

High Energy δ – ray Production by Muons in the *BABAR* Detector

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Abstract

In this note we present a measurement of the δ – ray production cross section by muons in the *BABAR* detector [1], for δ – rays between 50 and 300 MeV. The *BABAR* experiment recorded over 500M $e^+e^- \rightarrow \mu^+\mu^-$ events near the upilon resonances over its 7 year run, producing over 1G muons with a lab frame energies between 3 and 9 GeV. These muons produced δ – rays when traversing the beampipe and other material in the BaBar inner tracker region. We identify δ – ray candidates by associating a reconstructed low-momentum negatively charged track with a nearby reconstructed high-momentum track in events consistent with the $e^+e^- \rightarrow \mu^+\mu^-$ process, where the two tracks point of closest approach (POCA) is consistent with a known piece of dense material. A detailed simulation predicts the primary δ – ray background comes from electrons produced in the conversion of photons radiated by the muons, which we measure by applying the same reconstruction algorithm to positively charged tracks. Monte Carlo simulation is used to estimate the remaining backgrounds, which are estimated to be less than 1% of the predicted signal. Simulation is also used to estimate the reconstruction efficiency, the δ – ray energy resolution, and to predict the mis-identification of the radiator material. We extract a cross-section by normalizing the corrected production rate to the estimated material traversed by the associated reconstructed muon candidate track, as a function of the reconstructed electron candidate track energy. The cross-section errors are dominated by uncertainty in the radiator material mass, which is roughly 5% and coherent across the energy points measured. The resulting cross-section

is found to be consistent with theoretical predictions within errors.

Keywords: **Material Interactions**

1. Introduction

Atomic electrons ejected by high-energy charged particles passing through matter (δ -rays) can be a background for experiments using electrons as signals. An example is the $\mu \rightarrow e$ experiment [2], which will search for muon to electron
5 conversion in the field of an aluminum nucleus. In that experiment the signal signature is the observation of an isolated 105 MeV electron. Simulations of the $\mu \rightarrow e$ experiment based on the Geant [3] Monte Carlo indicate δ -rays produced by cosmic ray muons mimic the signal process at a rate many times the experimental sensitivity goal, motivating the need for an active cosmic ray
10 veto. Validating the $\mu \rightarrow e$ experiment design and estimating the resulting experimental sensitivity thus depends crucially on understanding the δ -ray production cross section.

A quantitative theoretical explanation of δ -ray production was derived in the 1930s from quantum electrodynamics [4], and is incorporated in the δ -ray
15 production model used by the Geant Monte Carlo [3]. This theory and has been shown to agree with measurements for low energy (< 1 MeV) electrons [5], however it has not been verified experimentally for δ -rays in the energy range relevant to the $\mu \rightarrow e$ experiment.

2. The *BABAR* Detector and Data Sample

20 The *BABAR* experiment operated for 7 years at the PEP-II collider at SLAC [6]. This provided approximately ** accumulated luminosity at the $\Upsilon(4S)$. Additional off-resonance samples and samples at other Υ resonances were recorded but not used in this analysis. The process $e^+e^- \rightarrow \mu^+\mu^-$ produced a sample of over 10^9 muons between roughly 3 and 9 GeV.

25 The *BABAR* detector was designed to record the results of *epem* collisions at a center of mass energy near the Upsilon resonances. The *BABAR* beampipe, Silicon Vertex Tracker (SVT) layers, Carbon fiber Support Tube (CST), and Drift Chamber (DCH) inner wall each presented a few % of a radiation length of thin material to normal incident muons, providing radiators for δ - ray pro-
30 duction. A detailed description of the material composition of these structures is given in table ??, while more general information about the *BABAR* detector can be found in [1].

Charged particles in *BABAR* produced signals in the SVT and DCH, which were reconstructed as tracks using standard pattern recognition algorithms.
35 High transverse momentum ($p_t > 100 \text{ MeV}/c$) were identified primarily in the DCH, while low transverse momentum tracks were found primarily in the SVT. Tracks identified in the DCH (SVT) were extrapolated into the SVT (DCH) to produce complete tracks. Final tracks were fit with a Kalman filter that compensated for the effects of the energy lost in traversed materials and the
40 inhomogeneous *BABAR* magnetic field. Electrons with transverse momentum above 50 MeV/c (100 MeV/c) from the e^+e^- interaction point were reconstructed with $> 50\%$ efficiency by the SVT (DCH), respectively. Details of the *BABAR* track reconstruction algorithms and its performance can be found in [7].

3. Monte Carlo Simulation

45 To simulate the $e^+e^- \rightarrow \mu^+\mu^-$ process we use the *KK2F* [8] generator, which includes the effects of initial and final state radiation (** more is needed here). The initial electron and positron 4-momenta are selected randomly according to the distributions observed in the *BABAR* $\Upsilon(4S)$ runs, as described in ??.

Figure ?? plots the lab-frame energy of muons from $e^+e^- \rightarrow \mu^+\mu^-$ predicted by the *kk2f* generator. Only events in which both muons are produced in the active tracking volume are included. Overlaid on figure ?? is the expected momentum spectrum of cosmic ray muons that generate signals in the $\mu \rightarrow e$ tracker, as predicted by the Daya Bay experiment (need a real reference here**) [9] generator and the $\mu \rightarrow e$ simulation [2]. A reasonable overlap is seen.

55 The interaction of particles with the *BABAR* detector is simulated using a detailed model built using the Geant4 toolkit [3]. Detector components are modeled assuming the ideal (design) geometry with sim

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simplified shapes, using materials mixtures with relative elemental composition and masses set to agree with measurements made during construction, averaged over the relevant volumes. For instance, the beampipe in the region of the *epem* interaction point is modeled as a cylinder of 2.5 cm radius with 2 mm wall thickness, as a material mix composed mostly of Beryllium and water, along with trace elements from platings and coatings, as listed in table 1. Detailed features such as the interior vanes used to direct the cooling water flow were not modeled. The material mix was determined by integrating the volume of each components and dividing by the total volume. While the dimensions of the Beryllium components were well measured during detector construction, the coating thicknesses varied spatially, introducing some uncertainty. In addition, the exact composition of one of the coatings [10] was proprietary to the corporation which produced it, and our values were based on oral communication. Similarly, the SVT wafers are modeled as parallelepipeds with dimensions set to the average of those measured during construction [1]. The material composition of the flex circuits was estimated by dissolving a spare circuit in acid, and precipitating out and weighing the metals.

The Geant4 δ -ray production is simulated based on (more here). Figure ?? plots the δ -ray spectrum produced in a simple Geant4 Monte Carlo simulation of a mono-energetic 4 GeV muon traversing 2 mm of pure Beryllium. Overlaid

Table 1: This is a table

| | | |
|---|---|---|
| A | B | C |
|---|---|---|

is a theoretical calculation based on [4],

showing good agreement in the range covered. Note there are no free pa-
105 rameters in this comparison.

The *BABAR* detector response simulation was tuned to reproduce detailed
models of the low-level sensor response measurements, as described in [7] chapter
**. Dead channels in the detector were monitored and recorded coherently as
a function of time. Simulated data samples were produced suppressing a given
110 set of measured dead channels, with event counts proportional to the luminosity
recorded in the period in which that set was dead.

The same reconstruction algorithms were run on simulated data as on *BABAR*
data. The efficiency for reconstructing tracks in *BABAR* data and simulation
has been found to agree within the measured uncertainties over the range of
115 momentum relevant to this measurement using well-defined test samples [11].

5. δ – ray Reconstruction

5.1. $e^+e^- \rightarrow \mu^+\mu^-$ Event Selection

5.2. δ – ray Candidate Identification

6. δ – ray Corrections

120 6.1. *Backgrounds*

6.2. *Acceptance and Efficiency*

6.3. *Energy Correction*

6.4. *Material Transit*

6.5. *Material Misidentification*

¹²⁵ **7. δ – ray Cross Section**

In section ?? we discussed....”

in this paper we've shown blah blah.

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