

USV Modeling

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Abstract

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I. INTRODUCTION

II. BACKGROUND

A. [Sonnenburg *et al.*, 2010] and [Sonnenburg and Woolsey, 2013]

[Sonnenburg and Woolsey, 2013] and [Sonnenburg *et al.*, 2010] examine model for USV with steerable outboard motor (vectored thrust) where sideslip is a major concern. Uses notation and maneuvering model from [Fossen, 1994].

- Full vessel model is includes linear and quadratic damping
- All models are then linearized (perturbation dynamics) for the purpose of identification where the coefficients are parameterized by the states (surge, sway and yaw-rate)
- Actuation model thrust as a linear and quadratic with velocity dependence.

Model identification appears to throw out the physical model and rely on identification of linear models (perturbation models) and a set of discrete speeds.

- Model identification includes open-loop maneuvers to identify steady-state parameters and closed-loop maneuvers to identify dynamic parameters.
 - 1) Steady-state values are identified by as linear relationships between yaw-rate, side-slip angle and side-slip speed. The coefficients of these a relationships are determined for a set of discrete, constant speeds. The values of these coefficients change significantly over the range of forward speeds and now functional relationship is offered.
 - 2) Thruster model conflates both the thrust relationship and the inertia of the system. Model is identified by measuring initial acceleration during step changes in engine RPM and identifying

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the two linear coefficients of a linear model. This single model is constant over the range of speeds. Sparse data at higher speeds.

- 3) Speed/surge is modeled as first-order model parameterized over speed range.
- 4) Steering model (first order Nomoto model with sideslip) is indentified by minimizing a quadratic cost function of sideslip angle and yaw-rate over close-loop time histories. Again, the linearized models have coefficients that change with speed. The results again show significant changes over the speed envelope without offering a functional relationship.

B. [Caccia et al., 2008]

Uses the nonlinear model of Blanke [Blanke, 1981] as reported in [Fossen, 1994]. Speed/surge model:

- Neglects added mass term based on Blanke comments. Blanke suggests that the surge added mass term “will typically be less then 5%” [Fossen, 1994].
- Neglects r^2 terms, as suggested by Blanke, based on assuming reasonably low yaw-rate ($\dot{\psi}$ 10 degrees/s).
- Assumes linear+quadratic drag. The Blanke model reported in [Fossen, 1994] includes only the quadratic term.
- Neglects the Coriolis terms based on assuing negligible sway velocity. The experimental evidence from [Sonnenburg and Woolsey, 2013] and [Sonnenburg et al., 2010] indicate that there is a significant sideslip angle, hence a significant sway velocity.

Steering model

- Neglects added mass in surge and yaw.

Sway (sideslip) was not observable in experiments! Used only GPS and heading for identification making making many of the estimated quantities unreliable or unobservable.

Thrust model neglects speed of advance, assumes thrust is indepeneden of vessel speed.

III. MANEUVERING MODEL

In this section we follow the notation and process detailed in [Fossen, 2011], Chapter 7. The horiozontal-plane maneuvering model captures is formulated using state vector $\boldsymbol{\nu} = [u, v, r]^T$ where the velocities u , v and r are in the surge, sway and yaw directions respectively. The velocities are considered to be relative to an irrotational constant ocean current. The nonlinear maneuvering equations from [Fossen, 2011] are

$$\underbrace{M_{RB}\dot{\boldsymbol{\nu}} + C_{RB}(\boldsymbol{\nu})\boldsymbol{\nu}}_{\text{rigid-body forces}} + \underbrace{M_A\dot{\boldsymbol{\nu}}_r + C_A(\boldsymbol{\nu}_r)\boldsymbol{\nu}_r + D(\boldsymbol{\nu}_r)\boldsymbol{\nu}_r}_{\text{hydrodynamic forces}} = \boldsymbol{\tau} + \boldsymbol{\tau}_{wind} + \boldsymbol{\tau}_{waves} \quad (1)$$

where $\boldsymbol{\nu}_r$ is the velocity vector relative to an irrotational water current $\boldsymbol{\nu}_c$, i.e., $\boldsymbol{\nu} = \boldsymbol{\nu}_r + \boldsymbol{\nu}_c$. The rigid body kinetics are represented by the rigid body mass \mathbf{M}_{RB}

$$\mathbf{M}_{RB} = \begin{bmatrix} m & 0 & 0 \\ 0 & m & mx_g \\ 0 & mx_g & I_z \end{bmatrix}, \quad (2)$$

where m is the mass of the vehicle, I_z is the moment of inertia about the body-centered z-axis and x_g is distance, along the x-axis, from the origin of the body-centered frame to the center of gravity of the vessel, and by the rigid body Coriolis-centripetal matrix,

$$\mathbf{C}_{RB}(\boldsymbol{\nu}) = \begin{bmatrix} 0 & 0 & -m(x_g r + v) \\ 0 & 0 & mu \\ m(x_g r + v) & -mu & 0 \end{bmatrix}. \quad (3)$$

Noting that $\mathbf{C}_{RB}(\boldsymbol{\nu})$ is skew-symmetric, i.e., $\mathbf{C}_{RB}(\boldsymbol{\nu}) = -\mathbf{C}_{RB}^T(\boldsymbol{\nu})$. The hydrodynamic effects are represented by the added mass matrix

$$\mathbf{M}_A = \begin{bmatrix} -X_{\dot{u}} & 0 & 0 \\ 0 & -Y_{\dot{v}} & -Y_{\dot{r}} \\ 0 & -Y_{\dot{r}} & -N_{\dot{r}} \end{bmatrix}. \quad (4)$$

and the Coriolis-centripetal matrix for the added mass

$$\mathbf{C}_A(\boldsymbol{\nu}_r) = \begin{bmatrix} 0 & 0 & Y_{\dot{v}}v_r + Y_{\dot{r}}r \\ 0 & 0 & -X_{\dot{u}}u_r \\ -Y_{\dot{v}}v_r - Y_{\dot{r}}r & X_{\dot{u}}u_r & 0 \end{bmatrix}. \quad (5)$$

It is worth noting that \mathbf{C}_A includes the nonlinear Munk moment (see [Fossen, 2011] p.121). Following [Fossen, 2011] the SNAME notation for the hydrodynamic derivatives.

Use [Sonnenburg et al., 2010] to get linear + quadratic

$$\mathbf{D}(\boldsymbol{\nu}_r) = \begin{bmatrix} X_u + X_{u|u}|u| & 0 & 0 \\ 0 & Y_v + Y_{v|v}|v| & Y_r + Y_{r|r}|r| \\ 0 & N_v + N_{v|v}|v| & N_r + N_{r|r}|r| \end{bmatrix}. \quad (6)$$

A. Thrust Model

Two options:

Assume thrust is independent of speed (as done in the Caccia papers).

Or assume and unknown, linear decrease in thrust with speed.

$$T = T_o(1 - au)$$

where a is the linear speed reduction

B. Speed model

Consider the surge state of the model above where

$$\underbrace{m\dot{u}}_{\text{RB inertia}} - \underbrace{mx_g r^2}_{\text{RB centripetal}} - \underbrace{mvu}_{\text{RB Coriolis}} = \underbrace{X_{\dot{u}}\dot{u}}_{\text{AM inertia}} + \underbrace{Y_{\dot{v}}v_r r}_{\text{AM Coriolis}} + \underbrace{Y_{\dot{r}}r^2}_{\text{AM centripetal}} + \underbrace{X_u u + X_{u|u}|u|u}_{\text{Drag}} + \underbrace{T}_{\text{Thrust}} \quad (7)$$

Following [Caccia et al., 2008], based on [Fossen, 1994], we neglect the second-order centripetal terms

$$m\dot{u} - mvu = X_{\dot{u}}\dot{u} + Y_{\dot{v}}v_r r + X_u u + X_{u|u}|u|u + T \quad (8)$$

For steady state forward motion ($\dot{u} = v = r = 0$) in stationary water ($v_r = 0$)

$$0 = +X_u u + X_{u|u}|u|u + T \quad (9)$$

We can estimate X_u and $X_{u|u}$ from steady state forward motion trials with known thrust input by testing at a series of known forward speeds and measuring u .

Considering forward-only acceleration

$$m\dot{u} = X_{\dot{u}}\dot{u} + X_u u + X_{u|u}|u|u + T \quad (10)$$

we can identify the added mass ($X_{\dot{u}}$) by either estimating the initial acceleration (see [Sonnenburg et al., 2010]) or by examining the 'time constant' for such tests.

This leaves the coefficient $Y_{\dot{v}}$, related to the added-mass Coriolis force, as the single unknown.

C. Steering Model

IV. MODEL IDENTIFICATION TESTS

A. Physical Measurements

- Measure the mass (m) directly.
- Measure the moment of inertia (I_z) using a bifilar pendulum.

B. Thrust Characterization

Bollard tests in the tank to measure thrust force (at zero velocity) as a function of motor command.

C. Steady-State Tests

- Surge: Measure the steady-state speed at a variety of thrust inputs to identify the drag terms.
- Yaw: Measure the steady-state yaw rate at variety of torque inputs to identify the yaw drag terms.

D. Open-Loop Dynamic Tests

- Surge: Measure step response to forward thrust (with heading control?) to estimate added surge mass.
- Yaw: Measure step response to torque to estimate added mass/inertia in yaw.

E. Closed-Loop Dynamic Tests

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