

Resonant tunnelling example.

Resonant tunnelling from a piecewise flat potential:

$$V(x) = 0, \quad |x| > b \text{ and } |x| < a, \\ = v, \quad a < |x| < b,$$

with $b > a > 0$. Define

$$k = \sqrt{2mE} / \hbar, \\ \kappa = \sqrt{2m(v-E)} / \hbar.$$

Define the matrices for matching conditions at $x = \pm a, \pm b$, are

$$\begin{aligned} \mathbf{M1}[a_-] &= \left\{ \left\{ e^{ik a}, e^{-ik a} \right\}, \left\{ ik e^{ik a}, -ik e^{-ik a} \right\} \right\}; \\ \mathbf{M2}[a_-] &= \left\{ \left\{ e^{\kappa a}, e^{-\kappa a} \right\}, \left\{ \kappa e^{\kappa a}, -\kappa e^{-\kappa a} \right\} \right\}; \\ \mathbf{M1}[x] & // \text{MatrixForm} \\ \mathbf{M2}[x] & // \text{MatrixForm} \\ \begin{pmatrix} e^{ikx} & e^{-ikx} \\ ik e^{ikx} & -ik e^{-ikx} \end{pmatrix} \\ \begin{pmatrix} e^{\kappa x} & e^{-\kappa x} \\ \kappa e^{\kappa x} & -\kappa e^{-\kappa x} \end{pmatrix} \end{aligned}$$

Then, if the wave function for $|x| > b$ has the form

$$\begin{aligned} \psi(x) &= A e^{ikx} + B e^{-ikx}, \text{ for } x < -b, \\ &= C e^{ikx}, \text{ for } x > b, \end{aligned}$$

the matching conditions can be written as the matrix equation $\vec{a} = X \vec{c}$ where $\vec{a}^T = (A, B)$, $\vec{c}^T = (C, 0)$, and

$$\mathbf{X} = \text{Simplify}[\text{Inverse}[\mathbf{M1}[-b]] . \mathbf{M2}[-b] . \text{Inverse}[\mathbf{M2}[-a]] . \mathbf{M1}[-a] . \text{Inverse}[\mathbf{M1}[a]] . \mathbf{M2}[a] . \text{Inverse}[\mathbf{M2}[b]] . \mathbf{M1}[b]] ;$$

Then $C/A = 1/X_{11}$ where

$$\mathbf{X11} = \mathbf{X}[[1, 1]]$$

$$\frac{1}{16 k^2 \kappa^2} e^{-2ia\kappa + 2ib\kappa - 2a\kappa - 2b\kappa} \left(-e^{4b\kappa} (k + ik\kappa)^4 - e^{4a\kappa} (ik + \kappa)^4 + \right. \\ \left. 2 e^{2(a+b)\kappa} (k^2 + \kappa^2)^2 + e^{4a\kappa} (ik + \kappa) (k^2 + \kappa^2)^2 - 2 e^{4ia\kappa + 2a\kappa + 2b\kappa} (k^2 + \kappa^2)^2 + e^{4ia\kappa + 4b\kappa} (k^2 + \kappa^2)^2 \right)$$

so the transmission probability is given by $T = 1/|X_{11}|^2$:

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Xexp = ComplexExpand[X11];
Xconj = Xexp /. i -> -i;
Xsquared = Simplify[Expand[Xexp Xconj]]

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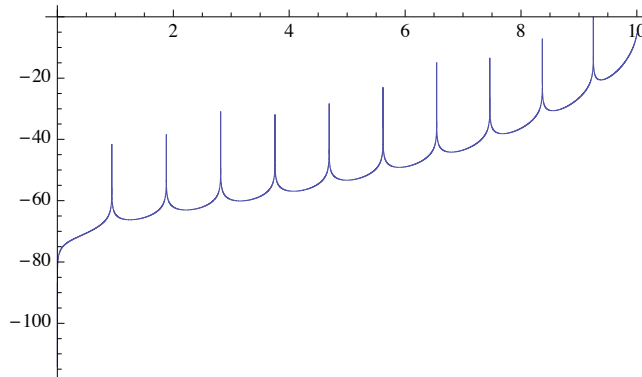
$$\begin{aligned}
& -\frac{1}{128 k^4 \kappa^4} \\
& e^{-4(a+b)\kappa} \left(-e^{8a\kappa} k^8 - e^{8b\kappa} k^8 - 6e^{4(a+b)\kappa} k^8 + 4e^{2(3a+b)\kappa} k^8 + 4e^{2(a+3b)\kappa} k^8 - 4e^{8a\kappa} k^6 \kappa^2 - 4e^{8b\kappa} k^6 \kappa^2 + \right. \\
& 8e^{4(a+b)\kappa} k^6 \kappa^2 - 6e^{8a\kappa} k^4 \kappa^4 - 6e^{8b\kappa} k^4 \kappa^4 - 100e^{4(a+b)\kappa} k^4 \kappa^4 - 8e^{2(3a+b)\kappa} k^4 \kappa^4 - \\
& 8e^{2(a+3b)\kappa} k^4 \kappa^4 - 4e^{8a\kappa} k^2 \kappa^6 - 4e^{8b\kappa} k^2 \kappa^6 + 8e^{4(a+b)\kappa} k^2 \kappa^6 - e^{8a\kappa} \kappa^8 - \\
& e^{8b\kappa} \kappa^8 - 6e^{4(a+b)\kappa} \kappa^8 + 4e^{2(3a+b)\kappa} \kappa^8 + 4e^{2(a+3b)\kappa} \kappa^8 + (e^{2a\kappa} - e^{2b\kappa})^2 (k^2 + \kappa^2)^2 \\
& \left. (-2e^{2(a+b)\kappa} (k^2 + \kappa^2)^2 + e^{4a\kappa} (k^4 - 6k^2 \kappa^2 + \kappa^4) + e^{4b\kappa} (k^4 - 6k^2 \kappa^2 + \kappa^4)) \cos[4ak] + \right. \\
& \left. 4(e^{2a\kappa} - e^{2b\kappa})^3 (e^{2a\kappa} + e^{2b\kappa}) k \kappa (-k^2 + \kappa^2) (k^2 + \kappa^2)^2 \sin[4ak] \right)
\end{aligned}$$

This is a mess, but plot it for some convenient values ($a=\pi/2$, $b=\pi$, $E=n^2$, $v=100$, $m/(2\hbar)=1$), so that the two barriers each have thickness $\pi/2$:

```

WellTransmission = 1 / Xsquared /. {a -> Pi / 2, b -> Pi, k -> n, kappa -> Sqrt[100 - n^2]};
welltranplot = Plot[Log[WellTransmission], {n, 0, 10}, PlotRange -> All, PlotPoints -> 100 000]

```



Note the peaks in T at $n \approx \{1,2,3,4,\dots\}$. These values of n correspond to the bound state energies of the square well potential in the middle ($x < |a|$). These peaks in the transmission probability are called "resonant tunnelling".

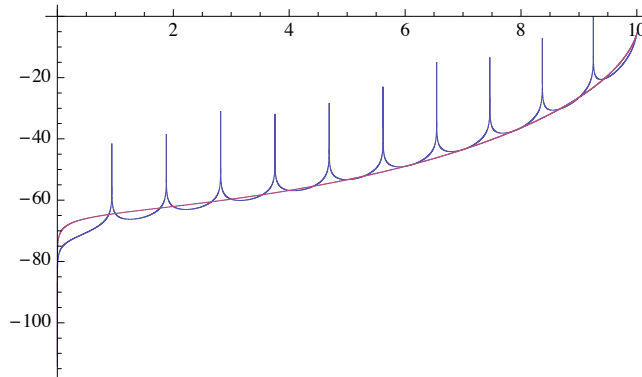
For comparison, consider instead the transmission probability through a single square barrier of thickness π (i.e., without an intervening potential well). Following the same steps as above gives ...

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B = Simplify[Inverse[M1[-b]].M2[-b].Inverse[M2[-a]].M1[-a]];
B11 = B[[1, 1]];
Bexp = ComplexExpand[B11];
Bconj = Bexp /. i -> -i;
Bsquared = Simplify[Expand[Bexp Bconj]]

BarrierTransmission = 1 / Bsquared /. {a -> 0, b -> pi, k -> n, kappa -> Sqrt[100 - n^2]};
Plot[{Log[WellTransmission], Log[BarrierTransmission]},
  {n, 0, 10}, PlotRange -> All, PlotPoints -> 100 000]

```

$$\frac{e^{-2(a+b)\kappa} \left(e^{4a\kappa} (k^2 + \kappa^2)^2 + e^{4b\kappa} (k^2 + \kappa^2)^2 - 2e^{2(a+b)\kappa} (k^4 - 6k^2\kappa^2 + \kappa^4) \right)}{16k^2\kappa^2}$$


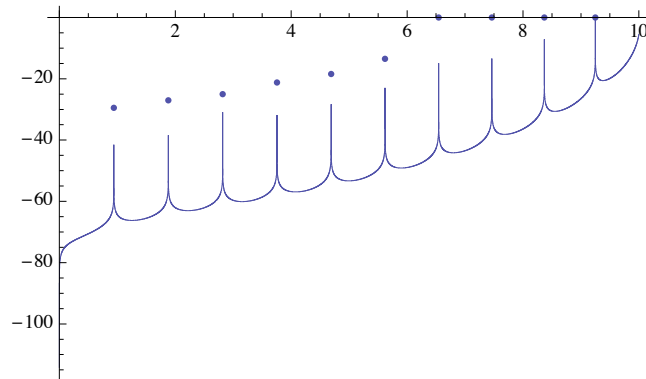
... which closely follows the transmission probability for the case with an intervening potential well, except for the absence of the resonant tunnelling peaks.

[Note: In the plot the resonant tunnelling peaks only seem to rise about 10 orders of magnitude (i.e., a factor of about e^{25}) above the background, but this is just because of the numerical resolution of *Mathematica* --- in fact, they rise to values much closer to 1 (i.e., $\log(T) \approx 0$). Numerical evidence of this is just to evaluate the peak values by searching for relevant extrema of T. This gives the points in the following plot...

```

Off[FindRoot::"lstol"]
Off[Power::"infy"]
peak[1] = FindRoot[D[Denominator[WellTransmission], n] == 0, {n, .95}];
val[1] = {y → Log[Abs[WellTransmission /. peak[1]]]};
peak[2] = FindRoot[D[Denominator[WellTransmission], n] == 0, {n, 1.9}];
val[2] = {y → Log[Abs[WellTransmission /. peak[2]]]};
peak[3] = FindRoot[D[Denominator[WellTransmission], n] == 0, {n, 2.8}];
val[3] = {y → Log[Abs[WellTransmission /. peak[3]]]};
peak[4] = FindRoot[D[Denominator[WellTransmission], n] == 0, {n, 3.75}];
val[4] = {y → Log[Abs[WellTransmission /. peak[4]]]};
peak[5] = FindRoot[D[Denominator[WellTransmission], n] == 0, {n, 4.7}];
val[5] = {y → Log[Abs[WellTransmission /. peak[5]]]};
peak[6] = FindRoot[D[Denominator[WellTransmission], n] == 0, {n, 5.6}];
val[6] = {y → Log[Abs[WellTransmission /. peak[6]]]};
peak[7] = FindRoot[D[Denominator[WellTransmission], n] == 0, {n, 6.55}];
val[7] = {y → Min[Log[Abs[WellTransmission /. peak[7]]], 0]};
peak[8] = FindRoot[D[Denominator[WellTransmission], n] == 0, {n, 7.4}];
val[8] = {y → Min[Log[Abs[WellTransmission /. peak[8]]], 0]};
peak[9] = FindRoot[D[Denominator[WellTransmission], n] == 0, {n, 8.35}];
val[9] = {y → Log[Abs[WellTransmission /. peak[9]]]};
peak[10] = FindRoot[D[Denominator[WellTransmission], n] == 0, {n, 9.24}];
val[10] = {y → Log[Abs[WellTransmission /. peak[10]]]};
peaks = Table[{n /. peak[i], y /. val[i]}, {i, 1, 10}];
peakplot = ListPlot[peaks];
Show[welltranplot, peakplot]

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... showing another factor of e^{10} increase in T over the line plot. Also, the fact that for the last 4 peaks the value is essentially $T=1$ ($\log(T)=0$), indicates that the values at the earlier peaks are probably just limited by numerical accuracy.]