

1)

*Proof.* Let  $w(x, t) = u(-x, t) - u(x, t)$ . Then,  $w(x, t)$  satisfies the differential equation and,

$$w(x, 0) = u(-x, 0) - u(x, 0) = \varphi(-x) - \varphi(x) = \varphi(x) - \varphi(x) = 0$$

Thus, we have the initial condition diffusion equation

$$w_t = kW_{xx} \quad w(x, 0) = 0$$

By uniqueness,  $w(x, 0) \equiv 0$ . Then,

$$\begin{aligned} u(-x, t) - u(x, t) &= 0 \\ u(-x, t) &= u(x, t) \end{aligned}$$

Thus,  $u(x, t)$  is an even function of  $x$ . Since  $u(x, t)$  is an even function, we know that  $u(x, t)$  will be symmetric around the graph representing the physical situation, meaning that there will be a situation where two different points in that domain will give you the same diffusion level and that at other points, the diffusion level will be zero.

□

2) Consider the  $Q$ -problem

$$\frac{\partial Q}{\partial t} = D \frac{\partial^2 Q}{\partial x^2} \quad \text{for } -\infty < x < \infty, t > 0$$

$$Q(x, 0) = \begin{cases} 0 & \text{if } x < 0 \\ -1 & \text{if } x > 0 \end{cases}$$

Now, suppose  $Q(x, t) = g\left(\frac{x}{\sqrt{4Dt}}\right)$ , for some one-variable function  $g$ .

Next, let  $\xi = \frac{x}{\sqrt{4Dt}}$ . Then,  $\frac{\partial Q}{\partial t} = g'(\xi) \left( -\frac{1}{2}x(4Dt)^{-\frac{3}{2}}(4D) \right)$ .

Looking more closely at  $g'(\xi)$ , we have

$$g'(\xi) = \left( -\frac{x4D}{2(4Dt)\sqrt{4Dt}} \right)$$

Returning to  $\frac{\partial Q}{\partial t}$  and simplifying, we have

$$\begin{aligned}\frac{\partial Q}{\partial t} &= g'(\xi) \left( -\frac{x}{2t\sqrt{4Dt}} \right) \\ \frac{\partial Q}{\partial t} &= g'(\xi) \left( -\frac{1}{2t}\xi \right)\end{aligned}$$

Now, if we look at derivatives with respect with to  $x$ , we get

$$\frac{\partial Q}{\partial x} = g'(\xi) \left( \frac{1}{\sqrt{4Dt}} \right) \quad \text{and} \quad \frac{\partial^2 Q}{\partial x^2} = g''(\xi) \left( \frac{1}{\sqrt{4Dt}} \right)$$

Then, if we substitute the derivatives we just calculated into the PDE, we see it transform into

$$-\frac{1}{2t}\xi g'(\xi) = D\frac{1}{4Dt}g''(\xi)$$

If we then multiply through by  $4t$ , we get

$$\begin{aligned}-\frac{4}{2}\xi g'(\xi) &= g''(\xi) \\ g''(\xi) + 2\xi g'(\xi) &= 0\end{aligned}$$

Now, let  $w(\xi) = g'(\xi)$  and  $w'(\xi) = g''(\xi)$ . Substituting these into our differential equation, we get

$$w'(\xi) + 2\xi w(\xi) = 0$$

So, if we remember our “training”, in order to solve this differential equation, the first thing we need to do is calculate the integrating factor

$$e^{\int_0^\xi 2s \, ds} = e^{\xi^2}$$

Then, if we multiply through by the integrating factor, we have

$$e^{\xi^2} w'(\xi) + 2\xi w(\xi) = 0$$

$$\frac{d}{d\xi} \left[ e^{\xi^2} w(\xi) \right] = 0$$

$$\int_0^\xi \frac{d}{ds} \left[ e^{s^2} w(s) \right] = 0$$

Continuing on, by the *Fundamental Theorem of Calculus*, we have

$$\left[ e^{s^2} w(s) \right] \Big|_{s=0}^{s=\xi}$$

Which, simplifies to the following if we evaluate it

$$e^{\xi^2} w(\xi) - w(0) = 0$$

$$w(\xi) = w(0)e^{-\xi^2}$$

Now, if we back substitute for  $w$ , we get

$$g'(\xi) = g'(0)e^{-\xi^2}$$

$$\int_0^\xi g'(s)ds = g'(0) \int_0^\xi e^{-s^2} ds$$

$$g(\xi) = g(0) + g'(0) \int_0^\xi e^{-s^2} ds$$

Now, if we recall

$$Q(x, 0) = \begin{cases} 0 & \text{if } x < 0 \\ -1 & \text{if } x > 0 \end{cases}$$

Then, for  $x > 0$ , we have

$$\begin{aligned} 1 = Q(x, 0) &= \lim_{x \rightarrow 0^+} Q(x, t) = g(0) + g'(0) \int_0^{+\infty} e^{-s^2} ds \\ 1 &= g(0) + g'(0) \int_0^{+\infty} e^{-s^2} ds = g(0) + g'(0) \frac{\sqrt{\pi}}{2} \\ 1 &= g(0) + g'(0) \frac{\sqrt{\pi}}{2} \end{aligned}$$

If we then look at the case where  $x < 0$ , we have

$$\begin{aligned} 0 = Q(x, 0) &= \lim_{x \rightarrow 0^-} Q(x, t) = g(0) + g'(0) \int_0^{+\infty} e^{-s^2} ds \\ 0 &= g(0) + g'(0) \int_0^{+\infty} e^{-s^2} ds = g(0) + g'(0) \frac{\sqrt{\pi}}{2} \\ 0 &= g(0) + g'(0) \frac{\sqrt{\pi}}{2} \end{aligned}$$

Finally, if we add our results for the above two cases we considered, we have

$$\begin{aligned} g(0) &= \frac{1}{2} \\ g'(0) &= \frac{1}{\sqrt{\pi}} \\ Q(x, t) &= g(\xi) = \frac{1}{2} + \frac{1}{\pi} \int_0^\xi e^{-s^2} ds \\ Q(x, t) &= g(\xi) = \frac{1}{2} + \frac{1}{\pi} \int_0^{\frac{x}{\sqrt{4Dt}}} e^{-s^2} ds \\ Q(x, t) &= g(\xi) = \frac{1}{2} + \frac{1}{2} \operatorname{erf}\left(\frac{x}{\sqrt{4Dt}}\right) \end{aligned}$$

Now, if we consider

$$\varphi(x) = \begin{cases} 0 & \text{if } x < -L \\ \sqrt{\pi} & \text{if } -L \leq x \leq L \\ 0 & \text{if } x > L \end{cases}$$

we can note a few things. The first being that whenever  $x$  is less than  $-L$ , then we have that  $\varphi$  is zero which indicates that the diffusion of kool-aid through the tube is also zero. Next, we note that when  $-L \leq x \leq L$ , we have that  $\varphi$  is equal to  $\sqrt{\pi}$ , so we know that the diffusion throughout the tube in this case is linear. Finally, whenever  $x$  is greater than  $L$ , we have that  $\varphi$  is equal to zero which indicates that the diffusion of kool-aid through the tube is, again, zero.

3) Suppose that if  $u_1(x, t)$  and  $u_2(x, t)$  are solutions to the given problem, then  $u_1(x, t) = u_2(x, t) \forall (x, t) \in \Omega_{L, \infty}$ , suppose

*Proof.* To begin, define  $w(x, t) = u_1(x, t) - u_2(x, t)$ .

At this point, we know that  $w$  satisfies the PDE  $\frac{\partial w}{\partial t} - D \frac{\partial^2 w}{\partial x^2} = 0$  with the initial condition of  $w(x, 0) = 0$  and the boundary conditions  $w(0, t) = 0$  and  $w(L, t) = 0$ . Okay, now, multiply  $\frac{\partial w}{\partial t} - D \frac{\partial^2 w}{\partial x^2} = 0$  by  $w$  to obtain

$$w \frac{\partial w}{\partial t} - Dw \frac{\partial^2 w}{\partial x^2} = 0$$

Now, before continuing, let us note the following:

$$\begin{aligned} 1) \quad & \frac{\partial}{\partial t} \left[ \frac{1}{2} w^2 \right] = w \left( \frac{\partial w}{\partial t} \right) \\ 2) \quad & \frac{\partial}{\partial x} \left[ -Dw \frac{\partial w}{\partial x} \right] = -Dw \frac{\partial^2 w}{\partial x^2} - D \left( \frac{\partial w}{\partial x} \right)^2 \\ 3) \quad & \frac{\partial}{\partial x} \left[ -Dw \frac{\partial w}{\partial x} \right] + D \left( \frac{\partial w}{\partial x} \right)^2 = -Dw \frac{\partial^2 w}{\partial x^2} \end{aligned}$$

Using the above points we just noted,  $w \frac{\partial w}{\partial t} - Dw \frac{\partial^2 w}{\partial x^2} = 0$  becomes

$$\frac{\partial}{\partial t} \left[ \frac{1}{2} w^2 \right] + \frac{\partial}{\partial x} \left[ -Dw \left( \frac{\partial w}{\partial x} \right) \right] + D \left( \frac{\partial w}{\partial x} \right)^2 = 0$$

Next, we want to integrate over the interval  $0 \leq x \leq L$ , giving us

$$\int_0^L \frac{\partial}{\partial t} \left[ \frac{1}{2} w^2 \right] dx + \int_0^L \frac{\partial}{\partial x} \left[ -Dw \left( \frac{\partial w}{\partial x} \right) \right] + D \int_0^L \left( \frac{\partial w}{\partial x} \right)^2 dx = 0$$

which simplifies to, with help from the *Fundamental Theorem of Calculus*,

$$\frac{d}{dt} \int_0^L \frac{1}{2} w^2 dx + \left[ -Dw \left( \frac{\partial w}{\partial x} \right) \right] \Big|_{x=0}^{x=L} + D \int_0^L \left( \frac{\partial w}{\partial x} \right)^2 dx = 0$$

Next, if we use the provided boundary conditions, the  $\left[ -Dw \left( \frac{\partial w}{\partial x} \right) \right] \Big|_{x=0}^{x=L}$  vanishes since, according to our boundary conditions

when we “plug-in” both  $x$  and  $L$  into our  $w$  function, we get zero each time. So, after doing some rearrangement, we come up with the following

$$\frac{d}{dt} \int_0^L \frac{1}{2} w^2 dx = -D \int_0^L \left( \frac{\partial w}{\partial x} \right)^2 dx \leq 0$$

Thus, we have that the function  $F(t) = \int_0^L \frac{1}{2} w(x, t)^2 dx$ , by the above function we just derived, is a decreasing function. So, since  $F(0) \geq F(t)$  for any  $t > 0$ , we know that  $\int_0^L w(x, 0)^2 dx \geq \int_0^L w(x, t)^2 dx$ . So,  $0 \geq \int_0^L w(x, t)^2 dx \forall t > 0$ . The previous statements imply that  $w(x, t) = 0 \forall (x, t) \in \Omega_{L, \infty}$ , which in turn implies that  $u_1(x, t) = u_2(x, t) \in \Omega_{L, \infty}$ . □