

计算物理作业 6

杨远青 22300190015

CompPhys 24

2024 年 11 月 6 日

起视四境，而作业又至矣。

1 题目 1：一维 Kronig-Penney 模型的本征值求解

1.1 题目描述

One-dimensional Kronig-Penney problem. Considering the Hamiltonian of the system as

$$\hat{H} = -\frac{\hbar^2}{2m_e} \frac{\partial^2}{\partial x^2} + V(x)$$

with a one-dimensional periodic potential $V(x) = V(x + a)$. The potential can be expressed as

$$V(x) = \begin{cases} 0, & \text{if } 0 \leq x < L_W, \\ U_0, & \text{if } L_W \leq x < a \end{cases}$$

and the period of the potential is $a = L_W + L_B$, which is also shown in the Figure below.

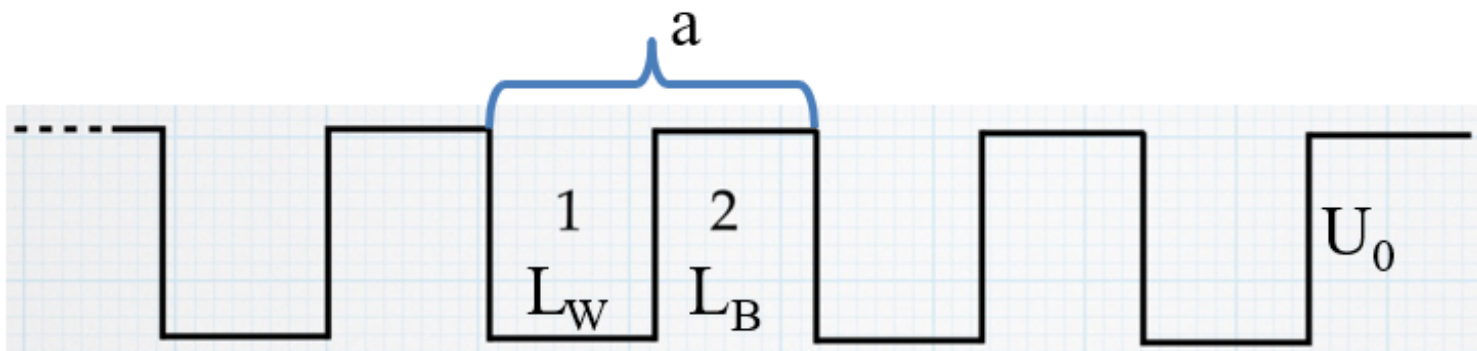


图 1: Kronig-Penney potential well

With parameters:

$$U_0 = 2 \text{ eV}, \quad L_W = 0.9 \text{ nm}, \quad L_B = 0.1 \text{ nm} \quad (a = 1 \text{ nm})$$

Using FFT, find the lowest three eigenvalues of the electric eigenstates that satisfy

$$\hat{H}\psi_i = E\psi_i \quad \text{and} \quad \psi_i(x) = \psi_i(x + a).$$

Explanation: Since the system is translation-invariant, i.e., $\psi_i(x) = \psi_i(x + a)$, we can use the plane wave basis expansion

$$\psi(x) = \frac{1}{\sqrt{a}} \sum_q C_q e^{iq \frac{2\pi}{a} x}, \quad q = -N, -N+1, \dots, -1, 0, 1, \dots, N-1, N.$$

In this basis set, the Hamiltonian can be represented in matrix form as

$$H_{pq} = \frac{1}{a} \langle e^{ip \frac{2\pi}{a} x} | \hat{H} | e^{iq \frac{2\pi}{a} x} \rangle_{\text{cell}} = \frac{1}{a} \int_0^a dx e^{-ip \frac{2\pi}{a} x} \hat{H} e^{iq \frac{2\pi}{a} x}.$$

To calculate $\hat{H} e^{iq \frac{2\pi}{a} x}$, the periodic potential $V(x)$ can be expanded in Fourier series as $V(x) \rightarrow V_q$, where

$$V(x) = \sum_{q'=-N}^N V_{q'} e^{iq' \frac{2\pi}{a} x}.$$

The basis wave function can then be written as:

$$\hat{H} e^{iq \frac{2\pi}{a} x} = (\hat{T} + \hat{V}) e^{iq \frac{2\pi}{a} x} = \frac{2\hbar^2 q^2 \pi^2}{m_e a^2} e^{iq \frac{2\pi}{a} x} + \sum_{q'=-N}^N V_{q'} e^{i(q'+q) \frac{2\pi}{a} x}$$

Try constructing the Hamiltonian matrix H_{pq} and solve the eigenvalue equation $\hat{H} \psi_i = E \psi_i$ to obtain the three lowest energy eigenvalues.

Special note: You can use built-in functions to simplify the eigenvalue calculations and FFT transformations.

1.2 程序描述

这道题的过程和结果都太神奇了。

我们先顺着题目未完成的思路进行下来，计算哈密顿矩阵元

$$H_{pq} = \frac{1}{a} \langle e^{ip \frac{2\pi}{a} x} | \hat{T} | e^{iq \frac{2\pi}{a} x} \rangle + \frac{1}{a} \langle e^{ip \frac{2\pi}{a} x} | \hat{V} | e^{iq \frac{2\pi}{a} x} \rangle$$

其中动能项是对角阵

$$\frac{1}{a} \langle e^{ip \frac{2\pi}{a} x} | \hat{T} | e^{iq \frac{2\pi}{a} x} \rangle = \frac{\hbar^2}{2m_e} \left(\frac{2\pi q}{a} \right)^2 \cdot \frac{1}{a} \int_0^a e^{-ip \frac{2\pi}{a} x} e^{iq \frac{2\pi}{a} x} dx = \frac{\hbar^2}{2m_e} \left(\frac{2\pi q}{a} \right)^2 \delta_{pq} = \frac{2\hbar^2 \pi^2 q^2}{m_e a^2} \delta_{pq}$$

利用

$$\langle e^{ip \frac{2\pi}{a} x} | e^{iq' \frac{2\pi}{a} x} | e^{iq \frac{2\pi}{a} x} \rangle = \int_0^a e^{-ip \frac{2\pi}{a} x} e^{iq' \frac{2\pi}{a} x} e^{iq \frac{2\pi}{a} x} dx = \int_0^a e^{i(q'+q-p) \frac{2\pi}{a} x} dx = a \delta_{p-q, q'}$$

可以发现，势能矩阵元很自然地承载势能的傅里叶系数

$$\frac{1}{a} \langle e^{ip \frac{2\pi}{a} x} | V(x) | e^{iq \frac{2\pi}{a} x} \rangle = \sum_{q'=-N}^N V_{q'} \cdot \frac{1}{a} \langle e^{ip \frac{2\pi}{a} x} | e^{iq' \frac{2\pi}{a} x} | e^{iq \frac{2\pi}{a} x} \rangle = \sum_{q'=-N}^N V_{q'} \delta_{p-q, q'} = V_{p-q}$$

注意到我们对实数域的势能展开时 $V_{-q'} = V_{q'}^*$ ，所以最后的哈密顿矩阵是厄米阵。

不过这里有个小漏洞，原本 $p - q$ 的取值应该在 $[-2N, 2N]$ ，但我们在傅里叶展开势能为 $V_{q'}$ 的时候只展开到了 $q' \in [-N, N]$ ，所以按题意需要将 $p - q$ 截断在 $p - q \in [-N, N]$ 之间（在赋值的时候加一组布尔掩码），这样才能保证 $\sum_{q'=-N}^N V_{q'} \delta_{p-q, q'} = V_{p-q}$ 。

我原本认为这样是不合理的，应该将势能展开到 $q' \in [-2N, 2N]$ ，而不是选择牺牲势能项的精度。但在助教提示下发现，本题的势能项本身影响就不大，所以这样的截断是可以接受的。在主程序中我也对这一点进行了对比论证，详见结果示例。

在最初阅读题干时，我最不明白的是，如此简单的阶梯函数已经可以解析计算出势能的傅里叶展开

$$V_{q'} = \begin{cases} \frac{U_0 L_B}{a}, & q = 0 \\ \frac{U_0}{iq'2\pi} \left(1 - e^{-iq'2\pi \frac{L_W}{a}}\right), & q \neq 0 \end{cases}$$

为何还需要进行 FFT 变换？随后想到这节课毕竟是在学习 FFT，故在主程序中加入了 FFT 变换（先将实空间离散化再传递给 `scipy.fft` 进行变换）的系数与解析计算的系数进行对比，发现两者差异在容许范围内，后面的计算则是基于解析计算的系数进行的。

源代码在 `src/KP_well.py` 中，请在该目录下运行 `python -u KP_well.py` 即可得到结果，需要安装 `numpy` 与 `scipy` 库辅助运算，`matplotlib` 库绘图。结果非常神奇，在大规模的绘制与计算中，我们发现了疑似二重简并的现象。在与 QM1 荣誉课的助教老师讨论后，合理怀疑是有限 N 的截断带来的边界效应，且在高能级时非局域化更明显，两侧的对称传播似乎是合理的。可惜在 `debug` 时我尝试过绘制其它相邻能级的波函数概率分布，未发现明显对称性。二重简并的原因有待进一步探讨，但能级的 n^2 增长是可以确认的。

主程序流程描述

1. **初始化物理常量与参数：**从 `const` 获取需要使用的物理常数（如 \hbar 、 m_e 等），定义参数，如 U_0 、 L_w 和 L_b 。
2. **傅里叶系数计算对比：**均采用自然顺序 $[0, 1, 2, \dots, N, -N, \dots, -1]$ ，便于切片负数索引
 - 使用 `compute_Vq_analytical` 方法计算势能傅里叶系数，基于解析表达式，
 - 使用 `compute_Vq_fft` 方法，通过 FFT 计算傅里叶系数，验证解析方法的精确性。
 - 比较两种方法的结果，输出最大和平均差异。
3. **小规模收敛性检查：**借助 `check_convergence` 函数，在较小 N 下进行能级收敛性测试。
4. **AB 方法对比：**
 - **方法 A：**在傅里叶系数范围 $Vq' \in [-N, N]$ 下进行求解，得到小规模能级，
 - **方法 B：**在扩展傅里叶系数范围 $Vq' \in [-2N, 2N]$ 下进行求解，
 - 比较两种方法的能级差异，在之后计算中遵循题意采用方法 A。
5. **中等规模 $N = 100$ 计算与绘制：**
 - 使用解析方法计算的傅里叶系数，构建哈密顿量矩阵，通过 `solve_eigenvalues` 方法得到能级和波函数，
 - 使用 `reconstruct_wavefunction_vectorized` 重构波函数，充分考虑基底正交性，归一化操作不需要积分，
 - 通过 `plot_wavefunctions_and_potential` 绘制势阱和前三个波函数，
 - 绘制能级图，发现在小规模 N 中未出现的二级简并现象。
6. **大规模计算与拟合（可选）：**在用户确认的情况下
 - 使用更大 $N = 500$ 计算较多能级，
 - 通过 `curve_fit` 函数拟合能级间距与量子数的平方增长关系，
 - 输出拟合系数并绘制拟合结果。

1.3 伪代码

Powered by **L^AT_EX** pseudocode generator

Algorithm 1: Main Process Flow

Input: $U_0, L_w, L_b, N_{initial}, num_levels_{initial}, medium_N, medium_num_levels, large_N, large_num_levels$

Output: Energy levels and wave functions, potential well plots, convergence analysis

- 1 **Step 1: Fourier Coefficients Comparison;**
- 2 $Vq_analytical, Vq_fft \leftarrow compare_fourier_methods(N_{initial}, U_0, L_w, L_b, a);$
- 3 **if** $compare_fourier_methods$ accuracy is high **then**
- 4 | Proceed to next step;
- 5 **end**
- 6 **Step 2: Convergence Check;**
- 7 $check_convergence(U_0, L_w, L_b, a, \hbar, m_e);$
- 8 **Step 3: Energy Levels Calculation with Method A and B;**
- 9 **for** $Method$ in $[A, B]$ **do**
- 10 | $eigenvalues, eigenvectors \leftarrow build_and_solve(Method\ parameters);$
- 11 | **if** degeneracies found **then**
- 12 | | Print "Degeneracies exist";
- 13 | **end**
- 14 | **else**
- 15 | | Print "No degeneracies found";
- 16 | **end**
- 17 **end**
- 18 **Step 4: Plotting for Medium N;**
- 19 $eigenvalues_medium, eigenvectors_medium \leftarrow build_and_solve(medium_N, medium_num_levels);$
- 20 $x_medium, wavefunctions_medium \leftarrow reconstruct_wavefunction_vectorized(eigenvectors_medium);$
- 21 $V_x_medium \leftarrow generate_potential_vectorized(x_medium);$
- 22 $plot_wavefunctions_and_potential(x_medium, wavefunctions_medium, V_x_medium);$
- 23 $plot_energy_levels(eigenvalues_medium, medium_N);$
- 24 **if** Proceed to large scale **then**
- 25 | Repeat plotting with $large_N$ and fit quadratic relation;
- 26 **end**

Algorithm 2: Compute Analytical Fourier Coefficients

Input: N, U_0, L_w, L_b, a

Output: Fourier coefficients Vq

```
1  $n\_list \leftarrow \text{get\_n\_list}(N)$ ;  
2 for Each non-zero  $n$  in  $n\_list$  do  
3   |   Compute  $Vq[n]$  based on analytical formula;  
4 end  
5  $Vq[0] \leftarrow U_0 \cdot L_b/a$ ;  
6 return  $Vq$ 
```

Algorithm 3: Build Kinetic Matrix

Input: q_values, a, \hbar, m_e

Output: Kinetic matrix T

```
1 for Each  $q$  in  $q\_values$  do  
2   |    $T[q, q] \leftarrow \frac{\hbar^2 \cdot (2\pi q/a)^2}{2m_e}$ ;  
3 end  
4 return  $T$  as a sparse diagonal matrix
```

Algorithm 4: Solve for Eigenvalues with Error Handling

Input: Hamiltonian matrix H , desired levels num_levels

Output: Eigenvalues and eigenvectors, or an error message if solving fails

```
1 if eigsh can solve for eigenvalues then  
2   |   (eigenvalues, eigenvectors)  $\leftarrow \text{eigsh}(H, k = num\_levels, which = "SA")$ ;  
3   |   Sort eigenvalues and eigenvectors by eigenvalue magnitude;  
4   |   return eigenvalues, eigenvectors;  
5 end  
6 else  
7   |   Print "Error in eigenvalue solver";  
8   |   return  $None, None$ ;  
9 end
```

Algorithm 5: Reconstruct Wavefunction

Input: Eigenvectors, q_values, a

Output: Wavefunctions over extended domain

```
1 for Each eigenvector  $v$  do  
2   |   for Each position  $x$  in domain do  
3     |    $\psi(x) \leftarrow v \cdot \exp(2\pi i \cdot q \cdot x/a)$ ;  
4     |   Normalize  $\psi$  over extended domain;  
5   |   end  
6 end
```

1.4 结果示例

第一步：比较解析方法和FFT方法计算傅里叶系数的差异

傅里叶系数计算方法比较：

傅里叶级数范围： $[-50, 50]$

解析方法用时： $0.00e+00$ 秒

FFT方法用时： $5.12e-04$ 秒

最大绝对差异： $2.058607e-02$ eV

平均绝对差异： $1.223039e-02$ eV

解析方法和FFT方法计算的傅里叶系数差异非常小。

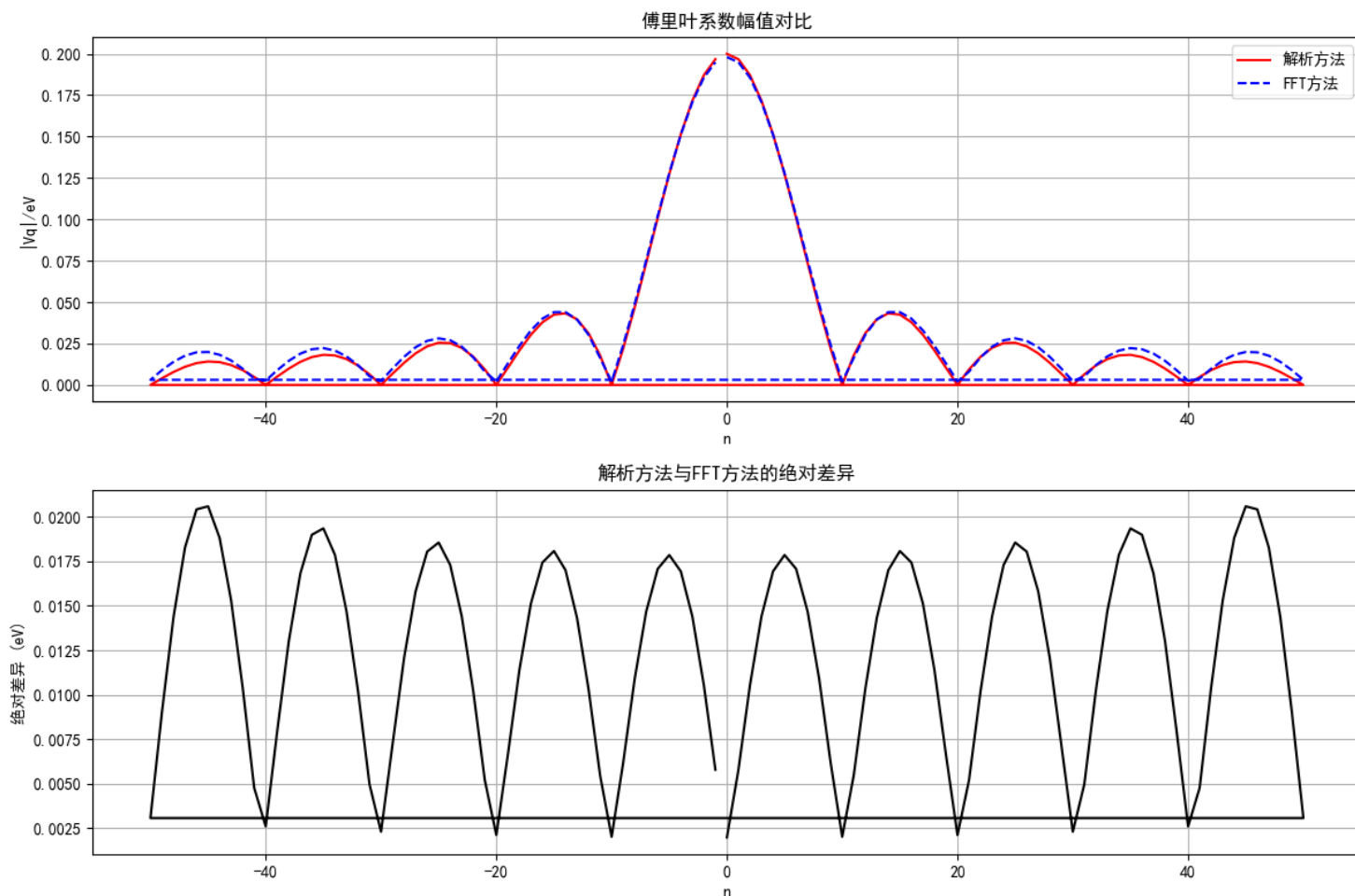


图 2: 第一步：比较解析计算与 FFT 得到的势能傅里叶系数

```
第二步：进行收敛性检查（仅对小规模N进行）
N=10, Lowest 3 eigenvalues (eV): [0.14464301 1.51649618 1.88295936]
No degeneracies found.
N=20, Lowest 3 eigenvalues (eV): [0.1446063 1.51647998 1.88283855]
No degeneracies found.
N=30, Lowest 3 eigenvalues (eV): [0.14460216 1.51647846 1.88282499]
No degeneracies found.
Converged at N=30
```

图 3: 第二步：收敛性检查

第三步：使用解析方法计算的傅里叶系数进行AB方法对比

```
方法A ( $V_{q'} \in [-50, 50]$ )
Lowest 3 eigenvalues (eV):
  Level 1: 1.4460e-01 eV
  Level 2: 1.5165e+00 eV
  Level 3: 1.8828e+00 eV
No degeneracies found.
```

```
方法B ( $V_{q'} \in [-100, 100]$ )
Lowest 3 eigenvalues (eV):
  Level 1: 1.4460e-01 eV
  Level 2: 1.5165e+00 eV
  Level 3: 1.8828e+00 eV
No degeneracies found.
```

方法A和方法B的前3个能级相对差异: [2.39109506e-06 7.59823372e-08 6.01067236e-07]

方法A和方法B的能级差异非常小。

图 4: 第三步：不同势能展开范围的 AB 方法对比

第四步：使用中等规模 $N=100$ 进行能级绘制

$N=100$, Lowest 20 eigenvalues (eV):

```
Level 1: 1.4460e-01 eV
Level 2: 1.5165e+00 eV
Level 3: 1.8828e+00 eV
Level 4: 6.0633e+00 eV
Level 5: 6.3848e+00 eV
Level 6: 1.3633e+01 eV
Level 7: 1.3855e+01 eV
Level 8: 2.4216e+01 eV
Level 9: 2.4325e+01 eV
Level 10: 3.7801e+01 eV
Level 11: 3.7810e+01 eV
Level 12: 5.4321e+01 eV
Level 13: 5.4379e+01 eV
Level 14: 7.3860e+01 eV
Level 15: 7.3946e+01 eV
Level 16: 9.6426e+01 eV
Level 17: 9.6503e+01 eV
Level 18: 1.2201e+02 eV
Level 19: 1.2206e+02 eV
Level 20: 1.5061e+02 eV
```

Degeneracies found.

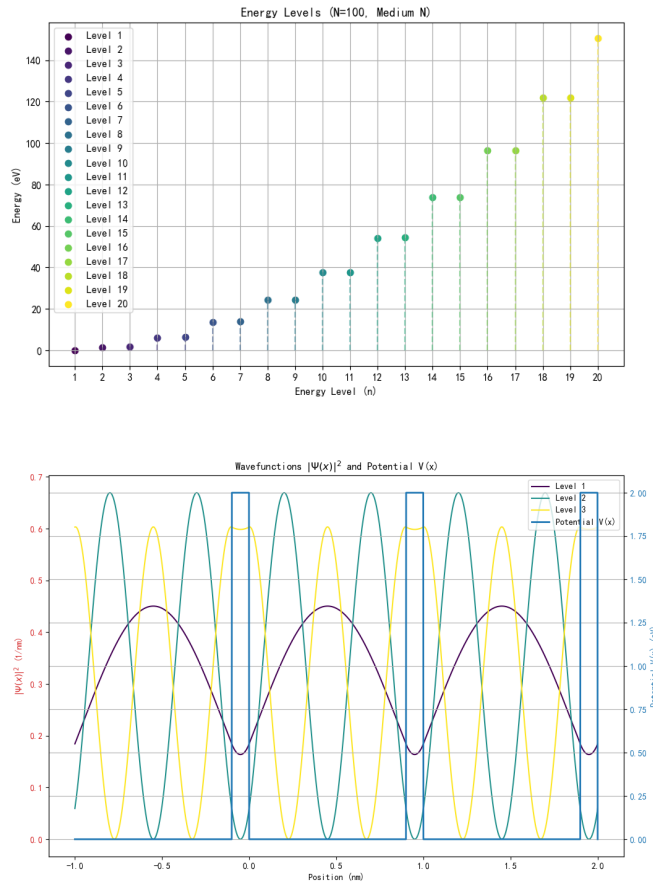


图 5: 第四步：进行中等规模 $N = 100$, $num_levels = 20$ 的计算与绘制，并绘制势阱与前三个能级的波函数


```
是否进行大规模N=500的能级绘制与拟合? (y/n) : y
继续进行大规模N的能级绘制与拟合
N=500, Lowest 100 eigenvalues (eV):
Level 1: 1.4460e-01 eV
Level 2: 1.5165e+00 eV
Level 3: 1.8828e+00 eV
Level 4: 6.0633e+00 eV
Level 5: 6.3848e+00 eV
Level 6: 1.3633e+01 eV
Level 7: 1.3855e+01 eV
Level 8: 2.4216e+01 eV
Level 9: 2.4325e+01 eV
Level 10: 3.7801e+01 eV
Level 11: 3.7810e+01 eV
Level 12: 5.4321e+01 eV
Level 13: 5.4379e+01 eV
Level 14: 7.3860e+01 eV
Level 15: 7.3946e+01 eV
Level 16: 9.6426e+01 eV
Level 17: 9.6503e+01 eV
Level 18: 1.2201e+02 eV
Level 19: 1.2206e+02 eV
Level 20: 1.5061e+02 eV
Level 21: 1.5061e+02 eV
Level 22: 1.8218e+02 eV
Level 23: 1.8222e+02 eV
Level 24: 2.1677e+02 eV
Level 25: 2.1682e+02 eV
Level 26: 2.5437e+02 eV
Level 27: 2.5442e+02 eV
Level 28: 2.9499e+02 eV
Level 29: 2.9502e+02 eV
Level 30: 3.3863e+02 eV
Level 31: 3.3863e+02 eV
Level 32: 3.8524e+02 eV
Level 33: 3.8527e+02 eV
Level 34: 4.3487e+02 eV
Level 35: 4.3487e+02 eV
Level 36: 4.8752e+02 eV
Level 37: 4.8755e+02 eV
Level 38: 5.4318e+02 eV
Level 39: 5.4320e+02 eV
Level 40: 6.0185e+02 eV
Level 41: 6.0185e+02 eV
Level 42: 6.6351e+02 eV
Level 43: 6.6353e+02 eV
Level 44: 7.2818e+02 eV
Level 45: 7.2821e+02 eV
Level 46: 7.9587e+02 eV
Level 47: 7.9589e+02 eV
Level 48: 8.6657e+02 eV
Level 49: 8.6658e+02 eV
Level 50: 9.4028e+02 eV
Level 51: 9.4028e+02 eV
Level 52: 1.0170e+03 eV
Level 53: 1.0170e+03 eV
Level 54: 1.0967e+03 eV
Level 55: 1.0967e+03 eV
Level 56: 1.1794e+03 eV
Level 57: 1.1794e+03 eV
Level 58: 1.2652e+03 eV
Level 59: 1.2652e+03 eV
Level 60: 1.3539e+03 eV
Level 61: 1.3539e+03 eV
Level 62: 1.4457e+03 eV
Level 63: 1.4457e+03 eV
Level 64: 1.5404e+03 eV
Level 65: 1.5404e+03 eV
Level 66: 1.6382e+03 eV
Level 67: 1.6382e+03 eV
Level 68: 1.7390e+03 eV
Level 69: 1.7390e+03 eV
Level 70: 1.8427e+03 eV
Level 71: 1.8427e+03 eV
Level 72: 1.9495e+03 eV
Level 73: 1.9495e+03 eV
Level 74: 2.0593e+03 eV
Level 75: 2.0593e+03 eV
Level 76: 2.1721e+03 eV
Level 77: 2.1722e+03 eV
Level 78: 2.2880e+03 eV
Level 79: 2.2880e+03 eV
Level 80: 2.4068e+03 eV
Level 81: 2.4068e+03 eV
Level 82: 2.5286e+03 eV
Level 83: 2.5286e+03 eV
Level 84: 2.6535e+03 eV
Level 85: 2.6535e+03 eV
Level 86: 2.7813e+03 eV
Level 87: 2.7813e+03 eV
Level 88: 2.9122e+03 eV
Level 89: 2.9122e+03 eV
Level 90: 3.0460e+03 eV
Level 91: 3.0460e+03 eV
Level 92: 3.1829e+03 eV
Level 93: 3.1829e+03 eV
Level 94: 3.3228e+03 eV
Level 95: 3.3228e+03 eV
Level 96: 3.4657e+03 eV
Level 97: 3.4657e+03 eV
Level 98: 3.6116e+03 eV
Level 99: 3.6116e+03 eV
Level 100: 3.7605e+03 eV
Degeneracies found.
拟合结果: k = 0.371510 eV
```

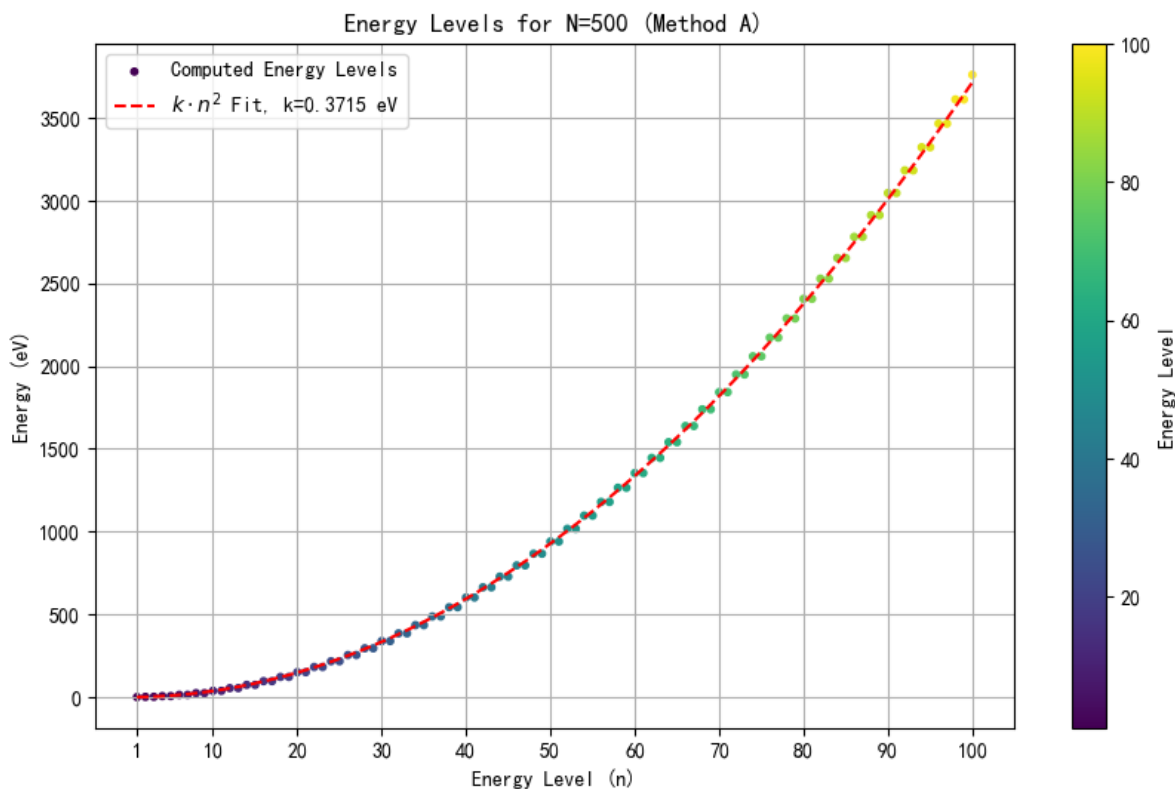


图 6: 第五步: 进行大规模 $N = 500, num_levels = 100$ 的计算与绘制, 并进行拟合

2 题目 2：太阳黑子周期性检测

2.1 题目描述

Use the file called `sunspots.txt`, which contains the observed number of sunspots on the Sun for each month since January 1749. Write a program to calculate the Fourier transform of the sunspot data and then make a graph of the magnitude squared $|c_k|^2$ of the Fourier coefficients as a function of k —also called the power spectrum of the sunspot signal. You should see that there is a noticeable peak in the power spectrum at a nonzero value of k . Find the approximate value of k to which the peak corresponds. What is the period of the sine wave with this value of k ?

Special note: You may use any built-in functions for the Fourier transform.

2.2 程序描述

使用 `scipy` 库的 `detrend` 函数去除线性趋势，也进行了使用 Hanning 窗平滑过渡原信号

$$w(n) = 0.5 \left(1 - \cos \left(\frac{2\pi n}{N-1} \right) \right)$$

的效果对比，该操作有效避免边缘频率泄露，最后使用 `np.fft` 进行快速傅里叶变换，得到频谱图并使用 `np.fft.ifft` 进行信号重构的对比。最终峰值频率对应周期为 $130.6m \approx 10.9a$ ，与太阳黑子爆发周期实验观测值相符。

源代码在 `src/analysis.py` 中，请在该目录下运行 `python -u analysis.py` 即可得到结果，需要安装 `numpy` 与 `scipy` 库辅助运算，`matplotlib` 库绘图，其他平台运行可能需要调整字体管理代码。

2.3 结果示例

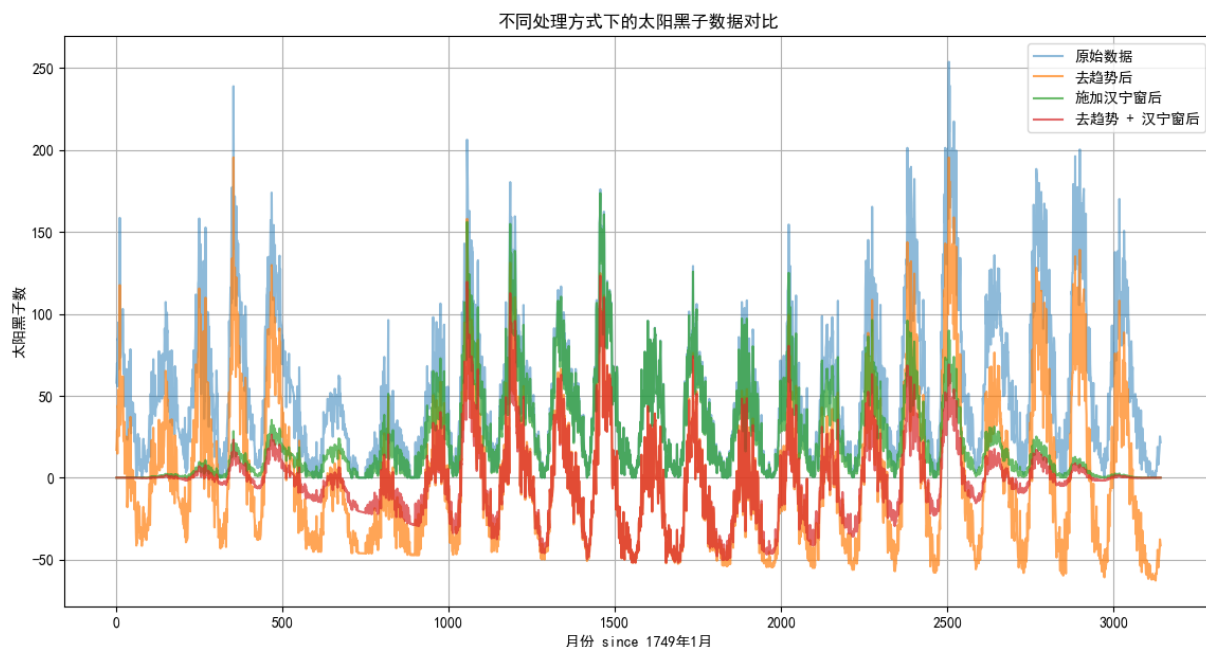


图 7: 不同方式处理的数据对比

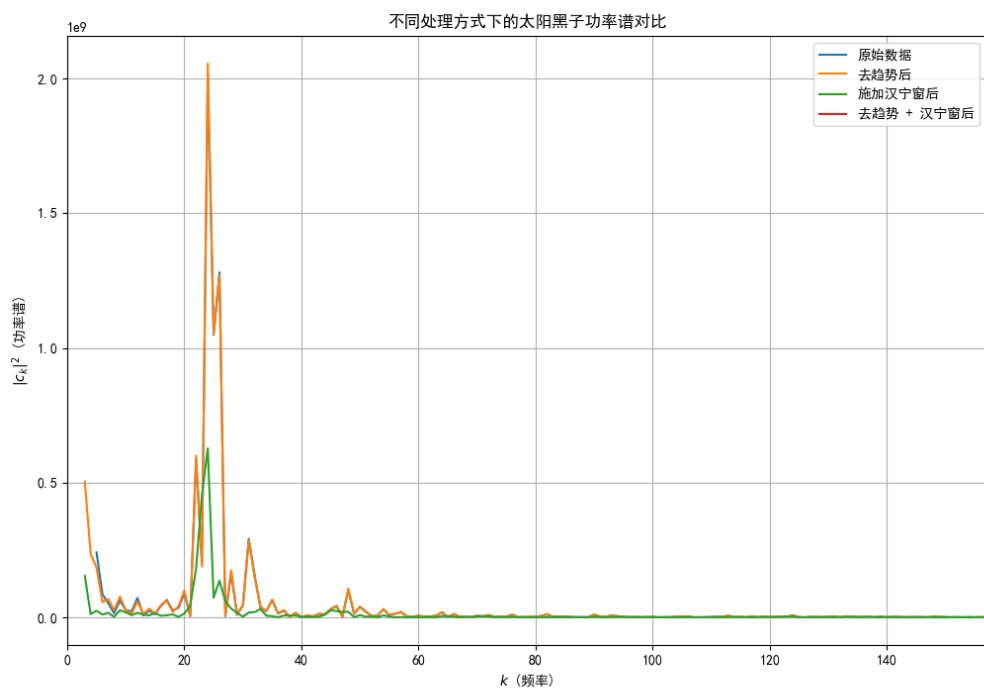


图 8: 去趋势 + 汉宁窗处理后的功率谱

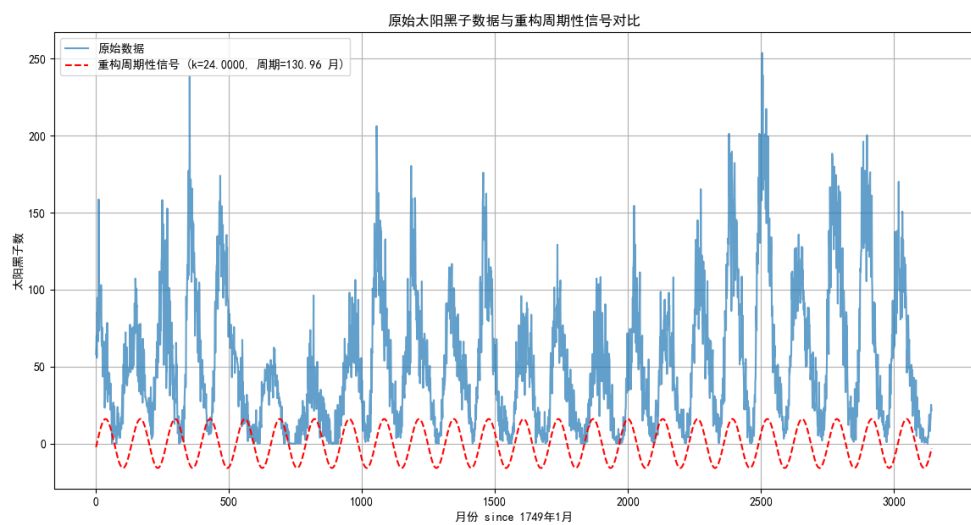


图 9: 使用峰值频率重构与原信号对比