

0.1 Model: ΛK_S^0 , ΛK^\pm , $\Xi^{ch} K_S^0$

The two-particle relative momentum correlation function may be written theoretically by the Koonin-Pratt equation [?, ?]:

$$C(\mathbf{k}^*) = \int S(\mathbf{r}^*) |\Psi_{\mathbf{k}^*}(\mathbf{r}^*)|^2 d^3 \mathbf{r}^* \quad (1)$$

In the absence of Coulomb effects, and assuming a spherically gaussian source of width R , the 1D femtoscopic correlation function can be calculated analytically using:

$$C(k^*) = 1 + \lambda [C_{QI}(k^*) + C_{FSI}(k^*)] \quad (2)$$

C_{QI} describes plane-wave quantum interference:

$$C_{QI}(k^*) = \alpha \exp(-4k^{*2}R^2) \quad (3)$$

where $\alpha = (-1)^{2j}/(2j+1)$ for identical particles with spin j , and $\alpha = 0$ for non-identical particles. Obviously, $\alpha = 0$ for all analyses presented in this note. C_{FSI} describes the s-wave strong final state interaction between the particles:

$$C_{FSI}(k^*) = (1 + \alpha) \left[\frac{1}{2} \left| \frac{f(k^*)}{R} \right|^2 \left(1 - \frac{d_0}{2\sqrt{\pi}R} \right) + \frac{2\Re f(k^*)}{\sqrt{\pi}R} F_1(2k^*R) - \frac{\Im f(k^*)}{R} F_2(2k^*R) \right] \quad (4)$$

$$f(k^*) = \left(\frac{1}{f_0} + \frac{1}{2} d_0 k^{*2} - i k^* \right)^{-1}; \quad F_1(z) = \int_0^z \frac{e^{x^2 - z^2}}{z} dx; \quad F_2(z) = \frac{1 - e^{-z^2}}{z}$$

where R is the source size, $f(k^*)$ is the s-wave scattering amplitude, f_0 is the complex scattering length, and d_0 is the effective range of the interaction.

The code developed to fit the data is called ‘‘LednickyFitter’’, and utilizes the ROOT TMinuit implementation of the MINUIT fitting package. In short, given a function with a number of parameters, the fitter explores the parameter space searching for the minimum of the equation. In this implementation, the function to be minimized should represent the difference between the measure and theoretical correlation functions. However, a simple χ^2 test is inappropriate for fitting correlation functions, as the ratio two Poisson distributions does not result in a Poisson distribution. Instead, a log-likelihood fit function of the following form is used [?]:

$$\chi_{PML}^2 = -2 \left[A \ln \left(\frac{C(A+B)}{A(C+1)} \right) + B \ln \left(\frac{A+B}{B(C+1)} \right) \right] \quad (5)$$

where A is the experimental signal distribution (numerator), B is the experimental background distribution (denominator), and C is the theoretical fit correlation function.

The LednickyFitter uses Equations 2 – 4 to build the theoretical fit, and Equation 5 as the statistic quantifying the quality of the fit. The parameters to be varied by MINUIT are: λ , R , f_0 ($\Re f_0$ and $\Im f_0$ separately), d_0 , and normalization N . The fitter currently includes methods to correct for momentum resolution and a non-flat background. These corrections are applied to the fit function, the data is never touched. The fitter is able to share parameters between different analyses and fit all simultaneously.

In a typical fit, a given pair is fit with its conjugate (ex. ΛK^+ with $\bar{\Lambda} K^-$) across all centralities (0-10%, 10-30%, 30-50%), for a total of 6 simultaneous analyses. Each analysis has a unique λ and normalization

parameter. The radii are shared between analyses of like centrality, as these should have similar source sizes. The scattering parameters ($\mathbb{R}f_0$, $\mathbb{I}f_0$, d_0) are shared amongst all.

Figures ??, ??, and ?? (??, ??, and ??, or ??, ??, and ??), in Section ??, show experimental data with fits for all studied centralities for ΛK_S^0 with $\bar{\Lambda} K_S^0$, ΛK^+ with $\bar{\Lambda} K^-$, and ΛK^- with $\bar{\Lambda} K^+$, respectively. In the figures, the black solid line represents the “raw” fit, i.e. not corrected for momentum resolution effects nor non-flat background. The green line shows the fit to the non-flat background. The purple points show the fit after momentum resolution, non-flat background, and residual correlations (if applicable) corrections have been applied. The initial values of the parameters is listed, as well as the final fit values with uncertainties.