

1. The stress power is defined to be  $S_p = \int_R \text{tr}(\boldsymbol{\sigma} \cdot \mathbf{d}) dV$  where  $\boldsymbol{\sigma}$  is the Cauchy stress and  $\mathbf{d} = \mathbf{L}_{sym}$ .

Other measures of stress are

$$\begin{aligned}\boldsymbol{\Sigma} &= \mathbf{R}^T \cdot \boldsymbol{\sigma} \cdot \mathbf{R} && \text{Rotated Cauchy stress} \\ \hat{\mathbf{P}} &= J \boldsymbol{\sigma} \cdot \mathbf{F}^{-T} && \text{Piola-Kirchoff stress of the first kind} \\ \mathbf{P} &= J \mathbf{F}^{-l} \cdot \boldsymbol{\sigma} \cdot \mathbf{F}^{-T} = \mathbf{F}^{-l} \cdot \hat{\mathbf{P}} && \text{Piola-Kirchoff stress of the second kind}\end{aligned}$$

and other rates of deformation are

$$\mathbf{D} = \mathbf{F}^T \cdot \mathbf{d} \cdot \mathbf{F} = \dot{\mathbf{E}} \quad \mathbf{D}^* = \mathbf{R}^T \cdot \mathbf{d} \cdot \mathbf{R}$$

Show that alternative expressions for the stress power are:

$$S_p = \int_R \text{tr}(\boldsymbol{\Sigma} \cdot \mathbf{D}^*) dV = \int_{R_o} \text{tr}(\hat{\mathbf{P}} \cdot \dot{\mathbf{F}}^T) dV_o = \int_{R_o} \text{tr}(\mathbf{P} \cdot \dot{\mathbf{E}}) dV_o$$

These combinations of stress and deformation rates are said to be "conjugate."

**Soln:**

We use  $\text{tr}(\mathbf{A} \cdot \mathbf{B} \cdot \mathbf{C}) = \text{tr}(\mathbf{B} \cdot \mathbf{C} \cdot \mathbf{A})$ ,  $dV = J dV_o$  and  $\text{tr}(\mathbf{A} \cdot \mathbf{B}) = \text{tr}(\mathbf{A} \cdot \mathbf{B}^T)$  if  $\mathbf{A} = \mathbf{A}^T$ .

$$\begin{aligned}\text{tr}(\boldsymbol{\sigma} \cdot \mathbf{d}) &= \text{tr}(\mathbf{R} \cdot \boldsymbol{\Sigma} \cdot \mathbf{R}^T \cdot \mathbf{R} \cdot \mathbf{D}^* \cdot \mathbf{R}^T) = \text{tr}(\mathbf{R} \cdot \boldsymbol{\Sigma} \cdot \mathbf{D}^* \cdot \mathbf{R}^T) \\ &= \text{tr}(\boldsymbol{\Sigma} \cdot \mathbf{D}^* \cdot \mathbf{R}^T \cdot \mathbf{R}) = \text{tr}(\boldsymbol{\Sigma} \cdot \mathbf{D}^*) \\ \text{tr}(\boldsymbol{\sigma} \cdot \mathbf{d}) &= \text{tr}(\boldsymbol{\sigma} \cdot \mathbf{L}) = \text{tr}(\boldsymbol{\sigma} \cdot \mathbf{L}^T) = \frac{1}{J} \text{tr}(\hat{\mathbf{P}} \cdot \mathbf{F}^T \cdot \mathbf{L}^T) = \frac{1}{J} \text{tr}(\hat{\mathbf{P}} \cdot \dot{\mathbf{F}}^T) \\ \text{tr}(\boldsymbol{\sigma} \cdot \mathbf{d}) &= \frac{1}{J} \text{tr}(\mathbf{F} \cdot \mathbf{P} \cdot \mathbf{F}^T \cdot \mathbf{F}^{-T} \cdot \mathbf{D} \cdot \mathbf{F}^{-l}) = \frac{1}{J} \text{tr}(\mathbf{F} \cdot \mathbf{P} \cdot \mathbf{D} \cdot \mathbf{F}^{-l}) \\ &= \frac{1}{J} \text{tr}(\mathbf{P} \cdot \mathbf{D} \cdot \mathbf{F}^{-l} \cdot \mathbf{F}) = \frac{1}{J} \text{tr}(\mathbf{P} \cdot \mathbf{D})\end{aligned}$$

2. Recall that in connection with the study of a continuum, tensors could be defined as one of four possibilities: m-m, s-s, s-m, m-s where "m" denotes "material" and "s" denotes "spatial". Recall the classifications of  $\mathbf{F}$  and  $\mathbf{R}$  from the notes. Assume  $\boldsymbol{\sigma}$  and  $\mathbf{d}$  are both s-s. Use the relations given in Prob. 1 to classify the tensors  $\boldsymbol{\Sigma}$ ,  $\hat{\mathbf{P}}$ ,  $\mathbf{P}$ ,  $\mathbf{D}$  and  $\mathbf{D}^*$ .

**Soln:**

Start with

$$\boldsymbol{\sigma} = \overset{s-s}{\boldsymbol{\sigma}}, \mathbf{d} = \overset{s-s}{\mathbf{d}}, \mathbf{F} = \overset{s-m}{\mathbf{F}}, \mathbf{F}^T = \overset{m-s}{\mathbf{F}^T}, \mathbf{F}^{-l} = \overset{m-s}{\mathbf{F}^{-l}}, \mathbf{F}^{-T} = \overset{s-m}{\mathbf{F}^{-l}}, \mathbf{R} = \overset{s-m}{\mathbf{R}} \text{ and } \mathbf{R}^T = \overset{m-s}{\mathbf{R}^T}$$

We look for the two outermost superscripts to identify the following tensors:

$$\begin{aligned}
\boldsymbol{\Sigma} &= \overset{m-s}{\mathbf{R}^T} \cdot \overset{s-s}{\boldsymbol{\sigma}} \cdot \overset{s-m}{\mathbf{R}} \quad \Rightarrow \quad \boldsymbol{\Sigma} = \overset{m-m}{\boldsymbol{\Sigma}} \quad \text{material-material} \\
\hat{\mathbf{P}} &= J \overset{s-s}{\boldsymbol{\sigma}} \cdot \overset{s-m}{\mathbf{F}^{-T}} \quad \Rightarrow \quad \hat{\mathbf{P}} = \overset{s-m}{\hat{\mathbf{P}}} \quad \text{spatial-material} \\
\mathbf{P} &= J \overset{m-s}{\mathbf{F}^{-1}} \cdot \overset{s-s}{\boldsymbol{\sigma}} \cdot \overset{s-m}{\mathbf{F}^{-T}} \quad \Rightarrow \quad \mathbf{P} = \overset{m-m}{\mathbf{P}} \quad \text{material-material} \\
\mathbf{D} &= \overset{m-s}{\mathbf{F}^T} \cdot \overset{s-s}{\mathbf{d}} \cdot \overset{s-m}{\mathbf{F}} \quad \Rightarrow \quad \mathbf{D} = \overset{m-m}{\mathbf{D}} \quad \text{material-material} \\
\mathbf{D}^* &= \overset{m-s}{\mathbf{R}^T} \cdot \overset{s-s}{\mathbf{d}} \cdot \overset{s-m}{\mathbf{R}} \quad \Rightarrow \quad \mathbf{D}^* = \overset{m-m}{\mathbf{D}^*} \quad \text{material-material}
\end{aligned}$$

3. A bar of original length  $L$  that is initially horizontal (Fig. 1) deforms in a plane as the result of a simultaneous stretch and rotation as indicated in Fig. 2. The end  $O$  is fixed in space. The rotation is defined by  $\theta = \omega t$  and the elongation of the end of the bar is  $\delta_A = \varepsilon L t$  with both  $\omega$  and  $\varepsilon$  constant.

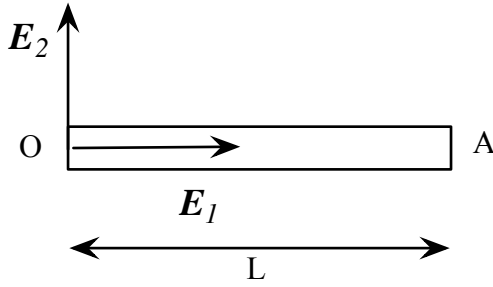


Fig. 1. Initial Position

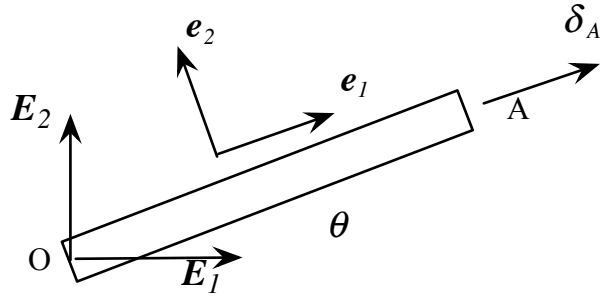


Fig. 2. Deformed position

The easiest way to describe the deformation is to use the two bases. Then the deformation is defined by

$$\begin{aligned}
\mathbf{r} &= x_i \mathbf{e}_i \quad \text{and} \quad \mathbf{R} = X_i \mathbf{E}_i \\
x_1 &= X_1(1 + \varepsilon t) \quad x_2 = X_2 \quad x_3 = X_3
\end{aligned} \tag{0-1}$$

The elongation at the end of the bar is  $\delta_A = (x_1 - X_1)|_{X_1=L} = (X_1 \varepsilon t)|_{X_1=L} = \varepsilon L t$ .

3.1 Perform the transformation so that the components of  $\mathbf{r}$  are given with respect to the basis  $\mathbf{E}_i$ . Let  $\mathbf{r} = x_i^E \mathbf{E}_i$ . Use this form to obtain  $\mathbf{F}, \mathbf{R}, \mathbf{U}, \dot{\mathbf{F}}, \dot{\mathbf{R}}, \dot{\mathbf{U}}$  and  $\boldsymbol{\Omega}$  all with the tensor basis  $\mathbf{E}_i \otimes \mathbf{E}_j$ .

**Soln:**

The deformed and undeformed position vectors are

$$\mathbf{r} = x_i^e \mathbf{e}_i \quad \text{and} \quad \mathbf{R} = X_i \mathbf{E}_i$$

For simplicity, we drop the superscript notation and let  $x_i \equiv x_i^e$ . The deformation is

$$\begin{aligned} x_1 &= X_1(1 + \epsilon t) & x_2 &= X_2 & x_3 &= X_3 \\ \mathbf{r} &= X_1(1 + \epsilon t)\mathbf{e}_1 + X_2\mathbf{e}_2 + X_3\mathbf{e}_3 \end{aligned}$$

and the two bases are related by

$$\begin{aligned} \{\mathbf{e}\} &= \overset{e-E}{[a]}\{\mathbf{E}\} & c &= \cos\theta & s &= \sin\theta \\ \overset{e-E}{[a]} &= \begin{bmatrix} c & s & 0 \\ -s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} \end{aligned}$$

To use a fixed basis for both position vectors, we transform  $\mathbf{e}_i$  to  $\mathbf{E}_i$ . Then

$$\begin{aligned} \mathbf{r} &= \langle x \rangle \{\mathbf{e}\} = \langle x \rangle \overset{e-E}{[a]}\{\mathbf{E}\} = \langle x^E \rangle \{\mathbf{E}\} \\ \langle x^E \rangle &= \langle x \rangle \overset{e-E}{[a]} \\ x_1^E &= cX_1(1 + \epsilon t) - sX_2 & x_2^E &= sX_1(1 + \epsilon t) + cX_2 & x_3^E &= X_3 \\ \mathbf{r} &= x_1^E\mathbf{E}_1 + x_2^E\mathbf{E}_2 + x_3^E\mathbf{E}_3 \end{aligned}$$

Now we have expressions for the deformed position vector where the components with respect to  $\mathbf{E}_i$  are explicitly denoted with a superscript “E” as  $x_i^E$ .

**Error Statement:** In the original problem statement I had stated the second component was

$$x_2^E = X_1(1 + \epsilon t)\sin\theta + X_2\sin\theta$$

and the above shows that it should have been

$$x_2^E = X_1(1 + \epsilon t)\sin\theta + X_2\cos\theta$$

**End of Error Statement**

Now we obtain  $\mathbf{F}, \mathbf{R}, \mathbf{U}, \dot{\mathbf{F}}, \dot{\mathbf{R}}, \dot{\mathbf{U}}$  and  $\boldsymbol{\Omega}$ .

Start with

$$\begin{aligned} x_1^E &= cX_1(1 + \epsilon t) - sX_2 & x_2^E &= sX_1(1 + \epsilon t) + cX_2 & x_3^E &= X_3 \\ \mathbf{F} &= \overset{E-E}{F}_{ij}\mathbf{E}_i \otimes \mathbf{E}_j & \overset{E-E}{F}_{ij} &= \partial x_i^E / \partial X_j \end{aligned}$$

Then

$${}^{E-E}[\mathbf{F}] = \begin{bmatrix} (1+\varepsilon t)c & -s & 0 \\ (1+\varepsilon t)s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad {}^{E-E}[\mathbf{F}^T] = \begin{bmatrix} (1+\varepsilon t)c & (1+\varepsilon t)s & 0 \\ -s & c & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

and  $\mathbf{U}^2 = \mathbf{F}^T \cdot \mathbf{F}$  or

$${}^{E-E}[\mathbf{U}^2] = \begin{bmatrix} (1+\varepsilon t)c & (1+\varepsilon t)s & 0 \\ -s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} (1+\varepsilon t)c & -s & 0 \\ (1+\varepsilon t)s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} (1+\varepsilon t)^2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

We have the principal basis  $\mathbf{N}_i = \mathbf{E}_i$  so that

$${}^{E-E}[\mathbf{U}] = \begin{bmatrix} (1+\varepsilon t) & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad {}^{E-E}[\mathbf{U}^{-1}] = \begin{bmatrix} \frac{1}{(1+\varepsilon t)} & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

It follows that

$$\mathbf{R} = \mathbf{F} \cdot \mathbf{U}^{-1} \quad {}^{E-E}[\mathbf{R}] = {}^{E-E}[\mathbf{F}] {}^{E-E}[\mathbf{U}^{-1}]$$

and

$${}^{E-E}[\mathbf{R}] = \begin{bmatrix} (1+\varepsilon t)c & -s & 0 \\ (1+\varepsilon t)s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \frac{1}{(1+\varepsilon t)} & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} c & -s & 0 \\ s & c & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Because  $\mathbf{E}_i$  are time independent, it follows that

$$\dot{\mathbf{F}} = \dot{F}_{ij} \mathbf{E}_i \otimes \mathbf{E}_j \quad \dot{\mathbf{U}} = \dot{U}_{ij} \mathbf{E}_i \otimes \mathbf{E}_j \quad \dot{\mathbf{R}} = \dot{R}_{ij} \mathbf{E}_i \otimes \mathbf{E}_j$$

Note that  $\dot{c} = -\omega s$  and  $\dot{s} = \omega c$ . Then it follows that

$${}^{E-E}[\dot{\mathbf{F}}] = \begin{bmatrix} \varepsilon c - \omega(1+\varepsilon t)s & -\omega c & 0 \\ \varepsilon s + \omega(1+\varepsilon t)c & -\omega s & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$${}^{E-E}[\dot{\mathbf{U}}] = \begin{bmatrix} \varepsilon & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad {}^{E-E}[\dot{\mathbf{R}}] = \omega \begin{bmatrix} -s & -c & 0 \\ c & -s & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Recall that  $\dot{\boldsymbol{\Omega}} = \dot{\mathbf{R}} \cdot \mathbf{R}^T$  so that

$${}^{E-E}[\Omega] = {}^{E-E}[\dot{R}] {}^{E-E}[R^T] = \omega \begin{bmatrix} -s & -c & 0 \\ c & -s & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} c & s & 0 \\ -s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} = \omega \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

**End of 3.1**

3.2 Now use the two bases and the deformation as given by (0-1).

(i) Determine  $\mathbf{F}$ ,  $\mathbf{R}$  and  $\mathbf{U}$

**Soln:** (i)

Start with

$$c = \cos \theta \quad s = \sin \theta \quad \dot{\theta} = \omega \quad \dot{c} = -\omega s \quad \dot{s} = \omega c$$

$$x_1 = X_1(1 + \varepsilon t) \quad x_2 = X_2 \quad x_3 = X_3$$

and

$$\mathbf{F} = {}^{e-E}F_{ij} \mathbf{e}_i \otimes \mathbf{E}_j \quad {}^{e-E}F_{ij} = \partial x_i / \partial X_j$$

Then

$${}^{e-E}[F] = \begin{bmatrix} (1 + \varepsilon t) & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

and

$${}^{E-E}[U^2] = {}^{E-E}[F^T] {}^{e-E}[F^T] = \begin{bmatrix} (1 + \varepsilon t)^2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad {}^{E-E}[U] = \begin{bmatrix} (1 + \varepsilon t) & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

with  $N_i = \mathbf{E}_i$ . Now  $\mathbf{R} = \mathbf{F} \cdot \mathbf{U}^{-1}$  so that

$$\mathbf{R} = {}^{e-E}R_{ij} \mathbf{e}_i \otimes \mathbf{E}_j \quad {}^{e-E}[R] = {}^{e-E}[F] {}^{E-E}[U^{-1}] = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

or  $\mathbf{R} = \mathbf{e}_i \otimes \mathbf{E}_i$  and  $\mathbf{R}^T = \mathbf{E}_i \otimes \mathbf{e}_i$ .

**End of (i).**

(ii) The basis  $\mathbf{e}_i$  is a function of time. Determine  $\dot{\mathbf{e}}_i$  and determine the tensor  $\boldsymbol{\Omega}^*$  such that  $\dot{\mathbf{e}}_i = \boldsymbol{\Omega}^* \cdot \mathbf{e}_i$ .

**Soln:** (ii)

Recall that

$$\{\mathbf{e}\} = \overset{e-E}{[a]} \{\mathbf{E}\} \quad \overset{e-E}{[a]} = \begin{bmatrix} c & s & 0 \\ -s & c & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Then

$$\{\dot{\mathbf{e}}\} = \overset{e-E}{[\dot{a}]} \{\mathbf{E}\} \quad \overset{e-E}{[\dot{a}]} = \omega \begin{bmatrix} -s & c & 0 \\ -c & -s & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Now transform to the basis  $\mathbf{e}_i$ :

$$\begin{aligned} \{\dot{\mathbf{e}}\} &= \overset{e-E}{[\dot{a}]} \overset{E-e}{[a]} \{\mathbf{e}\} & \overset{e-E}{[\dot{a}]} &= \omega \begin{bmatrix} -s & c & 0 \\ -c & -s & 0 \\ 0 & 0 & 0 \end{bmatrix} & \overset{E-e}{[a]} &= \begin{bmatrix} c & -s & 0 \\ s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} \\ \overset{e-E}{[\dot{a}]} \overset{E-e}{[a]} &= \omega \begin{bmatrix} -s & c & 0 \\ -c & -s & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} c & -s & 0 \\ s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} = \omega \begin{bmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \end{aligned}$$

It follows that

$$\dot{\mathbf{e}}_1 = \omega \mathbf{e}_2 \quad \dot{\mathbf{e}}_2 = -\omega \mathbf{e}_1$$

Let  $\boldsymbol{\Omega}^* = \overset{e-e}{\Omega}_{ij}^* \mathbf{e}_i \otimes \mathbf{e}_j$ . Then  $\boldsymbol{\Omega}^* \cdot \mathbf{e}_k = \overset{e-e}{\Omega}_{ij}^* \mathbf{e}_i \delta_{jk} = \overset{e-e}{\Omega}_{ik}^* \mathbf{e}_i$  and

$$\begin{aligned} \dot{\mathbf{e}}_i &= \boldsymbol{\Omega}^* \cdot \mathbf{e}_i \Rightarrow \{\dot{\mathbf{e}}\} = \{\mathbf{e}\} \cdot \boldsymbol{\Omega}^* = \{\mathbf{e}\} \cdot \langle \mathbf{e} \rangle [\overset{e-e}{\Omega}^*] \{\mathbf{e}\} = [\overset{e-e}{\Omega}^*] \{\mathbf{e}\} \\ \dot{\mathbf{e}}_1 &= \Omega_{11}^* \mathbf{e}_1 + \Omega_{12}^* \mathbf{e}_2 + \Omega_{13}^* \mathbf{e}_3 = \omega \mathbf{e}_2 \\ \dot{\mathbf{e}}_2 &= \Omega_{21}^* \mathbf{e}_1 + \Omega_{22}^* \mathbf{e}_2 + \Omega_{23}^* \mathbf{e}_3 = -\omega \mathbf{e}_1 \end{aligned}$$

It follows that

$$[\overset{e-e}{\Omega}^*] = \omega \begin{bmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad \boldsymbol{\Omega}^* = \omega (\mathbf{e}_1 \otimes \mathbf{e}_2 - \mathbf{e}_2 \otimes \mathbf{e}_1)$$

**End of (ii).**

(iii) Determine  $\dot{\mathbf{F}}, \dot{\mathbf{U}}, \dot{\mathbf{R}}$  and  $\boldsymbol{\Omega}$

**Soln:** (iii)

Start with

$$\dot{\mathbf{F}} = \overset{e-E}{\dot{F}}_{ij} \mathbf{e}_i \otimes \mathbf{E}_j + \overset{e-E}{F}_{ij} \dot{\mathbf{e}}_i \otimes \mathbf{E}_j = \langle \mathbf{e} \rangle \otimes [\overset{e-E}{\dot{F}}] \{ \mathbf{E} \} + \langle \mathbf{e} \rangle [\overset{e-e}{\boldsymbol{\Omega}^{*T}}] \otimes [\overset{e-E}{F}] \{ \mathbf{E} \}$$

where

$$[\overset{e-E}{\dot{F}}] = \begin{bmatrix} \varepsilon & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$[\overset{e-e}{\boldsymbol{\Omega}^{*T}}][\overset{e-E}{F}] = \omega \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} (1+\varepsilon t) & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} = \omega \begin{bmatrix} 0 & -1 & 0 \\ (1+\varepsilon t) & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Define  $\overset{e-E}{\dot{F}}_{ij}^{Tot}$  such that  $\dot{\mathbf{F}} = \overset{e-E}{\dot{F}}_{ij}^{Tot} \mathbf{e}_i \otimes \mathbf{E}_j = \langle \mathbf{e} \rangle \otimes [\overset{e-E}{\dot{F}^{Tot}}] \{ \mathbf{E} \}$ . Then

$$\dot{\mathbf{F}} = \langle \mathbf{e} \rangle \otimes [\overset{e-E}{\dot{F}^{Tot}}] \{ \mathbf{E} \} \quad [\overset{e-E}{\dot{F}^{Tot}}] = \begin{bmatrix} \varepsilon & -\omega & 0 \\ \omega(1+\varepsilon t) & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Convert components:

$$[\overset{E-E}{\dot{F}^{Tot}}] = [\overset{E-e}{a}][\overset{e-E}{\dot{F}^{Tot}}] = \begin{bmatrix} c & -s & 0 \\ s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \varepsilon & -\omega & 0 \\ \omega(1+\varepsilon t) & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$= \begin{bmatrix} c\varepsilon - s\omega(1+\varepsilon t) & -c\omega & 0 \\ s\varepsilon + c\omega(1+\varepsilon t) & -s\omega & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

which agrees with 3.1.

The right stretch is already expressed in terms of the fixed basis so that the result is the same as 3.1.

$$\dot{\mathbf{U}} = \langle \mathbf{E} \rangle [\overset{E-E}{\dot{U}}] \otimes \{ \mathbf{E} \} \quad [\overset{E-E}{\dot{U}}] = \begin{bmatrix} \varepsilon & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Recall that  $\mathbf{R} = \mathbf{e}_i \otimes \mathbf{E}_i$  and  $\mathbf{R}^T = \mathbf{E}_i \otimes \mathbf{e}_i$ . Then the spin is

$$\boldsymbol{\Omega} = \dot{\mathbf{R}} \cdot \mathbf{R}^T = (\dot{\mathbf{e}}_i \otimes \mathbf{E}_i) \cdot (\mathbf{E}_j \otimes \mathbf{e}_j) = \dot{\mathbf{e}}_i \otimes \mathbf{e}_i = \langle \dot{\mathbf{e}} \rangle \otimes \{\mathbf{e}\}$$

Recall that we have derived

$$\{\dot{\mathbf{e}}\} = [\overset{e-e}{\boldsymbol{\Omega}^*}] \{\mathbf{e}\} \Rightarrow \langle \dot{\mathbf{e}} \rangle = \langle \mathbf{e} \rangle [\overset{e-e}{\boldsymbol{\Omega}^{*T}}]$$

Therefore, for this simple problem  $\boldsymbol{\Omega} = \boldsymbol{\Omega}^{*T}$  and  $\mathbf{n}_i = \mathbf{e}_i$ .

As a check, do the transformation

$$\begin{aligned} \overset{E-E}{[\boldsymbol{\Omega}]} &= \overset{E-e}{[a]} \overset{e-e}{[\boldsymbol{\Omega}]} \overset{e-E}{[a]} = \omega \begin{bmatrix} c & -s & 0 \\ s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & -I & 0 \\ I & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} c & s & 0 \\ -s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} \\ &= \omega \begin{bmatrix} -s & -c & 0 \\ c & -s & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} c & s & 0 \\ -s & c & 0 \\ 0 & 0 & 1 \end{bmatrix} = \omega \begin{bmatrix} 0 & -I & 0 \\ I & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \end{aligned}$$

which agrees with 3.1.

**End of solution to 3.2**



3.3 Obtain  $\mathbf{v} = \mathbf{v}(\mathbf{r}, t)$  and  $\mathbf{L}$ , and show that  $\dot{\mathbf{F}} = \mathbf{L} \cdot \mathbf{F}$

**Soln.** Start with

$$\begin{aligned} \mathbf{r} &= x_i \mathbf{e}_i & \mathbf{v} &= \dot{x}_i \mathbf{e}_i + x_i \dot{\mathbf{e}}_i & \{\dot{\mathbf{e}}\} &= [\overset{e-e}{\Omega^*}] \{\mathbf{e}\} \\ x_1 &= X_1(1 + \varepsilon t) & x_2 &= X_2 & x_3 &= X_3 \\ \mathbf{v} &= \left\langle \begin{matrix} \varepsilon X_1 & 0 & 0 \end{matrix} \right\rangle \{\mathbf{e}\} + \omega \left\langle \begin{matrix} x_1 & x_2 & x_3 \end{matrix} \right\rangle \left[ \begin{matrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{matrix} \right] \{\mathbf{e}\} \end{aligned}$$

Now substitute  $X_1 = \frac{x_1}{(1 + \varepsilon t)}$  to obtain

$$\mathbf{v} = v_i^e \mathbf{e}_i = \left\langle \begin{matrix} \frac{\varepsilon x_1}{(1 + \varepsilon t)} - \omega x_2 & \omega x_1 & 0 \end{matrix} \right\rangle \{\mathbf{e}\}$$

Then

$$\mathbf{L} = \overset{e-e}{L}_{ij} \mathbf{e}_i \otimes \mathbf{e}_j \quad \overset{e-e}{L}_{ij} = \frac{\partial v_i^e}{\partial x_j} \quad [\overset{e-e}{L}] = \left[ \begin{matrix} \frac{\varepsilon}{(1 + \varepsilon t)} & -\omega & 0 \\ \omega & 0 & 0 \\ 0 & 0 & 0 \end{matrix} \right]$$

and

$$\begin{aligned} \dot{\mathbf{F}} &= \mathbf{L} \cdot \mathbf{F} & \dot{\mathbf{F}} &= \overset{e-E}{\dot{\mathbf{F}}}_{ij}^{Tot} & [\overset{e-E}{\dot{\mathbf{F}}}_{Tot}] &= [\overset{e-e}{L}][\overset{e-E}{F}] \\ &= \left[ \begin{matrix} \frac{\varepsilon}{(1 + \varepsilon t)} & -\omega & 0 \\ \omega & 0 & 0 \\ 0 & 0 & 0 \end{matrix} \right] \left[ \begin{matrix} (1 + \varepsilon t) & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{matrix} \right] = \left[ \begin{matrix} \varepsilon & -\omega & 0 \\ \omega(1 + \varepsilon t) & 0 & 0 \\ 0 & 0 & 0 \end{matrix} \right] \end{aligned}$$

which agrees with the result in 3.2(iii).

3.4 Consider an element  $d\mathbf{X} = dX_1 \mathbf{E}_1$ . Determine the vector  $\mathbf{U} \cdot d\mathbf{X}$  and then the vector  $\mathbf{R} \cdot (\mathbf{U} \cdot d\mathbf{X})$ .

**Soln:** Use the approach of 3.2.

$$\mathbf{U} \cdot dX_1 \mathbf{E}_1 = \{(1 + \varepsilon t) \mathbf{E}_1 \otimes \mathbf{E}_1 + \mathbf{E}_2 \otimes \mathbf{E}_2 + \mathbf{E}_3 \otimes \mathbf{E}_3\} \cdot dX_1 \mathbf{E}_1 = (1 + \varepsilon t) dX_1 \mathbf{E}_1$$

which represents a stretch of the element with no rotation. Then

$$\mathbf{R} \cdot (\mathbf{U} \cdot d\mathbf{X}) = [\mathbf{e}_1 \otimes \mathbf{E}_1 + \mathbf{e}_2 \otimes \mathbf{E}_2 + \mathbf{e}_3 \otimes \mathbf{E}_3] \cdot (1 + \varepsilon t) dX_1 \mathbf{E}_1 = (1 + \varepsilon t) dX_1 \mathbf{e}_1$$

which represents a pure rotation of the stretched bar.