Candidate for the crystal $B(-\infty)$ for the queer Lie superalgebra

Ben Salisbury and Travis Scrimshaw

Abstract It is shown that the direct limit of the semistandard decomposition tableau model for polynomial representations of the queer Lie superalgebra exists, which is believed to be the crystal for the upper half of the corresponding quantum group. An extension of this model to describe the direct limit combinatorially is given. Furthermore, it is shown that the polynomial representations may be recovered from the limit in most cases.

1. Introduction

In the 1990s, Kashiwara began the study of crystals, a combinatorial skeleton of a quantum group representation $U_q(\mathfrak{g})$, where \mathfrak{g} is a symmetrizable Kac–Moody algebra. (See also Lusztig's canonical basis in [30] and Littelmann's path model in [28], [29].) Kashiwara [22], [23] showed that irreducible highest-weight representations have crystal bases $B(\lambda)$ and that the lower half of the quantum group has a crystal basis $B(\infty)$. The global crystal basis (general q version of a crystal basis) is known (see [17]) to equal the canonical basis introduced by Lusztig [30]. Kashiwara also proved that the direct limit of $B(\lambda)$ is isomorphic to $B(\infty)$, and one can recover $B(\lambda)$ by cutting a part of $B(\infty)$ by taking the tensor product with a specific crystal \mathcal{R}_{λ} . Using the direct limit, numerous combinatorial models for $B(\infty)$ have been developed, such as (marginally) large tableaux (see [12], [20]) and rigged configurations (see [33]–[35]).

For Lie superalgebras, there are two natural analogues of $\mathfrak{gl}(n)$. The first is the general linear Lie superalgebra $\mathfrak{gl}(m|n)$, where crystal bases have been constructed for the polynomial representations (see [2]), Kac modules (see [27]), and $B(\infty)$ for $\mathfrak{gl}(m|1)$ (see [11]). Furthermore, the character theory for $\mathfrak{gl}(m|n)$ has been well studied (see [4], [10], [37], [39], [41]), with connections to quasisymmetric functions (see [26]). The other is the queer superalgebra $\mathfrak{q}(n)$. The tensor powers of the fundamental representation form a semisimple category (see [16]), the irreducible representations are called the *polynomial representations*, and crystal bases of these irreducible representations have been constructed using semistandard decomposition tableaux (see [14], [15]). The character theory of

 $\mathfrak{q}(n)$ has also been studied (see [5], [8], [40]). In particular, the characters of the polynomial representations are Schur P-functions, which (along with the closely related Schur Q-functions) are of broad interest with other crystal connections (see [3], [18], [19], [21], [36]). Recently, a local characterization of the crystals for polynomial representations of $\mathfrak{q}(n)$ was given (see [1], [13]) in analogy to the Stembridge axioms (see [38]).

The goal of this paper is to construct the direct limit of the crystals of polynomial representations. We call this limit $B(-\infty)$, as we are considering the polynomial representations as lowest-weight representations and taking the corresponding limit. We believe this to be the crystal basis of the contragredient dual of the q-Weyl module W(0) defined in [16, Section 4],¹ in parallel to the case for $\mathfrak{gl}(n)$, where $B(-\infty)$ is the crystal basis of the contragredient dual of the Verma module M(0). Subsequently, this should be a combinatorial model for the crystal basis of the upper half of the corresponding quantum group $U_q(\mathfrak{q}(n))$ and also for the contragredient dual of the q-Weyl module $W(\lambda)$ after shifting the weight by λ . This claim is partially supported by a character calculation and an explicit combinatorial characterization of the lowest-weight elements in $B(-\infty)$. (See Proposition 3.15 and Proposition 3.18.)

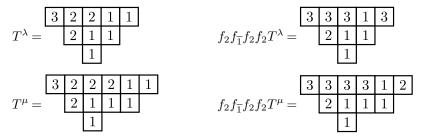
Our method follows the construction of (marginally) large tableaux for semi-standard decomposition tableaux SDT(λ) by showing that one can construct a directed system and that the $\mathfrak{q}(n)$ -crystal operators respect enlarging the shape. (See Lemma 3.12 and Corollary 3.13.) We then identify elements in each SDT(λ) based on their distance from the lowest-weight element and take a distinguished representative. Thus, our model is this limit crystal SDT($-\infty$) of distinguished representatives that we call dual marginally large semistandard decomposition tableaux. (See Theorem 3.14.) Our other main result (see Theorem 3.21) describes how we can recover SDT(λ) from SDT($-\infty$) using a dual version of \mathcal{R}_{λ} in the case when λ corresponds to a strict partition that has maximal length. We expect this to generalize to the case when λ does not have maximal length by a modification of the tensor product rule, which we also expect to construct crystals for dual polynomial representations.

It is noteworthy that one cannot take the limit of the polynomial representations considered as highest-weight representations. Indeed, if we consider the shapes $\lambda = (5,3,1)$ and $\mu = (6,4,1)$, then, for the direct limit, we must have an inclusion $B(\lambda) \hookrightarrow B(\mu)$. However, if $T^{\lambda} \in B(\lambda)$ and $T^{\mu} \in B(\mu)$ are the unique highest-weight elements of weight λ and μ , respectively, then

$$f_{\overline{1}}f_2f_{\overline{1}}f_2f_2T^\lambda\neq\mathbf{0},\qquad f_{\overline{1}}f_2f_{\overline{1}}f_2f_2T^\mu=\mathbf{0}.$$

Roughly speaking, the reason this inclusion fails is because adding entries to go from T^{λ} to T^{μ} does not behave well with respect to the crystal operators $e_{\overline{1}}$ and $f_{\overline{1}}$. Indeed, we have

¹The modules as $q \to 1$ were called *Verma modules* in [9, Section 2.1.6].



and $f_{\overline{1}}$ wants to change a 1 to a 2 in the rightmost entry of the topmost row. However, the 2 that was added to the top row of T^{μ} led to an extra 2 in the top row of $f_2f_{\overline{1}}f_2f_2T^{\mu}$, which breaks the inclusion. Thus there is no direct limit of $B(\lambda)$ as highest-weight crystals, in sharp contrast to $\mathfrak{gl}(n)$, where we can take both limits.

2. Background

Partitions and tableaux will be written using English convention. For brevity, write a column of height h as

$$\begin{bmatrix} x_1 \\ x_2 \\ \vdots \\ \vdots \\ x_h \end{bmatrix} = [x_1, x_2, \dots, x_h]^\top$$

2.1. Crystals for the superalgebra q(n)

Let $I_0 = \{1, \ldots, n-1\}$ and $I = I_0 \sqcup \{\overline{1}\}$. Denote the standard basis vectors of \mathbb{Z}^n by $\epsilon_1, \ldots, \epsilon_n$, and define the *simple root* $\alpha_i = \epsilon_i - \epsilon_{i+1}$ for each $i \in I_0$. Let $Q = \bigoplus_{i \in I_0} \mathbf{Z}\alpha_i$ be the root lattice, and let $Q^+ = \bigoplus_{i \in I_0} \mathbf{Z}_{\geq 0}\alpha_i$ and $Q^- = \bigoplus_{i \in I_0} \mathbf{Z}_{\leq 0}\alpha_i$ be the positive and negative root cones. Let $\Phi^{\pm} = \{\pm(\epsilon_i - \epsilon_j) : 1 \leq i < j \leq n\}$ be the set of positive and negative roots. Set

$$\Lambda^{-} = \left\{ \lambda = -\lambda_{1} \epsilon_{1} - \dots - \lambda_{n} \epsilon_{n} \in \mathbf{Z}_{\leq 0}^{n} : \lambda_{i} \geq \lambda_{i+1} \text{ and } \lambda_{i} = \lambda_{i+1} \text{ implies} \\ \lambda_{i} = \lambda_{i+1} = 0 \text{ for all } i = 1, \dots, n-1 \right\}.$$

Equip Λ^- with a partial order $\lambda \leq \mu$ if and only if $\mu - \lambda \in \Lambda^-$. An element $\lambda = -\lambda_1 \epsilon_1 - \dots - \lambda_n \epsilon_n \in \Lambda^-$ will be henceforth be identified with the strict partition $w_0 \lambda = (\lambda_n, \dots, \lambda_1)$.

DEFINITION 2.1 ([24, Definition 1.2.1])

An abstract $\mathfrak{gl}(n)$ -crystal is a set \mathcal{B} together with maps $e_i, f_i : \mathcal{B} \longrightarrow \mathcal{B} \sqcup \{0\}$, $\varphi_i, \varepsilon_i : \mathcal{B} \longrightarrow \mathbf{Z} \sqcup \{-\infty\}$, for $i \in I_0$, and wt: $\mathcal{B} \longrightarrow \mathbf{Z}^n$ satisfying the following conditions:

²This is not the usual (opposite) dominance order, where the result should be in Q^- , but instead should be thought of as the lex order.

- (1) for any $i \in I_0$ such that $e_i b \neq \mathbf{0}$, we have $\operatorname{wt}(e_i b) = \operatorname{wt}(b) + \alpha_i$;
- (2) for any $i \in I_0$ such that $f_i b \neq \mathbf{0}$, we have $\operatorname{wt}(f_i b) = \operatorname{wt}(b) \alpha_i$;
- (3) for any $i \in I_0$ and $b \in \mathcal{B}$, we have $\varphi_i(b) = \varepsilon_i(b) + \operatorname{wt}_i(b)$;
- (4) for any $i \in I_0$ and $b, b' \in \mathcal{B}$, we have $f_i b = b'$ if and only if $b = e_i b'$;
- (5) for any $i \in I_0$ and $b \in \mathcal{B}$ such that $e_i b \neq \mathbf{0}$, we have $\varepsilon_i(e_i b) = \varepsilon_i(b) 1$ and $\varphi_i(e_i b) = \varphi_i(b) + 1$;
- (6) for any $i \in I_0$ and $b \in \mathcal{B}$ such that $f_i b \neq \mathbf{0}$, we have $\varepsilon_i(f_i b) = \varepsilon_i(b) + 1$ and $\varphi_i(f_i b) = \varphi_i(b) 1$;
 - (7) for any $i \in I_0$ and $b \in \mathcal{B}$ such that $\varphi_i(b) = -\infty$, we have $e_i b = f_i b = \mathbf{0}$.

In the above definition (and throughout), the notation $\operatorname{wt}_i(b)$ means $\operatorname{wt}_i(b) = \mu_i - \mu_{i+1}$ provided that $\operatorname{wt}(b) = \sum_{i \in I_0} \mu_i \epsilon_i$.

DEFINITION 2.2 ([15, Definition 1.9])

An abstract $\mathfrak{gl}(n)$ -crystal is an abstract $\mathfrak{gl}(n)$ -crystal \mathcal{B} together with maps $e_{\overline{1}}, f_{\overline{1}} \colon \mathcal{B} \longrightarrow \mathcal{B} \sqcup \{\mathbf{0}\}$ such that

- (1) $\operatorname{wt}(\mathcal{B}) \subset \mathbf{Z}_{\geq 0}^n$;
- (2) $\operatorname{wt}(e_{\overline{1}}b) = \operatorname{wt}(b) + \alpha_1$ provided that $e_{\overline{1}}b \neq \mathbf{0}$;
- (3) $\operatorname{wt}(f_{\overline{1}}b) = \operatorname{wt}(b) \alpha_1$ provided that $f_{\overline{1}}b \neq \mathbf{0}$;
- (4) for any $b, b' \in \mathcal{B}$, $f_{\overline{1}}b = b'$ if and only if $b = e_{\overline{1}}b'$;
- (5) if $3 \le i \le n-1$, then we have that
- (a) the operators $e_{\overline{1}}$ and $f_{\overline{1}}$ commute with e_i and f_i , and
- (b) if $e_{\overline{1}}b \in \mathcal{B}$, then $\varepsilon_i(e_{\overline{1}}b) = \varepsilon_i(b)$ and $\varphi_i(e_{\overline{1}}b) = \varphi_i(b)$.

Let \mathcal{B} and \mathcal{C} be abstract $\mathfrak{q}(n)$ -crystals. A crystal morphism is a map $\psi \colon \mathcal{B} \longrightarrow \mathcal{C} \sqcup \{\mathbf{0}\}$ such that

- (1) if $b \in \mathcal{B}$ and $\psi(b) \in \mathcal{C}$, then for all $i \in I_0$, $\operatorname{wt}(\psi(b)) = \operatorname{wt}(b), \qquad \varepsilon_i(\psi(b)) = \varepsilon_i(b), \qquad \varphi_i(\psi(b)) = \varphi_i(b);$
- (2) for $b \in \mathcal{B}$ and $i \in I$, we have $\psi(e_i b) = e_i \psi(b)$ provided that $\psi(e_i b) \neq \mathbf{0}$ and $e_i \psi(b) \neq \mathbf{0}$;
- (3) for $b \in \mathcal{B}$ and $i \in I$, we have $\psi(f_i b) = f_i \psi(b)$ provided that $\psi(f_i b) \neq \mathbf{0}$ and $f_i \psi(b) \neq \mathbf{0}$.

A morphism ψ is called *strict* if ψ commutes with e_i and f_i for all $i \in I$. Moreover, a morphism $\psi \colon \mathcal{B} \longrightarrow \mathcal{C} \sqcup \{\mathbf{0}\}$ is called an *embedding* (resp., *isomorphism*) if the induced map $\psi \colon \mathcal{B} \longrightarrow \mathcal{C} \sqcup \{\mathbf{0}\}$ is injective (resp., bijective).

Again, let \mathcal{B} and \mathcal{C} be abstract $\mathfrak{q}(n)$ -crystals. The tensor product $\mathcal{B} \otimes \mathcal{C}$ is defined to be the Cartesian product $\mathcal{B} \times \mathcal{C}$ equipped with crystal operations defined, for $i \in I_0$, by

$$e_{i}(b \otimes c) = \begin{cases} e_{i}b \otimes c & \text{if } \varphi_{i}(c) < \varepsilon_{i}(b), \\ b \otimes e_{i}c & \text{if } \varphi_{i}(c) \geq \varepsilon_{i}(b), \end{cases}$$

$$e_{\overline{1}}(b \otimes c) = \begin{cases} b \otimes e_{\overline{1}}c & \text{if } e_{\overline{1}}b = f_{\overline{1}}b = \mathbf{0}, \\ e_{\overline{1}}b \otimes c & \text{otherwise}, \end{cases}$$

$$f_{i}(b \otimes c) = \begin{cases} f_{i}b \otimes c & \text{if } \varphi_{i}(c) \leq \varepsilon_{i}(b), \\ b \otimes f_{i}c & \text{if } \varphi_{i}(c) > \varepsilon_{i}(b), \end{cases}$$

$$f_{\overline{1}}(b \otimes c) = \begin{cases} b \otimes f_{\overline{1}}c & \text{if } e_{\overline{1}}b = f_{\overline{1}}b = \mathbf{0}, \\ f_{\overline{1}}b \otimes c & \text{otherwise}. \end{cases}$$

REMARK 2.3

This is equivalent to the rule in [13, Remark 2.4] in the situations we consider. Moreover, this is the reverse convention of the tensor product to that given in [14].

The remaining crystal structure of $\mathcal{B} \otimes \mathcal{C}$ is defined as

$$\varepsilon_i(b \otimes c) = \max(\varepsilon_i(c), \varepsilon_i(b) - \operatorname{wt}_i(c)),$$

$$\varphi_i(b \otimes c) = \max(\varphi_i(b), \varphi_i(c) + \operatorname{wt}_i(b)),$$

$$\operatorname{wt}(b \otimes c) = \operatorname{wt}(b) + \operatorname{wt}(c).$$

Following the method of [25, p. 74], one can construct direct limits in the category of abstract $\mathfrak{q}(n)$ -crystals. Indeed, let $\{\mathcal{B}_j\}_{j\in J}$ be a directed system of crystals, and let $\psi_{k,j}\colon \mathcal{B}_j\longrightarrow \mathcal{B}_k,\ j\leq k$, be a crystal embedding (with $\psi_{j,j}$ being the identity map on \mathcal{B}_j) such that $\psi_{k,j}\psi_{j,i}=\psi_{k,i}$. Let $\vec{\mathcal{B}}=\varinjlim \mathcal{B}_j$ be the direct limit of this system, and let $\psi_j\colon \mathcal{B}_j\longrightarrow \vec{\mathcal{B}}$. Then $\vec{\mathcal{B}}$ has a crystal structure induced from the crystals $\{\mathcal{B}_j\}_{j\in J}$. Indeed, for $\vec{b}\in \vec{\mathcal{B}}$ and $i\in I$, define $e_i\vec{b}$ to be $\psi_j(e_ib_j)$ if there exists $b_j\in \mathcal{B}_j$ such that $\psi_j(b_j)=\vec{b}$ and $e_i(b_j)\neq 0$. This definition does not depend on the choice of b_j . If there is no such b_j , then set $e_i\vec{b}=0$. The definition of $f_i\vec{b}$ is similar. Moreover, the functions wt, ε_i , and φ_i on \mathcal{B}_j extend to functions on $\vec{\mathcal{B}}$.

2.2. Semistandard decomposition tableaux

This section summarizes the results of [14] using the conventions of [13].

DEFINITION 2.4

Let $\eta = (\eta_n, \dots, \eta_1)$ be a strict partition. Define $|\eta| = \eta_1 + \dots + \eta_n$ and $\ell(\eta)$ to be the number of $1 \le i \le n$ such that $\eta_i \ne 0$.

- (1) The shifted Young diagram of shape η is an array of cells in which the ith row has η_{n+1-i} cells and is shifted i-1 units to the right with respect to the top row.
 - (2) A word $u = u_1 u_2 \cdots u_N$ is a *hook word* if there exists $1 \le k \le N$ such that

$$u_1 \ge u_2 \ge \cdots \ge u_k < u_{k+1} < \cdots < u_N.$$

- (3) A semistandard decomposition tableau of shifted shape η is a filling T of η with letters from $\{1, 2, ..., n\}$ such that
- (a) the word v_i formed by reading the *i*th row from left to right is a hook word of length η_{n-i+1} , and
- (b) v_i is a hook subword of maximal length in $v_{i+1}v_i$ for $1 \le i \le \ell(\eta) 1$, where the entries of v_i are taken consecutively.
- (4) Set $\operatorname{read}(T)$ to be the word obtained by reading T in rows from right to left starting at the top.³ Let $\operatorname{read}_{i}(T)$ denote the subword of $\operatorname{read}(T)$ consisting of only the letters i and i+1. Let $\overline{\operatorname{read}}_{i}(T)$ denote the subword of $\operatorname{read}(T)$ formed by successively removing all pairs of letters (i+1,i) (in that order).
- (5) For $\lambda \in \Lambda^-$, let $SDT(\lambda)$ denote the set of all semistandard decomposition tableaux of shape $w_0\lambda$.

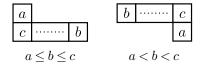
Note that $\overline{\text{read}}_i(T)$ is a word consisting of i's following by (i+1)'s.

We will require the following characterization of semistandard decomposition tableaux.

PROPOSITION 2.5 ([14, Proposition 2.3])

A tableau T is a decomposition tableau if

- (1) every row is a hook word,
- (2) the leftmost entry of a given row is strictly larger than every entry in the row below it, and
 - (3) neither of the following configurations occur:



The prohibited configurations in Proposition 2.5 will be referred to as $type\ L$ and $type\ U$, respectively.

DEFINITION 2.6

Let T be a semistandard decomposition tableau of shape $w_0\lambda$.

- (1) Suppose $i \in I_0$.
- (a) If there is no i+1 in $\overline{\text{read}}_i(T)$, then $e_iT = \mathbf{0}$. Otherwise, e_iT is the tableau obtained from T by changing the (i+1)-box corresponding to the leftmost i+1 in the subword above to an i-box.

³The choice of the reading word and the opposite order of the tensor product means that we obtain the same crystal as in [14]; i.e., the two reversals nullify the effects of each other.

- (b) If there is no i in $\overline{\text{read}}_i(T)$, then $f_iT = \mathbf{0}$. Otherwise, f_iT is the tableau obtained from T by changing the i-box corresponding to the rightmost i remaining in the subword above to an (i+1)-box.
 - (2) Suppose $i = \overline{1}$.
- (a) If the leftmost letter in $\operatorname{read}_1(T)$ is 1, then $e_{\overline{1}}T = 0$. Otherwise, $e_{\overline{1}}T$ is the tableau obtained from T by changing the 2-box corresponding to the leftmost 2 in $\operatorname{read}_1(T)$ to a 1-box.
- (b) If the leftmost letter in $\operatorname{read}_1(T)$ is 2, then $f_{\overline{1}}T = 0$. Otherwise, $f_{\overline{1}}T$ is the tableau obtained from T by changing the 1-box corresponding to the leftmost 1 in $\operatorname{read}_1(T)$ to a 2-box.

For a $\lambda \in \Lambda^-$ with $\ell(\lambda) = N$, define $L^{\lambda} \in \mathrm{SDT}(\lambda)$ to be the tableau whose ith row from the bottom contains only the letter i. As described in the next theorem, L^{λ} is a generator of $\mathrm{SDT}(\lambda)$, and a generalization of L^{λ} will serve as a generator for the candidate for $B(-\infty)$ in the next section.

EXAMPLE 2.7

Let n = 5 and $w_0 \lambda = (7, 4, 3, 2, 1)$. Then

	5	5	5	5	5	5	5
		4	4	4	4		
$L^{\lambda} =$			3	3	3		
<i>L</i> –				2	2		
					1		

An element $T \in \text{SDT}(\lambda)$ is called lowest weight (resp., highest weight) if $f_i T = \mathbf{0}$ (resp., $e_i T = \mathbf{0}$) for all $i \in I$.

THEOREM 2.8 ([14, Theorem 2.5])

For $\lambda \in \Lambda^-$, the set SDT(λ) together with the operators defined in Definition 2.6 form an abstract $\mathfrak{q}(n)$ -crystal isomorphic to the crystal of the irreducible highest-weight $\mathfrak{q}(n)$ -module with highest weight $w_0\lambda$. Moreover, SDT(λ) is generated by L^{λ} (and, hence, connected).

Let us briefly remark on the notation $w_0\lambda$. Typically w_0 denotes the long element of the symmetric group \mathfrak{S}_n , which acts naturally on \mathbf{Z}^n . From [14, Theorem 2.5], the lowest weight that appears in the character of $B(w_0\lambda)$ has weight $w_0 \cdot \lambda$, where now w_0 is acting on λ as an element of the Weyl group of $\mathfrak{gl}(n)$, up to shifting by $(1,\ldots,1)$. However, this does not change the structure of the crystals (up to this shift in the weight).

3. Main results

In this section, we prove our main results. To do so, we first give $\mathfrak{q}(n)$ -analogues of some auxiliary crystals used in the construction of $B(\infty)$. Next, an explicit

description of the direct limit $SDT(-\infty)$ is given, and this object is our proposed $B(-\infty)$. It is then shown that $\{SDT(\lambda) : \lambda \in \Lambda^-\}$ forms a directed system and is isomorphic to $SDT(-\infty)$. We conclude with a method to recover $SDT(\lambda)$ from $SDT(-\infty)$ in some cases.

3.1. Auxiliary crystals

There are two abstract crystals which are necessary for later work. The first of the two crystals defined next simply shift the weights of a given crystal when tensored, and the second will be used to construct $SDT(\lambda)$ from $SDT(-\infty)$.

DEFINITION 3.1

Let $\lambda = (\lambda_1, \dots, \lambda_n) \in \mathbf{Z}^n$.

(1) Define $\mathcal{T}_{\lambda} = \{t_{\lambda}\}$ with operations

$$\begin{split} e_i t_\lambda &= e_{\overline{1}} t_\lambda = f_i t_\lambda = f_{\overline{1}} t_\lambda = \mathbf{0}, \\ \varepsilon_i(t_\lambda) &= \varphi_i(t_\lambda) = -\infty, \\ \mathrm{wt}(t_\lambda) &= \lambda, \end{split}$$

for $i \in I_0$.

(2) Define $\mathcal{R}_{\lambda}^{\vee} = \{r_{\lambda}^{\vee}\}$ with operations

$$e_{i}r_{\lambda}^{\vee} = e_{\overline{1}}r_{\lambda}^{\vee} = f_{i}r_{\lambda}^{\vee} = f_{\overline{1}}r_{\lambda}^{\vee} = \mathbf{0},$$

$$\varepsilon_{i}(r_{\lambda}^{\vee}) = 0, \qquad \varphi_{i}(r_{\lambda}^{\vee}) = \lambda_{i} - \lambda_{i+1},$$

$$\operatorname{wt}(r_{\lambda}^{\vee}) = \lambda,$$

for $i \in I_0$.

The following is clear from the definitions.

LEMMA 3.2

If $\lambda \in \mathbf{Z}_{\geq 0}^n$, then \mathcal{T}_{λ} and $\mathcal{R}_{\lambda}^{\vee}$ are abstract $\mathfrak{q}(n)$ -crystals.

3.2. Candidate for $B(-\infty)$

DEFINITION 3.3

A semistandard decomposition tableau T for $\mathfrak{q}(n)$ is called *dual large* if

- (1) T has n rows, and
- (2) for all $1 \le i \le n$, the number of leftmost *i*-boxes in row n-i+1 is strictly greater than the total number of boxes in row n-i+2.

EXAMPLE 3.4

Consider the following semistandard decomposition tableaux of shape (4,3,1) for $\mathfrak{q}(3)$:

	3	3	3	3		3	3	3	3
dual large:		2	2	2	not dual large:		2	1	2
			1					1	

DEFINITION 3.5

A semistandard decomposition tableau T for $\mathfrak{q}(n)$ is called *dual marginally large* if

- (1) it has n rows, and
- (2) for all $1 \le i \le n$, the number of leftmost *i*-boxes in row n-i+1 is greater than the total number of boxes in row n-i+2 by exactly one.

Denote the set of all dual marginally large semistandard tableaux for $\mathfrak{q}(n)$ by $\mathrm{SDT}(-\infty)$.

We note that a dual marginally large tableau is dual large.

EXAMPLE 3.6

Consider the following semistandard decomposition tableaux for q(3):

	3	3	3	3		3	3	3
not dual marginally large:		2	2	2	dual marginally large:		2	2
			1					1

A column C of height h is called *trivial* if the column is $C = [n, n-1, \ldots, n+1-h]^{\top}$ and every entry to the left of an $i \in C$ is i (if it exists). We will say that we push in (resp., push out) a trivial column of height h from C if we add (resp., remove) a box with an i to the left of row n+1-i and sliding the remaining boxes back to a shifted shape. Furthermore, in the case that a tableau is dual large, the act of pushing in (resp., out) a column of height r is equivalent to adding (resp., removing) a trivial column of height r. However, we will need to consider pushing in/out columns for the case when the tableau is not dual large.

EXAMPLE 3.7

In the following, a trivial column of height 2 is pushed in to the given semistandard decomposition tableaux, where the pushed-in trivial column is represented by the shaded boxes:

	3	3	3	3		3	3	3	3	3		3	3	3	3	3
•		2	2	2	\longrightarrow		2	2	2	2	\longrightarrow		2	2	2	2
			1		•				1		•			1		

The following lemma is straightforward from the definitions, where the uniqueness follows from the diamond lemma.

LEMMA 3.8

Given any dual large tableau T, we can construct a unique dual marginally large tableau T_M by pushing out to the left trivial columns.

EXAMPLE 3.9

In Example 3.6, the tableau on the left has one trivial column: the column $[3,2]^{\top}$ on the right. Removing this column yields the dual marginally large tableau on the right in the same example.

DEFINITION 3.10

Let $T \in \text{SDT}(-\infty)$ for $\mathfrak{q}(n)$.

- (1) Suppose $i \in I_0$.
- (a) Let T' be the tableau obtained from T by changing the (i+1)-box corresponding to the leftmost $i+1 \in \overline{\mathrm{read}}_i(T)$ to an i-box. If T' is dual marginally large, then $T' = e_i T$. Otherwise, let T'' be the tableau obtained from T' by adding a (n-k+1)-box in row k, for each $1 \le k \le n-i+1$. Then $T'' = e_i T$.
- (b) If there is no such $i \in \overline{\mathrm{read}}_i(T)$, then $f_iT = \mathbf{0}$. Otherwise, let T' be the tableau obtained from T by changing the i-box corresponding to the rightmost $i \in \overline{\mathrm{read}}_i(T)$ to an (i+1)-box. If T' is dual marginally large, then $T' = f_iT$. Otherwise, let T'' be the tableau obtained from T' by removing a (n-k+1)-box in row k, for each $1 \le k \le n-i+1$. Then $T'' = f_iT$.
- (2) Both $e_{\overline{1}}$ and $f_{\overline{1}}$ are defined exactly as in Definition 2.6, except for the need to maintain the dual marginally large condition as in (1a) and (1b) above.

EXAMPLE 3.11

Let n=3 and

$$T = \begin{array}{c|cccc} 3 & 3 & 3 & 3 & 2 \\ \hline 2 & 2 & 1 & \in SDT(-\infty) \\ \hline & 1 & & \end{array}$$

Then read(T) = 233331221. After pairing off all possible (2,1), there is no 1 remaining and the leftmost 2 remaining corresponds to the underlined $\underline{2}$ in $233331\underline{2}21$. Hence $f_1T = \mathbf{0}$, but

Note that a 3-box needed to be added to the first row and a 2-box needed to be added to the second row to maintain the dual marginally large condition in e_1T . In other words, we pushed in a trivial column of height 2 (from the left).

After pairing off all possible (3,2) in read(T), the leftmost 3 and rightmost 2 remaining are the underlined letters in $\underline{23}3331221$. Hence

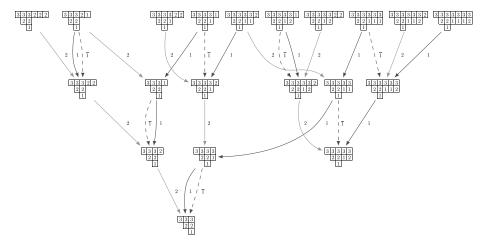


Figure 1. (Color online) A bottom portion of the q(3)-crystal SDT $(-\infty)$ containing the lowest-weight element $L^{-\infty}$ created using SAGEMATH [32].

Note that a 3-box needed to be added to the first row of e_2T and removed from the first row of f_2T to maintain the dual marginally large condition. Note that a 3-box is a trivial column of height 1.

A diagram of the crystal graph $SDT(-\infty)$ up to height 3 is included in Figure 1.

Now we show that the standard decomposition tableaux form a directed system.

LEMMA 3.12

Suppose $\lambda, \mu \in \Lambda^-$. Then there exists a $\mathfrak{q}(n)$ -crystal embedding

$$v_{\lambda,\lambda+\mu} : \mathrm{SDT}(\lambda) \otimes \mathcal{T}_{-\lambda} \hookrightarrow \mathrm{SDT}(\lambda+\mu) \otimes \mathcal{T}_{-\lambda-\mu}$$

such that $L^{\lambda} \otimes t_{-\lambda} \mapsto L^{\lambda+\mu} \otimes t_{-\lambda-\mu}$.

Proof

For $T \in \text{SDT}(\lambda)$, define E(T) to be the tableau obtained from T by adding μ_{n+1-i} *i*-boxes on the left to row n+1-i, for each $1 \le i \le n$. It is evident that the result has shape $w_0(\lambda + \mu)$. Define

$$v_{\lambda,\lambda+\mu} : \mathrm{SDT}(\lambda) \otimes \mathcal{T}_{-\lambda} \hookrightarrow \mathrm{SDT}(\lambda+\mu) \otimes \mathcal{T}_{-\lambda-\mu}$$

by $(T \otimes t_{-\lambda}) \mapsto E(T) \otimes t_{-\lambda-\mu}$. Immediately, we have $v_{\lambda,\lambda+\mu}(L^{\lambda} \otimes t_{-\lambda}) = L^{\lambda+\mu} \otimes t_{-\lambda-\mu}$ and

$$\varepsilon_i(L^{\lambda}\otimes t_{-\lambda})=\varepsilon_i(L^{\lambda+\mu}\otimes t_{-\lambda-\mu}), \qquad \varphi_i(L^{\lambda}\otimes t_{-\lambda})=\varphi_i(L^{\lambda+\mu}\otimes t_{-\lambda-\mu}),$$

(which implies $\operatorname{wt}_i(L^{\lambda} \otimes t_{-\lambda}) = \operatorname{wt}_i(L^{\lambda+\mu} \otimes t_{-\lambda-\mu})$ for $i \in I_0$, by the crystal axiom (3)). Since the crystal $\operatorname{SDT}(\nu) \otimes t_{-\nu}$ is generated by $L^{\nu} \otimes t_{-\nu}$, the crystal axioms imply it suffices to show that $E(x_iT) = x_iE(T)$ for all $x_i = e_i$, f_i and $T \in \operatorname{SDT}(\lambda)$ such that $x_iT \neq 0$.

Fix some $T \in \mathrm{SDT}(\lambda)$ and $i \in I_0$. From the definition of the crystal operators, we need to show that the difference between $\overline{\mathrm{read}}_i(T)$ and $\overline{\mathrm{read}}_i(E(T))$ is possibly some additional i's on the left (resp., (i+1)'s on the right) corresponding to one of the added i-boxes (resp., (i+1)-boxes) in E(T). By induction and the construction, it is sufficient to consider $w_0\mu$ being a column of height h. The case when h < n - i is trivial, as no i nor i + 1 is added to T.

We note that Proposition 2.5(2) implies that every entry in row n+1-i cannot have a value larger than i, as otherwise the first row would have a value of n+1 or larger. Hence the added i+1 (to row n-i) will either be the rightmost uncanceled i+1 or will cancel with some i. If the added i+1 in row n-i cancels with an i (which must be in row n-i+1) in E(T), then there must exist an i+1 in the first position of row n-i in i, as otherwise this would violate Proposition 2.5(2). If there is no other i in row i in row i in the added i in i in row i in

$$\operatorname{read}_{i}(T) = w' \underbrace{i+1, \dots, i+1}_{k} \underbrace{i, \dots, i}_{k},$$
$$\operatorname{read}_{i}(E(T)) = w' \underbrace{i+1, \dots, i+1, \underline{i+1}}_{k+1} \underbrace{i, \dots, i, \underline{i}}_{k+1},$$

where the underlined entries are those added to form E(T) and w' is some word unaffected when constructing E(T). Thus we cannot obtain an additional i on the right of read_i(E(T)) coming from the i-box added to row n-i+1.

To show $E(f_{\overline{1}}T)=f_{\overline{1}}E(T)$, note that the only problem that can arise is when pushing in a trivial column that contains a 2 to some $T\in \mathrm{SDT}(\lambda)$ such that the only 1 is in row n but no 2 in row n-1 of T (as this would make $f_{\overline{1}}E(T)=\mathbf{0}$). However, all entries in row n-1 must be at most 2 as a consequence of Proposition 2.5(2) as shown above. Thus there must be a 2 in row n-1, which shows we must have $E(f_{\overline{1}}T)=f_{\overline{1}}E(T)$. For $E(e_{\overline{1}}T)=e_{\overline{1}}E(T)$, note that the only way this could not happen is if there exists a 2 in row n and we push in a trivial column containing a 2 (as this would come before the 2 we act upon for $e_{\overline{1}}T$). However, this is impossible, as it follows from Proposition 2.5(2) that there cannot exist a 2 in row n.

COROLLARY 3.13

The collection $\{SDT(\lambda) \otimes \mathcal{T}_{-\lambda}\}_{\lambda \in \Lambda^-}$ together with the inclusion maps from Lemma 3.12 form a directed system.

To prove the corollary, one makes repeated use of Lemma 3.12 applied to diagrams of the following form:

$$SDT(\lambda) \otimes \mathcal{T}_{-\lambda} \xrightarrow{\upsilon_{\lambda+\mu,\lambda}} SDT(\lambda+\mu) \otimes \mathcal{T}_{-\lambda-\mu}$$

$$\downarrow \upsilon_{\lambda+\mu+\xi,\lambda} \qquad \qquad \qquad \downarrow \upsilon_{\lambda+\mu+\xi,\lambda+\mu}$$

$$SDT(\lambda+\mu+\xi) \otimes \mathcal{T}_{-\lambda-\mu-\xi}.$$

Define $L^{-\infty}$ to be the decomposition tableau $L^{-n\epsilon_1-(n-1)\epsilon_2-\cdots-\epsilon_n}$.

THEOREM 3.14

The set $SDT(-\infty)$ together with e_i , f_i from Definition 3.10 is an abstract $\mathfrak{q}(n)$ crystal such that

$$\mathrm{SDT}(-\infty) \cong \varinjlim_{\lambda \in \Lambda^-} \mathrm{SDT}(\lambda) \otimes \mathcal{T}_{-\lambda}.$$

Proof

Let \overrightarrow{SDT} denote the direct limit. We claim the map $\psi \colon SDT(-\infty) \longrightarrow \overrightarrow{SDT}$ given by $\psi(T) = [T]$ is the desired crystal isomorphism.

For any $T \in SDT(-\infty)$ of shape $w_0\lambda$, we note that there exists a projection

$$\pi_T : \mathrm{SDT}(-\infty) \longrightarrow \mathrm{SDT}(\lambda) \otimes \mathcal{T}_{-\lambda}$$

such that for any $T' = f_{i_1} \cdots f_{i_\ell} T$, we can form $\pi_T(T')$ by adding in suitably many trivial columns to T' until we obtain the shape $w_0 \lambda$ (and adding the inconsequential tensor factor $t_{-\lambda}$). The latter follows by induction since $\pi_T(f_i T') = f_i \pi_T(T')$, where this equality was shown in the proof of Lemma 3.12 as f_i commutes with the procedure of adding or removing trivial columns. Therefore, we have $\psi(f_i T) = [f_i T] = f_i [T] = f_i \psi(T)$.

Next, it was shown in the proof of Lemma 3.12 that any two $\tau, \tau' \in [T]$ differ only by trivial columns. Hence, we have $\psi(e_iT) = e_i\psi(T)$. It is clear that ψ satisfies the rest of the properties of being a crystal morphism. Since there is a unique dual marginally large seminstandard decomposition tableau in each class $[T] \in \overrightarrow{\mathrm{SDT}}$, we have that ψ is a bijection.

3.3. Comparison of characters

Let us provide some evidence that $SDT(-\infty)$ is a combinatorial model of $B(-\infty)$. We show that the character of $SDT(-\infty)$ agrees with that of the upper half of the quantum group $U_q^+(\mathfrak{q}(n))$.

PROPOSITION 3.15

We have

$$\operatorname{ch} \operatorname{SDT}(-\infty) = \operatorname{ch} U_q^+(\mathfrak{q}(n)).$$

Proof

From [9, Section 2.3.1], it is straightforward to see that

(3.1)
$$\operatorname{ch} U_q^+(\mathfrak{q}(n)) = \operatorname{ch} U_q^+(\mathfrak{gl}(n)) \prod_{\alpha \in \Phi^+} (1 + e^{\alpha}) = \frac{\prod_{\alpha \in \Phi^+} (1 + e^{\alpha})}{\prod_{\alpha \in \Phi^+} (1 - e^{\alpha})},$$

where Φ^+ is set of positive roots associated to $\mathfrak{gl}(n)$ (see also [6, Section 3]).⁴ Next, for $\lambda \in \Lambda^-$ from [7] and [31], we have

$$\operatorname{ch} \operatorname{SDT}(\lambda) = \frac{\prod_{\alpha \in \Phi^+} (1 + e^{-\alpha})}{\prod_{\alpha \in \Phi^+} (1 - e^{-\alpha})} \sum_{w \in \mathfrak{S}_n} (-1)^{\ell(w)} w \left(\frac{e^{w_0 \lambda}}{\prod_{\alpha \in \Phi^+(w_0 \lambda)} (1 + e^{-\alpha})} \right),$$

where $\Phi^+(\mu) := \{ \epsilon_i - \epsilon_j : \mu_i = \mu_j \ (i < j) \}$ and \mathfrak{S}_n is the symmetric group on n letters. For the direct limit, it is sufficient to take $\lambda_1 > \lambda_2 > \cdots > \lambda_n > 0$, and thus we have $\Phi^+(w_0\lambda) = \emptyset$. It follows that

$$\operatorname{ch}(\operatorname{SDT}(\lambda) \otimes \mathcal{T}_{-\lambda}) = \frac{\prod_{\alpha \in \Phi^{+}} (1 + e^{-\alpha})}{\prod_{\alpha \in \Phi^{+}} (1 - e^{-\alpha})} e^{-\lambda} \sum_{w \in \mathfrak{S}_{n}} (-1)^{\ell(w)} w e^{w_{0} \lambda}$$
$$= \frac{\prod_{\alpha \in \Phi^{+}} (e^{\alpha} + 1)}{\prod_{\alpha \in \Phi^{+}} (e^{\alpha} - 1)} \sum_{w' \in \mathfrak{S}_{n}} (-1)^{\ell(w'w_{0})} e^{w'\lambda - \lambda},$$

where we rewrite the sum by $w' = ww_0$. Next, we recall that $|\Phi^+| = \ell(w_0)$ since the length of a permutation equals the number of inversions (which are positive roots sent to a negative root), we have

$$\operatorname{ch}(\operatorname{SDT}(\lambda) \otimes \mathcal{T}_{-\lambda}) = \frac{\prod_{\alpha \in \Phi^+} (1 + e^{\alpha})}{\prod_{\alpha \in \Phi^+} (1 - e^{\alpha})} \sum_{w \in \mathfrak{S}_n} (-1)^{\ell(w)} e^{w\lambda - \lambda}.$$

Define $\rho = \sum_{i \in I_0} \Lambda_i$. Now for any $\mu \in Q^+$, we note that there exists a k such that for all $\lambda > -k\rho$ the coefficient of e^{μ} is 0 in

$$\frac{\prod_{\alpha \in \Phi^+} (1 + e^{\alpha})}{\prod_{\alpha \in \Phi^+} (1 - e^{\alpha})} \sum_{w \in \mathfrak{S}_n \setminus \{1\}} (-1)^{\ell(w)} e^{w\lambda - \lambda},$$

or, alternatively, μ is in the interior of the convex hull of the points $\mathfrak{S}_n(-k\rho)$. Hence, we have

$$\operatorname{ch} \operatorname{SDT}(-\infty) = \operatorname{ch}\left(\varinjlim_{\lambda \in \Lambda^{-}} \operatorname{SDT}(\lambda) \otimes \mathcal{T}_{-\lambda}\right) = \frac{\prod_{\alpha \in \Phi^{+}} (1 + e^{\alpha})}{\prod_{\alpha \in \Phi^{+}} (1 - e^{\alpha})} = \operatorname{ch} U_{q}^{+}(\mathfrak{q}(n)). \quad \Box$$

⁴Note that the character of a Verma module of weight 0 after substituting $\epsilon_i \mapsto -\epsilon_i$ is equal to $\operatorname{ch} U_q^+(\mathfrak{q}(n))$.

We can prove a stronger statement, that the lowest-weight vectors, and hence the $U_q(\mathfrak{gl}(n))$ -highest-weight decomposition, of $SDT(-\infty)$ is given by (3.1).

DEFINITION 3.16

Call the leftmost j-boxes of row j the trivial boxes. Let $\operatorname{read}_{nt}(T)$ denote the subword of $\operatorname{read}(T)$ coming from all nontrivial boxes. Let T be a marginally large tableau and $1 \leq i \leq k \leq n$. Define an (i,k)-consecution of T to be a subword of the $\operatorname{read}_{nt}(T)$ of the form

$$k(k-1)\ldots(i+1)i$$

such that $i \in \overline{\text{read}}_{i-1}(T)$ and j is paired with j+1 for $i \leq j < k$.

EXAMPLE 3.17

Consider

5	5	15	15	5	5	5	5	15	5	5
	4	4	4	4	4	4	4	2	3	4
		3	3	3	2	1	3			
			2	2						
				1						

which has a (2,3)-consecution consisting of the underlined entries and a (1,4)-consecution consisting of the bold entries. Note that there is also a (2,2)-consecution contained in the (2,3)-consecution, and similarly a (1,k)-consecution for k=1,2,3 contained in the (1,4)-consecution (which are also nested).

PROPOSITION 3.18

There exists a bijection Ξ between subsets of Φ^+ and lowest-weight elements of $SDT(-\infty)$. Furthermore, Ξ is weight-preserving: for $B = \{\beta_1, \ldots, \beta_k\} \subseteq \Phi^+$, we have $wt(\Xi(B)) = \beta_1 + \cdots + \beta_k$.

Proof

For a fixed $B \subseteq \Phi^+$, write $B = X_2 \cup \cdots \cup X_n$, where $X_j = \{\epsilon_i - \epsilon_j : i < j\} \subseteq \Phi^+$. We construct the element $\Xi(B)$ using the following algorithm. We start with j = 2 and $T_1 = L^{-\infty}$ and proceed inductively on j. On step j, we perform the following procedure.

- (1) Order X_j such that $\epsilon_i \epsilon_j \prec \epsilon_{i'} \epsilon_j$ if and only if i > i', and enumerate the elements in X_j by β_1 to β_{ℓ_j} .
- (2) Start with k = 1, and define $T_{j,0}$ by adding a trivial column of height n + 1 j to T_{j-1} . Proceed inductively on k as follows.
 - (a) Suppose $\beta_k = \epsilon_i \epsilon_j$.

- (b) Push in a trivial column of height n+2-j. Change the entry of the (j+1)th entry of the jth row⁵ of $T_{i,k-1}$ to j-k.
 - (c) Define b = j k.
- (d) If b = i, then terminate this subprocess and denote the resulting tableau by $T_{j,k}$.
 - (e) We have the following:
- (i) Suppose there is a b+1 in $\overline{\text{read}}_{b+1}(T_{j,k-1})$ in cell c in $T_{j,k-1}$. If there is a (b+1)-box c^{\rightarrow} immediate to the right of c in $T_{j,k-1}$, then replace the b+1 in c^{\rightarrow} with b. Otherwise, replace b+1 in cell c with b. Set a=b.
- (ii) Otherwise, find the leftmost (a,b)-consecution of T such that a > i. Replace this subword with $(a-1)a \dots (b-1)$.
- (iii) If there is no such (a,b)-consecution, set a=b, push in a trivial n+1-a column, and replace the rightmost trivial a with an a-1.
 - (f) Redefine b = a 1 and repeat from step (2d).
 - (3) Set $T_j = T_{j,\ell_j}$ to be the result of the previous subprocess.

Define $\Xi(B) = T_n$.

Note that the increasing word⁶ in row j has size $|X_j|$; in particular, for $|X_j| > 0$, the nontrivial entries of the jth row of T_j are

$$j - |X_j|$$
 \cdots $j - 1$ j

by the result of applying (2b). Furthermore, there exists precisely one (i', j)-consecution in $T_{j,k}$ after adding the root $\beta_k = \epsilon_i - \epsilon_j$ with $i' \leq i$.

It is clear that the algorithm does not fail. So in order to show that Ξ is well-defined, we need to show that $\Xi(B)$ is a lowest-weight element. Therefore, we have that $\Xi(B)$ is a I_0 -lowest-weight element since the b-1 after applying step (2e) pairs with the a'-1 from the previous iteration with k-1 (recall from step (2f) that b=a'-1), and in an (a,b)-consecution, all letters except a and possibly b are paired. Note that if no such (a,b)-consecution exists, we can consider the rightmost trivial a=b as an (a,a)-consecution for the above proof. Furthermore, any 1 will always have 2 before it in the reading word, and so $f_{\overline{1}}T_{j,k}=0$. Hence, the map Ξ is well-defined. Moreover, Ξ is weight-preserving as $\operatorname{wt}(T_{j,k})-\operatorname{wt}(T_{j,k-1})=\beta_k$.

Now we show that Ξ is invertible, and so consider a lowest-weight element $T \in \mathrm{SDT}(-\infty)$. Let j be maximal such that the jth row of T has a nontrivial entry. If such j exists, then $T = L^{-\infty}$, and we terminate, as this is the unique element of weight 0 in $\mathrm{SDT}(-\infty)$. Otherwise, let i be minimal such that there is an (i,j)-consecution. Increase i to i' until the subword $i'w = i'(i'+1) \dots j$ of this (i,j)-consecution does not contain an (a,b)-consecution, for some $i < a \le b < j$. Define \mathfrak{c} to be the cell in T containing the i' if the cell immediately to its left is

 $^{{}^5}$ In our convention, the jth row of a tableau T is the jth row from the bottom.

⁶Following [14], the increasing word of a hook word $x_1 \ge \cdots \ge x_\ell < y_1 < \cdots < y_m$ is $y_1 < \cdots < y_m$.

not an i'; otherwise, c is this cell immediately to the left. Define T' by increasing the letter in c and every letter in w by 1 and removing the rightmost column $[n,\ldots,j+1,j+1]^{\top}$ and any trivial columns so that the result is marginally large. It is straightforward to see that this is the inverse of adding the root $\beta_{\ell_j} = \epsilon_{i'} - \epsilon_j$ to T' under step (2). By the definition of an (i',j)-consecution, T' is a lowest-weight element, and, furthermore, the maximality of j and the minimality of i and i' ensure that the process of doing the aforementioned procedure again yields $\beta_{\ell_j-1} \prec \beta_{\ell_j}$. Hence, the map Ξ is a bijection.

EXAMPLE 3.19

Consider the roots $B = \{\epsilon_2 - \epsilon_3, \epsilon_2 - \epsilon_4, \epsilon_1 - \epsilon_5\}$ for n = 5. At each step, we will shade the nontrivial entries and bold the changed letters. The first step is adding $\epsilon_2 - \epsilon_3$, which results in

$$T_2 = T_{2,1} = \begin{bmatrix} 5 & 5 & 5 & 5 & 5 & 5 & 5 \\ & 4 & 4 & 4 & 4 & 4 & 4 \\ & & 3 & 3 & 3 & 2 & 3 \\ & & & 2 & 2 & \\ & & & 1 & \\ \end{bmatrix}$$

Next we add in $\epsilon_2 - \epsilon_4$, and we note that there is no (3,3)-consecution, so we change the rightmost (trivial) 3 in the third row to a 2. Hence, we have

	5	5	5	5	5	5	5	5	5	5
	0	4	4	H	H	4	4	4	0	4
		4	4	4	4	4	4	4	3	4
$T_{3,1} =$			3	3	3	2	2	3		
- 3,1				2	2					
					1					

When we add in the $\epsilon_1 - \epsilon_4$, we note that there is a (2,2)-consecution, which is inside the (2,3)-consecution in the third row. Hence, we split this into a (1,1)-consecution and a (3,3)-consecution, resulting in

Now when we add in $\epsilon_1 - \epsilon_5$, there are no (a, 4)-consecutions, so we change the rightmost trivial 4 to a 3. Next, by Example 3.17, there is a (2, 3)-consecution in the third row, which then decrease to obtain

Now let us calculate one step of the inverse operation for T_4 . There exists a (1,5)-consecution that is marked in bold:

5	5	5	5	5	5	5	5	5	5	5	5	4	5
	4	4	4	4	4	4	4	3	2	3	4		
		3	3	3	1	1	2						
			2	2									
				1									

and note that it does not contain the (1,4)-consecution. However, the cell immediately to the left of the 1 in the (1,5)-consecution is also a 1, so the letters we increase are precisely those changed in equation (3.2).

EXAMPLE 3.20

Consider the roots $B = \{\epsilon_1 - \epsilon_3, \epsilon_2 - \epsilon_5, \epsilon_1 - \epsilon_5\}$. Adding the root $\epsilon_1 - \epsilon_3$ results in

$$T_{3} = T_{3,1} = \begin{bmatrix} 5 & 5 & 5 & 5 & 5 & 5 & 5 \\ 4 & 4 & 4 & 4 & 4 & 4 & 4 \\ & 3 & 3 & 3 & 3 & 2 & 3 \\ & & 2 & 2 & 1 & \\ & & & 1 & & \end{bmatrix}$$

and adding the root $\epsilon_2 - \epsilon_5$ results in

Finally, we add $\epsilon_1 - \epsilon_5$, and, in this case, when we add the 3 to the top row, the 3 in the subsequent row is in $\overline{\text{read}}_3(T_{5,1})$, and so we must change that to a 2. Denote the result $T'_{5,1}$, and then the leftmost 2 in the next row becomes unpaired in $\overline{\text{read}}_2(T'_{5,1})$, so we must additionally change that to a 1. However, the cell immediately to its right is also a 2; hence, we obtain

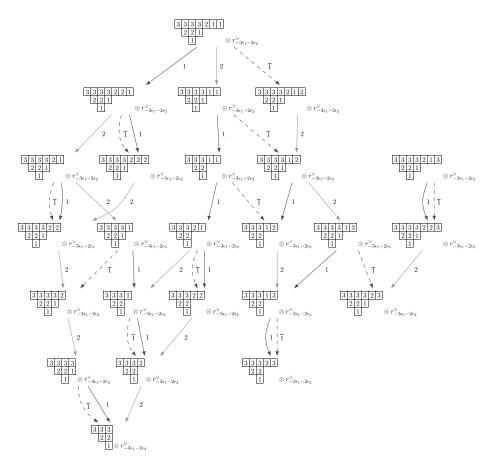
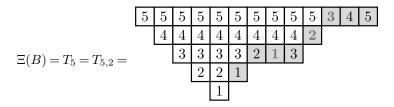


Figure 2. (Color online) The $\mathfrak{q}(3)$ -crystal SDT(λ) with $\lambda = -3\epsilon_1 - \epsilon_2$ created using SAGEMATH [32].



3.4. Recovering $SDT(\lambda)$ from $SDT(-\infty)$

Our construction is parallel to the $\mathfrak{gl}(n)$ -crystal construction of $B(\lambda)$ from $B(\infty)$ by essentially undoing the direct limit construction and adjusting $\varepsilon_i(b)$ to be the number of times that we can apply e_i before getting $\mathbf{0}$. See Figure 2 for an example. However, we note that whenever $\lambda_1 - \lambda_2 = 0$, we obtain a connected component that is too large, as we should have $e_{\overline{1}}(L^{-\infty} \otimes r_{\lambda}^{\vee}) = \mathbf{0}$. Thus, we would require a modification to the tensor product rule, but we can obtain $\mathrm{SDT}(\lambda)$ when $\lambda_i > \lambda_{i+1}$ for all $i \in I_0$.

THEOREM 3.21

Let $\lambda \in \Lambda^-$ be such that $\lambda_i > \lambda_{i+1}$ for all $i \in I_0$. Let

$$\mu = \sum_{i=1}^{n} (\lambda_i - k)\epsilon_{n+1-i} \quad for \ some \ k \ge -\min\{\lambda_i : i \in I_0\}.$$

The connected component C of $\mathrm{SDT}(-\infty) \otimes \mathcal{R}_{\mu}^{\vee}$ generated by $L^{-\infty} \otimes r_{\mu}^{\vee}$, using the tensor product rule in equation (2.1), is isomorphic to $\mathrm{SDT}(\lambda)$ as $\mathfrak{q}(n)$ -crystals.

Proof

We define a map $\psi \colon \mathcal{C} \longrightarrow \mathrm{SDT}(\lambda)$ by $\psi(T \otimes r_{\mu}^{\vee}) = E(T)$, where E(T) is the semistandard decomposition tableau formed by pushing in/out trivial columns. The tensor product rule in equation (2.1) says that $e_{\overline{1}}(T \otimes r_{\mu}^{\vee}) = (e_{\overline{1}}T) \otimes r_{\mu}^{\vee}$ and similarly for $f_{\overline{1}}$. Moreover, we have in fact shown in Lemma 3.12 that E commutes with all $e_{\overline{1}}$ and $f_{\overline{1}}$ crystal operators, and hence $e_{\overline{1}}$ and $f_{\overline{1}}$ commute with ψ . Next, we note that from the $\mathfrak{gl}(n)$ -case, we have that the I_0 -component generated by $L^{-\infty} \otimes r_{\mu}$ is isomorphic to the I_0 -component of L^{λ} and ψ is a crystal isomorphism when restricted to these I_0 -components. Thus, since both \mathcal{C} and $\mathrm{SDT}(\lambda)$ are $\mathfrak{q}(n)$ -crystals and the computation of $\overline{\mathrm{read}}_i(T)$, we have that ψ commutes with e_i and f_i for $i \in I_0$. Hence, ψ is the desired crystal isomorphism.

Let us remark on the μ in Theorem 3.21. Recall that tensoring $B(\lambda)$ with the determinant representation of $\mathfrak{q}(n)/\mathfrak{gl}(n)$ (i.e., the one-dimensional representation of weight $(1,\ldots,1)$) does not change the crystal structure of $B(\lambda)$ up to a shift in the weight by $(1,\ldots,1)$. Therefore, to remove this dependence of the determinant representation, we could modify the condition $\mathrm{wt}(\psi(b)) = \mathrm{wt}(b)$ of a crystal morphism to instead be $\mathrm{wt}_i(\psi(b)) = \mathrm{wt}_i(b)$ for all $i \in I_0$. All of the theory given above holds under this modification, as our tensor product rule does not change under this shift in weight (unlike that given in [14]). This is analogous to considering the $\mathfrak{sl}(n)$ -crystals instead of $\mathfrak{gl}(n)$ -crystals.

The condition $\lambda_i > \lambda_{i+1}$ ensures that the weight λ is not on a wall of the dominant chamber. Furthermore, when $\lambda_i = \lambda_{i+1} > 0$, the corresponding irreducible representation is not finite-dimensional (see [9, Theorem 2.18]). Since the crystal structure does not change when shifting the weight by (1, 1, ..., 1) but the representation-theoretic information does (the dimensionality changes), we require all parts to be not equal to preserve finite-dimensionality for the corresponding irreducible representation.

More concretely, if $\nu = \lambda + k(1^n)$ for some $k \in \mathbf{Z}$, then $\mathcal{R}_{\lambda}^{\vee} \cong \mathcal{R}_{\nu}^{\vee}$ and $B(\lambda) \cong B(\nu)$ as $\mathfrak{q}(n)$ -crystals under this broader definition of a crystal isomorphism. Therefore, our choice of μ in Theorem 3.21 is so that we have $\mu \in \mathbf{Z}_{\geq 0}^n$, but it does not change the resulting crystal graphs, ε_i , and φ_i .

EXAMPLE 3.22

We consider $\mathfrak{q}(3)$ and $\lambda = -3\epsilon_1 - \epsilon_2$. Then we can take $\mu = 3\epsilon_3 + \epsilon_2 - k(\epsilon_1 + \epsilon_2 + \epsilon_3)$ for some $k \geq 3$. Thus, when we take the connected component generated by

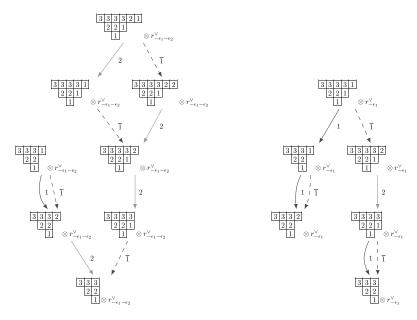


Figure 3. (Color online) The $\mathfrak{q}(3)$ -crystal connected component of $\mathrm{SDT}(-\infty) \otimes \mathcal{R}_{-\epsilon_1}^{\vee}$ (resp., $\mathrm{SDT}(-\infty) \otimes \mathcal{R}_{-\epsilon_1-\epsilon_2}^{\vee}$) generated by $L^{-\infty} \otimes r_{-\epsilon_1}^{\vee}$ (resp., $L^{-\infty} \otimes r_{-\epsilon_1-\epsilon_2}^{\vee}$).

 $L^{-\infty} \otimes r_{\mu}^{\vee}$, we obtain the crystal graph in Figure 2 (taken explicitly with k = 3, so $\mu = 0\epsilon_3 - 2\epsilon_2 - 3\epsilon_1$). Thus for k = 3, we obtain a crystal isomorphic to SDT(λ) given by [14, Figure 1].

Let us discuss how to extend Theorem 3.21 to more general cases. Consider the examples in Figure 3. For $\lambda = -\epsilon_1 - \epsilon_2$, we note that the connected component we obtain after also setting $e_{\overline{1}}(L^{-\infty} \otimes r_{\lambda}^{\vee}) = \mathbf{0}$ is isomorphic to $\mathrm{SDT}(\lambda)$. Therefore, a suitably modified tensor product rule should yield $\mathrm{SDT}(\lambda)$ when λ may also contain zero entries. Furthermore, we would expect a modified tensor product rule to yield dual representations. For instance, if we consider $\lambda = -\epsilon_1$, note that after setting $e_{\overline{1}}(e_2e_1L^{-\infty}\otimes r_{\lambda}^{\vee}) = \mathbf{0}$, we would obtain the dual version of $\mathrm{SDT}(-\epsilon_1 - \epsilon_2)$.

Appendix. SageMath code

We construct Figure 2:

```
sage: B = crystals.infinity.Tableaux(['Q',3])
sage: e = B.weight_lattice_realization().basis()
sage: R = crystals.elementary.R(['Q',3], -3*e[0]-2*e[1], dual=True)
sage: x = tensor([B.module_generators[0], R.module_generators[0]])
sage: view(x.subcrystal(max_depth=10,index_set=[-1,1,2]))
```

In Figure 2, we have done some mild postprocessing on the resulting crystal graph to shift the basis of Λ^- from being (e_0, e_1, e_2) to $(\epsilon_0, \epsilon_1, \epsilon_2)$. The construction of Figure 3 is similar.

Acknowledgments. T.S. would like to thank Central Michigan University for its hospitality during his visit in November 2018 and January 2020. The authors thank the referees for helpful comments. This work benefited from computations using SageMath [32].

B.S. was partially supported by Simons Foundation grant 429950. T.S. was partially supported by Australian Research Council grant DP170102648.

References

- S. Assaf and E. Kantarci Oğuz, Toward a local characterization of crystals for the quantum queer superalgebra, Ann. Comb. 24 (2020), no. 1, 3–46.
 MR 4078137. DOI 10.1007/s00026-019-00477-0.
- [2] G. Benkart, S.-J. Kang, and M. Kashiwara, Crystal bases for the quantum superalgebra $U_q(\mathfrak{gl}(m,n))$, J. Amer. Math. Soc. **13** (2000), no. 2, 295–331. MR 1694051. DOI 10.1090/S0894-0347-00-00321-0.
- [3] S. Billey and M. Haiman, Schubert polynomials for the classical groups, J. Amer.
 Math. Soc. 8 (1995), no. 2, 443–482. MR 1290232. DOI 10.2307/2152823.
- [4] J. Brundan, Kazhdan–Lusztig polynomials and character formulae for the Lie superalgebra $\mathfrak{gl}(m|n)$, J. Amer. Math. Soc. **16** (2003), no. 1, 185–231. MR 1937204. DOI 10.1090/S0894-0347-02-00408-3.
- [5] _____, Kazhdan-Lusztig polynomials and character formulae for the Lie superalgebra q(n), Adv. Math. 182 (2004), no, 1, 28–77. MR 2028496.
 DOI 10.1016/S0001-8708(03)00073-2.
- [6] S.-J. Cheng, Supercharacters of queer Lie superalgebras. J. Math. Phys. 58
 (2017), no. 6, art. 061701, 9 pp. MR 3657145. DOI 10.1063/1.4984594.
- S.-J. Cheng and J.-H. Kwon, Finite-dimensional half-integer weight modules over queer Lie superalgebras, Comm. Math. Phys. 346 (2016), no. 3, 945–965.
 MR 3537341. DOI 10.1007/s00220-015-2544-0.
- [8] S.-J. Cheng, J.-H. Kwon, and W. Wang, Character formulae for queer Lie superalgebras and canonical bases of types A/C, Comm. Math. Phys. 352 (2017), no. 3, 1091–1119. MR 3631400. DOI 10.1007/s00220-016-2809-2.
- [9] S.-J. Cheng and W. Wang, Dualities and Representations of Lie Superalgebras, Grad. Stud. in Math. 144, Amer. Math., Soc., Providence, 2012. MR 3012224. DOI 10.1090/gsm/144.
- [10] M. Chmutov, C. Hoyt, and S. Reif, The Kac-Wakimoto character formula for the general linear Lie superalgebra, Algebra Number Theory 9 (2015), no. 6, 1419–1452. MR 3397407. DOI 10.2140/ant.2015.9.1419.
- [11] S. Clark, Canonical bases for the quantum enveloping algebra of $\mathfrak{gl}(m|1)$ and its modules, preprint, 1605.04266.
- [12] G. Cliff, Crystal bases and Young tableaux, J. Algebra 202 (1998), no. 1, 10–35.
 MR 1614241. DOI 10.1006/jabr.1997.7244.

- [13] M. Gillespie, G. Hawkes, W. Poh, and A. Schilling, Characterization of queer supercrystals, J. Combin. Theory Ser. A 173 (2020), art. 105235, 53 pp. MR 4065438. DOI 10.1016/j.jcta.2020.105235.
- [14] D. Grantcharov, J. H. Jung, S.-J. Kang, M. Kashiwara, and M. Kim, Crystal bases for the quantum queer superalgebra and semistandard decomposition tableaux, Trans. Amer. Math. Soc. 366 (2014), no. 1, 457–489. MR 3118402. DOI 10.1090/S0002-9947-2013-05866-7.
- [15] ______, Crystal bases for the quantum queer superalgebra, J. Eur. Math. Soc. (JEMS) 17 (2015), no. 7, 1593–1627. MR 3361724. DOI 10.4171/JEMS/540.
- [16] D. Grantcharov, J. H. Jung, S.-J. Kang, and M. Kim, Highest weight modules over quantum queer superalgebra $U_q(\mathfrak{q}(n))$, Comm. Math. Phys. **296** (2010), no. 3, 827–860. MR 2628823. DOI 10.1007/s00220-009-0962-6.
- [17] I. Grojnowski and G. Lusztig, "A comparison of bases of quantized enveloping algebras" in *Linear Algebraic Groups and Their Representations (Los Angeles, CA, 1992)*, Contemp. Math. 153, Amer. Math. Soc., Providence, 11–19. MR 1247495. DOI 10.1090/conm/153/01304.
- [18] Z. Hamaker, E. Marberg, and B. Pawlowski, Schur P-positivity and involution Stanley symmetric functions, Int. Math. Res. Not. IMRN 2019, no. 17, 5389–5440. MR 4023763. DOI 10.1093/imrn/rnx274.
- [19] G. Hawkes, K. Paramonov, and A. Schilling, Crystal analysis of type C Stanley symmetric functions, Electron. J. Combin. 24 (2017), no. 3, art. 3.51, 32 pp. MR 3711093. DOI 10.37236/6952.
- [20] J. Hong and H. Lee, Young tableaux and crystal B(∞) for finite simple Lie algebras, J. Algebra 320 (2008), no. 10, 3680–3693. MR 2457716.
 DOI 10.1016/j.jalgebra.2008.06.008.
- [21] T. Józefiak, Schur Q-functions and cohomology of isotropic Grassmannians, Math. Proc. Cambridge Philos. Soc. 109 (1991), no. 3, 471–478. MR 1094746. DOI 10.1017/S0305004100069917.
- [22] M. Kashiwara, Crystalizing the q-analogue of universal enveloping algebras, Comm. Math. Phys. 133 (1990), no. 2, 249–260. MR 1090425.
- [23] ______, On crystal bases of the q-analogue of universal enveloping algebras, Duke Math. J. 63 (1991), no. 2, 465–516. MR 1115118. DOI 10.1215/S0012-7094-91-06321-0.
- [24] ______, The crystal base and Littlemann's refined Demazure character formula,
 Duke Math. J. 71 (1993), no. 3, 839–858. MR 1240605.
 DOI 10.1215/S0012-7094-93-07131-1.
- [25] M. Kashiwara, Bases cristallines des groupes quantiques, Cours Spéc. 9, Soc. Math. France, Paris, 2002. MR 1997677.
- [26] J.-H. Kwon, Crystal graphs for general linear Lie superalgebras and quasi-symmetric functions, J. Combin. Theory Ser. A, 116 (2009), no. 7, 1199–1218. MR 2527604. DOI 10.1016/j.jcta.2009.03.007.

- [27] _____, Crystal bases of q-deformed Kac modules over the quantum superalgebras $U_q(\mathfrak{gl}(m|n))$, Int. Math. Res. Not. IMRN **2014**, no. 2, 512–550. MR 3159079. DOI 10.1093/imrn/rns221.
- [28] P. Littelmann, A Littlewood–Richardson rule for symmetrizable Kac–Moody algebras, Invent. Math. 116 (1994), nos. 1–3, 329–346. MR 1253196. DOI 10.1007/BF01231564.
- [29] _____, Paths and root operators in representation theory, Ann. of Math. (2) **142** (1995), no. 3, 499–525. MR 1356780. DOI 10.2307/2118553.
- [30] G. Lusztig, Canonical bases arising from quantized enveloping algebras, J. Amer. Math. Soc. 3 (1990), no. 2, 447–498. MR 1035415. DOI 10.2307/1990961.
- [31] I. Penkov and V. Serganova, Characters of irreducible G-modules and cohomology of G/P for the Lie supergroup G=Q(N), J. Math. Sci. (New York) 84 (1997), no. 5, 1382–1412. Algebraic geometry, 7. MR 1465520. DOI 10.1007/BF02399196.
- [32] Sage Developers, SageMath Software, version 9.0, Sage Development Team, 2020, http://www.sagemath.org.
- [33] B. Salisbury and T. Scrimshaw, A rigged configuration model for $B(\infty)$, J. Combin. Theory Ser. A **133** (2015), 29–57. MR 3325627. DOI 10.1016/j.jcta.2015.01.008.
- [34] _____, Rigged configurations for all symmetrizable types, Electron. J. Combin. **24** (2017), no. 1, art. 30, 13 pp. MR 3625907.
- [35] ______, Rigged configurations and the *-involution, Lett. Math. Phys. 108 (2018), no. 9, 1985–2007. MR 3832140. DOI 10.1007/s11005-018-1063-2.
- [36] J. Schur. Über die Darstellung der symmetrischen und der alternierenden Gruppe durch gebrochene lineare Substitutionen, J. Reine Angew. Math. 139 (1911), 155–250. MR 1580818. DOI 10.1515/crll.1911.139.155.
- [37] V. Serganova, "Kazhdan–Lusztig polynomials for Lie superalgebra $\mathfrak{gl}(m|n)$ " in *I. M. Gel'fand Seminar*, Adv. Soviet Math. **16**, Amer. Math. Soc., Providence, 1993, 151–165. MR 1237837.
- [38] J. R. Stembridge, A local characterization of simply-laced crystals, Trans.
 Amer. Math. Soc. 355 (2003) no. 12, 4807–4823. MR 1997585.
 DOI 10.1090/S0002-9947-03-03042-3.
- [39] Y. Su and R. B. Zhang, Character and dimension formulae for general linear superalgebra, Adv. Math. 211 (2007), no. 1, 1–33. MR 2313526. DOI 10.1016/j.aim.2006.07.010.
- [40] _____, Character and dimension formulae for queer Lie superalgebra, Comm. Math. Phys. 333 (2015), no. 3, 1465–1481. MR 3302639. DOI 10.1007/s00220-014-2209-4.
- [41] J. Van der Jeugt, J. W. B. Hughes, R. C. King, and J. Thierry-Mieg, *Character formulas for irreducible modules of the Lie superalgebras* $\mathfrak{sl}(m/n)$, J. Math. Phys. **31** (1990), no. 9, 2278–2304. MR 1067852. DOI 10.1063/1.528637.

 $Salisbury: \ Department \ of \ Mathematics, \ Central \ Michigan \ University, \ Mount \ Pleasant, \\ Michigan, \ USA; \ salis1bt@cmich.edu; \ https://bsalisbury1.github.io$

Scrimshaw: School of Mathematics and Physics, The University of Queensland, St. Lucia, Queensland, Australia; tcscrims@gmail.com; https://tscrim.github.io