from this we know where the decays are happening.

Assuming the following quantities (TODO add pdg data link oder wikipedia link): [TODO find values] [TODO show gamma, beta claculation...]

$$m_{K^+} =$$

$$m_{pi^+} =$$

$$m_{pi^0} =$$

$$\beta =$$

$$\gamma =$$

In the  $K^+$  rest frame we now have a momentum of 0 and a rest energy  $m_K^+$ . The  $K^+$  now decays into a  $\pi^+$  and a  $\pi^0$  with momenta  $p_{\pi^+}$ ,  $p_{\pi^0}$  in oposing directions. The momenta must have the same magnitude and  $p_{\pi^+} = -p_{\pi^0}$ .

The magnitude of the momentum is calculated with the energy of lost mass in the decay. In the following equations p is used as absolute scalar.

$$E_{\pi^{+}} = \frac{m_{K^{+}}^{2} + m_{pi^{+}}^{2} - m_{pi^{0}}^{2}}{2m_{K^{+}}}$$

$$p_{\pi^{+}} = \sqrt{E_{\pi^{+}}^{2} - m_{pi^{+}}^{2}}$$
(1)

$$p_{\pi^+} = \sqrt{E_{\pi^+}^2 - m_{pi^+}^2} \tag{2}$$

$$p_{\pi^+} = p_{\pi^0} \tag{3}$$

(4)

We use isotropically distributed unit vetors in the simulation (see AW). We multiply them with  $p_{\pi^0}$  to get the momenta for  $\pi^0$  in the  $K^+$  rest frame. Tho get the momenta of  $\pi^+$  we take the same distribution and flip signs. We then build the four vectors  $P_{\pi}^*$ .

$$P_{\pi}^{*} = \begin{bmatrix} E_{\pi}^{*} \\ p_{\pi,x}^{*} \\ p_{\pi,y}^{*} \\ p_{\pi}^{*} z \end{bmatrix} = \begin{bmatrix} \sqrt{m_{\pi}^{2} + ||p_{\pi}||^{2}} \\ p_{\pi,x}^{*} \\ p_{\pi,y}^{*} \\ p_{\pi}^{*} z \end{bmatrix}$$
(5)

Back to the lab rest frame, we boost them with the boost matrix. This boost matrix performs a boost along the z-axis.

$$P_{\pi} = \begin{bmatrix} \gamma & 0 & 0 & \beta \gamma \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ \beta \gamma & 0 & 0 & \gamma \end{bmatrix} \cdot P_{\pi}^{*}$$
 (6)

 $P_{\pi}$  is the four-vector in the lab frame.  $P_{\pi}^*$  is the four-vector in the kaon rest frame.

## 0.1 With scattering

When we use a beam that is not parallel to the z-axis, we need to perform a rotation to the momentum of the four-vector before and after the boost. The rotation after the boost is the same that would rotate  $e_z$  to align with the position vector. the one before is the inverse rotation.

Since the momentum of the four-vectors  $P_{\pi}$  of the pions are isotropically distributed in the  $K^+$  rest frame, we can omit the first rotation and still end up with a an isotropical distribution.

The operation performed with the angle  $\alpha$  around y-axis and  $\beta$  around x-axis:

$$R(\alpha, \beta) = \begin{bmatrix} \cos \alpha & 0 & \sin \alpha \\ 0 & 1 & 0 \\ -\sin \alpha & 0 & \cos \alpha \end{bmatrix} \cdot \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos \beta & -\sin \beta \\ 0 & \sin \beta & \cos \beta \end{bmatrix}$$
(7)

We apply  $R(\alpha, \beta)$  to the momentum after we have extracted the momentum of the four-vector  $P_{\pi}$ .

## 1 Hits

To determine the number of decays for which both pions hit the detector, we first define a distance  $z_d$  from source to detector. And a radius r of the detector. The impulses in the lab frame are  $\overrightarrow{p_{\pi^+}}$  and  $\overrightarrow{p_{\pi^0}}$ . And decay position  $\overrightarrow{x} = (x, y, z)$ . We want to find the number of double hits.

- 1. Drop all the decay entries where  $z > z_d$ .
- 2. Do the following for  $\overrightarrow{p} = \overrightarrow{p_{\pi^+}}$  or  $\overrightarrow{p_{\pi^0}}$ .
  - (a) Calculate s such that  $\overrightarrow{q} = \overrightarrow{x} + s \cdot \overrightarrow{p}$  lands in the plane of  $z_d$ .
  - (b) Drop the pion if it's outside the detector.  $(x+q_x)^2+(y+q_y)^2>r^2$ .
- 3. Drop the decay if at least one of the two pions was dropped.
- 4. Count the number of decays left.

## 2 Determine The optimal detector distance $z_d$

We will now determine the optimal distance such that the acceptance of is maximized. Acceptance is defined as follows:

$$P(z_d) = \frac{\text{n decays for which both pions hit}}{\text{n decays simulated}}$$
(8)

To maximize the acceptance we run the experiment about 500 times with  $N=10^6$  kaons generated. First we simulate a beam aligned to the z axis. Second we generate a beam that scatters at the source in a gaussian distribution with  $\sigma_{\theta}=0.001$  rad.

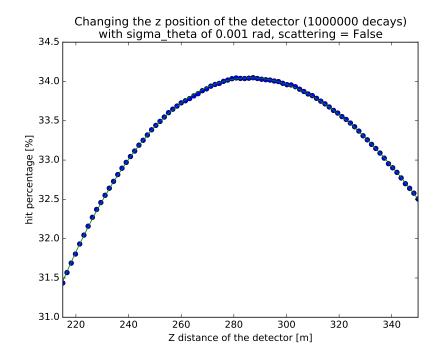


Figure 1: Sample output of decay\_simulation\_6.py

To determine the optimal z distance, we implemented the algorithm described in 1. We let it run for z values in a range from 215 to 350 meters in 85 evenly sized steps. The algorithm generates 85 double hit counts in an array. We divide by N to get the probabilities. Then we fit it with a polynom of degree 3 and determine the maximum of that polynom. To show it graphically, we will show a plot. The simulation was run once with  $N=10^6$  decays.

We run the simulation about 500 times. Fitting the maximum and saving it. The maxima are then put into a histogram. the average and standard deviation of the maximas are listed below. The standard deviation gives a measure of the statistical uncertainty of the simulation.

[TODO add units]

Aligned Beam

$$\langle z \rangle = 285.5$$
  
 $\sigma_z : 0.7$ 

## **Scattered Beam**

$$\langle z \rangle = 280.7$$
  
 $\sigma_z : 0.7$ 

To calculate the final uncertainty we used the relative error on the average decay length from chapter [TODO add reference]. The relative errors get quadratically added.

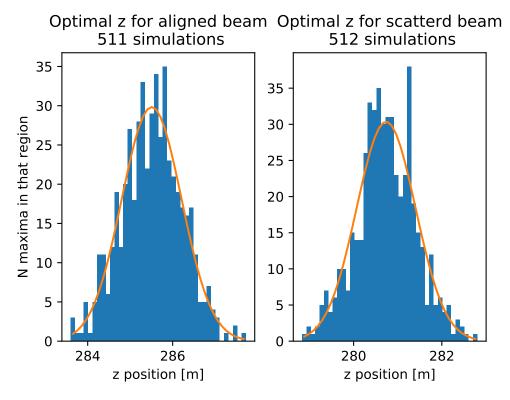


Figure 2: [TODO find good title]

$$m_z = \sigma_z \cdot \sqrt{\left(\frac{\text{uncertanty on adl}}{\text{adl}}\right)^2 + \left(\frac{\sigma_z}{z}\right)^2}$$
 (9)

The final optimal values for z are:

$$z_{
m aligned} = (286 \pm 5) m$$
  
 $z_{
m scattered} = (281 \pm 5) m$ 

The error does not matter that much, since a change of 5 meters around the maximum does maximally change the acceptance by 0.5 percent as seen in figure 1.