

CS IN CHIRAL GAUGE

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Fix a holomorphic volume form on Σ . Typically, when $\Sigma = \mathbb{C}$ we will pick the form dz .

In this section, we will work on flat space \mathbb{R}^3 which we will identify with $\mathbb{C} \times \mathbb{R}$. Choose holomorphic coordinates

$$(z, \bar{z}, t) \in \mathbb{C} \times \mathbb{R}$$

The fields of critical Chern-Simons theory are

$$(\gamma, \beta) \in \Omega^{0,*}(\mathbb{C}) \hat{\otimes} \Omega^*(\mathbb{R}) \otimes \mathfrak{g}[1] \oplus \Omega^{0,*}(\mathbb{C}) \hat{\otimes} \Omega^*(\mathbb{R}) \otimes \mathfrak{g}.$$

The action is

$$S(\gamma, \beta) = \int_{\mathbb{C} \times \mathbb{R}} dz \wedge \left(\langle \beta, d\gamma \rangle + \frac{1}{2} \langle \beta, [\gamma, \gamma] \rangle \right).$$

1. THE CHIRAL GAUGE

We use the gauge fixing condition

$$Q^{GF} = \bar{\partial}_{\mathbb{C}}^* + d_{\mathbb{R}}^*.$$

BW: might need to adjust factors in this equation so that $[Q, Q^{GF}]$ is the standard flat Laplacian in each degree. In this gauge, we can write the heat kernel as

$$K_T((z, \bar{z}, t), (w, \bar{w}, s)) = k_T^{an}((z, \bar{z}, t), (w, \bar{w}, s))(d\bar{z} - d\bar{w})(dt - ds) \otimes \text{Cas}_{\mathfrak{g}},$$

where $k_T^{an} \in C^\infty((\mathbb{C} \times \mathbb{R}) \times (\mathbb{C} \times \mathbb{R}))$, $T > 0$, is the heat kernel for the flat Laplacian on \mathbb{R}^3 . Explicitly

$$k_T^{an}((z, \bar{z}, t), (w, \bar{w}, s)) = \frac{1}{(4\pi T)^{3/2}} e^{-|z-w|^2/4T - |t-s|^2/4T}.$$

2. ONE-LOOP QUANTIZATION

Theorem 2.1. *Critical Chern-Simons admits a one-loop exact and finite quantization on $\mathbb{R}^3 = \mathbb{C} \times \mathbb{R}$. The space of one-loop quantizations is a torsor for the group $H_{\text{Lie}}^3(\mathfrak{g})$.*

Proposition 2.2. *The following limit*

$$\lim_{\epsilon \rightarrow 0} W(P_{\epsilon < L}, I)$$

exists for all $L > 0$. Thus, in the holomorphic gauge, critical Chern-Simons admits a finite prequantization.

Lemma 2.3. *In the holomorphic gauge, only diagrams of at most one-loop are nonzero in the expansion of $W(P_{\epsilon < L}, I)$.*

The strategy is to analyze the one-loop diagram by the order of the number of vertices. There is a special diagram that could potentially give rise to issues.

Lemma 2.4. *Let Γ_2 be the wheel with two trivalent vertices. Then $W_{\Gamma_2}(P_{\epsilon < L}, I) = 0$ identically as a distribution for all $\epsilon, L > 0$.*

Proof. Note that the propagator is of the form

$$P_{\epsilon < L}((z, \bar{z}, t), (w, \bar{w}, s)) = p_{\epsilon < L}^{an} \cdot ((d\bar{z} - d\bar{w}) - (dt - ds))$$

where $p_{\epsilon < L}^{an}$ is some smooth function on $(\mathbb{C} \times \mathbb{R})^2$. For any fields α_1, α_2 , the weight $W_{\Gamma_2}(P_{\epsilon < L}, I)(\alpha_1, \alpha_2)$ is given by some integral over $(\mathbb{C} \times \mathbb{R})^2 = \mathbb{C}_z \times \mathbb{R}_t \times \mathbb{C}_w \times \mathbb{R}_s$, but the integrand is proportional to

$$P_{\epsilon < L}((z, \bar{z}, t), (w, \bar{w}, s)) \wedge P_{\epsilon < L}((w, \bar{w}, s), (z, \bar{z}, t)).$$

This expression is identically zero being the wedge product of proportional one-forms on the same space. \square

Lemma 2.5. *Let Γ_k be the wheel with k trivalent vertices. Then, for $k > 2$ the limit*

$$\lim_{\epsilon \rightarrow 0} W_{\Gamma_k}(P_{\epsilon < L}, I)$$

exists for all $L > 0$.

When \mathfrak{g} is semi-simple, one can see in a similar way to the observations made for ordinary Chern-Simons that the obstruction deformation complex guarantees the existence of a quantization of critical Chern-Simons. However, we can see that the anomaly for critical Chern-Simons vanishes directly using the holomorphic gauge.

Lemma 2.6. *As above, let Γ_k denote the wheel with k trivalent vertices. When $k = 2$, one has*

$$W_{\Gamma_k}(P_{\epsilon < L}, K_\epsilon, I) = 0$$

identically as a distribution. When $k > 2$, then

$$\lim_{\epsilon \rightarrow 0} W_{\Gamma_k}(P_{\epsilon < L}, K_\epsilon, I) = 0.$$

This lemma shows that the one-loop obstruction to solving the quantum master equation of critical Chern-Simons on $\mathbb{C} \times \mathbb{R}$ in the holomorphic gauge vanishes for any Lie algebra \mathfrak{g} . Since the quantization is one-loop exact, this is enough to guarantee a full quantization.