

21/04/25



## Lecture - 22

### Stress - strain Tensors :

06/11/20

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శాస్త్రానుష్ఠాన విభాగము

It follows that the ratio between the components of the elastic tensor  $\Phi_{uw}^{ij}$  and a length is homogeneous with  $\rho \ddot{\psi}_u$ , whose dimensions are those of a force per unit volume. In conclusion, the dimensions of the elastic tensor are those of a force per unit area, that is, a stress. Defining the *stress tensor* as

$$\sigma_{uw} = \sum_{i,j=1}^3 \Phi_{uw}^{ij} \epsilon_{ij},$$

whose dimensions are the same as those of  $\Phi_{uw}^{ij}$ , one may recast the long-wavelength relation for the acoustic modes as

$$\rho \ddot{\psi}_u = \sum_{u=1}^3 \frac{\partial \sigma_{uw}}{\partial x_w},$$

which provides the elasticity law for a continuous medium. The stress tensor is symmetric because  $\Phi_{uw}^{ij}$  is invariant when  $u$  and  $w$  are exchanged. The relation between the strain and stress tensors is the generalization of *Hooke's law* to an anisotropic medium.

R. Hooke (1635-1703) first mentioned the law in 1676 as an anagram, *ceiiinosssttuu*, which he explained in 1678 as *ut tensio sic vis*.

# Advanced Solid-State Sensors M

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# Course Program 2019-2020

1. Introduction – Classification of sensors
2. Main physical effects used in solid-state sensors  
and simple readout circuits
3. Resolution & Noise
4. Image sensors
5. Mechanical sensors

# Mechanical Sensors

Characterization of mechanical sensors

- the measurand and the physical structures
- 1. Most relevant mechanical signals, their units in the S.I., and ranges for specific applications
- 2. Definition of the stress and strain tensors
- 3. Description of the structural elements mainly used in solid-state sensors (beams & plates)

# List of mechanical measurands:

subset of six signals covering the most important  
classes of mechanical microsensors

- Pressure/Stress
- Acceleration/Deceleration
- Displacement
- Flow rate
- Force/Torque
- Position/Angle

# Automotive accelerometers

Application	Range
Frontal airbags	+/- 50 g
Lateral airbags	+/- 100 – 250 g
Suspensions	+/- 2 g
ABS (Antilock Braking System)	+/- 1 g

**Units are referred to g, with  $g = 9,8 \text{ m/s}^2$**

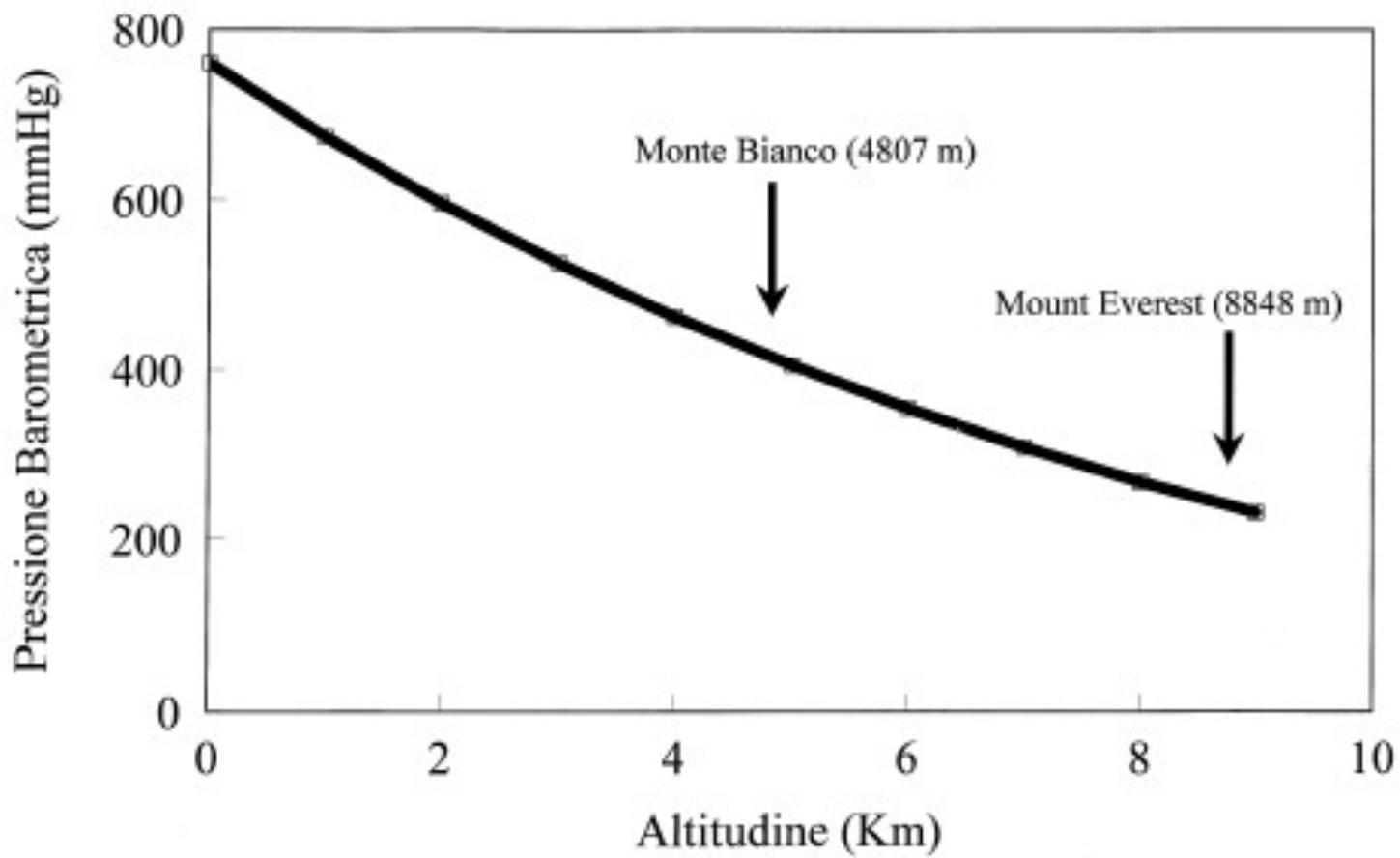
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# Pressure units

<b>1 Pa</b>	International System
<b>1 KPa</b>	$1\text{e}3 \text{ Pa}$
<b>1 kgf/m<sup>2</sup></b>	$9,8 \text{ Pa}$ $(g = 9,8 \text{ m/s}^2)$
<b>1 psi</b> (pound-force per square inch)	$6,89\text{e}3 \text{ Pa}$ $(1\text{inch} = 25.4\text{mm}; 1\text{Kg} = 2,2046\text{p})$
<b>1 mmHg</b> $(= 1 \text{ Torr})$	$1,33\text{e}2 \text{ Pa}$ $(p_{\text{Hg}} = 13,59 \text{ g/cm}^3, \text{ liquid})$
<b>1 atm</b>	$760 \text{ Torr} = 1,01\text{e}5 \text{ Pa}$ $(\text{equivalent to } 10,33 \text{ m of water})$
<b>1 bar</b>	$1\text{e}5 \text{ Pa}$

# Applications of silicon pressure sensors

Application	Pressure Range
Manifold pressure	15 - 250 kPa
Fuel pressure	15 - 400 kPa
Tyre pressure	1,8 –2,5 bar
Blood pressure (max 120 – 130 mmHg; min 70 – 80 mmHg)	30 - 300 mmHg
Barometric pressure	in hectopascal = mbar o in mmHg



**Figura 1.** Modificazione della pressione barometrica in relazione all'altitudine secondo le tabelle dell'Organizzazione dell'Aviazione Civile Internazionale.

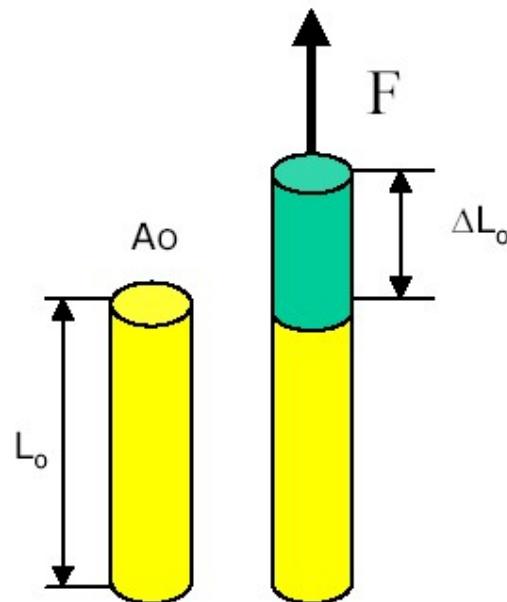
# Mechanical Sensors

## Characterization of mechanical sensors

- the measurand and the physical structures
  1. Most relevant mechanical signals, their units in the S.I., and ranges for specific applications
  2. **Definition of the stress and strain tensors**
  3. Description of the structural elements mainly used in solid-state sensors (beams & plates)

## Definition of stress and strain tensors

- one dimensional case of Hooke's law



Stress:  $\sigma = F/A_0$  [Pa = N/m<sup>2</sup>]

Strain:  $\varepsilon = \Delta L_0/L_0$

HOOKE's LAW:

$E = \sigma / \varepsilon$  Young's modulus [Pa]

\* we will also see Poisson coefficient that also describes the Thinning of the Bar.

List of hypothesis we already started examining!

## Assumptions

- 1) • We will consider elastic deformations and linear response → small-deformation hypothesis.
- 2) • In equilibrium, the resultant force over the body is 0, the resultant moment over the body is 0. (everything is static)
- 3) • The linear dimensions of the physical system are large enough to consider it as a continuous body (example: lattice distance in silicon  $a=5,431 \text{ \AA}$ , MEMS structures with  $L_{\min} \approx 50 \text{ nm}$ ).
- 4) • The surface elements becomes more and more relevant in small bodies:

$$l = 1 \text{ m} \rightarrow \text{Surface}/\text{Volume} = 6 \text{ m}^{-1};$$

$$l = 1 \text{ cm} \rightarrow \text{Surface}/\text{Volume} = 6 \text{ cm}^{-1} = 600 \text{ m}^{-1}$$

Gravity is a Volume Force

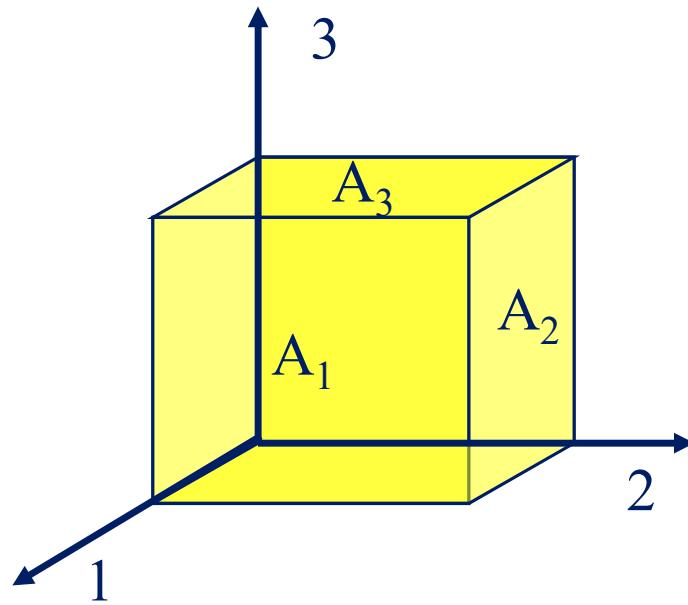
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\* when the size of the object becomes the surface forces dominate the volume forces. 11

→ we consider in our discussions materials of cubic type.

- Forces applied to the body (e.g, gravity)
- Forces applied at the surfaces (e.g., pressure, bending, twisting)

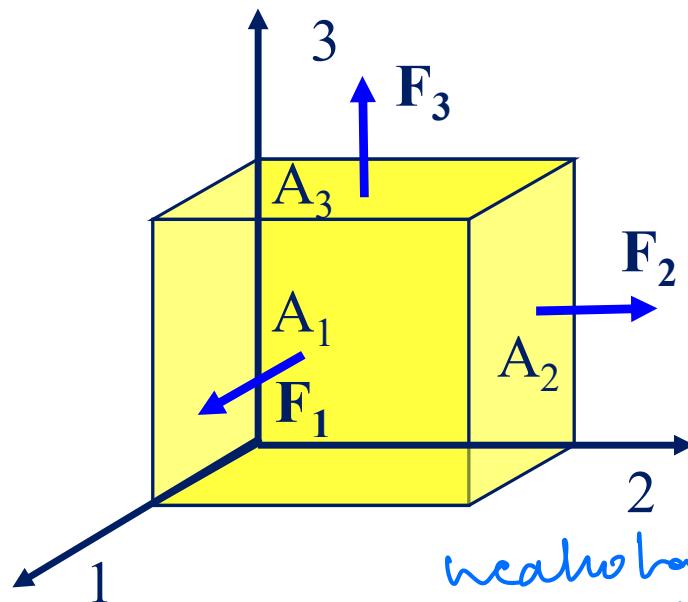
$A_1, A_2, A_3$  are the Areas  
of the Faces



- Remember when we apply a force, there is force & Torque but the object becomes smaller the force becomes more relevant.

- Normal stresses (or axial)

$F_1, F_2, F_3$  are Normal components of the Applied force.



we also have forces that are acting on the other 3 faces

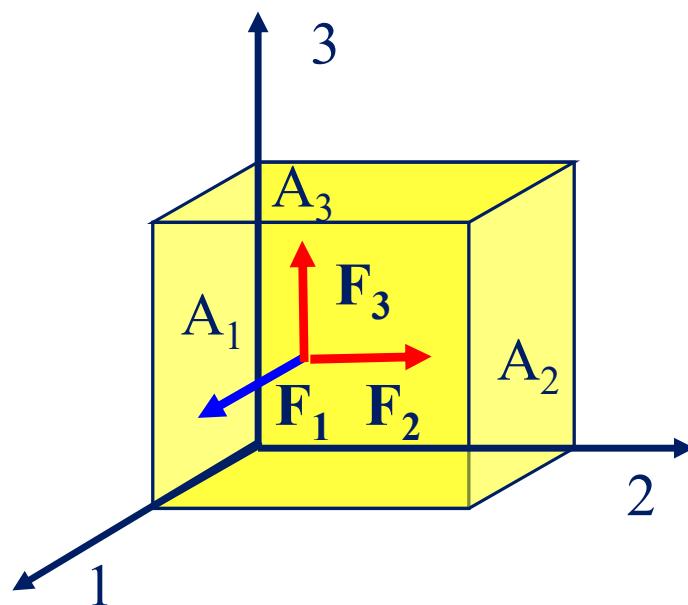
$$\sigma_{11} = F_1/A_1$$

$$\sigma_{22} = F_2/A_2$$

$$\sigma_{33} = F_3/A_3$$

In this figure we are considering only face  $A_1$

- Shear stresses



$$\sigma_{31} = \frac{F_3}{A_1}$$
$$\sigma_{21} = \frac{F_2}{A_2}$$
$$\sigma_{ij}$$

This Tensor is  
symmetric  
to Balance Rotational  
MOMENTS

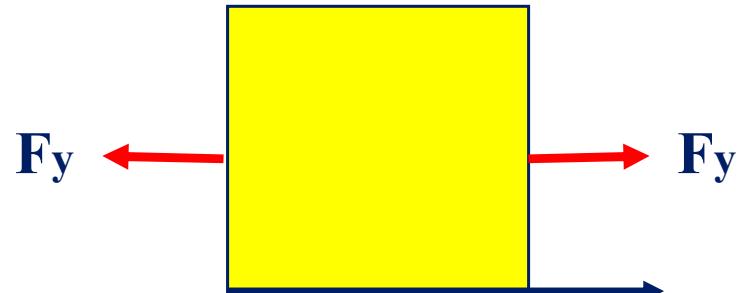
 Equilibrium condition: in order to have no accelerations of the body, all forces and moments must balance:

- i) The surface axial forces need to be equal on two opposite planes:

$$\sum F_x = 0$$

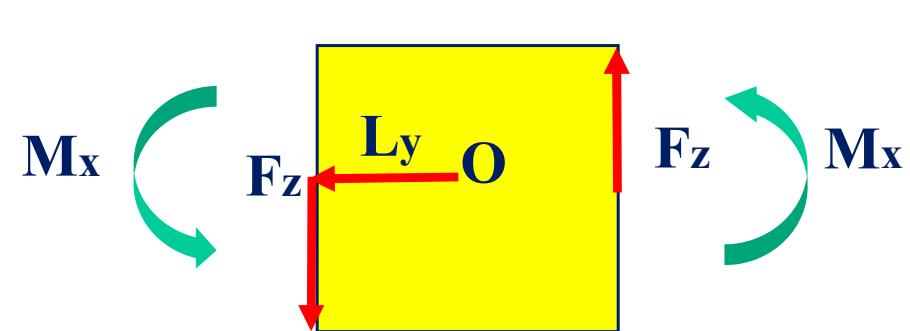
$$\sum F_y = 0$$

$$\sum F_z = 0$$



- ii) The shear forces, even if balanced, would induce a rotation moment:

$$M_x = 2 * L_y F_z$$



## Equilibrium condition:

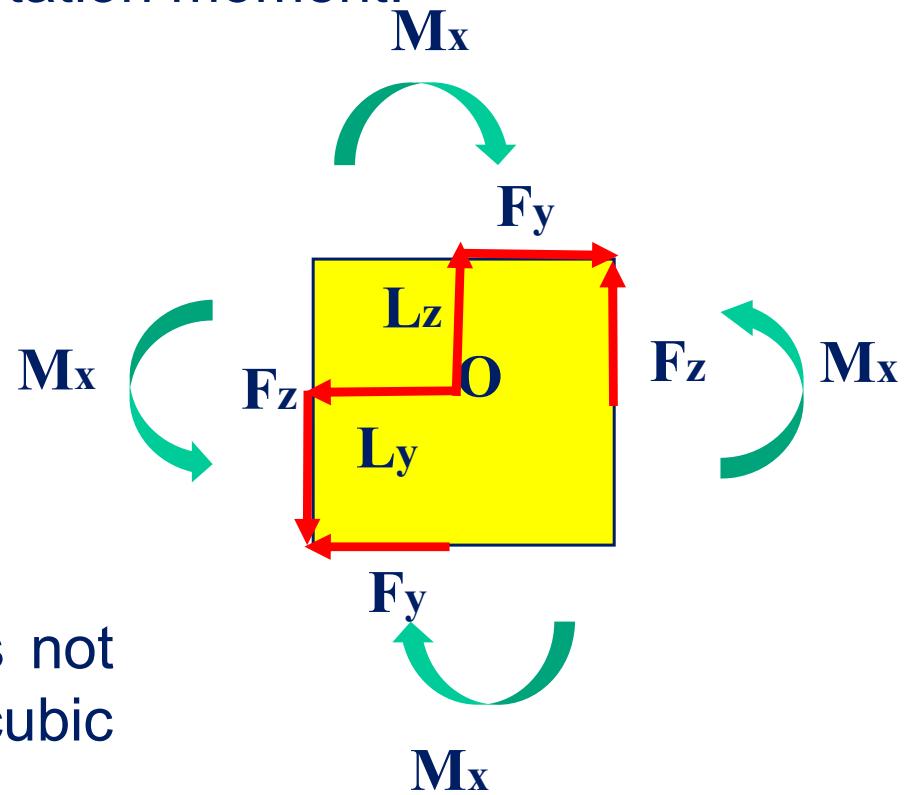
In order to have no accelerations of the body one must consider the action of a second couple of shear forces on the dual planes, balancing the rotation moment:

$$\Sigma M_x = 2 * L_y F_z - 2 * L_z F_y = 0$$

Balancing Rotational Moments

→  $\sigma_{ij} = \sigma_{ji}$

(for the above to apply it is not necessary to consider a cubic block)

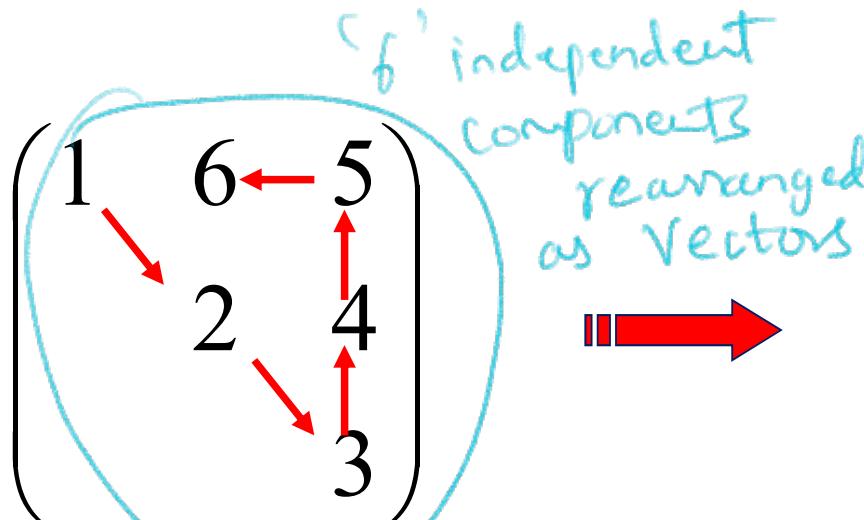


Note: Q) Does this argument for Equilibrium work only for cubic block ? 16

Any It doesn't depend on the shape, we will see the Answer later.

## STRESS TENSOR

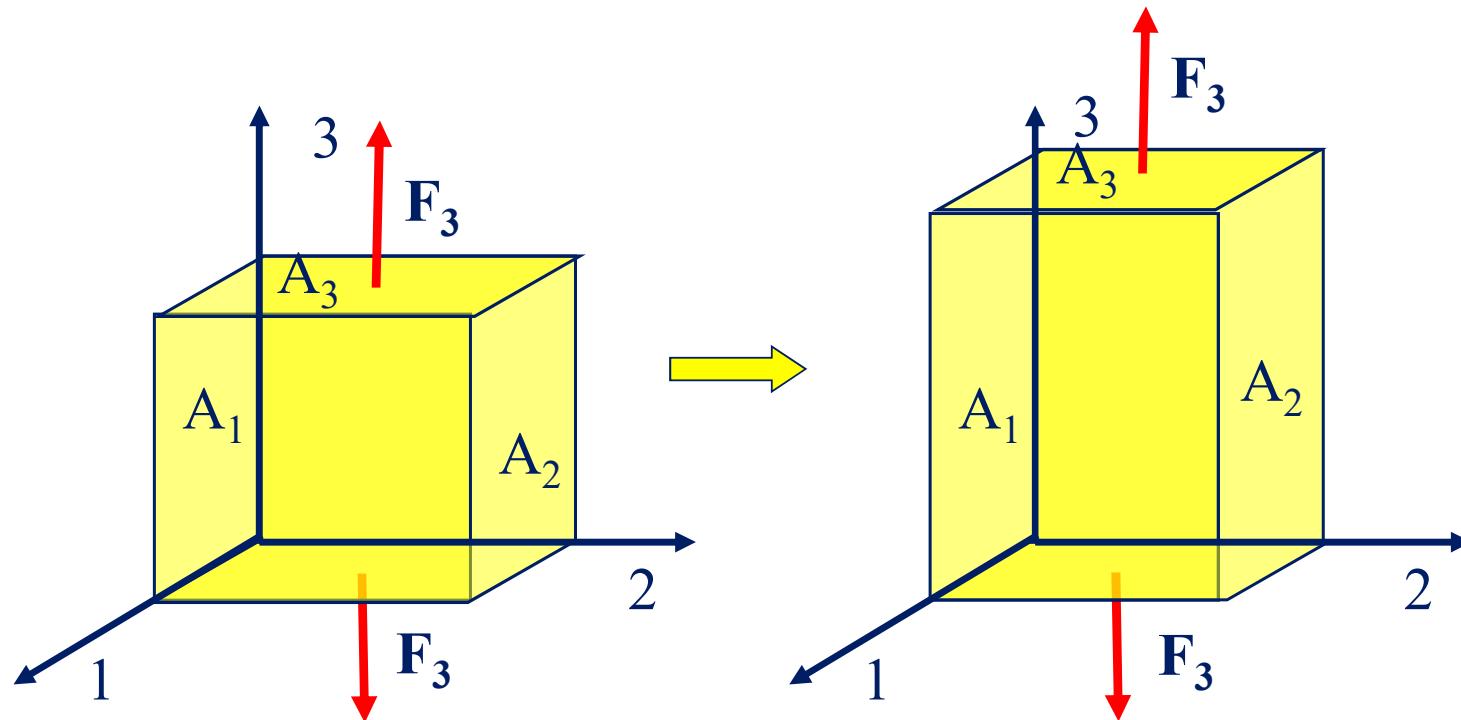
$$\begin{pmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{pmatrix} = \begin{pmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{12} & \sigma_{22} & \sigma_{23} \\ \sigma_{13} & \sigma_{23} & \sigma_{33} \end{pmatrix}$$



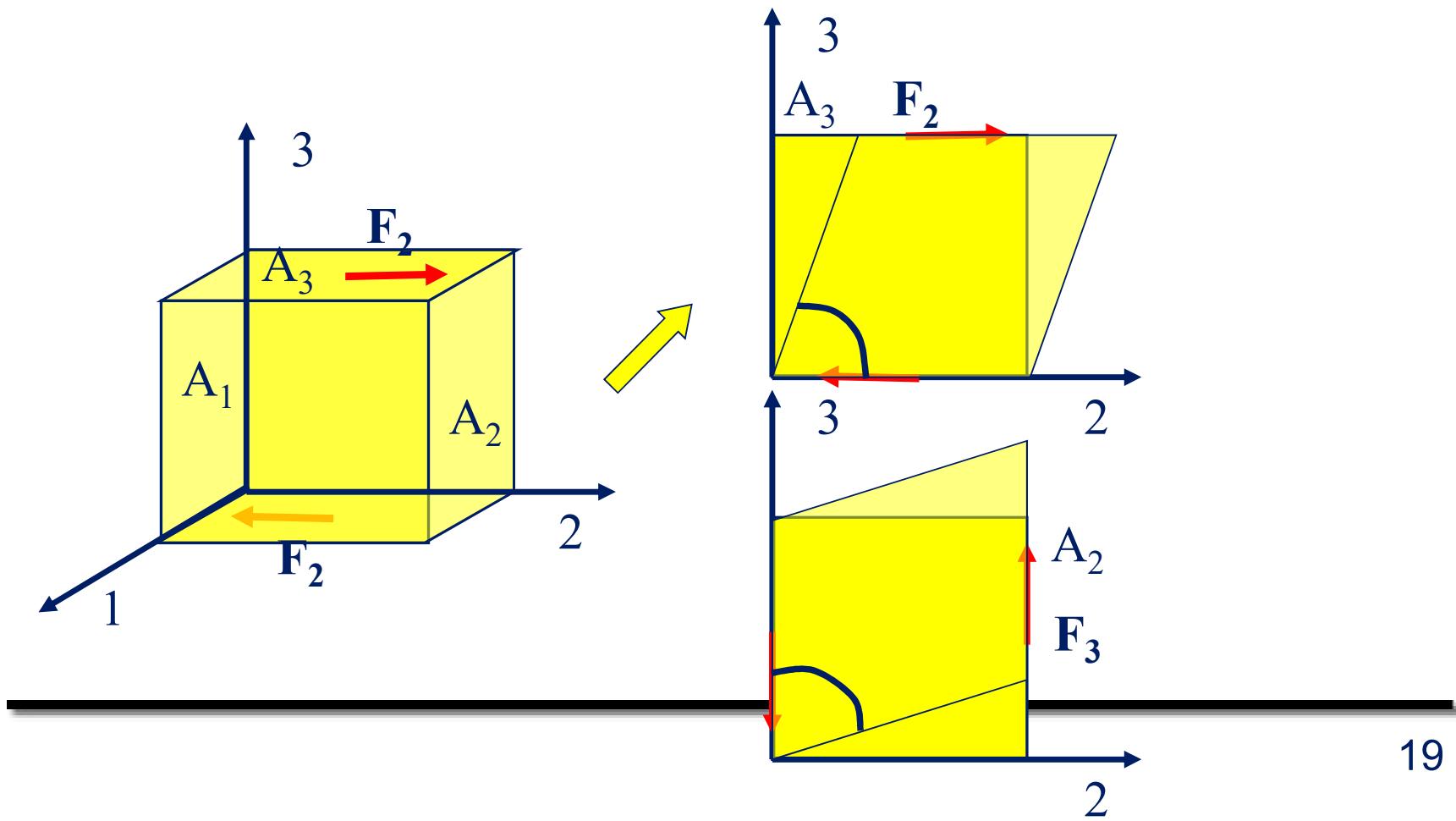
$$\begin{pmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{33} \\ \sigma_{23} \\ \sigma_{13} \\ \sigma_{12} \end{pmatrix}$$

# Deformations

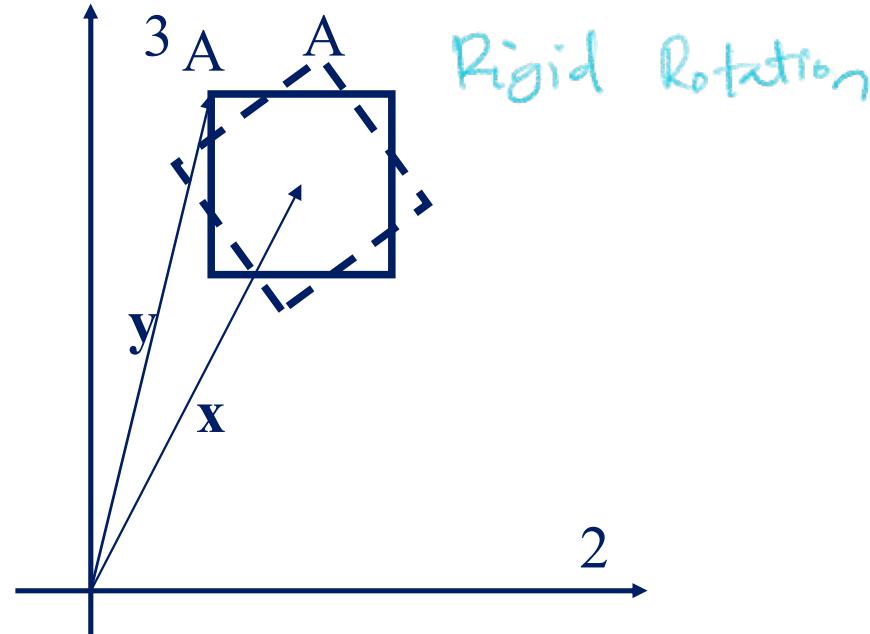
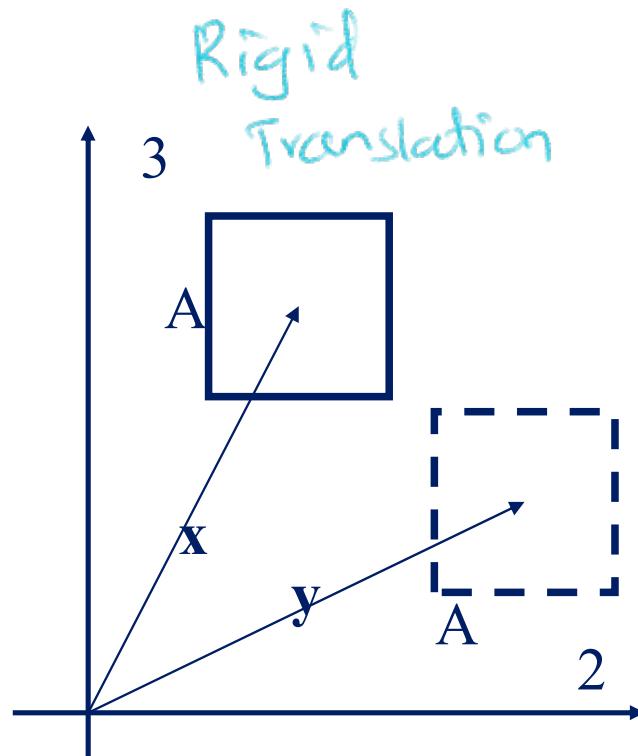
- Deformations due to axial stresses



- Deformations due to shear stresses



## Strain tensor



$$t = (y - x)$$

$$r = \omega \wedge (y - x)$$

→ We ignore this slide because we assumed that the

Block of material besides the Deformation the 20

Block of material also undergoes a rigid Translation  
or rigid rotation but we Neglect them.

We consider only local Deformations of the Body.

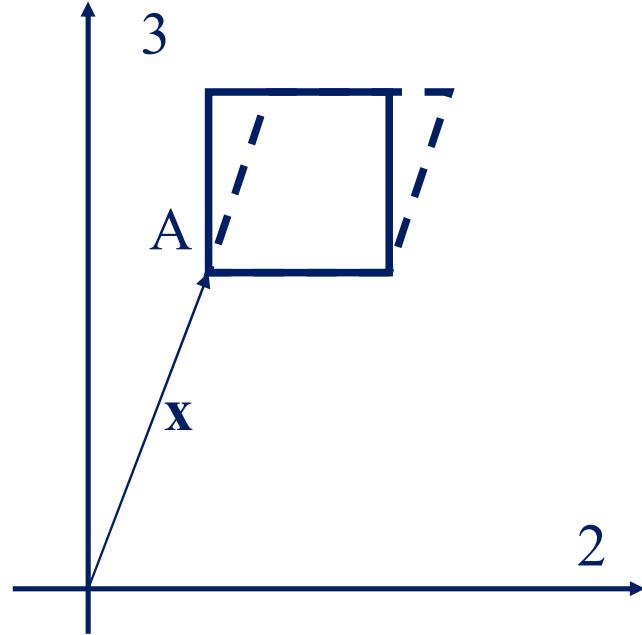
## Traslation and Rotation without deformation

$$\bar{\omega} = \begin{pmatrix} 0 & -\omega_3 & \omega_2 \\ \omega_3 & 0 & -\omega_1 \\ -\omega_2 & \omega_1 & 0 \end{pmatrix}$$

$$u = dx + \bar{\omega} dx$$



## Deformation → Strain tensor



$$\mathbf{u} = \mathbf{u}_0 + (\nabla_{\mathbf{x}} \mathbf{u}) d\mathbf{x}$$

$$(\nabla_{\mathbf{x}} \mathbf{u})_{ij} = \frac{\partial u_i}{\partial x_j} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) + \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial x_i} \right)$$

Anti symmetric Tensor  
that describes  
rotation  
it is  
neglected

$$(\nabla_x \mathbf{u})_{ij} = \frac{\partial u_i}{\partial x_j} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) + \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial x_i} \right)$$

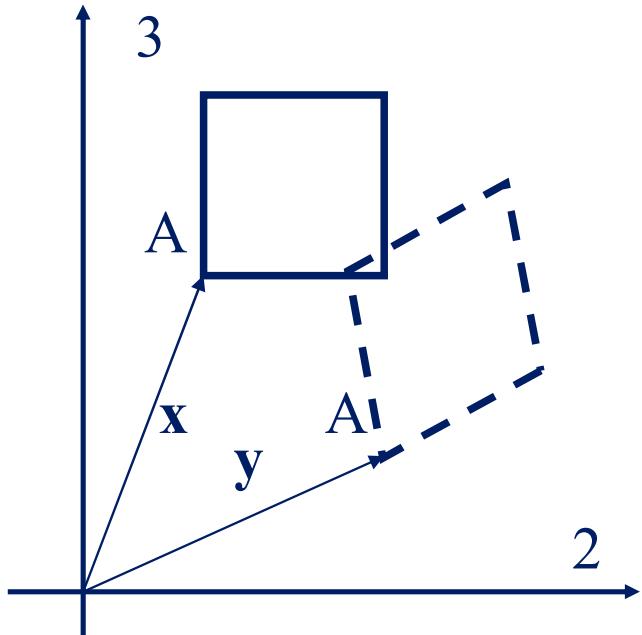
$$(\nabla_x \mathbf{u})_{ij} = \underline{\varepsilon_{ij}} + \omega_{ij}$$

$$\boxed{\varepsilon_{ij} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)}$$

Stress Tensor is also symmetric  
 $\therefore$  we consider only 6 components

$$\begin{pmatrix} \varepsilon_{11} \\ \varepsilon_{22} \\ \varepsilon_{33} \\ \varepsilon_{23} \\ \varepsilon_{13} \\ \varepsilon_{12} \end{pmatrix}$$

## Traslation, Rotation and Deformation



$$u = u_0 + \omega dx + \epsilon dx$$

Note: Remember initially when we introduced Elastic Tensor  
The total no. of components were 81 we reduced it to  
 $21 \Rightarrow 15 \Rightarrow$  finally considering the case of Cubic

like Crystal ex: Ge down to '3'

\* A Vector is equal to a Matrix multiplied by another vector.

## Generalized Hooke's Law (compliance matrix for isotropic materials)

$$\varepsilon_{hk} = \sum_{ij} b_{hk,ij} \sigma_{ij} \quad \text{red arrow}$$

$$\begin{pmatrix} \varepsilon_{11} \\ \varepsilon_{22} \\ \varepsilon_{33} \\ \varepsilon_{23} \\ \varepsilon_{13} \\ \varepsilon_{12} \end{pmatrix} = \begin{pmatrix} b_{11} & b_{12} & b_{12} & 0 & 0 & 0 \\ b_{12} & b_{11} & b_{12} & 0 & 0 & 0 \\ b_{12} & b_{12} & b_{11} & 0 & 0 & 0 \\ 0 & 0 & 0 & b_{44} & 0 & 0 \\ 0 & 0 & 0 & 0 & b_{44} & 0 \\ 0 & 0 & 0 & 0 & 0 & b_{44} \end{pmatrix} \begin{pmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{33} \\ \sigma_{23} \\ \sigma_{13} \\ \sigma_{12} \end{pmatrix}$$

Here, we have only Three independent components for Elastic Tensor 25

## Generalized Hooke's Law

- The inverse of the compliance matrix is called *stiffness matrix*.
- Elastic constants in isotropic materials:

$$b_{11} = 1 / E,$$

$$b_{12} = -v / E,$$

$$b_{44} = 1 / G$$

○  $E$  = Young's modulus

○  $v$  = Poisson's coefficient

○  $G$  = shear modulus

## The equations of linear elasticity

$$F_i^b = \int_V f_i^b dV$$

$$F_i^s = \int_S \sigma_{ij} n_j dS = \int_V \sum_{j=1}^3 \frac{\partial \sigma_{ij}}{\partial x_j} dV$$

$$\sum_{j=1}^3 \frac{\partial \sigma_{ij}}{\partial x_j} + f_i^b = \rho \dot{v}_i$$

# Mechanical Sensors

## Characterization of mechanical sensors

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1. Most relevant mechanical signals, their units in the S.I., and ranges for specific applications
  2. Definition of the stress and strain tensors
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Q)

How do we apply The Elastic Tensor concepts  
in Micro Electronics?

A)

Of course, we must be able to find Micro 28  
Electronic structures in which we can realize

Mechanical things that can realise Mechanical  
things - that can Bend (or) Deform.

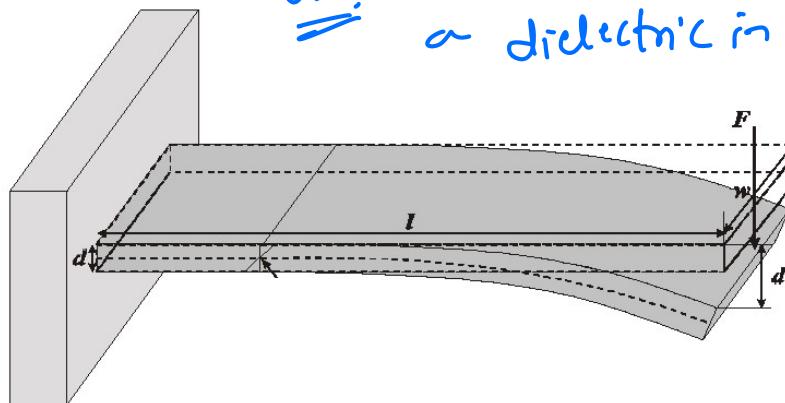
- (i) Cantilever Beams
- (ii) Membranes / Plates



## Mechanical structures used in MEMS sensors

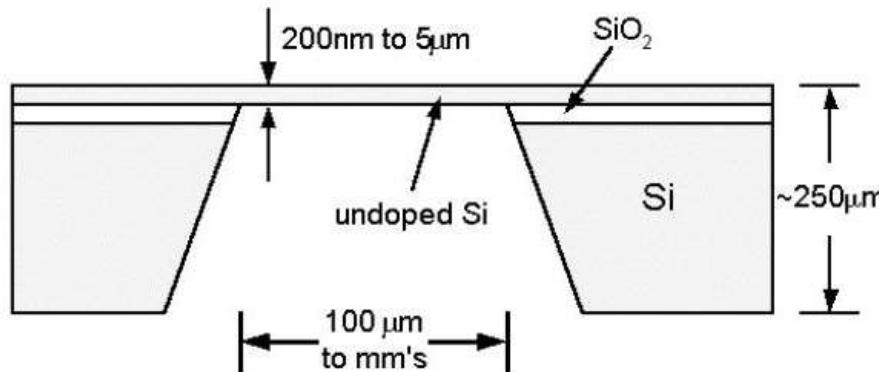
### (i) Cantilever beams

eg: cantilever beam used as  
a dielectric in capacitor; its slope  
change is reflected in  
its capacitance



### (ii) Membranes/Plates

Both these mechanical  
structures deform  
under a force





## Cantilever beam (fixed end beam)

Type of loads

- a. concentrated load (single force)
- b. distributed load (measured by their intensity) :
  - uniformly distributed load (uniform load)
  - linearly varying load

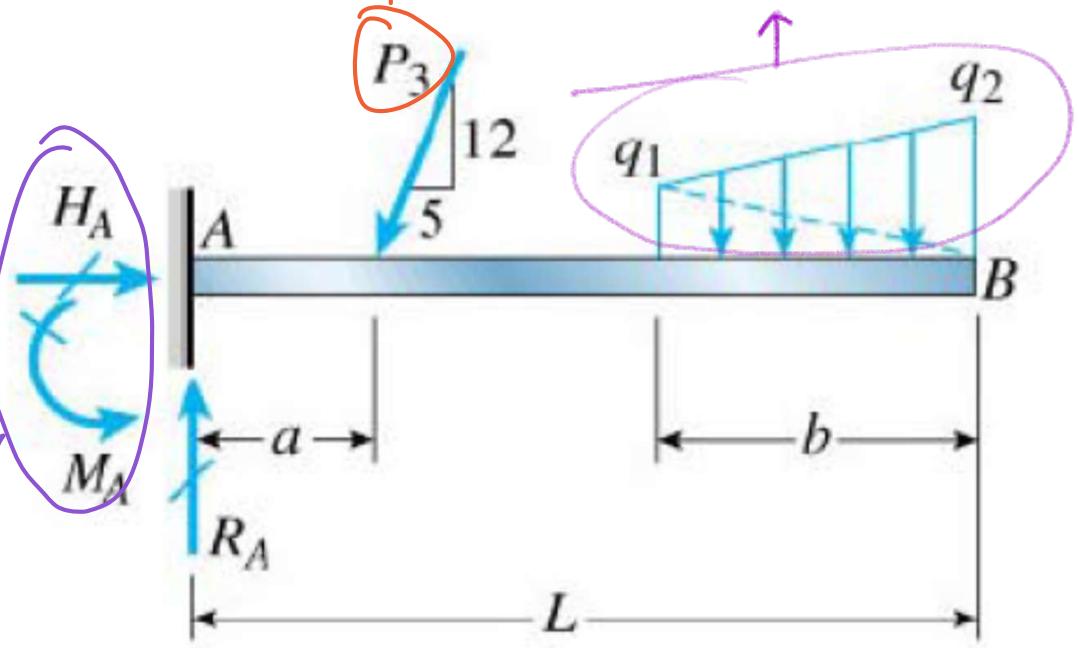
Reactions:

The fixed end prevents translations and rotations

No Normal Force  
and NO Rotation:  
‘A’ is fixed point

; it is a concentrated load

Distributed Load



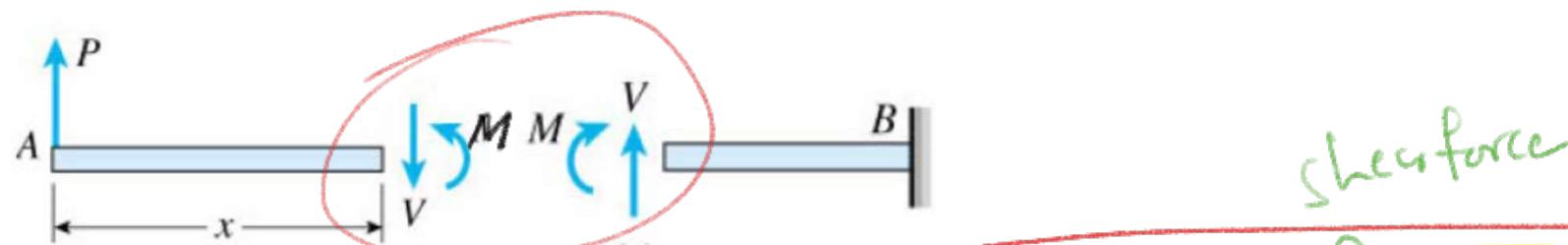
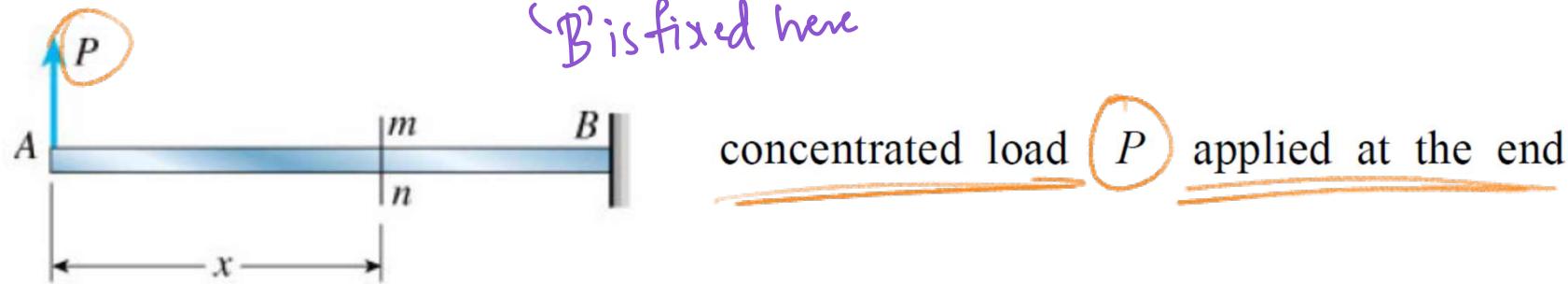
Same for  $R_A$ .

Q) How do we calculate the Shear forces & Moments inside a cantilever?

Any

## Cantilever beam (fixed end beam)

### Shear Forces and Bending Moments



To study what happens at 'x' distance from position 'A', we simply cut ideally the cantilever at position 'x' and to each end of separated we apply force

$$\sum F_y = 0$$

$$\sum M = 0$$

$$V = P$$

$$M = P x$$

A Moment & Reaction & Force & Moment compensate for the part we have removed.

\* Considering finite extension where  $\nu = \mu$  everywhere



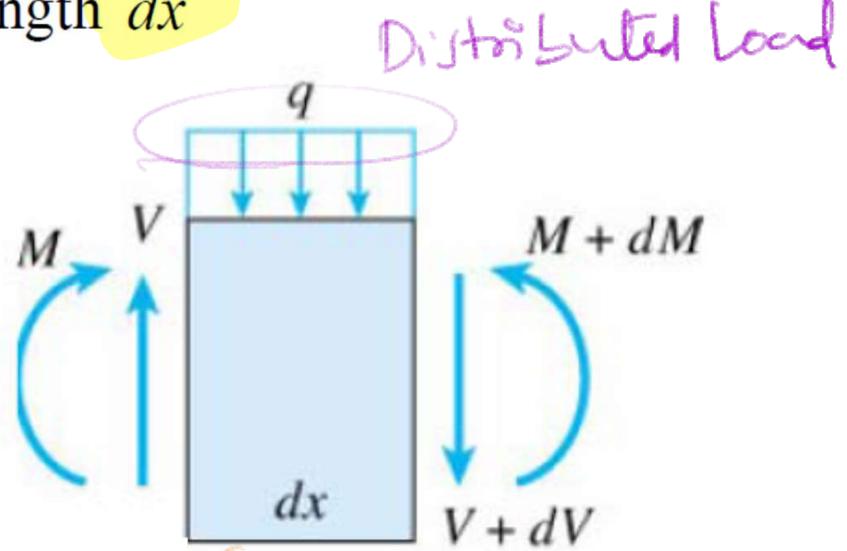
## Cantilever beam (fixed end beam)

### Relationships Between Loads, Shear Forces, and Bending Moments

consider an element of a beam of length  $dx$

subjected to distributed loads  $q$

Force  
length



$$\sum F_y = 0$$

or

$$V - q dx - (V + dV) = 0$$

$$dV/dx = -q$$

The extracted Block  
we consider to be in  
Equilibrium.

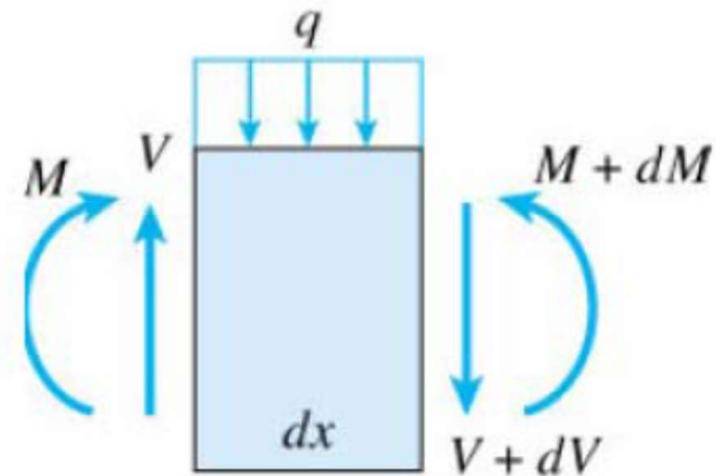
## Cantilever beam (fixed end beam)

### Relationships Between Loads, Shear Forces, and Bending Moments

consider an element of a beam of length  $dx$

subjected to distributed loads  $q$

we take the origin on the left face in  
 $v(0) = v(0) = 0$   
here



$$\Sigma M = 0 \quad -M - q dx (dx/2) - (V + dV) dx + M + dM = 0$$

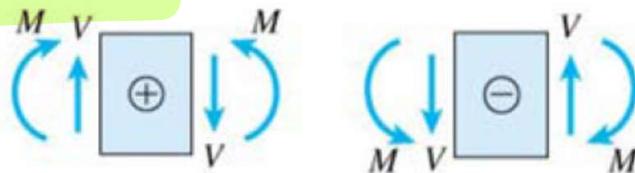
or

$$dM/dx = V$$

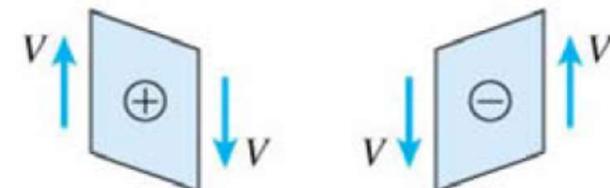
we eliminate the 2<sup>nd</sup> order terms.

$$\frac{d^2M}{dx^2} = \frac{dv}{dx} = -q$$

## Cantilever beam (fixed end beam) Deformations

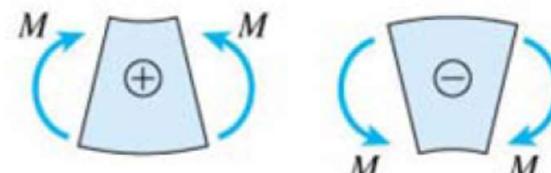


- the shear force tends to rotate the material clockwise is defined as positive



(a)

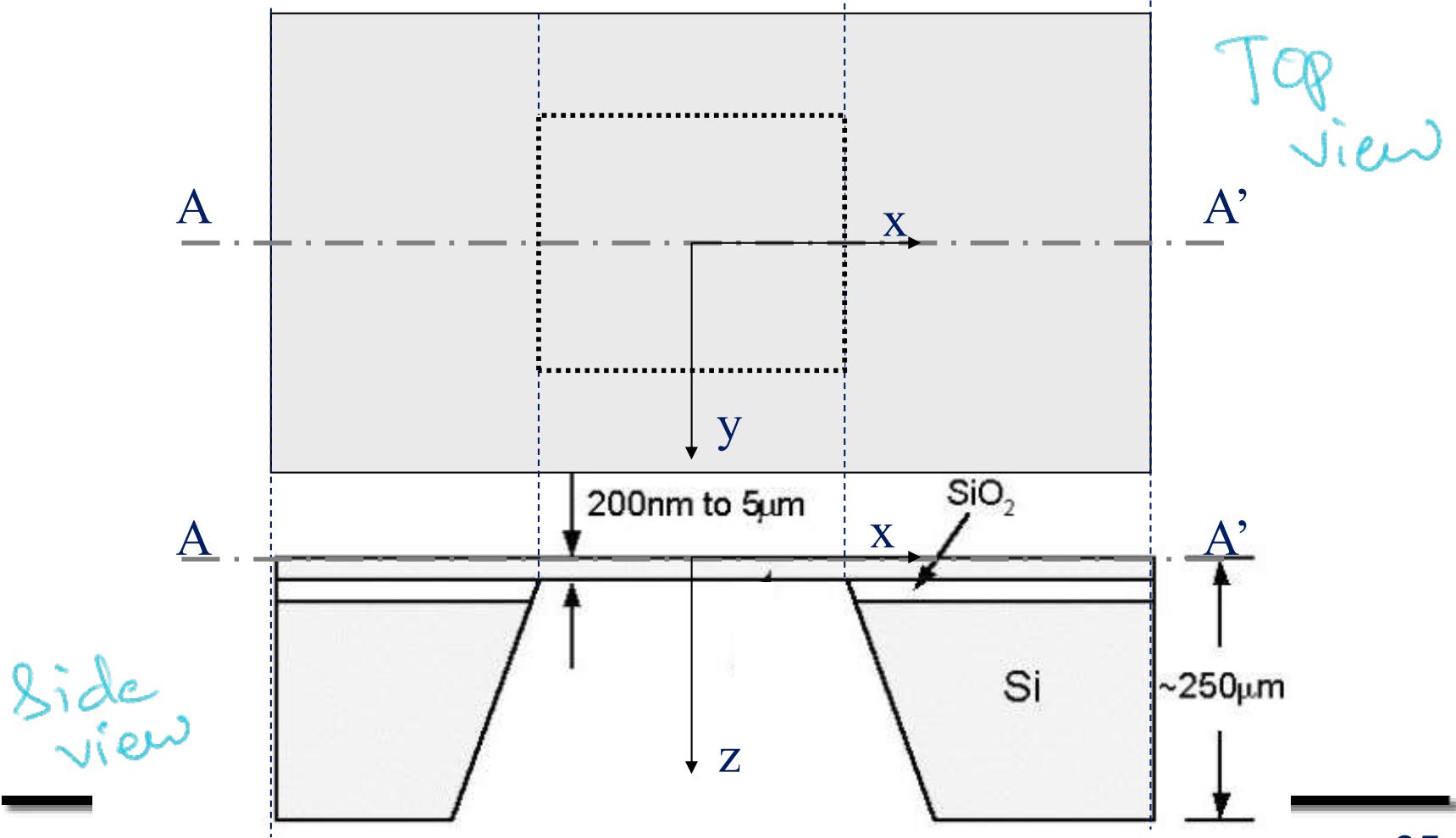
- the bending moment tends to compress the upper part of the beam and elongate the lower part is defined as positive



(b)



## MEMS plates in silicon



## Plates (basic concepts of plate bending)

- ④ A plate can be considered the two-dimensional extension of a beam in simple bending.

Both plates and beams support loads transverse or perpendicular to their plane and through bending action.

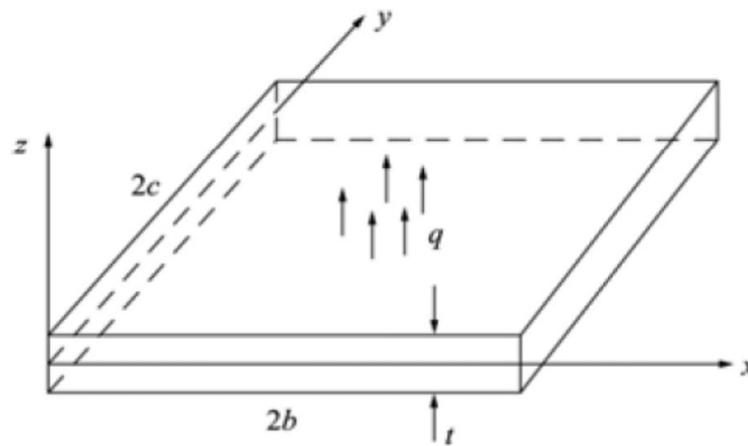
A plate is a flat (if it were curved, it would be a shell).

A beam has a single bending moment resistance, while a plate resists bending about two axes and has a twisting moment.

We will consider the classical thin-plate theory or Kirchhoff plate theory.

## Plates (basic concepts of plate bending)

Consider the thin plate in the  $x$ - $y$  plane of thickness  $t$  measured in the  $z$  direction shown in the figure below:



The plate surfaces are at  $z = \pm t/2$ , and its midsurface is at  $z = 0$ .

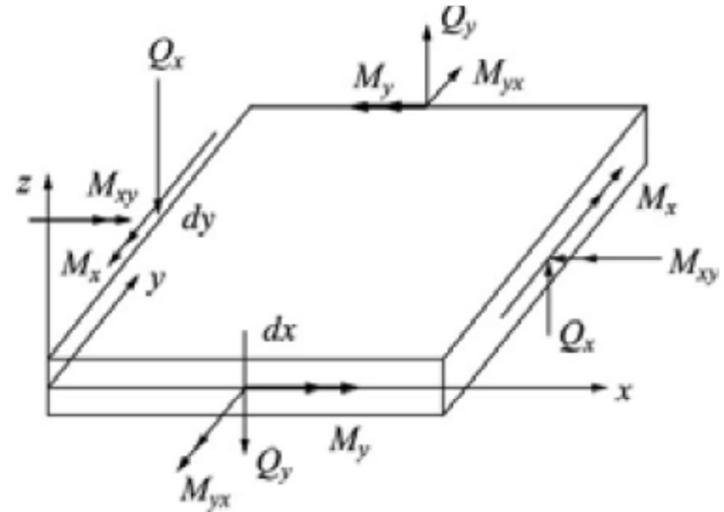
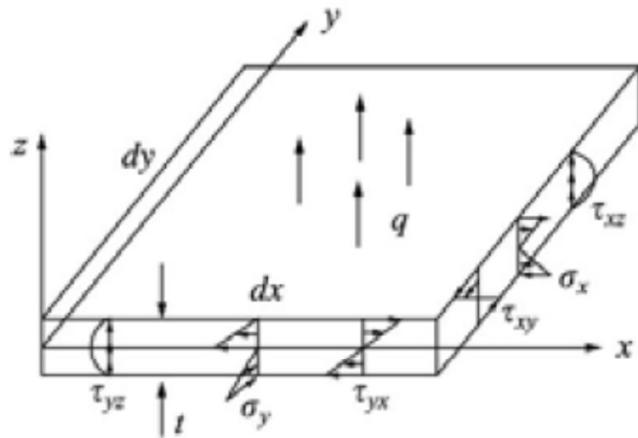
1. The plate thickness is much smaller than its inplane **dimensions**  $b$  and  $c$  (that is,  $t \ll b$  or  $c$ )

# Plates (basic concepts of plate bending)

## Kirchhoff Assumptions

1. Normals remain normal. This implies that transverse shears strains  $\gamma_{yz} = 0$  and  $\gamma_{xz} = 0$ . However  $\gamma_{xy}$  does not equal to 0.
2. Thickness changes can be neglected and normals undergo no extension. This means that  $\varepsilon_z = 0$ .
3. Normal stress  $\sigma_z$  has no effect on in-plane strains  $\varepsilon_x$  and  $\varepsilon_y$  in the stress-strain equations and is considered negligible.
4. Therefore, the in-plane deflections in the  $x$  and  $y$  directions at the midsurface,  $t = 0$ , are assumed to be zero;  $u(x, y, 0) = 0$  and  $v(x, y, 0) = 0$ .

# Plates (basic concepts of plate bending)



The governing differential equations are:

$$\frac{\partial Q_x}{\partial x} + \frac{\partial Q_y}{\partial y} + q = 0$$

$$\frac{\partial M_x}{\partial x} + \frac{\partial M_{xy}}{\partial y} - Q_x = 0$$

$$\frac{\partial M_y}{\partial y} + \frac{\partial M_{xy}}{\partial x} - Q_y = 0$$

where  $q$  is the transverse distributed loading and  $Q_x$  and  $Q_y$  are the transverse shear line loads.

## Solution of the plate bending problem

(S.K.Clark & K.D. Wise, IEEE Tr. ED 26, 1979)

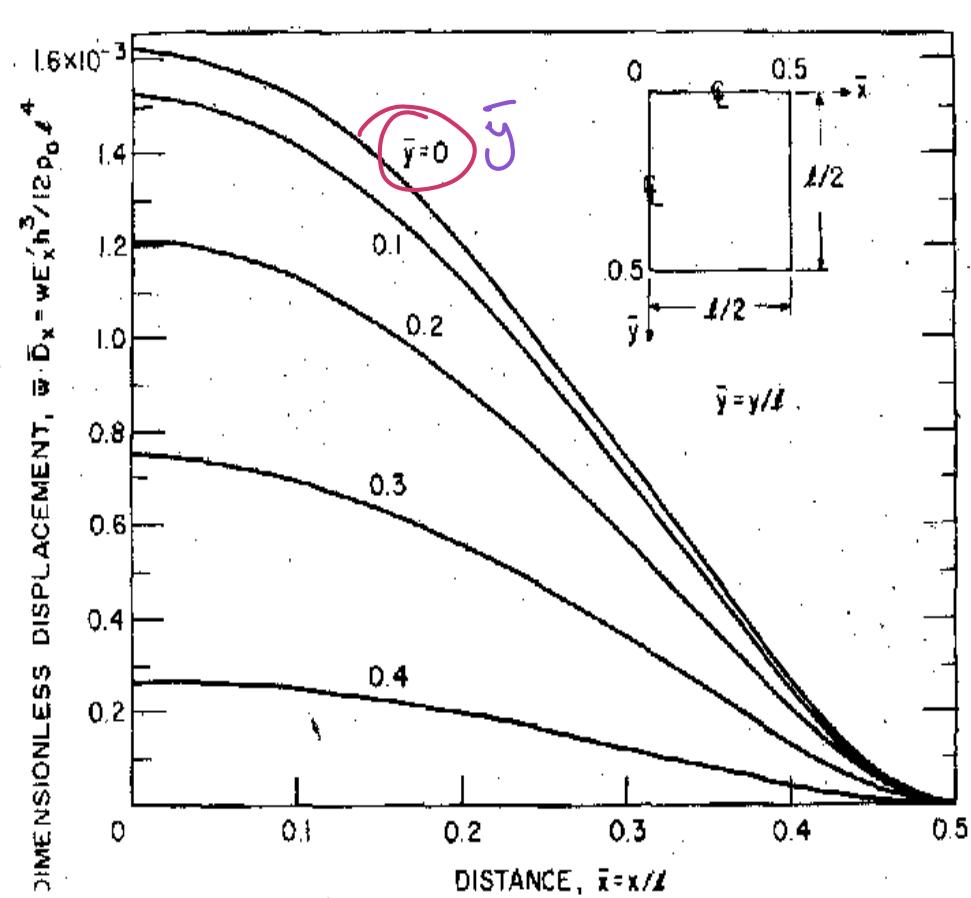
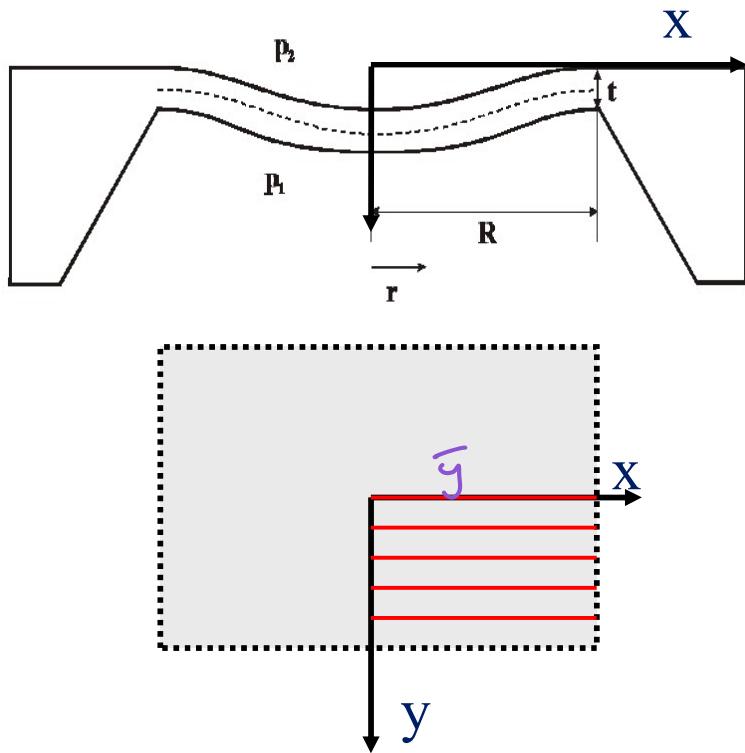


Fig. 4. Dimensionless displacement of a square silicon diaphragm having built-in edges as a function of position on the diaphragm.

# Solution of the plate bending problem

(S.K.Clark & K.D. Wise, IEEE Tr. ED 26, 1979)

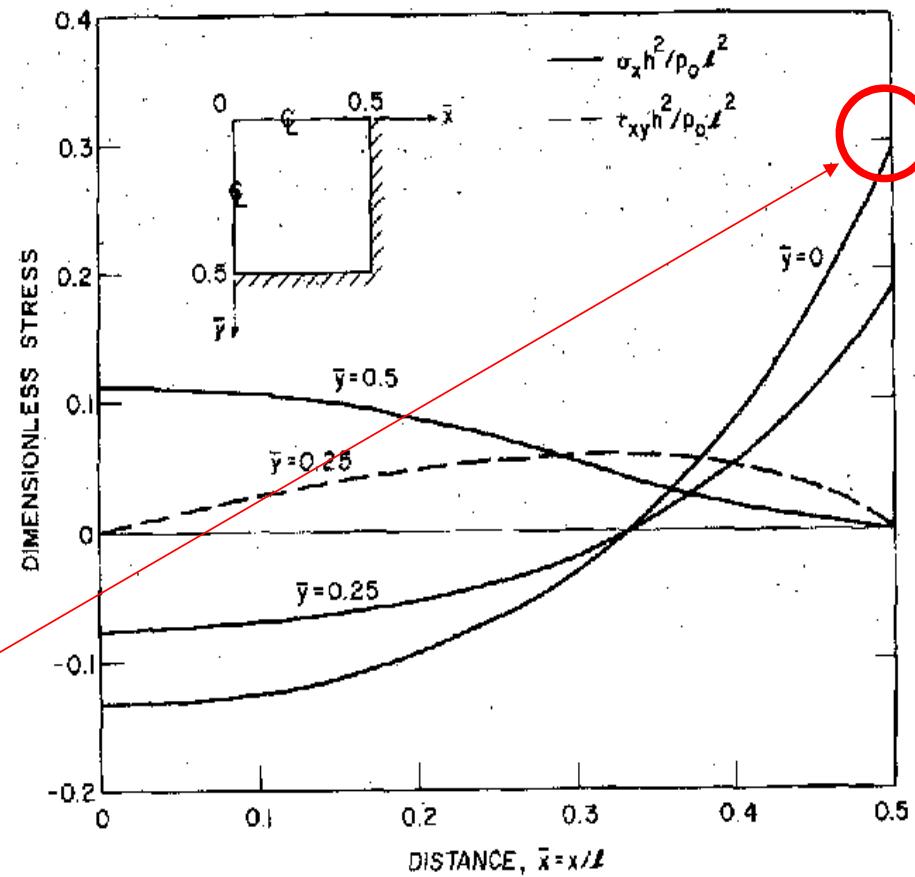
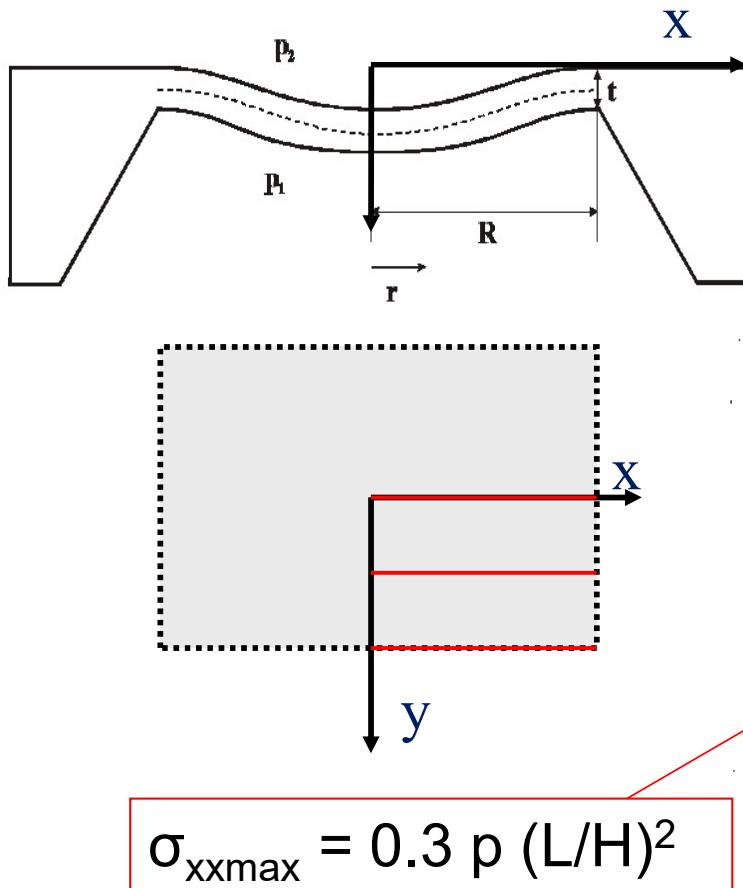


Fig. 5. Dimensionless stress distributions on a silicon diaphragm having built-in edges.

# Solution of the plate bending problem

(S.K.Clark & K.D. Wise, IEEE Tr. ED 26, 1979)

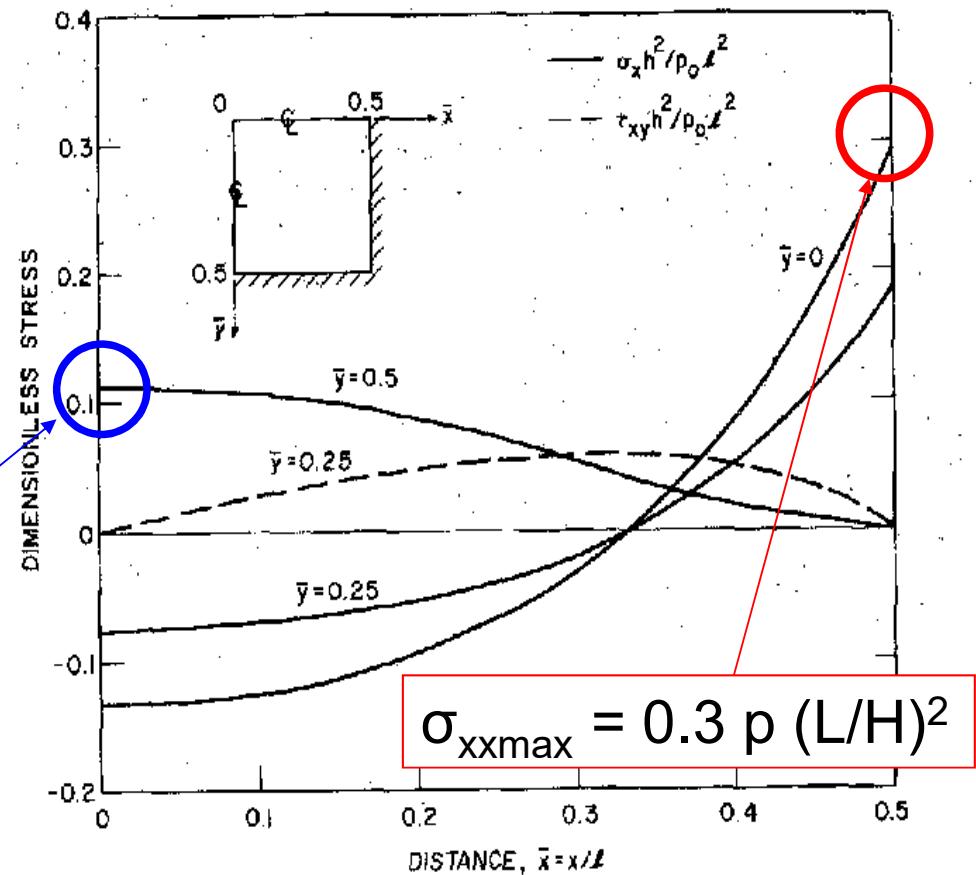
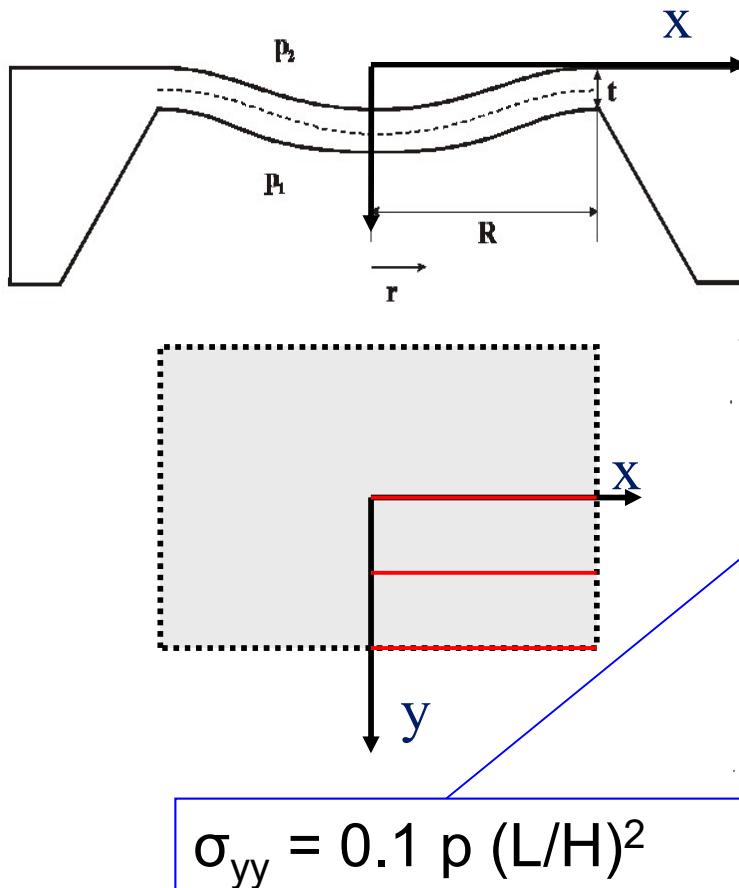


Fig. 5. Dimensionless stress distributions on a silicon diaphragm having built-in edges.



# Advanced Solid-State Sensors M

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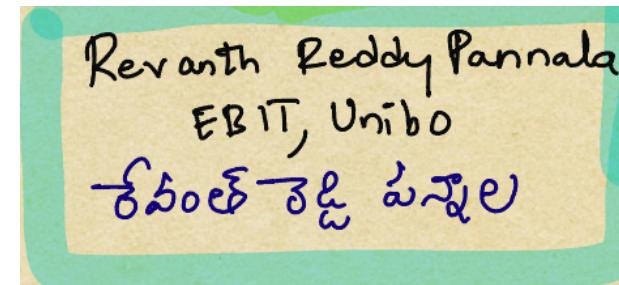
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Microelectronics Laboratory (main building - third floor)

# Course Program 2019-2020

1. Introduction – Classification of sensors
2. Main physical effects used in solid-state sensors and simple readout circuits
3. Resolution & Noise
4. Image sensors
5. Mechanical sensors



# Mechanical Sensors

Characterization of mechanical sensors

- MEMS sensors and the different transduction approaches:
  1. Capacitive accelerometers
  2. Piezoelectric sensors
  3. Piezoresistive pressure sensors

22|04|25

## Mechanical Sensors

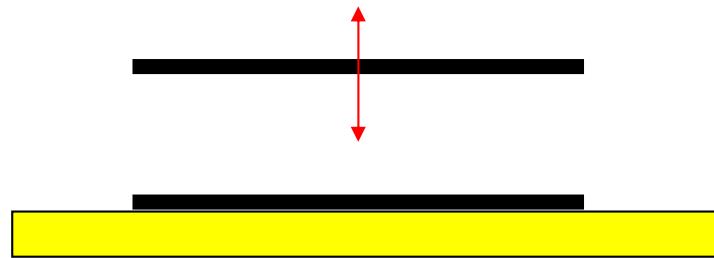
Characterization of mechanical sensors

- MEMS sensors and the different transduction approaches:
  1. Capacitive accelerometers
  2. Piezoelectric sensors
  3. Piezoresistive pressure sensors



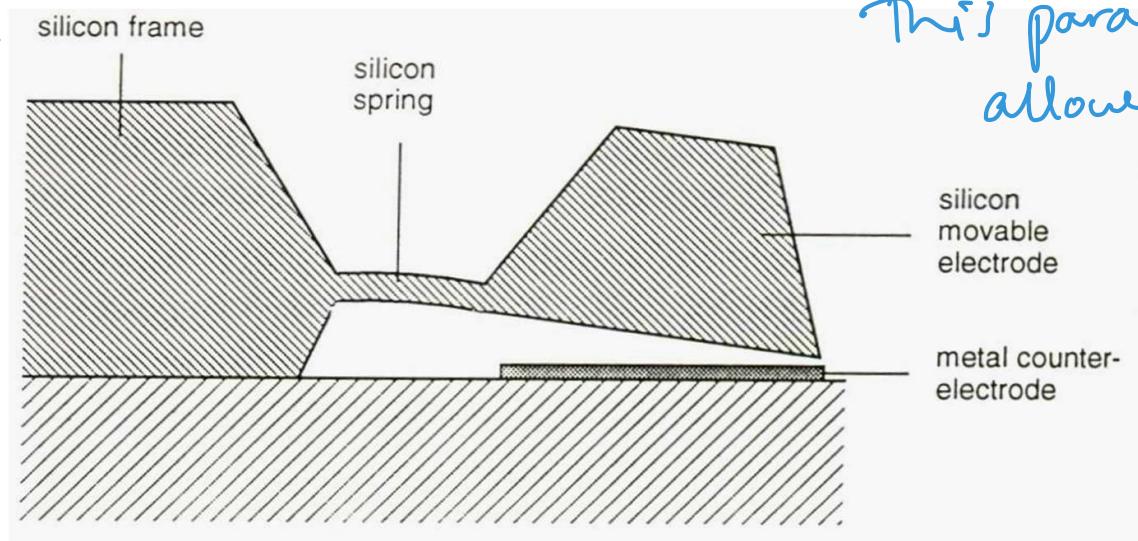
## Capacitive sensors

expression of parallel plate capacitor



$$C_0 = \epsilon_0 \epsilon \frac{A}{d} = \epsilon_A \frac{1}{\bar{d}}$$

\* Example of capacitor in which one of the plates moves



This parameter is allowed to change

Why don't we use THREE PLATES!

Differential capacitance system

upper & lower plates are fixed.



$C_1$

$C_2$



Movable plate



$$C_1 = \epsilon_A \frac{1}{x_1} = \epsilon_A \frac{1}{d+x} = C_0 - \Delta C,$$

$$C_2 = \epsilon_A \frac{1}{x_2} = \epsilon_A \frac{1}{d-x} = C_0 + \Delta C.$$

$$C_2 - C_1 = 2\Delta C = 2\epsilon_A \frac{x}{d^2 - x^2}$$

$$\Delta C x^2 + \epsilon_A x - \Delta C d^2 = 0$$

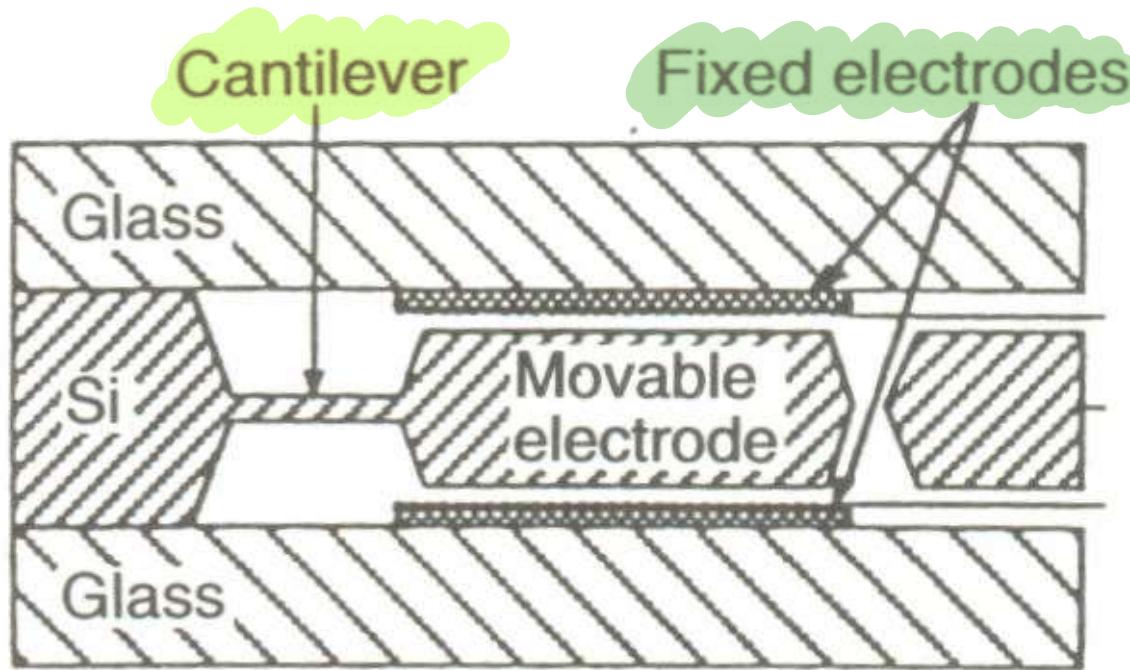
$\Delta C x^2$  is negligible

$$x \approx \frac{d^2}{\epsilon_A} \Delta C = d \frac{\Delta C}{C_0}$$

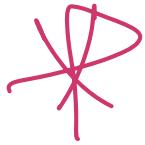
Conclusion:

By measuring  $\underline{\Delta C}$  we measure  
 $x$

## Differential capacitance system: vertical structure

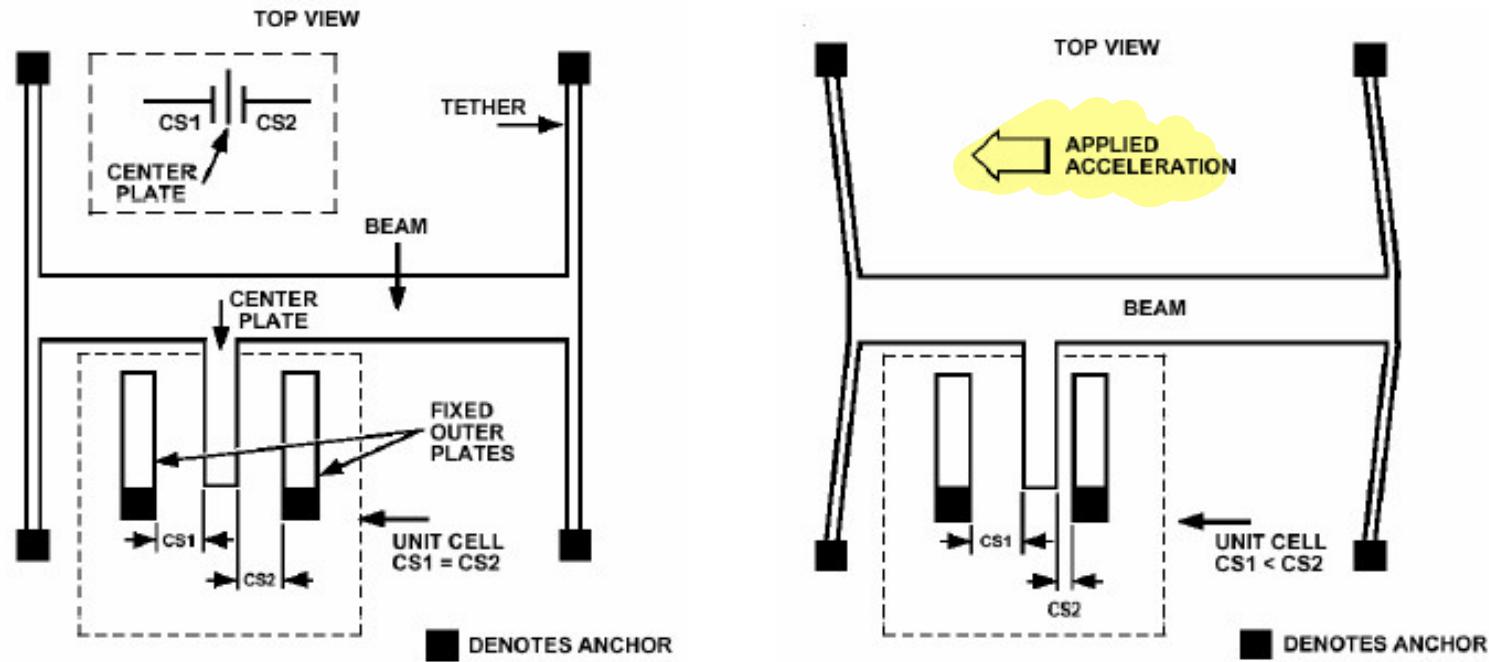


(HITACHI)

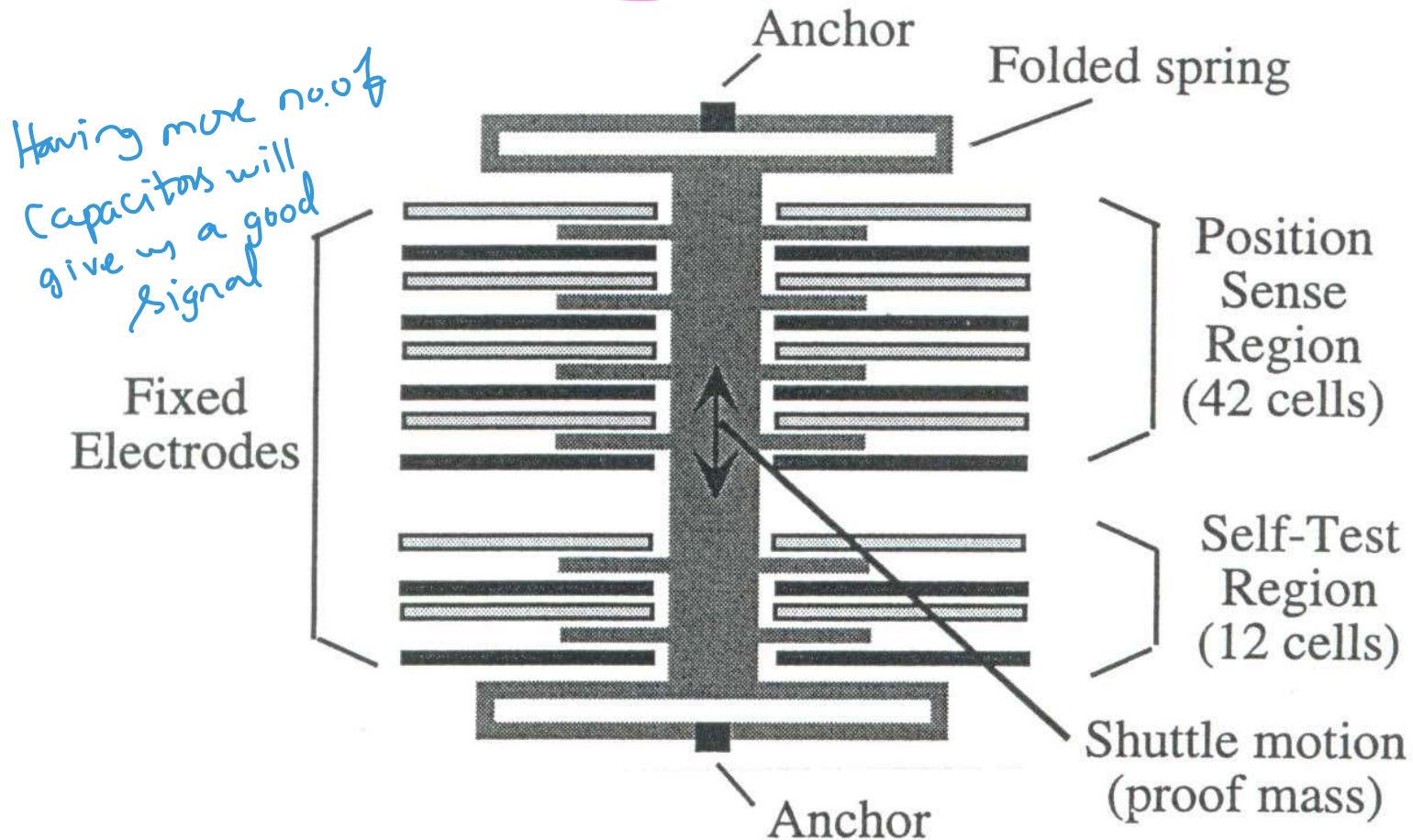


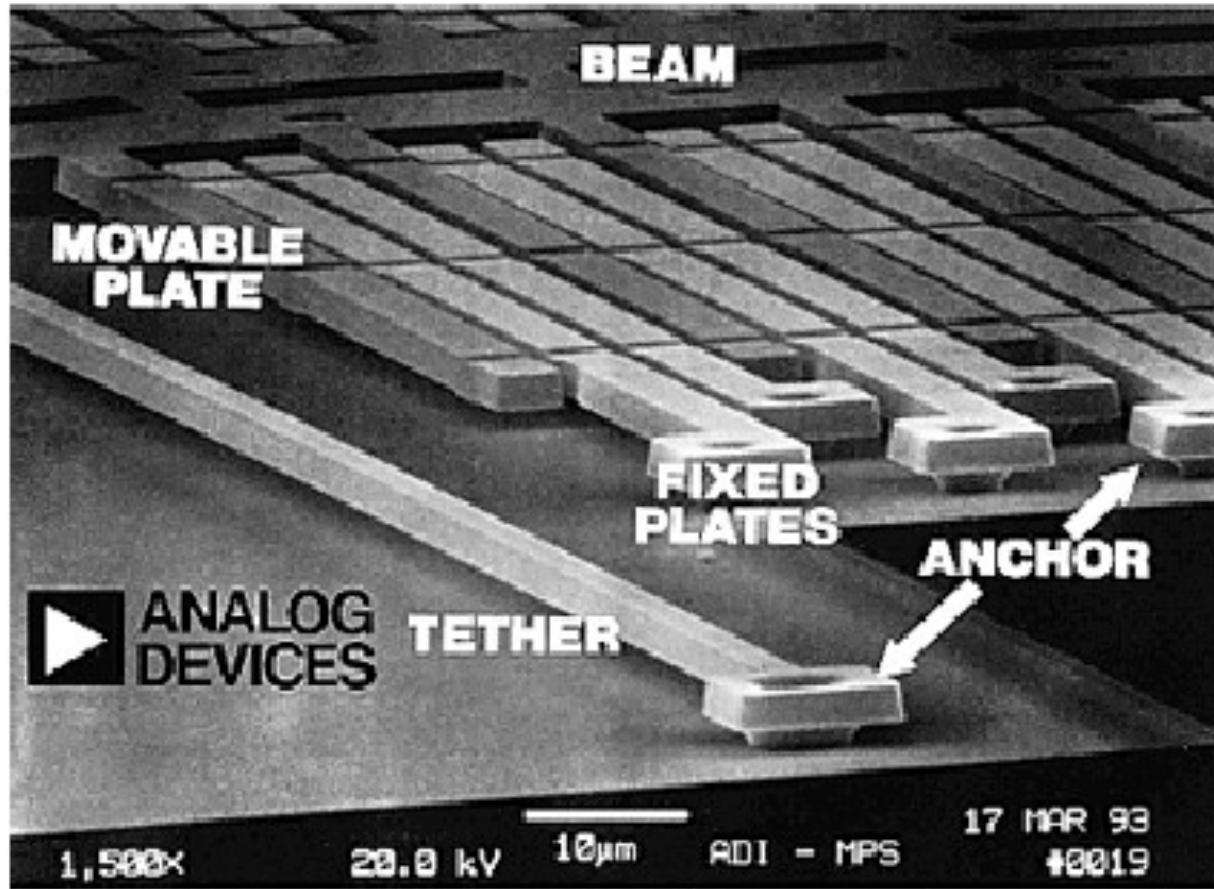
## Differential capacitance system: lateral structure

### ADXL Differential Capacitive Sensor



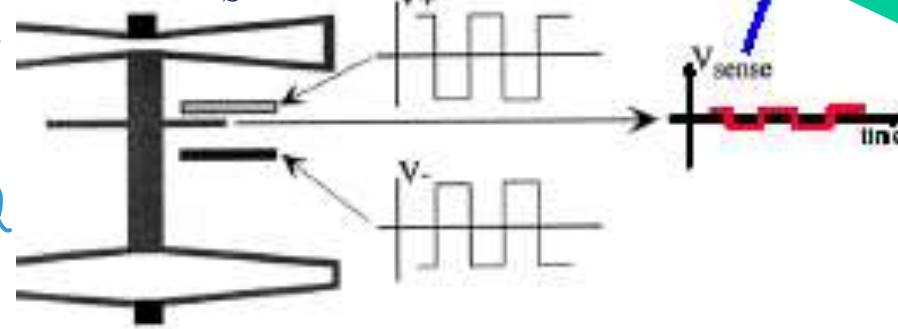
## ANALOG DEVICES ADXL



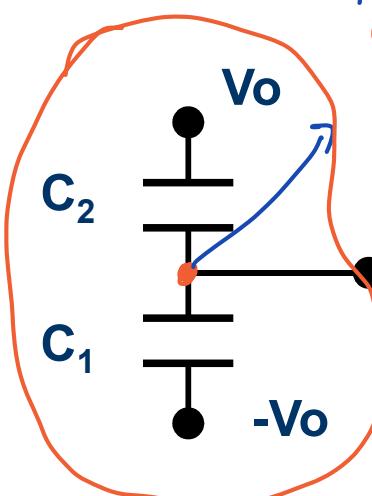


## READOUT CIRCUIT (1)

Spring,  $k_s$



Total charge  $Q$   
Globally cannot change because we are measuring with ' $\infty$ ' impedance



$V_x$  Voltage of middle plate

$k_s$  is the spring constant  $\rightarrow F_s = k_s x$

$$(V_x + V_0)C_1 + (V_x - V_0)C_2 = 0$$

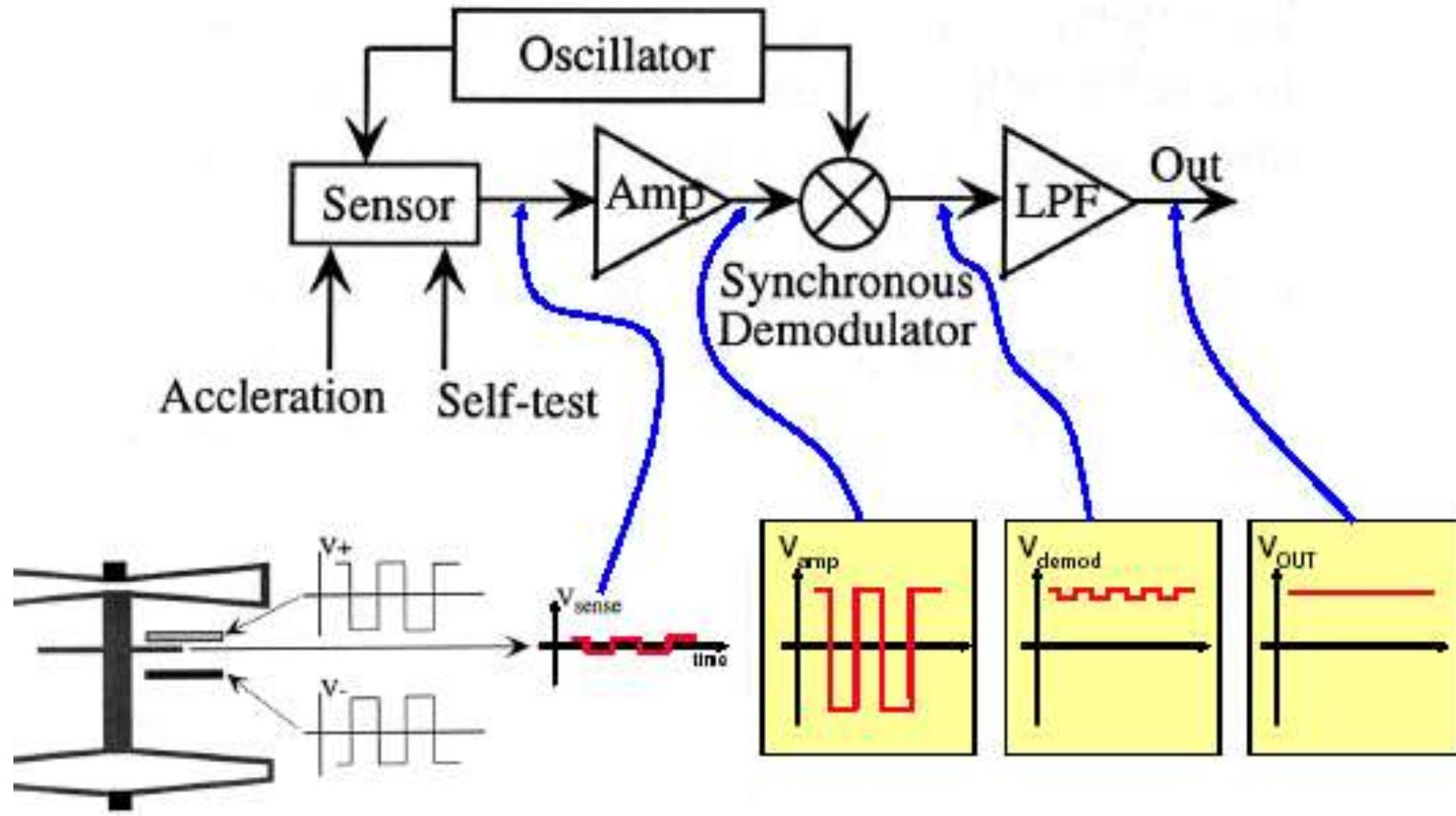
$$V_x = V_0 \frac{C_2 - C_1}{C_2 + C_1} = \frac{x}{d} V_0$$

$$a = \frac{k_s}{m} x$$

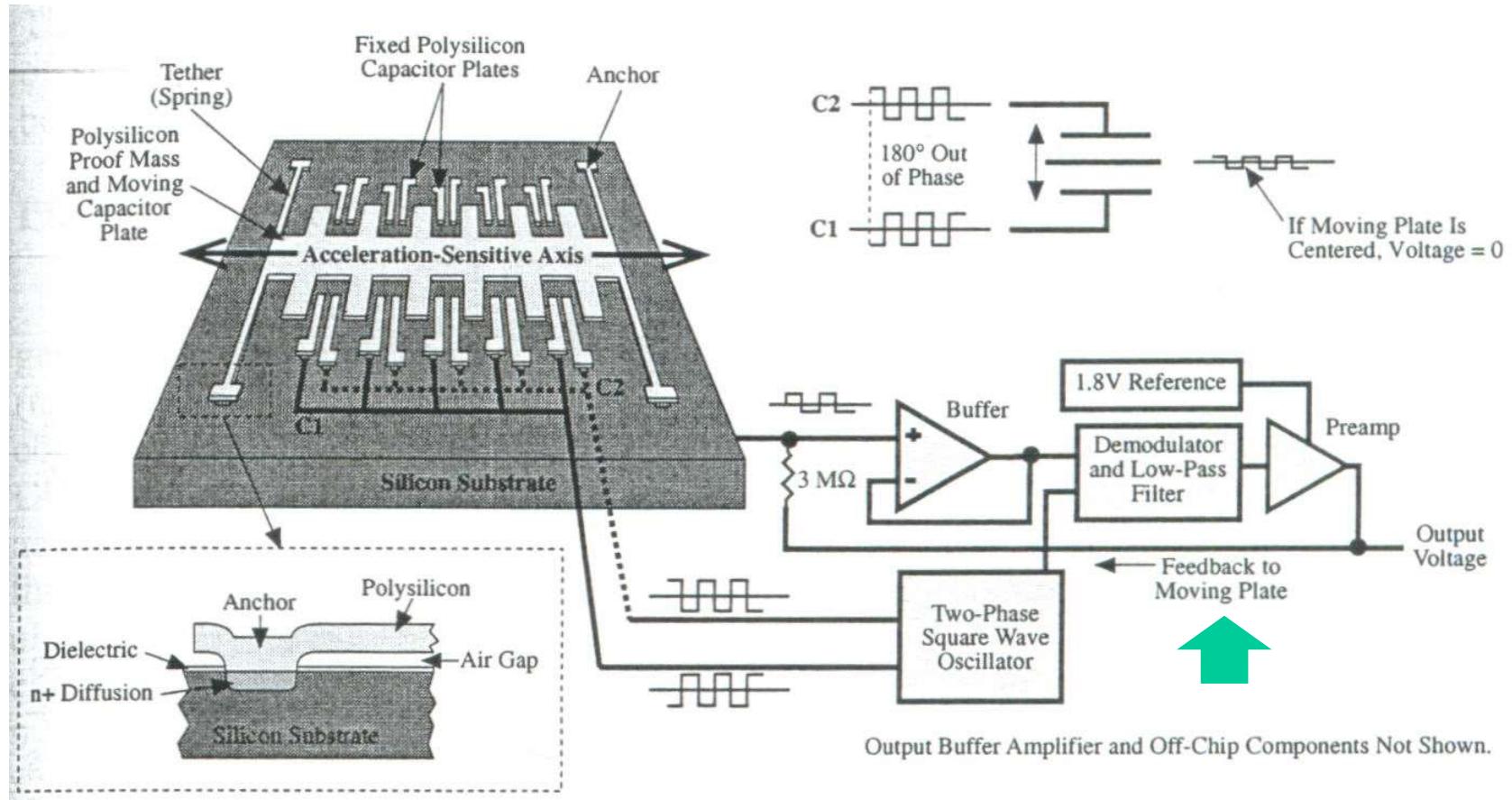
$$a = \frac{k_s d}{m V_0} V_x$$

This is how Accelerometers work!

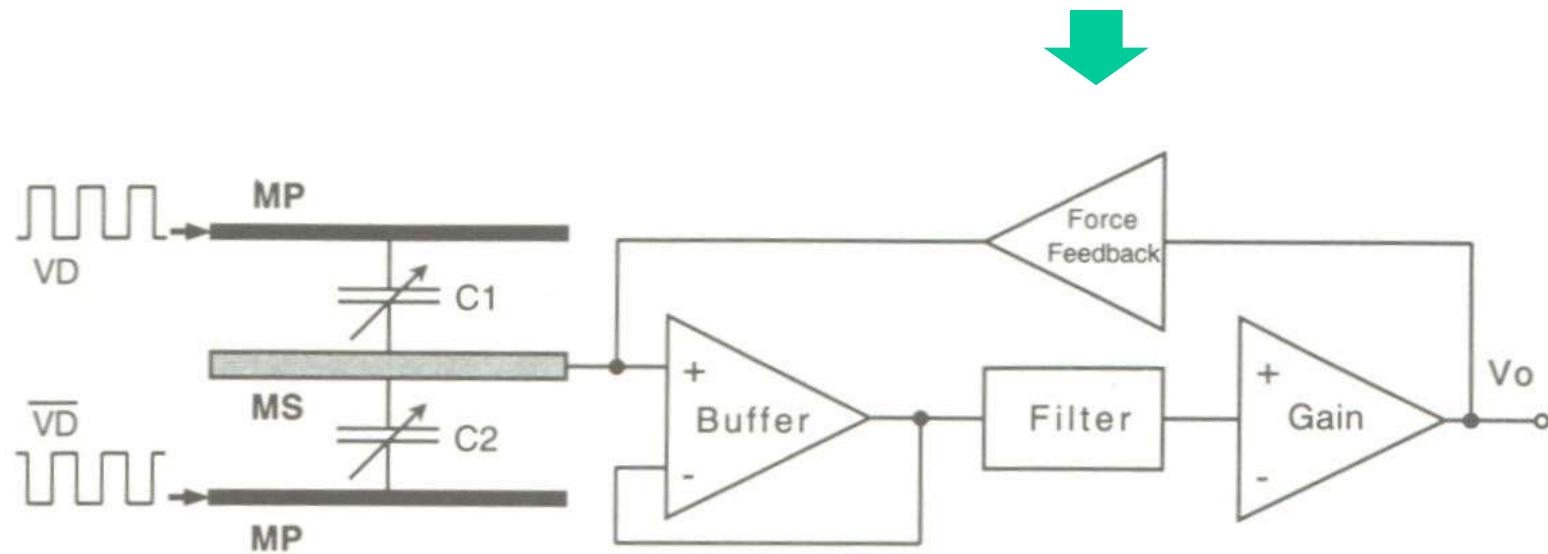
## Readout circuit (1)



## (2) Readout with feedback



(2) Readout with feedback: closed-loop technique with a force-balancing effect



# Noise: «Measurements of noise characteristics of MEMS accelerometers» by F. Mohd-Yasin et al., George Washington University, Washington D.C., Solid-State Electronics 2003.

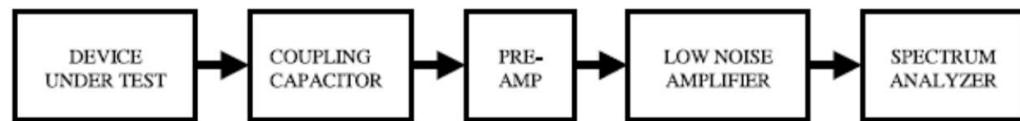


Fig. 2. Noise measurement system for MEMS accelerometers.

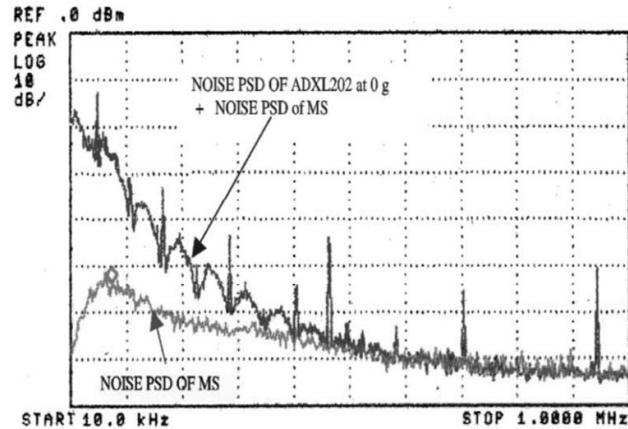


Fig. 3. ADXL202 noise PSD.

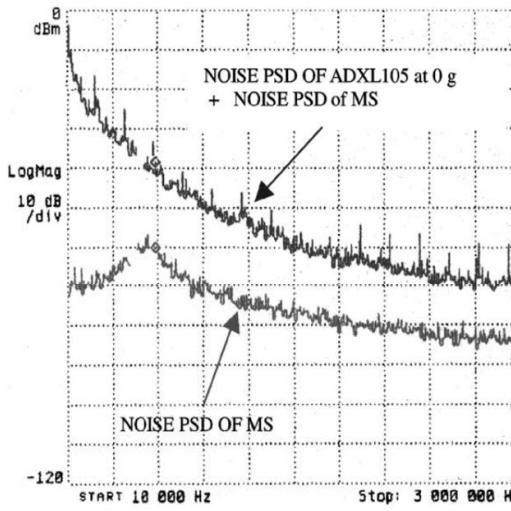


Fig. 4. ADXL105 noise PSD.

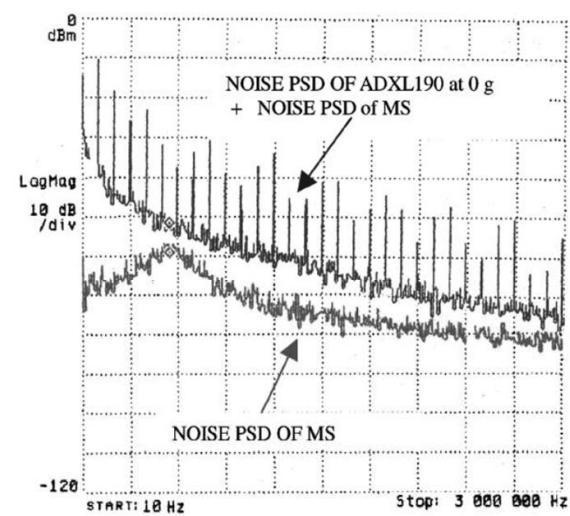


Fig. 5. ADXL190 noise PSD.

# Mechanical Sensors

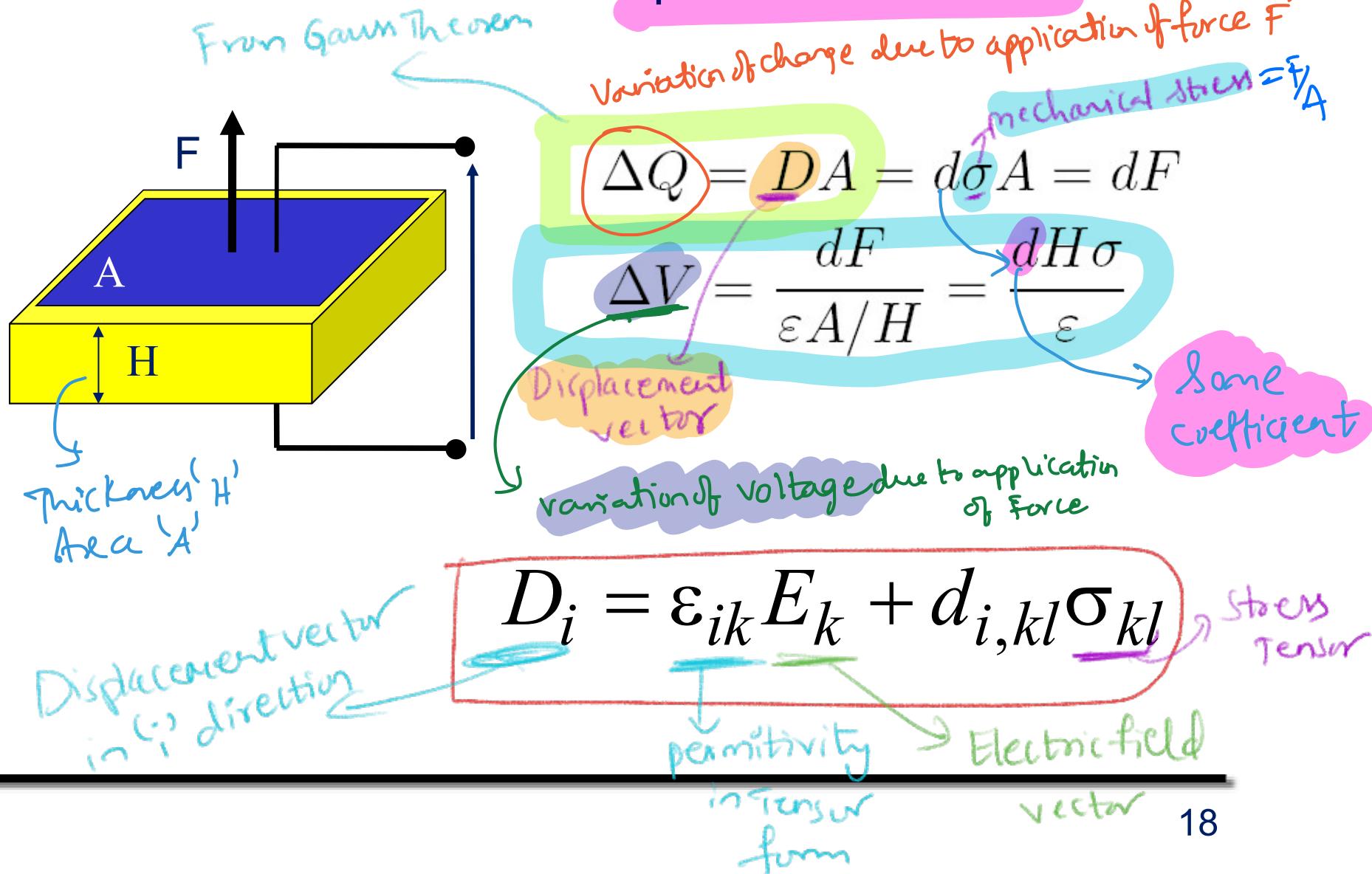
Characterization of mechanical sensors

- MEMS sensors and the different transduction approaches:
  1. Capacitive accelerometers
  2. Piezoelectric sensors
  3. Piezoresistive pressure sensors

## Piezoelectric sensors:

- Generalization of the piezoelectric effect: piezoelectric tensor
- Accelerometer PI-FET proposed by Chen et al. in 1980
- Piezoelectric oxide semiconductor field effect transistor touch sensing devices proposed by R. Dahiya et al. in 2009

# Generalization of the piezoelectric effect



## Generalization of the piezoelectric effect

$$D_i = \epsilon_{ik} E_k + d_{i,kl} \sigma_{kl}$$

$$\epsilon = \begin{bmatrix} \epsilon_{11} & 0 & 0 \\ 0 & \epsilon_{22} & 0 \\ 0 & 0 & \epsilon_{33} \end{bmatrix}$$

	$\epsilon_{11} = \epsilon_{22}$	$\epsilon_{33}$
$\text{SiO}_2$	$4.5 \epsilon_0$	$4.6 \epsilon_0$
$\text{BaTiO}_3$	$1268 \epsilon_0$	$1419 \epsilon_0$
$\text{ZnO}$	$10.8-11 \epsilon_0$	$10.8-11 \epsilon_0$

## Generalization of the piezoelectric effect

$$D_i = \epsilon_{ik} E_k + d_{i,kl} \sigma_{kl}$$

$$\sigma_{kl} \rightarrow \sigma_m, \quad m=1, \dots, 6 \rightarrow d_{i,kl} \rightarrow d_{im}$$

6 independent components

$$d = \begin{bmatrix} d_{11} & d_{12} & d_{13} & d_{14} & d_{15} & d_{16} \\ d_{21} & d_{22} & d_{23} & d_{24} & d_{25} & d_{26} \\ d_{31} & d_{32} & d_{33} & d_{34} & d_{35} & d_{36} \end{bmatrix}$$

# Generalization of the piezoelectric effect

$$\mathbf{d} = \begin{bmatrix} d_{11} & d_{11} & 0 & d_{14} & 0 & 0 \\ 0 & 0 & 0 & 0 & -d_{14} & -2d_{11} \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix} \quad \text{SiO}_2$$

$d_{11} = -2.31 \text{ pC/N}$

$d_{14} = -0.727 \text{ pC/N}$

$$\mathbf{d} = \begin{bmatrix} 0 & 0 & 0 & 0 & d_{15} & 0 \\ 0 & 0 & 0 & d_{15} & 0 & 0 \\ d_{31} & d_{31} & d_{33} & 0 & 0 & 0 \end{bmatrix} \quad \text{BaTiO}_3$$

$d_{31} = -79 \text{ pC/N}$

$d_{33} = 191 \text{ pC/N}$

$d_{15} = 270 \text{ pC/N}$

$$\mathbf{d} = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & d_{15} & 0 & 0 \\ d_{31} & d_{31} & d_{33} & 0 & 0 & 0 \end{bmatrix} \quad \text{ZnO}$$

$d_{31} = -5.1 \text{ pC/N}$

$d_{33} = 12.3 \text{ pC/N}$

$d_{15} = -8.3 \text{ pC/N}$

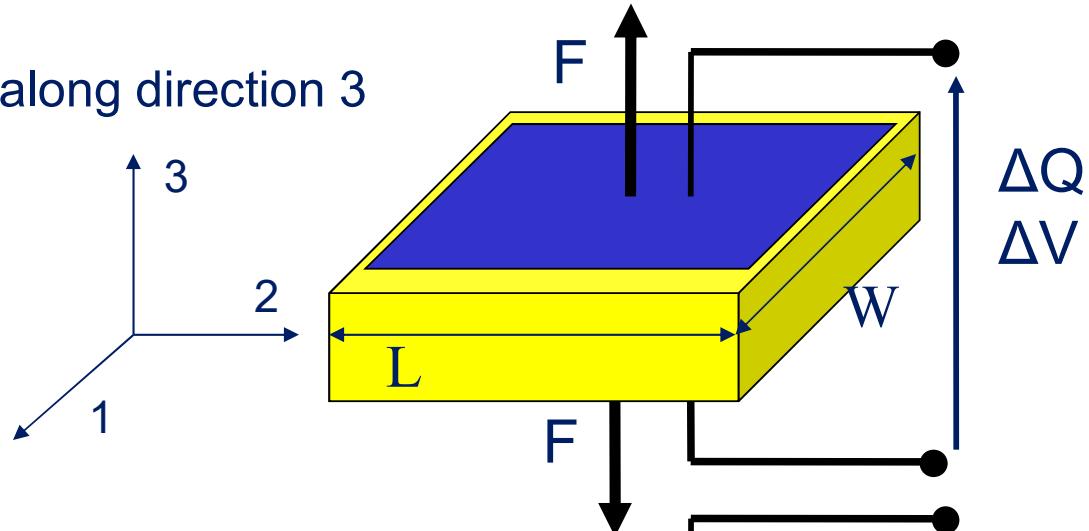
Q) How do we measure the coefficients?

## Characterization of the piezoelectric coefficients

### 1) Longitudinal coefficient

Example: ZnO measured along direction 3  
shown in figure

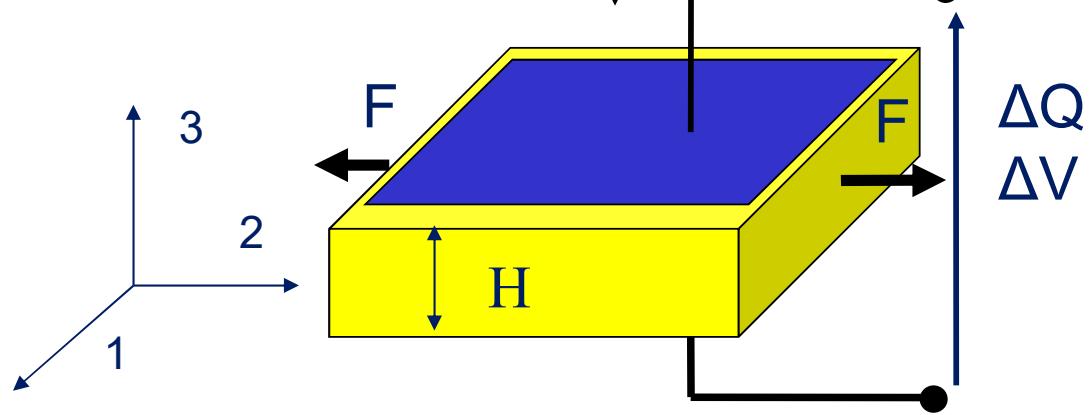
$$D_3 = d_{33} F/(WL)$$



### 2) Transverse coefficient

Example: ZnO measured  
along direction 3  
shown in figure

$$D_3 = d_{31} F/(HW)$$

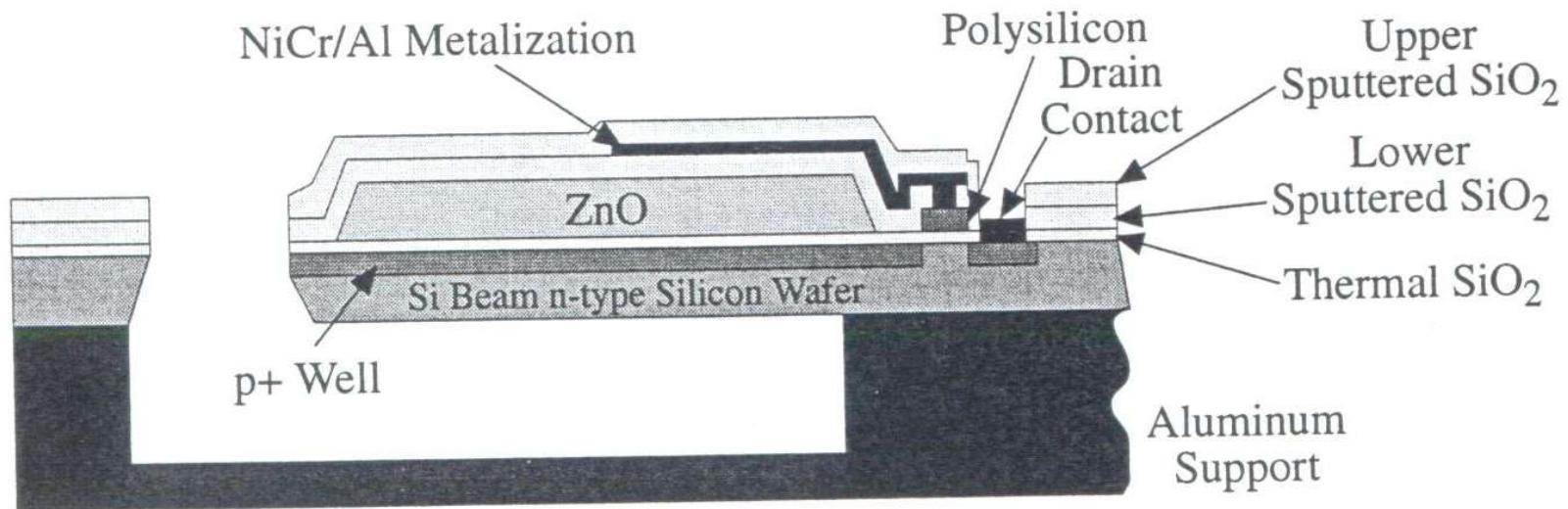


Q) How can we exploit the piezoelectric effect?

## Integrated silicon microbeam PI-FET accelerometer

P. Chen et al., IEEE Tr. ED 29, 1982

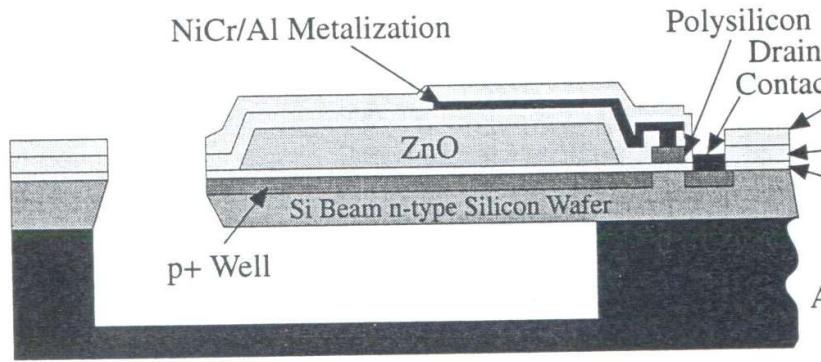
Any



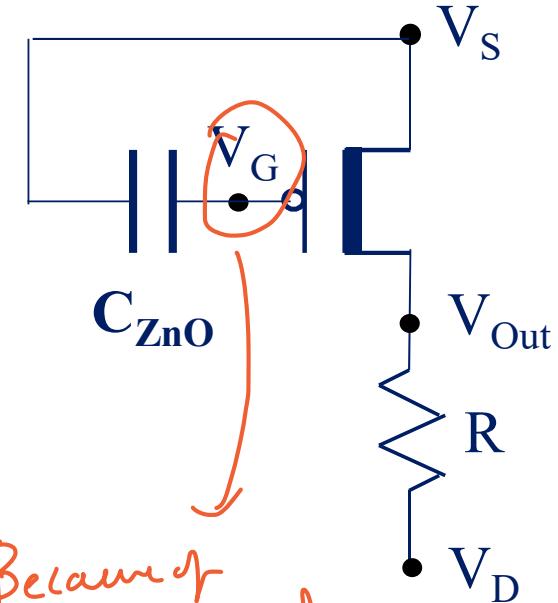
The structure may become a cantilever

## Readout circuit:

1-stage amplifier (depletion p-channel MOSFET)



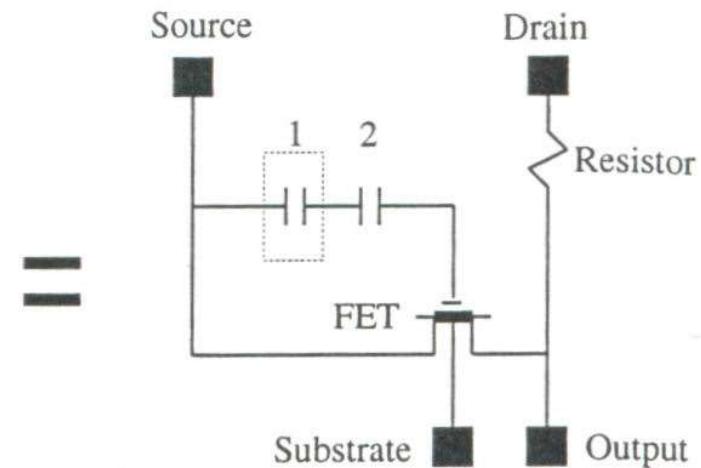
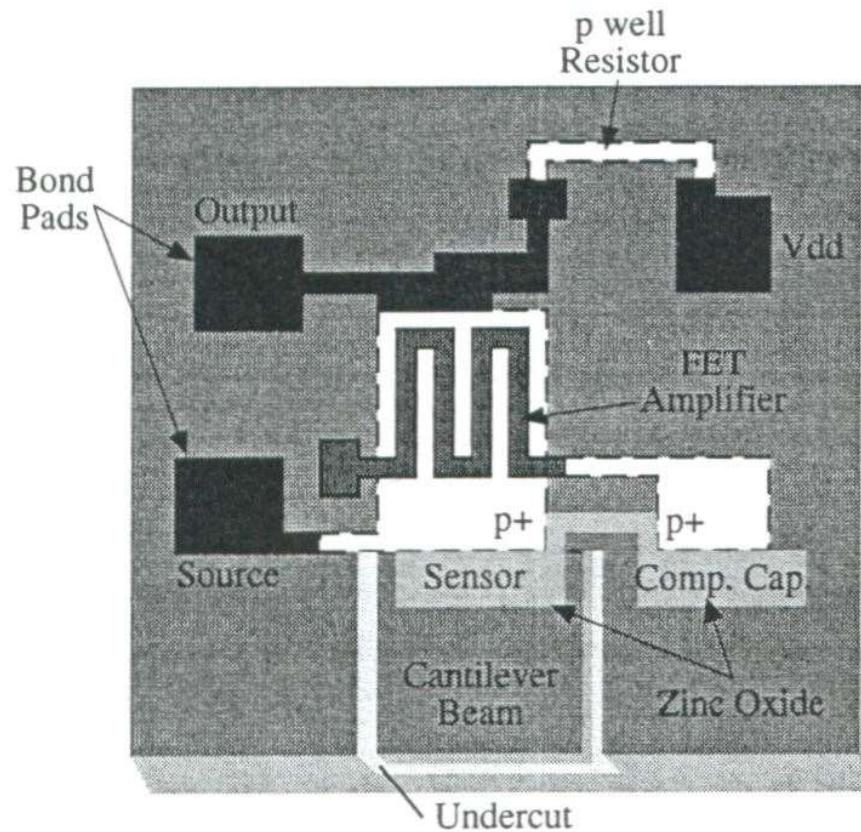
$$\text{Voltage gain: } A_v = -R g_{mp}$$



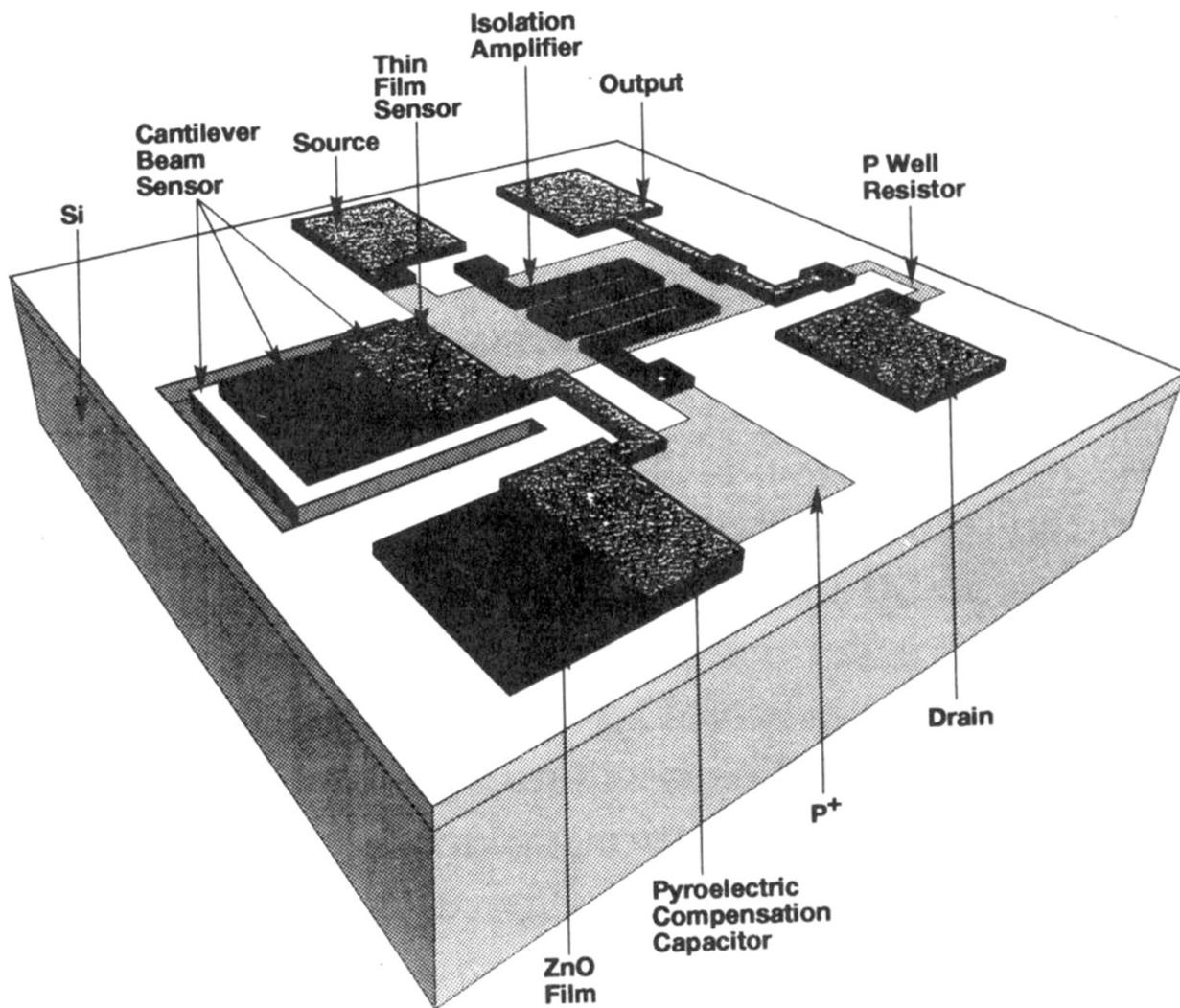
The force applied  
on the ZnO layer  
the capacitance vary  
thus by varying  $V_G$

## Readout circuit:

### 2. Pyroelectric compensation capacitor



Piezoelectric thin film capacitors:  
1 - Strained capacitor on beam  
2 - Unstrained compensation capacitor



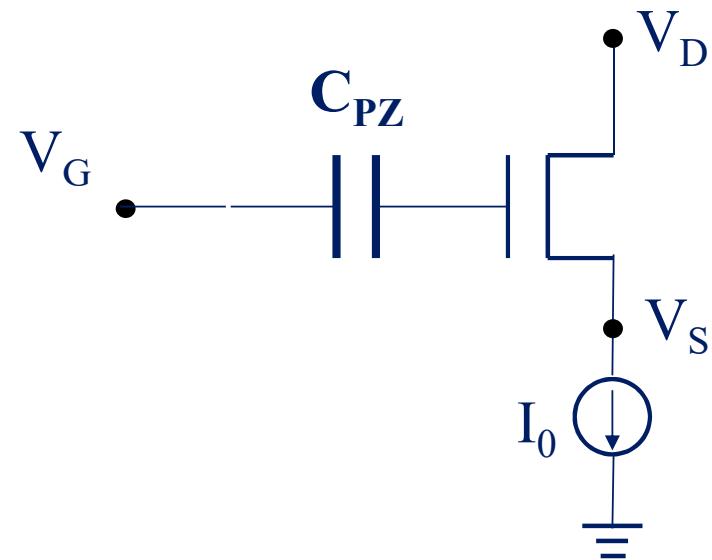
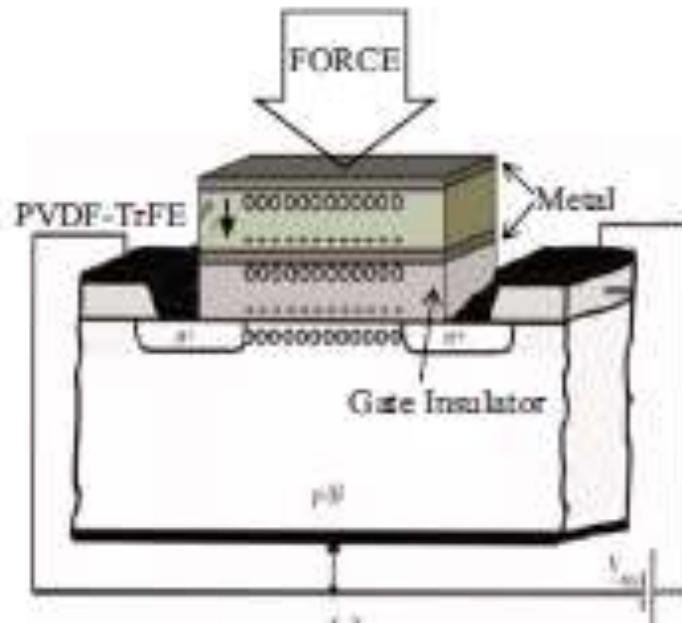
# Piezoelectric oxide semiconductor field effect transistor touch sensing devices

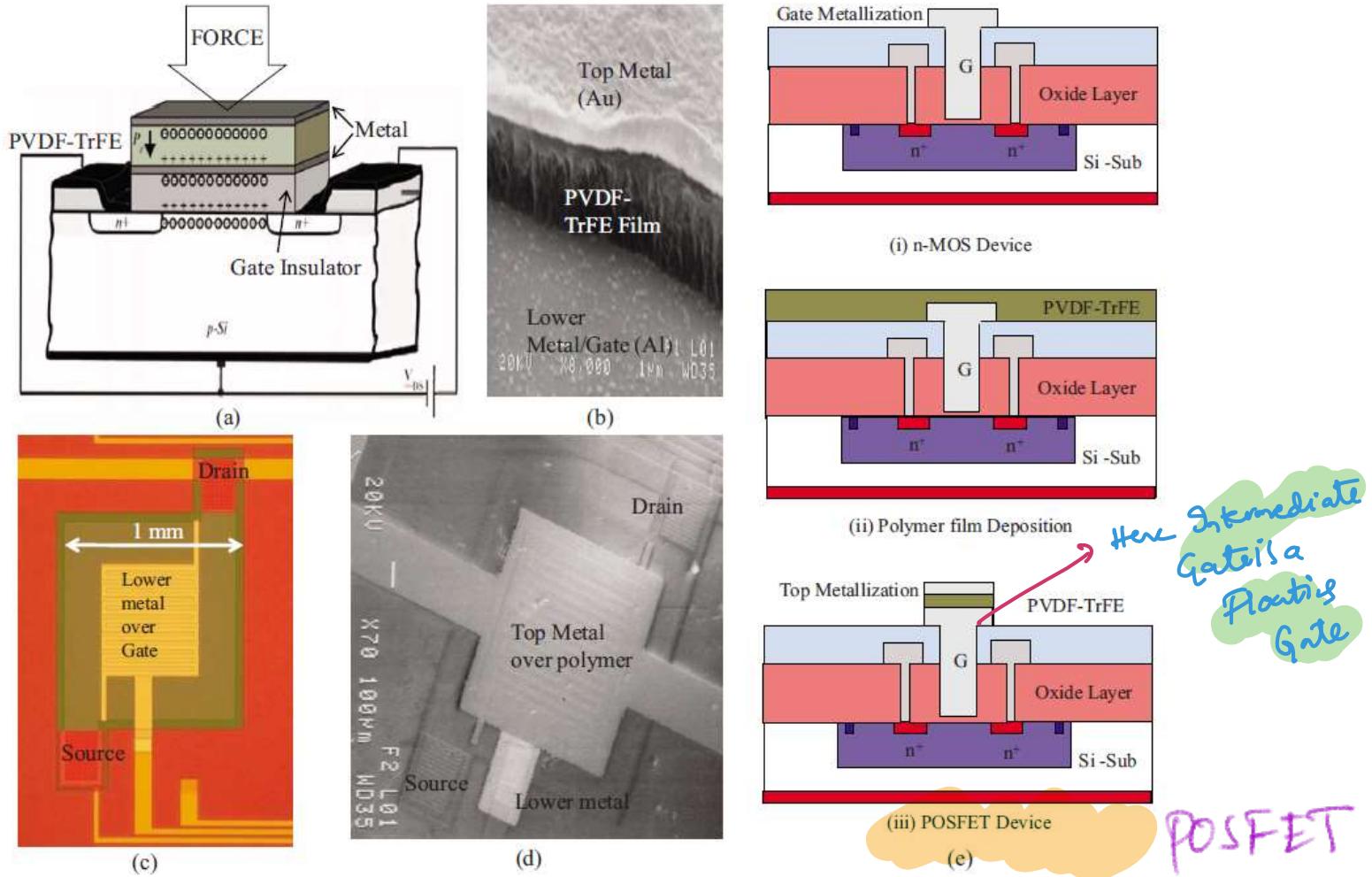
Ravinder S. Dahiya,<sup>1,a)</sup> Giorgio Metta,<sup>1,b)</sup> Maurizio Valle,<sup>2</sup> Andrea Adami,<sup>3</sup> and Leandro Lorenzelli<sup>3</sup>

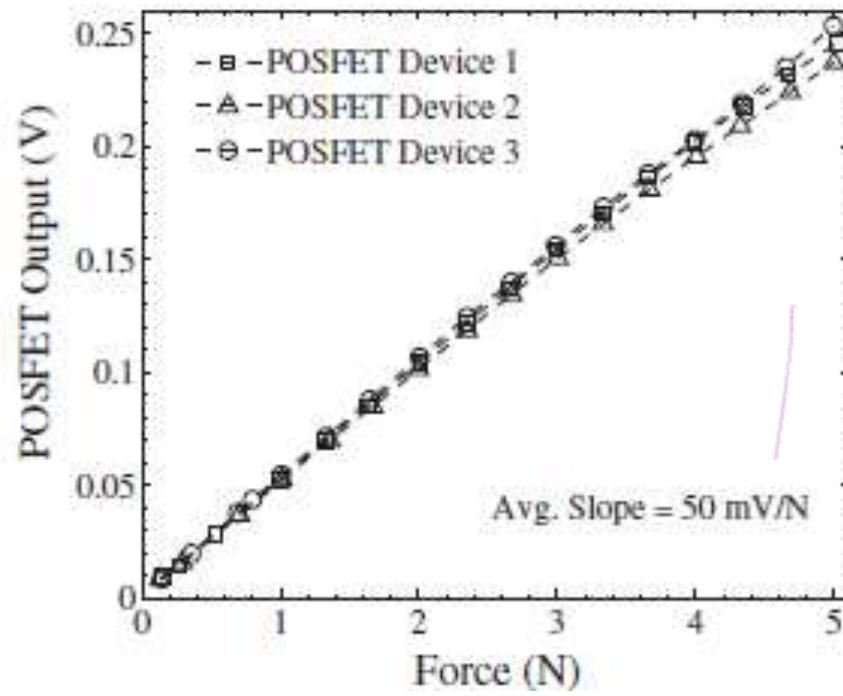
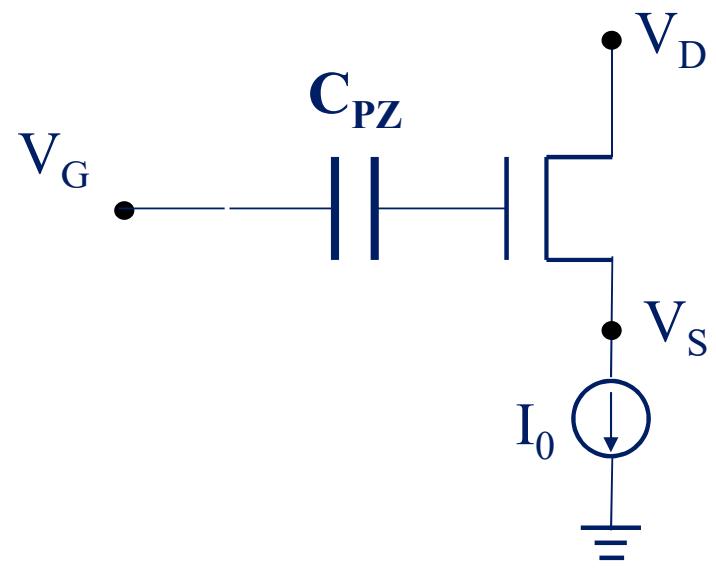
<sup>1</sup>*RBCS, Italian Institute of Technology, Genoa 16163, Italy*

<sup>2</sup>*DIBE, University of Genoa, 16145, Italy*

<sup>3</sup>*Bio-MEMS Group, Fondazione Bruno Kessler, Povo, Trento 38050, Italy*







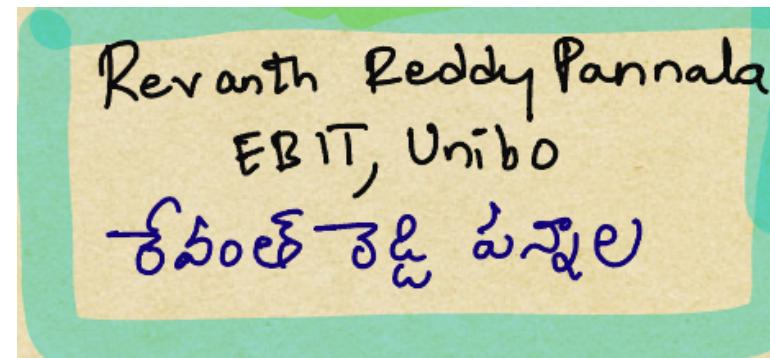
# Mechanical Sensors

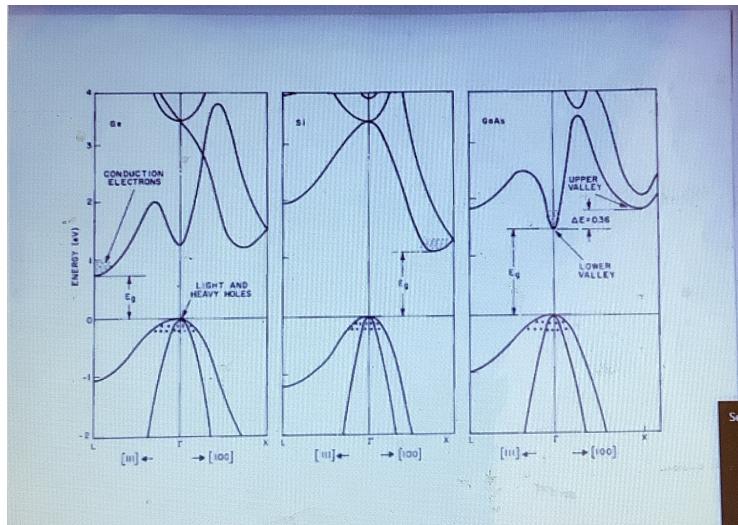
Characterization of mechanical sensors

- MEMS sensors and the different transduction approaches:
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  2. Piezoelectric sensors
  3. Piezoresistive pressure sensors

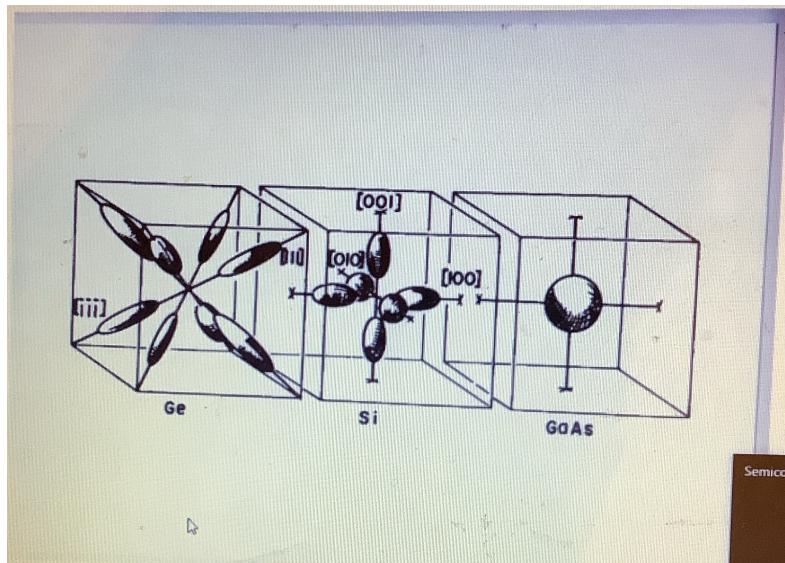
## Piezoresistive pressure sensors:

- Generalization of the piezoresistive effect:  
piezoresistive tensor
- Graphical representation of the piezoresistive  
coefficients in silicon (Kanda, 1982)
- The MAP sensor by MOTOROLA





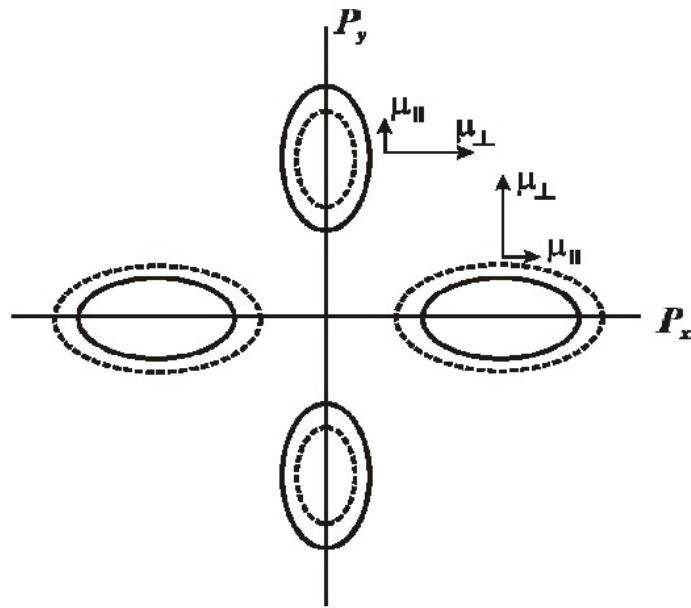
- This is a slide showing bands of Si, Ge & GaAs for Si has 'S' minima in the CB & they are aligned 2x2 in [100] [010] & [001] directions



- This figure shows the surfaces of equal energy for Si the surfaces of equal energy are ellipsoids of rotation ie two axis are equal and third axis different

- We remember that two axis that are equal are associated to the Transverse Mass or the Mass tensor. The other axis of the ellipsoid is longitudinal Mass.
- In order to describe the dynamics of electrons of silicon we describe the relation between Force & acceleration using a scalar mass that we called Effective Mass. This mass is a combination of Transversal & Longitudinal Mass.

## Piezoresistive effect



$$\Delta\rho/\rho_0 = \pi \sigma$$

$\rho$ - silicon resistivity

$\sigma$ - stress

$$\mathbf{E} = \rho_e [1 + \Pi \sigma] \mathbf{J}$$

Electric field      Resistivity tensor      Stress      Current density

## Piezoresistive effect

$$\mathbf{E} = \rho_e [1 + \boxed{\boldsymbol{\Pi}} \boldsymbol{\sigma}] \mathbf{J}$$

$$\begin{pmatrix} (\Delta\rho/\rho)_{11} \\ (\Delta\rho/\rho)_{22} \\ (\Delta\rho/\rho)_{33} \\ (\Delta\rho/\rho)_{23} \\ (\Delta\rho/\rho)_{13} \\ (\Delta\rho/\rho)_{12} \end{pmatrix} = \begin{pmatrix} \boldsymbol{\Pi}_{11} & \boldsymbol{\Pi}_{12} & \boldsymbol{\Pi}_{12} & 0 & 0 & 0 \\ \boldsymbol{\Pi}_{12} & \boldsymbol{\Pi}_{11} & \boldsymbol{\Pi}_{12} & 0 & 0 & 0 \\ \boldsymbol{\Pi}_{12} & \boldsymbol{\Pi}_{12} & \boldsymbol{\Pi}_{11} & 0 & 0 & 0 \\ 0 & 0 & 0 & \boldsymbol{\Pi}_{44} & 0 & 0 \\ 0 & 0 & 0 & 0 & \boldsymbol{\Pi}_{44} & 0 \\ 0 & 0 & 0 & 0 & 0 & \boldsymbol{\Pi}_{44} \end{pmatrix} \begin{pmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{33} \\ \sigma_{23} \\ \sigma_{13} \\ \sigma_{12} \end{pmatrix}$$

## Piezoresistive effect

$$\mathbf{E} = \rho_e [1 + \Pi \sigma] \mathbf{J}$$

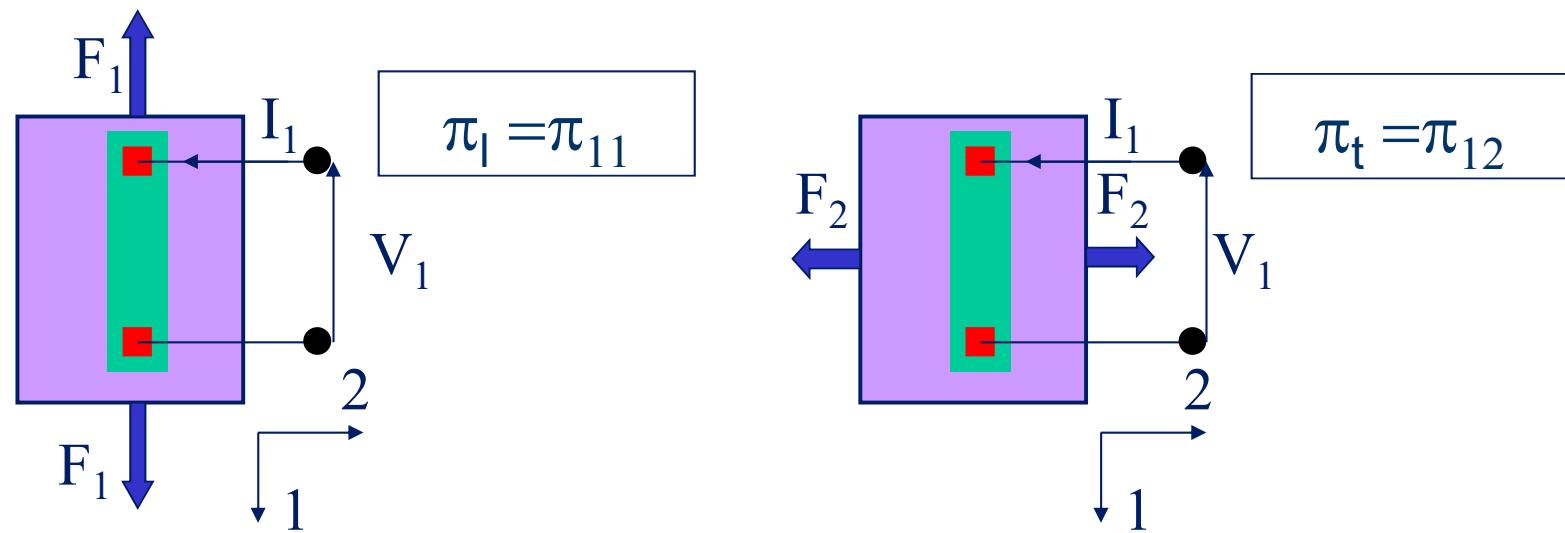
$$\frac{E_1}{\rho_e} = [1 + \pi_{11}\sigma_{11} + \pi_{12}(\sigma_{22} + \sigma_{33})] J_1 + \pi_{44}(\sigma_{12}J_2 + \sigma_{13}J_3)$$

$$\frac{E_2}{\rho_e} = [1 + \pi_{11}\sigma_{22} + \pi_{12}(\sigma_{11} + \sigma_{33})] J_2 + \pi_{44}(\sigma_{12}J_1 + \sigma_{23}J_3)$$

$$\frac{E_3}{\rho_e} = [1 + \pi_{11}\sigma_{33} + \pi_{12}(\sigma_{11} + \sigma_{22})] J_3 + \pi_{44}(\sigma_{13}J_1 + \sigma_{23}J_2)$$

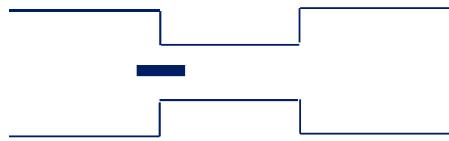
$$\frac{E_1}{\rho_e} = [1 + \pi_{11}\sigma_{11} + \pi_{12}(\sigma_{22} + \sigma_{33})] J_1 + \pi_{44}(\sigma_{12}J_2 + \sigma_{13}J_3)$$

Type	Resistivity	$\pi_{11}$	$\pi_{12}$	$\pi_{44}$
Units	$\Omega \text{-cm}$	$10^{-11} \text{ Pa}^{-1}$	$10^{-11} \text{ Pa}^{-1}$	$10^{-11} \text{ Pa}^{-1}$
<b>n-type</b>	11.7	-102.2	53.4	-13.6
<b>p-type</b>	7.8	6.6	-1.1	138.1



# Concept of a piezoresistive sensing scheme

If piezo-resistor is along [110]:



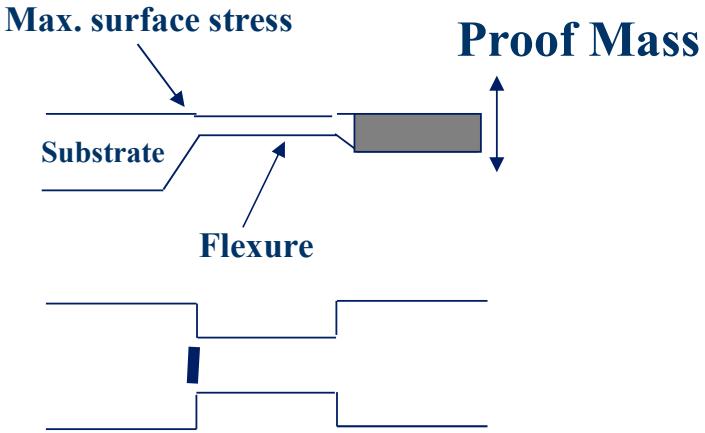
Longitudinal  
direction

[100]

$\pi_l$

[110]

$1/2 (\pi_{11} + \pi_{12} + \pi_{44})$



Transverse  
direction

[010]

$\pi_t$

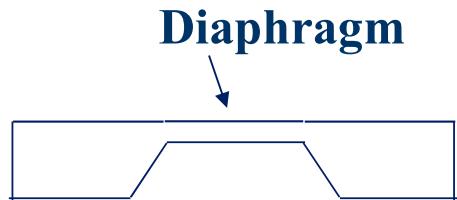
[110]

$1/2 (\pi_{11} + \pi_{12} - \pi_{44})$

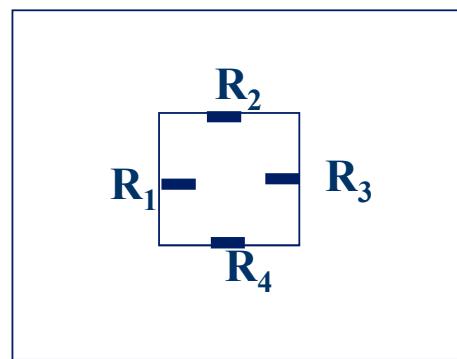
## Concept of a piezoresistive sensing scheme

Type	$\pi_l$	$\pi_t$	orientation
Units	$10^{-11} \text{ Pa}^{-1}$	$10^{-11} \text{ Pa}^{-1}$	
n-type	<b>-102.2</b>	<b>53.4</b>	[100]
p-type	<b>6.6</b>	<b>-1.1</b>	[100]
n-type	<b>-31.2</b>	<b>-17.6</b>	[110]
p-type	<b>71.8</b>	<b>-66.3</b>	[110]

CROSS-SECTION



TOP VIEW



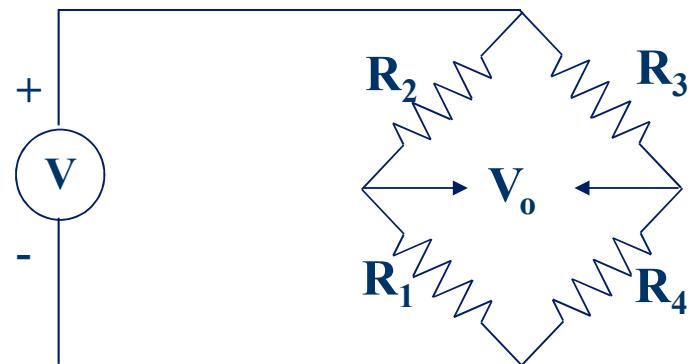
$$R_1 = R_3 = (1 + \alpha_1) R_o$$

$$R_2 = R_4 = (1 - \alpha_2) R_o$$

$$\frac{\Delta R_1}{R_1} = \pi_l \sigma_l + \pi_t \sigma_t$$

$$\frac{\Delta R_2}{R_2} = \pi_t \sigma_l + \pi_l \sigma_t$$

WHEATSTONE BRIDGE



$$\frac{\Delta R}{R} = \pi_l \sigma_l + \pi_t \sigma_t \quad l: \text{longitudinal}, t: \text{transverse}$$

**Longitudinal  
direction**

[100]

$\pi_l$

[110]

$\frac{1}{2} (\pi_{11} + \pi_{12} + \pi_{44})$

**Transverse  
direction**

[010]

$\pi_t$

$[\bar{1}10]$

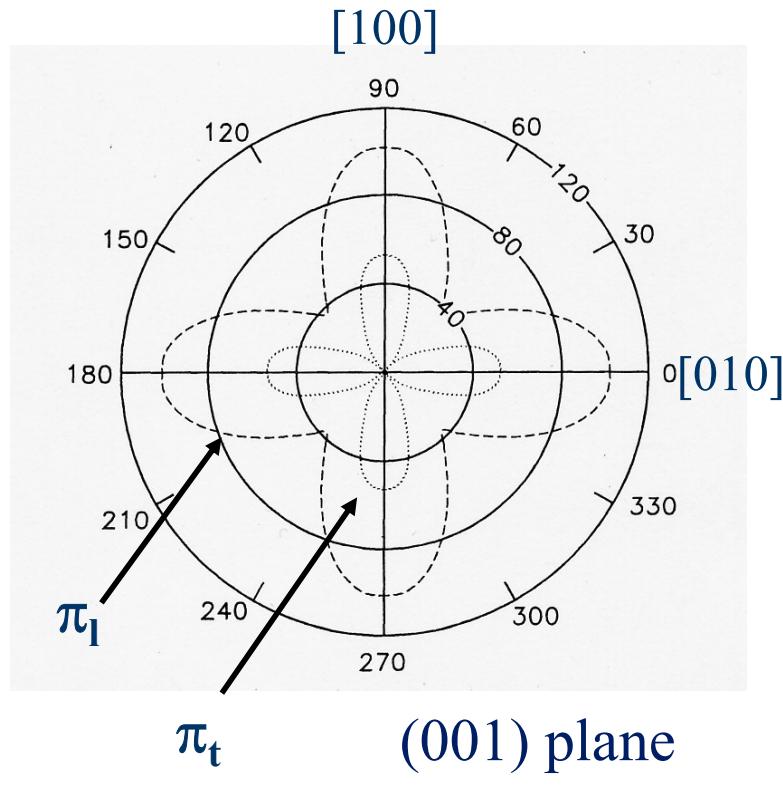
$\frac{1}{2} (\pi_{11} + \pi_{12} - \pi_{44})$



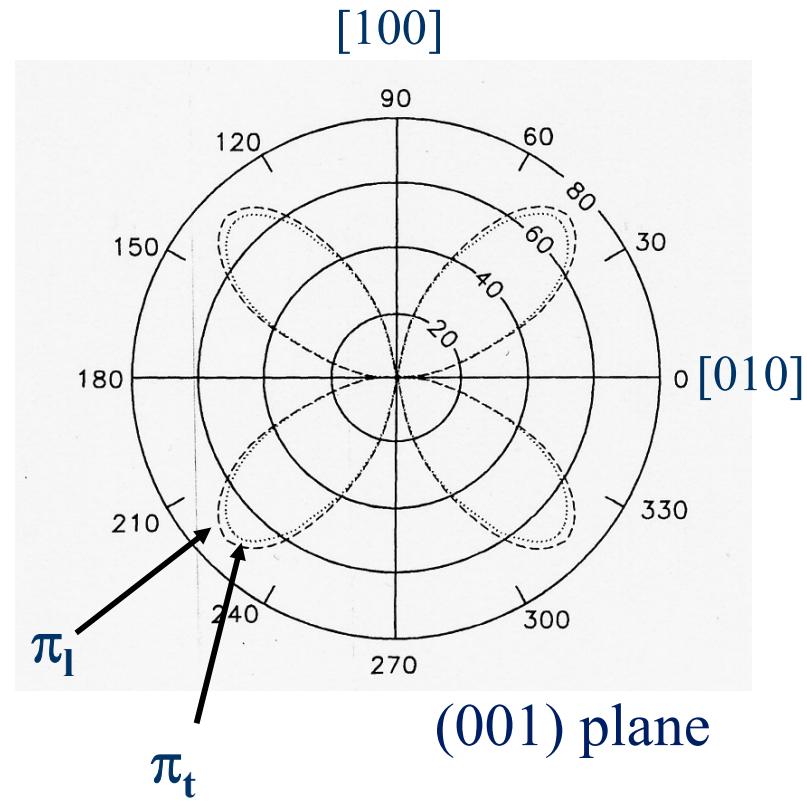
**GRAPHICAL REPRESENTATION**  
(Y. Kanda, IEEE Tr. ED 29, 1982)

## Longitudinal & Transverse piezoresistance coefficients

(Y. Kanda, IEEE Tr. ED 29, 1982)

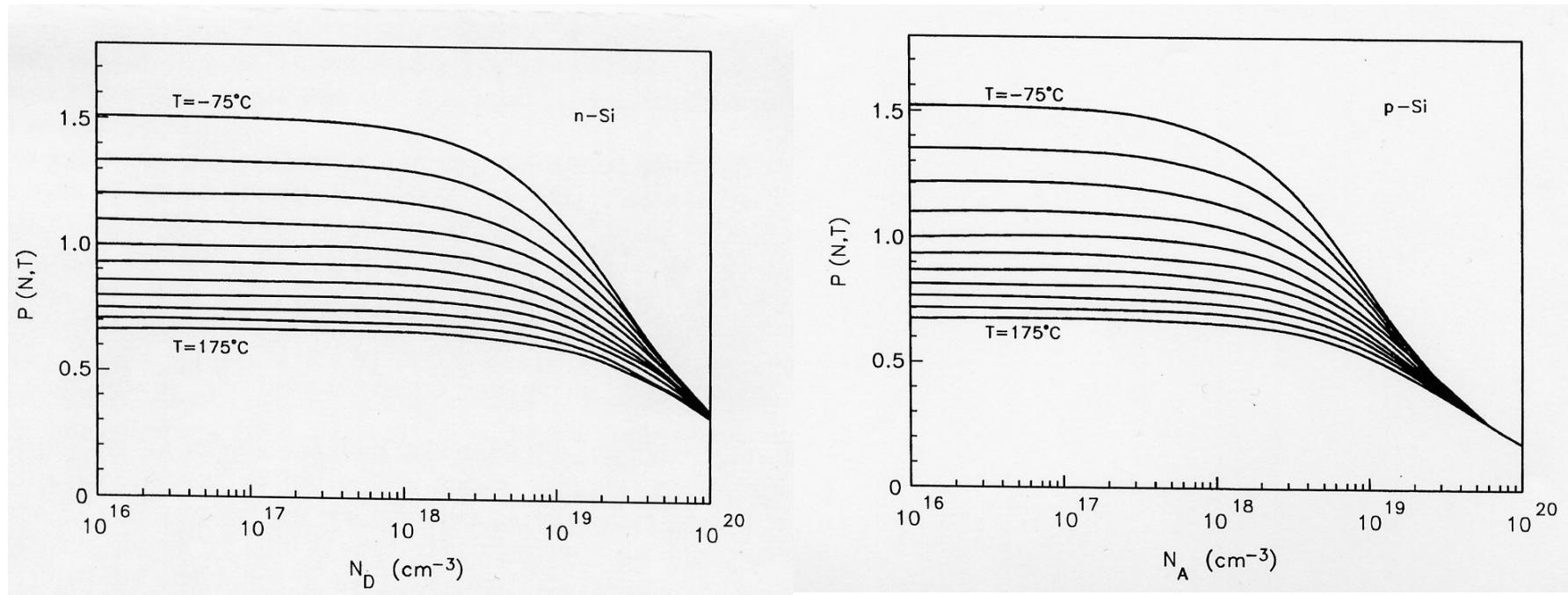


n-type Si



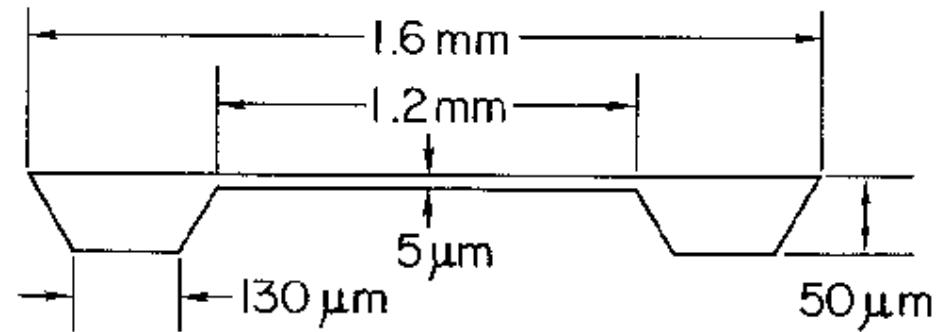
p-type Si

Piezoresistance coefficients as a function of impurity concentration and temperature for n-Si and p-Si  
(Y. Kanda, IEEE Tr. ED 29, 1982)

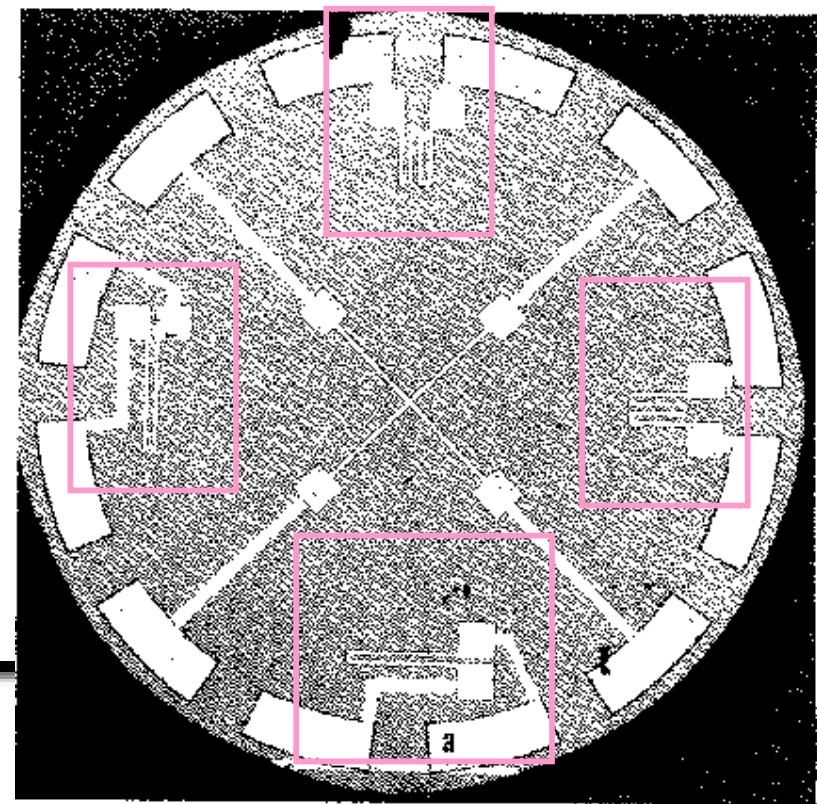


# An IC Piezoresistive Pressure Sensor for Biomedical Instrumentation

(S. Samaun et al., ISSCC 1971)

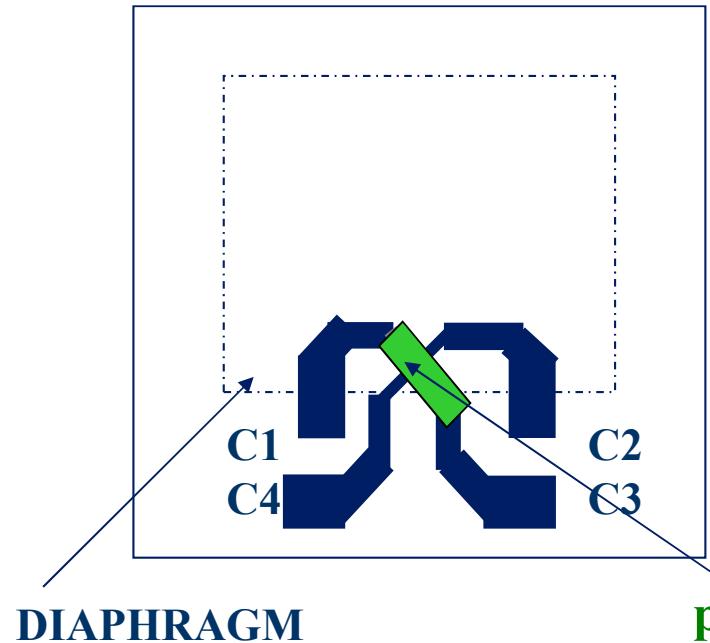


Layout →



# The Motorola X-ducer™ piezoresistor

TOP VIEW



piezoresistor

