

1. Intro

Define variational free energy with a set of external auxiliary fields (Lagrange multipliers):

$$-\beta G = \ln \sum_{\mathbf{s}} \exp \left(\beta \sum_{(ij)} w_{ij} s_i s_j + \beta \sum_i a_i s_i + \sum_i \lambda_i(\beta) (s_i - m_i) \right)$$

We are interested in expanding $-\beta G(\beta, \mathbf{m})$ around $\beta = 0$ where the spins are entirely controlled by their auxiliary fields:

$$-\beta G = -(\beta G)_{\beta=0} - \left(\frac{\partial(\beta G)}{\partial \beta} \right)_{\beta=0} \beta - \left(\frac{\partial^2(\beta G)}{\partial \beta^2} \right)_{\beta=0} \frac{\beta^2}{2} - \dots$$

At $\beta = 0$ the spins are uncorrelated (where \mathbb{E} refers to the average configuration under the Boltzmann measure):

$$m_i = \mathbb{E}_{\beta=0}(s_i) = \frac{\exp(\lambda_i(0))}{\exp(\lambda_i(0)) + 1} = \text{sigmoid}(\lambda_i(0)) \quad (1)$$

and

$$\begin{aligned} -(\beta G)_{\beta=0} &= \ln \sum_{\mathbf{s}} \exp \left(\sum_i \lambda_i(0) (s_i - m_i) \right) \\ &= \ln \left\{ \sum_{s_1} \exp(\lambda_1(0)(s_1 - m_1)) \dots \sum_{s_n} \exp(\lambda_n(0)(s_n - m_n)) \right\} \\ &= \ln \{ (\exp(\lambda_i(0)) + 1) \exp(-\lambda_i(0)m_1) \dots (\exp(\lambda_i(0)) + 1) \exp(-\lambda_i(0)m_n) \} \\ &= \sum_i \left(\ln \left[1 + \left(\frac{m_i}{1 - m_i} \right) \right] - \lambda_i(0)m_i \right) \end{aligned} \quad (2)$$

Using 1, we obtain:

$$\lambda_i(0) = \text{logit}(m_i) = \ln \left(\frac{m_i}{1 - m_i} \right)$$

and thus:

$$\begin{aligned} -(\beta G)_{\beta=0} &= \sum_i \left\{ \ln \left(\frac{1}{1 - m_i} \right) - m_i \ln \left(\frac{m_i}{1 - m_i} \right) \right\} \\ &= - \sum_i [m_i \ln(m_i) + (1 - m_i) \ln(1 - m_i)] \end{aligned}$$

Yedida and Georges showed how to create the Taylor expansion beyond $O(\beta^2)$ (TODO referring to MF and TAP variant) (derivation in Appendix):

$$\begin{aligned} -\beta G &= - \sum_i [m_i \ln(m_i) + (1 - m_i) \ln(1 - m_i)] \\ &\quad + \beta \sum_{(ij)} w_{ij} m_i m_j + \beta \sum_i a_i m_i \\ &\quad + \frac{\beta^2}{2} \sum_{(ij)} w_{ij}^2 (m_i - m_i^2)(m_j - m_j^2) \\ &\quad + \frac{2\beta^3}{3} \sum_{(ij)} w_{ij}^3 (m_i - m_i^2) \left(\frac{1}{2} - m_i \right) (m_j - m_j^2) \left(\frac{1}{2} - m_j \right) \end{aligned}$$

Appendix

Lets define energy as:

$$E = - \sum_{ij} w_{ij} s_i s_j - \sum_i a_i s_i \quad (3)$$

and introduce the following operator:

$$U \equiv E - \mathbb{E}(E) - \sum_i \frac{\partial \lambda_i(\beta)}{\partial \beta} (s_i - m_i) \quad (4)$$

which poses useful property:

$$\mathbb{E}(U) = 0$$

For any other operator O we then have:

$$\frac{\partial \mathbb{E}(O)}{\partial \beta} = \mathbb{E} \left(\frac{\partial O}{\partial \beta} \right) - \mathbb{E}(OU) \quad (5)$$

Now, the first derivative from the Taylor expansion is:

$$\begin{aligned} \frac{\partial(\beta G)}{\partial \beta} &= \frac{\sum_{\mathbf{s}} \exp \left(\beta \sum_{(ij)} w_{ij} s_i s_j + \beta \sum_i a_i s_i + \sum_i \lambda_i(\beta) (s_i - m_i) \right) \left(\sum_{(ij)} w_{ij} s_i s_j + \sum_i a_i s_i + \sum_i \frac{\partial \lambda_i(\beta)}{\partial \beta} (s_i - m_i) \right)}{\sum_{\mathbf{s}} \exp \left(\beta \sum_{(ij)} w_{ij} s_i s_j + \beta \sum_i a_i s_i + \sum_i \lambda_i(\beta) (s_i - m_i) \right)} \\ &= \sum_{(ij)} w_{ij} \mathbb{E}(s_i s_j) + \sum_i a_i \mathbb{E}(s_i) + \frac{\partial \lambda_i(\beta)}{\partial \beta} \sum_i \mathbb{E}(s_i - m_i) \end{aligned}$$

In the case of $\beta = 0$ we have:

$$\left. \frac{\partial(\beta G)}{\partial \beta} \right|_{\beta=0} = \sum_{(ij)} w_{ij} m_i m_j + \sum_i a_i m_i$$

Using 5 we obtain:

$$\frac{\partial m_i}{\partial \beta} = 0 = \frac{\partial \mathbb{E}(s_i)}{\partial \beta} = \mathbb{E} \left(\frac{\partial s_i}{\partial \beta} \right) - \mathbb{E}(s_i U) = -\mathbb{E}(s_i U) = -\mathbb{E}(U(s_i - m_i)) \quad (6)$$

Thus:

$$\begin{aligned} \frac{\partial U}{\partial \beta} &= \frac{\partial H}{\partial \beta} - \frac{\partial \mathbb{E}(H)}{\partial \beta} - \sum_i \frac{\partial^2 \lambda_i(\beta)}{\partial \beta^2} (s_i - m_i) \\ &= \mathbb{E}(U^2) - \sum_i \frac{\partial^2 \lambda_i(\beta)}{\partial \beta^2} (s_i - m_i) \end{aligned} \quad (7)$$

The second derivative has the form:

$$\begin{aligned} \frac{\partial^2 U}{\partial \beta^2} &= 2\mathbb{E} \left(\frac{\partial U}{\partial \beta} U \right) - \mathbb{E}(U^3) - \sum_i \frac{\partial^3 \lambda_i(\beta)}{\partial \beta^3} (s_i - m_i) \\ &= -\mathbb{E}(U^3) - \sum_i \frac{\partial^3 \lambda_i(\beta)}{\partial \beta^3} (s_i - m_i) \end{aligned} \quad (8)$$

Coming back to our expansion of free energy using formulas derived above we now can calculate:

$$\frac{\partial(\beta G)}{\partial \beta} = \mathbb{E}(E) - \sum_i \frac{\partial \lambda_i(\beta)}{\partial \beta} \mathbb{E}(s_i - m_i) = \mathbb{E}(E) \quad (9)$$

and:

$$\frac{\partial^2(\beta G)}{\partial \beta^2} = \mathbb{E} \left(\frac{\partial E}{\partial \beta} \right) - \mathbb{E}(EU) = -\mathbb{E}(U^2)$$

$$\frac{\partial^3(\beta G)}{\partial \beta^3} = -2\mathbb{E} \left(U \frac{\partial U}{\partial \beta} \right) + \mathbb{E}(U^3) = \mathbb{E}(U^3) \quad (10)$$

$$\frac{\partial^4(\beta G)}{\partial \beta^4} = 3\mathbb{E} \left(U^2 \frac{\partial U}{\partial \beta} \right) - \mathbb{E}(U^4) = 3 \left(\mathbb{E}(U^2) \right)^2 - 3 \sum_i \frac{\partial^2 \lambda_i(\beta)}{\partial \beta^2} \mathbb{E}(U^2(s_i - m_i)) - \mathbb{E}(U^4)$$

TS was considered around point $\beta = 0$. Using derivations from above we obtain again a 'naive' term:

$$\left. \frac{\partial(\beta G)}{\partial \beta} \right|_{\beta=0} = \mathbb{E}_{\beta=0}(E) = - \sum_{(ij)} w_{ij} m_i m_j - \sum_i a_i m_i = -\frac{1}{2} \sum_i \sum_j w_{ij} m_i m_j - \sum_i a_i m_i \quad (11)$$

Consider now

$$\left. \frac{\partial(\beta G)}{\partial m_i \partial \beta} \right|_{\beta=0} = - \sum_{j \neq i} w_{ij} m_j - a_i \quad (12)$$

On the other hand:

$$\left. \frac{\partial(\beta G)}{\partial m_i \partial \beta} \right|_{\beta=0} = \left. \frac{\partial(\beta G)}{\partial \beta \partial m_i} \right|_{\beta=0} = \frac{\partial}{\partial \beta} \mathbb{E}(\lambda_i(\beta)) = \left. \frac{\partial \lambda_i(\beta)}{\partial \beta} \right|_{\beta=0} \quad (13)$$

Substituting 13 into 4 gives us:

$$\begin{aligned} U_{\beta=0} &= - \sum_{(ij)} w_{ij} s_i s_j - \sum_i a_i s_i + \frac{1}{2} \sum_i \sum_j w_{ij} m_i m_j + \sum_i a_i m_i + \sum_i \left(\sum_{j \neq i} w_{ij} m_j + a_i \right) (s_i - m_i) \\ &= - \sum_{(ij)} w_{ij} s_i s_j - \frac{1}{2} \sum_{(ij)} w_{ij} m_i m_j + \sum_i \sum_{j \neq i} w_{ij} s_i m_j \\ &= - \sum_{(ij)} w_{ij} (s_i - m_i)(s_j - m_j) = - \sum_l w_l y_l \end{aligned} \quad (14)$$

where y_l stands for the 'link' operator $w_l = w_{ij}$ and $y_l = (s_i - m_i)(s_j - m_j)$ which poses useful properties:

$$\begin{aligned} \mathbb{E}(y_l)_{\beta=0} &= \mathbb{E}(s_i s_j) - m_j \mathbb{E}(s_i) - m_i \mathbb{E}(s_j) + m_i m_j = 0 \\ \mathbb{E}(y_l(s_i - m_i))_{\beta=0} &= m_j - m_j - m_i^2 m_j + m_i^2 m_j \\ &\quad - m_i^2 m_j + m_i^2 m_j + m_i^2 m_j - m_i^2 m_j \\ &= 0 \end{aligned} \quad (15)$$

Finally, if $k \neq l$ then:

$$\mathbb{E}(y_k y_l) = \mathbb{E}(y_k) \mathbb{E}(y_l) = 0$$

while for $k = l$ we have:

$$\begin{aligned} \mathbb{E}((s_i - m_i)^2 (s_j - m_j)^2) &= m_i m_j - 2m_i m_j^2 + m_i m_j^2 - 2m_i^2 m_j + 4m_i^2 m_j^2 \\ &\quad - 2m_i^2 m_j^2 + m_i^2 m_j - 2m_i^2 m_j^2 + m_i^2 m_j^2 \\ &= (m_i - m_i^2)(m_j - m_j^2) \end{aligned} \quad (16)$$

Using properties from 16 in equations 10 we can obtain:

$$\begin{aligned}
\left. \frac{\partial^2(\beta G)}{\partial \beta^2} \right|_{\beta=0} &= -\mathbb{E}(U^2)_{\beta=0} \\
&= -\sum_{l_1 l_2} w_{l_1} w_{l_2} \mathbb{E}_{\beta=0}(y_{l_1} y_{l_2}) \\
&= -\sum_{(i,j)} w_{ij}^2 (m_i - m_i^2)(m_j - m_j^2)
\end{aligned}$$

which yields the TAP-Onsager term.

To obtain the next term for the Taylor expansion we need to compute $\mathbb{E}(y_{l_1} y_{l_2} y_{l_3})$ term and by definition the structure of the RBM model doesn't admit triangles in its corresponding factor graphs. Thus, we need to consider only the case when $l_1 = l_2 = l_3$:

$$\begin{aligned}
\mathbb{E}((s_i - m_i)^3 (s_j - m_j)^3) &= m_i m_j - 3m_i m_j^2 + 2m_i m_j^2 + 2m_i m_j^3 - 3m_i^2 m_j + 2m_i^3 m_j \\
&\quad + 9m_i^2 m_j^2 - 6m_i^3 m_j^2 - 6m_i^2 m_j^3 + 4m_i^3 m_j^3 \\
&= 4(m_i - m_i^2) \left(\frac{1}{2} - m_i\right) (m_j - m_j^2) \left(\frac{1}{2} - m_j\right),
\end{aligned} \tag{17}$$

which gives the third term:

$$\left. \frac{\partial^3(\beta G)}{\partial \beta^3} \right|_{\beta=0} = \frac{2\beta^3}{3} \sum_{(ij)} w_{ij}^3 (m_i - m_i^2) \left(\frac{1}{2} - m_i\right) (m_j - m_j^2) \left(\frac{1}{2} - m_j\right).$$