

April 10 (M, g) oriented, Riemannian manifold.

Recall we have defined the twistor space

$$\begin{aligned}\tau(M) &= Fr_M^{so} \times_{so(2n)} P(Ps^+) \\ &\cong Fr_M^{so} \times_{so(2n)} (so(2n)/u(n)) \\ &\cong Fr_M^{so} / u(n).\end{aligned}$$

• Fix an almost cplx str J_0 of \mathbb{R}^{2n} . Then any $A \in so(2n) \cong \wedge^2 \mathbb{R}^{2n}$ can be written

$$A = \frac{1}{2} (A - J_0 A J_0) + \frac{1}{2} (A + J_0 A J_0)$$

$$\Rightarrow so(2n) = u(n) \oplus \mathfrak{m}$$

where

$$\mathfrak{m} = \left\{ A \in so(2n) \mid A J_0 = -J_0 A \right\}$$

Thus $T_{J_0} (so(2n)/u(n)) \cong \mathfrak{m}$.

$$j : A \hookrightarrow J_0 A$$

endows m , and hence $SO(2n)/U(n)$ w/ an almost cplx structure.

$$m_{\mathbb{C}} = m \otimes_{\mathbb{R}} \mathbb{C} = m_+ \oplus m_-$$

where J_0 acts on m_{\pm} as $\pm i$.

$$m_{\pm} = J^{\pm} m, \quad J^{\pm} = \frac{1}{2} (1 \mp i J_0).$$

- The Levi-Civili connection determines a decomposition

$$T_z(M) = \mathcal{H} \oplus \mathcal{V}$$

where $\mathcal{H} = \pi^* TM$ and \mathcal{V} = vertical tangent bundle.

Have constructed an almost cplx str. on \mathcal{V} .

Since $\mathcal{H}_z \cong T_{\pi(z)} M$, there is also a tautological almost cplx str. on \mathcal{H} .

• Weyl tensors : Let (M, g) be Riem.

manifold, ∇ is the L.C. connection.

Then $R = F_{\nabla} \in \mathcal{N}^2(M, \text{End}(T_M))$.

Symmetrical: So, a priori, the Riemann curvature is a section

$$R \in \Gamma(M, T_M^* \otimes T_M^* \otimes T_M^* \otimes T_M)$$

$$\stackrel{112}{\Gamma(M, (T_M^*)^{\otimes 4})}$$

$$\cup$$

In fact.

$$R \in \Gamma(M, S^2(\wedge^2 T_M^*))$$

Locally, denote $V = T_x M$, so we are looking at

$$S^2(\wedge^2 V).$$

The Bianchi identity further constrains the symmetry:

$$b : S^2(\Lambda^2 V) \longrightarrow \Lambda^4 V$$

\leadsto

$$R_{\alpha} \in \ker b \subset S^2(\Lambda^2 V).$$

• Ricci contraction $r : \ker b \longrightarrow S^2 V$

Thm: When $\dim^n M \geq 4$ have

$$\ker b = \mathbb{R} \oplus S^2_0 V \oplus \ker r.$$

\uparrow
 $O(n)$

(Ricci scalar, Ricci trace-free, Weyl).

Thm: The Weyl tensor is conformally invariant.

Moreover if $n \geq 4$ then M is conformally

flat $\Leftrightarrow W \equiv 0$.

When $\dim M = 4$ we have

$$\Lambda^2 T^* M = \Lambda^2_+ \oplus \Lambda^2_-.$$

This induces a conformally invt decomposition

$$W = W_+ + W_-.$$

\mathbb{E}_X : $\dim M = 4k$. Then $H^{2k}(M, \mathbb{R})$ is
a symmetric vector space.

\leadsto symmetric bilinear form

$$B : H^{2k} \times H^{2k} \longrightarrow H^{4k} \cong \mathbb{R}$$

$$\leadsto Q(x) = B(x, x).$$

$$\sigma(M) = \text{signature of } Q$$

Thm: $\dim M = 4$:

$$\sigma(M) = \frac{1}{12\pi^2} \int_M (|W_+|^2 - |W_-|^2) \, \text{dvol}.$$

Rek: When $\dim M = 4$ and M is Kähler
then W_+ is completely determined by
the Ricci scalar curvature.

Thm: [Atiyah - Hitchin - Singer] M oriented, Riem.
manifold of $\dim_{\mathbb{R}} = 4$.

- Then the almost cplx str. of $\tau(M)$ is
integrable $(\Rightarrow) W_+ \equiv 0$.
- This cplx str. is Kähler $(\Rightarrow) M$ is conformally
equivalent to S^4 or \mathbb{CP}^2 .

[When $\dim > 4$ the only Kähler twistor
spaces arise from even spheres $M = S^{2n}$].

Ex: • $M = \mathbb{R}^4$. We will see how it is natural to view

$$\tau(M) \cong \text{Tot} \left(\begin{array}{c} \mathcal{O}(1) \oplus \mathcal{O}(1) \\ \downarrow \\ \mathbb{CP}^1 \end{array} \right)$$

• $M = S^4$. $\tau(S^4) \cong \mathbb{CP}^3$.

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• Let's turn to the question of integrability for sections of $\tau(M)$.

A "pure spinor field" is a section of $\tau(M)$
 \longleftrightarrow almost cplx str. of M
 cplx w/ metric + or...

Prop: A pure spinor field σ determines an integrable cplx str. (\Rightarrow)

$$v \cdot \nabla_w \sigma - w \cdot \nabla_v \sigma = 0$$

$$\forall v, w \in T(T^0, 1)_M.$$

Pf : $T_{\mathcal{H}}^{0,1}$ is defined in terms of the pure spinor field σ by

$$T_{\mathcal{H}}^{0,1} = \left\{ v \in \Gamma(T_{\mathcal{H}} \otimes \mathbb{C}) \mid v \cdot \sigma = 0 \right\}.$$

Take $w \cdot \sigma = 0$ and apply ∇_v :

$$0 = \nabla_v (w \cdot \sigma) = \nabla_v w \cdot \sigma + w \cdot \nabla_v \sigma.$$

Similarly $0 = \nabla_w v + v \cdot \nabla_w \sigma$.

$$\Rightarrow 0 = \underbrace{(\nabla_v w - \nabla_w v)}_{\nabla[v,w]} \cdot \sigma + v \cdot \nabla_w \sigma - w \cdot \nabla_v \sigma.$$

$$\text{So } v \cdot \nabla_w \sigma - w \cdot \nabla_v \sigma = 0$$

$$\Leftrightarrow [v, w] \cdot \sigma = 0 \Leftrightarrow [v, w] \in \Gamma(T_{\mathcal{H}}^{0,1}).$$

\square

Thm: Let M be or. Riem. of $\dim = 2m$.

Then the almost cplx str. determined by

$$s \in \Gamma(T(X))$$

is integrable (\Rightarrow) s is (almost) holomorphic.

Pf: Pick $[\sigma] = s$. Then s is holomorphic

$$(\Rightarrow) \quad \nabla_v \sigma = \lambda_v \sigma \quad \forall v \in \Gamma(T^{0,1})$$

where λ_v is some function depending on v .

Our prop \Rightarrow

$$s = \bigcap_i \ker(\mu_{\bar{\varepsilon}_i}) \quad , \quad \{\bar{\varepsilon}_i\} \text{ hermitian form for } T^{0,1}$$

Since $\bar{\varepsilon}_j^2 = 0$, we can use the above

prop to see that:

$$-\bar{\varepsilon}_i \bar{\varepsilon}_j \nabla_{\bar{\varepsilon}_j} \sigma = \bar{\varepsilon}_j \bar{\varepsilon}_i \nabla_{\bar{\varepsilon}_j} \sigma = \bar{\varepsilon}_j \bar{\varepsilon}_j \nabla_{\bar{\varepsilon}_i} \sigma = 0$$

$$\Rightarrow \bar{\varepsilon}_j \nabla_{\bar{\varepsilon}_i} \sigma \in s \subset s^\perp \quad \forall i, j.$$

But also $\bar{\varepsilon}_j \nabla_{\bar{\varepsilon}_i} \sigma \in S^-$ by type reasons.

\Rightarrow

$$(*) \quad \bar{\varepsilon}_j \nabla_{\bar{\varepsilon}_i} \sigma = 0 \quad \forall j.$$

Now: Consider the symmetric $C^\infty(X)$ -bilinear

$$\text{form } \beta(v, w) = w \nabla_v \sigma, \quad v, w \in \Gamma(T^{0,1}_H).$$

$(*) \Rightarrow \beta \equiv 0$. So $\nabla_v \sigma \in S = [\sigma]$ for

all $v \in \Gamma(T^{0,1})$.

Conversely, if $\nabla_v \sigma = \lambda_v \sigma$ then

automatically $w \cdot \nabla_v \sigma = 0$. \square