

Lec 2

Ex. Let's look at what these ~~q~~ quantities are in thermal equilib.
(We haven't defined that, yet, but we'll see what it implies)

\bar{I} is uniform (no r depend.)
isotropic (no \hat{n} depend.)
~~time~~ stationary (no t -depend)

Given by Planck function

$$I_\nu(T) = \frac{2h\nu^3}{c^2} \left(e^{\frac{h\nu}{kT} - 1} \right)^{-1} \equiv B_\nu(T)$$

$$J_\nu = B_\nu(T)$$

→ Monochromatic energy density is

$$\mathcal{E}_R(\nu) = \frac{4\pi}{c} J = \frac{4\pi}{c} B_\nu$$

→ Total energy density $\int d\nu \mathcal{E}_R(\nu)$

$$\mathcal{E} = \frac{8\pi h}{c^3} \int_0^\infty \left(e^{\frac{h\nu}{kT} - 1} \right)^{-1} \nu^3 d\nu = a T^4$$

$$a = \frac{8\pi^5 k^4}{15c^3 h^3} \equiv \text{rad. constant}$$

In spherical sym.

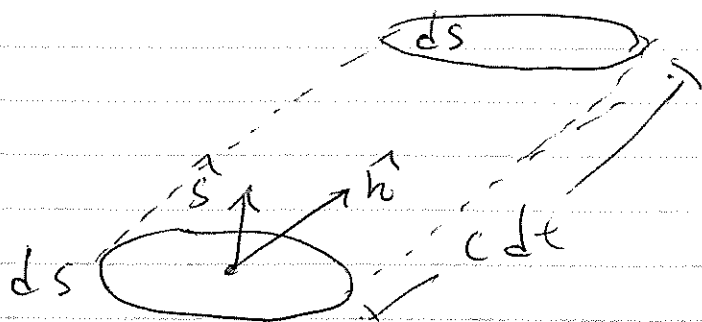
$$\bar{J}(r, \nu, t) = \frac{1}{2} \int_{-1}^1 I(r, \mu, \nu, t) d\mu$$

The mean intensity is closely related to the energy density of radiation.

Consider a volume V . What is the energy contained in V ?

What is the energy contained in volume dV ?

Since $v=c$ for photons, they will cover a distance $c \cdot dt$ in time dt , so let's make up our volume by moving a surface in the direction of propagation.



Then

$$dV = c \, dt \cdot \hat{n} \cdot d\mathbf{s} \\ = c \, dt \, ds \cdot \cos \theta$$

Since

$$dE = I_\nu \, ds \cos \theta \, d\Omega \, d\nu \, dt \\ = \frac{I_\nu}{c} \, d\Omega \, d\nu \, dV$$

First Moment: ^{Energy} Flux

$$\vec{F} = \int I \hat{n} d\Omega$$

Is this the flux as defined earlier?

$$\vec{F} = \frac{dE}{dt ds dv} \hat{n} \quad [F] = \frac{\text{erg}}{\text{s cm}^2 \text{Hz}}$$

Recall: $dE = I_v ds \cdot \cos\theta d\Omega dv dt$

And $= I_v \cdot \hat{n} \cdot d\vec{s} d\Omega dv dt$

$$F = \frac{dE}{dt ds dv} \Rightarrow \int I \hat{n} \cdot d\Omega = \vec{F}$$

All good!

Example: An isotropic rad. field is given by $I(\hat{n}) = I_0$ [same in every direction]

What is $F = ?$ Solve in spherical coords.

$$\begin{aligned} F &= \int I_0 \hat{n} d\Omega = I_0 \int \cos\theta \sin\theta d\theta d\phi \\ &= I_0 \int_0^{2\pi} \int_0^\pi \cos\theta \sin\theta d\theta d\phi = 0 \end{aligned}$$

But, how can this be? In the centers of stars, I should be isotropic

How can there be flux from a star?
Shouldn't there be a net outward flux?

Example: I is nearly isotropic, but not completely

What is F, J , if the rad. field has small deviation from isotropy that is connected to a variable called optical depth, as

$$I_\nu(\tau, \mu) = a_\nu(\tau) + b_\nu(\tau)\mu$$

$$J = \frac{1}{2} \int_{-1}^1 I d\mu = a$$

$$\bar{F} = 2\pi \int_{-1}^1 I \mu d\mu = 4\pi \frac{b}{3}$$

In the \vec{r} direction because τ changes with r more on that sector

Notice that we defined J and F differently.

Why does one have $\frac{1}{4\pi}$ but not the other?

$$J = \frac{1}{4\pi} \int I d\Omega, \quad \vec{F} = \int I \hat{n} d\Omega$$

We can define the "Eddington flux"

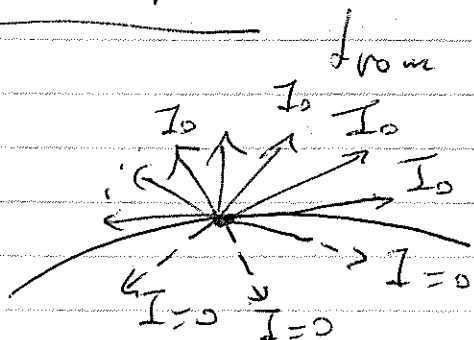
$$\vec{H} = \frac{1}{4\pi} \int I \vec{n} d\Omega$$

We'll use both \vec{F} (the astrophysical flux) and \vec{H} .

In the above example

$$H = \frac{1}{2} \int_{-1}^1 I \mu d\mu = \frac{b}{3}$$

Example: Consider radiation leaving from a boundary, isotropically



All outgoing rays have I_0 . All inbound rays have $I=0$. "no radiation is coming into" the domain from that boundary.

What is H ?

$$\begin{aligned} H &= \frac{1}{2} \int_{-1}^1 I \mu d\mu = \frac{1}{2} \int_0^1 I \mu d\mu = \\ &= \frac{1}{2} I_0 \int_0^1 d\left(\frac{\mu^2}{2}\right) = \frac{1}{4} I_0 \end{aligned}$$

$$O. F = 4\pi r^2 H = \pi I_0$$

So, this is the case pertaining to the earlier problem with isotropic rad. all around: for an object of finite size radiating luminosity L "isotropically", on the surface $r=R$, there is no inward specific intensity, and at $r=R$, $F = \pi I$

For $r > R$ "free-streaming" starts, I stays constant, but $F \sim \frac{1}{r^2}$.

Example. Calculate, I, H for a "free-streaming" rad field (specific intensity in a direction)

$$I(\mu) = I_0 \delta(\mu - \mu_0)$$

$$J = \frac{1}{2} \int I_0 \delta(p - p_0) dp = \frac{I_0}{2}$$

$$H = \frac{1}{2} \int I_0 p \delta(p - p_0) = \frac{p_0 I_0}{2}$$

Second moment: Radiation Pressure

J is scalar, \vec{H}/E are vectors ^{Tensor}

P_{ij} is a tensor!

$$P(\vec{r}, \nu, t) = \frac{1}{c} \int I(\vec{r}, \hat{n}, \nu, t) \hat{n} \hat{n} d\Omega$$

or $P_{ij} = \frac{1}{c} \int I \cdot n_i n_j d\Omega$

al: P_{ii} are related to energy density ^[Lecture 3]

eg. for a fluid $\int \rho v_i v_j d\Omega \Rightarrow$
~~radiation~~ energy density / pressure

For rad. the trace of P_{ij}

$$\sum P_{ii} = P_{xx} + P_{yy} + P_{zz} =$$

$$= \frac{1}{c} \int I (\underbrace{n_x^2 + n_y^2 + n_z^2}) d\Omega$$