

OBSTRUCTION FOR FLAT STRING

Obstruction calculation for the flat string. Throughout we work on $\Sigma = \mathbb{C}$.

1. DEFORMATION COMPLEX

The full deformation complex is

$$\text{Def} = (\mathcal{O}_{\text{loc}}(\mathcal{E}_{\text{string}}), \bar{\partial} + \{I, -\})$$

There is a subcomplex $\text{Def}' \subset \text{Def}$ consisting of those functionals that only depend on $\mathcal{E}_{\text{grav}}$.

Proposition 1.1. *The obstruction lies in Def' .*

1.1. Heat Kernels. [this might fit better in an appendix](#) The heat kernel is a degree 1 element $\mathbb{K}_t \in (\mathcal{E}_{\text{string}})^{\otimes 2}$ that depends smoothly on $t \geq 0$. It splits as

$$\mathbb{K}_t = (\mathbb{K}_t^{\beta\gamma} \otimes 1) \oplus (1 \otimes \mathbb{K}_t^{\text{grav}}) \in \left((\mathcal{E}_{\beta\gamma}^{\oplus 13})^{\otimes 2} \otimes 1 \right) \oplus (1 \otimes \mathcal{E}_{\text{grav}})$$

For they single free $\beta\gamma$ -theory we have the explicit form

$$\mathbb{K}_t^{\beta\gamma}(z, w) =$$

2. ANOMALY CANCELATION

The machinery of [C1] allows us to define a pre-theory

$$I[L] := \lim_{\epsilon \rightarrow 0} W(\mathbb{P}_\epsilon^L, I - I^{\text{CT}}(\epsilon)).$$

Lemma 2.1. *The counterterms vanish. In other words, the limit*

$$\lim_{\epsilon \rightarrow 0} W(\mathbb{P}_\epsilon^L, I)$$

exists, and we denote it by $I[L]$.

The obstruction has the form

$$\Theta[L] := \bar{\partial}I[L] + \{I[L], I[L]\}_L + \Delta_L I[L].$$

Lemma 2.2. *T*

he obstruction is computed by

$$\Theta[L] = \lim_{\epsilon} W(\mathbb{P}_\epsilon^L, \mathbb{K}_\epsilon, I).$$

Lemma 2.3. *Only wheels (of valency 3) contribute to the obstruction. Further, tadpole and wheels with vertices greater than 2 don't contribute. That leaves two wheels of with two vertices each. One of them has inputs $\beta\gamma$ the other has inputs vector fields.*

Thus, the obstruction is

$$\begin{aligned} \Theta[L] &= \lim_{\epsilon \rightarrow 0} W_{\gamma_2}(\mathbb{P}_\epsilon^L, \mathbb{K}_\epsilon, I) \\ &= 13 \cdot \lim_{\epsilon \rightarrow 0} W_{\gamma_2}(\mathbb{P}_\epsilon^{\beta\gamma, L}, \mathbb{K}_\epsilon^{\beta\gamma}, I) + \lim_{\epsilon \rightarrow 0} W_{\gamma_2}(\mathbb{P}_\epsilon^{\text{string}, L}, \mathbb{K}_\epsilon^{\text{string}}, I). \end{aligned}$$

We already know from the DEFORMATION section that these functionals are only nonzero on vector fields. Denote by $\Theta_\epsilon^{\beta\gamma}[L](\xi_0, \xi_1)$ the $\beta\gamma$ weight applied to vector fields ξ_0, ξ_1 . Similarly, define $\Theta_\epsilon^{\text{grav}}[L](\xi_0, \xi_1)$.

We have

$$\Theta_\epsilon^{\beta\gamma}[L](\xi_0, \xi_1) = \int_{(z_0, z_1) \in \mathbb{C}^2} (\xi_0(z_0) \cdot \mathbb{P}_\epsilon^{\beta\gamma, L}(z_0, z_1)) (\xi_1 \cdot \mathbb{K}_\epsilon^{\beta\gamma}(z_0, z_1)).$$

We assume that ξ_0, ξ_1 are homogenous. Since the obstruction is a degree one element it is nonzero provided $\deg(\xi_0) + \deg(\xi_1) = 1$. For concreteness, suppose

$$\xi_0 = f(z, \bar{z})\partial_z, \quad \xi_1 = g(z, \bar{z})d\bar{z}\partial_z.$$

Then

$$\Theta_\epsilon^{\beta\gamma}[L](\xi_0, \xi_1) = \int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} (f(z_0)\partial_{z_0} P_\epsilon^L(z_0, z_1)) (g(z_1)\partial_{z_1} K_\epsilon(z_0, z_1)).$$

After making the change of coordinates this is

$$\int_{y_0, y_1} \int_\epsilon^L f g \epsilon^{-2} t^{-3} y_0^3 \exp\left(-\left(\frac{1}{t} + \frac{1}{\epsilon}\right) y_0\right) dt d\text{vol}_{\mathbf{y}}.$$

Taylor expanding fg in y_0 and performing Wick in the y_0 -variable this reduces to

$$\int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \int_\epsilon^L \epsilon^{-2} t^{-3} \frac{\epsilon t}{\epsilon + t} \left[\left(\frac{\epsilon t}{\epsilon + t} \right)^3 + \text{higher order terms} \right].$$

Taking the limit as $\epsilon \rightarrow 0$ we see that

$$\Theta^{\beta\gamma}[L] = \frac{1}{12} \int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1.$$

3. COMPUTING Θ^{grav}

Note that $\nu \wedge [\xi, \xi] = \nu L_\xi \xi$ where L_ξ denotes Lie derivative.

$$(1) \quad \Theta_\epsilon^{\text{grav}}[L](\xi_0, \xi_1) = \int_{(z_0, z_1) \in \mathbb{C}^2} (L_{\xi_0} \mathbb{P}_\epsilon^{\beta\gamma, L}(z_0, z_1)) (L_{\xi_1} \cdot \mathbb{K}_\epsilon^{\beta\gamma}(z_0, z_1)).$$

Again, we assume that ξ_0, ξ_1 are homogenous and we must have $\deg(\xi_0) + \deg(\xi_1) = 1$. For concreteness, suppose

$$\xi_0 = f(z, \bar{z})\partial_z, \quad \xi_1 = g(z, \bar{z})d\bar{z}\partial_z.$$

Expanding out, the integral (??) splits into four parts

$$(2) \quad \Theta_\epsilon^{\text{grav}}[L](\xi_0, \xi_1) = \int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} f(z_0) [\partial_{z_0} P_\epsilon^L(z_0, z_1)] g(z_1) [\partial_{z_1} K_\epsilon(z_0, z_1)]$$

$$(3) \quad - \int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} [\partial_{z_0} f(z_0)] P_\epsilon^L(z_0, z_1) g(z_1) [\partial_{z_1} K_\epsilon(z_0, z_1)]$$

$$(4) \quad - \int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} f(z_0) [\partial_{z_0} P_\epsilon^L(z_0, z_1)] [\partial_{z_1} g(z_1)] K_\epsilon(z_0, z_1)$$

$$(5) \quad + \int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} [\partial_{z_0} f(z_0)] P_\epsilon^L(z_0, z_1) [\partial_{z_1} g(z_1)] K_\epsilon(z_0, z_1).$$

We make the change of coordinates $y_0 = z_0 - z_1, y_1 = z_1$. Before evaluating the obstruction we compute the three basic integrals.

(I) First consider

$$\int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} (f(z_0) \partial_{z_0} P_\epsilon^L(z_0, z_1)) (g(z_1) \partial_{z_1} K_\epsilon(z_0, z_1)).$$

After making the change of coordinates this is

$$\int_{y_0, y_1} \int_\epsilon^L f g \epsilon^{-2} t^{-3} y_0^3 \exp\left(-\left(\frac{1}{t} + \frac{1}{\epsilon}\right) y_0\right) dt d\text{vol}_{\mathbf{y}}.$$

Taylor expanding fg in y_0 and performing Wick in the y_0 -variable this reduces to

$$\int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \int_\epsilon^L \epsilon^{-2} t^{-3} \frac{\epsilon t}{\epsilon + t} \left[\left(\frac{\epsilon t}{\epsilon + t} \right)^3 + \text{higher order terms} \right].$$

(II) Next we evaluate

$$\int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} (f(z_0) \partial_{z_0} \partial_{z_1} P_\epsilon^L(z_0, z_1)) (g(z_1) K_\epsilon(z_0, z_1)) = \int_{z_0, z_1} fg(\partial_{z_0} \partial_{z_1} PK).$$

After changing coordinates and Wick expanding as above this has the form

$$\int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \int_\epsilon^L \epsilon^{-1} t^{-4} \frac{\epsilon t}{\epsilon + t} \left[\left(\frac{\epsilon t}{\epsilon + t} \right)^3 + \text{higher order terms} \right].$$

(III) Finally we evaluate

$$\int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} (f(z_0) P_\epsilon^L(z_0, z_1)) (g(z_1) \partial_{z_0} \partial_{z_1} K_\epsilon(z_0, z_1)) = \int_{z_0, z_1} fg(P \partial_{z_0} \partial_{z_1} K).$$

Changing coordinates and Wick expanding this reduces to

$$\int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \int_\epsilon^L \epsilon^{-3} t^{-2} \frac{\epsilon t}{\epsilon + t} \left[\left(\frac{\epsilon t}{\epsilon + t} \right)^3 + \text{higher order terms} \right]$$

Since only the lowest order terms in the above integrals contribute in the limit as $\epsilon \rightarrow 0$ we will forget about the higher order terms in the following calculation. Consider the first integral in the second line above. We integrate by parts to get

$$\int (\partial_{z_0} f P)(g \partial_{z_1} K) = - \left(\int fg \partial_0 P \partial_{z_1} K + \int fg P \partial_{z_0} \partial_{z_1} K \right).$$

Integration by parts applied to the second line gives

$$\int (f \partial_{z_0} P)(\partial_{z_1} g K) = - \left(\int fg \partial_{z_0} \partial_{z_1} PK + \int fg \partial_{z_0} P \partial_{z_1} K \right).$$

Finally, we can integrate by parts twice for the term in the last line to get

$$\begin{aligned} \int (\partial_{z_0} f P)(\partial_{z_1} g K) &= - \left(\int f \partial_{z_0} P(\partial_{z_1} g K) + \int (f P)(\partial_{z_1} g \partial_{z_0} K) \right) \\ &= 2 \int fg \partial_{z_0} P \partial_{z_1} K + \int fg \partial_{z_0} \partial_{z_1} P + \int fg P \partial_{z_0} \partial_{z_1} K. \end{aligned}$$

Thus, we see the obstruction has the form

$$\begin{aligned} \Theta_\epsilon^{\text{grav}}[L](\xi_0, \xi_1) &= 5\text{I} + 2(\text{II} + \text{III}) + [\text{higher } \epsilon \text{ terms}] \\ &= \int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \int_{t=\epsilon}^L \frac{\epsilon}{(\epsilon + t)^4} (\epsilon t + 2(2\epsilon t + \epsilon^2 + t^2)) dt + [\text{higher } \epsilon \text{ terms}] \\ &= \int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \int_{t=\epsilon}^L \frac{\epsilon}{(\epsilon + t)^4} (\epsilon t + 2(\epsilon + t)^2) dt + [\text{higher } \epsilon \text{ terms}]. \end{aligned}$$

In the limit as $\epsilon \rightarrow 0$ we get

$$\Theta^{\text{grav}}[L](\xi_0, \xi_1) = \int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \left(\frac{1}{12} + 1 \right) = \frac{13}{12} \int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1.$$