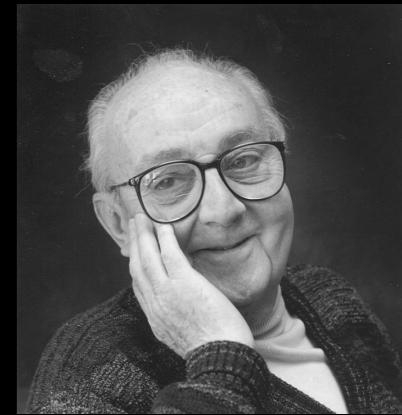


“All models are wrong”

“All models are wrong”

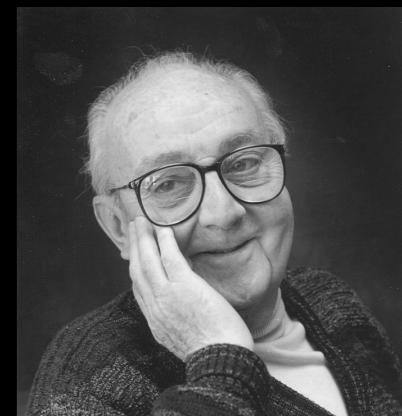
– *George Box, in a 1976 paper published in the Journal of the American Statistical Association*



George Box
1919–2013

“All models are wrong”

– *George Box, in a 1976 paper published in the Journal of the American Statistical Association*



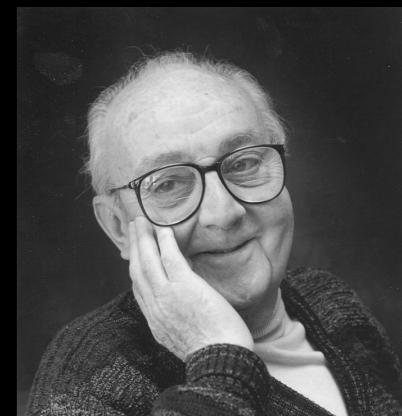
George Box
1919–2013

“All models are wrong, but some are useful”

– *George Box, in the proceedings of a 1978 statistics workshop*

“All models are wrong”

– *George Box, in a 1976 paper published in the Journal of the American Statistical Association*



George Box
1919–2013

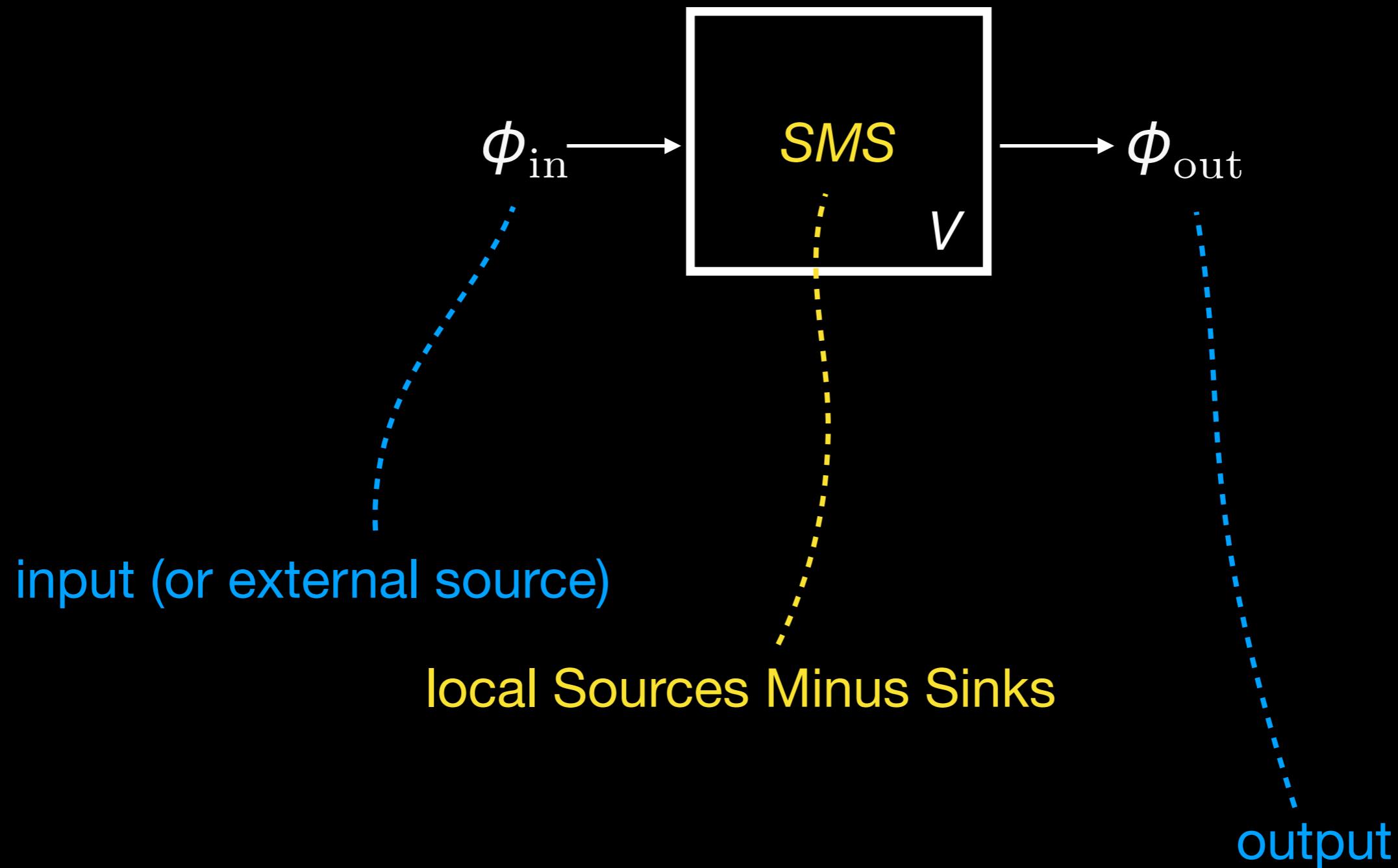
“All models are wrong, but some are useful”

– *George Box, in the proceedings of a 1978 statistics workshop*

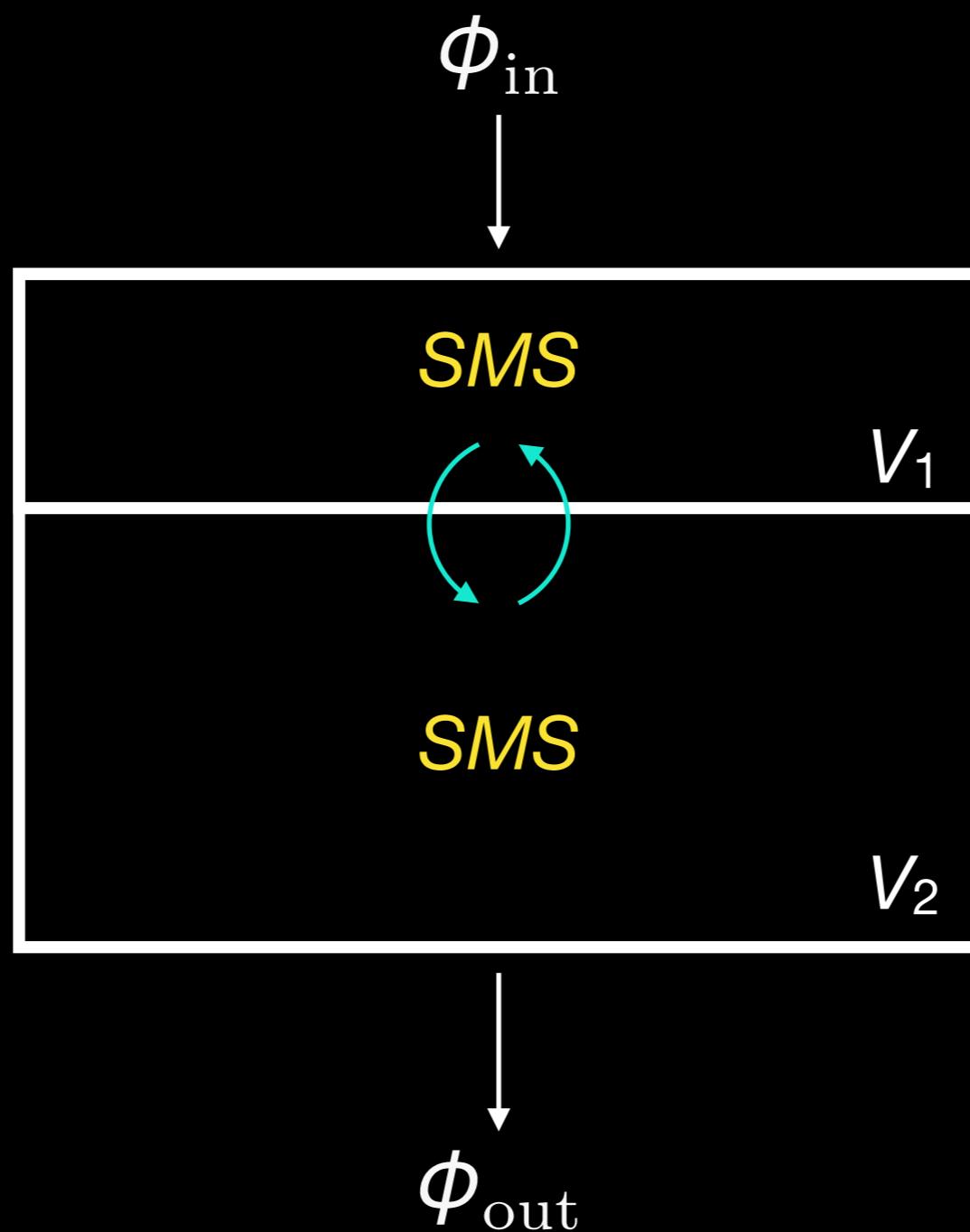
“Box models are useful despite being approximations”

– *Benoit Pasquier, in a 2018 presentation at the Francois Primeau group meeting*

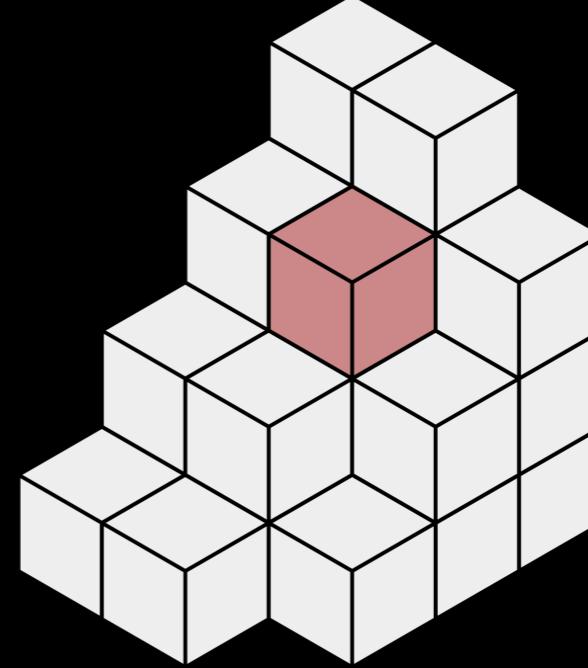
From a single-box model...



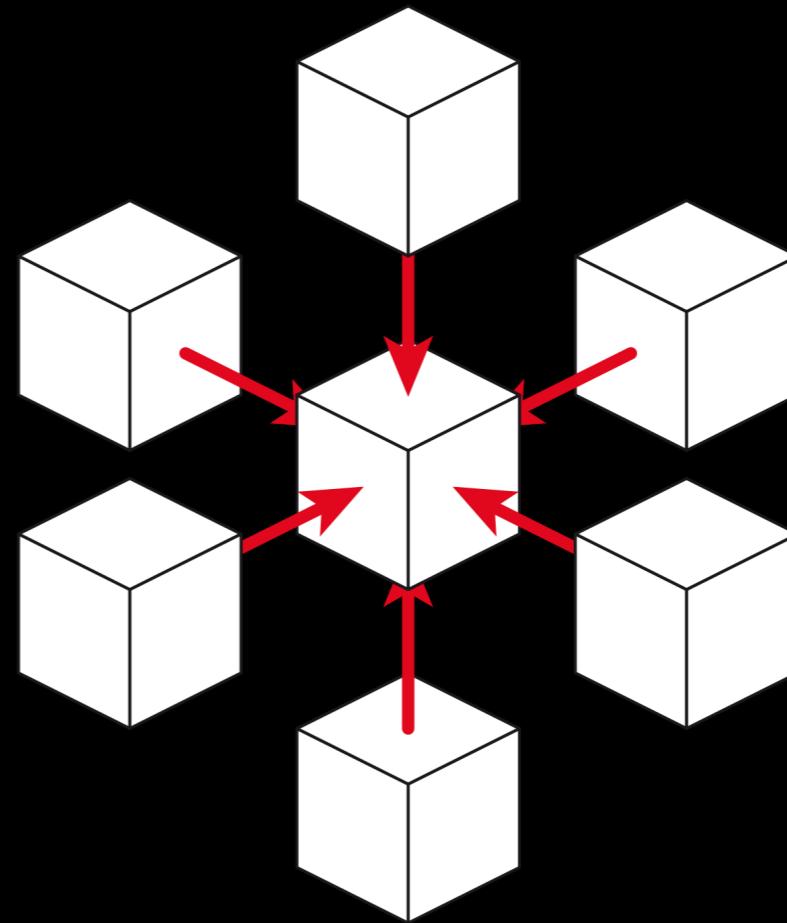
to a two-boxes model...



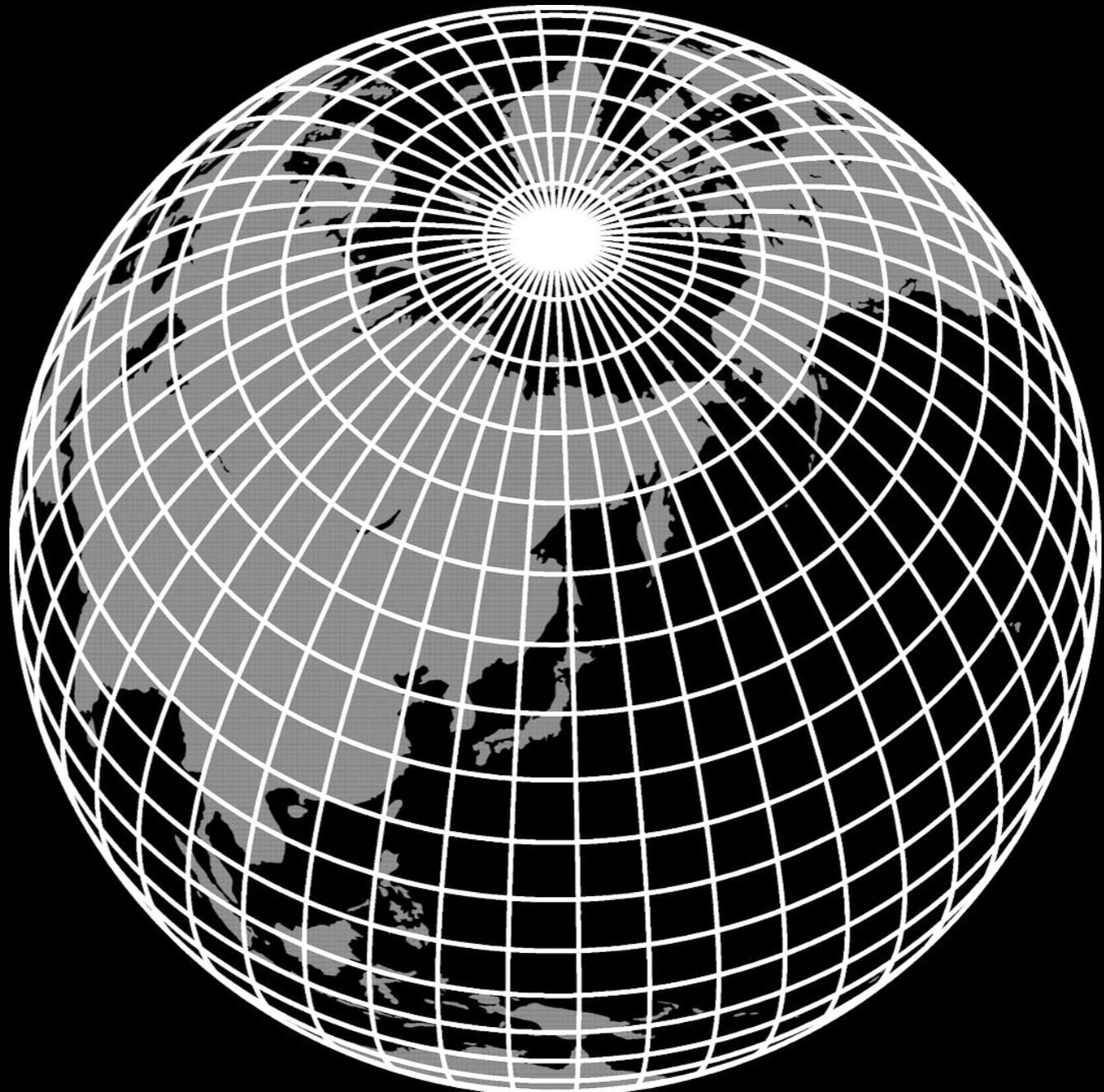
to 3D models with many boxes



where boxes exchange tracers with each other (not only neighbors).



We have global models that cover the entire earth with a grid...



Proof of the existence of such grids can be seen in the photograph below:



More seriously, the awesome OCIM (Ocean Circulation Inverse Model) provide the advective-diffusive transport on such global 3D grids.

More seriously, the awesome OCIM (Ocean Circulation Inverse Model) provide the advective-diffusive transport on such global 3D grids.

The ocean circulation is described by the transport operator, \mathcal{T} :

$$\mathcal{T} = \nabla \cdot \mathbf{u} - \nabla \cdot \mathbf{K} \nabla$$

More seriously, the awesome OCIM (Ocean Circulation Inverse Model) provide the advective-diffusive transport on such global 3D grids.

The ocean circulation is described by the transport operator, \mathcal{T} :

$$\mathcal{T} = \boxed{\nabla \cdot \mathbf{u}} - \boxed{\nabla \cdot \mathbf{K} \nabla}$$

advection diffusion

More seriously, the awesome OCIM (Ocean Circulation Inverse Model) provide the advective-diffusive transport on such global 3D grids.

The ocean circulation is described by the transport operator, \mathcal{T} :

$$\mathcal{T} = \boxed{\nabla \cdot \mathbf{u}} - \boxed{\nabla \cdot \mathbf{K} \nabla}$$

advection diffusion

We extensively use \mathcal{T} in the tracer equation

$$\left(\frac{\partial}{\partial t} + \mathcal{T} \right) \chi = \boxed{\text{SMS}(\chi)}$$

local Sources Minus Sinks

More seriously, the awesome OCIM (Ocean Circulation Inverse Model) provide the advective-diffusive transport on such global 3D grids.

The ocean circulation is described by the transport operator, \mathcal{T} :

$$\mathcal{T} = \boxed{\nabla \cdot \mathbf{u}} - \boxed{\nabla \cdot \mathbf{K} \nabla}$$

advection diffusion

We extensively use \mathcal{T} in the tracer equation

$$\left(\frac{\partial}{\partial t} + \mathcal{T} \right) \chi = \boxed{\text{SMS}(\chi)}$$

local Sources Minus Sinks

OCIM provide a matrix \mathbf{T} that corresponds to the discrete version of \mathcal{T}

$$\mathcal{T} \longleftrightarrow \mathbf{T}$$

More seriously, the awesome OCIM (Ocean Circulation Inverse Model) provide the advective-diffusive transport on such global 3D grids.

The ocean circulation is described by the transport operator, \mathcal{T} :

$$\mathcal{T} = \boxed{\nabla \cdot \mathbf{u}} - \boxed{\nabla \cdot \mathbf{K} \nabla}$$

advection diffusion

We extensively use \mathcal{T} in the tracer equation

$$\left(\frac{\partial}{\partial t} + \mathcal{T} \right) \chi = \boxed{\text{SMS}(\chi)}$$

local Sources Minus Sinks

OCIM provide a matrix \mathbf{T} that corresponds to the discrete version of \mathcal{T}

continuous advection and diffusion

$$\boxed{\mathcal{T}} \longleftrightarrow \boxed{\mathbf{T}}$$

OCIM
transport
matrix

We thus have a discrete version of the tracer equation:

$$\left(\frac{\partial}{\partial t} + \mathbf{T} \right) \mathbf{x} = \mathbf{SMS}(\mathbf{x})$$

which we can rearrange into

$$\frac{\partial \mathbf{x}}{\partial t} = \mathbf{f}(\mathbf{x})$$

where $\mathbf{f}(\mathbf{x}) = -\mathbf{T}\mathbf{x} + \mathbf{SMS}(\mathbf{x})$

Because I am looking for a useful (although wrong) model,
I assume **steady state**, so that I am looking for \mathbf{x} such that

$$\mathbf{f}(\mathbf{x}) = 0$$

Why is OCIM awesome?

$$f(\mathbf{x}) = -\mathbf{T}\mathbf{x} + \mathbf{S}\mathbf{M}\mathbf{S}(\mathbf{x})$$

so if $\mathbf{S}\mathbf{M}\mathbf{S}(\mathbf{x})$ is affine, then f is affine...

af·fine

/ə'fīn, 'afīn/ 

adjective MATHEMATICS

1. allowing for or preserving parallel relationships.

noun ANTHROPOLOGY

1. a relative by marriage.

$$f(\mathbf{x}) = \mathbf{M}\mathbf{x} + \mathbf{b}$$

And if f is affine, then finding $f(\mathbf{x}) = 0$ is very easy:

$$\mathbf{x} = -\mathbf{M}^{-1}\mathbf{b}$$

or

$$\mathbf{x} = -\mathbf{M} \setminus \mathbf{b};$$

in MATLAB

Example: the ideal mean age of water, \mathbf{x} .

$$\mathbf{f}(\mathbf{x}) = \boxed{-\mathbf{T}\mathbf{x}} + \boxed{1} - \boxed{\mathbf{D}_s\mathbf{x}/\tau}$$

advection and diffusion

increasing age (one second per second)

The age at the surface is quickly relaxed to zero
 \mathbf{D}_s is a diagonal matrix of 1's in the surface only
and τ is a fast restoring timescale (e.g., $\tau = 1$ s).

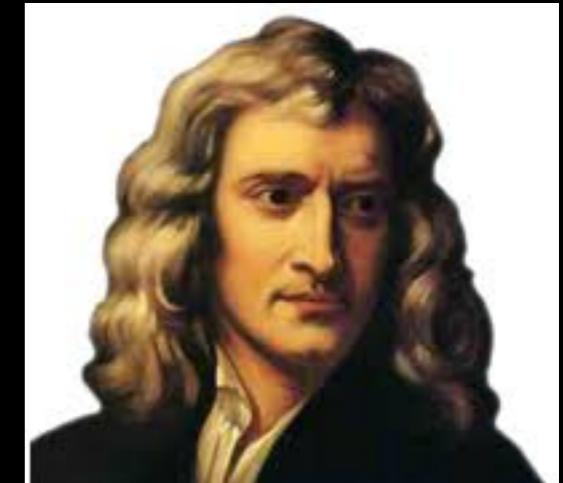
The solution is $\mathbf{x} = -\mathbf{M}^{-1}\mathbf{b}$

where $\mathbf{M} = -\mathbf{T} - \mathbf{D}_s/\tau$

Computation time: a few seconds on my laptop!

Why is OCIM awesome?

$$\mathbf{f}(\mathbf{x}) = -\mathbf{T}\mathbf{x} + \mathbf{SMS}(\mathbf{x})$$

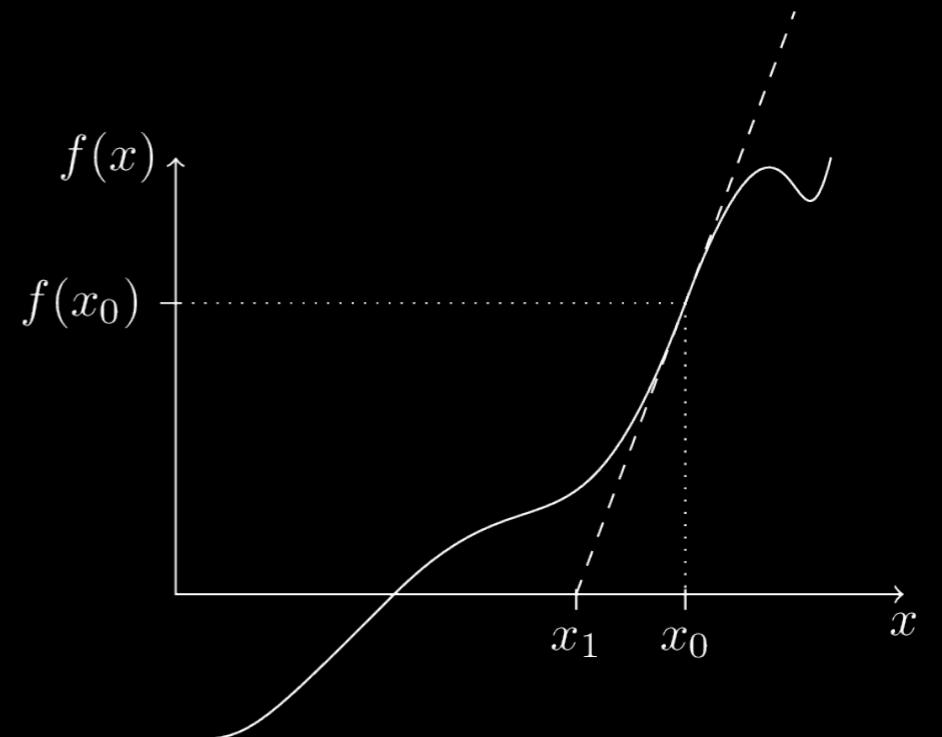


If $\mathbf{SMS}(\mathbf{x})$ is not affine, then \mathbf{f} is not affine either
(most people say that \mathbf{f} is nonlinear - including me)

1643–1727

But we can use **Newton's method!**

$$\mathbf{x}_{n+1} = \mathbf{x}_n - \left[\frac{\partial \mathbf{f}}{\partial \mathbf{x}}(\mathbf{x}_n) \right]^{-1} \mathbf{f}(\mathbf{x}_n)$$



because we want $\mathbf{f}(\mathbf{x}_{n+1}) = 0$ and because

$$\mathbf{f}(\mathbf{x}_{n+1}) \simeq \mathbf{f}(\mathbf{x}_n) + \left[\frac{\partial \mathbf{f}}{\partial \mathbf{x}}(\mathbf{x}_n) \right] (\mathbf{x}_{n+1} - \mathbf{x}_n)$$

Example: the silicate concentration, \mathbf{x} .

$$f(\mathbf{x}) = -T\mathbf{x} + (S - 1)(\mathbf{x} - \mathbf{x}^{\text{obs}})_{\text{eup}}^+ / \tau_{\text{res}} - (\mathbf{x} - \bar{\mathbf{x}}^{\text{obs}}) / \tau_{\text{geo}}$$

Example: the silicate concentration, \mathbf{x} .

$$f(\mathbf{x}) = \boxed{-\mathbf{T}\mathbf{x}} + (S - 1)(\mathbf{x} - \mathbf{x}^{\text{obs}})_{\text{eup}}^+ / \tau_{\text{res}} - (\mathbf{x} - \bar{\mathbf{x}}^{\text{obs}}) / \tau_{\text{geo}}$$

Advection and diffusion

Example: the silicate concentration, \mathbf{x} .

$$f(\mathbf{x}) = \boxed{-\mathbf{T}\mathbf{x}} + \boxed{(S - 1)(\mathbf{x} - \mathbf{x}^{\text{obs}})^+_{\text{eup}} / \tau_{\text{res}}} - (\mathbf{x} - \bar{\mathbf{x}}^{\text{obs}}) / \tau_{\text{geo}}$$

Advection and diffusion

Biological uptake and remineralization at depth,
where S is the particle flux divergence.

\mathbf{x}^{obs} is the observed Si(OH)_4 concentration

$\tau_{\text{res}} \simeq 30$ days

This term reproduces the biological pump mechanism.

Example: the silicate concentration, \mathbf{x} .

$$f(\mathbf{x}) = \boxed{-\mathbf{T}\mathbf{x}} + \boxed{(S - 1)(\mathbf{x} - \mathbf{x}^{\text{obs}})^+_{\text{eup}} / \tau_{\text{res}}} - \boxed{(\mathbf{x} - \bar{\mathbf{x}}^{\text{obs}}) / \tau_{\text{geo}}}$$

Advection and diffusion

Biological uptake and remineralization at depth,
where S is the particle flux divergence.

\mathbf{x}^{obs} is the observed Si(OH)_4 concentration

$\tau_{\text{res}} \simeq 30$ days

This term reproduces the biological pump mechanism.

The geological restoring term,
which slowly restores the global inventory.

$\tau_{\text{geo}} = 10^6$ years

$\bar{\mathbf{x}}^{\text{obs}} \simeq 92$ μM is the mean observed Si(OH)_4 concentration

Example: the silicate concentration, \mathbf{x} .

$$f(\mathbf{x}) = \boxed{-\mathbf{T}\mathbf{x}} + \boxed{(S - 1)(\mathbf{x} - \mathbf{x}^{\text{obs}})^+_{\text{eup}} / \tau_{\text{res}}} - \boxed{(\mathbf{x} - \bar{\mathbf{x}}^{\text{obs}}) / \tau_{\text{geo}}}$$

Advection and diffusion

Biological uptake and remineralization at depth,
where S is the particle flux divergence.

\mathbf{x}^{obs} is the observed Si(OH)_4 concentration

$\tau_{\text{res}} \simeq 30$ days

This term reproduces the biological pump mechanism.

The geological restoring term,
which slowly restores the global inventory.

$\tau_{\text{geo}} = 10^6$ years

$\bar{\mathbf{x}}^{\text{obs}} \simeq 92$ μM is the mean observed Si(OH)_4 concentration

Computation time: a few minutes on my laptop!

Why is OCIM awesome?

Because OCIM affords
very fast computations!

OK... But is it useful for anything in particular?

Yes: OCIM allows
parameter optimization!

Wait a minute... What parameters?

Optimization example: the Si cycle

$$f(\mathbf{x}) = \boxed{-\mathbf{T}\mathbf{x}} + \boxed{(S - 1)(\mathbf{x} - \mathbf{x}^{\text{obs}})_{\text{eup}}^+ / \tau_{\text{res}}} - \boxed{(\mathbf{x} - \bar{\mathbf{x}}^{\text{obs}}) / \tau_{\text{geo}}}$$

Advection and diffusion

Biological pump

Geological restoring

Optimization example: the Si cycle

$$f(\mathbf{x}) = \boxed{-\mathbf{T}\mathbf{x}} + \boxed{(\mathbf{S} - 1)(\mathbf{x} - \mathbf{x}^{\text{obs}})_{\text{eup}}^+ / \tau_{\text{res}}} - \boxed{(\mathbf{x} - \bar{\mathbf{x}}^{\text{obs}}) / \tau_{\text{geo}}}$$

Advection and diffusion Biological pump Geological restoring

There are some **optimizable parameters**, $\mathbf{p} = (p_1, \dots, p_n)$.
So now $f(\mathbf{p}, \mathbf{x})$ depends on those parameters.

Optimization example: the Si cycle

$$f(\mathbf{x}) = [-\mathbf{T}\mathbf{x}] + [(\mathbf{S} - 1)(\mathbf{x} - \mathbf{x}^{\text{obs}})_{\text{eup}}^+ / \tau_{\text{res}}] - [\mathbf{x} - \bar{\mathbf{x}}^{\text{obs}}] / \tau_{\text{geo}}$$

Advection and diffusion Biological pump Geological restoring

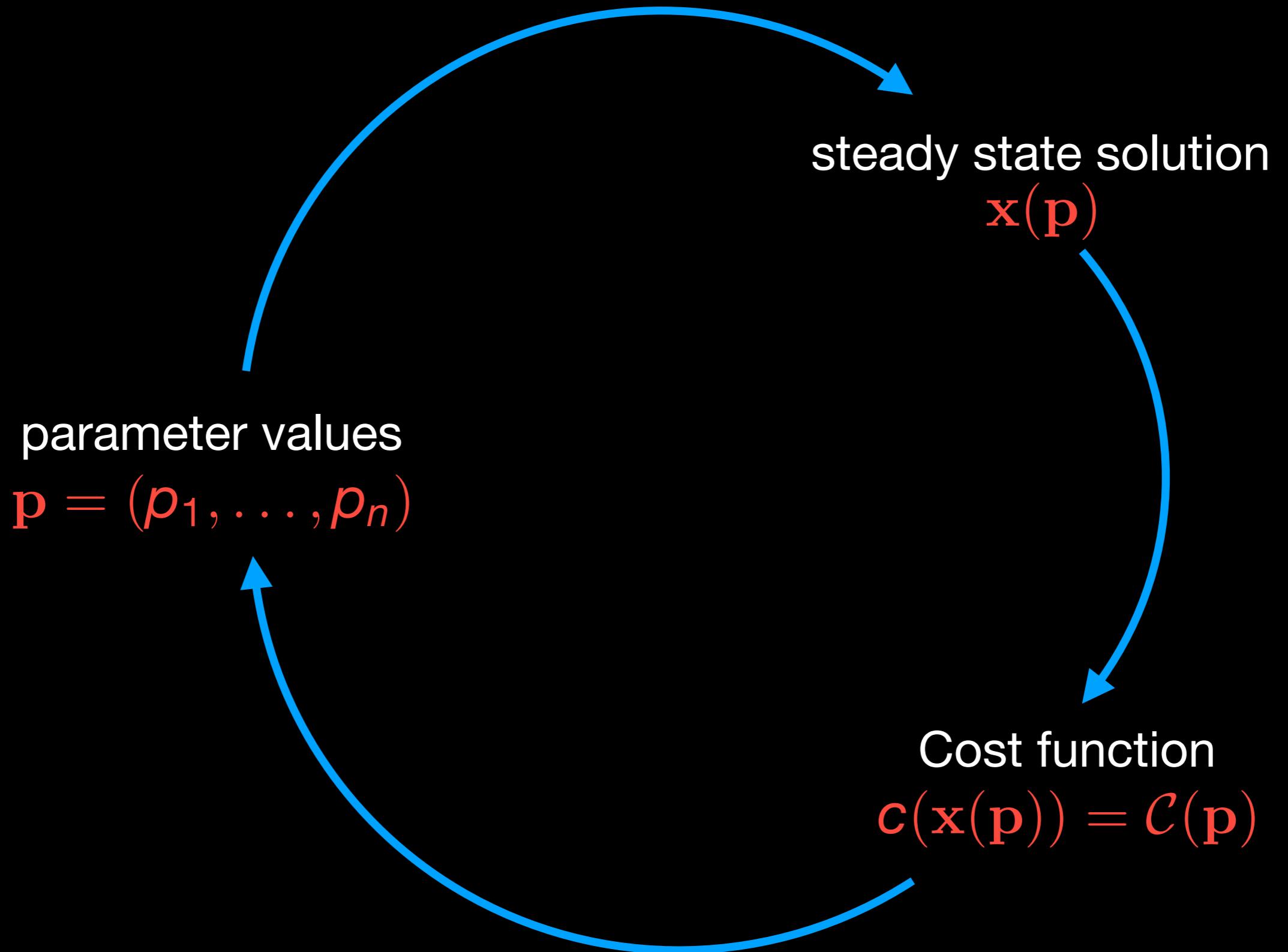
There are some **optimizable parameters**, $\mathbf{p} = (p_1, \dots, p_n)$.
So now $f(\mathbf{p}, \mathbf{x})$ depends on those parameters.

We want to minimize the error of our estimate when we compare it to observations, so we build a **cost function** that represents this error by a scalar:

$$c(\mathbf{x}) = \delta\mathbf{x}^\top \mathbf{V} \delta\mathbf{x} \longleftrightarrow \int d^3r (\chi^{\text{mod}} - \chi^{\text{obs}})^2$$

where $\delta\mathbf{x} = \mathbf{x} - \mathbf{x}^{\text{obs}}$ is the mismatch with observations, and \mathbf{V} is a diagonal matrix of the volumes of the grid boxes.

Optimization example: the Si cycle



Optimization example: the Si cycle

Find x such that $f(p, x) = 0$

e.g., using Newton's method

(`nsold.m`)

parameter values

$p = (p_1, \dots, p_n)$

steady state solution

$x(p)$

Cost function
 $c(x(p)) = \mathcal{C}(p)$

Optimization example: the Si cycle

Find x such that $f(p, x) = 0$

e.g., using Newton's method

(`nsold.m`)

parameter values
 $p = (p_1, \dots, p_n)$

steady state solution
 $x(p)$

Compare to observations

Cost function
 $c(x(p)) = \mathcal{C}(p)$

Optimization example: the Si cycle

Find x such that $f(p, x) = 0$

e.g., using Newton's method

(`nsold.m`)

parameter values

$p = (p_1, \dots, p_n)$

Optimize parameters
(e.g., `fminunc.m`)

steady state solution

$x(p)$

Compare to observations

Cost function
 $c(x(p)) = \mathcal{C}(p)$

Optimization example: the Si cycle

Find x such that $f(p, x) = 0$

e.g., using Newton's method

(`nsold.m`)

parameter values
 $p = (p_1, \dots, p_n)$

Optimize parameters
(e.g., `fminunc.m`)

steady state solution
 $x(p)$

Compare to observations

Cost function
 $c(x(p)) = \mathcal{C}(p)$

Optimization example: the Si cycle

Find x such that $f(p, x) = 0$

e.g., using Newton's method

(`nsold.m`)

parameter values

$p = (p_1, \dots, p_n)$

Optimize parameters
(e.g., `fminunc.m`)

This estimate depends
on the parameters

steady state solution

$\boxed{x(p)}$

Compare to observations

Cost function
 $c(x(p)) = \mathcal{C}(p)$

Optimization example: the Si cycle

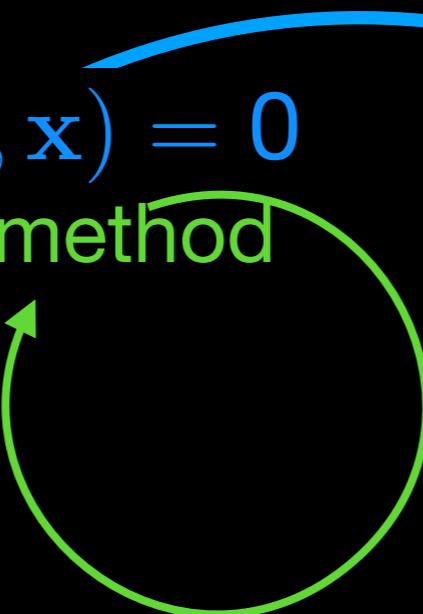
Find x such that $f(p, x) = 0$

e.g., using Newton's method

(`nsold.m`)

parameter values
 $p = (p_1, \dots, p_n)$

Optimize parameters
(e.g., `fminunc.m`)



This estimate depends
on the parameters

steady state solution

$x(p)$

Compare to observations

Cost function

$c(x(p)) = C(p)$

This is the cost function
used by the optimizer

Some notes on Newton's method

Newton's method requires the computation of the Jacobian of f with respect to x .

This can be done in many ways:

Analytically:

- that is, write the code for the Jacobian (fast and accurate)

numerically:

- using finite differences (slow and inaccurate)
- using the complex step method (fast, accurate)
- using algorithmic differentiation (fast, accurate)

Finite differences, CSD, and AD, example with Jacobian

Finite differences (what MATLAB uses internally)

$$\frac{\partial \mathbf{f}}{\partial \mathbf{x}} \simeq \frac{\mathbf{f}(\mathbf{x} + \varepsilon \mathbf{I}) - \mathbf{x}}{\varepsilon} \quad (\text{Abuse of notation here})$$

Complex Step Differentiation

$$\frac{\partial \mathbf{f}}{\partial \mathbf{x}} \simeq \frac{\Im[\mathbf{f}(\mathbf{x} + i\varepsilon \mathbf{I})]}{\varepsilon}$$

Algorithmic Differentiation (A.K.A. automatic differentiation)

No formula, but calculates the derivatives at every line of code that defines \mathbf{f} (using the composition rule recursively).

Some notes on the optimization

The optimization of the parameters, \mathbf{p} (e.g, using MATLAB's `fminunc`), can go faster if you provide (I think):

- the Jacobian of \mathcal{C} with respect to \mathbf{p}
- the Hessian of \mathcal{C} with respect to \mathbf{p}

This can also be done (I think) in the ways mentioned before:

Analytically:

- that is, write the code for the Jacobian (fast and accurate)

numerically:

- using finite differences (slow and inaccurate)
- using the complex step method (fast, accurate)
- using algorithmic differentiation (fast, accurate)

Some more notes on the optimization

One can derive the **Jacobian** from the steady state equation...

$$f(p, x(p)) = 0 \implies \frac{\partial f}{\partial p} + \frac{\partial f}{\partial x} \frac{dx}{dp} = 0$$

Use the composition rule:

$$\frac{dC}{dp} = \frac{d}{dp} [c(x(p))] = \frac{dc}{dx} \frac{dx}{dp}$$

And inject in the final result:

$$\frac{dC}{dp} = - \frac{dc}{dx} \frac{\partial f^{-1}}{\partial x} \frac{\partial f}{\partial p}$$

And even more notes on the optimization

One can derive the Hessian from the steady state solution too!
It is a bit more complicated, especially notation-wise... So this is
merely an attempt to tackle it, and François may be able to explain it to us!

The composition rule for Hessian can be written as

$$\frac{d^2C}{dp^2} = \frac{d^2C}{dx^2} \left(\frac{dx}{dp} \otimes \frac{dx}{dp} \right) + \frac{dc}{dx} \boxed{\frac{d^2x}{dp^2}}$$

This one might
be tough

$$\frac{d^2C}{dx^2} = 2V$$

Already explained
(easy!)

$$\frac{dc}{dx} = 2(x - x^{\text{obs}})^T V$$

And even more notes on the optimization (continued)

$\frac{d^2\mathbf{x}}{dp^2}$ size: $n_x \times n_p \times n_p$
can be found from something like this

$$f(p, \mathbf{x}(p)) = 0 \implies \frac{\partial^2 f}{\partial p^2} + \frac{\partial^2 f}{\partial \mathbf{x}^2} \left(\frac{d\mathbf{x}}{dp} \otimes \frac{d\mathbf{x}}{dp} \right) + \frac{\partial f}{\partial \mathbf{x}} \frac{d^2 \mathbf{x}}{dp^2} = 0$$

Could be full but smallish

$$n_x \times n_p \times n_p$$

Already explained
(easy!)

↑
Jacobian
for Newton

Is huge but should be sparse!

$$n_x \times n_x \times n_x$$

(would be zero anyway
in my example model)

But be reassured...

You do not have to go the analytical way!

fortuitous pun



Although it can probably save you a lot of time in the long run.

You can use the usual **finite differences** and let MATLAB's built-in functions (**fminsearch**, **fminunc**, etc.) do all the work.

You can use the **complex step differentiation (CSD)**, but there is a bit of work and some caveats (but I love it - much elegance).

And you can probably use the **algorithmic differentiation (AD)**, but I have not tried it myself, therefore I cannot recommend it.