

## GAUGING VERSUS THE BACKGROUND FIELD METHOD

### 0.1. Gauging = local BRST cohomology. [BW: recall ordinary BRST cohomology for symplectic actions](#)

We now explain, in more detail, the process of gauging a symmetry of a BV theory. We will find that the observables of the gauged theory introduce a “local” version of BRST cohomology. In our setup, of course, this will be a factorization algebra. As in the case of the equivariant observables, the process of gauging is defined whenever we have a symmetry of a theory by some local Lie algebra, but there are many key differences as we have explained in Table [BW: ref](#). First, we focus on the classical case.

0.1.1. Indeed, let  $\mathcal{E}$  be the fields of a classical BV theory on a manifold  $M$  described by some action functional

$$S = S^{free} + I \in \mathcal{O}_{loc}(\mathcal{E}),$$

satisfying the classical master equation  $\{S, S\} = 0$ . Let  $\mathcal{L}$  be a local Lie algebra that acts on the theory  $\mathcal{E}$  with corresponding Noether functional  $I^{\mathcal{L}} : \mathcal{L} \rightarrow \mathcal{O}_{loc}(\mathcal{E})$ . Recall, this means that the equivariant classical master equation  $(d_{\mathcal{L}} + Q)(I + I^{\mathcal{L}}) + \frac{1}{2}\{I + I^{\mathcal{L}}, I + I^{\mathcal{L}}\} = 0$ , where  $d_{\mathcal{L}}$  is the Chevalley–Eilenberg differential on  $\mathcal{L}$ , is satisfied.

We define a new BV theory built from this action of  $\mathcal{L}$  on  $\mathcal{E}$ . The space of fields is

$$\mathcal{E}_{BRST}^{\mathcal{L}} = \mathcal{E} \oplus \mathcal{L}[1] \oplus \mathcal{L}^![-2].$$

The  $(-1)$ -shifted symplectic pairing is given by the sum of the pairing defining the theory  $\mathcal{E}$  together with the natural pairing

$$\langle -, - \rangle_{\mathcal{L}} : \mathcal{L}[1] \times \mathcal{L}^![-2] \xrightarrow{ev} \text{Dens}_M \xrightarrow{\int_M} \mathbb{C}.$$

The first arrow is the density valued evaluation pairing between  $\mathcal{L}$  and its  $!$ -dual  $\mathcal{L}^! = \mathcal{L}^{\vee} \otimes \text{Dens}_M$ . Write  $X, X^!$  for elements of  $\mathcal{L}, \mathcal{L}^!$ , respectively. The action functional for this new theory is

$$S_{BRST}^{\mathcal{L}} = S + I^{\mathcal{L}} + \sum_{n \geq 0} \frac{1}{n!} \langle X^!, \ell_n(X, \dots, X) \rangle_{\mathcal{L}} \in \mathcal{O}_{loc}(\mathcal{E}_{BRST}^{\mathcal{L}}).$$

If we denote by  $\{\ell_n\}$  the  $L_{\infty}$  structure maps for  $\mathcal{L}$  we can think about this new theory as splitting up into a free part

$$S_{BRST}^{\mathcal{L}, free} = S^{free} + \langle X^!, \ell_1(X) \rangle_{\mathcal{L}}$$

and an interacting part

$$I_{BRST}^{\mathcal{L}} = I + I^{\mathcal{L}} + \sum_{n \neq 1} \langle X^!, \ell_n(X, \dots, X) \rangle_{\mathcal{L}}.$$

**Lemma 0.1.** *Let  $\{-, -\}_{BRST}$  denote the bracket on  $\mathcal{O}_{loc}(\mathcal{E}^{BRST})$  induced by the  $(-1)$ -shifted symplectic structure on  $\mathcal{E}^{BRST}$  defined above. The functional  $S_{BRST}^{\mathcal{L}}$  satisfies the curved classical master equation*

$$\{S_{BRST}^{\mathcal{L}}, S_{BRST}^{\mathcal{L}}\} = \ell_0.$$

**Definition 0.2.** The *local BRST complex* for the action of  $\mathcal{L}$  on  $\mathcal{E}$  is the factorization algebra  $\text{Obs}_{BRST}^{cl}$  of classical observables for the theory  $(\mathcal{E}_{BRST}^{\mathcal{L}}, S_{BRST}^{\mathcal{L}})$ .

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*Example 0.3.* Suppose that  $\Sigma$  is a Riemann surface with fixed metric  $g_0$ . We have already seen in Example ?? how the local Lie algebra controlling deformations of Riemannian metrics  $\mathcal{L}^{Riem}$  acts on the free boson on  $\Sigma$ . The corresponding BRST theory

$$\mathcal{E}_{BRST} = \mathcal{E}_{\Sigma}^{boson} \oplus \mathcal{L}_{(\Sigma, g_0)}^{Riem}[1] \oplus (\mathcal{L}_{(\Sigma, g_0)}^{Riem})^![-2],$$

is precisely the perturbative expansion of the Polchinski string ?? near the fixed metric  $(\Sigma, g_0)$ . Explicitly, the action functional is [BW: finish, copy from string](#)

*Example 0.4.* Suppose, again, that  $\Sigma$  is a Riemann surface. Consider local Lie algebra given by the Dolbeault complex of holomorphic vector fields  $\mathcal{L} = \mathcal{T}_{\Sigma} = \Omega^{0,*}(\Sigma, T_{\Sigma}^{1,0})$ .

0.1.2. [BW: Quantum story. Formula in the case that the theory is free.](#)

0.2. **Examples of BRST cohomology.**

0.3. **Chiral gravity.**