

Infinitely long solenoid

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Ideal (infinitely long) Solenoid

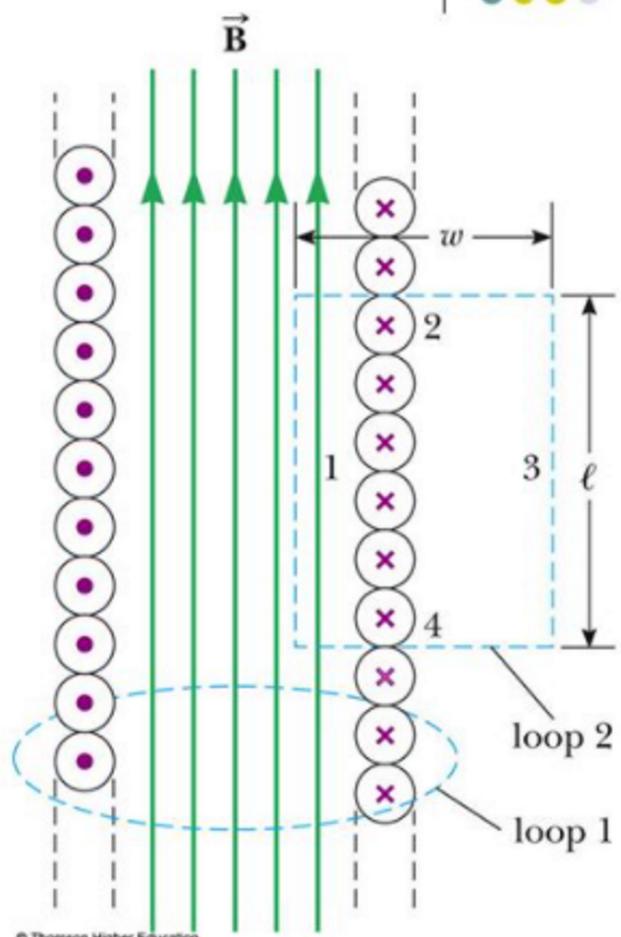


- An *ideal solenoid* is approached when:
 - the turns are closely spaced
 - the length is much greater than the radius of the turns
- Apply Ampere's Law to loop 2:

$$\oint \vec{B} \cdot d\vec{s} = \int_{\text{path 1}} \vec{B} \cdot d\vec{s} = B \int_{\text{path 1}} ds = B\ell$$

$$\oint \vec{B} \cdot d\vec{s} = B\ell = \mu_0 NI$$

$$B = \mu_0 \frac{N}{\ell} I = \mu_0 nI$$

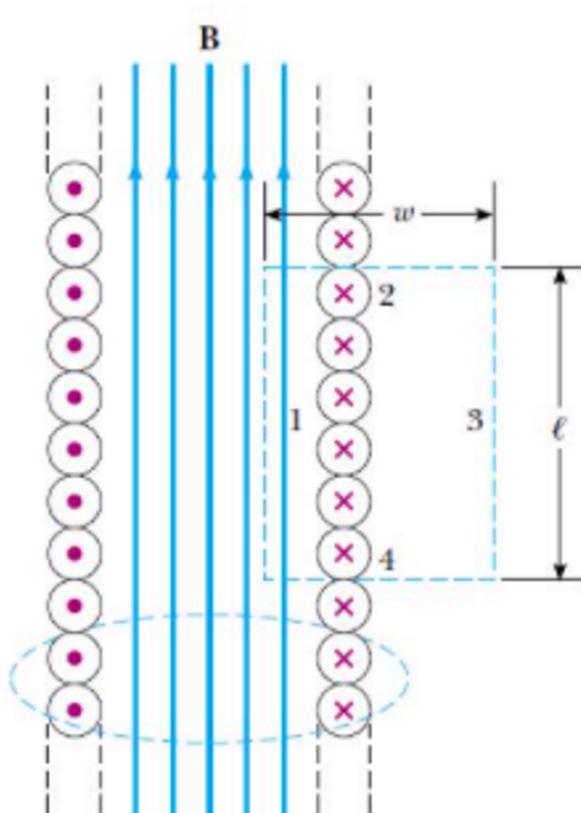


$n = N / \ell$ is the number of turns per unit length

Ampere's Law: Example, Infinitely long solenoid



Cross sectional view:



$$\oint \vec{B} \cdot d\vec{s} = \int_{(1)} \vec{B} \cdot d\vec{s} + \int_{(2)} \vec{B} \cdot d\vec{s} + \int_{(3)} \vec{B} \cdot d\vec{s} + \int_{(4)} \vec{B} \cdot d\vec{s}$$

$\vec{B} \parallel d\vec{s}, B_l$ $\vec{B} \perp d\vec{s}, 0$

$B=0$ outside

$$I_{\text{enc}} = \frac{N}{L} l I, \quad N = \# \text{ of turns}$$

$L = \text{total length of}$
 Solenoid

$I = \text{Current in wire}$

$$B_l = \mu_0 \left(\frac{N}{L} l I \right) \Rightarrow B = \mu_0 n I, \text{ RHR}$$

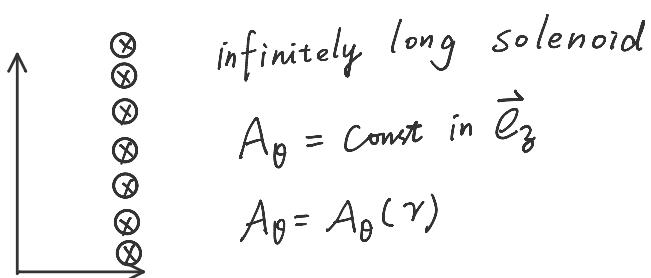
where $n = \# \text{ of turns per length}$

B is uniform inside solenoid !!

Magnetic field outside the solenoid
It was just proved that the magnetic field outside the solenoid is constant, independent of the distance to the

solenoid. Applying the physically plausible boundary condition that the field must go to zero for large distances, one sees that the field is equal to zero everywhere outside the solenoid.

$$\frac{1}{\gamma} \frac{\partial}{\partial r} \left(\gamma \frac{\partial A_\theta}{\partial \gamma} \right) A_\theta - \frac{A_\theta}{\gamma^2} + \frac{\partial}{\partial z^2} A_\theta = -\mu_0 (\mathcal{J}_{coil} + \mathcal{J}_{plasma})$$



$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \right) A_\theta - \frac{A_\theta}{r^2} = -\mu_0 I_{\text{coil}}$$

$$\frac{\partial^2}{\partial r^2} A_\theta + \frac{1}{r} \frac{\partial}{\partial r} A_\theta - \frac{A_\theta}{r^2} = -\mu_0 J_{coil}$$

$$\frac{\partial^2}{\partial \gamma^2} A_\theta = \frac{A_\theta(\gamma_{i+1}) - 2A_\theta(\gamma_i) + A_\theta(\gamma_{i-1})}{\Delta \gamma^2} \rightarrow \frac{1}{\Delta \gamma^2} \begin{bmatrix} -2 & 1 & & & \\ 1 & -2 & 1 & & \\ & \ddots & \ddots & \ddots & \\ & & 1 & -2 & 1 \\ & & & 1 & -2 \end{bmatrix} \begin{bmatrix} A_\theta(\gamma_1) \\ A_\theta(\gamma_2) \\ \vdots \\ A_\theta(\gamma_i) \\ A_\theta(\gamma_{i+1}) \end{bmatrix} = \frac{1}{\Delta \gamma^2} M_i A_\theta$$

$$\frac{1}{\gamma} \frac{\partial}{\partial \gamma} A_\theta = \frac{A_\theta(\gamma_i) - A_\theta(\gamma_{i-1})}{\gamma_i \Delta \gamma} = \frac{A_\theta(\gamma_i) - A_\theta(\gamma_{i-1})}{i \Delta \gamma^2} \rightarrow \frac{1}{\Delta \gamma^2} \begin{bmatrix} \frac{1}{1} & & & \\ -\frac{1}{2} & \frac{1}{2} & & \\ & -\frac{1}{3} & \frac{1}{3} & \\ & & \ddots & \\ & & -\frac{1}{n-1} & \frac{1}{n-1} \\ & & & -\frac{1}{n} & \frac{1}{n} \end{bmatrix} \begin{bmatrix} A_\theta(\gamma_1) \\ A_\theta \\ \vdots \\ A_\theta(\gamma_n) \end{bmatrix} = \frac{1}{\Delta \gamma^2} M_2 A_\theta$$

$$-\frac{A_\theta}{r^2} = -\frac{A_\theta(\gamma_i)}{\gamma_i^2} = -\frac{A_\theta(\gamma_i)}{i^2 \Delta r^2} \rightarrow \frac{1}{\Delta r^2} \begin{bmatrix} -\frac{1}{1^2} & & & \\ & -\frac{1}{2^2} & & \\ & & -\frac{1}{3^2} & \\ & & & \ddots \\ & & & -\frac{1}{(n-1)^2} \\ & & & & -\frac{1}{n^2} \end{bmatrix} \begin{bmatrix} A_\theta(\gamma_1) \\ A_\theta \\ A_\theta(\gamma_n) \end{bmatrix} = \frac{1}{\Delta r^2} M_3 A_\theta$$

$$-\mu_0 J_{coil} = -\mu_0 J_{coil} \delta(r = \gamma_{coil}) \rightarrow \begin{bmatrix} 0 \\ \vdots \\ -\mu_0 J_{coil} \\ \vdots \\ 0 \end{bmatrix} \rightarrow \gamma_j = \gamma_{coil}$$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \right) A_\theta - \frac{A_\theta}{r^2} = -\mu_0 J_{coil} \rightarrow \boxed{\frac{1}{\Delta r^2} (M_1 + M_2 + M_3) A_\theta = -\mu_0 J_{coil}}$$

$$\vec{B} = \nabla \times \vec{A}$$

$$\nabla \times \vec{A} = \begin{aligned} & \left(\frac{1}{\rho} \frac{\partial A_z}{\partial \phi} - \frac{\partial A_\phi}{\partial z} \right) \hat{\rho} + \\ & \left(\frac{\partial A_\rho}{\partial z} - \frac{\partial A_z}{\partial \rho} \right) \hat{\phi} + \\ & \frac{1}{\rho} \left(\frac{\partial (\rho A_\phi)}{\partial \rho} - \frac{\partial A_\rho}{\partial \phi} \right) \hat{z} \end{aligned}$$

$$\vec{B} = B_z(r) \vec{e}_z = \frac{1}{r} \left(\frac{\partial (r A_\theta)}{\partial r} - \cancel{\frac{\partial A_r}{\partial \theta}} \right) \vec{e}_z$$

$$= \frac{1}{r} \frac{\partial (r A_\theta)}{\partial r} \vec{e}_z$$

$$= \left(\frac{\partial A_\theta}{\partial r} + \frac{A_\theta}{r} \right) \vec{e}_z$$

$$B_z(r) = \frac{\partial A_\theta}{\partial r} + \frac{A_\theta}{r}$$

$$B_z(\gamma_i) = \frac{A_\theta(\gamma_i) - A_\theta(\gamma_{i-1})}{\Delta r} + \frac{A_\theta(\gamma_i)}{i \Delta r} = \frac{1}{\Delta r} \left(1 + \frac{1}{i} \right) A_\theta(\gamma_i) - \frac{1}{\Delta r} A_\theta(\gamma_{i-1})$$

$$\begin{bmatrix} B_3 \end{bmatrix} = \frac{1}{\Delta\gamma} \begin{bmatrix} 1+\frac{1}{1} & & & & A_0(\gamma_1) \\ -1 & 1+\frac{1}{2} & & & \\ & -1 & 1+\frac{1}{3} & & \\ & & \ddots & & \\ & & & -1 & 1+\frac{1}{n-1} \\ & & & & -1 & 1+\frac{1}{n} \end{bmatrix} \begin{bmatrix} A_0 \\ A_0 \\ \vdots \\ A_0 \end{bmatrix}$$

$$B(\gamma_1) = 2A_0(\gamma_1)$$

$$B(\gamma_2) = \left(1 + \frac{1}{2}\right)A_0(\gamma_2) - A_0(\gamma_1)$$

$$B(\gamma_2) = B(\gamma_1) \Rightarrow 1.5A_0(\gamma_2) - A_0(\gamma_1) = 2A_0(\gamma_1)$$

$$1.5A_0(\gamma_2) = 3A_0(\gamma_1)$$

$$A_0(\gamma_2) = 2A_0(\gamma_1)$$

$$B(\gamma_3) = B(\gamma_1) \Rightarrow \left(1 + \frac{1}{3}\right)A_0(\gamma_3) - A_0(\gamma_2) = 2A_0(\gamma_1)$$

$$\left(1 + \frac{1}{3}\right)A_0(\gamma_3) - 2A_0(\gamma_1) = 2A_0(\gamma_1)$$

$$\frac{4}{3}A_0(\gamma_3) = 4A_0(\gamma_1)$$

$$A_0(\gamma_3) = 3A_0(\gamma_1)$$



$$A_0(\gamma_i) = iA_0(\gamma_1)$$