

Approximation of Integrals near

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Superconducting state

$$\delta\chi_{\parallel}(\mathbf{q}) = -\frac{1}{\chi_0} \sum_{\mathbf{k},s} \frac{(f(\epsilon_{k-s}) - f(\epsilon_{k+s})) (u_{k+} u_{k-} + v_{k+} v_{k-})^2}{\epsilon_{k-s} - \epsilon_{k+s}} - \frac{(1 - f(\epsilon_{k-s}) - f(\epsilon_{k+s})) (u_{k+} v_{k-} - v_{k+} u_{k-})^2}{\epsilon_{k-s} + \epsilon_{k+s}} - \frac{f(\xi_{k-s}) - f(\xi_{k+s})}{\xi_{k-s} - \xi_{k+s}} \quad (1)$$

$$\delta\chi_{\perp}(\mathbf{q}) = -\frac{1}{\chi_0} \sum_{\mathbf{k},s} \frac{(f(\epsilon_{k-s}) - f(\epsilon_{k+s})) (u_{k+} u_{k-} + v_{k+} v_{k-})^2}{\epsilon_{k-s} - \epsilon_{k+s}} - \frac{(1 - f(\epsilon_{k-s}) - f(\epsilon_{k+s})) (u_{k+} v_{k-} - v_{k+} u_{k-})^2}{\epsilon_{k-s} + \epsilon_{k+s}} - \frac{f(\xi_{k-s}) - f(\xi_{k+s})}{\xi_{k-s} - \xi_{k+s}} \quad (2)$$

We make the following assumptions: $\xi_{k\pm} = \pm v_f k_x, \quad \Delta_{k\pm} = \pm v_d k_y$.

The excitation energies are then equal $\epsilon_{k+} = \epsilon_{k-} = -\sqrt{v_f^2 k_x^2 + v_d^2 k_y^2}$. Now we can plug these into the formula for the u 's and v 's.:

$$u_{k\pm} = \text{sgn}(\Delta_{k\pm}) \sqrt{\frac{1}{2} (1 + \xi_{k\pm}/\epsilon_{k\pm})} = \text{sgn}(\Delta_{k\pm}) \sqrt{\frac{1}{2} (1 \pm v_f k_x / \sqrt{v_f^2 k_x^2 + v_d^2 k_y^2})}$$

$$v_{k\pm} = \sqrt{\frac{1}{2} (1 - \xi_{k\pm}/\epsilon_{k\pm})} = \sqrt{\frac{1}{2} (1 \mp v_f k_x / \sqrt{v_f^2 k_x^2 + v_d^2 k_y^2})}$$

Now we can see that $u_{k+} = \text{sgn}(\Delta_{k+}) v_{k-}$ and $u_{k-} = \text{sgn}(\Delta_{k-}) v_{k+}$. Plugging these into the formula for $\delta\chi$ causes the first term to be zero: $u_{k+} u_{k-} + v_{k+} v_{k-} = \text{sgn}(\Delta_{k+}) \text{sgn}(\Delta_{k-}) v_{k+} v_{k-} + v_{k+} v_{k-} = 0$.

The second term, however survives:

$$(u_{k+} v_{k-} - v_{k+} u_{k-})^2 = (\text{sgn}(\Delta_{k+}) v_{k-}^2 - \text{sgn}(\Delta_{k-}) v_{k+}^2)^2 = (v_{k-}^2 + v_{k+}^2)^2 = 1$$

So

$$\delta\chi_{\parallel}(\mathbf{q}) = -\frac{1}{\chi_0} \sum_{\mathbf{k},s} -\frac{1 - f(\epsilon_{k-s}) - f(\epsilon_{k+s})}{2\epsilon_{k+}} - \frac{f(\xi_{k-s}) - f(\xi_{k+s})}{kq \cos(\theta)} \quad (3)$$

$$\delta\chi_{\perp}(\mathbf{q}) = -\frac{1}{\chi_0} \sum_{\mathbf{k},s} -\frac{1 - f(\epsilon_{k-s}) - f(\epsilon_{k+s})}{2\epsilon_{k-} + 2s\mu_e H} - \frac{f(\xi_{k-s}) - f(\xi_{k+s})}{kq \cos(\theta) + 2s\mu_e H} \quad (4)$$

Now we need to find the regions of integration for both the normal state component and the superconducting component. The function $1 - f(\epsilon_{k-}) - f(\epsilon_{k+})$ which appears in the superconducting term is approximately equal to 1 in the region near the intersections of the fermi surfaces Γ_{k+} and Γ_{k-} . We take advantage of this by switching to an elliptic coordinate system and cut off the integration at some energy Λ . Now we can write the sums as integrals, using a normalized area of integration ($A/(2\pi\hbar)^2 = 1$). The superconducting terms are:

$$\sum_{\mathbf{k},s} \frac{1}{2\epsilon_{k+}} = \int dk_x dk_y \frac{1}{\sqrt{v_f^2 k_x^2 + v_d^2 k_y^2}} = \frac{1}{v_d v_f} \int dx dy \frac{1}{\sqrt{x^2 + y^2}} = \frac{2\pi}{v_d v_f} \int_0^{\sqrt{\Lambda}} dr \frac{r}{\sqrt{r^2}} \quad (5)$$

$$= \frac{2\pi\sqrt{\Lambda}}{v_d v_f} \quad (6)$$

$$\sum_{\mathbf{k},s} \frac{1}{2\epsilon_{k-} + 2s\mu_e H} = \frac{1}{2} \sum_s \int dk_x dk_y \frac{1}{\sqrt{v_f^2 k_x^2 + v_d^2 k_y^2 + s\mu_e H}} = \frac{1}{2v_d v_f} \sum_s \int dx dy \frac{1}{\sqrt{x^2 + y^2 + s\mu_e H}} \quad (7)$$

$$= \frac{\pi}{v_d v_f} \sum_s \int_0^{\sqrt{\Lambda}} dr \frac{r}{r + s\mu_e H} = \frac{\pi}{v_d v_f} \sum_s \left[r - s\mu_e H \ln |r + s\mu_e H| \right]_{r=0}^{\sqrt{\Lambda}} \quad (8)$$

$$= \frac{\pi}{v_d v_f} \sum_s \sqrt{\Lambda} - s\mu_e H \ln |\sqrt{\Lambda}/(s\mu_e H) + 1| \quad (9)$$

$$= \frac{2\pi\sqrt{\Lambda}}{v_d v_f} - \frac{\pi}{v_d v_f} \mu_e H \ln \left| \frac{\sqrt{\Lambda} + \mu_e H}{-\sqrt{\Lambda} + \mu_e H} \right| \quad (10)$$

Now we calculate the normal state components:

$$\sum_{\mathbf{k},s} \frac{f(\xi_{k-s}) - f(\xi_{k+s})}{kq \cos(\theta)} = \sum_s \int dk_x dk_y \frac{f(\xi_{k-s}) - f(\xi_{k+s})}{k_x q} = 4 \sum_s \int_0^{\sqrt{\Lambda}} dk_x \int_0^{\alpha_{\parallel s} k_x} dk_y \frac{1}{k_x q} \quad (11)$$

$$= 4 \sum_s \frac{\alpha_{\parallel s} \sqrt{\Lambda}}{q} \quad (12)$$

$$\sum_{\mathbf{k},s} \frac{f(\xi_{k-s}) - f(\xi_{k+\bar{s}})}{kq \cos(\theta) + 2s\mu_e H} = \sum_s \int dk_x dk_y \frac{f(\xi_{k-s}) - f(\xi_{k+\bar{s}})}{k_x q + 2s\mu_e H} = 4 \sum_s \int_0^{\sqrt{\Lambda}} dk_x \int_0^{\alpha_{\perp} k_x} dk_y \frac{1}{k_x q + 2s\mu_e H} \quad (13)$$

$$= (4\alpha_{\perp}/q) \sum_s \sqrt{\Lambda} - (2s\mu_e H/q) \ln |\sqrt{\Lambda}/(2s\mu_e H/q) + 1| \quad (14)$$

$$= (8\alpha_{\perp} \sqrt{\Lambda}/q) - (8\mu_e H \alpha_{\perp}/q^2) \ln \left| \frac{\sqrt{\Lambda} + (2\mu_e H/q)}{-\sqrt{\Lambda} + (2\mu_e H/q)} \right| \quad (15)$$