

## CRITICAL CASE

We compute the critical wheel. In dimension  $d$  this is the wheel with  $d+1$  vertices. On  $\mathbb{C}^d$  we have seen the obstruction lives in

$$H_{\text{loc}}^1(\mathcal{L}^{\mathbb{C}^d})$$

which we can identify with  $H^{2d+2}(\text{BU}(d))$ . Write the generators of this space as  $c_\sigma$  where  $\sigma$  labels degree  $2d+2$  characteristic classes in  $d$ -(complex) dimensions. For instance when  $d=2$  we have two admissible indices and they correspond to

$$c_{\sigma_1} = c_1^3, \quad c_{\sigma_2} = c_1 c_2.$$

For coefficients in the tensor bundle  $E = T_X^{\otimes n} \otimes (T_X^\vee)^{\otimes m}$  we will find polynomials  $p_\sigma(n, m) \in \mathbb{Q}[n, m]$  such that under the above isomorphism the obstruction has the form

$$\Theta(d) = \sum_{\sigma} p_\sigma(n, m) c_\sigma.$$

$$1. \dim_{\mathbb{C}} = 1$$

We take coefficients in the tensor bundle  $T_{\mathbb{C}}^{\otimes n}$ , i.e.  $m=0$  in the above notation.

Recall that

$$f(z) \partial_z \cdot (\phi(z) \partial_z^{\otimes n}) := f(z) [\partial_z \phi(z)] \partial_z^{\otimes n} - n \phi(z) [\partial_z f(z)] \partial_z^{\otimes n},$$

extending the adjoint action to the  $n$ -fold power.

The obstruction has the form  $\Theta(1) = \lim_{\epsilon \rightarrow 0} \Theta_\epsilon(1)$  where  $\Theta_\epsilon(1)$  is the functional that sends  $(f \partial_z, g d\bar{z} \partial_z)$  to

$$\begin{aligned} & \int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} f(z_0) [\partial_{z_0} P_\epsilon^L(z_0, z_1)] g(z_1) [\partial_{z_1} K_\epsilon(z_0, z_1)] \\ & - n \left( \int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} [\partial_{z_0} f(z_0)] P_\epsilon^L(z_0, z_1) g(z_1) [\partial_{z_1} K_\epsilon(z_0, z_1)] - \int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} f(z_0) [\partial_{z_0} P_\epsilon^L(z_0, z_1)] [\partial_{z_1} g(z_1)] K_\epsilon(z_0, z_1) \right) \\ & + n^2 \left( \int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} [\partial_{z_0} f(z_0)] P_\epsilon^L(z_0, z_1) [\partial_{z_1} g(z_1)] K_\epsilon(z_0, z_1) \right). \end{aligned}$$

Recall that

$$K_t(z_0, z_1) = \frac{1}{t} e^{-|z_0 - z_1|/t}, \quad P_\epsilon^L(z_0, z_1) = \int_\epsilon^L (\bar{\partial}^* \otimes 1) K_t(z_0, z_1).$$

As in the previous sections we will utilize the change of coordinates  $y_0 = z_0 - z_1$ ,  $y_1 = z_1$ . We compute the three basic integrals.

(I) First consider

$$\int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} (f(z_0) \partial_{z_0} P_\epsilon^L(z_0, z_1)) (g(z_1) \partial_{z_1} K_\epsilon(z_0, z_1)).$$

After making the change of coordinates this is

$$\int_{y_0, y_1} \int_\epsilon^L f g \epsilon^{-2} t^{-3} y_0^3 \exp\left(-\left(\frac{1}{t} + \frac{1}{\epsilon}\right) y_0\right) dt d\text{vol}_{\mathbf{y}}.$$

Taylor expanding  $fg$  in  $y_0$  and performing Wick in the  $y_0$ -variable this reduces to

$$\int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \int_{\epsilon}^L \epsilon^{-2} t^{-3} \frac{\epsilon t}{\epsilon + t} \left[ \left( \frac{\epsilon t}{\epsilon + t} \right)^3 + \text{higher order terms} \right].$$

(II) Next we evaluate

$$\int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} (f(z_0) \partial_{z_0} \partial_{z_1} P_{\epsilon}^L(z_0, z_1)) (g(z_1) K_{\epsilon}(z_0, z_1)) = \int_{z_0, z_1} fg(\partial_{z_0} \partial_{z_1} PK).$$

After changing coordinates and Wick expanding as above this has the form

$$\int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \int_{\epsilon}^L \epsilon^{-1} t^{-4} \frac{\epsilon t}{\epsilon + t} \left[ \left( \frac{\epsilon t}{\epsilon + t} \right)^3 + \text{higher order terms} \right].$$

(III) Finally we evaluate

$$\int_{(z_0, z_1) \in \mathbb{C} \times \mathbb{C}} (f(z_0) P_{\epsilon}^L(z_0, z_1)) (g(z_1) \partial_{z_0} \partial_{z_1} K_{\epsilon}(z_0, z_1)) = \int_{z_0, z_1} fg(P \partial_{z_0} \partial_{z_1} K).$$

Changing coordinates and Wick expanding this reduces to

$$\int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \int_{\epsilon}^L \epsilon^{-3} t^{-2} \frac{\epsilon t}{\epsilon + t} \left[ \left( \frac{\epsilon t}{\epsilon + t} \right)^3 + \text{higher order terms} \right]$$

Since only the lowest order terms in the above integrals contribute in the limit as  $\epsilon \rightarrow 0$  we will forget about the higher order terms in the following calculation. Consider the first integral in the second line above. We integrate by parts to get

$$\int (\partial_{z_0} f P)(g \partial_{z_1} K) = - \left( \int fg \partial_0 P \partial_{z_1} K + \int fg P \partial_{z_0} \partial_{z_1} K \right).$$

Integration by parts applied to the second line gives

$$\int (f \partial_{z_0} P)(\partial_{z_1} g K) = - \left( \int fg \partial_{z_0} \partial_{z_1} PK + \int fg \partial_{z_0} P \partial_{z_1} K \right).$$

Finally, we can integrate by parts twice for the term in the last line to get

$$\begin{aligned} \int (\partial_{z_0} f P)(\partial_{z_1} g K) &= - \left( \int f \partial_{z_0} P(\partial_{z_1} g K) + \int (f P)(\partial_{z_1} g \partial_{z_0} K) \right) \\ &= 2 \int fg \partial_{z_0} P \partial_{z_1} K + \int fg \partial_{z_0} \partial_{z_1} P + \int fg P \partial_{z_0} \partial_{z_1} K. \end{aligned}$$

Thus we see the obstruction has the form

$$\begin{aligned} \Theta_{\epsilon}(1)(f \partial_z, g d\bar{z} \partial_z) &= (1 + 2n + 2n^2)I + (n^2 + n)(II + III) \\ &= \int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \int_{t=\epsilon}^L \frac{\epsilon}{(\epsilon + t)^4} (\epsilon t + (n + n^2)(2\epsilon t + \epsilon^2 + t^2)) dt \\ &= \int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \int_{t=\epsilon}^L \frac{\epsilon}{(\epsilon + t)^4} (\epsilon t + (n + n^2)(\epsilon + t)^2) dt. \end{aligned}$$

In the limit as  $\epsilon \rightarrow 0$  we get

$$\Theta(1) = \int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 \left( \frac{1}{12} + \frac{1}{2}(n^2 + n) \right) = \frac{1}{12} \int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1 (1 + 6n + 6n^2).$$

Identifying the cocycle

$$\frac{1}{12} \int_{y_1} \partial_{y_1}^3 (fg) dy_1 d\bar{y}_1$$

with the generator for the obstruction group we conclude that the polynomial for  $d = 1$  and  $m = 0$  is

$$p(n, 0) = 1 + 6n + 6n^2.$$

I know this isn't exactly the argument you had in mind but this is a slightly different way of using integration by parts and shows that we really only need to know how to evaluate TWO rational integrals. I think something similar should happen in higher dimensions.

In another direction, I think you are right in that we could try to just look at polynomial inputs and utilize a "recursive" Wick's lemma. I'm working on writing that up now for dimension 1.

## 2. GENERAL DIMENSION

Now, we work on  $\mathbb{C}^d$ . We modify a calculation of Si to calculate the obstruction. As we have already seen, only the wheel with  $d+1$  vertices contributes nontrivially. Thus, the obstruction is described by a functional

$$\Theta : \Omega^{0,*}(\mathbb{C}, T\mathbb{C})^{\otimes(d+1)} \rightarrow \mathbb{C}$$

which is explicitly described by the weight  $W_{\gamma_d}(P_\epsilon^L, K_\epsilon)$ . On a  $d+1$ -tuple of vector fields  $(\vec{\xi}_0, \dots, \vec{\xi}_d)$  its value is

$$\Theta(\vec{\xi}_0, \dots, \vec{\xi}_d) = \int_{(z_\alpha) \in (\mathbb{C}^d)^{d+1}} \left( \vec{\xi}_d \cdot K_\epsilon \right) \prod_{\alpha=0}^{d-1} \left( \vec{\xi}_\alpha \cdot P_{\alpha, \alpha+1} \right)$$

where  $P_{\alpha, \alpha+1} = P(z_\alpha, z_{\alpha+1})$  and  $K_\epsilon = K_\epsilon(z_0, z_d)$ .

We have a canonical framing

$$\Omega^{0,*}(\mathbb{C}, T\mathbb{C}) = \Omega^{0,*}(\mathbb{C}) \otimes \mathbb{C} \langle \partial_{z^1}, \dots, \partial_{z^d} \rangle.$$

Fix a collection of integers  $i_0, \dots, i_d \in \{1, \dots, d\}$ . We will compute

$$\Theta(A_0 \partial_{z_0^{i_0}}, \dots, A_d \partial_{z_d^{i_d}})$$

where  $A_\alpha \in \Omega^{0,*}(\mathbb{C})$  for  $0 \leq \alpha \leq d$ .

First, we compute  $\partial_{z_\alpha^{i_\alpha}} P_{\alpha, \alpha+1} = (\bar{z}_\alpha^{i_\alpha} - \bar{z}_{\alpha+1}^{i_\alpha}) \tilde{P}_{\alpha, \alpha+1}$  where

$$\tilde{P}_{\alpha, \alpha+1} = \int_{t_\alpha \in [\epsilon, L]} t_\alpha^{-1} \bar{\partial}_{z_\alpha}^* K_{t_\alpha}(z_\alpha, z_{\alpha+1}) dt_\alpha.$$

Similarly,  $\partial_{z_d^{i_d}} K_\epsilon = (\bar{z}_d^{i_d} - \bar{z}_0^{i_d}) \tilde{K}_\epsilon$  where  $\tilde{K}_\epsilon = \epsilon^{-1} K_\epsilon$ .

We also employ the following change of coordinates

$$w_\alpha = w_\alpha - w_{\alpha+1} \quad , \quad 0 \leq \alpha < d$$

and  $w_d = z_d$ .

Up to a sign, this weight is

$$\int_{(w_\alpha) \in (\mathbb{C}^d)^{d+1}} \left( \prod_{\alpha=0}^{d-1} \bar{w}_\alpha^{i_\alpha} A_\alpha \wedge \text{dvol}_\alpha^{\text{hol}} \right) (\bar{w}_0^{i_d} + \dots + \bar{w}_{d-1}^{i_d}) \tilde{K}_\epsilon \int_{\mathbf{t} \in [\epsilon, L]^d} \prod_{\alpha=0}^d \tilde{P}_{\alpha, \alpha+1} \text{dvol}_{\mathbf{t}}.$$

We compute the term

$$\begin{aligned} \tilde{K}_\epsilon \int_{\mathbf{t} \in [\epsilon, L]^d} \prod_{\alpha=0}^d \tilde{P}_{\alpha, \alpha+1} \text{dvol}_{\mathbf{t}} &= \pm \epsilon^{-d-1} \int_{\mathbf{t} \in [\epsilon, L]^d} \left( \prod_{\alpha=0}^{d-1} t_\alpha^{-d} \right) \\ &\quad \sum_{\mathbf{k}=(k_0, \dots, k_{d-1})} \epsilon_{\mathbf{k}} \left( \prod_{\alpha=0}^{d-1} t_\alpha^{-1} \bar{w}_\alpha^{k_\alpha} \right) \exp(-\bar{w}^T M w) \left( \prod_{\alpha=0}^{d-1} d^d \bar{w}_\alpha \right) \text{dvol}_{\mathbf{t}}. \end{aligned}$$

Here,  $\bar{w}^T M w$  denotes the quadratic form

$$\sum_{\alpha=0}^{d-1} |w_\alpha|^2 / t_\alpha - |w_0 + \cdots + w_{d-1}| / \epsilon.$$

Thus, when we write  $A_\alpha$  in the new  $w$ -coordinates only terms with  $d\bar{w}_d^j$  factors will contribute.