

Problem set 1

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1 How far apart are two quantum states?

Consider two quantum states described by density operators ρ and $\tilde{\rho}$ in an N -dimensional Hilbert space, and consider the complete orthogonal measurement $\{E_a, a = 1, 2, 3, \dots, N\}$, where the E_a 's are one-dimensional projectors satisfying

$$\sum_{a=1}^N E_a = I \quad (1)$$

When the measurement is performed, outcome a occurs with probability $p_a = \text{tr}(\rho E_a)$ if the state is ρ and with probability $\tilde{p}_a = \text{tr}(\tilde{\rho} E_a)$ if the state is $\tilde{\rho}$. The L^1 distance between two probability distributions is defined as

$$d(p, \tilde{p}) \equiv ||p - \tilde{p}||_1 \equiv \frac{1}{2} \sum_{a=1}^N |p_a - \tilde{p}_a|; \quad (2)$$

this distance is zero if the two distributions are identical, and attains its maximum value one if the two distributions have support on disjoint sets.

(a) Show that

$$d(p, \tilde{p}) \leq \frac{1}{2} \sum_{i=1}^N |\lambda_i| \quad (3)$$

where the λ_i 's are the eigenvalues of the Hermitian operator $\rho - \tilde{\rho}$

Solution:

Proof. Define $\Delta = \rho - \tilde{\rho}$. Then, $p_a - \tilde{p}_a = \text{tr}((\rho - \tilde{\rho})(E_a)) = \text{tr}(E_a \Delta)$ by linearity of trace.

$$\implies d(p, \tilde{p}) = \frac{1}{2} \sum_{a=1}^N |\text{tr}(E_a \Delta)| \quad (1)$$

Since Δ is Hermitian, we can diagonalize $\Delta = \sum_{i=1}^N \lambda_i |i\rangle \langle i|$.

$$= \frac{1}{2} \sum_{a=1}^N |\text{tr}(E_a \sum_{i=1}^N \lambda_i |i\rangle \langle i|)| = \frac{1}{2} \sum_{a=1}^N \left| \sum_{i=1}^N \lambda_i \text{tr}(E_a |i\rangle \langle i|) \right| \quad (2)$$

$$= \sum_{a=1}^N \left| \sum_{i=1}^N \lambda_i \langle i | E_a | i \rangle \right| \leq \frac{1}{2} \sum_{i=1}^N \sum_{a=1}^N |\lambda_i \langle i | E_a | i \rangle| \quad (3)$$

Since E_a is positive, $\langle i | E_a | i \rangle > 0$, therefore

$$= \frac{1}{2} \sum_{i=1}^N \sum_{a=1}^N |\lambda_i| \langle i | E_a | i \rangle = \frac{1}{2} \sum_{i=1}^N |\lambda_i| \langle i | \sum_{a=1}^N E_a | i \rangle \quad (4)$$

Using the completeness property, $\sum_{a=1}^N E_a = I$, therefore

$$= \frac{1}{2} \sum_{i=1}^N |\lambda_i| \langle i | I | i \rangle = \frac{1}{2} \sum_{i=1}^N |\lambda_i| \quad (5)$$

as desired. \square

- (b) Find a choice for the orthogonal projector $\{E_a\}$ that saturates the upper bound eq. (3).

Solution: For scalars $\lambda_i, |\sum_i \lambda_i| = \sum_i |\lambda_i| \iff$ all λ_i are the same sign. Then, construct operators E_+ and E_- that project onto the eigenspaces corresponding to positive/negative eigenvalues of $\rho - \tilde{\rho}$. We see that $E_+ + E_- = I$, since $\rho - \tilde{\rho}$ is positive and thus has no zero eigenvalues. Then,

$$\begin{aligned} d(p, \tilde{p}) &= \frac{1}{2} \sum_a^N \left| \sum_i^N \lambda_i \langle i | E_a | i \rangle \right| \\ &= \frac{1}{2} \left(\left| \sum_i^N \lambda_i \langle i | E_+ | i \rangle \right| + \left| \sum_i^N \lambda_i \langle i | E_- | i \rangle \right| \right) \\ &= \frac{1}{2} \left| \sum_k \lambda_k \right| + \left| \sum_l \lambda_l \right| \end{aligned}$$

where $\lambda_k > 0, \lambda_l < 0$.

$$= \frac{1}{2} \sum_i^N |\lambda_i|$$

- (c) Define distance $d(\rho, \tilde{\rho})$ between density operators as the maximal L^1 distance between corresponding probability distributions that can be achieved by any orthogonal measurement. From (a), (b), we have found that

$$d(\rho, \tilde{\rho}) = \frac{1}{2} \sum_{i=1}^N |\lambda_i|.$$

The L^1 norm $\|A\|_1$ of an operator A is defined as:

$$\|A\|_1 \equiv \text{tr}[(AA^\dagger)^{\frac{1}{2}}]$$

How can the distance $d(\rho, \tilde{\rho})$ be expressed as the L^1 norm of an operator?

Solution: AA^\dagger is PSD ($x^\dagger AA^\dagger x = (A^\dagger x)^\dagger (A^\dagger x)$, which is inner product, nonnegative) and thus $\sqrt{AA^\dagger}$ is PSD. Therefore, by the spectral theorem $\sqrt{AA^\dagger} = U = \sum_i^N \lambda_i |i\rangle \langle i|$ for $\lambda_i > 0$.

$$\begin{aligned} \text{tr}(U) &= \text{tr}\left(\sum_i^N \lambda_i |i\rangle \langle i|\right) = \sum_i^N \lambda_i \text{tr}(|i\rangle \langle i|) \\ &= \sum_i^N \lambda_i = \|A\|_1 \iff d(\rho, \tilde{\rho}) = \frac{1}{2} \|\rho - \tilde{\rho}\|_1 \end{aligned}$$

- (d) Suppose the states $\rho, \tilde{\rho}$ are pure states $\rho = |\psi\rangle \langle \psi|$ and $\tilde{\rho} = |\tilde{\psi}\rangle \langle \tilde{\psi}|$. If we adopt a suitable basis in the space spanned by the two vectors, and appropriate phase conventions, then these vectors can be expressed as:

$$|\psi\rangle = \begin{pmatrix} \cos \frac{\theta}{2} \\ \sin \frac{\theta}{2} \end{pmatrix}, |\tilde{\psi}\rangle = \begin{pmatrix} \sin \frac{\theta}{2} \\ \cos \frac{\theta}{2} \end{pmatrix}$$

where $\langle \psi | \tilde{\psi} \rangle = \sin \theta$

Express the distance $d(\rho, \tilde{\rho})$ in terms of the angle θ

Solution:

$$\rho - \tilde{\rho} = |\psi\rangle \langle \psi| - |\tilde{\psi}\rangle \langle \tilde{\psi}| \quad (6)$$

$$= \begin{pmatrix} \cos^2(\frac{\theta}{2}) - \sin^2(\frac{\theta}{2}) & 0 \\ 0 & \sin^2(\frac{\theta}{2}) - \cos^2(\frac{\theta}{2}) \end{pmatrix} \quad (7)$$

$$\iff \lambda = \cos^2(\frac{\theta}{2}) - \sin^2(\frac{\theta}{2}), \sin^2(\frac{\theta}{2}) - \cos^2(\frac{\theta}{2})$$

$$\implies d(\rho, \tilde{\rho}) = \frac{1}{2} \sum_i^N \lambda_i$$

$$\equiv d(\theta) = \frac{1}{2} (|\cos^2(\frac{\theta}{2}) - \sin^2(\frac{\theta}{2})| + |\sin^2(\frac{\theta}{2}) - \cos^2(\frac{\theta}{2})|)$$

Since $\cos(2\theta) = \cos^2(\theta) - \sin^2(\theta)$ (double angle identity):

$$= \frac{1}{2} |\cos \theta| + |-\cos \theta|$$

$$= |\cos \theta|$$

- (e) Express $\| |\psi\rangle - |\tilde{\psi}\rangle \|^2$ in terms of θ and by comparing the result of (d), derive the bound

$$d(|\psi\rangle \langle \psi|, |\tilde{\psi}\rangle \langle \tilde{\psi}|) \leq \| |\psi\rangle - |\tilde{\psi}\rangle \|.$$

Solution:

$$\begin{aligned} \| |\psi\rangle - |\tilde{\psi}\rangle \|^2 &= 2\cos^2(\frac{\theta}{2}) - 4\sin(\frac{\theta}{2})\cos(\frac{\theta}{2}) + 2\sin^2(\frac{\theta}{2}) \\ &= 2(\cos(\frac{\theta}{2}) - \sin(\frac{\theta}{2}))^2 \\ \implies \| |\psi\rangle - |\tilde{\psi}\rangle \| &= \sqrt{2}|\cos(\frac{\theta}{2}) - \sin(\frac{\theta}{2})| \end{aligned}$$

Then, we can see that $d(|\psi\rangle \langle \psi|, |\tilde{\psi}\rangle \langle \tilde{\psi}|)$

$$\begin{aligned} &= |\cos x| = |\cos^2(\frac{\theta}{2}) - \sin^2(\frac{\theta}{2})| \quad (\text{double angle}) \\ &= |(\cos(\frac{\theta}{2}) + \sin(\frac{\theta}{2}))(\cos(\frac{\theta}{2}) - \sin(\frac{\theta}{2}))| \\ &= |(\cos(\frac{\theta}{2}) + \sin(\frac{\theta}{2}))| |(\cos(\frac{\theta}{2}) - \sin(\frac{\theta}{2}))| \\ &\leq (\sqrt{2}\sqrt{\cos^2(\frac{\theta}{2}) + \sin^2(\frac{\theta}{2})}) |\cos(\frac{\theta}{2}) - \sin(\frac{\theta}{2})| \quad (\text{Cauchy-Schwarz}) \\ &= \sqrt{2}|\cos(\frac{\theta}{2}) - \sin(\frac{\theta}{2})| \end{aligned}$$

as desired.

- (f) Bob thinks that the norm $\| |\psi\rangle - |\tilde{\psi}\rangle \|$ should be a good measure of the distinguishability of the pure quantum states ρ and $\tilde{\rho}$. Explain why Bob is wrong.

Solution: Two indistinguishable states, for instance considering the case where $|\psi\rangle = |\tilde{\psi}\rangle$, should have a normed difference of zero, since. Since a quantum state is a ray, $|\psi\rangle \equiv e^{i\theta}|\psi\rangle$ and therefore should be indistinguishable. However $\| |\psi\rangle - e^{i\theta}|\psi\rangle \| \neq 0$ for nonzero θ .

2 Which state did Alice make?

Consider a game in which Alice prepares one of two possible states: either ρ_1 with *a priori* probability p_1 or ρ_2 with *a priori* probability $p_2 = 1 - p_1$. Bob

is to perform a measurement and on the basis on the outcome to guess which state Alice prepared. If Bob's guess is right, he wins; if he guesses wrong, Alice wins.

In this exercise you will find Bob's best strategy, and determine his optimal probability of error.

Let's suppose (for now) that Bob performs a POVM with two possible outcomes, corresponding to the two nonnegative Hermitian operators E_1 and $E_2 = I - E_1$. If Bob's outcome is E_1 , he guesses that Alice's state was ρ_1 , and if it is E_2 , he guesses ρ_2 . Then the probability that Bob guesses wrong is

$$p_{\text{error}} = p_1 \text{tr}(\rho_1 E_2) + p_2 \text{tr}(\rho_2 E_1)$$

(a) Show that

$$p_{\text{error}} = p_1 + \sum_i \lambda_i \langle i | E_1 | i \rangle$$

where $\{|i\rangle\}$ denotes the orthonormal basis of eigenstates of the Hermitian operator $p_2\rho_2 - p_1\rho_1$, and the λ_i 's are the corresponding eigenvalues.

Solution:

$$\begin{aligned} p_{\text{error}} &= \text{tr}(p_1 \rho_1 E_2 + p_2 \rho_2 E_1) && (\text{linearity of trace}) \\ &= \text{tr}((p_1 \rho_1)(I - E_1) + p_2 \rho_2 E_1) \\ &= p_1 \text{tr}(\rho_1 I) + \text{tr}(p_2 \rho_2 E_1 - p_1 \rho_1 E_1) && (\text{lin. of trace}) \\ &= p_1 + \text{tr}((p_2 \rho_2 - p_1 \rho_1) E_1) \end{aligned}$$

Since $p_2 \rho_2 - p_1 \rho_1$ is Hermitian, using spectral theorem, $p_2 \rho_2 - p_1 \rho_1 = \sum_i \lambda_i |i\rangle \langle i|$, where $\{|i\rangle\}$ form an orthonormal basis.

$$\begin{aligned} &= p_1 + \text{tr}\left(\left(\sum_i \lambda_i |i\rangle \langle i|\right) E_1\right) && (\text{spectral thm.}) \\ &= p_1 + \sum_i \lambda_i \text{tr}(|i\rangle \langle i| E_1) \\ &= p_1 + \sum_i \lambda_i \langle i | E_1 | i \rangle \end{aligned}$$

as desired.

- (b) Bob's best strategy is to perform the two-outcome POVM that minimizes this error probability. Find the nonnegative operator E_1 that minimizes p_{error} , and show that the error probability when Bob performs this optimal two-outcome POVM is

$$(p_{\text{error}})_{\text{optimal}} = p_1 + \sum_{\text{neg}} \lambda_i.$$

where \sum_{neg} denotes the sum over all of the *negative* eigenvalues of $p_2\rho_2 - p_1\rho_1$.

Solution: E_1 is operator that projects to negative eigenspace of $p_2\rho_2 - p_1\rho_1$, $E_1 = \sum_{\lambda_i < 0} |i\rangle\langle i|$.

$$p_1 + \sum_i \lambda_i \langle i | E_1 | i \rangle = p_1 + \sum_j \lambda_j$$

such that $\lambda_j < 0$. This is optimal since it must be the smallest probability, as if any λ_j in the sum was positive then it would be greater than this value.

- (c) It is convenient to express this optimal error probability in terms of the L^1 norm of the operator $p_2\rho_2 - p_1\rho_1$,

$$\|p_2\rho_2 - p_1\rho_1\|_1 = \text{tr}(|p_2\rho_2 - p_1\rho_1|) = \sum_{\text{pos}} \lambda_i - \sum_{\text{neg}} \lambda_i,$$

the difference between the sum of positive eigenvalues and the sum of negative eigenvalues. Use the property $\text{tr}(p_2\rho_2 - p_1\rho_1) = p_2 - p_1$ to show that

$$(p_{\text{error}})_{\text{optimal}} = \frac{1}{2} - \frac{1}{2} \|p_2\rho_2 - p_1\rho_1\|_1$$

Check whether the answer makes sense in the case where $\rho_1 = \rho_2$ and in the case where ρ_1 and ρ_2 have support on orthogonal subspaces

Solution:

$$\begin{aligned}
tr(p_2\rho_2 - p_1\rho_1) &= p_2 - p_1 = \sum_{pos} \lambda_i + \sum_{neg} \lambda_i \\
\implies \sum_{pos} \lambda_i &= p_2 - p_1 - \sum_{neg} \lambda_i \\
\implies \|p_2\rho_2 - p_1\rho_1\|_1 &= p_2 - p_1 - 2 \sum_{neg} \lambda_i \\
&= 1 - 2p_1 - 2 \sum_{neg} \lambda_i
\end{aligned}$$

Substituting into the expression for $(p_{error})_{optimal}$:

$$\begin{aligned}
(p_{error})_{optimal} &= \frac{1}{2} - \frac{1}{2} \|p_2\rho_2 - p_1\rho_1\|_1 \\
&= \frac{1}{2} - \frac{1}{2} + p_1 + \sum_{neg} \lambda_i \\
&= p_1 + \sum_{neg} \lambda_i
\end{aligned}$$

as desired.

- (d) Now suppose that Alice decides at random (with $p_1 = p_2 = \frac{1}{2}$) to prepare one of two pure states $|\psi_1\rangle, |\psi_2\rangle$ of a single qubit, with

$$|\langle\psi_1|\psi_2\rangle| = \sin(2\alpha), 0 \leq \alpha \leq \frac{\pi}{4}$$

With a suitable choice of basis, the two states can be expressed as:

$$|\psi_1\rangle = \begin{pmatrix} \cos \alpha \\ \sin \alpha \end{pmatrix}, |\psi_2\rangle = \begin{pmatrix} \sin \alpha \\ \cos \alpha \end{pmatrix},$$

Find Bob's optimal two-outcome measurement, and compute the optimal error probability.

Solution: The density operator for a pure state $|\psi\rangle$ is $\rho = |\psi\rangle\langle\psi|$.

$$\begin{aligned}\implies \rho_1 &= |\psi_1\rangle\langle\psi_1| = \begin{pmatrix} \cos^2\alpha & \cos\alpha\sin\alpha \\ \sin\alpha\cos\alpha & \sin^2\alpha \end{pmatrix}, \\ \rho_2 &= |\psi_2\rangle\langle\psi_2| = \begin{pmatrix} \sin^2\alpha & \cos\alpha\sin\alpha \\ \sin\alpha\cos\alpha & \cos^2\alpha \end{pmatrix} \\ \implies p_1\rho_1 - p_2\rho_2 &= \frac{1}{2} \begin{pmatrix} \cos^2\alpha - \sin^2\alpha & 0 \\ 0 & \sin^2\alpha - \cos^2\alpha \end{pmatrix} \\ &= \frac{1}{2} \begin{pmatrix} \cos(2\alpha) & 0 \\ 0 & -\cos(2\alpha) \end{pmatrix} \\ \implies \lambda_{neg} &= -\frac{1}{2}\cos(2\alpha) \\ \implies (p_{error})_{optimal} &= \frac{1}{2} - \frac{1}{2}\cos(2\alpha) \\ \implies E_1 &= |1\rangle\langle 1| \end{aligned}$$

- (e) Bob wonders whether he can find a better strategy if his POVM $\{E_i\}$ has more than two possible outcomes. Let $p(a|i)$ denote the probability that a state a was prepared, given that the measurement outcome was i ; it can be computed using the relations:

$$p_ip(1|i) = p_1p(i|1) = p_1\text{tr}\rho_1E_i,$$

$$p_ip(2|i) = p_2p(i|2) = p_2\text{tr}\rho_2E_i;$$

here $p(i|a)$ denotes the probability that Bob finds measurement outcome i if Alice prepared the state ρ_a , and p_i denotes the probability that Bob finds measurement outcome i , averaged over Alice's choice of state. For each outcome i , Bob will make his decision according to which of the two quantities

$$p(1|i), p(2|i)$$

is the larger; the probability that he makes a mistake is the smaller of these two quantities. This probability of error, given that Bob obtains outcome i , can be written as

$$p_{error}(i) = \min(p(1|i), p(2|i)) = \frac{1}{2} - \frac{1}{2}|p(2|i) - p(1|i)|.$$

Show that the probability of error, averaged over the measurement outcomes, is

$$p_{error} = \sum_i p_ip_{error}(i) = \frac{1}{2} - \frac{1}{2} \sum_i |\text{tr}(p_2\rho_2 - p_1\rho_1)E_i|.$$

Solution:

$$\begin{aligned}
p_{error} &= \sum_i p_i p_{error}(i) \\
&= \sum_i p_i \left(\frac{1}{2} - \frac{1}{2} |p(2|i) - p(1|i)| \right) \\
&= \frac{1}{2} \left(\sum_i p_i \right) - \frac{1}{2} \left(\sum_i p_i |p(2|i) - p(1|i)| \right)
\end{aligned}$$

Since $\sum_i p_i = 1$ and since $p_i > 0$,

$$\begin{aligned}
&= \frac{1}{2} - \frac{1}{2} \sum_i |p_i p(2|i) - p_i p(1|i)| \\
&= \frac{1}{2} - \frac{1}{2} \sum_i |tr(p_2 \rho_2 E_i) - tr(p_1 \rho_1 E_i)| \\
&= \frac{1}{2} - \frac{1}{2} \sum_i |tr(p_2 \rho_2 - p_1 \rho_1) E_i|
\end{aligned}$$

(f) By expanding in terms of the basis of eigenstates of $p_2 \rho_2 - p_1 \rho_1$, show that

$$p_{error} \geq \frac{1}{2} - \frac{1}{2} \|p_2 \rho_2 - p_1 \rho_1\|_1$$

Solution: Suppose we have $E_1 \dots E_n$ s.t. $n > 2$. Then, using spectral decomposition of $p_2 \rho_2 - p_1 \rho_1 = \sum_n \lambda_n |n\rangle \langle n|$

$$\begin{aligned}
p_{error} &= \frac{1}{2} - \frac{1}{2} \sum_i \left| \sum_n \lambda_n tr(|n\rangle \langle n| E_i) \right| \\
&= \frac{1}{2} - \frac{1}{2} \left(\sum_i \left| \sum_n \lambda_n \langle n | E_i | n \rangle \right| \right) \\
&\geq \frac{1}{2} - \frac{1}{2} \left(\left| \sum_n \lambda_n \langle n | \sum_i E_i | n \rangle \right| \right) = \frac{1}{2} - \frac{1}{2} \left(\left| \sum_n \lambda_n \right| \right) \\
&= \frac{1}{2} - \frac{1}{2} \|p_2 \rho_2 - p_1 \rho_1\|_1
\end{aligned}$$