

# Observational determination of surface radiative forcing by CO<sub>2</sub> from 2000 to 2010

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The climatic impact of CO<sub>2</sub> and other greenhouse gases is usually quantified in terms of radiative forcing<sup>1</sup>, calculated as the difference between estimates of the Earth's radiation field from pre-industrial and present-day concentrations of these gases. Radiative transfer models calculate that the increase in CO<sub>2</sub> since 1750 corresponds to a global annual-mean radiative forcing at the tropopause of  $1.82 \pm 0.19 \text{ W m}^{-2}$  (ref. 2). However, despite widespread scientific discussion and modelling of the climate impacts of well-mixed greenhouse gases, there is little direct observational evidence of the radiative impact of increasing atmospheric CO<sub>2</sub>. Here we present observationally based evidence of clear-sky CO<sub>2</sub> surface radiative forcing that is directly attributable to the increase, between 2000 and 2010, of 22 parts per million atmospheric CO<sub>2</sub>. The time series of this forcing at the two locations—the Southern Great Plains and the North Slope of Alaska—are derived from Atmospheric Emitted Radiance Interferometer spectra<sup>3</sup> together with ancillary measurements and thoroughly corroborated radiative transfer calculations<sup>4</sup>. The time series both show statistically significant trends of  $0.2 \text{ W m}^{-2}$  per decade (with respective uncertainties of  $\pm 0.06 \text{ W m}^{-2}$  per decade and  $\pm 0.07 \text{ W m}^{-2}$  per decade) and have seasonal ranges of  $0.1\text{--}0.2 \text{ W m}^{-2}$ . This is approximately ten per cent of the trend in downwelling longwave radiation<sup>5,6</sup>. These results confirm theoretical predictions of the atmospheric greenhouse effect due to anthropogenic emissions, and provide empirical evidence of how rising CO<sub>2</sub> levels, mediated by temporal variations due to photosynthesis and respiration, are affecting the surface energy balance.

Even though Northern Hemisphere atmospheric CO<sub>2</sub> mixing ratios have recently exceeded 400 parts per million (ppm), few investigations have directly explored the effect of changes in well-mixed greenhouse gases on spectral and broadband radiation fluxes. In fact, doing so has proved difficult. Using data from two satellite instruments launched 26 years apart, Harries *et al.*<sup>8</sup> attributed systematically decreased emission in CO<sub>2</sub> and CH<sub>4</sub> spectral bands to increased opacity from, and rising concentrations in, these well-mixed greenhouse gases. However, this effort was complicated by uncertainties in instrument performance, short measurement records from each instrument, and cloud contamination<sup>9</sup>.

In principle, CO<sub>2</sub> forcing can be predicted from knowledge of the atmospheric state assuming exact spectroscopy and accurate radiative transfer. Forcing can then be estimated using radiative transfer calculations with atmospheric temperature, the concentrations of radiatively active constituents including water vapour, O<sub>3</sub>, CH<sub>4</sub>, N<sub>2</sub>O, and less prominent well-mixed greenhouse gases, and changes in CO<sub>2</sub>. However, experimental validation of this forcing is needed outside the laboratory because CO<sub>2</sub> spectroscopy is an area of active research<sup>10–13</sup>. Furthermore, the fast radiative-transfer algorithms that drive regional and global climate models approximate spectroscopic absorption line-by-line calculations with errors of about  $0.6 \text{ W m}^{-2}$  (ref. 14), an amount comparable to the forcing by anthropogenic CH<sub>4</sub> and N<sub>2</sub>O.

Surface forcing represents a complementary, underutilized resource with which to quantify the effects of rising CO<sub>2</sub> concentrations on

downwelling longwave radiation. This quantity is distinct from stratosphere-adjusted radiative forcing at the tropopause, but both are fundamental measures of energy imbalance caused by well-mixed greenhouse gases<sup>15</sup>. The former is less than, but proportional to, the latter owing to tropospheric adjustments of sensible and latent heat<sup>16</sup>, and is a useful metric for localized aspects of climate response<sup>2</sup>. We focus here on clear-sky flux changes because models predict most of the CO<sub>2</sub> surface forcing to occur under clear-sky conditions<sup>15</sup>.

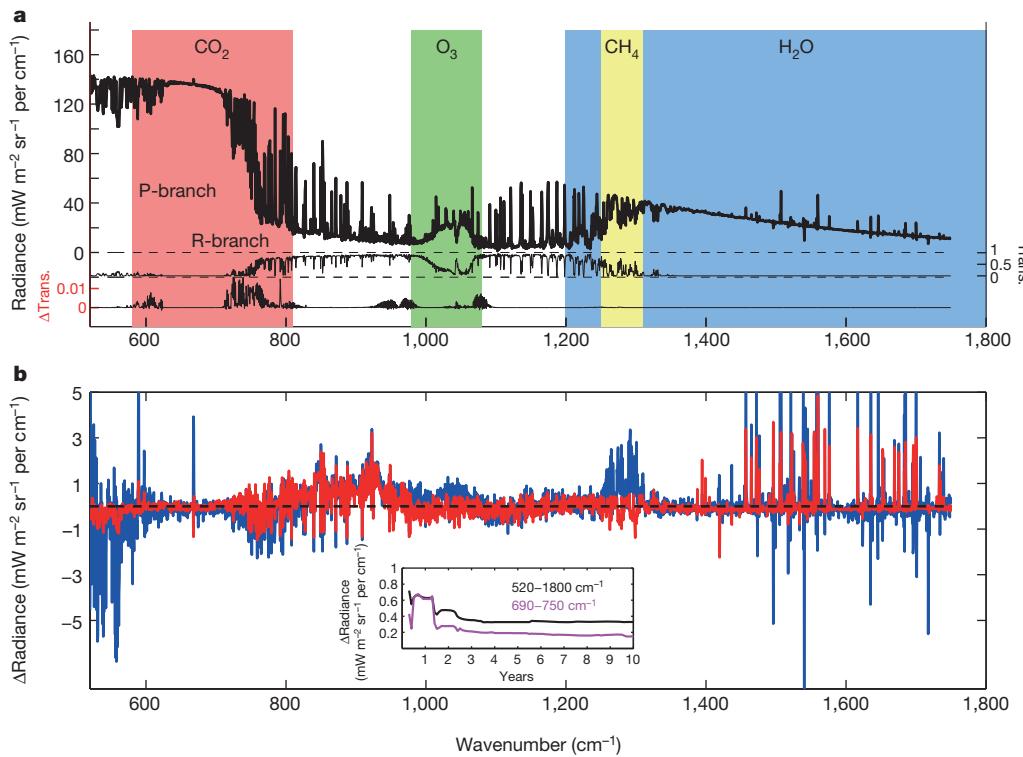
Specialized atmospheric observations at experimental sites in the mid-latitude continental Southern Great Plains (SGP) and the Arctic marine North Slope of Alaska (NSA) sites by the US Department of Energy Atmospheric Radiation Measurement (ARM) programme<sup>17</sup> produce the integrated data sets required for an independent diagnosis of the surface radiative effects of CO<sub>2</sub>. We used spectroscopic measurements from the Atmospheric Emitted Radiance Interferometer (AERI) instrument and atmospheric state data at these two sites to test whether the impact of rising CO<sub>2</sub> on downwelling longwave radiation can be rigorously detected. By basing this analysis on sets of independent measurements at high temporal frequency over long duration, we empirically established how anthropogenic emissions, mediated by variations in photosynthetic activity and respiration, are altering the Earth's surface energy balance.

However, AERI spectral measurements and trends are sensitive to many different components of the atmospheric state. To interpret these measurements and attribute specific signals to rising CO<sub>2</sub> requires an accurate radiative-transfer model that reproduces these spectra on the basis of an independent assessment of the state of the atmosphere. The model must capture instantaneous signals and long-term trends in the spectra to determine the effects of CO<sub>2</sub> on diurnal to decadal timescales.

We used the Line-by-Line Radiative Transfer Model (LBLRTM)<sup>18</sup>, which is continuously compared against other line-by-line models<sup>4</sup> and observations<sup>19</sup>. A sample clear-sky measured AERI spectrum is shown in Fig. 1a. Figure 1b shows residual spectra produced from the measurement ("obs"), minus spectra calculated ("calc") using (1) CO<sub>2</sub> concentrations from CarbonTracker 2011 (CT2011)<sup>20</sup>, which is a greenhouse gas assimilation system based on measurements and modelled emission and transport; (2) methane (CH<sub>4</sub>) profiles from CarbonTracker-CH4 (ref. 21); (3) ozone (O<sub>3</sub>) profiles from NASA's Modern-Era Retrospective Analysis for Research and Applications (MERRA)<sup>22</sup>; and (4) temperature and water-vapour profiles from radiosondes (see Methods). The measured spectrum in Fig. 1a shows Planck function behaviour near the centre of the fundamental ( $v_2$ ) CO<sub>2</sub> band and exhibits a departure from a Planck curve in the P- and R-branches of this feature, indicating that the emission in these branches is sub-saturated and could increase with increasing CO<sub>2</sub>. Water-vapour features, continuum emission, and O<sub>3</sub> emission are seen in the infrared window between  $800 \text{ cm}^{-1}$  and  $1,200 \text{ cm}^{-1}$ , and lesser features from CH<sub>4</sub> are seen around  $1,300 \text{ cm}^{-1}$ . Calculated transmission and the change in transmission with a 22 ppm CO<sub>2</sub> increase are also shown, indicating that weak vibration-rotation features in the far wings of the

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**Figure 1 | AERI spectrum and residual features.** **a**, Sample clear-sky AERI (channel 1) spectrum measured at SGP on 14 March 2001 2330Z, transmission calculation (Trans.) from LBLRTM, and the difference in transmission ( $\Delta$ Trans.) calculated for a 22 ppm change in column-averaged CO<sub>2</sub> (370–392 ppm). **b**, The

instantaneous spectral residual (obs – calc) (blue) and a spectral residual from observations for March 2001 (red). The inset indicates the running average of spectral residual rms of AERI Channel 01 (520–1,800 cm<sup>-1</sup>) and the CO<sub>2</sub> R-branch (690–750 cm<sup>-1</sup>).

fundamental and in the infrared window dominate surface radiative forcing from rising CO<sub>2</sub>.

The agreement between a single measured spectrum and the LBLRTM calculation (that is, the residual) is generally within 1.5 mW m<sup>-2</sup> sr<sup>-1</sup> per cm<sup>-1</sup>, with notable exceptions. These include the region below 550 cm<sup>-1</sup> where instrumental thermal noise is present, the centre of the v<sub>2</sub> CO<sub>2</sub> absorption feature, which is sensitive to the temperature within the instrument container, and sporadic features between 1,400 cm<sup>-1</sup> and 1,800 cm<sup>-1</sup> that are due to opaque water vapour lines. Nevertheless, over 90% of the residual is within a 3 $\sigma$  envelope established by the noise-effective radiance of 0.2 mW m<sup>-2</sup> sr<sup>-1</sup> per cm<sup>-1</sup>. Noise in the residual spectrum does not preclude long-term spectral analysis<sup>23</sup> because it has no long-term bias. As highlighted in the inset of Fig. 1 the spectral root mean square (rms) of the residual decreases with the length of time over which the spectra are averaged, with asymptotes at 0.3 mW m<sup>-2</sup> sr<sup>-1</sup> per cm<sup>-1</sup> and 0.1 mW m<sup>-2</sup> sr<sup>-1</sup> per cm<sup>-1</sup> across the AERI band-pass and in the CO<sub>2</sub> R-branch, respectively.

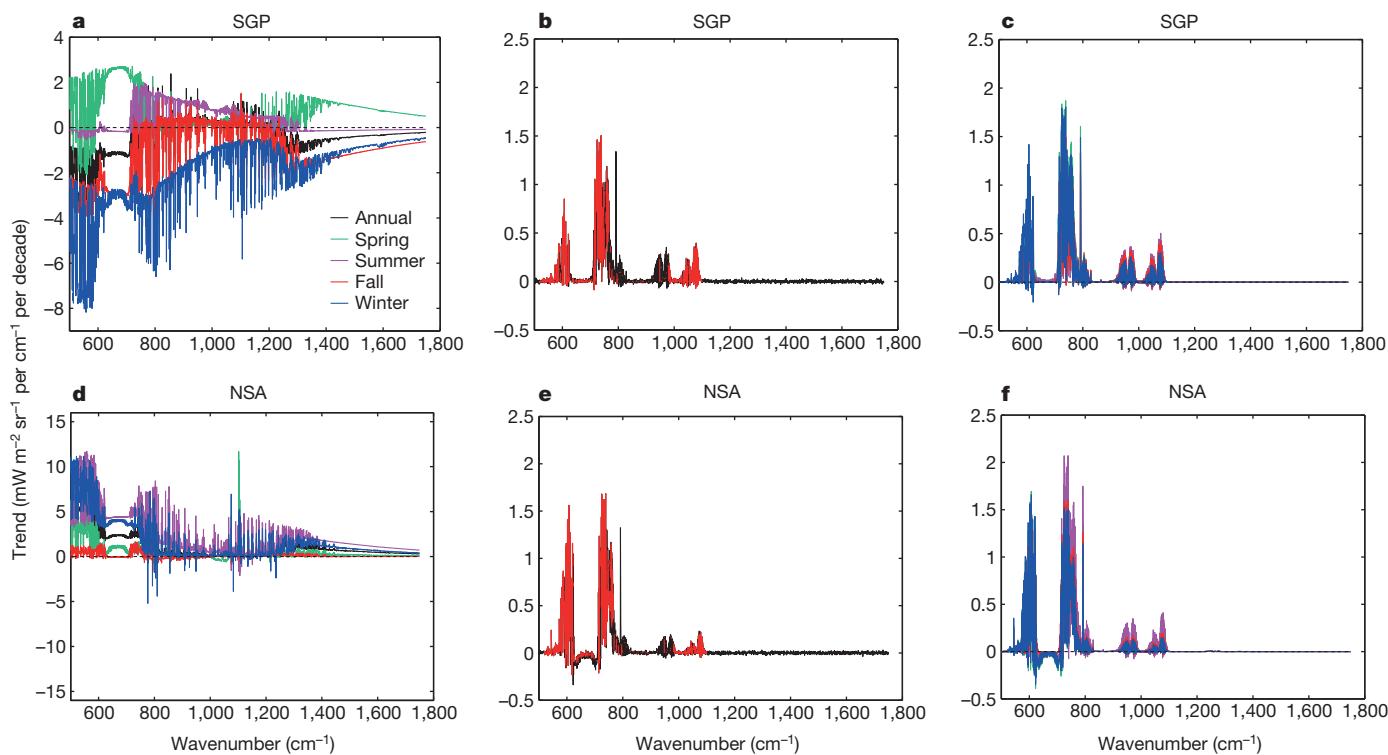
Over the length of the observation period (2000–2010), the modelled spectra at both SGP and NSA are dominated by trends associated with the temperature and humidity structure of the atmosphere rather than the smaller signal from CO<sub>2</sub>. The seasonal and annual trends in calculated clear-sky spectra at SGP (Fig. 2a) and NSA (Fig. 2d) are dominated by changes in the atmospheric thermodynamic state and are of opposite sign depending on the season. These signals arise from seasonally dependent clear-sky trends in temperature profiles and water vapour concentrations, as determined by radiosondes (see Methods) and must be taken into account to determine the forcing from CO<sub>2</sub>. We therefore construct counterfactual spectra (such spectra are produced from models that keep the CO<sub>2</sub> concentration fixed) to simulate spectra with time-invariant CO<sub>2</sub>, whereby we use temperature and water-vapour estimates from concurrent radiosondes to remove the thermodynamically derived radiometric signals from AERI spectra and isolate the signature of CO<sub>2</sub>. Since most CO<sub>2</sub> surface forcing occurs in the absence of clouds<sup>16</sup>,

we focus on clear-sky conditions, identified using the Active Remote-Sensing of Clouds (ARSCL) Value-Added Product<sup>24</sup> data set.

Differences between counterfactual spectra and coincident AERI measurements show structure in the major CO<sub>2</sub> absorption features as shown in Fig. 2b and e, at an order of magnitude greater than the long-term residual rms in Fig. 1b. Also shown are those spectral features for which the trend is non-zero at the 3 $\sigma$  level<sup>25</sup>. These panels show the unmistakable spectral fingerprint of CO<sub>2</sub>. The trends in forcing are significantly ( $P < 0.003$ ) different from zero only in the P- and R- branches of the v<sub>2</sub> CO<sub>2</sub> band.

We can exclude alternative explanations for the change in these measurements, such as instrument calibration or the temperature, water vapour, or condensate structure of the atmosphere because they would produce significant ( $P < 0.003$ ) trends in other spectral regions outside the CO<sub>2</sub> absorption bands—see Fig. 2b and e. Moreover, the spectral forcing from CO<sub>2</sub> is a strong function of changes in the CO<sub>2</sub> column concentration, and nonlinear interactions between temperature and water vapour were weak, as indicated by the lack of statistically significant differences in the seasonal and annual spectral trends in the CO<sub>2</sub> P- and R-branches. Therefore, the atmospheric structure of temperature and water vapour does not strongly affect CO<sub>2</sub> surface forcing, which is consistent with the findings of others<sup>26,27</sup>.

Finally, when we compare counterfactual spectra with radiative transfer calculations where CO<sub>2</sub> is derived from CT2011, as shown in Fig. 2c and f, we again find spectral structure only in the CO<sub>2</sub> absorption features, thereby confirming that the model is a reliable tool for determining the surface forcing from CO<sub>2</sub>. The utility of the counterfactual approach is further demonstrated by the probability distribution function of the difference in the rms spectral residual of AERI and LBLRTM with time-varying and fixed CO<sub>2</sub> concentrations for 2010 (Fig. 3). The two probability distribution functions differ substantially (two-sided  $t$ -test,  $P < 0.00001$ ). The distribution of the mean residuals also differ significantly (two-sided  $t$ -test,  $P < 0.00001$ ).



**Figure 2 | Measured and modelled spectral trends for 2000–2010.**

a, Calculated (simulated) SGP AERI clear-sky seasonal and annual spectral radiance trends for 2000–2010 with temperature, water vapour, and CO<sub>2</sub> changes. b, SGP AERI annually averaged clear-sky spectral residual trends, where the residual is the difference between the measurement and a

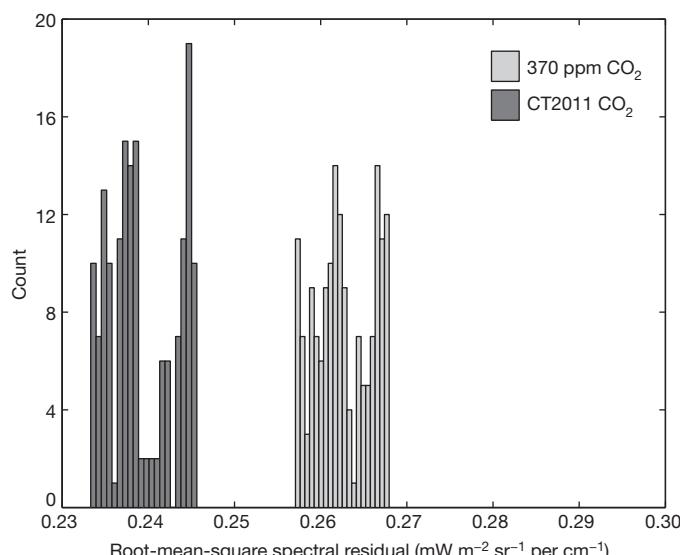
radiative-transfer calculation with CO<sub>2</sub> fixed at 370 ppm (in black). Red overlay indicates channel trends  $3\sigma$  different from zero. c, As for b but showing residual trends from differencing calculations based on CT2011 CO<sub>2</sub> concentrations and CO<sub>2</sub> fixed at 370 ppm. d, As for a but for NSA. e, As for b but for NSA. f, As for c but for NSA.

The time series of CO<sub>2</sub> surface forcing, derived from differencing AERI measurements and counterfactual calculations at the SGP (Fig. 4a) and spectrally integrating and converting to flux, shows clear and increasing trends in radiative surface forcing and seasonal variability. The least-squares trend in the long-term forcing is  $0.2 \pm 0.06 \text{ W m}^{-2}$  per decade and differs significantly ( $P < 0.003$ ) from zero. The seasonal amplitude of the forcing is  $0.1\text{--}0.2 \text{ W m}^{-2}$ , closely tracking the independently assessed pattern in the average CO<sub>2</sub> concentration in the

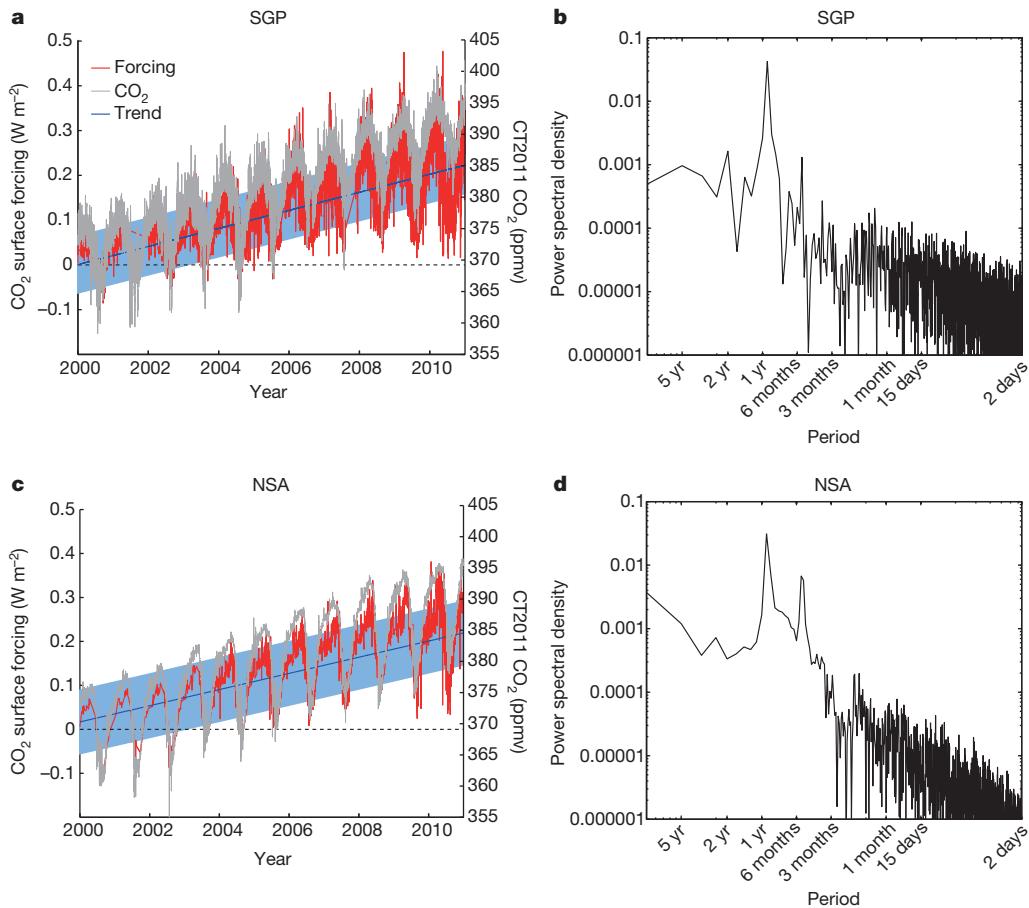
lowest 2 km of the atmosphere<sup>19</sup>. The variation in the power spectral density function of surface forcing with frequency (Fig. 4b) shows the largest peak associated with springtime photosynthesis and autumn respiration.

The time-series for NSA (Fig. 4c) also has a pronounced seasonal cycle and secular trend. The range of CO<sub>2</sub> surface forcing is similar to that of SGP, although the higher frequency variability is much less prominent. In addition, the time series of surface forcing at NSA shows increasing variability in the latter part of the 11-year analysis record. This variability results from increased numbers of samples and fewer outages (when not all of the necessary data streams are available to derive the CO<sub>2</sub> forcing) at NSA from 2004 onwards. Nevertheless, the least-squares trend in CO<sub>2</sub> surface forcing at NSA is also  $0.2 \pm 0.07 \text{ W m}^{-2}$  per decade, with a seasonal range of  $0.1 \text{ W m}^{-2}$  and differs significantly from zero ( $P < 0.02$ ).

Increasing atmospheric CO<sub>2</sub> concentrations between 2000 and 2010 have led to increases in clear-sky surface radiative forcing of over  $0.2 \text{ W m}^{-2}$  at mid- and high-latitudes. Fossil fuel emissions and fires contributed substantially to the observed increase<sup>20</sup>. The climate perturbation from this surface forcing will be larger than the observed effect, since it has been found that the water-vapour feedback enhances greenhouse gas forcing at the surface by a factor of three<sup>28</sup> and will increase, largely owing to thermodynamic constraints<sup>29</sup>. The evolving roles of atmospheric constituents, including water vapour and CO<sub>2</sub> (ref. 30), in their radiative contributions to the surface energy balance can be tracked with surface spectroscopic measurements from stand-alone (or networks of) AERI instruments. If CO<sub>2</sub> concentrations continue to increase at the current mean annual rate of 2.1 ppm per year, these spectroscopic measurements will continue to provide robust evidence of radiative perturbations to the Earth's surface energy budget due to anthropogenic climate change, but mediated by annual variations in photosynthetic activity. These perturbations will probably influence other energy fluxes and key properties of the Earth's surface and should be explored further.



**Figure 3 | Distributions of residual rms values in 2010.** Histograms of rms spectral residuals ( $520\text{--}1,800 \text{ cm}^{-1}$ ) at SGP with time-varying CO<sub>2</sub> (dark grey) and a uniform 370 ppm CO<sub>2</sub> profile (light grey). Each count corresponds to a separate spectrum measured by AERI in 2000 or 2010.



**Figure 4 | Time-series of surface forcing.** **a**, Time series of observed spectrally integrated ( $520\text{--}1,800\text{ cm}^{-1}$ ) CO<sub>2</sub> surface radiative forcing at SGP (in red) with overlaid CT2011 estimate of CO<sub>2</sub> concentration from the surface to an

altitude of 2 km (grey), and a least-squares trend of the forcing and its uncertainty (blue). **b**, Power spectral density of observed CO<sub>2</sub> surface radiative forcing at SGP. **c**, As for **a** but for the NSA site. **d**, As for **b** but for the NSA site.

**Online Content** Methods, along with any additional Extended Data display items and Source Data, are available in the online version of the paper; references unique to these sections appear only in the online paper.

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**Author Contributions** D.R.F. implemented the study design, performed the analysis of all measurements from the ARM sites, and wrote the manuscript. W.D.C. proposed the study design and oversaw its implementation. P.J.G. is the AERI instrument mentor and ensured the proper use of spectral measurements and quality control. M.S.T. mentored the implementation of the study and oversaw its funding. E.J.M. and T.R.S. performed calculations and analysis to determine fair-weather bias. All authors discussed the results and commented on and edited the manuscript.

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## METHODS

The methodology for this investigation focused on analysing time series of downwelling infrared spectra to determine the effects of CO<sub>2</sub> on these measurements and thereby to estimate its surface radiative forcing. The analysis also used temporally coincident measurements of atmospheric temperature and water vapour, retrievals of cloud occurrence, and assimilation estimates of CO<sub>2</sub> to construct simulated counterfactual measurements where CO<sub>2</sub> is held fixed.

**Code and data availability.** The measurement data sets used for this analysis are freely available through the ARM data repository (<http://www.arm.gov>). The radiative transfer codes are also freely available at <http://rtweb.aer.com>. CarbonTracker results were provided by NOAA/ESRL (<http://carbontracker.noaa.gov>). CarbonTracker-CH<sub>4</sub> results were provided by NOAA/ESRL (<http://www.esrl.noaa.gov/gmd/ccgg/carbontracker-ch4/>). The MERRA data used in this analysis are freely available for download at <ftp://goldsmr3.sci.gsfc.nasa.gov/data/s4pa/MERRA/MAI3CPASM.5.2.0/>. The Broadband Heating Rate Profile (BBHRP) data files, used to assess fair-weather bias, are freely available on the ARM BBHRP web page at <http://www.arm.gov/data/eval/24> under <http://dx.doi.org/10.5439/1163296> for the time-varying data stream and <http://dx.doi.org/10.5439/1163285> for the fixed CO<sub>2</sub> data stream. The computer routines used in this analysis will be made available upon request.

**Schematic.** The schematic (Extended Data Fig. 1) shows how CO<sub>2</sub> surface forcing is derived from differencing Atmospheric Emitted Radiance Interferometer (AERI) observations with a model calculation based on coincidental temperature and water vapour profiles. This produces a radiance difference (called a spectral residual) that is converted to flux units through a conversion factor calculated with a radiative-transfer model based on local meteorological conditions. The spectral residual, which is driven by observations, shows features only in the CO<sub>2</sub> absorption bands, meaning that this is a signal just from CO<sub>2</sub>. The integration of forcing over frequency gives a broadband forcing term and contains a secular trend due largely to anthropogenic emissions.

**Data sets.** Spectra observed by the AERI instrument channel 1 were analysed over the period from 2000 through the end of 2010 at the ARM SGP and NSA sites. The AERI instrument contains two channels that cover a spectral range of 520–3,020 cm<sup>-1</sup> (19.2–3.3 μm) with a resolution of 0.5 cm<sup>-1</sup>. Calibrated spectra are recorded every 8 min. Channel 1 observes the 520–1,800 cm<sup>-1</sup> spectral region, which covers at least 99.97% of the longwave CO<sub>2</sub> surface forcing. The absolute calibration is traceable to NIST standards<sup>31</sup> and has been transferred from blackbodies with uncertainties better than 0.05 K (3σ). The resulting combined absolute uncertainty of each AERI spectrum is 1% (3σ)<sup>32</sup>. A summary of the AERI instrument's spectral stability can be seen in Extended Data Fig. 2, which shows the time series of effective laser wavenumber that is used to define the instrument's spectral scale. This time series exhibits a trend in the date of 4 ppm (relative) per decade around the nominal 15,799 cm<sup>-1</sup> laser wavenumber. Given that the uncertainty in the AERI spectral calibration is ±5 ppm (relative) (3σ), we have confidence that there is no statistically significant drift in the AERI spectral calibration over the course of the observation period (2000–2010).

Although the AERI record extends from 1995 to the present, temporally resolved estimates of the CO<sub>2</sub> concentration profiles from CT2011<sup>20</sup>, which were necessary for validation of the radiative transfer model (see below), were only available beginning in 2000 and extending through 2010, resulting in an analysis period of 11 years. Overall, the AERI is a very robust instrument that inherently produces continuous, reliable operational data. Quality control for the spectra was achieved by ensuring that valid sky spectra were being observed by the instrument with the hatch open, and that unphysical radiance values, anomalously low variances in brightness temperature across the infrared window band (800–1,200 cm<sup>-1</sup>) and scenes with substantial variability in the view of the hot blackbody were removed.

CO<sub>2</sub> profile best estimates were obtained from CT2011, which reported eight daily profiles in the vicinity of, and which were spatially interpolated to, the SGP and NSA sites at the European Centre for Medium-Range Weather Forecasting forecast model levels on a 40° latitude by 66° longitude grid covering North America and may have spin-up effects<sup>33</sup>. The CT2011 assimilates data from many sources, including the average 12:00 to 16:00 local time CO<sub>2</sub> *in situ* measurement from the ARM Precision Carbon Dioxide Mixing Ratio System tower at 60 m altitude at SGP<sup>34,35</sup>. Comparisons of CT2011 against aircraft profiles collected by the ARM Carbon Measurement Experiment<sup>36</sup> at SGP, which are not assimilated into CT2011, exhibit a bias of less than 0.5 ppm up to 3,600 m above sea level, with a standard deviation ranging from 1 ppm in winter to 3 ppm in summer. Between 2000 and 2011, the global CO<sub>2</sub> concentration at the surface increased from about 369 ppm to about 392 ppm, as measured by ARM-NOAA Earth Science Research Laboratory (ESRL) flasks, while the CO<sub>2</sub> concentrations at SGP and NSA increased at 1.95 ppm versus 1.98 ppm per decade, respectively. Time series of CT2011 profiles at SGP and fossil fuel and fire components are shown in Extended Data Fig. 3. Output from the CT2011 data set were used to create a baseline CO<sub>2</sub> value of

370 ppm from which surface-forcing calculations were determined. These outputs enabled us to demonstrate that spectral residuals derived from the Line-By-Line Radiative Transfer Model (LBLRTM) version 12.2 (ref. 18) with CT2011 had a lower rms difference than spectral residuals derived from the CT2011 baseline.

Temperature and humidity profiles were gathered from the ARM Balloon-Borne Sounding System (BBSS). At the SGP site, radiosondes were launched four times per day throughout the period of analysis, while at the NSA site, radiosondes were launched two times per day over the analysis period. Three different sounding technologies have been used at both locations. At the SGP site, the Vaisala RS-80 technology was used until April 2001; the Vaisala RS-90 was used between May 2001 and February 2005; and the Vaisala RS-92 was used from February 2005 to present. At the NSA site, RS-80 technology was used through April 2002, the RS90 was used from April 2002 until January 2005, and the RS-92 technology was used since January 2005. To account for the dry bias that has been exhibited particularly by the RS-80 sensors<sup>37</sup>, we scaled the radiosonde-reported humidity profile by the ratio of the total column water retrieved from the microwave radiometer<sup>38,39</sup> to the radiosonde-derived total column water. We estimate that the corrected radiosonde uncertainty is under 0.5 °C and 5% relative humidity (3σ), given the technology's stated precision and accuracy specifications. The change in radiosonde technology in the middle of the investigation could greatly complicate this analysis, and we would expect that a change in technology could introduce a jump-discontinuity in the forcing. The technology change occurred near the start of the analysis period in 2001, and the water-column amounts across the technology change differ at the *P* = 0.1 level for a two-sided *t*-test. However, we analysed the probability distribution function of the microwave radiometer to BBSS column water vapour both before and after the RS90 to RS92 change and did not find statistically significant differences (Extended Data Fig. 4).

Cloud clearing, which was essential to this analysis, was achieved through the ARM ARSCL Value-Added Product, which combines the Micropulse Lidar, Millimeter Wavelength Cloud Radar, and Vaisala and Belfort Ceilometers to produce six different cloud masks of varying sensitivity<sup>24,40</sup>. We analysed spectra collected exclusively during clear-sky conditions, as identified by the absence of clouds in all of the coincident ARSCL masks.

**Thermodynamic trends.** Trends in the thermodynamic state at SGP and NSA are shown in Extended Data Fig. 5. The radiosonde-derived annually averaged least-squared trend in the clear-sky temperature profile at SGP showed an increase in spring in the lowest 2 km of the atmosphere, but decreases elsewhere. Lower-atmosphere water concentrations showed trends of opposite sign depending on season and altitude. At NSA, annually averaged clear-sky temperature profile trends were positive in the lower troposphere except in spring. Lower-tropospheric humidity also increased at NSA, except in the autumn. It is important to note that these thermodynamic trends are based on clear-sky conditions only and are distinct from all-sky trends.

Still, temperature trends can have a non-zero impact on spectral downwelling longwave radiation, which can be seen in Fig. 2e, where spectral trends at wavenumbers in the near-wings of the fundamental (v<sub>2</sub>) CO<sub>2</sub> band are negative. The cause of the negative trends arises from the nonlinear interactions between CO<sub>2</sub> and temperature trends, especially where there is a contrast in temperature trends between the boundary layer and the free troposphere under temperature inversion conditions, which are common at the NSA site. The effects of CO<sub>2</sub> and temperature generally produce separable impacts on the spectral downwelling longwave radiation, especially in the weak vibration-rotation bands that dominate the observed CO<sub>2</sub> surface radiative forcing. However, the effects interact nonlinearly near the centre of the fundamental (v<sub>2</sub>) CO<sub>2</sub> band<sup>27</sup>. This nonlinearity complicates the separability of temperature and CO<sub>2</sub> trends over this narrow spectral range.

**CarbonTracker-CH<sub>4</sub> and MERRA data use.** All radiative transfer calculations performed in this analysis used CarbonTracker-CH<sub>4</sub> profile information. CarbonTracker-CH<sub>4</sub> (ref. 21) provides eight daily profiles on a 40° latitude by 60° longitude global grid of component contributions of CH<sub>4</sub> to the atmospheric concentration (a background component, a component due to coal and oil/gas production, a component due to animals, rice cultivation, and waste, a component due to wetlands, soil, oceans, and insects/wild animals, a component due to emissions from fires, and a component due to emissions from oceans). We sum those components and linearly interpolate them to the ARM SGP and NSA sites in space and time to coincide with the radiosonde observation.

We also used NASA's Modern-Era Retrospective Analysis for Research and Applications (MERRA)<sup>22</sup> for O<sub>3</sub> profiles for all radiative transfer calculations. MERRA provides eight daily O<sub>3</sub> profiles at 42 pressure levels at 1.25° horizontal spacing and these were linearly interpolated to the ARM SGP and NSA sites in space and time to coincide with the radiosonde observation.

**Radiative transfer.** LBLRTM version 12.2 was used to simulate AERI spectra with and without CO<sub>2</sub> increases. LBLRTM was run on the 200 levels that are linearly spaced in log(pressure) onto which the radiosonde temperature and humidity

profiles have been interpolated. The conversion from spectral radiance to spectral flux was performed using spectrally dependent radiance-to-flux conversion factors determined through three-point quadrature flux calculations for each profile of atmospheric thermodynamic state corresponding to the AERI measurements following ref. 41. Histograms of the calculated distributions of the spectral conversion factor between radiance and flux are shown in Extended Data Fig. 6. These demonstrate that the conversion from radiance to flux depends on wavenumber, with stronger bands showing an Angular Distribution Model value very close to unity, indicating that the radiance is isotropic, and weaker bands having lower Angular Distribution Model values, indicating that the radiance is anisotropic.

To determine the contribution of CO<sub>2</sub> to surface forcing, we removed the signal associated with the other radiatively active constituents. To do this, we performed a radiative-transfer calculation using LBLRTM with observed temperature and water-vapour profiles from radiosondes (with water vapour scaled by microwave radiometer column water retrievals) and with a fixed CO<sub>2</sub> concentration set to 370 ppm (the CO<sub>2</sub> atmospheric concentration at the beginning of the analysis period) for cloud-free conditions when the AERI instrument produced calibrated spectra. We differenced these calculations against the actual AERI spectral measurements to produce a time-series of residual spectra between 2000 and the end of 2010. Since the fixed CO<sub>2</sub> concentration is the primary difference between the inputs to each calculation and the actual time-varying atmospheric states, the residual spectra contains the signal of changing CO<sub>2</sub>. We integrated the spectra over frequency (wavenumber) and converted radiance to flux, based on the quadrature calculations described above, to produce a forcing value at each time step.

We used the Rapid Radiative Transfer Model<sup>18</sup> to perform the calculations of fair-weather bias, as described below. It calculates lower radiative forcing values than line-by-line models for the same atmospheric state<sup>14</sup>.

**Error budget.** The error budget for each surface-forcing value contains contributions from several sources, including radiometric noise, instrument spectral and radiometric calibration, and residual removal. Instrument radiometric calibration is achieved through a suite of diagnostics, including detector nonlinearity characterization, electronics calibration, field-of-view testing, and routine views of the instrument blackbodies. The achieved radiometric uncertainty is better than 1% ( $3\sigma$ ) of ambient radiance for a single observation and decreases with the number of observations. The instantaneous noise-effective spectral radiance specification for the AERI instrument is  $<0.2 \text{ mW m}^{-2} \text{ sr}^{-1} \text{ per cm}^{-1}$  for 670–1,400 cm<sup>-1</sup>, for the rms of a 2-min ambient blackbody view. Spectral calibration is achieved through routine comparisons with stable atmospheric lines and has an uncertainty of better than 0.08 cm<sup>-1</sup> (corresponding to 5 ppm of the laser wavenumber) ( $3\sigma$ ). Residual removal uncertainty consists of two subcomponents: (1) uncertainty in the atmospheric structure as provided by the BBSS and (2) uncertainty in the spectroscopy that informs LBLRTM. The former is derived from the radiosonde uncertainty propagated in spectral space with LBLRTM; the latter is derived from the HITRAN database uncertainty codes<sup>10</sup>. The effects of spectroscopic uncertainty are derived from forcing estimates with LBLRTM calculations where line parameters are modified according to stated HITRAN error codes, which report uncertainties in intensity, half-width, temperature dependence, and pressure shift for each line.

**Trend determination.** Both spectral trends, as shown in Fig. 2 and broadband trends, as shown in Fig. 4, are determined from a least-squares regression using Matlab's "polyfit" (<http://www.mathworks.com/help/matlab/ref/polyfit.html>) function of the estimate values based on clear-sky conditions only. Uncertainty in the trend was calculated from the formulae in Weatherhead *et al.*<sup>25</sup>. The trend in seasonal amplitude was determined from a least-squares regression from the peaks of the monthly averaged time-series of the forcing.

**Fair-weather bias.** One key question associated with the methods concerns the potential for fair-weather bias in this approach due to the screening of data based on the ARSCL cloud masks. A bias could occur if there is a relationship between surface forcing and cloud cover that is distinct from, and would alter the trends in, what was found under clear-sky conditions. To address this question, we have undertaken a set of calculations that follow the BBHRP<sup>42</sup> calculations based on the Radiatively Important Parameters Best Estimate (RIPBE)<sup>43</sup> Value Added Products (see <http://www.arm.gov/data/eval/24> for details). This product contains broadband profiles of fluxes and heating rates throughout the atmosphere. We have recalculated BBHRP profiles at SGP for 2010 based on fixed CO<sub>2</sub> concentrations and compared those to the original BBHRP profiles (that contained time-varying CO<sub>2</sub>) and subset the data that were identified as clear-sky in RIPBE (through the ARSCL flags) in order to estimate this bias. The results are shown in Extended Data Fig. 7 and indicate that the all-sky surface forcing is  $0.05 \text{ W m}^{-2}$  less than the clear-sky forcing, which represents a 25% difference in the two quantities. This finding is expected because there is non-negligible overlap between cloud absorption, which is broadband, and CO<sub>2</sub> absorption features, so clouds mask the forcing from CO<sub>2</sub>. We note that the difference in clear-sky and all-sky surface radiative forcing values is not statistically significant and that the extent of the fair-weather bias depends on

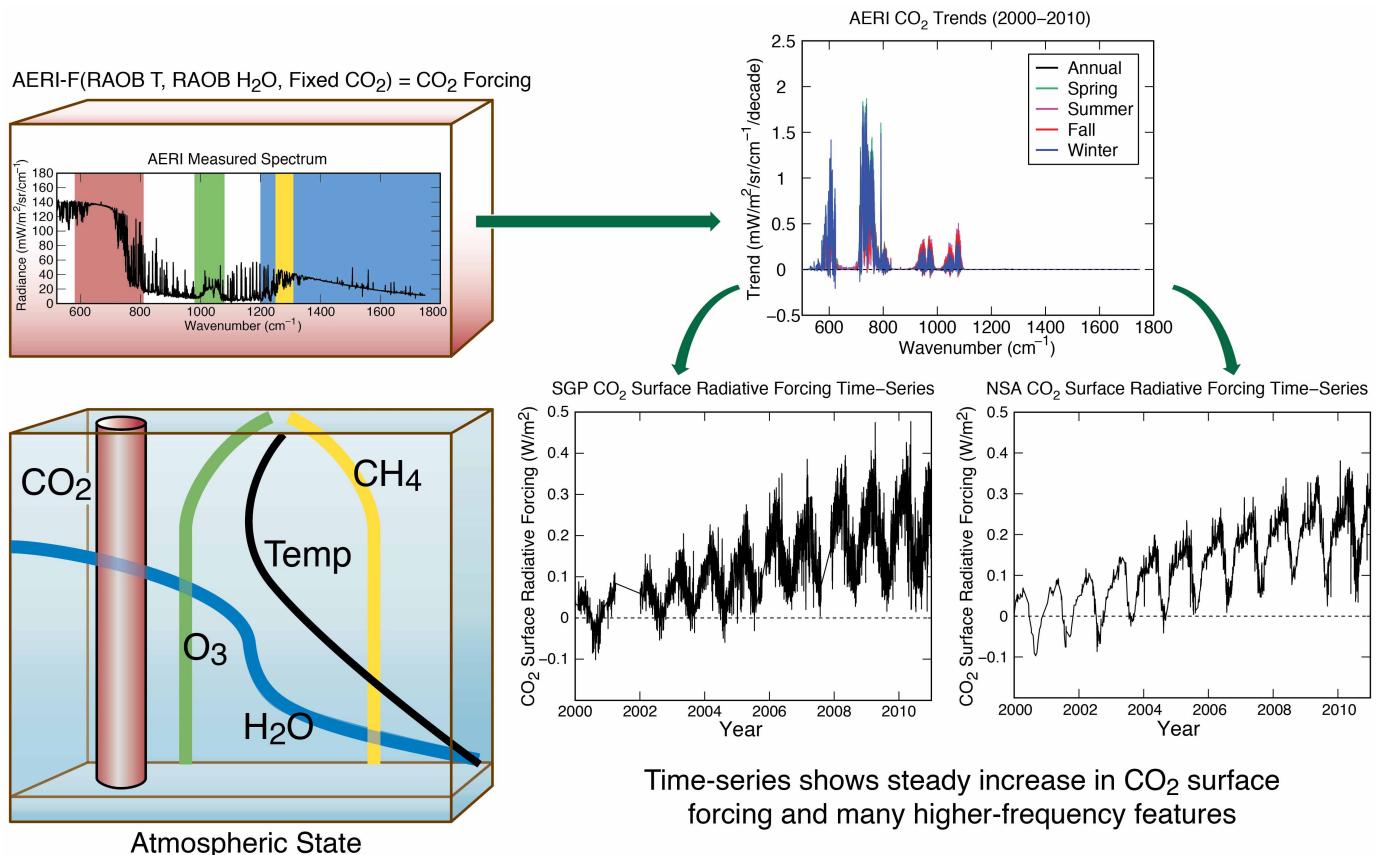
the occurrence frequency of cloudy conditions, which varies year-to-year at the SGP and NSA sites.

**Surface versus TOA forcing.** In general, several studies have found that there is a complicated relationship between surface and top-of-atmosphere (TOA) forcing, depending on the atmospheric thermodynamic and latent heating structure, with the former being less than the latter<sup>15,44</sup>. This is further underscored by other findings for an increase in CO<sub>2</sub> from 287 to 369 parts per million by volume (ppmv) which shows a forcing at the surface of  $0.57 \text{ W m}^{-2}$  and a forcing at 200 hPa of  $1.92 \text{ W m}^{-2}$  from line-by-line models<sup>45</sup>.

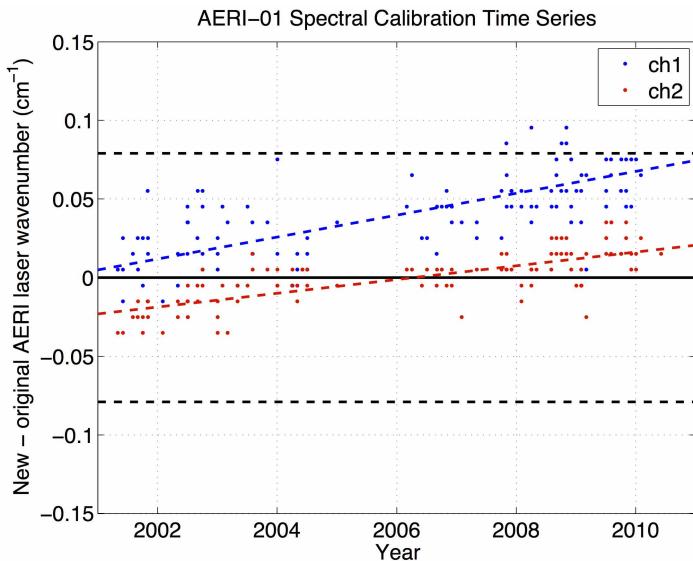
To specifically address the relationship between forcing at the surface and TOA at the ARM sites, we have also undertaken a set of calculations that use the BBHRP infrastructure described above to demonstrate the relationship between surface and TOA flux changes. Results are shown in Extended Data Fig. 8, and indicate that magnitudes of the flux changes at the surface and TOA are correlated, although the TOA flux changes are larger because the downwelling surface radiation of the former quantity is distributed into sensible and latent heat. We note that the flux change at the surface and TOA is calculated by differencing fluxes with time-varying CO<sub>2</sub> minus fluxes with fixed CO<sub>2</sub> at each altitude, and the decreased emission to space and increased emission to the surface leads to flux changes of opposite sign for TOA and surface flux changes, respectively.

**Vertical sensitivity.** Using LBLRTM, we calculated the vertical sensitivity of broadband longwave surface forcing due to perturbations in the atmospheric state profile. These calculations were based on model atmospheres, including the US standard, the tropical, the mid-latitude summer, the mid-latitude winter, the sub-Arctic summer, and the sub-Arctic winter atmospheres, which span a broad range of atmospheric thermodynamic states<sup>46</sup>. We separately perturbed the temperature by 1°K, 1% water vapour, 10 ppm CO<sub>2</sub>, 10% O<sub>3</sub>, and 1 ppm CH<sub>4</sub> profiles for a 1-km-thick layer of the atmosphere from the surface to 20 km. Results are shown in Extended Data Fig. 9 and indicate that, with the exception of O<sub>3</sub>, contributions to surface forcing are dominated by the bottom 5 km of the atmosphere, regardless of the thermodynamic state. O<sub>3</sub> perturbations have a more complicated vertical structure because O<sub>3</sub> is partially transparent in the troposphere and its concentration peaks in the stratosphere, meaning that upper-level O<sub>3</sub> concentration perturbations at the 10% level can have a modest impact on the surface energy balance.

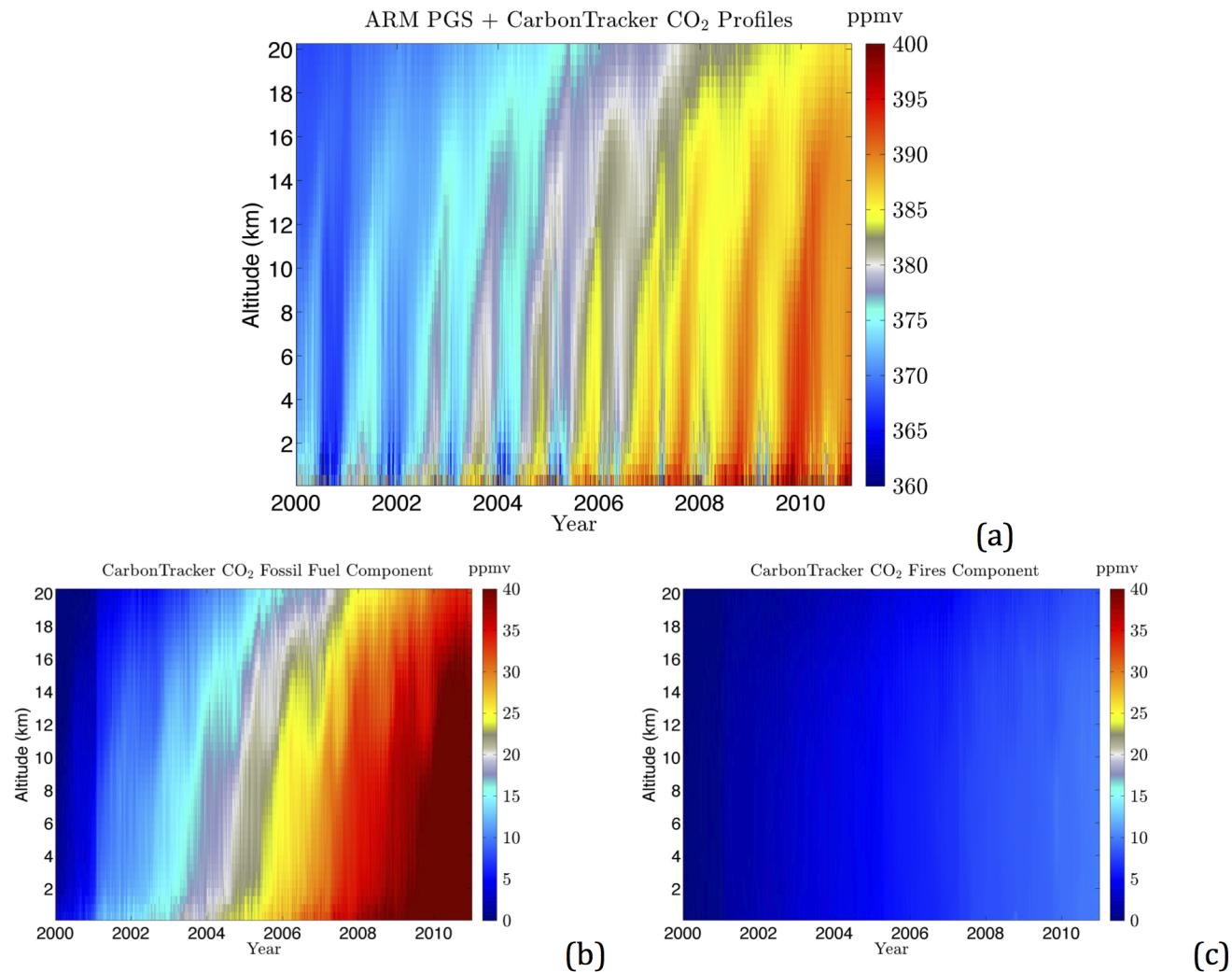
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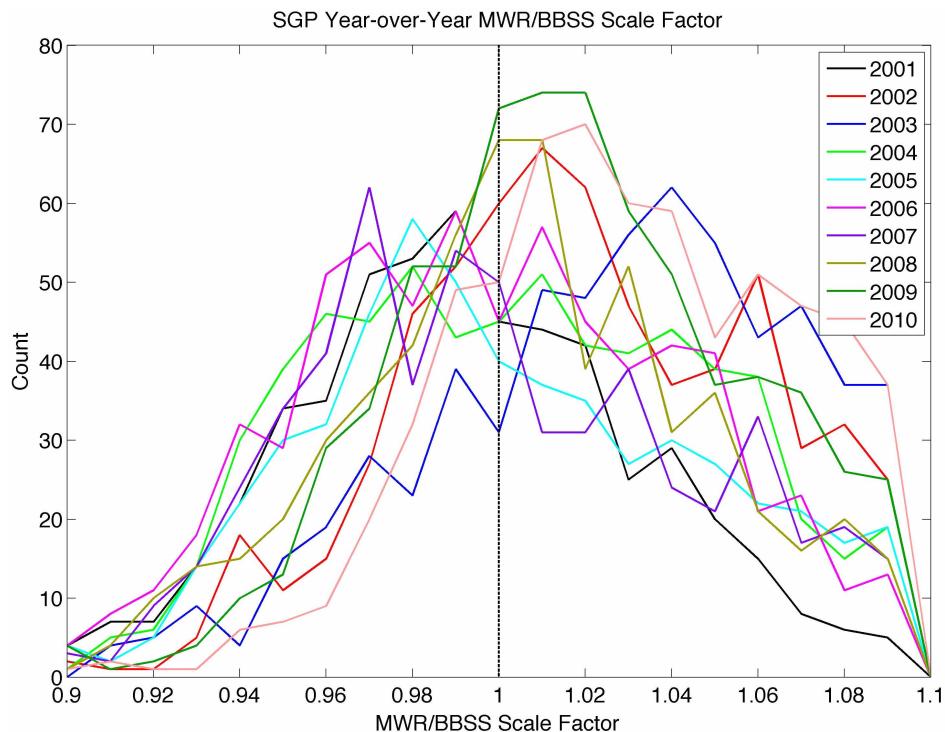
**Extended Data Figure 1 | Schematic.** Schematic of the derivation of surface forcing from AERI observations and calculations based on the atmospheric structure.



**Extended Data Figure 2 | AERI instrument stability.** Time series of the AERI-instrument-derived laser wavenumber around a nominal frequency of  $15,799 \text{ cm}^{-1}$ .

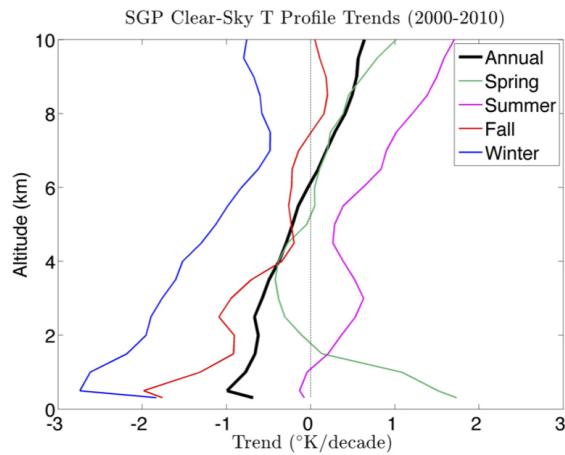


**Extended Data Figure 3 | CarbonTracker profiles.** **a**, CT2011 profile time series of CO<sub>2</sub> at the SGP site. **b**, CT2011 fossil fuel component of the CO<sub>2</sub> profile. **c**, CT2011 biomass burning component of the CO<sub>2</sub> profile. PGS, the ARM Precision Gas System Carbon Dioxide Mixing Ratio System.

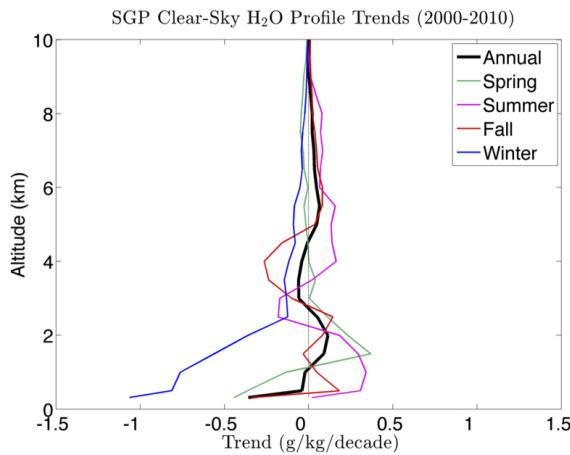

**Extended Data Figure 4 | Microwave radiometer radiosonde scaling.**

Distribution of microwave radiometer (MWR) precipitable water vapour to the precipitable water vapour derived from radiosondes for each year of the

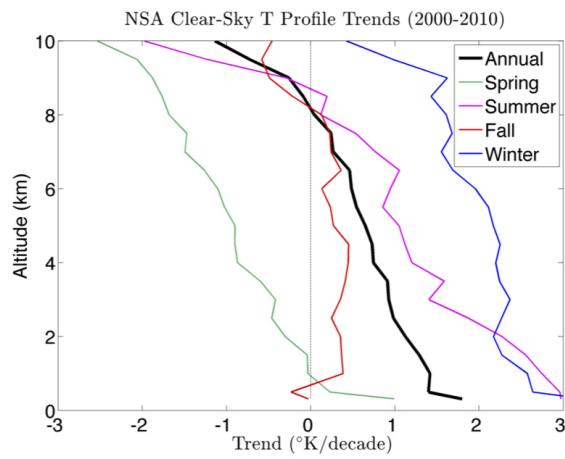
investigation at the ARM SGP site. Each count corresponds to the scaling between a collocated radiosonde and microwave radiometer retrieval.



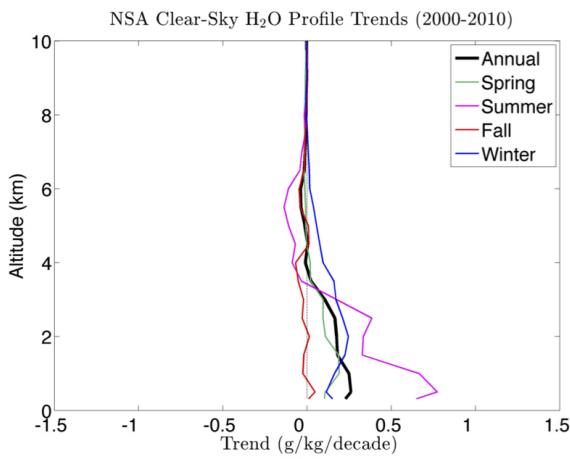
(a)



(b)



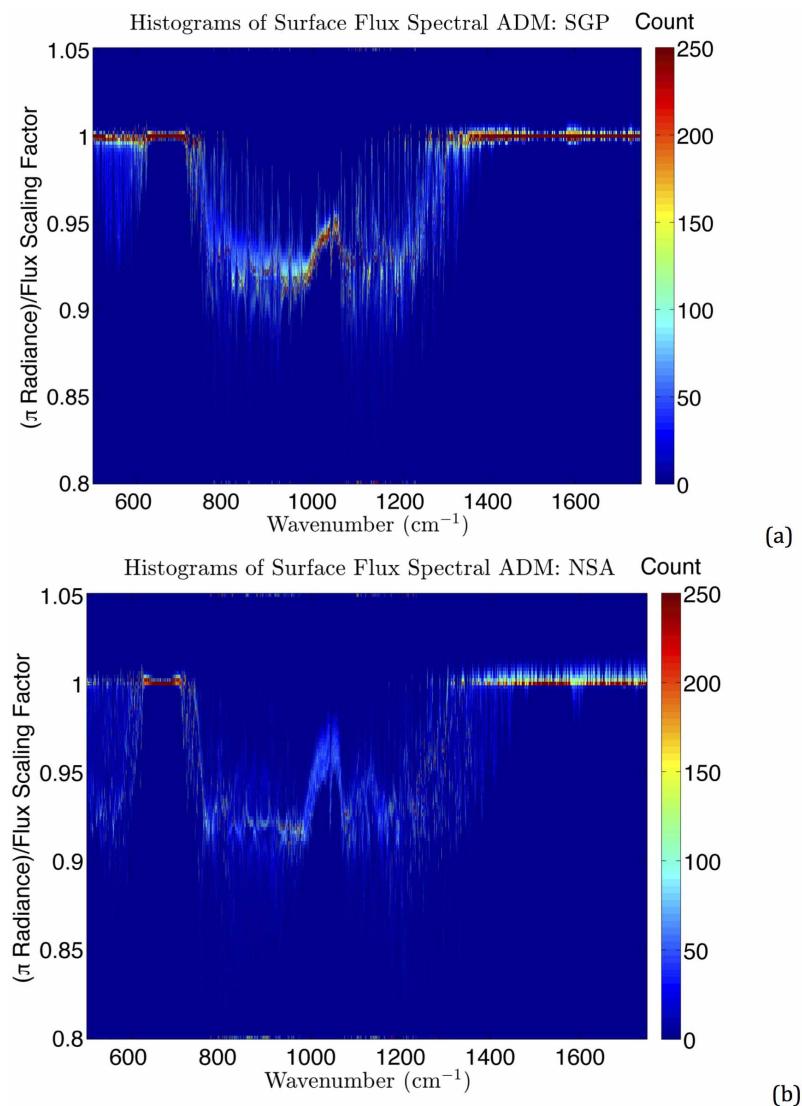
(c)



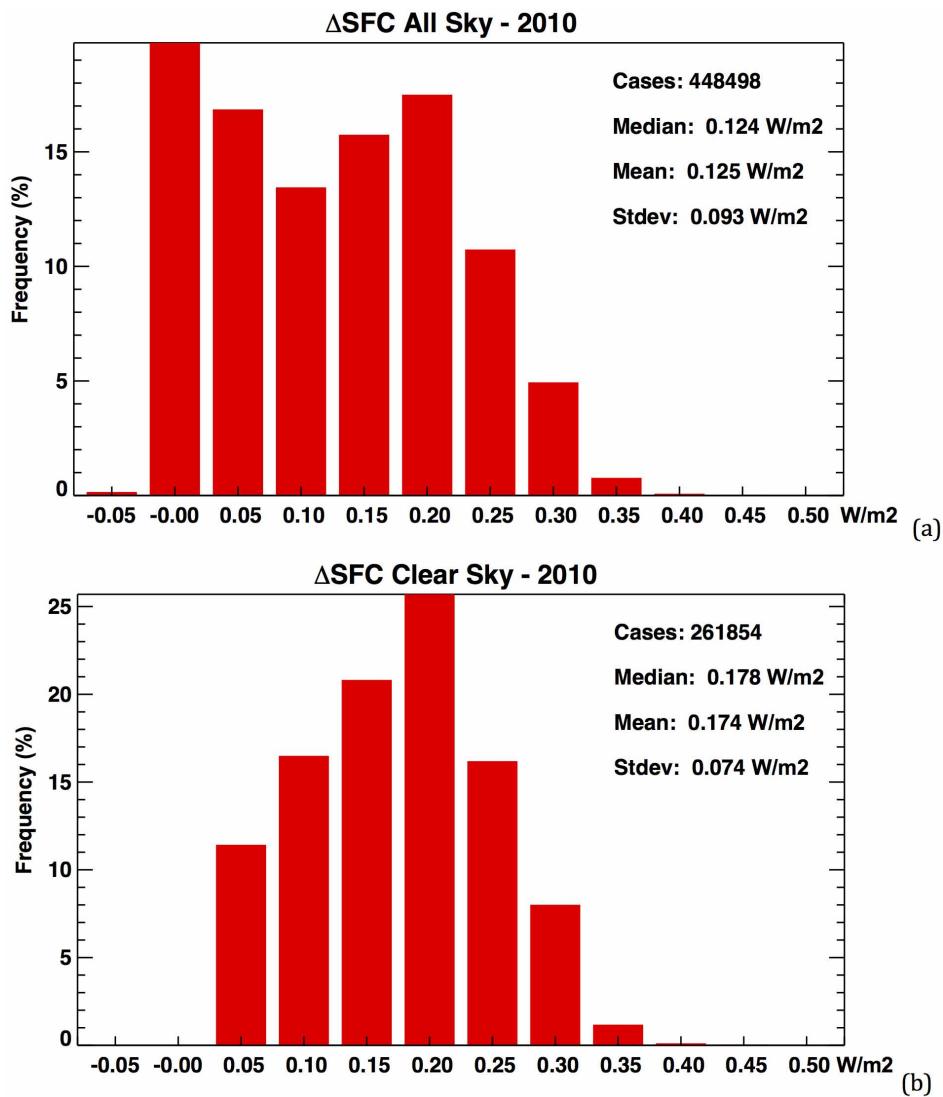
(d)

**Extended Data Figure 5 | Thermodynamic trends.** **a**, Annual and seasonal clear-sky temperature (T) profile trends derived from radiosondes and ARSCL data for cloud-clearing at SGP from 2000 to 2010. **b**, Same as **a** but for

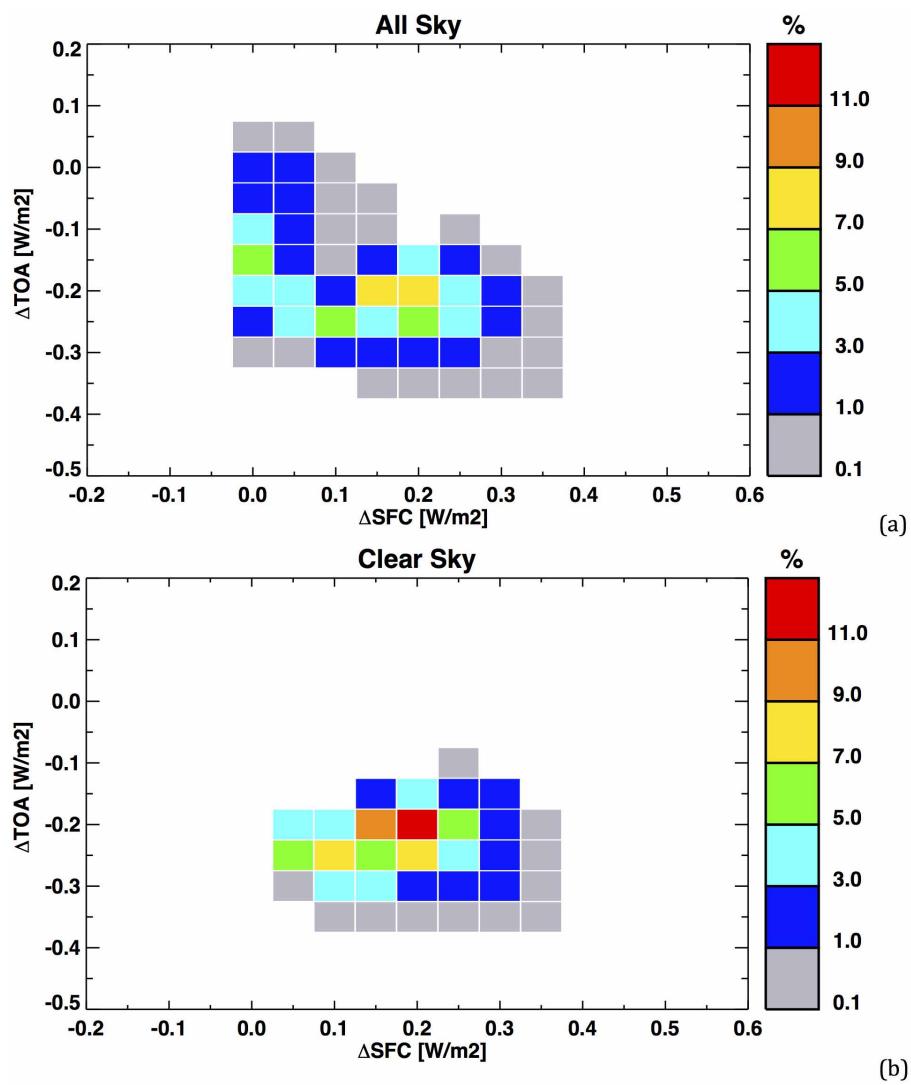
water vapour ( $\text{H}_2\text{O}$ ) profile trends. **c**, As for **a** but temperature profile trends at NSA. **d**, As for **b** but for water vapour profile trends (in grams of water vapour per kilogram of air per decade) at NSA.



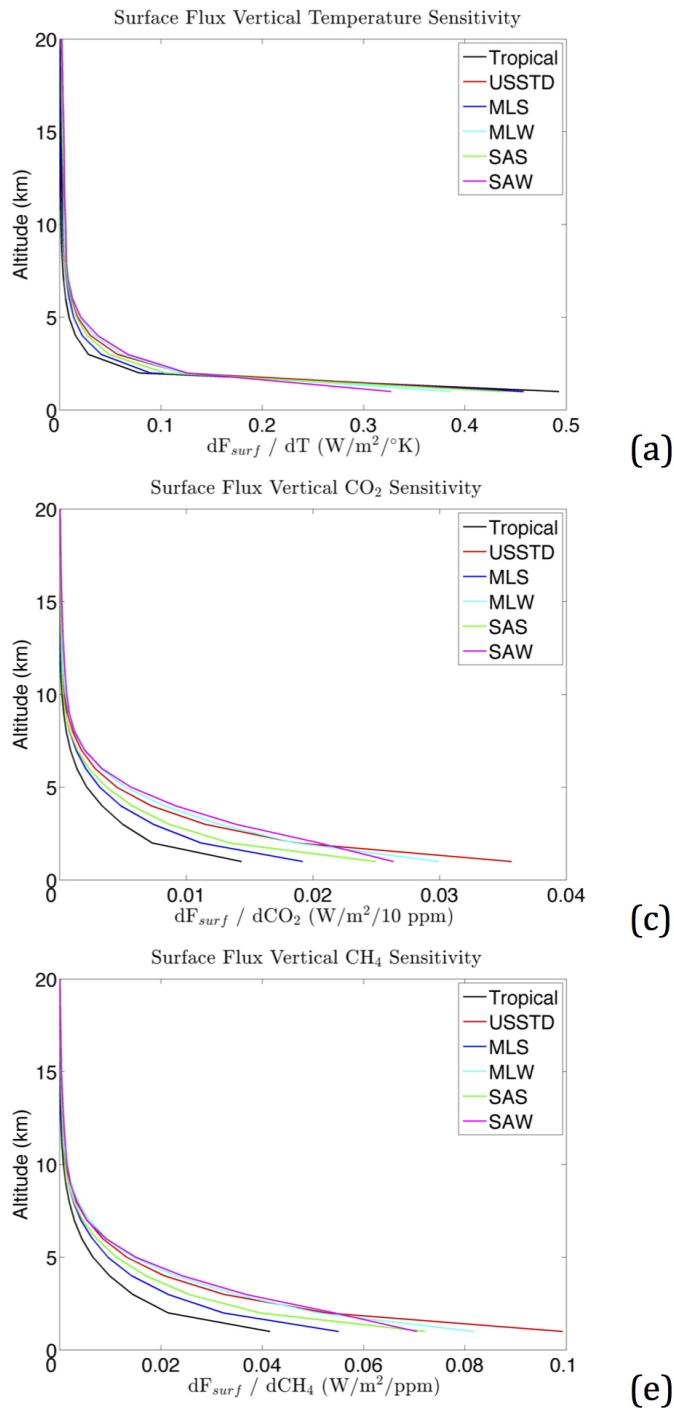
**Extended Data Figure 6 | Conversion from radiance to flux.** Histogram of zenith radiance to flux spectral conversion for AERI channel 1 spectral channels based on LBLRTM calculations based on the thermodynamic profiles from the ARM SGP site from 2000 to 2010. **b**, As for **a** but for the NSA site. ADM, Angular Distribution Model.



**Extended Data Figure 7 | Fair-weather bias.** **a**, Histogram of the difference in flux between BBHRP calculations with time-varying CO<sub>2</sub> and calculations where CO<sub>2</sub> = 370 ppmv for all profiles at 30-min resolution during 2010 at SGP. **b**, As for **a** but for the subset of data identified by the ARSCL as clear-sky.



**Extended Data Figure 8 | TOA and surface fluxes.** **a**, Occurrence frequency (in per cent) plot of tropopause versus surface forcing based on BBHRP calculations with time-varying CO<sub>2</sub> and where CO<sub>2</sub> = 370 ppmv for all profiles at 30-min resolution during 2010 at SGP for all-sky conditions as identified by ARSCL flags. **b**, As for **a** but for clear-sky conditions.



(a)

(b)

(c)

(d)

(e)

**Extended Data Figure 9 | Surface flux sensitivity to atmospheric profiles.**  
**a**, The sensitivity of the surface radiative flux ( $F_{surf}$ ) to the level of a 1°K perturbation in temperature for different model atmospheres including tropical, US standard (USSTD), mid-latitude summer (MLS), mid-latitude

winter (MLW), sub-Arctic summer (SAS), and sub-Arctic winter (SAW)<sup>46</sup>.  
**b**, As for **a** but level perturbations are given as percentage H<sub>2</sub>O. **c**, As for **a** but level perturbations are 10 ppm CO<sub>2</sub>. **d**, As for **a** but level perturbations are 10% O<sub>3</sub>. **e**, As for **a** but level perturbations are 1 ppm CH<sub>4</sub>.

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