计算物理作业6

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起视四境, 而作业又至矣。

1 题目 1: 一维 Kronig-Penney 模型的本征值求解

1.1 题目描述

One-dimensional Kronig-Penney problem. Considering the Hamiltonian of the system as

$$\hat{H} = -\frac{\hbar^2}{2m_o} \frac{\partial^2}{\partial x^2} + V(x)$$

with a one-dimensional periodic potential V(x) = V(x+a). The potential can be expressed as

$$V(x) = \begin{cases} 0, & \text{if } 0 \le x < L_W, \\ U_0, & \text{if } L_W \le x < a \end{cases}$$

and the period of the potential is $a = L_W + L_B$, which is also shown in the Figure below.

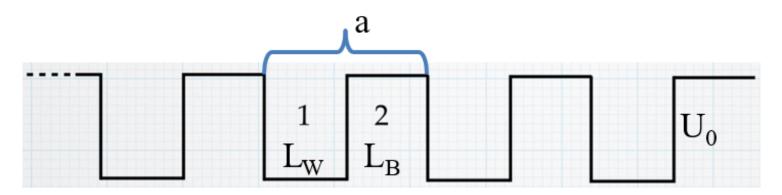


图 1: Kronig-Penney potential well

With parameters:

$$U_0 = 2 \,\text{eV}, \quad L_W = 0.9 \,\text{nm}, \quad L_B = 0.1 \,\text{nm} \quad (a = 1 \,\text{nm})$$

Using FFT, find the lowest three eigenvalues of the electric eigenstates that satisfy

$$\hat{H}\psi_i = E\psi_i$$
 and $\psi_i(x) = \psi_i(x+a)$.

Explanation: Since the system is translation-invariant, i.e., $\psi_i(x) = \psi_i(x+a)$, we can use the plane wave basis expansion

$$\psi(x) = \frac{1}{\sqrt{a}} \sum_{q} C_q e^{iq\frac{2\pi}{a}x}, \quad q = -N, -N+1, \dots, -1, 0, 1, \dots, N-1, N.$$

In this basis set, the Hamiltonian can be represented in matrix form as

$$H_{pq} = \frac{1}{a} \langle e^{ip\frac{2\pi}{a}x} | \hat{H} | e^{iq\frac{2\pi}{a}x} \rangle_{\text{cell}} = \frac{1}{a} \int_0^a dx e^{-ip\frac{2\pi}{a}x} \hat{H} e^{iq\frac{2\pi}{a}x}.$$

To calculate $\hat{H}e^{iq\frac{2\pi}{a}x}$, the periodic potential V(x) can be expanded in Fourier series as $V(x) \to V_q$, where

$$V(x) = \sum_{q'=-N}^{N} V_{q'} e^{iq'\frac{2\pi}{a}x}.$$

The basis wave function can then be written as:

$$\hat{H}e^{iq\frac{2\pi}{a}x} = (\hat{T} + \hat{V})e^{iq\frac{2\pi}{a}x} = \frac{2\hbar^2 q^2 \pi^2}{m_e a^2}e^{iq\frac{2\pi}{a}x} + \sum_{q'=-N}^{N} V_{q'}e^{i(q'+q)\frac{2\pi}{a}x}$$

Try constructing the Hamiltonian matrix H_{pq} and solve the eigenvalue equation $\hat{H}\psi_i = E\psi_i$ to obtain the three lowest energy eigenvalues.

Special note: You can use built-in functions to simplify the eigenvalue calculations and FFT transformations.

1.2 程序描述

这道题的过程和结果都太神奇了。

我们先顺着题目未完成的思路进行下来, 计算哈密顿矩阵元

$$H_{pq} = \frac{1}{a} \left\langle e^{ip\frac{2\pi}{a}x} \left| \hat{T} \right| e^{iq\frac{2\pi}{a}x} \right\rangle + \frac{1}{a} \left\langle e^{ip\frac{2\pi}{a}x} \left| \hat{V} \right| e^{iq\frac{2\pi}{a}x} \right\rangle$$

其中动能项是对角阵

$$\frac{1}{a} \left\langle e^{ip\frac{2\pi}{a}x} \left| \hat{T} \right| e^{iq\frac{2\pi}{a}x} \right\rangle = \frac{\hbar^2}{2m_e} \left(\frac{2\pi q}{a} \right)^2 \cdot \frac{1}{a} \int_0^a e^{-ip\frac{2\pi}{a}x} e^{iq\frac{2\pi}{a}x} dx = \frac{\hbar^2}{2m_e} \left(\frac{2\pi q}{a} \right)^2 \delta_{pq} = \frac{2\hbar^2 \pi^2 q^2}{m_e a^2} \delta_{pq}$$

利用

$$\left\langle e^{ip\frac{2\pi}{a}x} \left| e^{iq'\frac{2\pi}{a}x} \right| e^{iq'\frac{2\pi}{a}x} \right\rangle = \int_0^a e^{-ip\frac{2\pi}{a}x} e^{iq'\frac{2\pi}{a}x} e^{iq'\frac{2\pi}{a}x} dx = \int_0^a e^{i(q'+q-p)\frac{2\pi}{a}x} dx = a\delta_{p-q,q'}$$

可以发现,势能矩阵元很自然地承载势能的傅里叶系数

$$\frac{1}{a} \left\langle e^{ip\frac{2\pi}{a}x} \left| V(x) \right| e^{iq\frac{2\pi}{a}x} \right\rangle = \sum_{q'=-N}^{N} V_{q'} \cdot \frac{1}{a} \left\langle e^{ip\frac{2\pi}{a}x} \left| e^{iq'\frac{2\pi}{a}x} \right| e^{iq\frac{2\pi}{a}x} \right\rangle = \sum_{q'=-N}^{N} V_{q'} \delta_{p-q,q'} = V_{p-q} \delta_{p-q,q'} = V_{p-q} \delta_{p-q} \delta_{p-q,q'} = V_{p-q} \delta_{p-q,q'} = V_{p-q} \delta_{p-q,q'} \delta_{p-q,q'} = V_{p-q} \delta_{p-q,q'} \delta_{p-q,q'} = V_{p-q} \delta_{p-q,q'} \delta_{p-q,q'} \delta_{p-q,q'} \delta_{p-q,q'} = V_{p-q} \delta_{p-q,q'} \delta_{p-q,q'}$$

注意到我们对实数域的势能展开时 $V_{-q'}=V_{-q'}^*$,所以最后的哈密顿矩阵是厄米阵。

不过这里有个小漏洞,原本 p-q 的取值应该在 [-2N,2N],但我们在傅里叶展开势能为 $V_{q'}$ 的时候只展开到了 $q'\in [-N,N]$,所以按题意需要将 p-q 截断在 $p-q\in [-N,N]$ 之间(在赋值的时候加一组布尔掩码),这样才能保证 $\sum_{q'=-N}^{N}V_{q'}\delta_{p-q,q'}=V_{p-q}$.

我原本认为这样是不合理的,应该将势能展开到 $q' \in [-2N,2N]$,而不是选择牺牲势能项的精度。但在助教提示下发现,本题的势能项本身影响就不大,所以这样的截断是可以接受的。在主程序中我也对这一点进行了对比论证,详见结果示例。

在最初阅读题干时,我最不明白的是,如此简单的阶梯函数已经可以解析计算出势能的傅里叶展开

$$V_{q'} = \begin{cases} \frac{U_0 L_B}{a}, & q = 0\\ \frac{U_0}{iq'2\pi} \left(1 - e^{-iq'2\pi \frac{L_W}{a}}\right), & q \neq 0 \end{cases}$$

为何还需要进行 FFT 变换?随后想到这节课毕竟是在学习 FFT,故在主程序中加入了 FFT 变换(先将实空间离散 化再传递给 scipy.fft 进行变换)的系数与解析计算的系数进行对比,发现两者差异在容许范围内,后面的计算则 是基于解析计算的系数进行的。

源代码在 $src/KP_well.py$ 中,请在该目录下运行 $python -u \ KP_well.py$ 即可得到结果,需要安装 numpy 与 scipy 库辅助运算,matplotlib 库绘图。结果非常神奇,在大规模的绘制与计算中,我们发现了疑似二重简并的现象。在与 QM1 荣誉课的助教老师讨论后,合理怀疑是有限 N 的截断带来的边界效应,且在高能级时非局域化更明显,两侧的对称传播似乎是合理的。可惜在 debug 时我尝试过绘制其它相邻能级的波函数概率分布,未发现明显对称性。二重简并的原因有待进一步探讨,但能级的 n^2 增长是可以确认的。

主程序流程描述

- 1. **初始化物理常量与参数**: 从 const 获取需要使用的物理常数 (如 \hbar 、 m_e 等), 定义参数, 如 U_0 、 L_w 和 L_b 。
- 2. **傅里叶系数计算对比**: 均采用自然顺序 [0,1,2,...,N,-N,...,-1], 便于切片负数索引
- 使用 compute_Vq_analytical 方法计算势能傅里叶系数,基于解析表达式,
- 使用 compute_Vq_fft 方法,通过 FFT 计算傅里叶系数,验证解析方法的精确性。
- 比较两种方法的结果,输出最大和平均差异。
 - 3. 小规模收敛性检查: 借助 check convergence 函数, 在较小 N 下进行能级收敛性测试。
- 4. AB 方法对比:
- **方法 A**: 在傅里叶系数范围 $Vq' \in [-N, N]$ 下进行求解,得到小规模能级,
- 方法 B: 在扩展傅里叶系数范围 $Vq' \in [-2N, 2N]$ 下进行求解,
- 比较两种方法的能级差异,在之后计算中遵循题意采用方法 A.
- 5. 中等规模 N = 100 计算与绘制:
- 使用解析方法计算的傅里叶系数,构建哈密顿量矩阵,通过 solve_eigenvalues 方法得到能级和波函数,
- 使用 reconstruct_wavefunction_vectorized 重构波函数, 充分考虑基底正交性, 归一化操作不需要积分,
- 通过 plot_wavefunctions_and_potential 绘制势阱和前三个波函数,
- 绘制能级图,发现在小规模 N 中未出现的二级简并现象。
- 6. 大规模计算与拟合(可选): 在用户确认的情况下
- 使用更大 N = 500 计算较多能级,
- 通过 curve_fit 函数拟合能级间距与量子数的平方增长关系,
- 输出拟合系数并绘制拟合结果。

1.3 伪代码

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```
Algorithm 1: Main Process Flow
   Input: U_0, L_w, L_b, N_{initial}, num\_levels_{initial}, medium\_N, medium\_num\_levels, large\_N,
          large num levels
   Output: Energy levels and wave functions, potential well plots, convergence analysis
1 Step 1: Fourier Coefficients Comparison;
2 Vq\_analytical, Vq\_fft \leftarrow \text{compare\_fourier\_methods}(N_{initial}, U_0, L_w, L_b, a);
3 if compare_fourier_methods accuracy is high then
      Proceed to next step;
5 end
6 Step 2: Convergence Check;
7 check_convergence(U_0, L_w, L_b, a, hbar, m e);
8 Step 3: Energy Levels Calculation with Method A and B;
9 for Method in [A, B] do
      eigenvalues, eigenvectors \leftarrow \texttt{build\_and\_solve}(Method\ parameters);
10
      if degeneracies found then
11
         Print "Degeneracies exist";
12
      end
13
      else
14
         Print "No degeneracies found";
15
      end
16
17 end
18 Step 4: Plotting for Medium N;
19 eigenvalues\_medium, eigenvectors\_medium \leftarrow build\_and\_solve(medium\_N, medium\_num\_levels);
20 x medium, wavefunctions medium \leftarrow \texttt{reconstruct\_wavefunction\_vectorized}(eigenvectors medium);
21 V_x_medium \leftarrow generate_potential_vectorized(x_medium);
22 plot_wavefunctions_and_potential(x\_medium, wavefunctions\_medium, V\_x\_medium);
23 plot_energy_levels(eigenvalues_medium, medium_N);
24 if Proceed to large scale then
      Repeat plotting with large N and fit quadratic relation;
25
26 end
```

```
Algorithm 2: Compute Analytical Fourier Coefficients
```

```
Input: N, U_0, L_w, L_b, a
Output: Fourier coefficients Vq

1 n\_list \leftarrow \texttt{get\_n\_list}(N);
2 for Each \ non\text{-}zero \ n \ in \ n\_list \ do
3 | Compute Vq[n] based on analytical formula;
4 end
5 Vq[0] \leftarrow U_0 \cdot L_b/a;
6 return Vq
```

Algorithm 3: Build Kinetic Matrix

```
Input: q\_values, a, hbar, m\_e

Output: Kinetic matrix T

1 for Each\ q in q\_values do

2 \left|\begin{array}{c} T[q,q] \leftarrow \frac{\hbar^2 \cdot (2\pi q/a)^2}{2m_e}; \\ \end{array}\right|
3 end

4 return T as a sparse diagonal matrix
```

Algorithm 4: Solve for Eigenvalues with Error Handling

Input: Hamiltonian matrix H, desired levels num_levels

Output: Eigenvalues and eigenvectors, or an error message if solving fails

```
1 if eigsh can solve for eigenvalues then
```

```
    2 (eigenvalues, eigenvectors) ← eigsh(H, k = num_levels, which = "SA");
    3 Sort eigenvalues and eigenvectors by eigenvalue magnitude;
    4 return eigenvalues, eigenvectors;
```

5 end

6 else

Print "Error in eigenvalue solver";
return None, None;

9 end

6 end

Algorithm 5: Reconstruct Wavefunction

```
Input: Eigenvectors, q_values, a
```

Output: Wavefunctions over extended domain

```
1 for Each eigenvector v do
```

```
for Each position x in domain do
psi(x) \leftarrow v \cdot \exp(2\pi i \cdot q \cdot x/a);
Normalize psi over extended domain;
end
```

1.4 结果示例

第一步:比较解析方法和FFT方法计算傅里叶系数的差异

傅里叶系数计算方法比较:

傅里叶级数范围: [-50, 50]

解析方法用时: 0.00e+00秒

FFT方法用时: 5.12e-04秒

最大绝对差异: 2.058607e-02 eV

平均绝对差异: 1.223039e-02 eV

解析方法和FFT方法计算的傅里叶系数差异非常小。

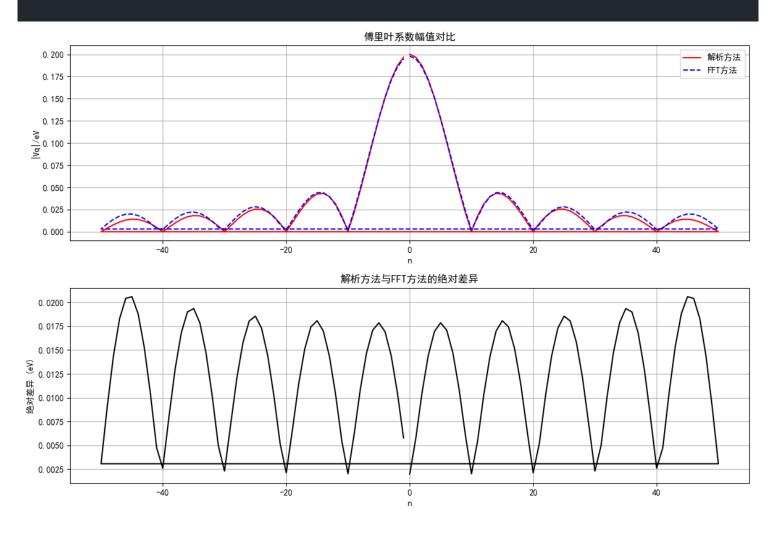


图 2: 第一步: 比较解析计算与 FFT 得到的势能傅里叶系数

```
第二步: 进行收敛性检查(仅对小规模N进行)
N=10, Lowest 3 eigenvalues (eV): [0.14464301 1.51649618 1.88295936]
No degeneracies found.
N=20, Lowest 3 eigenvalues (eV): [0.1446063 1.51647998 1.88283855]
No degeneracies found.
N=30, Lowest 3 eigenvalues (eV): [0.14460216 1.51647846 1.88282499]
No degeneracies found.
Converged at N=30
```

图 3: 第二步: 收敛性检查

```
第三步:使用解析方法计算的傅里叶系数进行AB方法对比

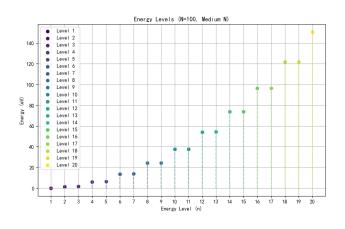
方法A(Vq' ∈[-50, 50])
Lowest 3 eigenvalues (eV):
    Level 1: 1.4460e-01 eV
    Level 2: 1.5165e+00 eV
    Level 3: 1.8828e+00 eV
No degeneracies found.

方法B(Vq' ∈[-100, 100])
Lowest 3 eigenvalues (eV):
    Level 1: 1.4460e-01 eV
    Level 2: 1.5165e+00 eV
    Level 3: 1.8828e+00 eV
No degeneracies found.

方法A和方法B的前3个能级相对差异: [2.39109506e-06 7.59823372e-08 6.01067236e-07]
方法A和方法B的能级差异非常小。
```

图 4: 第三步: 不同势能展开范围的 AB 方法对比

```
第四步: 使用中等规模N=100进行能级绘制
N=100, Lowest 20 eigenvalues (eV):
  Level 1: 1.4460e-01 eV
  Level 2: 1.5165e+00 eV
  Level 3: 1.8828e+00 eV
  Level 4: 6.0633e+00 eV
  Level 5: 6.3848e+00 eV
  Level 6: 1.3633e+01 eV
  Level 7: 1.3855e+01 eV
  Level 8: 2.4216e+01 eV
  Level 9: 2.4325e+01 eV
  Level 10: 3.7801e+01 eV
  Level 11: 3.7810e+01 eV
  Level 12: 5.4321e+01 eV
  Level 13: 5.4379e+01 eV
  Level 14: 7.3860e+01 eV
  Level 15: 7.3946e+01 eV
  Level 16: 9.6426e+01 eV
  Level 17: 9.6503e+01 eV
  Level 18: 1.2201e+02 eV
  Level 19: 1.2206e+02 eV
  Level 20: 1.5061e+02 eV
Degeneracies found.
```



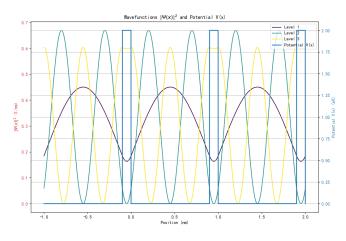
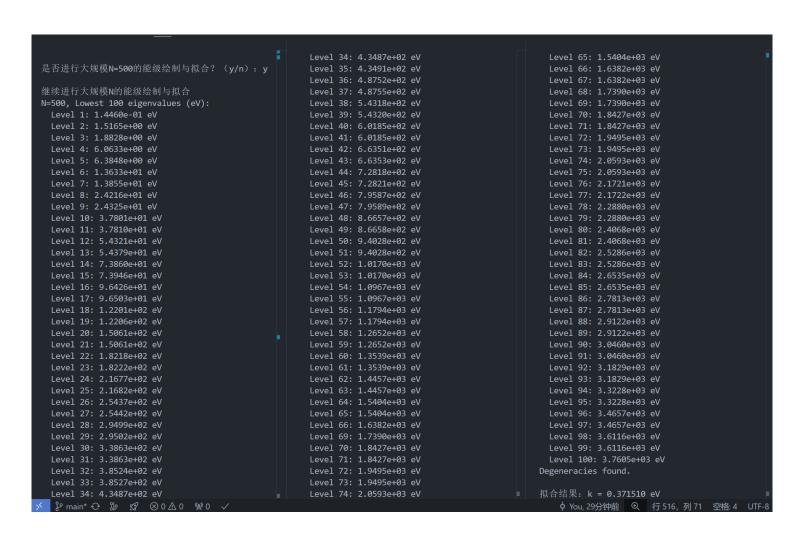


图 5: 第四步: 进行中等规模 $N=100, num_levels=20$ 的计算与绘制, 并绘制势阱与前三个能级的波函数



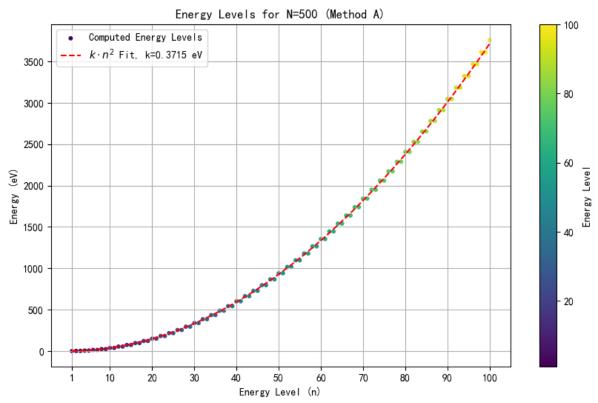


图 6: 第五步: 进行大规模 $N = 500, num_levels = 100$ 的计算与绘制,并进行拟合

2 题目 2: 太阳黑子周期性检测

2.1 题目描述

Use the file called sunspots.txt, which contains the observed number of sunspots on the Sun for each month since January 1749. Write a program to calculate the Fourier transform of the sunspot data and then make a graph of the magnitude squared $|c_k|^2$ of the Fourier coefficients as a function of k—also called the power spectrum of the sunspot signal. You should see that there is a noticeable peak in the power spectrum at a nonzero value of k. Find the approximate value of k to which the peak corresponds. What is the period of the sine wave with this value of k? Special note: You may use any built-in functions for the Fourier transform.

2.2 程序描述

2.3 伪代码

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2.4 结果示例