

Rust Quantum Library

The quantum library in Rust contains several useful measures of quantum entanglement that are calculated from the density matrix representing a quantum state. This library is generalized for any density matrix with dimension of 2^n . The Rust library uses `ndarray` and the `ndarray.linalg` packages.

We can run Rust from a Jupyter notebook by installing EvCxR Jupyter kernel. The Jupyter notebook runs with the **plotters** crate. Setup can be found from <https://datacrayon.com/posts/programming/rust-notebooks/setup-anaconda-jupyter-and-rust>. Miniconda is not necessary, Anaconda will work the same. Be careful when installing the actual kernel in the tutorial (shown below) - it is important to check on **"other installation methods"** in order for proper installation for Linux, Windows, and Mac.

Install the EvCxR Jupyter Kernel

Now we'll install the [EvCxR Jupyter Kernel](#). If you're wondering how it's pronounced, it's [been mentioned to be "Evic-ser"](#). This is what will allow us to execute Rust code in a Jupyter Notebook.

You can get [other installation methods](#) methods for EvCxR if you need then, but we will be using:

```
cargo install evcxr_jupyter --version 0.5.3
evcxr_jupyter --install
```

Below are the various quantum functions along with the functions that perform linear algebra operations.

1 Quantum functions

1.1 Create a density matrix ρ (*create_density_matrix*)

Input: The wavefunction expressed as a column vector of coefficients

Input type: pub type VecC64 = ndarray::Array1<c64>

Input Example: For the Bell state $\phi_+ = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$,

```
let norm_const = 1./2_f64.sqrt();
let bell_phi_plus = array![c64::new(norm_const , 0.0) , c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) ,
c64::new(norm_const , 0.0)];
```

Output: The density matrix, which is square and Hermitian.

Output type: pub type MatrixC64 = ndarray::Array2<c64>

Output Example: dens matrix for phi_+ =

```
[[0.4999999999999999+0i, 0+0i, 0+0i, 0.4999999999999999+0i],
[0+0i, 0+0i, 0+0i, 0+0i],
[0+0i, 0+0i, 0+0i, 0+0i],
[0.4999999999999999+0i, 0+0i, 0+0i, 0.4999999999999999+0i]]
```

The density matrix is a more general way of representing the state of a quantum system. It is represented as $\rho = |\psi\rangle \langle\psi|$, where the wavefunction $|\psi\rangle$ is represented as a column vector.

With this same example, the function performs the operation

$$\rho = |\phi_+\rangle \langle\phi_+| = \frac{1}{2} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 & 1 \end{pmatrix}^* = \frac{1}{2} \begin{pmatrix} 1 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 1 \end{pmatrix}$$

1.2 Find purity (*find_purity*)

Input: ρ (Density Matrix)

Input Type: pub type MatrixC64 = ndarray::Array2<c64>

Input Example:

```
array![ [c64::new(0.5 , 0.) , c64::new(0. , 0.) ],  
[c64::new(0. , 0.) , c64::new(0.5 , 0.)] ]
```

Output: Float

Output Type: f64

Output Example: 0.5

Simply returns the trace of ρ^2 as a f64 type. For an idea on where this purity lies within the range of possible values, print a statement such as `println!("The purity lies between {} and 1", 1./ (find_dim(rho.sqrd)))`. The *find_dim* function finds the dimension of the square matrix. For example, a 4 x 4 matrix has `dim = 4`.

1.3 Find fidelity (*find_fidelity*)

Input: Two density matrices ρ_1 and ρ_2

Input Type: pub type MatrixC64 = ndarray::Array2<c64>

Input Example: let rho_1 = array![

```
[c64::new(0.25 , 0.0) , c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) ],  
[c64::new(0.0 , 0.0) , c64::new(0.25 , 0.0) , c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) ],  
[c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) , c64::new(0.25 , 0.0) , c64::new(0.0 , 0.0) ],  
[c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) , c64::new(0.25 , 0.0)]  
];
```

let rho_2 = rho_1;

Output: A number between 0 and 1

Output Type: f64

Output Example: 1

This function is a distance measurement of two quantum states ρ_1 and ρ_2 . It is expressed as $F = \text{tr} \sqrt{\sqrt{\rho_1} \rho_2 \sqrt{\rho_1}}$. This function calls *find_sqr_root_of_matrix* to calculate $\sqrt{\rho_1}$ and the overall square root of $\sqrt{\rho_1} \rho_2 \sqrt{\rho_1}$.

1.4 Find concurrence (*find_concurrence*)

Input: The density matrix ρ

Input Type: pub type MatrixC64 = ndarray::Array2<c64>

Input Example: The state $\cos(\theta) |00\rangle + \sin(\theta) |11\rangle$ for $\theta = 30^\circ$ is represented by the vector:

```
pub const THETA: f64 = PI*(30./180.);  
let psi_part_entangled: VecC64 = array![c64::new(THETA.cos(), 0.0), c64::new(0.0, 0.0), c64::new(0.0,  
0.0), c64::new(THETA.sin(), 0.0)];
```

Output: A number between 0 and 1

Output Type: f64

Output Example: 0.866 (which is the rounded from $\sqrt{3}/2$)

For two qubits, concurrence $C(\rho) = \max(0, \lambda_1 - \lambda_2 - \lambda_3 - \lambda_4)$, where $\lambda_1, \dots, \lambda_4$ are the eigenvalues in decreasing order (ie λ_1 is the highest eigenvalue) of the matrix $R = \sqrt{\sqrt{\rho} \tilde{\rho} \sqrt{\rho}}$.

In this case, $\tilde{\rho} = (\sigma_y \otimes \sigma_y) \rho^* (\sigma_y \otimes \sigma_y)$. ρ^* is the complex conjugate of ρ , and σ_y is the Pauli-y spin matrix.

This function has been extended from 2 qubits (4 x 4 density matrix) to 3 or more qubits. The function will give wrong answers for odd numbers of qubits.

The extension uses $C(\rho) = \max(0, \lambda_1 - \sum \lambda_i)$, where i goes from 2 to the density matrix dimension. The dimension of the density matrix is 2^n , where n is the number of qubits. Also in this case $\tilde{\rho} = (\sigma_y \otimes^n \sigma_y) \rho^* (\sigma_y \otimes^n \sigma_y)$.

The tensor product operation for \otimes^n is called using the *find_tensor_product* function, and uses an iterator with the fold method to perform the tensor product n times.

1.5 Find trace norm (*find_trace_norm*)

Input: The density matrix ρ

Input Type: pub type MatrixC64 = ndarray::Array2<c64>

Input Example:

For the Bell state $\phi_+ = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$,

```
let norm_const = 1./2_f64.sqrt();
```

```
let bell_phi_plus = array![c64::new(norm_const , 0.0) , c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) ,
c64::new(norm_const , 0.0)];
```

Output: A float from 0 to 1

Output Type: f64

Output Example: 1

The trace norm is expressed as $\|\rho\|_1 = \text{tr} \sqrt{\rho^\dagger \rho}$. ρ^\dagger is the complex conjugate and transpose of the matrix ρ . The function calls *find_sqr_root_of_matrix*($\rho^\dagger \rho$). For Hermitian, normalized density matrices, the output is always 1.

1.6 Find negativity (*find_negativity*)

Input: The density matrix ρ

Input Type: pub type MatrixC64 = ndarray::Array2<c64>

Input Example: For a 5 qubit maximally mixed density matrix,

```
let rho_mixed_diag: VecC64 = Array::from_elem( 32, c64::new(1./32., 0.0) );
```

```
let rho_mixed: MatrixC64 = MatrixC64::from_diag(rho_mixed_diag);
```

Output: A number between 0 and 1

Output Type: f64

Output Example: 0.0

The negativity can be expressed as $N = \frac{\|\rho^{\Gamma_A}\|_1 - 1}{2}$ where ρ^{Γ_A} is the partial transpose of a substate of ρ . $\|\rho^{\Gamma_A}\|_1$ is the trace norm of that partial transpose. Although the trace norm for density matrices is 1, the trace norm of the partially transposed matrix can vary from 0 to 1. The Peres-Horodecki criterion separates ρ into the tensor product of two states A and B. ρ^{Γ_A} is called with the function *find_partial_transpose*(ρ).

1.7 Find log negativity (*find_log_negativity*)

Input: The density matrix ρ

Input Type: pub type MatrixC64 = ndarray::Array2<c64>

Input Example: For the Bell State $\phi_- = \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle)$:

```
pub const NORM_CONST: f64 = 1./std::f64::consts::SQRT_2;
```

```
let bell_phi_minus_vec: VecC64 = array![ c64::new(NORM_CONST , 0.0) , c64::new(0.0 , 0.0) ,
c64::new(0.0 , 0.0) , c64::new(-NORM_CONST , 0.0) ];
```

```
let rho_bell_psi_plus = create_density_matrix(BELL_PSI_PLUS_VEC);
```

Output: A number between 0 and 1

Output Type: f64

Output Example: 0.0

This function lets $N = \text{find_negativity}(\rho)$. Then it takes $\log_2(2N + 1)$.

1.8 Find Schmidt number (*find_schmidt_number*)

Input: The joint spectral intensity (JSI) matrix

Input Type: pub type MatrixF64 = ndarray::Array2<f64>

Input Example: A 200 x 200 JSI (located inside the examples folder)

Output: An entanglement measurement of the JSI.

Output Type: f64

Output Example: 6.74 (rounded)

This computation is done in several steps:

- 1) Take the square root of the elements of the JSI (joint spectral intensity) to find the JSA (joint spectral amplitude). The JSI is a matrix of floats from SPDCalc.
 - 2) Find the singular value decomposition (SVD) of the JSA. A SVD calculation decomposes the into three matrices, USV^T .
 - 3) A normalization constant is found by taking the inverse of the sum of the squares of the singular values of the S matrix (eigenvalues), i.e. $A = 1/\sum \lambda_i^2$.
 - 4) The S matrix is renormalized by multiplying the $\sqrt{A} * S$.
 - 5) From the normalized S matrix, take the sum of the eigenvalues to the fourth power, i.e. $\sum \lambda_{i,norm}^4$.
- The Schmidt number is $k = 1/\sum \lambda_{i,norm}^4$.

1.9 Create two source Hong-Ou-Mandel graph (*find_two_source_hom*)

Inputs: A signal vector of frequency modes ω , an idler vector of modes ω , a joint spectral amplitude (JSA) of modes ω , and a vector of time intervals dt in femtoseconds.

Input Types: signal: VecF64, idler: VecF64, jsa: MatrixC64, dt: VecF64

Output: A graph of time intervals (dt) on the x axis and coincidence probabilities on the y axis.

Output Example An example is located in the "two_source_hom_plot.ipynb" under qm/examples.

When calculating the two-source Hong-Ou-Mandel graph, we consider four photons from two identical crystals, i.e. two signal photons and two idler photons. The relevant integral to compute is

$$\int_{\omega_s} \int_{\omega_i} \int_{\omega'_s} \int_{\omega'_i} \psi(\omega_s, \omega_i, \omega'_s, \omega'_i) - e^{i\phi} \psi(\omega'_s, \omega_i, \omega_s, \omega'_i) d\omega'_i d\omega'_s d\omega_i d\omega_s.$$

The wave function ψ is equivalent to the joint spectral amplitude (JSA) for all four photons. Since the photons come from two separate sources, this joint amplitude is separable into two parts, so that

$$\psi(\omega_s, \omega_i, \omega'_s, \omega'_i) = JSA_1(\omega_s, \omega_i) \otimes JSA_2(\omega'_s, \omega'_i)$$

. Since the two sources are identical, we conclude that $JSA_1 = JSA_2$. So we only need to compute the JSA once and then simply reference different indices for the different signal and idler photons. Using this notation, the integrand becomes

$$JSA(\omega_s, \omega_i) \otimes JSA(\omega'_s, \omega'_i) - e^{i\phi} JSA(\omega'_s, \omega_i) \otimes JSA(\omega_s, \omega'_i)$$

.

We can simply iterate over many sums rather than explicitly calculating this integral. The variable ϕ depends on time and wavelength difference between two photons of interest. For signal-signal interference, $\phi_{ss} = 2\pi\Delta t(\frac{1}{\lambda_s} - \frac{1}{\lambda_i})$. For idler-idler interference, $\phi_{ii} = 2\pi\Delta t(\frac{1}{\lambda'_s} - \frac{1}{\lambda'_i})$. For signal-idler interference, $\phi_{si} = 2\pi\Delta t(\frac{1}{\lambda_s} - \frac{1}{\lambda'_i})$.

In the code, we let A, B, C, D stand for the different JSA configurations, between ω_s and ω_i , ω'_s and ω'_i , ω'_s and ω_i , and ω_s and ω'_i respectively.

The `find_two_source_hom_norm` function normalizes the coincidence probabilities.

1.10 Find two source Hong-Ou-Mandel norm (*find_two_source_hom_norm*)

Inputs: A signal vector of frequency modes ω , an idler vector of modes ω , and a joint spectral amplitude (JSA) of modes ω .

Input Types: signal: VecF64, idler: VecF64, jsa: MatrixC64

Output: A value that normalizes the two_source_hom for each time interval dt.

Output Type: f64

Output Example: 6.41e+61_f64 (rounded)

Normalizes the coincidence probabilities in the two source HOM calculation.

2 Matrix Operations

2.1 Find the symmetric square root of a matrix (*find_symmetric_square_root*)

Input: Any square, Hermitian matrix

Input Type: pub type MatrixC64 = ndarray::Array2<c64>

Input Example:

```
array![ [c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) ],
[c64::new(0.0 , 0.0) , c64::new(0.5 , 0.0) , c64::new(0.0 , 0.5) ],
[c64::new(0.0 , 0.0) , c64::new(0.0 , -0.5) , c64::new(0.5 , 0.0) ],];
```

Output: The square root of the matrix

Output Type: pub type MatrixC64 = ndarray::Array2<c64>

Output Example:

```
array![ [c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) , c64::new(0.0 , 0.0) ],
[c64::new(0.0 , 0.0) , c64::new(0.5 , 0.0) , c64::new(0.0 , 0.5) ],
[c64::new(0.0 , 0.0) , c64::new(0.0 , -0.5) , c64::new(0.5 , 0.0) ],];
```

Computes the symmetric square root of a Hermitian matrix. In the example above, the square root of the matrix is equivalent to the matrix itself.

***Bug in LAPACK involving numerical precision

For some of the entanglement measurements, the square root of the density matrix is required. Using a LAPACK function such as `matrix.ssqr(UPLO:Lower).unwrap()` is the easiest and fastest route, but LAPACK has a bug and fails to compute. It's suspected that the matrix needs to be semi-positive definite (eigenvalues need to be ≥ 0) in order for it to work.

So, the function was written by decomposing the Hermitian matrix as $M = SDS^{-1}$, where S and S^{-1} are complex unitary matrices and D is a real diagonal matrix. S is a matrix of the eigenvectors of M , and D has the eigenvalues of M in the diagonal. Since D is diagonal, one can take the square of the elements in order to find \sqrt{D} . So, $\sqrt{M} = S\sqrt{D}S^{-1}$.

But, one problem is that again, when doing this decomposition, the matrix M needs to be semi-positive definite. It usually is, but LAPACK will sometimes find that one eigenvalue is a very small negative number, such as 2.8×10^{-15} . A rescaling function is called in order to fix this issue.

The rescaling function takes matrix M as an input and finds its eigenvalues and eigenvectors, and sets any negative eigenvalues to 0. Then, matrix M can be decomposed into $M = SDS^{-1}$. The rescaling function returns (D, S) as a tuple. Then, the square root function returns $\sqrt{M} = S\sqrt{D}S^{-1}$.

Unfortunately after rescaling the function and calling LAPACK with `matrix.ssqr(UPLO:Lower).unwrap()` still fails in the computation. This is why (D, S) is returned as a tuple rather than a rescaled density matrix.

2.2 Find matrix dimension (*find_dim*)

Input: Square Matrix

Input Type: `pub type MatrixC64 = ndarray::Array2<c64>`

Input Example:

```
array![ [c64::new(3., 1.) , c64::new(-1., 1.) ],
[c64::new(2., -1.) , c64::new(-2., -1.)] ]
```

Output: Integer

Output Type: i32

Output Example: 2

A simple function that calculates the length of a square matrix. For example, if a matrix is 8 x 8, this will return 8 as a type i32 (integer 32 bit).

2.3 Find partial transpose (*find_partial_transpose*)

Input: Any square matrix

Input Type: pub type MatrixC64 = ndarray::Array2<c64>

Input Example: let matrix: MatrixC64 = array![

```
[c64::new(1.0 , 0.0) , c64::new(2.0 , 0.0) , c64::new(3.0 , 0.0) , c64::new(4.0 , 0.0) ] ,
[c64::new(5.0 , 0.0) , c64::new(6.0 , 0.0) , c64::new(7.0 , 0.0) , c64::new(8.0 , 0.0) ] ,
[c64::new(9.0 , 0.0) , c64::new(13.0 , 0.0) , c64::new(11.0 , 0.0) , c64::new(12.0 , 0.0)] ,
[c64::new(10.0 , 0.0) , c64::new(14.0 , 0.0) , c64::new(15.0 , 0.0) , c64::new(16.0 , 0.0)] ];
```

Output: A square matrix

Output Type: pub type MatrixC64 = ndarray::Array2<c64>

Output Example:

```
[c64::new(1.0 , 0.0) , c64::new(2.0 , 0.0) , c64::new(3.0 , 0.0) , c64::new(7.0 , 0.0) ] ,
[c64::new(5.0 , 0.0) , c64::new(6.0 , 0.0) , c64::new(4.0 , 0.0) , c64::new(8.0 , 0.0) ] ,
[c64::new(9.0 , 0.0) , c64::new(10.0 , 0.0) , c64::new(11.0 , 0.0) , c64::new(12.0 , 0.0)] ,
```

```
[c64::new(13.0 , 0.0) , c64::new(14.0 , 0.0) , c64::new(15.0 , 0.0) , c64::new(16.0 , 0.0)] ];
```

The partial transpose operation splits an even-numbered dimension matrix into four symmetrical blocks. The top right block and bottom left block are transposed, and then the matrix is put back together again.

2.4 Find tensor product (*find_tensor_product*)

Input: One square matrix with dim $m \times m$ and another with dim $n \times n$

Input Type: pub type MatrixC64 = ndarray::Array2<c64>

Input Example: let matrix: MatrixC64 = array![

```
[c64::new(1.0 , 0.0) , c64::new(1.0 , 0.0)] ,
[c64::new(1.0 , 0.0) , c64::new(-1.0 , 0.0)] ,
];
```

Output: A square matrix with dim $mn \times mn$

Output Type: pub type MatrixC64 = ndarray::Array2<c64>

Output Example:

```
[c64::new(1.0 , 0.0) , c64::new(1.0 , 0.0) , c64::new(1.0 , 0.0) , c64::new(1.0 , 0.0) ] ,
[c64::new(1.0 , 0.0) , c64::new(-1.0 , 0.0) , c64::new(1.0 , 0.0) , c64::new(-1.0 , 0.0) ] ,
[c64::new(1.0 , 0.0) , c64::new(1.0 , 0.0) , c64::new(-1.0 , 0.0) , c64::new(-1.0 , 0.0)] ,
[c64::new(1.0 , 0.0) , c64::new(-1.0 , 0.0) , c64::new(-1.0 , 0.0) , c64::new(1.0 , 0.0)] ];
```

Computes the tensor product of any matrix.