Appendix

Following [georges1991expand] lets define energy as:

$$E = -\sum_{ij} w_{ij} s_i s_j - \sum_i \theta_i s_i \tag{1}$$

and introduce the following operator:

$$U \equiv E - \mathbb{E}(E) - \sum_{i} \frac{\partial \lambda_{i}(\beta)}{\partial \beta} (s_{i} - m_{i})$$
(2)

which poses useful property – $\mathbb{E}(U) = 0$. For any other operator O we then have:

$$\frac{\partial \mathbb{E}(O)}{\partial \beta} = \mathbb{E}\left(\frac{\partial O}{\partial \beta}\right) - \mathbb{E}(OU). \tag{3}$$

Now, the first derivative from the Taylor expansion is:

$$\frac{\partial(\beta F)}{\partial\beta} = \frac{\sum_{\mathbf{s}} \exp\left(\beta \sum_{(ij)} w_{ij} s_i s_j + \beta \sum_i \theta_i s_i + \sum_i \lambda_i(\beta)(s_i - m_i)\right) \left(\sum_{(ij)} w_{ij} s_i s_j + \sum_i \theta_i s_i + \sum_i \frac{\partial \lambda_i(\beta)}{\partial\beta}(s_i - m_i)\right)}{\sum_{\mathbf{s}} \exp\left(\beta \sum_{(ij)} w_{ij} s_i s_j + \sum_i \theta_i s_i + \sum_i \lambda_i(\beta)(s_i - m_i)\right)}$$

$$= \sum_{(i,i)} w_{ij} \mathbb{E}(s_i s_j) + \sum_i a_i \mathbb{E}(s_i) + \frac{\partial \lambda_i(\beta)}{\partial\beta} \sum_i \mathbb{E}(s_i - m_i).$$

In the case of $\beta = 0$ we have:

$$\left. \frac{\partial (\beta F)}{\partial \beta} \right|_{\beta=0} = \sum_{(ij)} w_{ij} m_i m_j + \sum_i \theta_i m_i.$$

Using 3 we obtain:

$$\frac{\partial m_i}{\partial \beta} = 0 = \frac{\partial \mathbb{E}(s_i)}{\partial \beta} = \mathbb{E}\left(\frac{\partial s_i}{\partial \beta}\right) - \mathbb{E}(s_i U) = -\mathbb{E}(s_i U) = -\mathbb{E}(U(s_i - m_i)). \tag{4}$$

The first derivative of the operator U has the form:

$$\frac{\partial U}{\partial \beta} = \frac{\partial E}{\partial \beta} - \frac{\partial \mathbb{E}(E)}{\partial \beta} - \sum_{i} \frac{\partial^{2} \lambda_{i}(\beta)}{\partial \beta^{2}} (s_{i} - m_{i})$$

$$= \mathbb{E}(U^{2}) - \sum_{i} \frac{\partial^{2} \lambda_{i}(\beta)}{\partial \beta^{2}} (s_{i} - m_{i})$$
(5)

and the second derivative is:

$$\frac{\partial^2 U}{\partial \beta^2} = 2\mathbb{E}\left(\frac{\partial U}{\partial \beta}U\right) - \mathbb{E}(U^3) - \sum_i \frac{\partial^3 \lambda_i(\beta)}{\partial \beta^3}(s_i - m_i)$$

$$= -\mathbb{E}(U^3) - \sum_i \frac{\partial^3 \lambda_i(\beta)}{\partial \beta^3}(s_i - m_i)$$
(6)

The expansion of free energy using formulas derived above can now be reformulated in terms of the operator U:

$$\frac{\partial(\beta F)}{\partial \beta} = \mathbb{E}(E) - \sum_{i} \frac{\partial \lambda_{i}(\beta)}{\partial \beta} \mathbb{E}(s_{i} - m_{i}) = \mathbb{E}(E)$$
 (7)

and the higher orders:

$$\frac{\partial^{2}(\beta F)}{\partial \beta^{2}} = \mathbb{E}\left(\frac{\partial E}{\partial \beta}\right) - \mathbb{E}(EU) = -\mathbb{E}(U^{2}),
\frac{\partial^{3}(\beta F)}{\partial \beta^{3}} = -2\mathbb{E}\left(U\frac{\partial U}{\partial \beta}\right) + \mathbb{E}(U^{3}) = \mathbb{E}(U^{3}).$$
(8)

Taylor expansion was considered around point $\beta = 0$. Using derivations from 7 we obtain again a 'naive' term:

$$\frac{\partial(\beta F)}{\partial\beta}\Big|_{\beta=0} = \mathbb{E}_{\beta=0}(E) = -\sum_{(ij)} w_{ij} m_i m_j - \sum_i \theta_i m_i = -\frac{1}{2} \sum_i \sum_j w_{ij} m_i m_j - \sum_i \theta_i m_i. \tag{9}$$

Consider now:

$$\left. \frac{\partial(\beta F)}{\partial m_i \partial \beta} \right|_{\beta=0} = -\sum_{j \neq i} w_{ij} m_j - \theta_i. \tag{10}$$

On the other hand:

$$\frac{\partial(\beta F)}{\partial m_i \partial \beta}\bigg|_{\beta=0} = \frac{\partial(\beta F)}{\partial \beta \partial m_i}\bigg|_{\beta=0} = \frac{\partial}{\partial \beta} \mathbb{E}(\lambda_i(\beta)) = \frac{\partial \lambda_i(\beta)}{\partial \beta}\bigg|_{\beta=0}$$
(11)

Substituting 11 into 2 gives us:

$$U_{\beta=0} = -\sum_{(ij)} w_{ij} s_i s_j - \sum_i \theta_i s_i + \frac{1}{2} \sum_i \sum_j w_{ij} m_i m_j + \sum_i \theta_i m_i + \sum_i \left(\sum_{j \neq i} w_{ij} m_j + \theta_i \right) (s_i - m_i)$$

$$= -\sum_{(ij)} w_{ij} s_i s_j - \frac{1}{2} \sum_{(ij)} w_{ij} m_i m_j + \sum_i \sum_{j \neq i} w_{ij} s_i m_j$$

$$= -\sum_{(ij)} w_{ij} (s_i - m_i) (s_j - m_j) = -\sum_l w_l y_l$$
(12)

where $w_l = w_{ij}$ and $y_l = (s_i - m_i)(s_j - m_j)$ stands for the 'link' operator which poses useful properties:

$$\mathbb{E}(y_l)_{\beta=0} = \mathbb{E}(s_i s_j) - m_j \mathbb{E}(s_i) - m_i \mathbb{E}(s_j) + m_i m_j = 0$$

$$\mathbb{E}(y_l(s_i - m_i))_{\beta=0} = m_j - m_j - m_i^2 m_j + m_i^2 m_j - m_i^2 m_j + m_i^2 m_j - m_i^2 m_j = 0.$$
(13)

Finally, if $k \neq l$ then:

$$\mathbb{E}(y_k y_l) = \mathbb{E}(y_k) \mathbb{E}(y_l) = 0$$

while for k = l we have:

$$\mathbb{E}((s_i - m_1)^2 (s_j - m_j)^2) = m_i m_j - 2m_i m_j^2 + m_i m_j^2 - 2m_i^2 m_j + 4m_1^2 m_j^2 - 2m_i^2 m_j^2 + m_i^2 m_j - 2m_i^2 m_j^2 + m_i^2 m_j^2 = (m_i - m_i^2)(m_j - m_j^2).$$
(14)

Using properties from y_l in equations 8 we can derive:

$$\begin{split} \frac{\partial^{2}(\beta F)}{\partial \beta^{2}} \bigg|_{\beta=0} &= -\mathbb{E}(U^{2})_{\beta=0} \\ &= -\sum_{l_{i}l_{2}} w_{l_{i}} w_{l_{2}} \mathbb{E}_{\beta=0}(y_{l_{1}} y_{l_{2}}) \\ &= -\sum_{(i,j)} w_{ij}^{2} (m_{i} - m_{i}^{2}) (m_{j} - m_{j}^{2}) \end{split}$$

which yields the TAP-Onsager term. To obtain the next term in the Taylor expansion we need to compute $\mathbb{E}(y_{l_1}y_{l_2}y_{l_3})$ term and by definition the structure of the RBM model doesn't admit triangles in its corresponding factor graphs. Thus, we need to consider only the case when $l_1 = l_2 = l_3$:

$$\mathbb{E}((s_i - m_1)^3 (s_j - m_j)^3) = m_i m_j - 3m_i m_j^2 + 2m_i m_j^2 + 2m_i m_j^3 - 3m_i^2 m_j + 2m_i^3 m_j + 9m_i^2 m_j^2 - 6m_i^3 m_j^2 - 6m_i^2 m_j^3 + 4m_i^3 m_j^3 = 4(m_i - m_i^2)(\frac{1}{2} - m_i)(m_j - m_j^2)(\frac{1}{2} - m_j)$$
(15)

and the third-order term in the case of the RBM structure:

$$\frac{\partial^3(\beta F)}{\partial \beta^3}\bigg|_{\beta=0} = \frac{2\beta^3}{3} \sum_{(ij)} w_{ij}^3(m_i - m_i^2)(\frac{1}{2} - m_i)(m_j - m_j^2)(\frac{1}{2} - m_j).$$