-1	$\left(\frac{1}{2\sqrt{2-\sqrt{2}}} - \frac{1}{\sqrt{2(2-\sqrt{2})}}\right) 00\rangle + \left(\frac{1}{2\sqrt{2-\sqrt{2}}} - \frac{1}{\sqrt{2(2-\sqrt{2})}}\right) 01\rangle + \frac{ 10\rangle}{2\sqrt{2-\sqrt{2}}} + \frac{ 11\rangle}{2\sqrt{2-\sqrt{2}}}$
1	$\left(\frac{1}{2\sqrt{2+\sqrt{2}}} + \frac{1}{\sqrt{2(2+\sqrt{2})}}\right) 00\rangle + \left(\frac{1}{2\sqrt{2+\sqrt{2}}} + \frac{1}{\sqrt{2(2+\sqrt{2})}}\right) 01\rangle + \frac{ 10\rangle}{2\sqrt{2+\sqrt{2}}} + \frac{ 11\rangle}{2\sqrt{2+\sqrt{2}}}$
$\frac{1+i}{\sqrt{2}}$	$\frac{1}{2}i 00\rangle - \frac{1}{2}i 01\rangle - \frac{1}{2} 10\rangle + \frac{1}{2} 11\rangle$
$\frac{1-i}{\sqrt{2}}$	$-\frac{1}{2}i 00\rangle + \frac{1}{2}i 01\rangle - \frac{1}{2} 10\rangle + \frac{1}{2} 11\rangle$

One simple way to force numerical evaluation is to multiply the operator times 1.0 (with a decimal point!)

$$\text{TraditionalForm}\left[\text{QuantumEigensystemForm}\left[1.0 * C^{(\hat{1})}\left[NOT_2\right] \cdot \mathcal{H}_1\right]\right]$$

Eigenvalue	Eigenvector
-1.	$0.270598 00\rangle + 0.270598 01\rangle - 0.653281 10\rangle - 0.653281 11\rangle$
$0.707107 + 0.707107i$	$-0.5i 00\rangle + 0.5i 01\rangle + 0.5 10\rangle - 0.5 11\rangle$
$0.707107 - 0.707107i$	$0.5i 00\rangle - 0.5i 01\rangle + 0.5 10\rangle - 0.5 11\rangle$
1.	$-0.653281 00\rangle - 0.653281 01\rangle - 0.270598 10\rangle - 0.270598 11\rangle$

Here we obtain only the first eigenvalue, together with its eigenvector:

$$\text{QuantumEigensystem}\left[1.0 * C^{(\hat{1})}\left[NOT_2\right] \cdot \mathcal{H}_1, 1\right]$$

$$\{\{-1.\}, \{0.270598|0_1, 0_2\rangle + 0.270598|0_1, 1_2\rangle - 0.653281|1_1, 0_2\rangle - 0.653281|1_1, 1_2\rangle\}\}$$

Here we obtain only the last eigenvalue, together with its eigenvector:

$$\text{QuantumEigensystem}\left[1.0 * C^{(\hat{1})}\left[NOT_2\right] \cdot \mathcal{H}_1, -1\right]$$

$$\{\{1.\}, \{-0.653281|0_1, 0_2\rangle - 0.653281|0_1, 1_2\rangle - 0.270598|1_1, 0_2\rangle - 0.270598|1_1, 1_2\rangle\}\}$$

Here we extract the eigenvector corresponding to the last eigenvalue using standard *Mathematica* notation to access elements of lists:

$$\text{QuantumEigensystem}\left[1.0 * C^{(\hat{1})}\left[NOT_2\right] \cdot \mathcal{H}_1, -1\right][[2, 1]]$$

$$-0.653281|0_1, 0_2\rangle - 0.653281|0_1, 1_2\rangle - 0.270598|1_1, 0_2\rangle - 0.270598|1_1, 1_2\rangle$$

You can use the operator in Dirac bra-ket notation:

```
TraditionalForm[QuantumEigensystemForm[ $\frac{1}{2} \begin{vmatrix} |0_1, 1_2\rangle \cdot \langle 0_1, 0_2| + \\ \frac{1}{2} \begin{vmatrix} |1_1, 1_2\rangle \cdot \langle 0_1, 1_2| - \frac{1}{2} \begin{vmatrix} |0_1, 0_2\rangle \cdot \langle 1_1, 0_2| - \frac{1}{2} \begin{vmatrix} |1_1, 0_2\rangle \cdot \langle 1_1, 1_2| \end{vmatrix} \end{vmatrix} \end{vmatrix}$ 
```

Eigenvalue	Eigenvector
-1	$-\frac{1}{2} 00\rangle + \frac{1}{2} i 01\rangle + \frac{1}{2} i 10\rangle + \frac{1}{2} 11\rangle$
i	$\frac{1}{2} 00\rangle + \frac{1}{2} 01\rangle - \frac{1}{2} 10\rangle + \frac{1}{2} 11\rangle$
$-i$	$\frac{1}{2} 00\rangle - \frac{1}{2} 01\rangle + \frac{1}{2} 10\rangle + \frac{1}{2} 11\rangle$
1	$-\frac{1}{2} 00\rangle - \frac{1}{2} i 01\rangle - \frac{1}{2} i 10\rangle + \frac{1}{2} 11\rangle$

Partial Traces

Here we define a ket as a linear combination of two qubits:

$$|\psi\rangle = \alpha |0_1, 0_2\rangle + \beta |0_1, 1_2\rangle + \gamma |1_1, 0_2\rangle + \delta |1_1, 1_2\rangle$$

$$\alpha |0_1, 0_2\rangle + \beta |0_1, 1_2\rangle + \gamma |1_1, 0_2\rangle + \delta |1_1, 1_2\rangle$$

An operator can be obtained from the external product of the ket with its corresponding dual:

```
mydensityop = Expand[  $|\psi\rangle \cdot \langle\psi|$  ]
```

$$\begin{aligned} & \alpha \alpha^* |0_1, 0_2\rangle \cdot \langle 0_1, 0_2| + \beta \alpha^* |0_1, 1_2\rangle \cdot \langle 0_1, 0_2| + \\ & \gamma \alpha^* |1_1, 0_2\rangle \cdot \langle 0_1, 0_2| + \delta \alpha^* |1_1, 1_2\rangle \cdot \langle 0_1, 0_2| + \\ & \alpha \beta^* |0_1, 0_2\rangle \cdot \langle 0_1, 1_2| + \beta \beta^* |0_1, 1_2\rangle \cdot \langle 0_1, 1_2| + \gamma \beta^* |1_1, 0_2\rangle \cdot \langle 0_1, 1_2| + \\ & \delta \beta^* |1_1, 1_2\rangle \cdot \langle 0_1, 1_2| + \alpha \gamma^* |0_1, 0_2\rangle \cdot \langle 1_1, 0_2| + \beta \gamma^* |0_1, 1_2\rangle \cdot \langle 1_1, 0_2| + \\ & \gamma \gamma^* |1_1, 0_2\rangle \cdot \langle 1_1, 0_2| + \delta \gamma^* |1_1, 1_2\rangle \cdot \langle 1_1, 0_2| + \alpha \delta^* |0_1, 0_2\rangle \cdot \langle 1_1, 1_2| + \\ & \beta \delta^* |0_1, 1_2\rangle \cdot \langle 1_1, 1_2| + \gamma \delta^* |1_1, 0_2\rangle \cdot \langle 1_1, 1_2| + \delta \delta^* |1_1, 1_2\rangle \cdot \langle 1_1, 1_2| \end{aligned}$$

TraditionalForm gives a format closer to the format used in papers and textbooks:

```
TraditionalForm[mydensityop]
```

$$\begin{aligned} & \beta \alpha^* |01\rangle\langle 00| + \alpha \beta^* |00\rangle\langle 01| + \gamma \alpha^* |10\rangle\langle 00| + \alpha \gamma^* |00\rangle\langle 10| + \delta \alpha^* |11\rangle\langle 00| + \\ & \alpha \delta^* |00\rangle\langle 11| + \alpha \alpha^* |00\rangle\langle 00| + \gamma \beta^* |10\rangle\langle 01| + \beta \gamma^* |01\rangle\langle 10| + \delta \beta^* |11\rangle\langle 01| + \\ & \beta \delta^* |01\rangle\langle 11| + \beta \beta^* |01\rangle\langle 01| + \delta \gamma^* |11\rangle\langle 10| + \gamma \delta^* |10\rangle\langle 11| + \gamma \gamma^* |10\rangle\langle 10| + \delta \delta^* |11\rangle\langle 11| \end{aligned}$$

Here is the partial trace of the operator, with respect to qubit $\hat{1}$

```
mypartialtrace1 = QuantumPartialTrace[mydensityop,  $\hat{1}$ ]
```

$$\begin{aligned} & \alpha \alpha^* |0_2\rangle \cdot \langle 0_2| + \gamma \gamma^* |0_2\rangle \cdot \langle 0_2| + \beta \alpha^* |1_2\rangle \cdot \langle 0_2| + \delta \gamma^* |1_2\rangle \cdot \langle 0_2| + \\ & \alpha \beta^* |0_2\rangle \cdot \langle 1_2| + \gamma \delta^* |0_2\rangle \cdot \langle 1_2| + \beta \beta^* |1_2\rangle \cdot \langle 1_2| + \delta \delta^* |1_2\rangle \cdot \langle 1_2| \end{aligned}$$

Here is the partial trace of the operator, with respect to qubit $\hat{2}$

```
mypartialtrace2 = QuantumPartialTrace[mydensityop, 2]
```

$$\begin{aligned} & \alpha \alpha^* |0_1\rangle \langle 0_1| + \beta \beta^* |0_1\rangle \langle 0_1| + \gamma \alpha^* |1_1\rangle \langle 0_1| + \delta \beta^* |1_1\rangle \langle 0_1| + \\ & \alpha \gamma^* |0_1\rangle \langle 1_1| + \beta \delta^* |0_1\rangle \langle 1_1| + \gamma \gamma^* |1_1\rangle \langle 1_1| + \delta \delta^* |1_1\rangle \langle 1_1| \end{aligned}$$

Here is the trace of the operator. It is stored in the variable **mytrace**, in order to use it later in this document:

```
mytrace = QuantumPartialTrace[mydensityop, {1, 2}]
```

$$\alpha \alpha^* + \beta \beta^* + \gamma \gamma^* + \delta \delta^*$$

Partial Transposes

Here is the partial transpose of the operator defined above, with respect to qubit $\hat{1}$. Partial Transpose is used as a witness of entanglement in Quantum Computing.

```
mypartialtranspose1 = QuantumPartialTranspose[mydensityop, 1]
```

$$\begin{aligned} & \alpha \alpha^* |0_1, 0_2\rangle \langle 0_1, 0_2| + \beta \alpha^* |0_1, 1_2\rangle \langle 0_1, 0_2| + \\ & \alpha \gamma^* |1_1, 0_2\rangle \langle 0_1, 0_2| + \beta \gamma^* |1_1, 1_2\rangle \langle 0_1, 0_2| + \\ & \alpha \beta^* |0_1, 0_2\rangle \langle 0_1, 1_2| + \beta \beta^* |0_1, 1_2\rangle \langle 0_1, 1_2| + \alpha \delta^* |1_1, 0_2\rangle \langle 0_1, 1_2| + \\ & \beta \delta^* |1_1, 1_2\rangle \langle 0_1, 1_2| + \gamma \alpha^* |0_1, 0_2\rangle \langle 1_1, 0_2| + \delta \alpha^* |0_1, 1_2\rangle \langle 1_1, 0_2| + \\ & \gamma \gamma^* |1_1, 0_2\rangle \langle 1_1, 0_2| + \delta \gamma^* |1_1, 1_2\rangle \langle 1_1, 0_2| + \gamma \beta^* |0_1, 0_2\rangle \langle 1_1, 1_2| + \\ & \delta \beta^* |0_1, 1_2\rangle \langle 1_1, 1_2| + \gamma \delta^* |1_1, 0_2\rangle \langle 1_1, 1_2| + \delta \delta^* |1_1, 1_2\rangle \langle 1_1, 1_2| \end{aligned}$$

Here is the partial transpose of the operator, with respect to the qubit $\hat{2}$. Partial Transpose is used as a witness of entanglement in Quantum Computing.

```
mypartialtranspose2 = QuantumPartialTranspose[mydensityop, 2]
```

$$\begin{aligned} & \alpha \alpha^* |0_1, 0_2\rangle \langle 0_1, 0_2| + \alpha \beta^* |0_1, 1_2\rangle \langle 0_1, 0_2| + \\ & \gamma \alpha^* |1_1, 0_2\rangle \langle 0_1, 0_2| + \gamma \beta^* |1_1, 1_2\rangle \langle 0_1, 0_2| + \\ & \beta \alpha^* |0_1, 0_2\rangle \langle 0_1, 1_2| + \beta \beta^* |0_1, 1_2\rangle \langle 0_1, 1_2| + \delta \alpha^* |1_1, 0_2\rangle \langle 0_1, 1_2| + \\ & \delta \beta^* |1_1, 1_2\rangle \langle 0_1, 1_2| + \alpha \gamma^* |0_1, 0_2\rangle \langle 1_1, 0_2| + \alpha \delta^* |0_1, 1_2\rangle \langle 1_1, 0_2| + \\ & \gamma \gamma^* |1_1, 0_2\rangle \langle 1_1, 0_2| + \gamma \delta^* |1_1, 1_2\rangle \langle 1_1, 0_2| + \beta \gamma^* |0_1, 0_2\rangle \langle 1_1, 1_2| + \\ & \beta \delta^* |0_1, 1_2\rangle \langle 1_1, 1_2| + \delta \gamma^* |1_1, 0_2\rangle \langle 1_1, 1_2| + \delta \delta^* |1_1, 1_2\rangle \langle 1_1, 1_2| \end{aligned}$$

Here is the transpose of the operator. Notice that it is stored in the variable "mytranspose", so that we can use it later in this document:

```
mytranspose = QuantumPartialTranspose[mydensityop, {1, 2}]
```

$$\begin{aligned} & \alpha \alpha^* | 0_1, 0_2 \rangle \cdot \langle 0_1, 0_2 | + \alpha \beta^* | 0_1, 1_2 \rangle \cdot \langle 0_1, 0_2 | + \\ & \alpha \gamma^* | 1_1, 0_2 \rangle \cdot \langle 0_1, 0_2 | + \alpha \delta^* | 1_1, 1_2 \rangle \cdot \langle 0_1, 0_2 | + \\ & \beta \alpha^* | 0_1, 0_2 \rangle \cdot \langle 0_1, 1_2 | + \beta \beta^* | 0_1, 1_2 \rangle \cdot \langle 0_1, 1_2 | + \beta \gamma^* | 1_1, 0_2 \rangle \cdot \langle 0_1, 1_2 | + \\ & \beta \delta^* | 1_1, 1_2 \rangle \cdot \langle 0_1, 1_2 | + \gamma \alpha^* | 0_1, 0_2 \rangle \cdot \langle 1_1, 0_2 | + \gamma \beta^* | 0_1, 1_2 \rangle \cdot \langle 1_1, 0_2 | + \\ & \gamma \gamma^* | 1_1, 0_2 \rangle \cdot \langle 1_1, 0_2 | + \gamma \delta^* | 1_1, 1_2 \rangle \cdot \langle 1_1, 0_2 | + \delta \alpha^* | 0_1, 0_2 \rangle \cdot \langle 1_1, 1_2 | + \\ & \delta \beta^* | 0_1, 1_2 \rangle \cdot \langle 1_1, 1_2 | + \delta \gamma^* | 1_1, 0_2 \rangle \cdot \langle 1_1, 1_2 | + \delta \delta^* | 1_1, 1_2 \rangle \cdot \langle 1_1, 1_2 | \end{aligned}$$

Verifying traces and transposes with Matrix Notation

We can transform the operator, which was defined above, from Dirac Notation to standard *Mathematica* Matrix Notation (list of lists). Therefore, we can take advantage of the built-in *Mathematica* commands for matrices:

```
mymatrix = QuantumMatrix[mydensityop]
```

```
{ {α α*, α β*, α γ*, α δ*}, {β α*, β β*, β γ*, β δ*},  
  {γ α*, γ β*, γ γ*, γ δ*}, {δ α*, δ β*, δ γ*, δ δ*} }
```

Here we can visualizar the Matrix in a textbook-like format:

```
MatrixForm[mymatrix]
```

$$\begin{pmatrix} \alpha \alpha^* & \alpha \beta^* & \alpha \gamma^* & \alpha \delta^* \\ \beta \alpha^* & \beta \beta^* & \beta \gamma^* & \beta \delta^* \\ \gamma \alpha^* & \gamma \beta^* & \gamma \gamma^* & \gamma \delta^* \\ \delta \alpha^* & \delta \beta^* & \delta \gamma^* & \delta \delta^* \end{pmatrix}$$

This is the trace of the matrix. We store the result in the variable matrixtrace

```
matrixtrace = Tr[mymatrix]
```

```
α α* + β β* + γ γ* + δ δ*
```

The trace "matrixtrace" was obtained by transforming the original Dirac expression to a Matrix, then calculating the trace of the matrix. On the other hand, the trace "mytrace" was obtained by direct application of the command QuantumPartialTrace to the original Dirac expression.

As expected, both "matrixtrace" and "mytrace" are the same (notice the use of two equal symbols == in order to make the comparison)

```
matrixtrace == mytrace
```

```
True
```

This is the transpose of the matrix:

```
matrixtranspose = Transpose[mymatrix]
```

```
{ {α α*, β α*, γ α*, δ α*}, {α β*, β β*, γ β*, δ β*},  
  {α γ*, β γ*, γ γ*, δ γ*}, {α δ*, β δ*, γ δ*, δ δ*} }
```

The transpose is transformed to Dirac notation

```
secondtranspose = MatrixQuantum[matrixtranspose]
```

```
α α* | 01, 02⟩ · ⟨ 01, 02 | + α β* | 01, 12⟩ · ⟨ 01, 02 | +  
α γ* | 11, 02⟩ · ⟨ 01, 02 | + α δ* | 11, 12⟩ · ⟨ 01, 02 | +  
β α* | 01, 02⟩ · ⟨ 01, 12 | + β β* | 01, 12⟩ · ⟨ 01, 12 | + β γ* | 11, 02⟩ · ⟨ 01, 12 | +  
β δ* | 11, 12⟩ · ⟨ 01, 12 | + γ α* | 01, 02⟩ · ⟨ 11, 02 | + γ β* | 01, 12⟩ · ⟨ 11, 02 | +  
γ γ* | 11, 02⟩ · ⟨ 11, 02 | + γ δ* | 11, 12⟩ · ⟨ 11, 02 | + δ α* | 01, 02⟩ · ⟨ 11, 12 | +  
δ β* | 01, 12⟩ · ⟨ 11, 12 | + δ γ* | 11, 02⟩ · ⟨ 11, 12 | + δ δ* | 11, 12⟩ · ⟨ 11, 12 |
```

The operator "secondtranspose" was obtained by transforming the original Dirac expression to a Matrix, then calculating the transpose of the matrix, and finally transforming back to Dirac notation.

On the other hand, the operator "mytranspose" was obtained by direct application of the command QuantumPartialTranspose to the original Dirac expression.

As expected, both "secondtranspose" and "mytranspose" are the same:

```
secondtranspose == mytranspose
```

```
True
```

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