## 1 Correlation Functions

This analysis studies the momentum correlations of both  $\Lambda$ -K and  $\Xi$ -K pairs using the two-particle correlation function, defined as  $C(k^*) = A(k^*)/B(k^*)$ , where  $A(k^*)$  is the signal distribution,  $B(k^*)$  is the reference (or background) distribution, and  $k^*$  is the momentum of one of the particles in the pair rest frame. In practice,  $A(k^*)$  is constructed by binning in  $k^*$  pairs from the same event. Ideally,  $B(k^*)$  is similar to  $A(k^*)$  in all respects excluding the presence of femtoscopic correlations [?]; as such,  $B(k^*)$  is used to divide out the phase-space effects, leaving only the femtoscopic effects in the correlation function.

This analysis presents correlation functions for three centrality bins (0-10%, 10-30%, and 30-50%), and is currently pair transverse momentum ( $k_T = 0.5 |\mathbf{p}_{T,1} + \mathbf{p}_{T,2}|$ ) integrated (i.e. not binned in  $k_T$ ). The correlation functions are constructed separately for the two magnetic field configurations, and are combined using a weighted average:

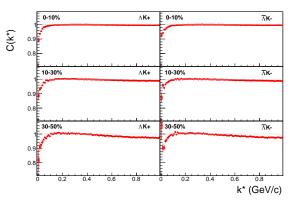
$$C_{combined}(k^*) = \frac{\sum_{i} w_i C_i(k^*)}{\sum_{i} w_i}$$
 (1)

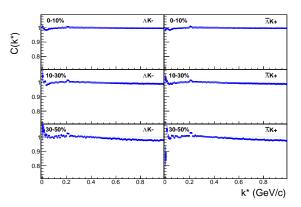
where the sum runs over the correlation functions to be combined, and the weight,  $w_i$ , is the number of numerator pairs in  $C_i(k^*)$ . Here, the sum is over the two field configurations.

## 1.1 Typical Correlation Function Construction

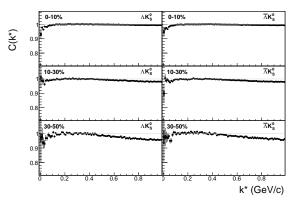
Typically, in practice,  $B(k^*)$  is obtained by forming mixed-event pairs, i.e. particles from a given event are paired with particles from  $N_{mix}(=5)$  other events, and these pairs are then binned in  $k^*$ . In forming the background distribution, it is important to mix only similar events; mixing events with different phase-spaces can lead to artificial signals in the correlaton function. Therefore, in this analysis, we mix events with primary vertices within 2 cm and centralities within 5% of each other. Also note, a vertex correction is also applied to each event, which essentially recenters the the primary vertices to z = 0.

Figures 1, 2, and 3 show the correlation functions for all centalities studied for  $\Lambda K_S^0(\bar{\Lambda}K_S^0)$ ,  $\Lambda K^+(\bar{\Lambda}K^-)$ , and  $\Lambda K^-(\bar{\Lambda}K^+)$ , respectively. All were normalized in the range  $0.32 < k^* < 0.4$  GeV/c.





- (a)  $\Lambda K^+$  (left) and  $\bar{\Lambda} K^-$  (right) correlations for 0-10% (top), 10-30%(middle), and 30-50%(bottom) centralities.
- (b)  $\Lambda K^-$  (left) and  $\bar{\Lambda} K^+$  (right) correlations for 0-10% (top), 10-30%(middle), and 30-50%(bottom) centralities. The peak at  $k^*\approx 0.2$  GeV/c is due to the  $\Omega^-$  resonance.



(c)  $\Lambda K_S^0$  (left) and  $\bar{\Lambda} K_S^0$  (right) correlations for 0-10% (top), 10-30%(middle), and 30-50%(bottom) centralities.

Fig. 1:  $\Lambda K$  and  $\bar{\Lambda} \bar{K}$  correlation functions for 0-10%, 10-30%, and 30-50% centralities. The lines represent the statistical errors, while the boxes represent the systematic errors.

## 1.2 Stavinsky Correlation Function Construction

Stavinsky is tight.

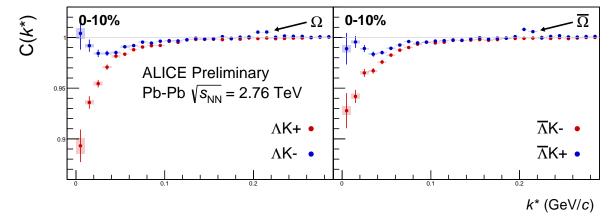


Fig. 2: Correlation Functions:  $\Lambda K^+$  vs  $\Lambda K^-$  ( $\bar{\Lambda} K^+$  vs  $\bar{\Lambda} K^-$ ) for 0-10% centrality. The peak in  $\Lambda K^-$ ( $\bar{\Lambda} K^+$ ) at  $k^* \approx 0.2$  GeV/c is due to the  $\Omega^-$  resonance. The lines represent the statistical errors. (NOTE: This figure is slightly dated, and a new one will be generated which includes both statistical and systematic uncertainties)