



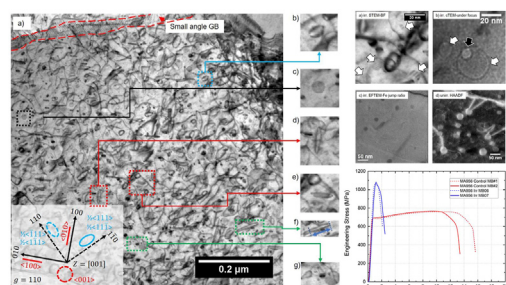
# Microstructure evolution in MA956 neutron irradiated in ATR at 328 °C to 4.36 dpa

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## GRAPHICAL ABSTRACT



## ARTICLE INFO

### Article history:

Received 5 December 2019

Received in revised form

3 March 2020

Accepted 4 March 2020

Available online 10 March 2020

## ABSTRACT

MA956 is an iron-chromium-aluminum (FeCrAl) based oxide dispersion strengthened (ODS) alloy produced by mechanical alloying. The alloy was irradiated in the Advanced Test Reactor (ATR) at 328 °C up to 4.36 dpa with both thermal and fast neutrons. The microstructures before and after irradiation were investigated by various TEM techniques. The size and number density of dislocation loops, voids, and oxides were analyzed. The results showed that both  $1/2\langle 111 \rangle$  and  $\langle 100 \rangle$  loops were generated in irradiated materials. The number density of dislocation loops reached up to  $7.17 \times 10^{21}/\text{m}^3$ . Cavities/voids were formed during irradiation. Argon bubbles were observed in unirradiated materials attached to the surface of oxide particles. The swelling rate was estimated to be 0.08% without subtracting pre-existed Ar bubbles. The oxide particles could maintain crystal structure during irradiation. The oxides had small decrease in size but half reduction in number density after irradiation. Calculation estimated that  $\langle 100 \rangle$  loops and dislocation lines contributed most to hardening. Alpha prime precipitates were suggested to be formed by comparing the calculated hardening. Dislocation lines are demonstrated not forming tangles in irradiated specimens.

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## 1. Introduction

Current designs for nuclear reactors like GenIV and fusion reactors require high performance of structure materials in extremely hostile environment, including high temperature, intense neutron fluxes and aggressive coolant [1,2]. Meanwhile accident tolerant fuel (ATF) cladding materials for light water reactors (LWR) need to

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have high oxidation resistance and good mechanical property at high temperature in beyond-design scenarios [3,4]. To this end, alternative candidates to traditional Zr-based alloys has been proposed, including modified/coating Zr-alloys, full ferritic Fe–Cr system steels, C-bearing ferritic/martensite (F/M) steels, refractory alloys, ceramics like SiC, and high entropy (HEA) concept alloys [5–7]. Among these candidates, the full ferritic steels are the most promising for in-core applications [4,7–9]. On one hand, the Fe–Cr–Al system could form a protective  $\text{Al}_2\text{O}_3$  layer in high temperature steam that meet the corrosion resistance demand for ATF materials. It has been proved that a commercial FeCrAl based steel APMT satisfies the X100 reduction in oxidation rate constant compared to Zr-alloys [10]. On the other hand, with oxide dispersion strengthening (ODS), the steels will enhance the creep and fatigue resistance over a wide range of temperature by pinning dislocations and grain boundaries [11]. At the same time, the dispersoids will improve the swelling tolerance by offering the sink site for vacancy and interstitial recombination under irradiation [11,12]. The uniformly distributed nanometer size oxide particles can also prevents embrittlement by trapping transmutant helium (He) into fine and less harmful bubbles [13]. As a result, the FeCrAl ODS steels have required excellent corrosion resistance, irradiation resistance and high temperature mechanical properties.

The operation temperature window for ODS ferritic steels is 300–700 °C [14]. Issues limiting the application of FeCrAl ODS steels under irradiation include 1) the irradiation embrittlement due to formation of dislocation loops and cavities, 2) phase stability and radiation induced segregation (RIS), and 3) the oxides stability, particularly the amorphization and dissolution. These issues are related to the chemical optimization and operation environment designs. For example, Cr could enhance the corrosion resistance and impede dislocation loops formation during irradiation, but high Cr content will cause aging embrittlement by  $\alpha'$  precipitation. Al could benefit to the oxidation resistance [15] and shift the miscibility gap of Fe–Cr to higher Cr direction [16] but trade off the tensile strength by coarsening the oxide particles [15]. The Ti-bearing (Al-free) ferritic ODS steels have lower swelling rate than Zr-bearing steels because of coarsen oxide particles owing to Al and Zr [17], however, Y–Ti–O oxides seems not as stable as Y–Zr–O or Y–Al–O oxides under irradiation environment [18]. Therefore, designs for ODS steels need considering the balance on the Cr and Al content, and proper bearing element like Ti, Zr, Hf and so on.

MA956 is an old commercial ferritic ODS steel and discontinued from market due to lack of industry demand [19]. The nominal composition of MA956 is Fe–20Cr–4.8Al–0.4Y<sub>2</sub>O<sub>3</sub>–0.4Ti. Detail composition was listed in Table 1. MA956 was considered only useful as a surrogate for an ATF FeCrAl ODS variant because optimized GenII ATF ODS FeCrAl have much improved ductility ensuring easier tube fabrication [20]. Mechanical properties of MA956 can be found in its commercial database [21]. Researches on unirradiated MA956 ranges from recrystallization, texture, DBTT and corrosion/oxidation [22–27]. Although these alloys are all called MA956, yet have four market shapes, e.g. sheet/plate and tube/bar, with different post-extrusion processing (rolling and annealing), that the microstructures and mechanical properties will be different to each other. The MA956 used in this study is in plate shape. It was produced from mechanical alloying (MA) with pre-alloyed powder (Fe–Cr–Al–Ti), yttria (Y<sub>2</sub>O<sub>3</sub>) and elementary

powder (Fe, Cr) in Ar shielding atmosphere [28–30]. After MA, the powders were consolidated by hot-extrusion. The consolidated billet was hot-rolled at 1050 °C to a thickness of 9.7 mm. The direction of hot-rolling was perpendicular to the direction of hot-extrusion. This cross-rolling yielded pancake-shape grains in the final plate.

Previous ion-irradiation in MA956 investigations were focusing on the oxide evolution and swelling behavior. The oxide particles of self-irradiated MA956 at 450 °C to 60–180 peak dpa tend to increase in size but reduce in number density [31]. This phenomenon might be explained as Ostwald ripening where dissolution of small nanoparticles (NPs) will redeposit onto or be absorbed by larger NPs [32]. The self-irradiation up to 2.4–60 peak dpa at 450 °C however, have constant oxide diameter but reduced number density [33]. Dual beam (Fe + He) simultaneous irradiation showed a helium shell around oxide core in MA956 at 425 °C [34]. Heavy Ne-ion implantation experiments revealed MA956 has smaller void growth at grain boundaries than F/M Grade 92 steel [35]. As for neutron irradiation, Gelles investigated various steels irradiated in the Fast Flux Test Facility (FFTF) at 420 °C to 200dpa [36]. The results showed the swelling of MA956 as well as MA957 measured by density change is only about 1%, indicating the commercial ferritic ODS steels are extremely resistant to radiation induced void swelling. Krumwiede et al. have reported the tensile test behavior of MA956 and other ODS steels after neutron irradiation [37]. The MA956 showed severe irradiation hardening and ductility reduction compared to Al-free 14YWT ODS steel.

In this study, we investigated the microstructure evolution of neutron irradiated MA956. The irradiation temperature was 328 °C which is in the typical light water reactor (LWR) operating temperature range. Three issues were addressed for the microstructure characterization. The first is statistical analysis of dislocation lines, loops and voids which contribute to the irradiation hardening and embrittlement. Particularly, the nature of dislocation loops, glissile  $\frac{1}{2}\langle 111 \rangle$  or sessile  $\langle 100 \rangle$  Burgers vectors were identified. The second is  $\alpha'$  precipitates which is essential in thermal aging but could be affected by irradiation. Even though we found no obvious  $\alpha'$  precipitates via current TEM techniques, they should be theoretically existed in this high Cr steel at the experiment temperature range. It is speculated that the resolution of TEM could not meet the minimum contrast requirement to recognize these fine precipitates. The third is oxide stability, including crystalline/amorphous and growth/dissolution behavior. These issues were characterized by TEM Moiré fringes and size/number density respectively. We also attached the tensile results of MA956 before and after irradiation. Finally, the irradiation hardening was calculated with Orowan equations and compared to the tensile yielding stress. The results will exhibit the microstructure morphology under irradiation and contribute to the database of neutron irradiated ODS steels.

## 2. Experimental

The irradiation in this paper is part of the large multipurpose program led by University of California Santa Barbara as part of the Advanced Test Reactor (ATR) Nuclear Science Users Facility (NSUF) Program [38]. As-received MA956 plates were obtained from Special Metals Corporation. The composition of the main elements is

**Table 1**  
Chemical composition of MA956 [21].

Fe	Cr	Al	Y <sub>2</sub> O <sub>3</sub>	Ti	Cu	Ni	Co	Mn	C	S
Bal	18.5–21.5	3.75–5.75	0.3–0.7	0.2–0.6	<0.15	<0.50	<0.3	<0.3	<0.1	<0.02

listed in Table 1. Fabrication of these plates has been introduced in Section 1. Neutron irradiation was carried out in ATR with pressurized light water coolant at Idaho National laboratory (INL) [39]. The neutron spectrum covered a wide range from  $1 \times 10^{-9}$  MeV–10 MeV. At the reactor power of 110MWth, the thermal neutron flux was about  $2 \times 10^{14}$  n/cm<sup>2</sup>-s while the fast neutron ( $E > 1$  MeV) flux was around  $2.3 \times 10^{14}$  n/cm<sup>2</sup>-s [40]. The samples were located at “A” capsule positions in ATR core and irradiated to accumulatively 4.36 dpa at  $328 \pm 10$  °C. The operating powers were between 108 and 116MWth. The total neutron fluences at each energy were given in Fig. 1.

TEM specimens were produced using an FEI Quanta and a Helios 600 DualBeam focused ion beam (FIB) micromachining. Both sides of the specimen were milled with 1 kV Ga<sup>+</sup> as a final step to minimize the damage from FIB. The microstructures were observed in FEI Tecnai F30 Super-Twin TEM. The electron beam source was provided by the field emission electron gun with the accelerating voltage at 300 kV. Both conventional TEM (cTEM) and Scanning TEM (STEM) mode were used to image the irradiation induced defects. In cTEM mode, the incident beam was parallel. The specimen was tilted to two beam kinematic condition with deviation parameter slightly positive and the dislocations were imaged closed to [001] zone axis with  $g$  equals (110). In STEM mode, a convergent beam was scanned on the specimen with the smallest convergent angle by inserting the smallest condenser aperture to avoid overlapping of systematic patterns which will generate phase contrast instead of diffraction contrast. The invisible criterion under STEM mode was demonstrated similar to that in cTEM [41]. During the practical operation, the specimen was first tilted to two beam condition in cTEM mode and then changed to STEM mode to make sure they have exactly the same diffraction condition. The imaging of dislocations in STEM mode could be on zone axis as well [42,43]. The details of the dislocation contrast imaging in STEM mode will be discussed and published elsewhere.

The thickness of TEM specimen ( $t$ ) was obtained from a relative thickness map using EFTEM utilizing equation (1) by considering the Poisson statistics of inelastic scattering. The  $I_0$  stands for the integral of zero-loss peak and  $I_t$  for the total integral of the EEL spectrum. The values of  $I_0$  and  $I_t$  corresponded to the intensity at each pixel in zero-loss image and unfiltered image respectively. The mean free path ( $\lambda$ ) was calculated with the collection (semi-) angle  $\beta$  set to 100 mrad because no objective aperture was inserted, that the scattering cross section will be independent of  $\beta$  above this value. The mean atomic number was 26 as it is Fe based alloy. The

calculation method could be found in Ref. [44].

$$t / \lambda = -\ln(I_t / I_0) \quad (1)$$

Line intercept method was used to measure dislocation density. In equation (2),  $\rho$  means the dislocation line density,  $N$  is the total intersections,  $l$  is the total line length,  $t$  is the thickness of specimen. The coefficient  $\alpha = 2$  by assuming the dislocations are spatially random oriented [45].

$$\rho = \alpha N / lt \quad (2)$$

Dislocation loops were measured under STEM two-beam like diffraction condition. The axis was chosen to be [001] and  $g$  vector parallel to (110). The diameter of dislocation loop is determined by the long axis of the projected eclipse.

Cavities were examined by out-of-focus method. The cavities will be bright at defocus condition in cTEM due to phase contrast. The voids were imaged with defocused beam where no diffraction pattern was strongly excited.

The  $\alpha'$  precipitation was examined by EFTEM jump ratio technique using one pre-edge and one post-edge of Cr to pick out the Cr concentrated or depleted part.

Oxides were examined via various TEM techniques. The oxides in unirradiated materials were imaged in STEM-HAADF mode where Z-contrast was dominate. The advantage of this mode is little moiré fringes shown compared to bright field that affect the precision of oxide size measurement. The diffraction contrast still remains owing to the relatively large collection angle. In this case, both the dislocations and oxides have bright contrast. As the beam intensity is much higher in STEM mode, a thicker area ( $t > 140$  nm) can be investigated to increase the accuracy of measurement. For the irradiated materials, we performed EFTEM to avoid the overlapping of contrast of high dense dislocation lines and oxides in diffraction-based method. The images were taken in thin area ( $t < 90$  nm) to ensure enough electrons collected.

Tensile tests were performed on 30 kN capacity Instron 5567 screw driven load frame in the Wing 9 hot cells at the Chemistry and Metallurgy Research Facility (CMR) at LANL. Tests were performed at constant cross head velocity of 0.15 mm/min corresponding to an initial strain rate of  $5 \times 10^{-4}$ /s. The tensile samples are typical S1/SSJ2 type, with 0.5 mm thickness, 5 mm gauge length and 1.2 mm width. The tensile experiments were performed at room temperature.

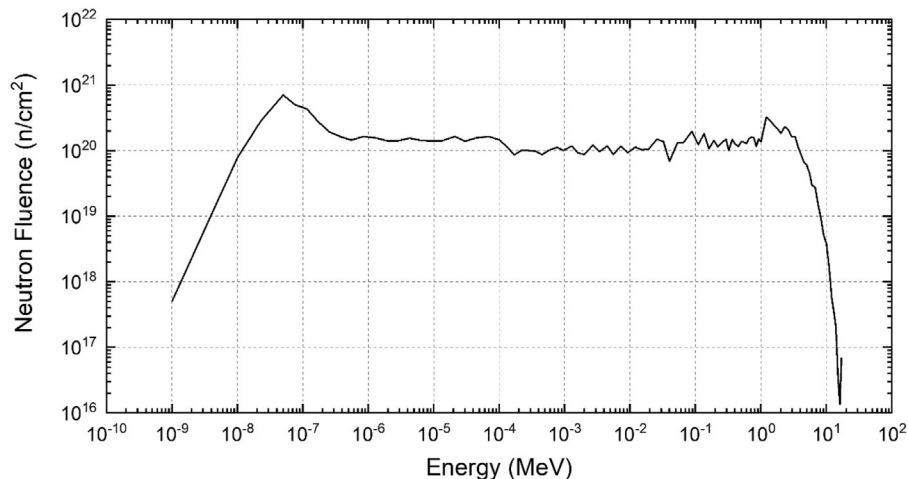


Fig. 1. Summation of neutron fluence vs energy at experiment reactor powers.

### 3. Results

#### 3.1. Overview of microstructures in unirradiated MA956

The overview of the microstructures of an unirradiated reference MA956 is shown in Fig. 2a. The grains are elongated in one direction because of mechanical processing. While the average width of the grains is about 0.5  $\mu\text{m}$ , the length is more than 1  $\mu\text{m}$  which leads to a grain aspect ratio (GAR) larger than two. To identify if there is any preferred grain orientation, the selected area aperture was inserted to the center of Fig. 2a with grains included as many as possible. The coincided diffraction pattern attached on the left top corner in Fig. 2a indicated these grains sharing a preferred crystal orientation. Coarse estimation showed the viewed plane (cross-section of the rolled plate) is close to (110)  $\langle 110 \rangle \langle 001 \rangle$  corresponding to a strong  $\langle 110 \rangle$   $\alpha$ -fiber to the elongated direction in cross-rolled ferritic steels [46]. The sub-grain boundaries could be seen across the elongated grains. The low angle sub-grain boundaries could be viewed as dislocation walls or networks for tilted or twisted boundaries respectively. Further investigation showed the dislocations inside the grains are at a very high density because of working hardening. Fig. 2b shows an enlarged region of interest (ROI) where Ar bubbles are attached to one side of the oxides. These bubbles were formed during MA processing with Ar as protection gas [47]. Existence of Ar bubbles indicates the reference materials did not reach to their full mass density.

#### 3.2. Irradiation induced dislocation loops and voids in MA956

##### 3.2.1. Dislocation loops

The irradiated specimen was observed under STEM bright field (BF) mode. The orientation was tilted close to [001] zone axis. Images were taken under two beam condition with reflecting vector  $g = (110)$ . In BCC materials, the burgers vectors of irradiation induced dislocation loops are  $1/2\langle 111 \rangle$  and  $\langle 100 \rangle$  types. Note that there are totally four types of  $1/2\langle 111 \rangle$  and three types of  $\langle 100 \rangle$  in crystallography. The irradiation induced dislocation loops were pre-assumed to be pure edge type that the habit plane could be read out directly from their Burgers vectors. Assuming the dislocation loops are perfect round circles, then the projected loop images could be calculated by equations (3) and (4).

$$\vec{l} = \vec{Z} \times \vec{b} \quad (3)$$

$$b/a = |\cos\theta| \quad (4)$$

Where  $\vec{l}$  is the orientation of projected loop (long axis),  $\vec{Z}$  is the observation zone axis,  $\vec{b}$  is the Burgers vector of dislocation loop,  $b$  and  $a$  stand for the short axis and long axis of the projected ellipse respectively,  $\theta$  is the angle between zone axis and burgers vector, which equals to the complementary angle between project plane and habit plane.

Yao et al. [48] has summarized the relationship between projected images of dislocation loops and their Burgers vectors. Here we recalculated the relationship under the [001] zone and re-indexed in the right-handed coordinate system. Fig. 3 shows the bright field images in STEM mode. In Fig. 3a, dislocation loops could be seen along a typical orientation. When viewed from [001] zone axis, all the  $1/2\langle 111 \rangle$  type dislocation loops will have the ratio  $b/a = 0.577$ . At the same time,  $1/2\langle 111 \rangle$  and  $1/2\langle 1\bar{1}1 \rangle$  will be projected along  $[1\bar{1}0]$  direction, as shown in Fig. 3b, and the  $1/2\langle \bar{1}11 \rangle$  and  $1/2\langle 1\bar{1}\bar{1} \rangle$  loops will be projected along  $[110]$  direction. The oxides are perfect round solid circles as shown in Fig. 3c. The  $\langle 100 \rangle$  and  $\langle 010 \rangle$  loops will have  $b/a = 0$  and should exhibit as a line or slim ellipse owing to sample tilting. They will extend to [010] and [100] respectively, as shown in Fig. 3d and e. The  $\langle 001 \rangle$  loops whose habit plane is perpendicular to the zone axis, will be a round unfilled circle.

When  $g = (110)$ , the  $1/2\langle \bar{1}11 \rangle$ ,  $1/2\langle 1\bar{1}1 \rangle$  and  $\langle 001 \rangle$  loops will be invisible. However, there are some other loops found in the specimen. Fig. 3f shows an example of a slim loop that extends along the  $[1\bar{1}0]$  direction which does not obey any of the aforementioned rules. The reason could be owing to the surface effect that the loop changed its orientation, or the inside contrast instead of outside contrast that made it looks slim. We do not expect it to be a  $\langle 110 \rangle$  type loop because they should not be induced by irradiation in BCC materials at high temperature.

The size distribution of dislocation loops is summarized in Fig. 4. A total of 176 loops were measured with a Java programmed tool ImageJ [49]. The curves were log-normal fitted instead of a Gaussian distribution. Most of the loops were distributed within 5–25 nm, while the largest loop size was over 60 nm. The distribution of the measured loops ignored invisible loops in TEM images. If considering the extinction criterion, the actual number density of loops corresponds to  $7.17 \times 10^{21}/\text{m}^3$ .

##### 3.2.2. Voids and swelling

Swelling occurs when vacancies start to migrate above the Stage

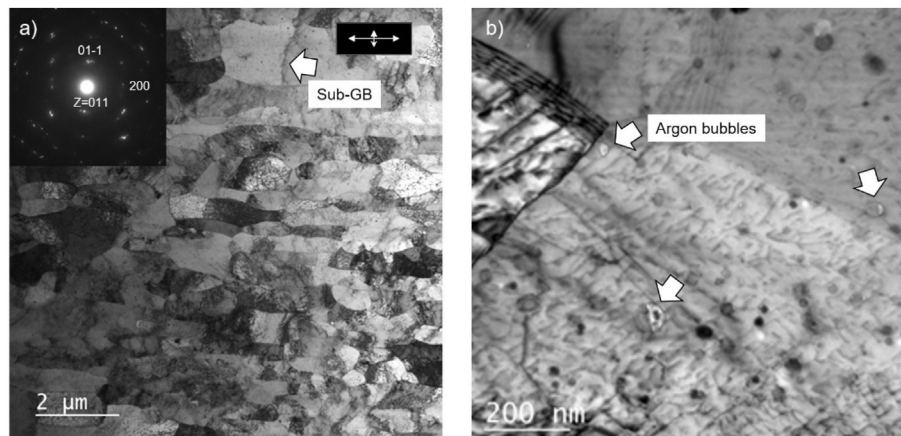
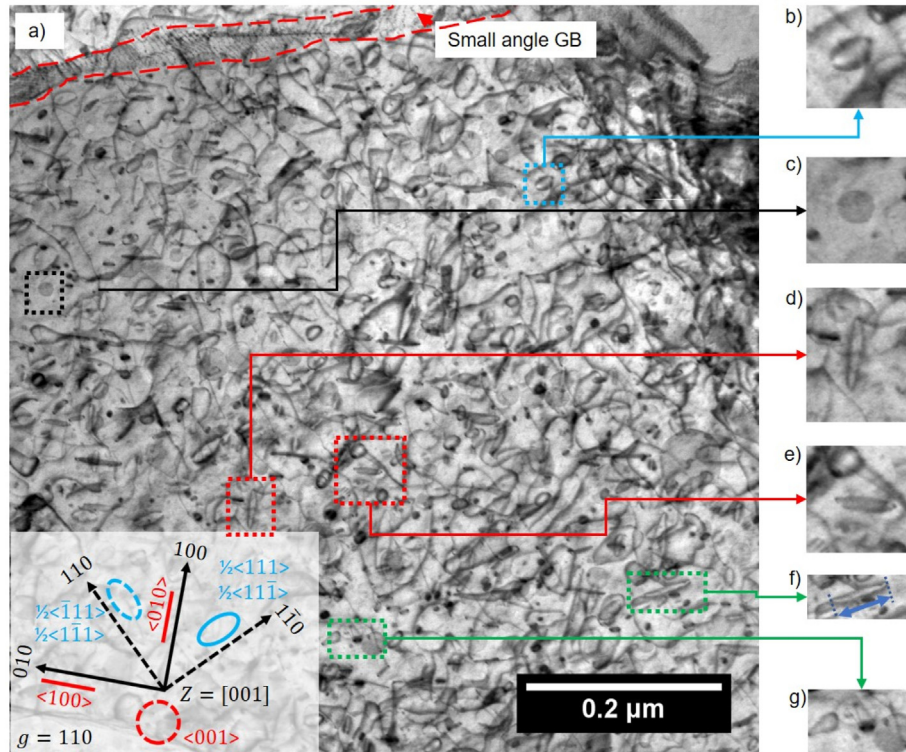
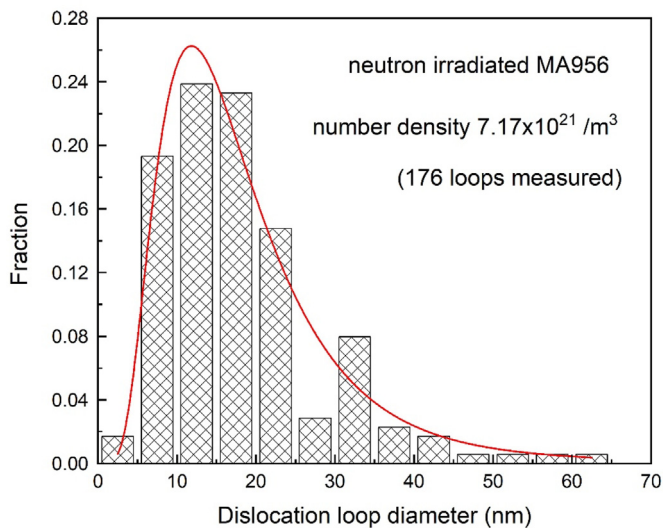


Fig. 2. The TEM images of reference MA956, a) the overview of grain morphology and b) the Ar bubbles attached to one side of oxide particles.





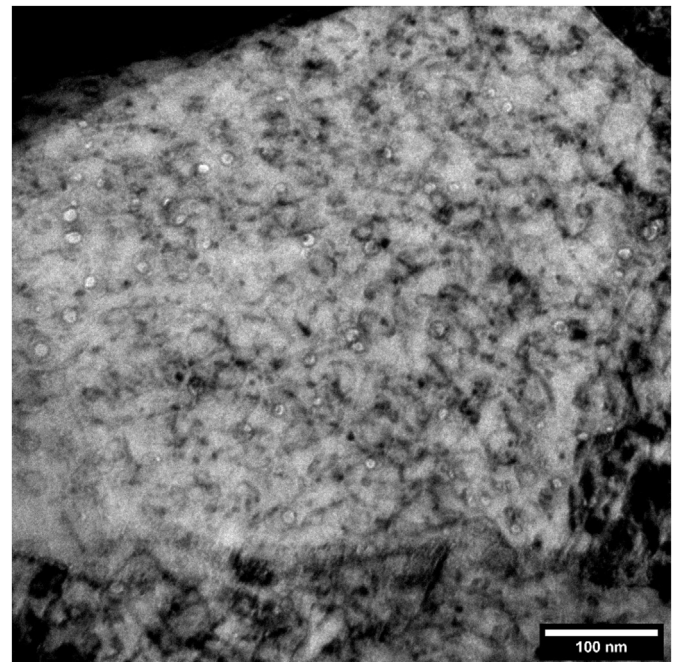
**Fig. 3.** a) the STEM bright field (BF) images of irradiated MA956, b)  $1/2\langle 111 \rangle$  or  $1/2\langle 11\bar{1} \rangle$  dislocation loop, c) oxide particle d)  $\langle 010 \rangle$  loop, e)  $\langle 100 \rangle$  loop, f) a loop with unknown Burgers vector, and g) a loop showed as black dots.



**Fig. 4.** Size distribution of dislocation loops in irradiated MA956.

III temperature range. In neutron irradiation, transmutant helium may combine with vacancies and form into stable helium-vacancy complexes. Typically, cavities without gas atoms are named voids, while the cavities containing inert gas are called bubbles. However, to distinguish from Argon bubbles which already existed in the reference materials, the cavities induced by neutron irradiation are called voids here, regardless of amount of helium in them.

Fig. 5 shows the void morphology in MA956 after neutron irradiation. Although some of those are Ar bubbles attached to the oxide particles as previously mentioned in Fig. 2b, most of the voids



**Fig. 5.** Under-focused TEM image shows the voids in MA956 after neutron irradiation.

are free and randomly distributed inside the grains. The average size of the voids is  $8.7 \pm 2.9$  nm, and the number density is  $1.80 \times 10^{21}/\text{m}^3$ . There is no preference of voids being attracted to grain boundaries or any void denuded zone. As it is difficult to distinguish the Ar bubbles and voids under current TEM method, the estimated size and number density of voids and further

swelling ratio are with Ar bubbles included. Even the density of Ar bubbles cannot compare to the voids, the swelling could be over-estimated by current method.

The way to calculate swelling via TEM images uses the following equation [50]:

$$\% \Delta V/V = 100 \left[ \Delta V / (V_f - \Delta V) \right] \quad (5)$$

$$\Delta V = \sum_i \frac{1}{6} \pi d_i^3 \quad (6)$$

$$V_f = At \quad (7)$$

Where  $\Delta V$  is the change in volume,  $V$  is the original volume of the material,  $V_f$  is the final volume of the material,  $d_i$  is the diameter of the  $i$ th void,  $A$  is the area of the TEM image,  $t$  is the sample thickness.

The swelling rate calculated by equation (5) is approximately 0.08%. The volume of Ar bubbles is estimated to be 0.007%. As there are few reports on neutron irradiated MA956, we cannot precisely evaluate the swelling per dpa at this temperature directly. The swelling in neutron irradiated Fe–Cr binary alloys has been summarized by Garner et al. [51] and 0.2%/dpa post transient swelling rate is reported after incubation. The concentration of Cr however, may suppress the swelling at high temperatures. When irradiated at ~450 °C, 15%Cr showed reduced swelling than 10%Cr, but when irradiated ~400 °C, swelling in 15%Cr was much greater than 10%Cr [51,52]. In Cr<sup>+</sup> irradiated MA957, peak swelling occurred between 420 °C and 450 °C [53]. The results showed MA957 has a higher onset swelling dose than EP450 and HT9. It is likely that the fine oxide dispersion suppressed void formation. In the case of the Fe–Cr–Al ternary system, Field and his coworkers found a small number of cavities in FeCrAl alloys at around 400 °C to 1.8 dpa, which may help demonstrate the effect of Al on swelling [54]. However, Song et al. showed that after helium implantations, Al-free Fe–Cr ODS steel showed less swelling than FeCrAl ODS steel in the temperature range from 300 °C to 700 °C [17]. In the current study, we only observed a small amount of swelling induced by neutron irradiation in MA956. More research needs be done on this type of Al-added Fe–Cr ODS steels in the future to understand the swelling behavior thoroughly.

### 3.3. Alpha-prime precipitates

Cr is an important element to raise the corrosion resistance in Fe based steels. However, distribution of Cr may be inhomogeneous. The Cr may form into Cr-rich  $\alpha'$  precipitates [55–59], form a shell structure around oxides [60,61], and deplete from or concentrate at grain boundaries owing to RIS [43,62–64]. In aging experiments at 475 °C after 2000 h, the FeCrAl-ODS steels with 12%–18% Cr (wt%) showed aging embrittlement due to  $\alpha'$  precipitate [60]. Irradiation accelerated formation of  $\alpha'$  is well established in both FeCr binary alloys [56] and FeCrAl ternary alloys [57]. Although high Cr content (>30 at%) will lead to spinodal decomposition, the  $\alpha'$  in this study's range should experience a nucleation/growth procedure, according to recent study on Fe–Cr phase diagram [65]. The additive Al however, can impede the formation of  $\alpha'$  during irradiation [66]. Irradiation mixing may induce concentration variation that enhance the precipitate process. Fig. 6 shows the EFTEM jump ratio images of Fe and Cr in the irradiated materials. The investigation did not indicate any  $\alpha'$  in the matrix or core-shell structures around oxides. However, a recent study of Fe-18 at% Cr irradiated in similar neutron environment shows  $\alpha'$  precipitates by atom probe

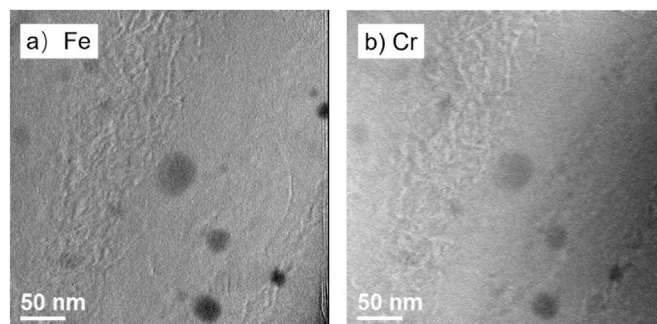


Fig. 6. EFTEM of a) Fe and b) Cr jump ratio images, which showed no obvious  $\alpha'$  precipitate in the matrix at this imaging technique.

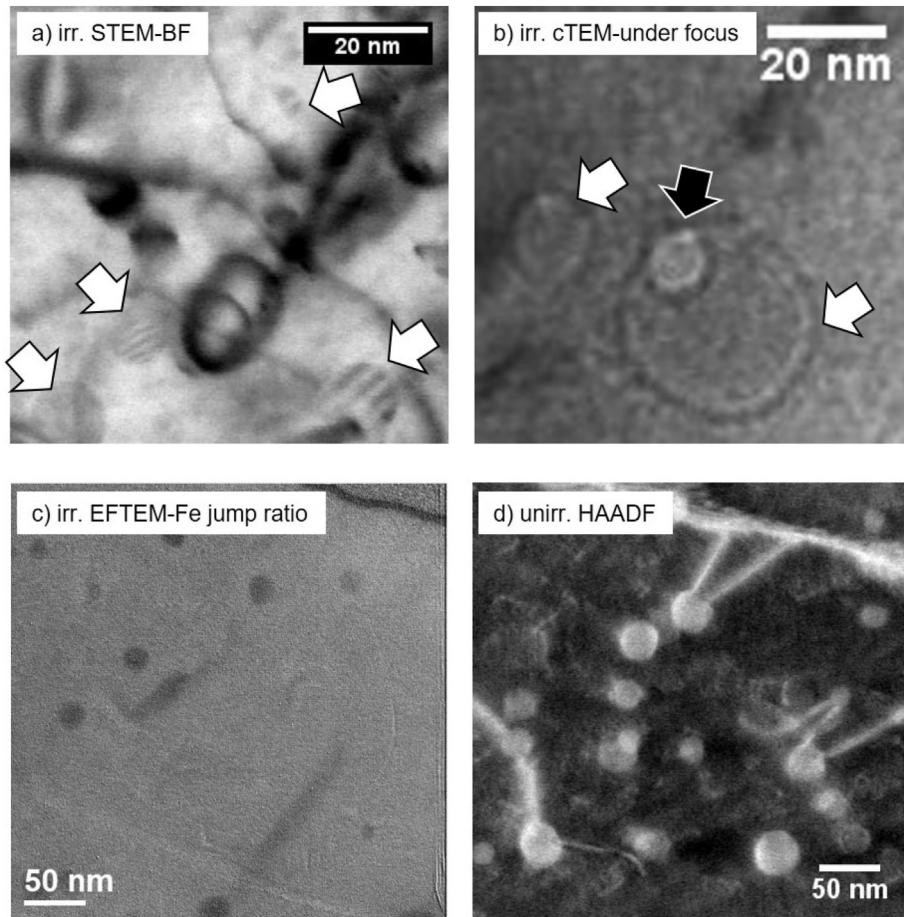
tomography (APT) method [58]. Another Al-free steel, 14YWT, also showed  $\alpha'$  precipitates after irradiations up to 7dpa at 360–370 °C [67]. Briggs et al. performed neutron irradiations to FeCrAl steels similar to MA956 [68]. The Fe-17.5Cr–3Al (wt%) steel was irradiated in High Flux Isotope Reactor (HFIR) at 320 °C to 7dpa. The results by small-angle neutron scattering (SANS) analysis showed the  $\alpha'$  precipitates have average size of 2.5 nm with number density of  $0.7 \times 10^{24} \text{ m}^{-3}$ . The reason that  $\alpha'$  was not observed in this paper could be owing to the small size of precipitates that exceeded the TEM technique limitation.

### 3.4. Oxides

There are various types of particles in Al-bearing ODS materials like MA956. The alumina particles ( $\text{Al}_2\text{O}_3$ ) could be around 500 nm while titanium carbonitrides  $\text{Ti}(\text{C},\text{N})$  are about 100–200 nm [69]. Both of the particles are at very low densities and are not interesting in this study. The Y–Al–O particles however, with an average size of 10 nm and a density of up to  $5 \times 10^{21} \text{ m}^{-3}$  [70], will contribute the most as trapping sites to the interstitials and vacancies induced by irradiation. The Y–Al–O particles could be YAP ( $\text{YAlO}_3$ , perovskite), YAG ( $\text{Y}_3\text{Al}_5\text{O}_{12}$ , garnet), YAM ( $\text{Y}_4\text{Al}_2\text{O}_9$ , monoclinic), YAP' ( $\text{YAlO}_3$ , pseudo-perovskite), YAH ( $\text{YAlO}_3$ , hexagonal), and YAT ( $\text{Y}_3\text{Al}_5\text{O}_{12}$ , tetragonal) [71]. The formation of oxide particles depends on annealing time and temperature [72], but may also be related to other factors such as Ti concentration [73]. Previous literatures showed lattice images of both YAT structure [35] and YAM structure [31] in MA956 using HRTEM. The X-ray diffraction (XRD) spectrum of extracted particles revealed that the oxides contain both YAP and YAG [74]. A study from diffraction pattern of dispersoids retrieved in carbon replica and STEM-EDS composition analysis claimed that most of the particles are YAP while about 10% are probably YAM [70].

Fig. 7 shows the oxide particle images using different TEM techniques. Fig. 7a shows the irradiated specimen using STEM-BF technique under two beam condition. One can see the coexistence of dislocation loops and oxide particles. Most of the oxide particles showed moiré fringes owing to the lattice misfit between particles and matrix. Fig. 7b shows the particles by under-focused cTEM. The specimen was tilted to avoid exciting any diffraction pattern. At this condition, diffraction contrast is suppressed, and voids will show up clearly. The black arrow in Fig. 7b indicates a void attached to the surface of an oxide particle, which can be considered as an evidence that the particles could act as a sink site for vacancy type defects. EFTEM was used to analyze particle size because dislocations will not appear in these images. Fig. 7c applied the jump ratio technique choosing one pre-edge and one post-edge of Fe to pick out the Fe depleted part, where the oxide particles will

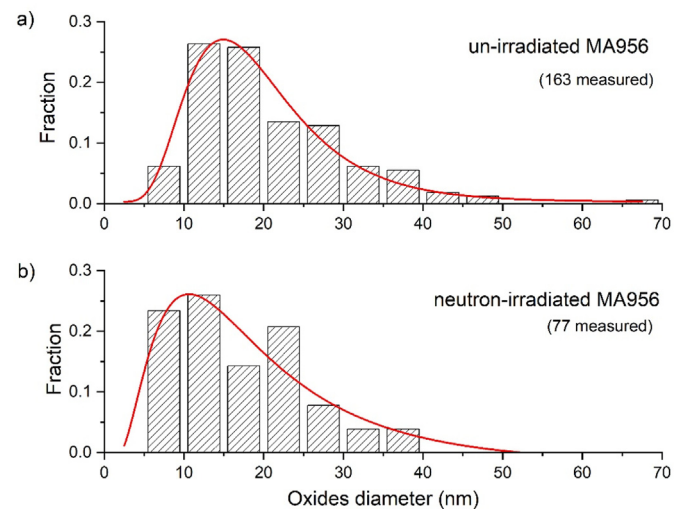




**Fig. 7.** The TEM images of oxide particles, the white arrows point to oxides and the black to voids, a) coexistence of dislocations and oxides, the moiré fringes indicate oxides remained crystal structure after irradiation, b) void attached on the surface of oxide, c) EFTEM image for oxides analysis after irradiation, where dislocation lines/loops did not show contrast, and d) STEM HAADF image for oxides characterization before irradiation, where dislocations and oxides showed strong bright contrast. (irr = irradiated, unirr = unirradiated).

show dark contrast. Compared with the irradiated specimen, an unirradiated one was imaged by STEM-HAADF, as shown in Fig. 7d. In this condition light elements will appear dark and heavy elements bright, following the Z-contrast rules. However, the dislocations will still appear bright when camera length is not small enough, where certain planes giving strong reflection absorbed by the HAADF detector that diffraction contrast will remain in the dark field image. In Fig. 7d, the dislocations could be seen folded when interacted with an oxide particle.

The oxides size distributions are shown in Fig. 8. Before irradiation, the diameter of oxides are mainly distributed between 10 and 30 nm, which is in agreement with Ref. [74]. After irradiation, the size distribution slightly moves to a lower range, which suggests that the oxide particles shrink or dissolve during irradiation. The dissolution process is irradiation dose, dose rate and temperature dependent. Generally, higher temperature and higher dose will cause heavier dissolution according to Ref. [75]. A recent study showed increased dissolution of pyrochlore  $\text{Y}_2\text{Ti}_2\text{O}_7$  under ion-irradiation at room temperature with increasing damage level (dpa dependence) [76]. In addition, Chen et al.'s work [77] suggests that at lower temperatures, the dissolution will be more significant (temperature dependence), and this behavior may also related to the coherency. The dissolution rate could be retarded by reducing the dose rate, that neutron irradiation might have a lower dissolution rate than ion irradiation. In the current study, the oxides in MA956 exhibit a slight reduction in size after neutron irradiation,



**Fig. 8.** The oxides diameter distribution of MA956 before and after irradiation.

from an average of 20.1 nm to that of 17.5 nm, but the number density reduced about a half, from  $\sim 16.2 \times 10^{20}/\text{m}^3$  before irradiation to  $\sim 8.4 \times 10^{20}/\text{m}^3$  after irradiation. The reason may be owing to different structure and/or different coherent relations between

matrix and oxides leading to different dissolution ability. In general, the oxide particles were considered stable because the order of magnitude didn't change after irradiation. In Fig. 7a, the moiré fringes indicate that the oxides maintained their crystal structure, which means these oxides have good irradiation resistance against amorphization. The number density and size of dislocation lines, loops, voids and oxides before and after irradiation are summarized in Table 2.

### 3.5. Tensile tests

Tensile tests were performed at CMR in LANL. Two samples were tested at each condition. The results were shown in Fig. 9. The yield stress, uniform tensile strength (UTS), uniform elongation ( $e_u$ ) and total elongation ( $e_t$ ) are summarized in Table 3. The unirradiated material has the total elongation ( $e_{tu}$ ) about  $13.70 \pm 1.56\%$ . After irradiation, the total elongation ( $e_{ti}$ ) reduced to  $0.90 \pm 0.21\%$ . The reduction of elongation ( $e_{ti}-e_{tu}$ ) corresponds to the yield strength increasing, from 690 MPa before to  $1072.5 \pm 10.6$  MPa after irradiation. The reduction rate is  $3.3\%/100$  MPa. In unirradiated materials, continuous strain hardening occurred after yielding. The post-necking ductility is about 4.7%. For the irradiated specimens, the engineering stress drop down immediately after reaching the yield point. This behavior is similar to several F/M steels but different to HT9, MA957 [78] and 14YWT [37]. Note that MA957 and 14YWT are both Al-free ODS steels with a high density of ultra-fine oxides ( $\sim 2$  nm) in MA957 [79] and 14YWT [67] that reduced the formation of irradiation induced defects, contrasting with the MA956 with coarser oxide particles presented here.

## 4. Discussion

### 4.1. Hardening model

Based on the statistical results in Table 2, the irradiation hardening caused by defects could be calculated by dispersed barrier hardening (DBH) using either linear method or root-sum-square (RSS) method.

$$\Delta\sigma_y^{linear} = \sum \Delta\sigma_{line, loop, void, oxide} \quad (8)$$

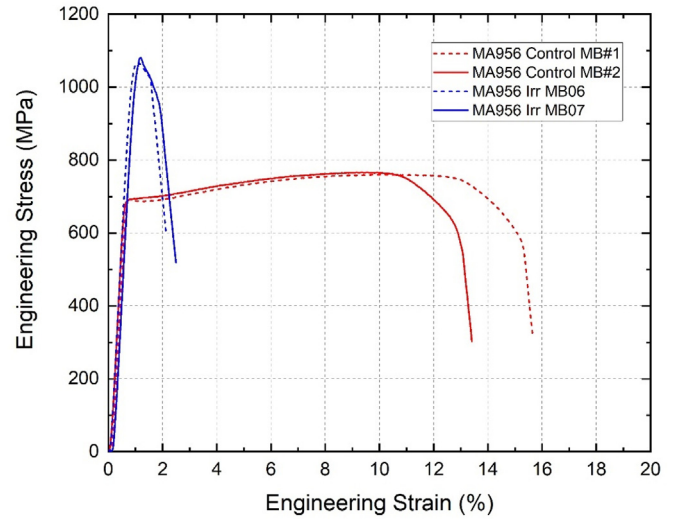
$$\Delta\sigma_y^{root-sum-square} = \sqrt{\sigma_{line}^2 + \sigma_{loop}^2 + \sigma_{void}^2 + \sigma_{oxide}^2} - \sqrt{\sigma_{unirr line, oxide}^2} \quad (9)$$

$$\sigma_{line} = M\alpha\mu b\sqrt{\rho} \quad (10)$$

$$\sigma_{loop, void, oxide} = M\alpha\mu b\sqrt{Nd} \quad (11)$$

**Table 2**  
Summary of defects in MA956 before and after neutron irradiation.

MA956		Average Size(nm)	Density
Irradiated	Dislocation lines	—	$10.58 \pm 0.94 \times 10^{14}/m^2$
	Dislocation loops	$17.6 \pm 10.2$	$7.17 \pm 10^{21}/m^3$
	$\frac{1}{2}\langle 111 \rangle$	$16.6 \pm 8.4$	$3.73 \times 10^{21}/m^3$
	$\langle 100 \rangle$	$18.6 \pm 19.1$	$3.44 \times 10^{21}/m^3$
	Oxides	$17.5 \pm 8.0$	$8.37 \pm 3.73 \times 10^{20}/m^3$
	Voids	$8.7 \pm 2.9$	$1.80 \times 10^{21}/m^3$
Unirradiated	Dislocation lines	—	$2.72 \pm 0.59 \times 10^{14}/m^2$
	Oxides	$20.1 \pm 9.6$	$16.15 \pm 0.25 \times 10^{20}/m^3$



**Fig. 9.** Tensile test results of control (as-received) and irradiated MA956.

$$\Delta\sigma_{loop}^{linear} = \sigma_{111} + \sigma_{100} \quad (12)$$

$$\Delta\sigma_{loop}^{root-sum-square} = \sqrt{\sigma_{111}^2 + \sigma_{100}^2} \quad (13)$$

Where  $\sigma_y$  is the yield strength,  $\Delta\sigma_{line, loop, void, oxide}$  represents the change of yield strength caused by dislocation lines, loops, voids and oxides, respectively.  $M$  is the Taylor factor which equals 3.06 for FCC and BCC materials. The hardening factor,  $\alpha$ , uses different values for different microstructures. The shear modulus  $\mu$  is 82 GPa, for FeCrAl ODS steels. The Burgers vector  $b$  is 0.249 nm for  $\frac{1}{2}\langle 111 \rangle$  moving dislocation lines. The dislocation line density  $\rho$ , number density  $N$  and diameter  $d$  of loops, voids and oxides are listed in Table 2.

As there are two types of dislocation loops,  $\frac{1}{2}\langle 111 \rangle$  and  $\langle 100 \rangle$ , in the irradiated MA956, it is important to know the fraction of each type of loop to apply equations 8–13. From Fig. 3, we can briefly identify the Burgers vectors of dislocation loops. However, as the sample is tilted to achieve a two-beam condition and some of the loops may change their directions due to the surface effect, the loops are not always along the calculated orientation. Therefore, we drew the distribution of the long axis direction of the loops to help analyze Burgers vectors. Note that according to the invisible criterion, the actual density should be two times the  $\frac{1}{2}\langle 111 \rangle$  type loops and 3/2 times the  $\langle 100 \rangle$  type loops in Fig. 3a.

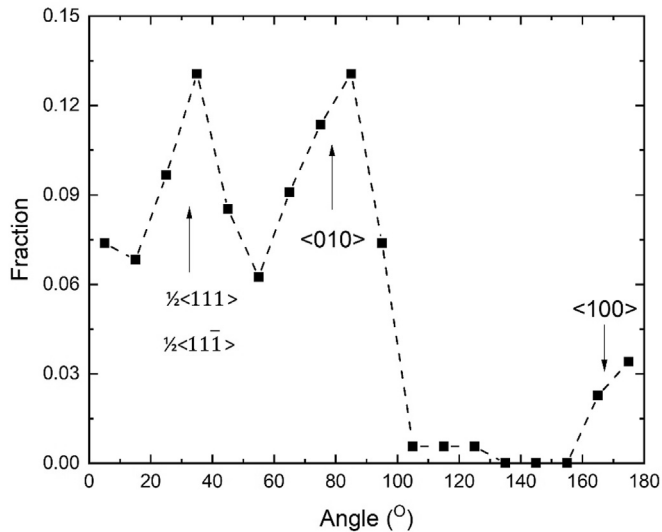
In Fig. 10, the arrows indicate the calculated angle of the dislocation loops with different Burgers vector. There is a small shift of the peak away from the calculated results. This shift is because of specimen tilting, as the calculation supposed the specimen is right on the [001] zone axis. Noticing that the peak of  $\langle 010 \rangle$  type loops is about four times the  $\langle 100 \rangle$  type loops, but physically they should have equal probability. There are also some dislocation loops lie along  $120^\circ$  which violated the invisible criterion. Here we take  $10^\circ \sim 60^\circ$  as the  $\frac{1}{2}\langle 111 \rangle$  type loops and  $0^\circ \sim 10^\circ$ ,  $60^\circ \sim 110^\circ$ ,  $150^\circ \sim 180^\circ$  as the  $\langle 100 \rangle$  type loops. The result shows fraction of  $\frac{1}{2}\langle 111 \rangle$  is 0.45 and fraction of  $\langle 100 \rangle$  is 0.55, which means the total fraction of  $\frac{1}{2}\langle 111 \rangle$  is 0.52 and  $\langle 100 \rangle$  0.48 if taking account for the invisible dislocation loops.

In equation (11), the hardening factor  $\alpha$  for different type loops are considered differently.  $\langle 100 \rangle$  type of dislocation loops has less mobility than  $\frac{1}{2}\langle 111 \rangle$  loops, thus the former should have a higher



**Table 3**  
Summary of tensile test results of MA956.

Type	ID	Yield (MPa)		UTS (MPa)		Uniform elongation (%)		Total elongation (%)		Test temp	Dose (dpa)	Irradiation Temp
Control	MB#1c	690	Avg. 690.0 ± 0.0	760	Avg. 763.0 ± 4.2	9.45	Avg. 9.08 ± 0.52	14.8	Avg. 13.70 ± 1.56	RT	—	—
	MB#2c	690		766		8.72		12.6				
Irradiated	MB06	1065	Avg. 1072.5 ± 10.6	1067	Avg. 1074.0 ± 9.9	0.23	Avg. 0.22 ± 0.01	0.75	Avg. 0.90 ± 0.21	RT	4.36	328 °C
	MB07	1080		1081		0.21		1.05				



**Fig. 10.** The distribution of long axis orientation of projected dislocation loops in Fig. 3. X axis stands for the angle between long axis of loop projection and horizontal line (x positive) and Y is the loop fraction.

hardening factor than the latter. In Refs. [54], the  $\alpha$  value for  $\frac{1}{2}\langle 111 \rangle$  is chosen as 0.17 and  $\langle 100 \rangle$  0.33 in the root-sum-square model while  $\frac{1}{2}\langle 111 \rangle$  0.05 and  $\langle 100 \rangle$  0.31 in linear model. A different set of  $\alpha$  was used in nano-indentation measured hardening [80], where the contribution of cavities was ignored which means the  $\alpha$  for cavities is 0, and the  $\alpha$  is 0.212 for all the other microstructures. In the following calculation, we adopted hardening factor 0.17 and 0.33 for  $\frac{1}{2}\langle 111 \rangle$  and  $\langle 100 \rangle$  loops, and 0.2 for all the others.

The calculated hardening was listed in Table 4. The  $\langle 100 \rangle$  loops and dislocation lines contribute the most to the total increase in yielding stress, with 165 MPa and 200 MPa respectively. The oxides due to shrink in size and decrease in density, have negative contribute to hardening. The calculated results are compared to tensile experiments. The total calculated hardening is 475 MPa by linear method and 234 MPa by RSS method. The underestimation by RSS method is owing to ignorance of  $\alpha'$  precipitates. Although  $\alpha'$  was considered to be weak barriers, the high density of  $\alpha'$  precipitate could yield dominant contribution to the hardening [54]. The compensation of hardening from  $\alpha'$  precipitate is estimated to be

**Table 4**  
Compare of experiment and calculated hardening.

	factor $\alpha$	Cal. $\Delta\sigma$ (MPa)	Cal. Linear (MPa)	Cal. RSS (MPa)	Exp. $\Delta\sigma_y$ (MPa)
$\frac{1}{2}\langle 111 \rangle$ loops	0.17	83.6	475	234	382
$\langle 100 \rangle$ loops	0.33	165			
voids	0.2	49.4			
Dis. lines	0.2	200			
Oxides	0.2	−23.4			
Unknown (fine precipitates)		Estimated 395 MPa			

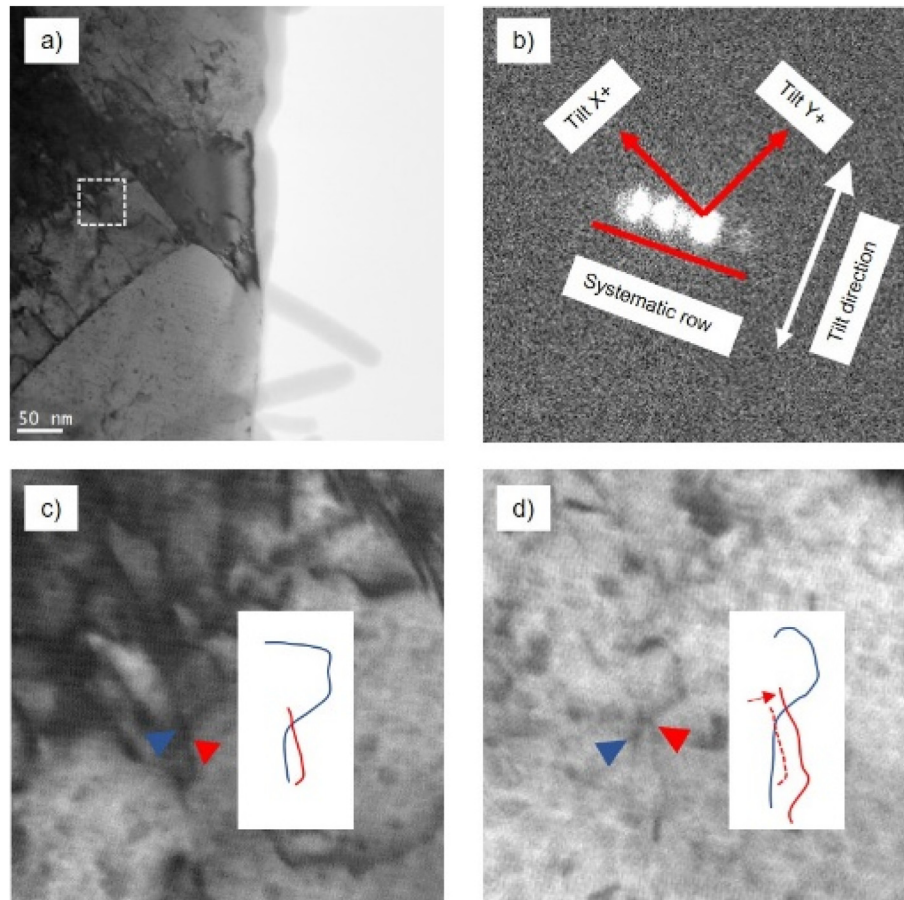
395 MPa. However, it might be overestimated again because of the hardening factor selection for other defects is not fully convincible by current limited data.

The linear method overestimated hardening, as shown in Table 4. This is not a favored method because doubled number density will not yield doubled hardening in equation (11). The linear method only works when a particular type of defects plays dominant role in the hardening. In Table 4, the hardening from  $\langle 100 \rangle$  loops and dislocation lines are pretty similar, resulting in the overestimation of linear-calculated hardening.

#### 4.2. Space distribution of dislocation lines

If a dislocation loop grows and meets a boundary of a grain surface, the loop circuit will open and become a dislocation line. Dislocation lines could act as a biased sink of point defects generated from irradiation. The dislocation line density is  $2.72 \times 10^{14}/\text{m}^2$  in unirradiated materials and  $10.58 \times 10^{14}/\text{m}^2$  after irradiation. The dislocation lines with high density could form into tangles during irradiation, however, it could be just an overlap of images such that the dislocation lines are separated in space indeed. To identify whether the dislocations are tangled, stereo scan of the specimen in TEM was performed. Fig. 11a shows the initial position and the scale bar of this field in STEM mode. The diffraction condition is given in Fig. 11b, where the systematic row was quickly scanned by a side-mount wide angle CCD camera. To keep the images under the same diffraction condition, the tilting was designed to be normal to  $g=(110)$  as much as possible. Nevertheless, while the tilting angle went larger, the Kikuchi line deviated far away from its original position, that the contrast of the dislocations became variable in the tilting direction.

Figs. 11c and d are zoomed ROI of the dashed box in Fig. 11a. The red and blue lines present for the overlapped dislocation respectively. After  $5^\circ$  tilted from Fig. 11c to d, the red line moved to the right and the point of interaction migrated to the upper right consequently. This phenomenon indicates these two lines are not intersected in space. Supplement S1 shows the animation of dislocations while tilting. Each frame in this GIF file was  $+1^\circ$  tilted then reverse displayed. In animation S1, the dislocation lines were found with changing the crossover position during tilting, which means these dislocations are seldom tangled with each other in 3D space.



**Fig. 11.** Changes in position of dislocation overlap respect to different tilt angles: a) initial position and b) the systematic row in STEM mode, c) the zoomed image of the dashed box and d) after 5° tilted. (Full animation is attached in the supplements).

## 5. Conclusions

MA956 is an iron-chromium-aluminum (FeCrAl) based oxide dispersion strengthened (ODS) alloy produced by mechanical alloying. The alloy was irradiated in Advanced Test Reactor (ATR) at 328 °C up to 4.36 dpa with both thermal and fast neutrons. The microstructures of MA956 before and after irradiation were investigated via various TEM techniques.

Dislocation loops were formed during neutron irradiation. The density of dislocation loops was  $7.17 \times 10^{21}/\text{m}^3$ . The mean size of dislocation loops was  $17.6 \pm 10.2$  nm, while the largest ones could reach to over 60 nm. The distribution obeys lognormal fitting. Both  $1/2\langle 111 \rangle$  and  $\langle 100 \rangle$  loops existed in irradiated materials according to the projection in STEM-BF images. Dislocation lines are demonstrated not forming tangles in irradiated specimens by stereo scanning. Voids were formed during irradiation. Argon bubbles were observed in unirradiated materials attached to the surface of oxide particles. The swelling rate was estimated to be 0.08% without subtracting the Ar bubbles.

The oxides showed small shrinkage in size with decrease in number density by 50% after irradiation. Dissolution may have occurred during irradiation. No amorphization was observed. The oxides are Y–Al–O types, with an average diameter of  $20.1 \pm 9.6$  nm before and  $17.5 \pm 8.0$  nm after irradiation.

Tensile test showed hardening and ductility reduction after irradiation. The reduction rate is 3.3%/100 MPa. Hardening was estimated by DBH model using both linear and root-sum-square method. Result showed  $\langle 100 \rangle$  loops and dislocation lines

contributed most to the irradiation hardening.  $\alpha'$  precipitates were considered as contributor to the hardening, even not observed in EFTEM method.

Based on the results we can conclude the oxides in MA956 could maintain crystal structure but partially dissolved in the experiment temperature range. Some suggestions are made here. First, as the swelling and irradiation embrittlement are issues for the irradiation, additional bearing element like Zr may mitigate this problem by refining size of oxides. Second, the high Cr concentration in MA956 have the alpha prime precipitate problem. Reduced Cr is recommended for cladding ferritic ODS steels design.

## Declaration of competing interest

There are no conflicts of interest to declare.

## CRediT authorship contribution statement

**Zhexian Zhang:** Methodology, Software, Validation, Formal analysis, Investigation, Data curation, Writing - original draft, Writing - review & editing. **Tarik A. Saleh:** Conceptualization, Methodology, Validation, Investigation, Resources, Data curation, Writing - review & editing. **Stuart A. Maloy:** Conceptualization, Methodology, Validation, Resources, Writing - review & editing, Supervision, Project administration, Funding acquisition. **Osman Anderoglu:** Conceptualization, Methodology, Validation, Investigation, Writing - original draft, Writing - review & editing.

## Acknowledgement

This research is supported by NEUP (grant No. DE-NE0008656) and NRC faculty development grants (No. NRC-HQ-60-17-G-0007). The authors also acknowledge partial funding from the Advanced Fuels Campaign of the Department of Energy's Nuclear Technology R&D Program. The MA-956 was irradiated in the UCSB ATR-1 irradiations that were carried out as part of the National Scientific Users Facility program.

## Appendix A. Supplementary data

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.jnucmat.2020.152094>.

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