

5.3 Nonlinear models
(with 4.10 material too)
a lesson for MATH F302 Differential Equations

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for textbook: D. Zill, *A First Course in Differential Equations with Modeling Applications*, 11th ed.

examples of nonlinear 2nd-order differential equations (DEs):

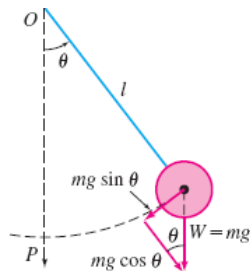
- pendulum (§5.3)
- using a numerical solver in MATLAB/OCTAVE (see §4.10)
- hard and soft springs (§5.3)
- non-constant gravity: from earth to high orbit (§5.3)
- dependent variable missing (§4.10)

nonlinear pendulum

- suppose a pendulum oscillates (swings back and forth) without resistance
- if you believe my §5.1 slides then it must be modeled by a 2nd-order linear DE
 - this is true for small oscillations
 - for bigger oscillations (more than 20°) a nonlinear model is more accurate
- the DE which comes from the diagram:

$$m\ell \frac{d^2\theta}{dt^2} = -mg \sin \theta$$

- you are not responsible for the derivation
- ... but $s = \ell\theta$ is arclength, so $\ell \frac{d^2\theta}{dt^2}$ is acceleration, and only the tangential force is relevant

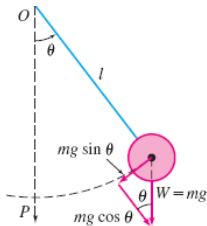


linear small angle model

- divide by $m\ell$ and move term: $\frac{d^2\theta}{dt^2} + \frac{g}{\ell} \sin \theta = 0$
- if $\omega = \sqrt{\frac{g}{\ell}}$ then $\boxed{\frac{d^2\theta}{dt^2} + \omega^2 \sin \theta = 0}$ for any angle
- recall $\sin \theta \approx \theta$ for small θ because $\sin z = z - \frac{z^3}{3!} + \frac{z^5}{5!} - \dots$
- a *small angle model*:

$$\boxed{\frac{d^2\theta}{dt^2} + \omega^2 \theta = 0}$$

- solution: $\theta(t) = c_1 \cos(\omega t) + c_2 \sin(\omega t)$



nonlinear versus linearized pendulum

nonlinear: any angles	linearized: small angles
$\theta'' + \omega^2 \sin \theta = 0$	$\theta'' + \omega^2 \theta = 0$
solution?	$\theta(t) = c_1 \cos(\omega t) + c_2 \sin(\omega t)$

- $\omega = \sqrt{g/\ell}$ in both DEs
- we don't know how to solve a nonlinear DE like a **pendulum**
 - the term " **$\sin \theta$** " is not linear: $\sin(a + b) \neq \sin(a) + \sin(b)$

what to do about a nonlinear DE?

- for example, the pendulum DE:

$$\theta'' + \omega^2 \sin \theta = 0$$

- **read section 4.10!** ← *gives advice, not a method*
- what to do about a nonlinear equation like this?
 - $\theta = e^{rt}$ is not a solution for any r (real or complex)
 - using the concept of *energy* makes progress (Mini-Project 3)
... but we get a hard-to-solve 1st-order equation
 - using infinite series can make progress too (Chapter 6) ... but basically only gives approximations
- numerical approximations can be used for an IVP!
 - Euler's method? ... inefficient but would work
 - *but* the equation is second order ... how does Euler even work?
 - next: using an efficient “black box” solver in MATLAB/OCTAVE

systems of 1st-order ODEs

idea: 2nd-order ODE is equivalent to a system of 1st-order ODEs

Example. convert into a 1st-order system:

$$x'' + 5(x')^2 + \sin x \stackrel{*}{=} \sqrt{t}$$

Solution. Second derivative $x''(t)$ is merely the derivative of $x'(t)$.
So give x' a name:

$$y = x'.$$

Now rewrite $*$ using y :

$$y' + 5y^2 + \sin x = \sqrt{t}.$$

Rearrange above two equations to a system:

$$x' = y$$

$$y' = -5y^2 - \sin x + \sqrt{t}$$

Summary: Ignore the complexity of $*$. Don't solve anything, just restate the problem.

pendulum as a 1st-order system

exercise. convert into a 1st-order system with initial conditions:

$$\theta'' + \omega^2 \sin \theta = 0, \quad \theta(0) = A, \quad \theta'(0) = B$$

solution.

$\begin{aligned} z_1' &= z_2 \\ z_2' &= -\omega^2 \sin(z_1) \end{aligned}$	$\begin{aligned} z_1(0) &= A \\ z_2(0) &= B \end{aligned}$
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using black-box solver ode45

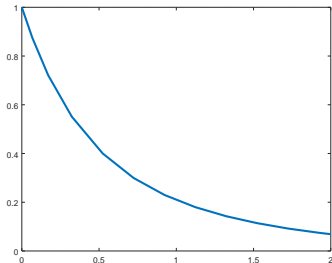
- before we get to numerical solutions of systems, let's do a single 1st-order IVP
- you can use [Octave Online](#) to do the following
- or use [MATLAB](#) or [OCTAVE](#) on your own computer

example. solve for $y(t)$ on $0 \leq t \leq 2$, and estimate $y(2)$:

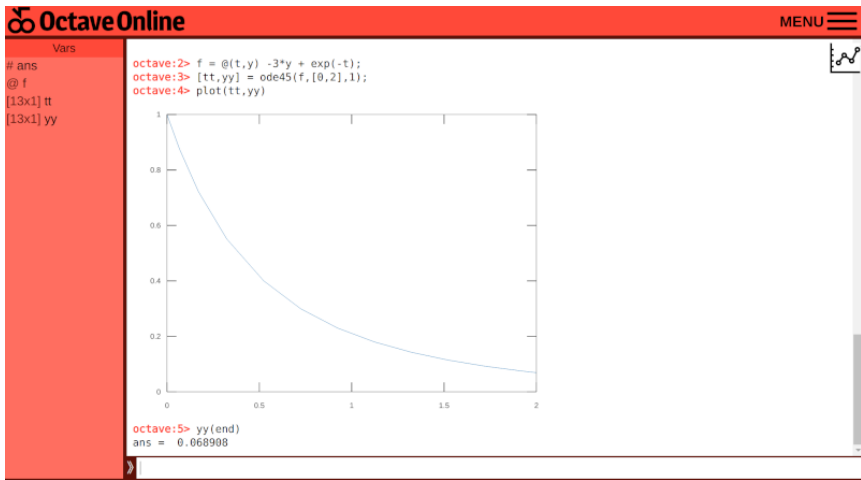
$$y' = -3y + e^{-t}, \quad y(0) = 1$$

solution. the DE is $y' = f(t, y)$ so

```
>> f = @(t,y) -3*y + exp(-t);  
>> [tt,yy] = ode45(f,[0,2],1);  
>> plot(tt,yy)  
>> yy(end)  
ans = 0.068908
```



in Octave Online



only 12 steps, but accurate

- the ode45 black-box is quite accurate
- *exercise.* solve *by hand* for the exact value $y(2)$:

$$y' = -3y + e^{-t}, \quad y(0) = 1$$

solution.

- compare to $y(\text{end})=y(13)$ on previous slides:

```
>> 0.5*(exp(-2)+exp(-6))  
ans = 0.068907
```

- Euler would need 10^5 or 10^6 steps for this accuracy

calling ode45

- from the [MATLAB documentation page on ode45](#):

```
[t,y] = ode45(odefun,tspan,y0),
```

where `tspan` = `[t0 tf]`, integrates the system of differential equations $y' = f(t,y)$ from `t0` to `tf` with initial conditions `y0`. Each row in the solution array `y` corresponds to a value returned in column vector `t`.

- see the [MATLAB page](#) for examples of functions $f(t,y)$ for the `odefun` argument
- note further fine print about the `tspan` argument:
 - If `tspan` has two elements `[t0 tf]` then the solver returns the solution evaluated at internal integration steps in the interval.
 - If `tspan` has more than two elements `[t0,t1,t2,...,tf]` then the solver returns the solution evaluated at the given points.

ode45 for pendulum

example. let $\omega = \sqrt{7}$. solve for $\theta(t)$ on the interval $t \in [0, 20]$:

$$\theta'' + \omega^2 \sin \theta = 0, \quad \theta(0) = 3, \quad \theta'(0) = 0$$

solution. $z_1 = \theta$ and $\omega^2 = 7$ so

$$z_1' = z_2$$

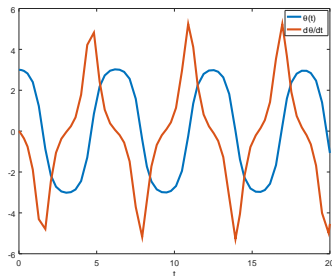
$$z_1(0) = 3$$

$$z_2' = -7 \sin(z_1)$$

$$z_2(0) = 0$$

This is $z' = f(t, z)$ so:

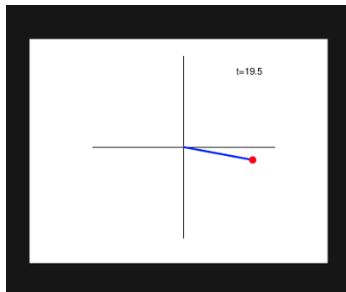
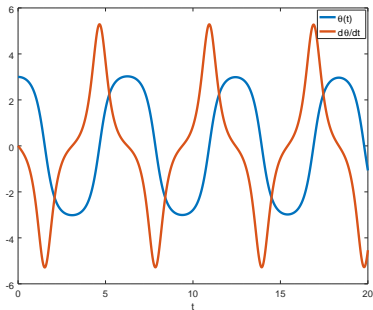
```
>> f = @(t,z) [z(2); -7*sin(z(1))];  
>> [tt,zz] = ode45(f,[0,20],[3;0]);  
>> plot(tt,zz)  
>> xlabel t  
>> legend('\theta(t)', 'd\theta/dt')
```



pendulum: better and movier

- the solution is more accurate than it looks!
- for better appearance, generate more points (left fig.):

```
>> [tt,zz] = ode45(f,[0:.01:20],[3;0]);  
>> plot(tt,zz), xlabel t
```
- one can also make a [movie](#)
 - see [pendmovie.m](#) at [Codes](#) at [bueler.github.io/math302](#)
 - right fig. is a frame



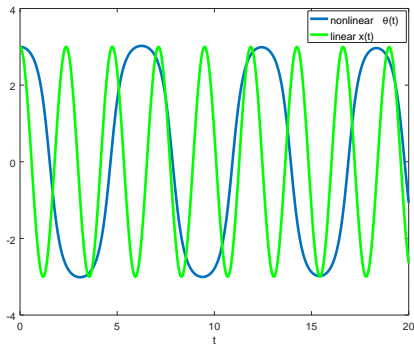
back to linear mass-spring

example. solve for $x(t)$ on the interval $t \in [0, 20]$:

$$x'' + 7x = 0, \quad x(0) = 3, \quad x'(0) = 0$$

exact solution.

$$x(t) = 3 \cos(\sqrt{7}t)$$



continuing previous code:

```
>> plot(tt,zz(:,1),'b',tt,3*cos(sqrt(7)*tt),'g')
>> xlabel t
>> legend('nonlinear \theta(t)', 'linear x(t)')
```

linear mass-spring: exact vs. numerical

- this is a good case on which to check accuracy
- example.* find $x(20)$:

$$x'' + 7x = 0, \quad x(0) = 3, \quad x'(0) = 0$$

exact solution. $x(20) = 3 \cos(\sqrt{7}(20)) = -2.6441$

numerical solution. $z_1 = x$ and $z_2 = x'$ so

$$\begin{aligned} z_1' &= z_2 & z_1(0) &= 3 \\ z_2' &= -7z_1 & z_2(0) &= 0 \end{aligned}$$

```
>> f1 = @(t,z) [z(2); -7*z(1)];  
>> [tt1,zz1] = ode45(f1,[0:.01:20],[3;0]);  
>> zz1(end,1)  
ans = -2.6492
```

- what about plots of the exact and numerical solutions?
 - you won't see difference: $x(t) = 3 \cos(\sqrt{7}t)$ versus $zz1(:,1)$

nonlinear springs

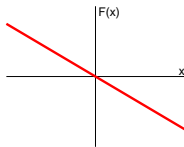
- springs are usually well-modeled by Hooke's law $F(x) = -kx$ for small displacements x from the equilibrium position
- ... *but* when they are over-extended, or closed coil, etc. then they need different models $m\ddot{x} = F(x)$



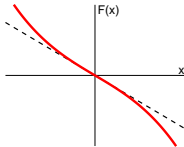
Open Coil



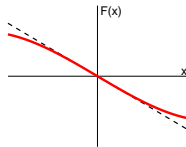
Closed Coil



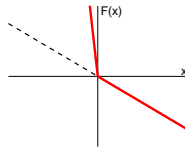
linear



hard



soft



closed

exercise #9: (numerical) nonlinear spring

- so $F(x) = -x - x^3$ is a hard spring model
- suppose we also have damping (thus $x(t) \rightarrow 0$ as $t \rightarrow \infty$)

exercise #9 in §5.3: numerically solve

$$\frac{d^2x}{dt^2} + \frac{dx}{dt} + x + x^3 = 0, \quad x(0) = -3, x'(0) = 8$$

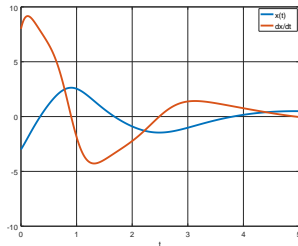
solution: write as system using $x = z_1$, $x' = z_2$:

$$z_1' = z_2 \quad z_1(0) = -3$$

$$z_2' = -z_2 - z_1 - z_1^3 \quad z_2(0) = 8$$

and use ode45:

```
>> f = @(t,z) [z(2); -z(2)-z(1)-z(1)^3];  
>> [tt,zz] = ode45(f,[0:.01:5],[-3;8]);  
>> plot(tt,zz), xlabel t, grid on  
>> legend('x(t)', 'dx/dt')
```



bullet to geosynchronous orbit

example. We want to use a bullet weighting 100 grams to destroy a satellite in geosynchronous (geostationary) orbit, approximately 36,000 km. What velocity is needed if we ignore air drag?

solution. Constant gravity g will *not* do.

The gravity decreases as the bullet rises.

From §5.3 we read Newton's law of gravitation:

$$my'' = -k \frac{Mm}{y^2} \text{ where } m = (\text{bullet mass}), M = (\text{earth mass})$$



After simplification (see text), and with initial conditions, this is

$$y'' = -g \frac{R^2}{y^2}, \quad y(0) = R, \quad y'(0) = V$$

We take $R = 6.4 \times 10^6$ m = (radius of earth) and $g = 9.8$. (*Note bullet mass does not matter. Earth's mass is built into g .*)

Question: Find V so that the maximum of $y(t)$ is 3.6×10^7 m.

bullet to geosynchronous orbit 2

question: Find V so $\max y(t) = 3.6 \times 10^7$, given

$$y'' = -g \frac{R^2}{y^2}, \quad y(0) = R, \quad y'(0) = V$$

and $R = 6.4 \times 10^6$ m =(radius of earth) and $g = 9.8$

solution?: as system with $y = z_1$, $y' = z_2$ and $C = gR^2$:

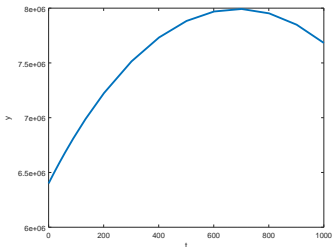
$$z_1' = z_2$$

$$z_1(0) = R$$

$$z_2' = -Cz_1^{-2}$$

$$z_2(0) = V$$

```
>> g = 9.8; R = 6.4e6; C = g*R^2;  
>> f = @(t,z) [z(2); -C/z(1)^2];  
>> V = 5000;  
>> [tt,zz] = ode45(f,[0,1000],[R;V]);  
>> plot(tt,zz(:,1))  
>> xlabel t, ylabel y  
>> max(zz(:,1))  
ans = 7.9924e+06
```



bullet to geosynchronous orbit 3

- trial and error needed!
- I finished with:

```
>> V = 10157; [tt,zz] = ode45(f,[0,20000],[R;V]);  
>> [max(zz(:,1)) zz(end,1)]  
ans =  
    3.60120e+07    2.36604e+07
```

a bit of hard-earned **extra credit** for any of these:

- ① energy methods allow you to solve the above problem *by hand*; see Mini-Project 3 and do so
- ② one can add air drag by a reasonable model and use same numerical methods; do so
- ③ given air drag from ②, will the bullet survive the heating? (ceramic bullet?)
 - this will need another DE coupled to the first

how the black box works

- how does the black box `ode45` work?
 - good question!
- *basically*: it is just a fancier form of Euler's method
- *more thoroughly*:
 - it uses a pair of [Runge-Kutta](#) methods
 - ... so it can adaptively choose its step size
 - see the [MATLAB reference page for ode45](#)
 - covered in Chapter 9 (= [Week 10](#))

dependent variable missing

- there are by-hand solvable nonlinear 2nd-order DEs:

DE	technique	first integral
$y'' = f(t, y, y')$	<i>too general</i>	
$y'' = f(t)$	just antidifferentiate	$y' = F(t) + c$ where $F(t) = \int f(t) dt$
$y'' = f(y)$	compute energy [Mini-Project 3]	$\frac{1}{2}(y')^2 + P(y) = c$ where $P(z) = -\int f(z) dz$
$y'' = f(y')$	substitute $u = y'$ [§4.10]	$Q(y') = t + c$ where $Q(u) = \int \frac{du}{f(u)}$

- last category called “dependent variable y is missing” (§4.10)
- you can often solve by the substitution $u = y'$
 - this can sometimes work for $y'' = f(t, y')$ too

exercise #6 in §4.10

exercise. find the general solution:

$$e^{-t}y'' = (y')^2$$

expectations

- just watching this video is *not* enough!
 - see “found online” videos at bueler.github.io/math302/week8.html
 - *read* section 4.10 in the textbook
 - skip the “Use of Taylor series” material ... we'll get to it later
 - *read* section 5.3 in the textbook
 - you can safely skip the material on “Telephone wires” (a boundary value problem ... not in Math 302)
 - *do* the WebAssign exercises for section 5.3, which include some problems from 4.10
 - compare Mini-Project 3 to these slides ... different but some overlap of ideas