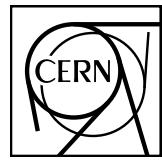


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ΛK and $\Xi^- K^\pm$ Femtoscopy in Pb-Pb Collisions at $\sqrt{s_{NN}} = 2.76$ TeV from the LHC ALICE Experiment

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Abstract

I present the first femtoscopic analysis of ΛK correlations in Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV measured by the ALICE experiment at the Large Hadron Collider (LHC). All pair combinations of Λ and $\bar{\Lambda}$ with K^+ , K^- and K_S^0 are analyzed. The femtoscopic correlations result from strong final-state interactions, and are fit with a parametrization allowing us to both characterize the emission source and measure the scattering parameters for the particle pairs. We observe a large difference in the ΛK^+ and ΛK^- correlations in pairs with low relative momenta ($k^* \lesssim 100$ MeV/c). This might suggest an effect arising from different quark-antiquark interactions within the pairs ($s\bar{s}$ in ΛK^+ and $u\bar{u}$ in ΛK^-), or from different net strangeness for each system. To gain further insight into this hypothesis, we are currently conducting a $\Xi^- K^\pm$ femtoscopic analysis, for which I present preliminary results.

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10 1 Introduction

11 We present results from a femtoscopic analysis of ΛK and $\Xi^- K^\pm$ correlations in Pb-Pb collisions at
 12 $\sqrt{s_{NN}} = 2.76$ TeV by the ALICE experiment at the LHC. All pair combinations of Λ and $\bar{\Lambda}$ with K^+ ,
 13 K^- and K_S^0 are analyzed. The femtoscopic correlations are the result of strong final-state interactions,
 14 and are fit with a parametrization based on a model by R. Lednicky and V. L. Lyuboshitz [1]. This
 15 allows us to both characterize the emission source and measure the scattering parameters for the particle
 16 pairs. We observe a large difference in the $\Lambda K^+(\bar{\Lambda} K^-)$ and $\Lambda K^-(\bar{\Lambda} K^+)$ correlations in pairs with low
 17 relative momenta ($k^* \lesssim 100$ MeV). The results suggest an effect arising from different quark-antiquark
 18 interactions in the pairs, i.e. $s\bar{s}$ in $\Lambda K^+(\bar{\Lambda} K^-)$ and $u\bar{u}$ in $\Lambda K^-(\bar{\Lambda} K^+)$, or from different net strangeness for
 19 each system. To gain further insight into this hypothesis, we currently are conducting a ΞK femtoscopic
 20 analysis.

21 2 Data Sample and Software

22 2.1 Data Sample

23 The analysis used “pass 2” reconstructed Pb-Pb data from LHC11h (AOD145). The runlist was selected
 24 from runs with global quality tag “1” in the ALICE Run Condition Table. Approximately 40 million
 25 combined central, semi-central, and minimum bias events were analyzed. Runs from both positive (++)
 26 and negative (--) magnetic field polarity settings were used.

27 Run list: 170593, 170572, 170388, 170387, 170315, 170313, 170312, 170311, 170309, 170308, 170306,
 28 170270, 170269, 170268, 170230, 170228, 170207, 170204, 170203, 170193, 170163, 170159, 170155,
 29 170091, 170089, 170088, 170085, 170084, 170083, 170081, 170040, 170027, 169965, 169923, 169859,
 30 169858, 169855, 169846, 169838, 169837, 169835, 169591, 169590, 169588, 169587, 169586, 169557,
 31 169555, 169554, 169553, 169550, 169515, 169512, 169506, 169504, 169498, 169475, 169420, 169419,
 32 169418, 169417, 169415, 169411, 169238, 169167, 169160, 169156, 169148, 169145, 169144, 169138,
 33 169099, 169094, 169091, 169045, 169044, 169040, 169035, 168992, 168988, 168826, 168777, 168514,
 34 168512, 168511, 168467, 168464, 168460, 168458, 168362, 168361, 168342, 168341, 168325, 168322,
 35 168311, 168310, 168315, 168108, 168107, 168105, 168076, 168069, 167988, 167987, 167985, 167920,
 36 167915

37 Analysis was also performed on the LHC12a17a_fix (AOD149) Monte Carlo HIJING events for certain
 38 checks. THERMINATOR2 was also used for certain aspects, such as generation of transform matrices
 39 describing feed-down contributions, estimation of λ parameters, and non-femtoscopic background
 40 modeling.

41 2.2 Software

42 The analysis was performed on the PWGCF analysis train using AliRoot v5-09-29-1 and AliPhysics
 43 vAN-20180505-1.

44 The main classes utilized include: AliFemtoVertexMultAnalysis, AliFemtoEventCutEstimators, AliFemto
 45 toESDTrackCutNSigmaFilter, AliFemtoV0TrackCutNSigmaFilter, AliFemtoXiTrackCut, AliFemtoV0PairCut,
 46 AliFemtoV0TrackPairCut, AliFemtoXiTrackPairCut, and AliFemtoAnalysisLambdaKaon. All of these
 47 classes are contained in /AliPhysics/PWGCF/FEMTOSCOPY/AliFemto and .../AliFemtoUser.

48 3 Data Selection

49 3.1 Event Selection and Mixing

50 The events used in this study were selected with the class AliFemtoEventCutEstimators according to the
 51 following criteria:

- 52 – Triggers
 - 53 – minimum bias (kMB)
 - 54 – central (kCentral)
 - 55 – semi-central (kSemiCentral)
- 56 – z-position of reconstructed event vertex must be within 10 cm of the center of the ALICE detector
- 57 – the event must contain at least one particle of each type from the pair of interest

58 The event mixing was handled by the AliFemtoVertexMultAnalysis class, which only mixes events with
 59 like vertex position and centrality. The following criteria were used for event mixing:

- 60 – Number of events to mix = 5
- 61 – Vertex position bin width = 2 cm
- 62 – Centrality bin width = 5

63 The AliFemtoEventReaderAODChain class is used to read the events. Event flattening is not currently
 64 used. FilterBit(7). The centrality is determined by the “V0M” method of AliCentrality, set by calling
 65 AliFemtoEventReaderAOD::SetUseMultiplicity(kCentrality). We utilize the SetPrimaryVertexCorrec-
 66 tionTPCPoints switch, which causes the reader to shift all TPC points to be relative to the event vertex.

67 3.2 K^\pm Track Selection

68 Charged kaons are identified using the AliFemtoESDTrackCutNSigmaFilter class. The specific cuts used
 69 in this analysis are as follows:

70 Track Selection:

- 71 – Kinematic range:
 - 72 – $0.14 < p_T < 1.5 \text{ GeV}/c$
 - 73 – $|\eta| < 0.8$
- 74 – FilterBit(7)
 - 75 – TPC tracks
- 76 – Track Quality
 - 77 – Minimum number of clusters in the TPC (fminTPCncls) = 80
 - 78 – Maximum allowed χ^2/N_{DOF} for ITS clusters = 3.0
 - 79 – Maximum allowed χ^2/N_{DOF} for TPC clusters = 4.0

- 80 – Primary Particle Selection:
 - 81 – Maximum XY impact parameter = 2.4 cm
 - 82 – Maximum Z impact parameter = 3.0 cm
- 83 – Remove particles with any kink labels (fRemoveKinks = true)
- 84 – Maximum allowed sigma to primary vertex (fMaxSigmaToVertex) = 3.0

85 K $^\pm$ Identification:

86 – PID Probabilities:

- 87 – K: > 0.2
- 88 – π : < 0.1
- 89 – μ : < 0.8
- 90 – p: < 0.1

91 – Most probable particle type must be Kaon (fMostProbable=3)

92 – TPC and TOF N $_\sigma$ cuts:

- 93 – $p < 0.4 \text{ GeV}/c$: N $_{\sigma K, \text{TPC}} < 2$
- 94 – $0.4 < p < 0.45 \text{ GeV}/c$: N $_{\sigma K, \text{TPC}} < 1$
- 95 – $0.45 < p < 0.8 \text{ GeV}/c$: N $_{\sigma K, \text{TPC}} < 3 \& N_{\sigma K, \text{TOF}} < 2$
- 96 – $0.8 < p < 1.0 \text{ GeV}/c$: N $_{\sigma K, \text{TPC}} < 3 \& N_{\sigma K, \text{TOF}} < 1.5$
- 97 – $p > 1.0 \text{ GeV}/c$: N $_{\sigma K, \text{TPC}} < 3 \& N_{\sigma K, \text{TOF}} < 1$

98 – Electron Rejection: Reject if N $_{\sigma e^-, \text{TPC}} < 3$

99 – Pion Rejection: Reject if:

- 100 – $p < 0.65 \text{ GeV}/c$
 - 101 * if TOF and TPC available: N $_{\sigma \pi, \text{TPC}} < 3 \& N_{\sigma \pi, \text{TOF}} < 3$
 - 102 * else
 - 103 · $p < 0.5 \text{ GeV}/c$: N $_{\sigma \pi, \text{TPC}} < 3$
 - 104 · $0.5 < p < 0.65 \text{ GeV}/c$: N $_{\sigma \pi, \text{TPC}} < 2$
- 105 – $0.65 < p < 1.5 \text{ GeV}/c$: N $_{\sigma \pi, \text{TPC}} < 5 \& N_{\sigma \pi, \text{TOF}} < 3$
- 106 – $p > 1.5 \text{ GeV}/c$: N $_{\sigma \pi, \text{TPC}} < 5 \& N_{\sigma \pi, \text{TOF}} < 2$

107 The purity of the K $^\pm$ collections was estimated using the HIJING MC data, for which the true identity
108 of each reconstructed K $^\pm$ particle is known. Therefore, the purity may be estimated as:

$$\text{Purity}(K^\pm) = \frac{N_{\text{true}}}{N_{\text{reconstructed}}} \quad (1)$$

109 Purity(K $^+$) \approx Purity(K $^-$) \approx 97%

110 3.3 V0 Selection

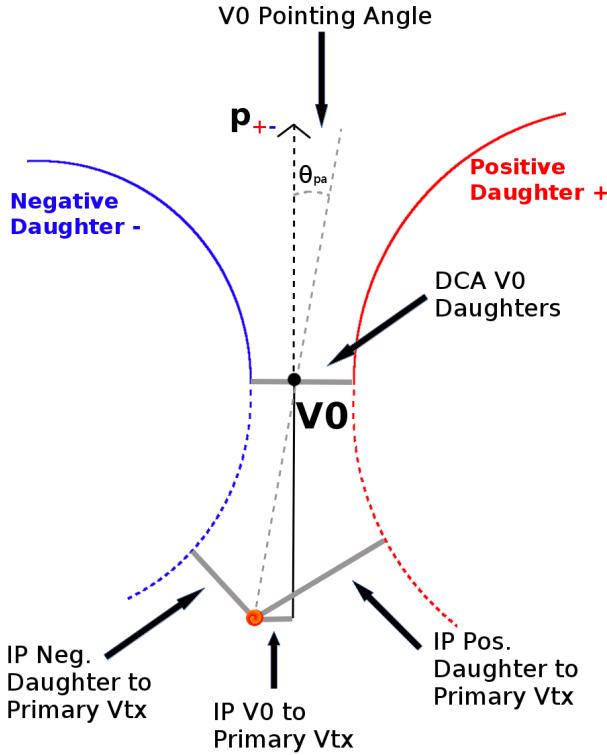
111 $\Lambda(\bar{\Lambda})$ and K $_S^0$ are neutral particles which cannot be directly detected, but must instead be reconstructed
112 through detection of their decay products, or daughters. This process is illustrated in Figure 1. In
113 general, particles which are topologically reconstructed in this fashion are called V0 particles. The
114 class AliFemtoV0TrackCutNSigmaFilter (which is an extension of AliFemtoV0TrackCut) is used to
115 reconstruct the V0s.

116 In order to ensure a true and reliable signal, one must ensure good purity of the V0 collection. The purity
117 of the collection is calculated as:

$$\text{Purity} = \frac{\text{Signal}}{\text{Signal} + \text{Background}} \quad (2)$$

| K^\pm selection | |
|--|--|
| $ \eta $ | < 0.8 |
| p_T | $0.14 < p_T < 1.5 \text{ GeV}/c$ |
| FilterBit | 7 |
| Min. number of clusters in the TPC | 80 |
| Max. allowed χ^2/N_{DOF} for ITS clusters | 3.0 |
| Max. allowed χ^2/N_{DOF} for TPC clusters | 4.0 |
| Maximum XY impact parameter | 2.4 cm |
| Maximum Z impact parameter | 3.0 cm |
| Remove particles with any kink labels | true |
| Maximum allowed sigma to primary vertex | 3.0 |
| PID Probabilities | |
| K | > 0.2 |
| π | < 0.1 |
| μ | < 0.8 |
| p | < 0.1 |
| Most probable particle type | Kaon (fMostProbable=3) |
| TPC and TOF N σ Cuts | |
| $p < 0.4 \text{ GeV}/c$ | $N_{\sigma K, \text{TPC}} < 2$ |
| $0.4 < p < 0.45 \text{ GeV}/c$ | $N_{\sigma K, \text{TPC}} < 1$ |
| $0.45 < p < 0.80 \text{ GeV}/c$ | $N_{\sigma K, \text{TPC}} < 3 \& N_{\sigma K, \text{TOF}} < 2$ |
| $0.80 < p < 1.0 \text{ GeV}/c$ | $N_{\sigma K, \text{TPC}} < 3 \& N_{\sigma K, \text{TOF}} < 1.5$ |
| $p > 1.0 \text{ GeV}/c$ | $N_{\sigma K, \text{TPC}} < 3 \& N_{\sigma K, \text{TOF}} < 1$ |
| Electron Rejection | Reject if $N_{\sigma e^-, \text{TPC}} < 3$ |
| Pion Rejection: Reject if: | |
| $p < 0.65 \text{ GeV}/c$ | if TOF and TPC available $N_{\sigma \pi, \text{TPC}} < 3 \& N_{\sigma \pi, \text{TOF}} < 3$ |
| $0.65 < p < 1.5 \text{ GeV}/c$ | $p < 0.5 \text{ GeV}/c$ $N_{\sigma \pi, \text{TPC}} < 3$ |
| $p > 1.5 \text{ GeV}/c$ | $0.5 < p < 0.65 \text{ GeV}/c$ $N_{\sigma \pi, \text{TPC}} < 2$ |
| | $N_{\sigma \pi, \text{TPC}} < 5 \& N_{\sigma \pi, \text{TOF}} < 3$ |
| | $N_{\sigma \pi, \text{TPC}} < 5 \& N_{\sigma \pi, \text{TOF}} < 2$ |

Table 1: K^\pm selection

**Fig. 1:** V0 Reconstruction

118 To obtain both the signal and background, the invariant mass distribution (m_{inv}) of all V0 candidates must
 119 be constructed immediately before the final invariant mass cut. Examples of such distributions can be
 120 found in Figures 3 and 5. It is vital that the distribution be constructed immediately before the final m_{inv}
 121 cut, otherwise, it would be impossible to estimate the background. As demonstrated in Figures 3 and 5,
 122 the background is fit with a polynomial outside of the peak region of interest in order to extrapolate an
 123 estimate for the background within the region. Within the m_{inv} cut limits, the background is the region
 124 below the fit while the signal is the region above the fit.

125 3.3.1 **Λ Reconstruction**

126 The following cuts were used to select good $\Lambda(\bar{\Lambda})$ candidates:

127 1. Daughter Particle Cuts

128 (a) Cuts Common to Both Daughters

- 129 i. $|\eta| < 0.8$
- 130 ii. SetTPCnclsDaughters(80)
- 131 iii. SetStatusDaughters(AliESDtrack::kTPCrefit)
- 132 iv. DCA πp Daughters < 0.4 cm

133 (b) Pion Specific Daughter Cuts

- 134 i. $p_T > 0.16$ GeV/c
- 135 ii. DCA to prim vertex > 0.3 cm
- 136 iii. TPC and TOF N σ Cuts
 - 137 A. $p < 0.5$ GeV/c : N $\sigma_{\text{TPC}} < 3$
 - 138 B. $p > 0.5$ GeV/c :

- 139 – if TOF & TPC available: $N\sigma_{\text{TPC}} < 3$ & $N\sigma_{\text{TOF}} < 3$
 140 – else $N\sigma_{\text{TPC}} < 3$

141 (c) Proton Specific Daughter Cuts

- 142 i. $p_T > 0.5(p)[0.3(\bar{p})] \text{ GeV}/c$
 143 ii. DCA to prim vertex $> 0.1 \text{ cm}$
 144 iii. TPC and TOF $N\sigma$ Cuts
 145 A. $p < 0.8 \text{ GeV}/c : N\sigma_{\text{TPC}} < 3$
 146 B. $p > 0.8 \text{ GeV}/c :$
 147 – if TOF & TPC available: $N\sigma_{\text{TPC}} < 3$ & $N\sigma_{\text{TOF}} < 3$
 148 – else $N\sigma_{\text{TPC}} < 3$

149 2. V0 Cuts

- 150 (a) $|\eta| < 0.8$
 151 (b) $p_T > 0.4 \text{ GeV}/c$
 152 (c) $|m_{\text{inv}} - m_{\text{PDG}}| < 3.8 \text{ MeV}$
 153 (d) DCA to prim. vertex $< 0.5 \text{ cm}$
 154 (e) Cosine of pointing angle > 0.9993
 155 (f) OnFlyStatus = false
 156 (g) Decay Length $< 60 \text{ cm}$

157 3. Shared Daughter Cut for V0 Collection

- 158 – Iterate through V0 collection to ensure that no daughter is used in more than one V0 candidate

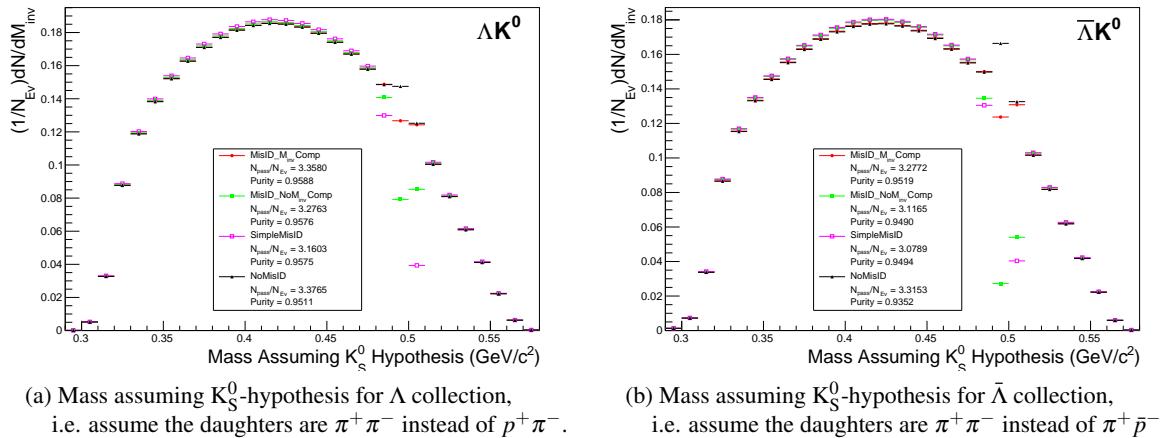


Fig. 2: Mass assuming K_S^0 -hypothesis for V0 candidates passing all Λ (2a) and $\bar{\Lambda}$ (2b) cuts. The “NoMisID” distribution (black triangles) uses the V0 finder without any attempt to remove misidentified K_S^0 . The slight peak in the “NoMisID” distribution around $m_{\text{inv}} = 0.5 \text{ GeV}/c^2$ contains misidentified K_S^0 particles in our $\Lambda(\bar{\Lambda})$ collection. “SimpleMisID” (pink squares) simply cuts out the entire peak, which throws away some good Λ and $\bar{\Lambda}$ particles. “MisID_NoM_{inv}Comp” (green squares) uses the misidentification cut outlined in the text, but does not utilize the final invariant mass comparison step. “MisID_M_{inv}Comp” (red circles) utilizes the full misidentification methods, and is currently used for this analysis. “ $N_{\text{pass}}/N_{\text{ev}}$ ” is the total number of $\Lambda(\bar{\Lambda})$ particles found, normalized by the total number of events. The purity of the collection is also listed.

| Λ selection | |
|---|--|
| $ \eta $ | < 0.8 |
| p_T | > 0.4 GeV/c |
| $ m_{\text{inv}} - m_{\text{PDG}} $ | < 3.8 MeV |
| DCA to prim. vertex | < 0.5 cm |
| Cosine of pointing angle | > 0.9993 |
| OnFlyStatus | false |
| Decay Length | < 60 cm |
| Shared Daughter Cut | true |
| Misidentification Cut | true |
| Daughter Cuts (π and p) | |
| $ \eta $ | < 0.8 |
| SetTPCnclsDaughters | (80) |
| SetStatusDaughters | (AliESDtrack::kTPCrefit) |
| DCA πp Daughters | < 0.4 cm |
| π-specific cuts | |
| p_T | > 0.16 GeV/c |
| DCA to prim vertex | > 0.3 cm |
| TPC and TOF N σ Cuts | |
| $p < 0.5 \text{ GeV}/c$ | $N\sigma_{\text{TPC}} < 3$ |
| $p > 0.5 \text{ GeV}/c$ | if TOF & TPC available else |
| | $N\sigma_{\text{TPC}} < 3 \& N\sigma_{\text{TOF}} < 3$ $N\sigma_{\text{TOF}} < 3$ |
| p-specific cuts | |
| p_T | > 0.5(p) [0.3(\bar{p})] GeV/c |
| DCA to prim vertex | > 0.1 cm |
| TPC and TOF N σ Cuts | |
| $p < 0.8 \text{ GeV}/c$ | $N\sigma_{\text{TPC}} < 3$ |
| $p > 0.8 \text{ GeV}/c$ | if TOF & TPC available else |
| | $N\sigma_{\text{TPC}} < 3 \& N\sigma_{\text{TOF}} < 3$ $N\sigma_{\text{TOF}} < 3$ |

Table 2: Λ selection

Figure 2a shows the mass assuming K_S^0 hypothesis for the Λ collection, i.e. assume the daughters are $\pi^+\pi^-$ instead of $p^+\pi^-$. Figure 2b is a similar plot, but is for the $\bar{\Lambda}$ collection, i.e. assume the daughters are $\pi^+\pi^-$ instead of $\pi^+\bar{p}^-$. The K_S^0 contamination is visible, although not profound, in both, in the slight peaks around $m_{\text{inv}} = 0.497 \text{ GeV}/c^2$. If one simply cuts out the entire peak, good Λ particles will be lost. Ideally, the Λ selection and K_S^0 misidentification cuts are selected such that the peak is removed from this plot while leaving the underlying distribution continuous. To attempt to remove these K_S^0 contaminations without throwing away good Λ and $\bar{\Lambda}$ particles, the following misidentification cuts are imposed; a $\Lambda(\bar{\Lambda})$ candidate is rejected if all of the following criteria are satisfied:

- $|m_{\text{inv}, K_S^0 \text{ Hypothesis}} - m_{\text{PDG}, K_S^0}| < 9.0 \text{ MeV}/c^2$
- Positive and negative daughters pass π daughter cut implemented for K_S^0 reconstruction
- $|m_{\text{inv}, K_S^0 \text{ Hypothesis}} - m_{\text{PDG}, K_S^0}| < |m_{\text{inv}, \Lambda(\bar{\Lambda}) \text{ Hypothesis}} - m_{\text{PDG}, \Lambda(\bar{\Lambda})}|$

Figure 3 shows the invariant mass (m_{inv}) distribution of all $\Lambda(\bar{\Lambda})$ candidates immediately before the final invariant mass cut. These distributions are used to calculate the collection purities. The Λ and $\bar{\Lambda}$ purities are found to be: Purity(Λ) \approx Purity($\bar{\Lambda}$) $\approx 95\%$.

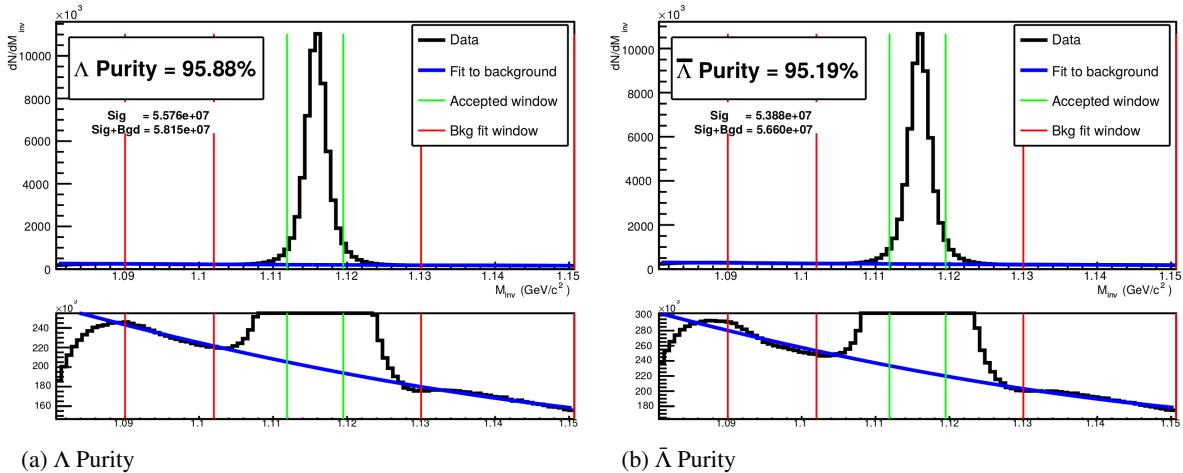


Fig. 3: Invariant mass (m_{inv}) distribution for all Λ (a) and $\bar{\Lambda}$ (b) candidates immediately before the final invariant mass cut. The bottom figures are zoomed to show the background with fit. The vertical green lines represent the m_{inv} cuts used in the analyses, the red vertical lines delineate the regions over which the background was fit, and the blue line shows the background fit. These distributions are used to calculate the collection purities, Purity(Λ) \approx Purity($\bar{\Lambda}$) $\approx 95\%$.

3.3.2 K_S^0 Reconstruction

The following cuts were used to select good K_S^0 candidates:

1. Pion Daughter Cuts

- (a) $|\eta| < 0.8$
- (b) SetTPCnclsDaughters(80)
- (c) SetStatusDaughters(AliESDtrack::kTPCrefic)
- (d) DCA $\pi^+\pi^-$ Daughters $< 0.3 \text{ cm}$

- 180 (e) $p_T > 0.15 \text{ GeV}/c$
 181 (f) DCA to prim vertex $> 0.3 \text{ cm}$
 182 (g) TPC and TOF $N\sigma$ Cuts
 183 i. $p < 0.5 \text{ GeV}/c : N\sigma_{\text{TPC}} < 3$
 184 ii. $p > 0.5 \text{ GeV}/c :$
 185 – if TOF & TPC available: $N\sigma_{\text{TPC}} < 3 \text{ & } N\sigma_{\text{TOF}} < 3$
 186 – else $N\sigma_{\text{TPC}} < 3$

187 2. K_S^0 Cuts

- 188 (a) $|\eta| < 0.8$
 189 (b) $p_T > 0.2 \text{ GeV}/c$
 190 (c) $m_{\text{PDG}} - 13.677 \text{ MeV} < m_{\text{inv}} < m_{\text{PDG}} + 2.0323 \text{ MeV}$
 191 (d) DCA to prim. vertex $< 0.3 \text{ cm}$
 192 (e) Cosine of pointing angle > 0.9993
 193 (f) OnFlyStatus = false
 194 (g) Decay Length $< 30 \text{ cm}$

195 3. Shared Daughter Cut for V0 Collection

- 196 – Iterate through V0 collection to ensure that no daughter is used in more than one V0 candidate

| K_S^0 selection | | | | | |
|--|--|------------------------|--|------|----------------------------|
| $ \eta $ | < 0.8 | | | | |
| p_T | $> 0.2 \text{ GeV}/c$ | | | | |
| $m_{\text{PDG}} - 13.677 \text{ MeV} < m_{\text{inv}} < m_{\text{PDG}} + 2.0323 \text{ MeV}$ | | | | | |
| DCA to prim. vertex | $< 0.3 \text{ cm}$ | | | | |
| Cosine of pointing angle | > 0.9993 | | | | |
| OnFlyStatus | false | | | | |
| Decay Length | $< 30 \text{ cm}$ | | | | |
| Shared Daughter Cut | true | | | | |
| Misidentification Cut | true | | | | |
| π^\pm Daughter Cuts | | | | | |
| $ \eta $ | < 0.8 | | | | |
| SetTPCnclsDaughters | (80) | | | | |
| SetStatusDaughters | (AliESDtrack::kTPCrefit) | | | | |
| DCA $\pi^+ \pi^-$ Daughters | $< 0.3 \text{ cm}$ | | | | |
| p_T | $> 0.15 \text{ GeV}/c$ | | | | |
| DCA to prim vertex | $> 0.3 \text{ cm}$ | | | | |
| TPC and TOF $N\sigma$ Cuts | | | | | |
| $p < 0.5 \text{ GeV}/c$ | $N\sigma_{\text{TPC}} < 3$ | | | | |
| $p > 0.5 \text{ GeV}/c$ | <table border="1"> <tr> <td>if TOF & TPC available</td><td>$N\sigma_{\text{TPC}} < 3 \text{ & } N\sigma_{\text{TOF}} < 3$</td></tr> <tr> <td>else</td><td>$N\sigma_{\text{TOF}} < 3$</td></tr> </table> | if TOF & TPC available | $N\sigma_{\text{TPC}} < 3 \text{ & } N\sigma_{\text{TOF}} < 3$ | else | $N\sigma_{\text{TOF}} < 3$ |
| if TOF & TPC available | $N\sigma_{\text{TPC}} < 3 \text{ & } N\sigma_{\text{TOF}} < 3$ | | | | |
| else | $N\sigma_{\text{TOF}} < 3$ | | | | |

Table 3: K_S^0 selection

197 As can be seen in Figure 4, some misidentified Λ and $\bar{\Lambda}$ particles contaminate our K_S^0 sample. Figure
 198 4a shows the mass assuming Λ -hypothesis for the K_S^0 collection, i.e. assume the daughters are $p^+ \pi^-$

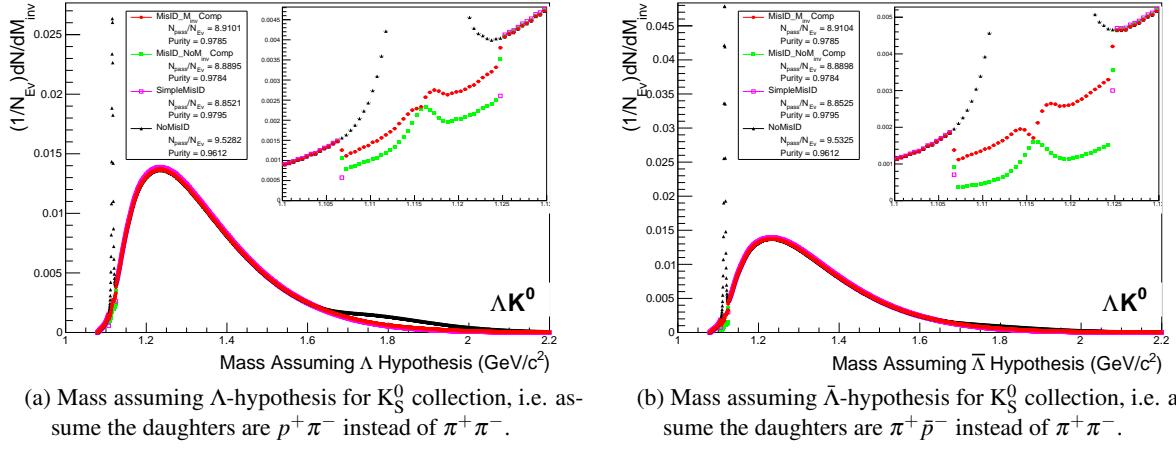


Fig. 4: Mass assuming Λ -hypothesis (4a) and $\bar{\Lambda}$ -hypothesis (4b) for K_S^0 collection. The “NoMisID” distribution (black triangles) uses the V0 finder without any attempt to remove misidentified Λ and $\bar{\Lambda}$. The peak in the “NoMisID” distribution around $m_{\text{inv}} = 1.115 \text{ GeV}/c^2$ contains misidentified Λ (4a) and $\bar{\Lambda}$ (4b) particles in our K_S^0 collection. “SimpleMisID” (pink squares) simply cuts out the entire peak, which throws away some good K_S^0 particles. “MisID_NoM_{inv}Comp” (green squares) uses the misidentification cut outlined in the text, but does not utilize the final invariant mass comparison step. “MisID_M_{inv}Comp” (red circles) utilizes the full misidentification methods, and is currently used for this analysis. “N_{pass}/N_{ev}” is the total number of K_S^0 particles found, normalized by the total number of events. The purity of the collection is also listed. Also note, the relative excess of the “NoMisID” distribution around $1.65 < m_{\text{inv}} < 2.1 \text{ GeV}/c^2$ shows misidentified $\bar{\Lambda}$ (4a) and Λ (4b) particles in our K_S^0 collection.

instead of $\pi^+\pi^-$. Figure 4b is similar, but shows the mass assuming $\bar{\Lambda}$ -hypothesis for the collection, i.e. assume the daughters are $\pi^+\bar{p}^-$ instead of $\pi^+\pi^-$. The Λ contamination can be seen in 4a, and the $\bar{\Lambda}$ contamination in 4b, in the peaks around $m_{\text{inv}} = 1.115 \text{ GeV}/c^2$. Additionally, the $\bar{\Lambda}$ contamination is visible in Figure 4a, and the Λ contamination visible in Figure 4b, in the region of excess around $1.65 < m_{\text{inv}} < 2.1 \text{ GeV}/c^2$. This is confirmed as the number of misidentified Λ particles in the sharp peak of Figure 4a (misidentified $\bar{\Lambda}$ particles in the sharp peak of Figure 4b) approximately equals the excess found in the $1.65 < m_{\text{inv}} < 2.1 \text{ GeV}/c^2$ region of Figure 4a (Figure 4b).

The peaks around $m_{\text{inv}} = 1.115 \text{ GeV}/c^2$ in Figure 4 contain both misidentified $\Lambda(\bar{\Lambda})$ particles and good K_S^0 . If one simply cuts out the entire peak, some good K_S^0 particles will be lost. Ideally, the K_S^0 selection and $\Lambda(\bar{\Lambda})$ misidentification cuts can be selected such that the peak is removed from this plot while leaving the underlying distribution continuous. To attempt to remove these Λ and $\bar{\Lambda}$ contaminations without throwing away good K_S^0 particles, the following misidentification cuts are imposed; a K_S^0 candidate is rejected if all of the following criteria are satisfied (for either Λ - or $\bar{\Lambda}$ -hypothesis):

- $|m_{\text{inv}, \Lambda(\bar{\Lambda}) \text{ Hypothesis}} - m_{\text{PDG}, \Lambda(\bar{\Lambda})}| < 9.0 \text{ MeV}/c^2$
- Positive daughter passes $p^+(\pi^+)$ daughter cut implemented for $\Lambda(\bar{\Lambda})$ reconstruction
- Negative daughter passes $\pi^-(\bar{p}^-)$ daughter cut implemented by $\Lambda(\bar{\Lambda})$ reconstruction
- $|m_{\text{inv}, \Lambda(\bar{\Lambda}) \text{ Hypothesis}} - m_{\text{PDG}, \Lambda(\bar{\Lambda})}| < |m_{\text{inv}, K_S^0 \text{ Hypothesis}} - m_{\text{PDG}, K_S^0}|$

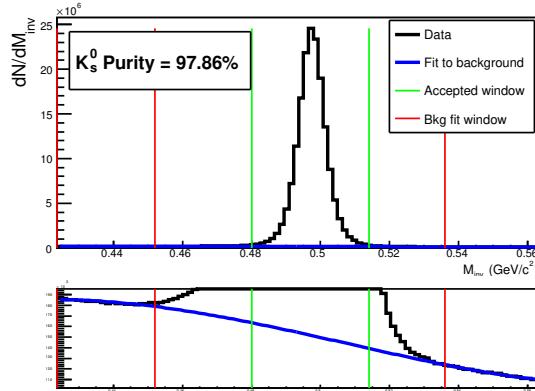


Fig. 5: Invariant mass (m_{inv}) distribution for all K_s^0 candidates immediately before the final invariant mass cut. The bottom figure is zoomed to show the background with fit. The vertical green lines represent the m_{inv} cut used in the analyses, the red vertical lines delineate the regions over which the background was fit, and the blue line shows the background fit. This distribution is used to calculate the collection purity, $\text{Purity}(K_s^0) \approx 98\%$.

216 3.3.3 V0 Purity Background Estimation

217 As previously stated, the backgrounds in the m_{inv} distributions are modeled by a polynomial which is
 218 fit outside of the final cut region in an attempt to estimate the background within the cut region. As
 219 this estimate of the background under the mass peak is vital for our estimate of our V0 purity, it is
 220 important for us to ensure that our estimate is accurate. More specifically, it is necessary that we ensure
 221 the background is well described by a polynomial fit within the cut region.

222 To better understand our background, we studied V0 candidates reconstructed with daughters from dif-
 223 ferent events. These mixed-event V0s certainly do not represent real, physical V0s (a single V0 cannot
 224 have daughters living in two different events!), but, rather, represent a large portion of the background
 225 creeping into our analysis.

226 The standard AliFemto framework is not equipped to handle this situation, as most are not interested in
 227 these fake-V0s. Therefore, we built the `AliFemtoV0PurityBgdEstimator` class to handle our needs. In ad-
 228 dition to finding fake-V0s using mixed-event daughters, we also used our `AliFemtoV0PurityBgdEstimator`
 229 class to find real-V0s using same-event daughters. The purpose here was to compare our simple V0 finder
 230 (in `AliFemtoV0PurityBgdEstimator`) to the established V0 finder used in standard AliFemto analyses.

231 Figure 6 shows the results of our study. In the figures, the black points, marked "Data", correspond to
 232 V0s found using the standard V0-finder, and to the V0s used in my analyses. The red points show real
 233 V0s reconstructed with our personal V0-finder (in `AliFemtoV0PurityBgdEstimator`) using same-event
 234 daughters, and the blue points show fake-V0s reconstructed with our personal V0-finder using mixed-
 235 event daughters. Both the red and blue points have been scaled by different factors (listed in the figure's
 236 legends) to nicely align all three data on a single plot.

237 Figure 6 shows that our personal V0-finder does a good, but not perfect, job of matching the shape of the
 238 m_{inv} plots obtained from the data. The scale factor listed in the legend reveals that we are only finding
 239 1/3 - 1/2 of the V0s found by the standard V0-finder. These two points are not of concern, as our purpose
 240 here was to gain a sense of the broad shape of the background. It is revealed in Fig. 6, when studying
 241 the red and blue points, that the background distribution within the mass peak region is simply a smooth
 242 connection of the backgrounds outside of the cut region, as we assumed. Therefore, our method of fitting
 243 the background outside of the cut region, fitting with a smooth polynomial, and extrapolating to the cut
 244 region is justified.

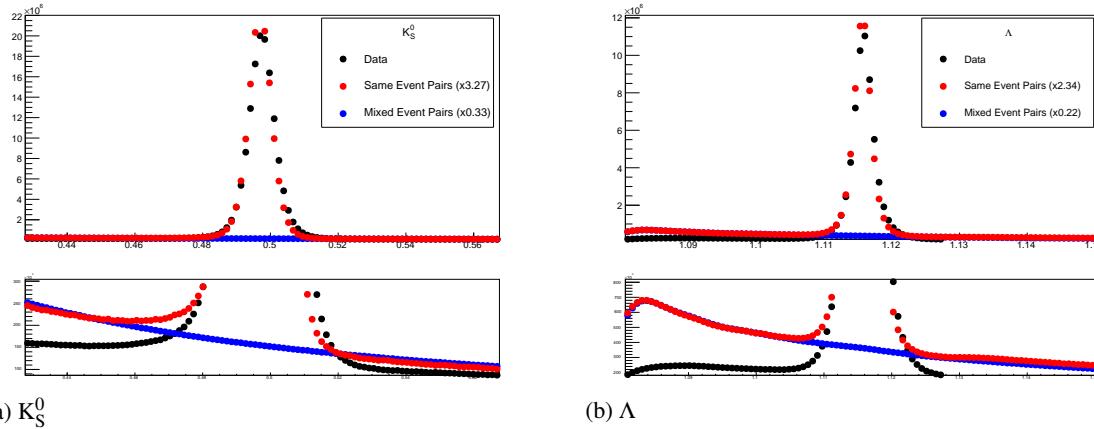


Fig. 6: V0 Purity Background Estimation. The black points, marked "Data", correspond to real V0s found using the standard V0-finder (i.e. the V0s used in my analyses). The red points, marked "Same Event Pairs", show real V0s reconstructed with our personal V0-finder in AliFemtoV0PurityBgdEstimator. These data are scaled by a factor (listed in the legend) to match their *Signal + Background* value in the cut region with that of the data. The blue points, marked "Mixed Event Pairs", show fake-V0s reconstructed with our personal V0-finder using mixed-event daughters. The blue points are scaled by a factor (listed in the legend) to closely match the red points in the side-band region.

3.4 Σ Reconstruction

Our motivation for studying Ξ^-K^\pm systems is to attempt to better understand the striking difference in the ΛK^+ and ΛK^- data at low k^* (Figure 13).

The reconstruction of Ξ particles is one level above V0 reconstruction. V0 particles are topologically reconstructed by searching for the charged daughters' tracks into which they decay. With Ξ particles, we search for the V0 particle and charged daughter into which the Ξ decays. In the case of Ξ^- , we search for the Λ (V0) and π^- (track) daughters. We will refer to this π as the “bachelor π ”.

The following cuts were used to select good Ξ^- ($\bar{\Xi}^+$) candidates:

1. V0 Daughter Reconstruction

(a) V0 Daughter Particle Cuts

i. Cuts Common to Both Daughters

- A. $|\eta| < 0.8$
 - B. SetTPCnclsDaughters(80)
 - C. SetStatusDaughters(AliESDtrack::kTPCrefic)
 - D. SetMaxDcaV0Daughters(0.4)

ii. Pion Specific Daughter Cuts

- A. $p_T > 0.16$
 - B. DCA to prim vertex > 0.3

iii. Proton Specific Daughter Cuts

- A. $p_T > 0.5(p) [0.3(\bar{p})]$ GeV/c
 - B. DCA to prim vertex > 0.1

(b) V0 Cuts

- i. $|\eta| < 0.8$

- 268 ii. $p_T > 0.4 \text{ GeV}/c$
 269 iii. $|m_{inv} - m_{PDG}| < 3.8 \text{ MeV}$
 270 iv. DCA to prim. vertex $> 0.2 \text{ cm}$
 271 v. Cosine of pointing angle to Ξ decay vertex > 0.9993
 272 vi. OnFlyStatus = false
 273 vii. Decay Length $< 60 \text{ cm}$
 274 viii. The misidentification cuts described in Section 3.3.1 are utilized

275 2. Bachelor π Cuts

- 276 (a) $|\eta| < 0.8$
 277 (b) $p_T < 100 \text{ GeV}/c$
 278 (c) DCA to prim vertex $> 0.1 \text{ cm}$
 279 (d) SetTPCnclsDaughters(70)
 280 (e) SetStatusDaughters(AliESDtrack::kTPCrefic)

281 3. Ξ Cuts

- 282 (a) $|\eta| < 0.8$
 283 (b) $0.8 < p_T < 100 \text{ GeV}/c$
 284 (c) $|m_{inv} - m_{PDG}| < 3.0 \text{ MeV}$
 285 (d) DCA to prim. vertex $< 0.3 \text{ cm}$
 286 (e) Cosine of pointing angle > 0.9992

287 4. Shared Daughter Cut for Ξ Collection

- 288 – Iterate through Ξ collection to ensure that no daughter is used in more than one Ξ candidate
- 289 – Remove any candidate in which the bachelor π is also a daughter of the Λ (implemented in AliFemtoXiTrackPairCut class)

291 The purity of our Ξ and $\bar{\Xi}$ collections are calculated just as those of our V0 collections 3.3. Figure 8,
 292 which is used to calculate the purity, shows the m_{inv} distribution of our $\Xi(\bar{\Xi})$ candidates just before the
 293 final m_{inv} cut. Currently, we have Purity(Ξ^-) $\approx 90\%$ and Purity($\bar{\Xi}^+$) $\approx 92\%$.

294 **3.5 Pair Selection**

295 It is important to obtain true particle pairs in the analysis. In particular, contamination from pairs con-
 296 structed with split or merged tracks, and pairs sharing daughters, can introduce an artificial signal into
 297 the correlation function, obscuring the actual physics. We impose the following pair cuts to combat these
 298 issues:

299 1. Shared Daughter Cut for Pairs

- 300 (a) V0-V0 Pairs (i.e. ΛK_S^0 analyses)
 - 301 – Remove all pairs which share a daughter
 - 302 – Ex. Λ and K_S^0 particles which share a π^- daughter are not included

303 (b) V0-Track Pairs (i.e. ΛK^\pm analyses)

 - 304 – Remove pairs if Track is also used as a daughter of the V0

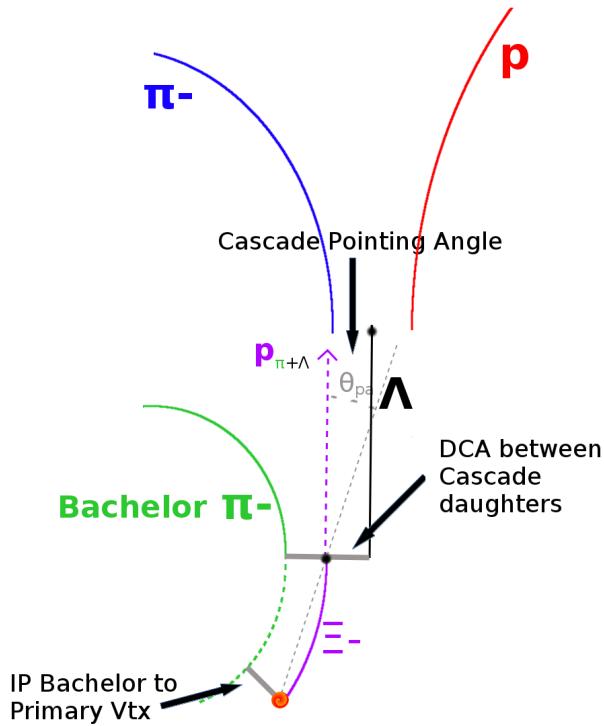


Fig. 7: Ξ Reconstruction

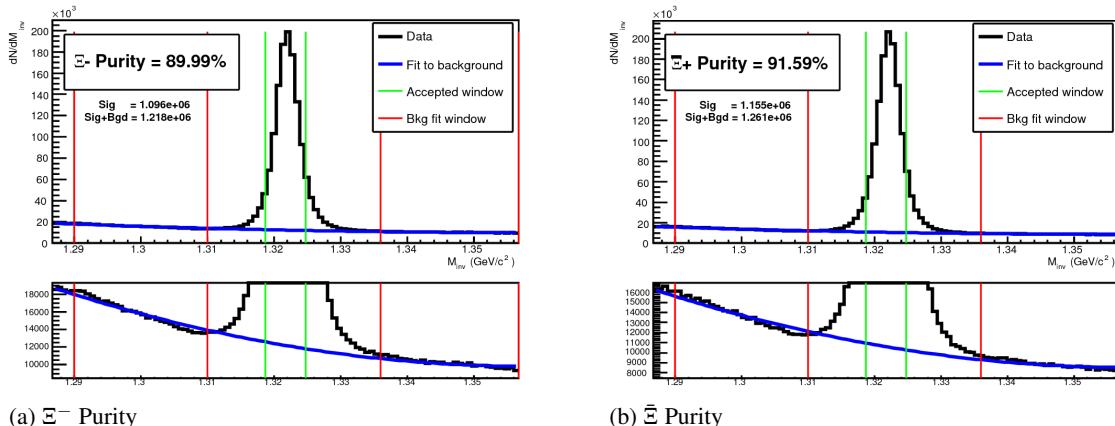


Fig. 8: Invariant mass (m_{inv}) distribution for all Ξ^- (a) and Ξ^+ (b) candidates immediately before the final invariant mass cut. The bottom figures are zoomed to show the background with fit. The vertical green lines represent the m_{inv} cuts used in the analyses, the red vertical lines delineate the regions over which the background was fit, and the blue line shows the background fit. These distributions are used to calculate the collection purities, $Purity(\Xi^-) \approx 90\%$ and $Purity(\Xi^+) \approx 92\%$.

- In these analyses, this could only occur if, for instance, a K is misidentified as a π or p in the V0 reconstruction

307 (c) Ξ -Track Pairs

- Remove pairs if Track is also used as a daughter of the Ξ
- In these analyses, this could only occur if, for instance, a K is misidentified as a π or p in the V0 reconstruction, or misidentified as bachelor π .

311 2. Average Separation Cuts (AvgSep)

- 312 – Used to cut out splitting and merging effects
 313 – The motivation for these cuts can be seen in Figures 9, 10, and 11, in which average separa-
 314 tion correlation functions are presented
- 315 (a) ΛK_S^0 Analyses
 316 – AvgSep > 6.0 cm for like charge sign daughters
 317 – ex. p daughter of Λ and π^+ daughter of K_S^0
 318 – No cut for unlike-sign daughters
- 319 (b) ΛK^\pm Analyses
 320 – AvgSep > 8.0 cm for daughter of $\Lambda(\bar{\Lambda})$ sharing charge sign of K^\pm
 321 – ex. in ΛK^+ analysis, p daughter of Λ with K^+
 322 – No cut for unlike signs
- 323 (c) $\Xi^- K^\pm$ Analyses
 324 – AvgSep > 8.0 cm for any daughter of Ξ sharing charge sign of K^\pm
 325 – ex. in $\Xi^- K^-$ analysis, π^- daughter of Λ daughter with K^- , and bachelor π^- daugh-
 326 ter with K^-
 327 – No cut for unlike signs

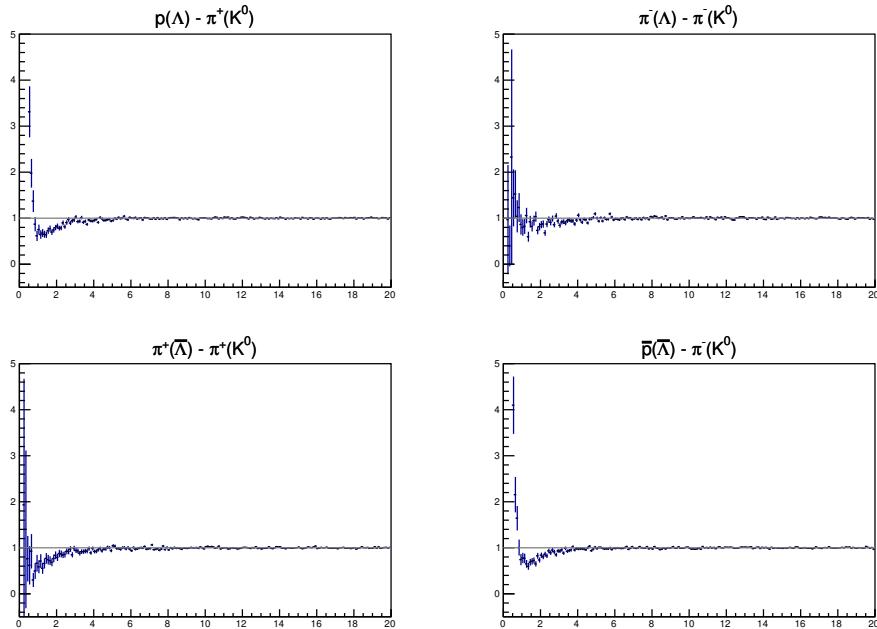


Fig. 9: Average separation (cm) correlation functions of $\Lambda(\bar{\Lambda})$ and K_S^0 Daughters. Only like-sign daughter pairs are shown (the distributions for unlike-signs were found to be flat). The title of each subfigure shows the daughter pair, as well as the mother of each daughter (in “()”), ex. top left is p from Λ with π^+ from K_S^0 .

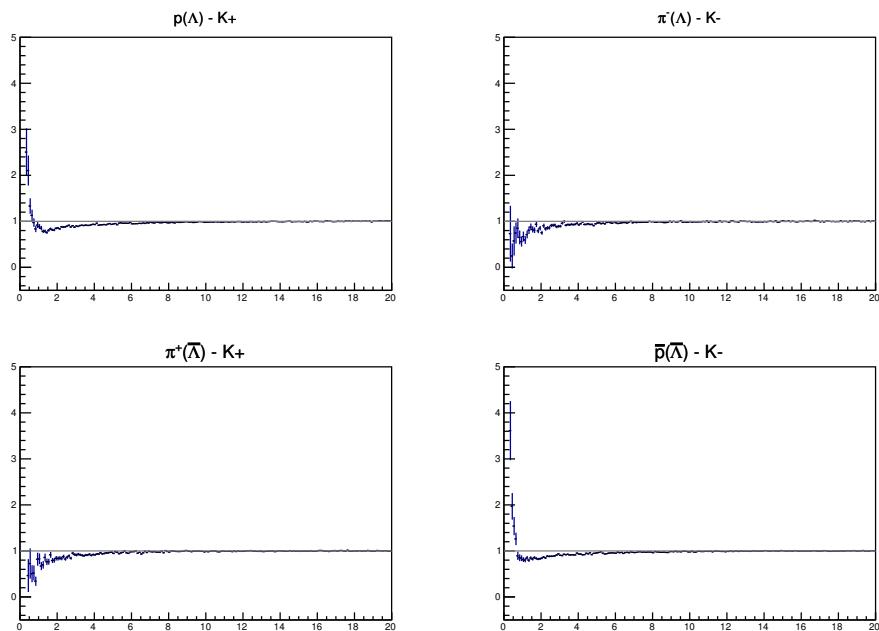


Fig. 10: Average separation (cm) correlation functions of $\Lambda(\bar{\Lambda})$ Daughter and K^\pm . Only like-sign pairs are shown (unlike-signs were flat). In the subfigure titles, the particles in “()” represent the mothers, ex. top left is p from Λ with K^+ .

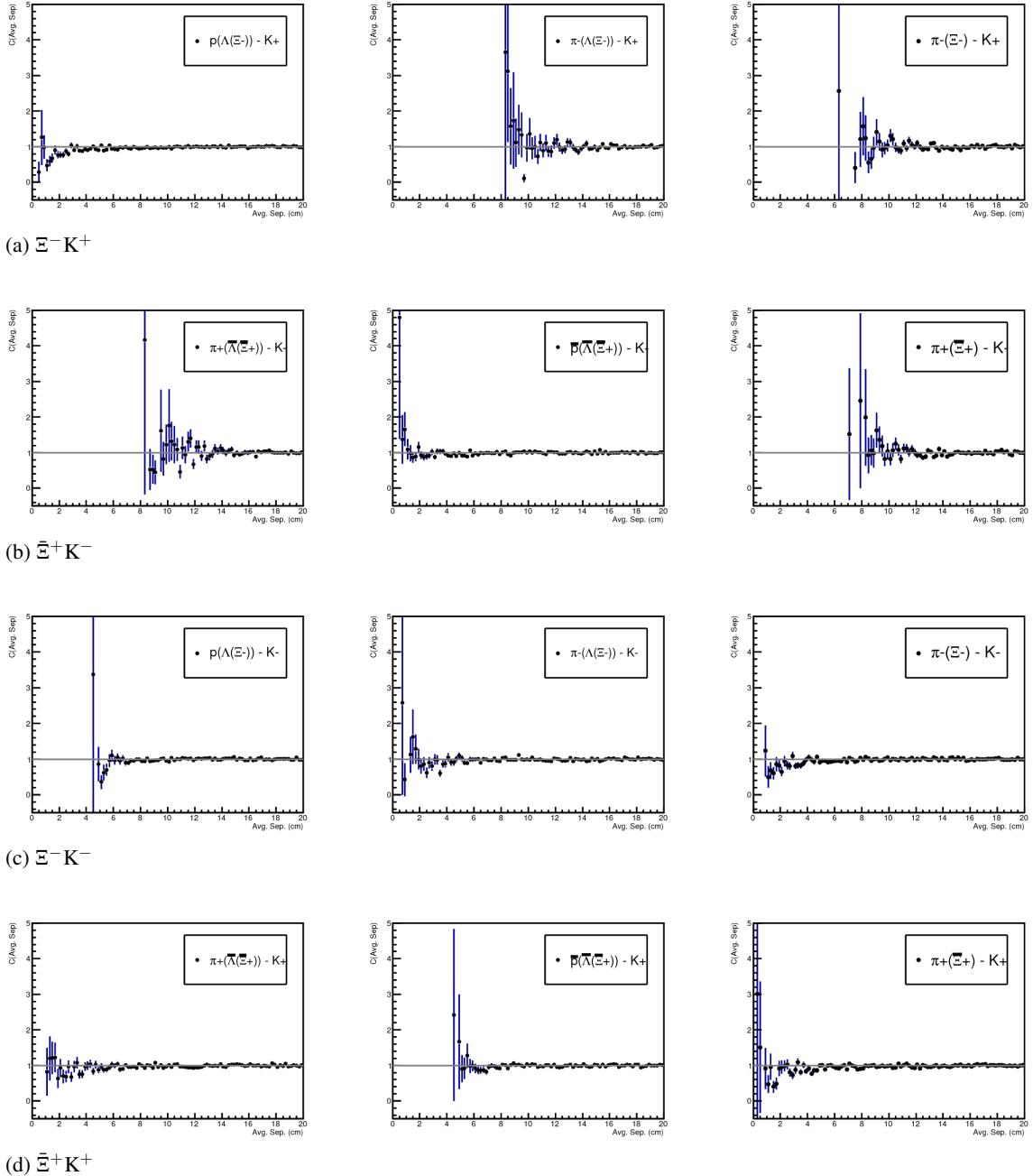


Fig. 11: Average separation (cm) correlation functions of Ξ Daughter and K^\pm . In the subfigure titles, the particles in “()” represent the mothers, ex. top left is p from Λ from Ξ^- with K^+ .

328 **4 Correlation Functions**

329 This analysis studies the momentum correlations of both ΛK and $\Xi^- K^\pm$ pairs using the two-particle
 330 correlation function, defined as $C(k^*) = A(k^*)/B(k^*)$, where $A(k^*)$ is the signal distribution, $B(k^*)$ is the
 331 reference (or background) distribution, and k^* is the momentum of one of the particles in the pair rest
 332 frame. In practice, $A(k^*)$ is constructed by binning in k^* pairs from the same event. Ideally, $B(k^*)$ is
 333 similar to $A(k^*)$ in all respects excluding the presence of femtoscopic correlations [2]; as such, $B(k^*)$
 334 is used to divide out the phase-space effects, leaving only the femtoscopic effects in the correlation
 335 function.

336 This analysis presents correlation functions for three centrality bins (0-10%, 10-30%, and 30-50%), and
 337 is currently pair transverse momentum ($k_T = 0.5|\mathbf{p}_{T,1} + \mathbf{p}_{T,2}|$) integrated (i.e. not binned in k_T). The
 338 correlation functions are constructed separately for the two magnetic field configurations, and, after
 339 assuring consistency, are combined using a weighted average:

$$C_{combined}(k^*) = \frac{\sum_i w_i C_i(k^*)}{\sum_i w_i} \quad (3)$$

340 where the sum runs over the correlation functions to be combined, and the weight, w_i , is the number of
 341 numerator pairs in $C_i(k^*)$. Here, the sum is over the two field configurations (++ and - -).

342 **4.1 Typical Correlation Function Construction**

343 In practice, $B(k^*)$ is typically obtained by forming mixed-event pairs, i.e. particles from a given event
 344 are paired with particles from $N_{mix}(= 5)$ other events, and these pairs are then binned in k^* . In forming
 345 the background distribution, it is important to mix only similar events; mixing events with different
 346 phase-spaces can result in an unreliable background distribution, and can introduce artificial signals in
 347 the correlation function. Therefore, in this analysis, we bin our events both in primary vertex location (2
 348 cm bin width) and in centrality (5% bin width), and we only mix events within a given bin; i.e. we only
 349 mix events of like centrality and of like primary vertex location. Also note, a vertex correction is also
 350 applied to each event, which essentially recenters the the primary vertices to $z = 0$.

351 Figures 12a, 12b, 12c show the correlation functions for all centralities studied for $\Lambda K^+(\bar{\Lambda} K^-)$, $\Lambda K^-(\bar{\Lambda} K^+)$,
 352 and $\Lambda(\bar{\Lambda})K_S^0$, respectively. All were normalized in the range $0.32 < k^* < 0.4$ GeV/c. It is interesting to
 353 note that the average of the $\Lambda K^+(\bar{\Lambda} K^-)$ and $\Lambda K^-(\bar{\Lambda} K^+)$ correlation functions is consistent with our
 354 $\Lambda K_S^0(\bar{\Lambda} K_S^0)$ measurement.

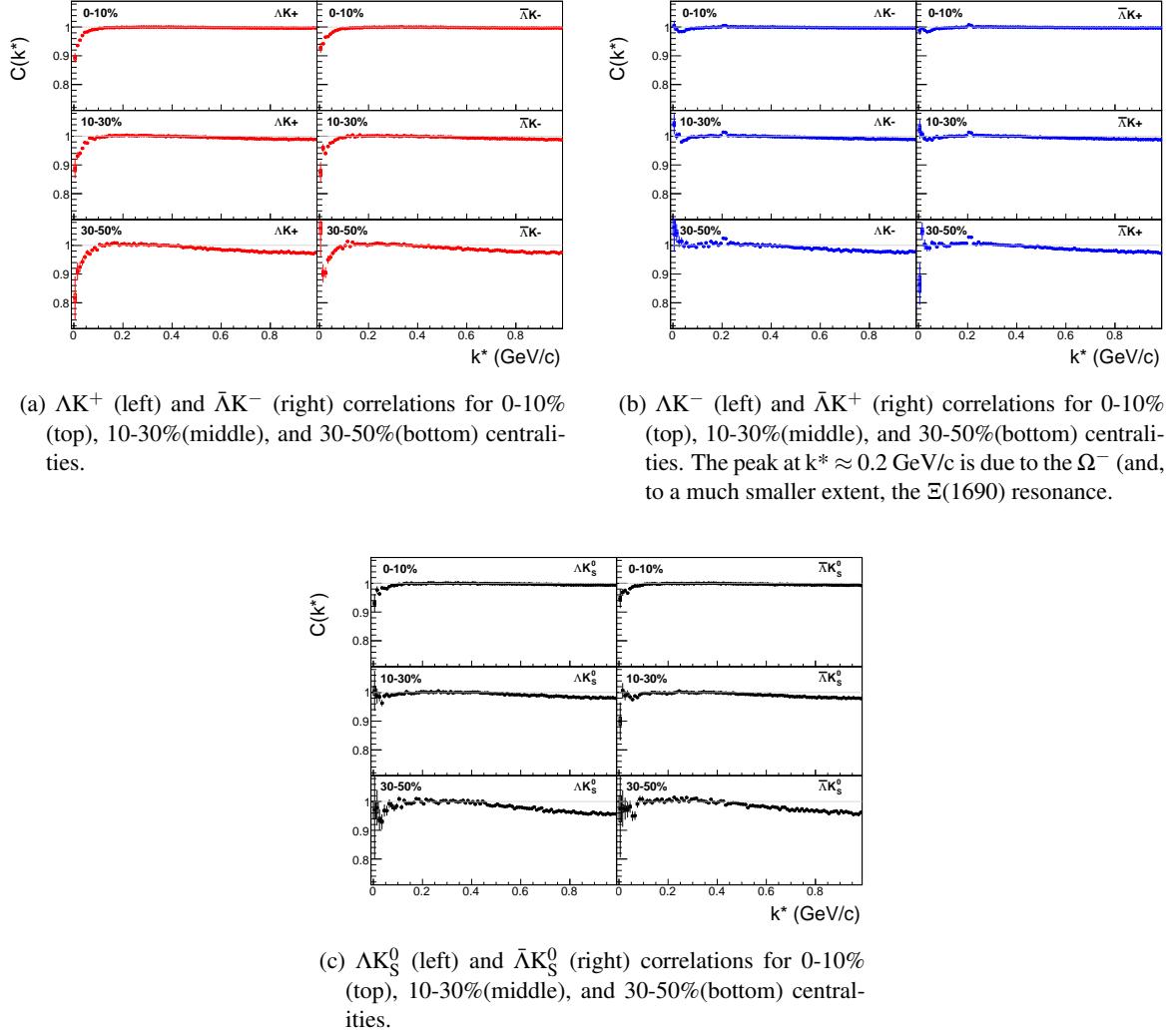


Fig. 12: ΛK and $\bar{\Lambda}K$ correlation functions for 0-10%, 10-30%, and 30-50% centralities. The lines represent the statistical errors, while the boxes represent the systematic errors.

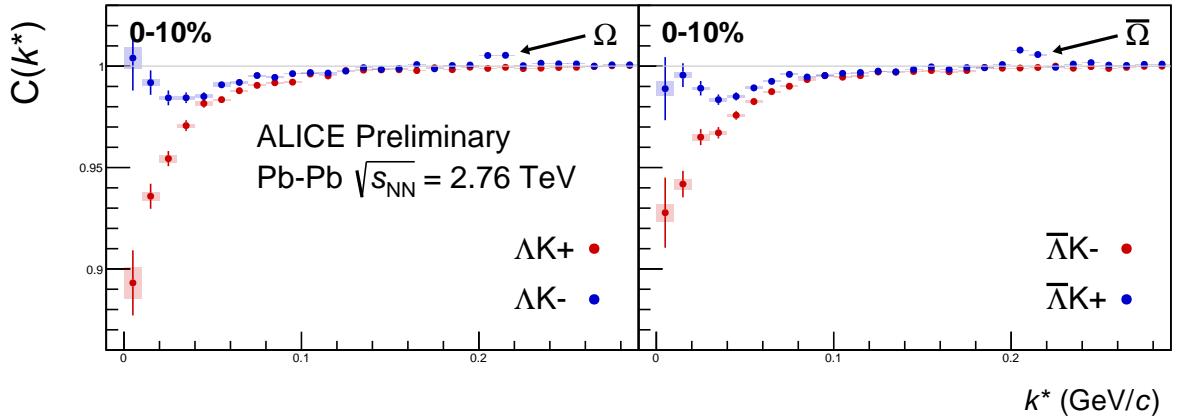


Fig. 13: Correlation Functions: ΛK^+ vs ΛK^- ($\bar{\Lambda}K^+$ vs $\bar{\Lambda}K^-$) for 0-10% centrality. The peak in $\Lambda K^- (\bar{\Lambda}K^+)$ at $k^* \approx 0.2$ GeV/c is due to the Ω^- (and, to a much smaller extent, the $\Xi(1690)$) resonance. The lines represent the statistical errors, while boxes represent systematic errors.

355 4.2 Stavinsky Correlation Function Construction

356 The purpose of the Stavinsky method is to rid the correlation functions of the non-femtoscopic back-
357 ground. More specifically, this method is intended to handle background contributions from elliptic
358 flow, and other sources having reflection symmetry in the transverse plane. With the Stavinsky method,
359 mixed-event pairs are not used for the background ($B(k^*)$); instead, same-event pseudo-pairs, formed
360 by rotating one particle in a real pair by 180° in the transverse plane, are used as a background. This
361 rotation rids the pairs of any femtoscopic correlation, while maintaining correlations due to elliptic flow
362 (and other suitably symmetric contributors).

363 The results of correctly implementing such a procedure are shown in Figure 14. The figure shows the
364 Stavinsky method does a very good job of ridding the ΛK^\pm correlations of their non-femtoscopic back-
365 grounds. We also see the procedure does not work as well on the ΛK_S^0 system.

366 Now, one must be somewhat careful when applying this Stavinsky method. We found that, in order to
367 obtain correct results, we had to run our pseudo-pairs through the same pair cuts used in our analyses.
368 In an ideal world, our pair cut would only remove truly bad pairs results from splitting, merging, etc. In
369 the real world, the pair cut always throws out some of the good with the bad. For the pseudo-pairs to
370 form a reliable background, they too must experience the pair cut, and the loss of “good” pseudo-pairs.
371 We found this issue affected mainly our ΛK^+ & $\bar{\Lambda} K^-$ analysis, as can be seen in Figure 15, which shows
372 both a correct implementation of the Stavinsky method, and an incorrect implementation lacking the
373 additional pair cut on the pseudo-pairs.

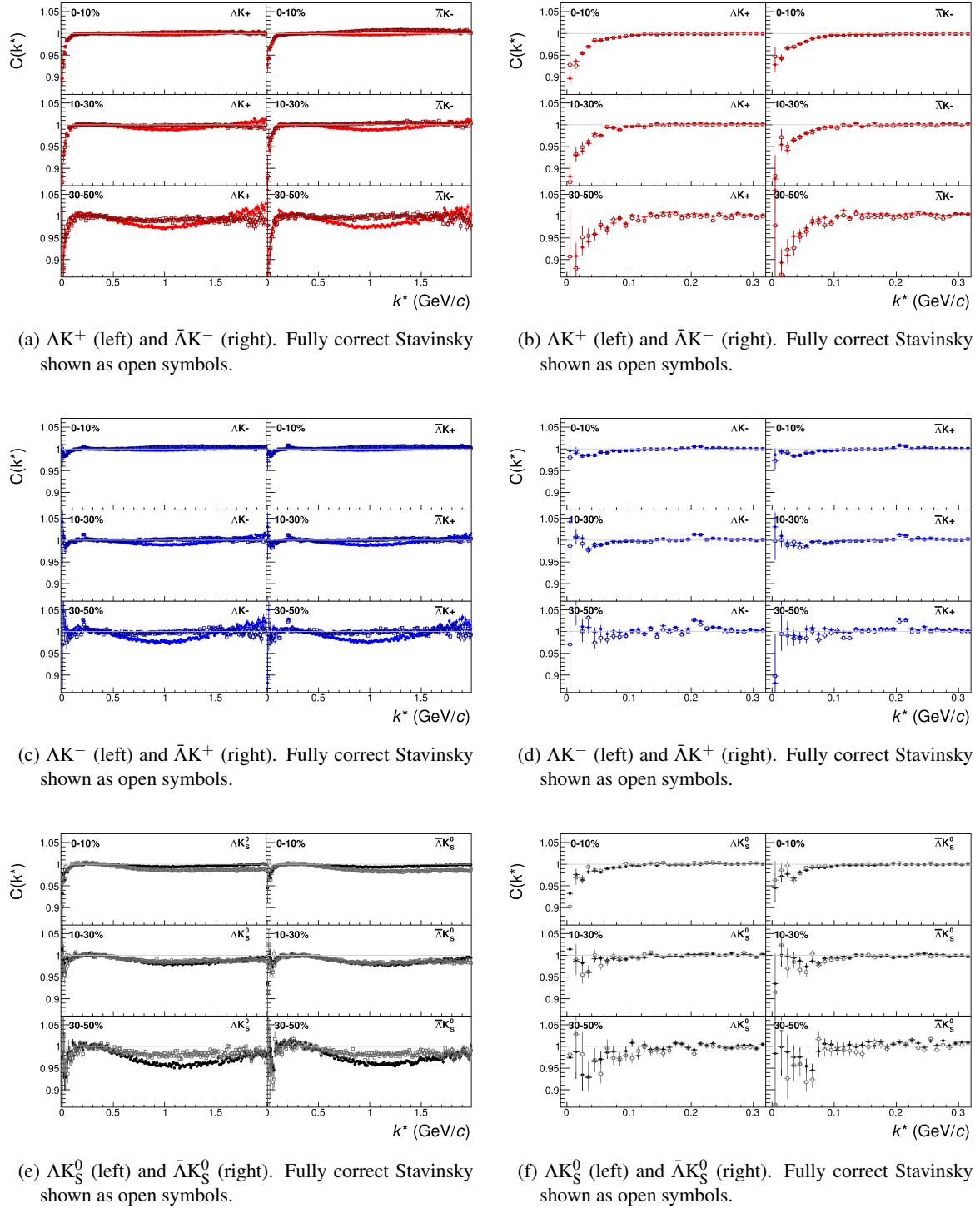
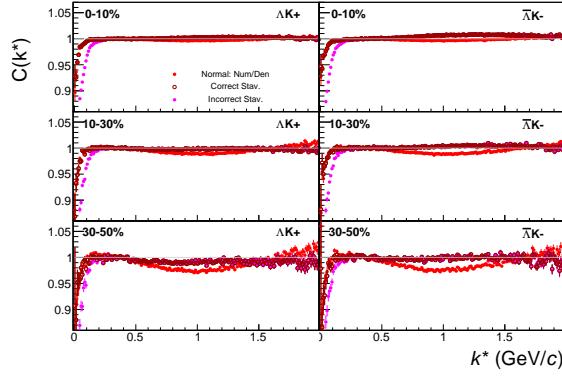
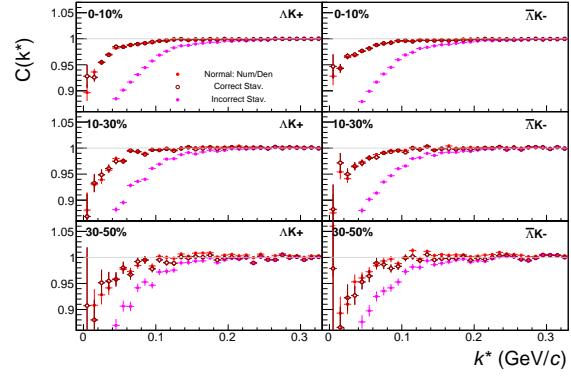


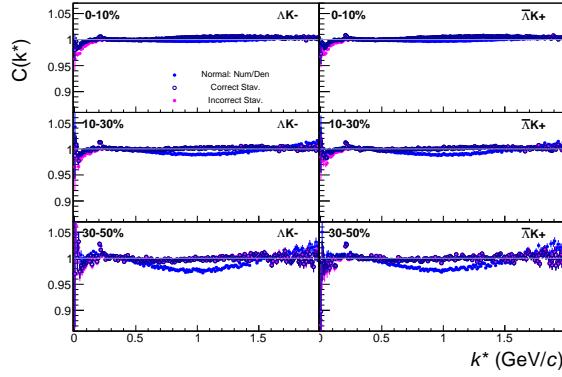
Fig. 14: ΛK and $\bar{\Lambda} \bar{K}$ correlation functions built using the fully correct Stavinsky method for 0-10%, 10-30%, and 30-50% centralities. In the fully correct method, the pseudo-pairs (same-event pairs with one particle rotated by 180° in the transverse plane) are also run through the pair cuts used in the analysis (an example of an incorrect implementation is shown in Fig. 15). Closed symbols represent correlations built using the normal mixed-event background, while open symbols represent correlations formed using the Stavinsky same-event pseudo-pairs as a background. Figures in the right column are zoomed-in versions of figures in the left column.



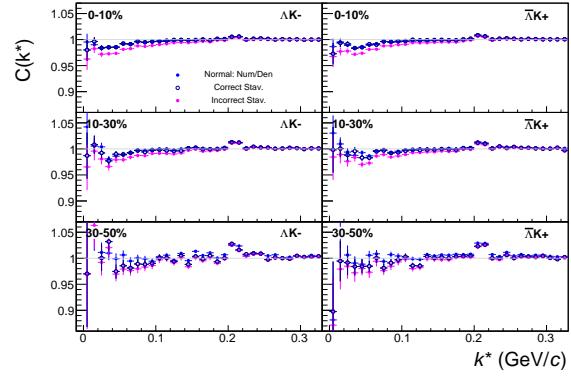
(a) ΛK^+ (left) and $\bar{\Lambda}K^-$ (right). Incorrect Stavinsky shown as open symbols.



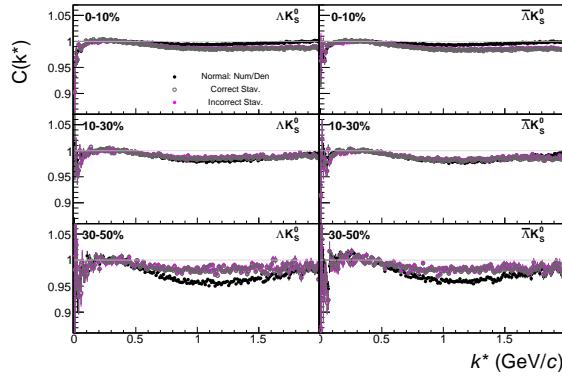
(b) ΛK^+ (left) and $\bar{\Lambda}K^-$ (right). Incorrect Stavinsky shown as open symbols.



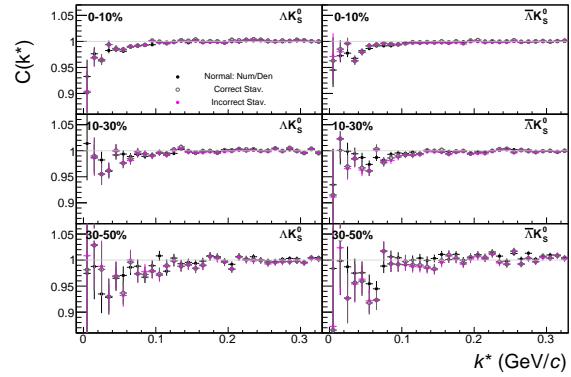
(c) ΛK^- (left) and $\bar{\Lambda}K^+$ (right). Incorrect Stavinsky shown as open symbols.



(d) ΛK^- (left) and $\bar{\Lambda}K^+$ (right). Incorrect Stavinsky shown as open symbols.



(e) ΛK_S^0 (left) and $\bar{\Lambda}K_S^0$ (right). Incorrect Stavinsky shown as open symbols.



(f) ΛK_S^0 (left) and $\bar{\Lambda}K_S^0$ (right). Incorrect Stavinsky shown as open symbols.

Fig. 15: ΛK and $\bar{\Lambda}K$ correlation functions built, both correctly and incorrectly, using the Stavinsky method for 0-10%, 10-30%, and 30-50% centralities. This figure is the same as Fig. 14, but with results from the incorrect Stavinsky implementation shown in magenta. The closed, (red, blue, black) symbols represent correlation functions formed using the normal method with mixed-event background pairs. The open, cyan symbols represent correlation functions formed using the correct Stavinsky method. The closed, magenta symbols represent correlation functions formed using the incorrect Stavinsky method. In the correct method, the pseudo-pairs (same-event pairs with one particle rotated by 180° in the transverse plane) are also run through the pair cuts used in the analysis; in the incorrect method, they are not. Figures in the right column are zoomed-in versions of figures in the left column.

374 **5 Fitting**

375 **5.1 Model: $\Lambda K_S^0, \Lambda K^\pm, \Xi^- K_S^0$**

376 The two-particle relative momentum correlation function may be written theoretically by the Koonin-
 377 Pratt equation [3, 4]:

$$C(\mathbf{k}^*) = \int S(\mathbf{r}^*) |\Psi_{\mathbf{k}^*}(\mathbf{r}^*)|^2 d^3 \mathbf{r}^* \quad (4)$$

378 where $S(\mathbf{r}^*)$ is the pair source distribution, $\Psi_{\mathbf{k}^*}(\mathbf{r}^*)$ is the two-particle wave-function, and k^* is the
 379 momentum of one particle in the pair rest frame. In the absence of Coulomb effects, and assuming a
 380 spherically Gaussian source of width R , and s-wave scattering, the 1D femtoscopic correlation function
 381 can be calculated analytically using:

$$C(k^*) = 1 + C_{QI}(k^*) + C_{FSI}(k^*) \quad (5)$$

382 C_{QI} describes plane-wave quantum interference:

$$C_{QI}(k^*) = \alpha \exp(-4k^{*2}R^2) \quad (6)$$

383 where $\alpha = (-1)^{2j}/(2j+1)$ for identical particles with spin j , and $\alpha = 0$ for non-identical particles. For
 384 all analyses presented in this note, $\alpha = 0$. C_{FSI} describes the s-wave strong final state interaction between
 385 the particles:

$$\begin{aligned} C_{FSI}(k^*) &= (1 + \alpha) \left[\frac{1}{2} \left| \frac{f(k^*)}{R} \right|^2 \left(1 - \frac{d_0}{2\sqrt{\pi}R} \right) + \frac{2\Re f(k^*)}{\sqrt{\pi}R} F_1(2k^*R) - \frac{\Im f(k^*)}{R} F_2(2k^*R) \right] \\ f(k^*) &= \left(\frac{1}{f_0} + \frac{1}{2} d_0 k^{*2} - ik^* \right)^{-1}; \quad F_1(z) = \int_0^z \frac{e^{x^2-z^2}}{z} dx; \quad F_2(z) = \frac{1-e^{-z^2}}{z} \end{aligned} \quad (7)$$

386 where R is the source size, $f(k^*)$ is the s-wave scattering amplitude, f_0 is the complex scattering length,
 387 and d_0 is the effective range of the interaction.

388 An additional parameter λ is typically included in the femtoscopic fit function to account for the purity
 389 of the pair sample. In the case of no residual correlations (to be discussed in Section 5.4), the fit function
 390 becomes:

$$C(k^*) = 1 + \lambda [C_{QI}(k^*) + C_{FSI}(k^*)] \quad (8)$$

391 **5.2 Model: $\Xi^- K^\pm$**

392 The two-particle correlation function may be written as:

$$C(\mathbf{k}^*) = \sum_S \rho_S \int S(\mathbf{r}^*) |\Psi_{\mathbf{k}^*}^S(\mathbf{r}^*)|^2 d^3 \mathbf{r}^* \quad (9)$$

393 where ρ_S is the normalized emission probability of particles in a state with spin S , $S(\mathbf{r}^*)$ is the pair
 394 emission source distribution (assumed to be Gaussian), and $\Psi_{\mathbf{k}^*}^S(\mathbf{r}^*)$ is the two-particle wave-function
 395 including both strong and Coulomb interactions [5]:

$$\Psi_{\mathbf{k}^*}(\mathbf{r}^*) = e^{i\delta_c} \sqrt{A_c(\eta)} [e^{i\mathbf{k}^*\cdot\mathbf{r}^*} F(-i\eta, 1, i\xi) + f_c(k^*) \frac{\tilde{G}(\rho, \eta)}{r^*}] \quad (10)$$

where $\rho = k^*r^*$, $\eta = (k^*a_c)^{-1}$, $\xi = \mathbf{k}^* \cdot \mathbf{r}^* + k^*r^* \equiv \rho(1 + \cos\theta^*)$, and $a_c = (\mu z_1 z_2 e^2)^{-1}$ is the two-particle Bohr radius (including the sign of the interaction). δ_c is the Coulomb s-wave phase shift, $A_c(\eta)$ is the Coulomb penetration factor, $\tilde{G} = \sqrt{A_c}(G_0 + iF_0)$ is a combination of the regular (F_0) and singular (G_0) s-wave Coulomb functions. $f_c(k^*)$ is the s-wave scattering amplitude:

$$f_c(k^*) = \left[\frac{1}{f_0} + \frac{1}{2} d_0 k^{*2} - \frac{2}{a_c} h(\eta) - ik^* A_c(\eta) \right]^{-1} \quad (11)$$

where, the “h-function”, $h(\eta)$, is expressed through the digamma function, $\psi(z) = \Gamma'(z)/\Gamma(z)$ as:

$$h(\eta) = 0.5[\psi(i\eta) + \psi(-i\eta) - \ln(\eta^2)] \quad (12)$$

In this case, the λ parameter may be included as:

$$C(\mathbf{k}^*) = (1 - \lambda) + \lambda \sum_S \rho_S \int S(\mathbf{r}^*) |\Psi_{\mathbf{k}^*}^S(\mathbf{r}^*)|^2 d^3\mathbf{r}^* \quad (13)$$

5.3 Momentum Resolution Corrections

Finite track momentum resolution causes the reconstructed momentum of a particle to smear around the true value. This, of course, also holds true for V0 particles. The effect is propagated up to the pairs of interest, which causes the reconstructed relative momentum (k_{Rec}^*) to differ from the true momentum (k_{True}^*). Smearing of the momentum typically will result in a suppression of the signal. More specifically, the smearing will broaden the signal, which would cause a decrease in the extracted radius of the system.

The effect of finite momentum resolution can be investigated using the HIJING MC data, for which both the true and reconstructed momenta are available. Figure 16 shows sample k_{True}^* vs. k_{Rec}^* plots for ΛK^\pm 0-10% analyses; Figure 16a was generated using same-event pairs, while Figure 16b was generated using mixed-event pairs (with $N_{\text{mix}} = 5$).

If there are no contaminations in our particle collection, the plots in Figure 16 should be smeared around $k_{\text{True}}^* = k_{\text{Rec}}^*$; this is mostly true in our analyses. However, there are some interesting features of our results which demonstrate a small (notice the log-scale on the z-axis) contamination in our particle collection. The structure around $k_{\text{Rec}}^* = k_{\text{True}}^* - 0.15 \text{ GeV}/c$ is mainly caused by K_S^0 contamination in our $\Lambda(\bar{\Lambda})$ sample. The remaining structure not distributed about $k_{\text{Rec}}^* = k_{\text{True}}^*$ is due to π and e contamination in our K^\pm sample. These contaminations are more visible in Figure 17, which show k_{Rec}^* vs. k_{True}^* plots (for a small sample of the ΛK^\pm 0-10% central analysis), for which the MC truth (i.e. true, known identity of the particle) was used to eliminate misidentified particles in the $K^+(a)$ and $\Lambda(b)$ collections. (NOTE: This is an old figure and is for a small sample of the data. A new version will be generated shortly. It, nonetheless, demonstrates the point well).

Information gained from looking at k_{Rec}^* vs k_{true} can be used to apply corrections to account for the effects of finite momentum resolution on the correlation functions. A typical method (“Ratio” method) involves using the MC HIJING data to build two correlation functions, $C_{\text{Rec}}(k^*)$ and $C_{\text{True}}(k^*)$, using the generator-level momentum (k_{True}^*) and the measured detector-level momentum (k_{Rec}^*). The data is then corrected by multiplying by the ratio, $C_{\text{True}}/C_{\text{Rec}}$, before fitting. This essentially unsmears the data, which then can be compared directly to theoretical predictions and fits. Although this is conceptually simple, there are a couple of big disadvantages to this method. First, HIJING does not incorporate final-state interactions, so weights must be used when building same-event (numerator) distributions. These weights account for the interactions, and, in the absence of Coulomb interactions, can be calculated using Eq. 5. Of course, these weights are valid only for a particular set of fit parameters. Therefore, in the fitting process, during which the fitter explores a large parameter set, the corrections will not remain valid. As

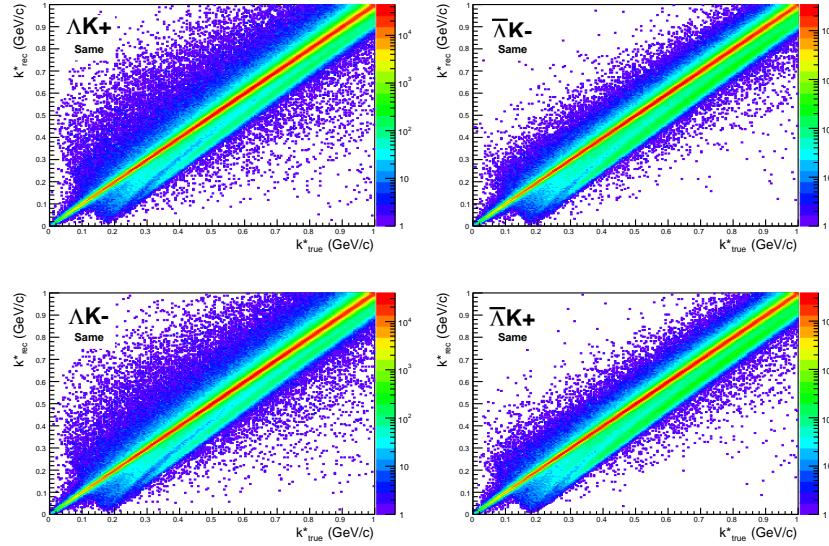
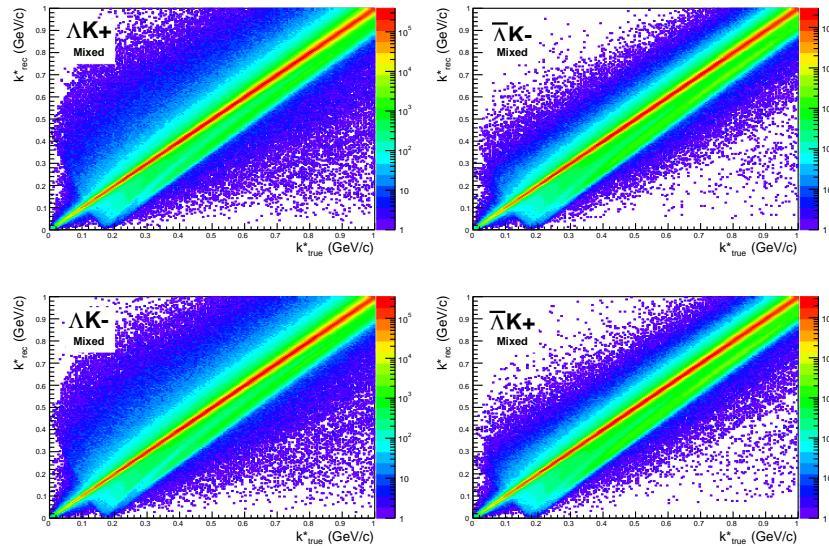
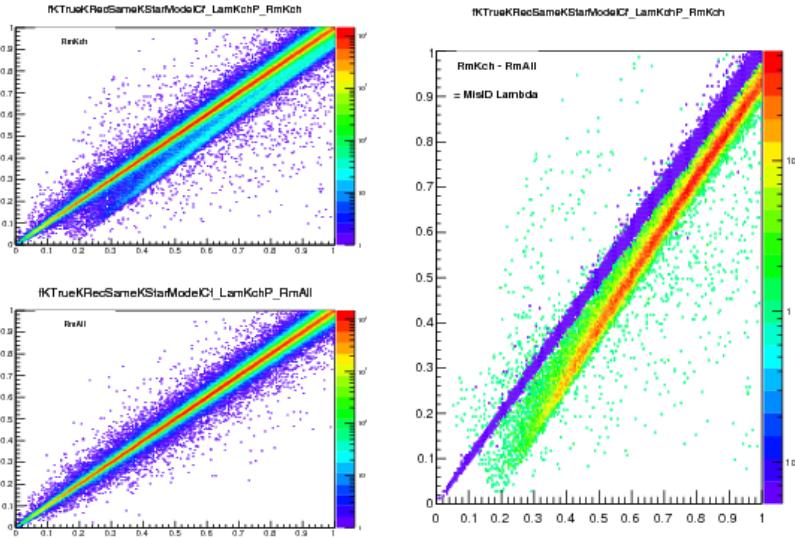
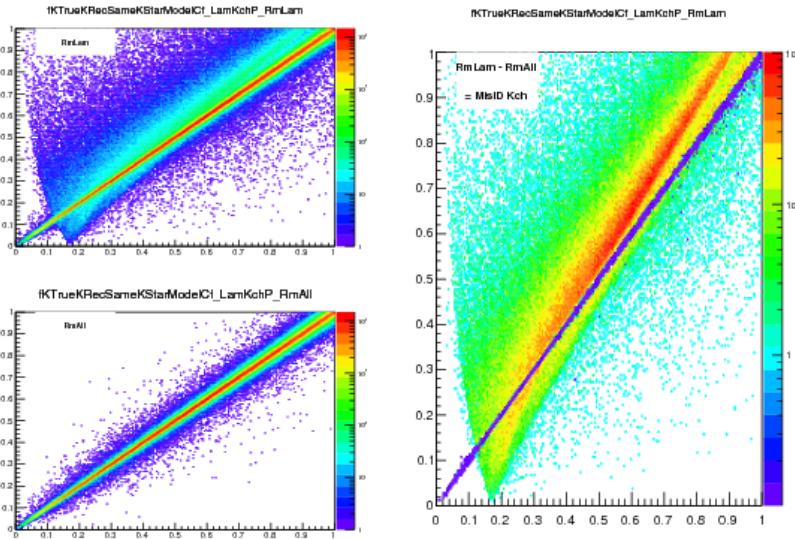
(a) Same Event Pairs (ΛK^\pm , 0-10% Centrality)(b) Mixed Event Pairs (ΛK^\pm , 0-10% Centrality)

Fig. 16: Sample k_{True}^* vs. k_{Rec}^* plots from MC HIJING events for ΛK^\pm 0-10% analyses. The structure which appears around $k_{\text{Rec}}^* = k_{\text{True}}^* - 0.15 \text{ GeV}/c$ is mainly caused by K_S^0 contamination in our $\Lambda(\bar{\Lambda})$ sample. The remaining structure not distributed about $k_{\text{Rec}}^* = k_{\text{True}}^*$ is due to π and e contamination in our K^\pm sample. These contaminations are more clearly visible in Figure 17

such, applying the momentum resolution correction and fitting becomes a long and drawn out iterative process. An initial parameter set is obtained (through fitting without momentum resolution corrections, theoretical models, or a good guess), then the MC data is analyzed to obtain correlation functions needed to calculate the correction factor, the data is fit using the correction factor, a refined parameter set is extracted, the MC data is analyzed again to obtain the new correction factor, etc. This process continues until the parameter set stabilizes. The second issue concerns statistics. With the MC data available on the grid, we were not able to generate the statistics necessary to use the raw $C_{\text{True}}/C_{\text{Rec}}$ ratio. The ratio was not stable, and when applied to the data, obscured the signal. Attempting to fit the ratio to use to generate the corrections also proved problematic. However, as HIJING does not include final-state interactions, the same-event and mixed-event pairs are very similar (with the exception of things like



(a) (Top Left) All misidentified K^+ excluded. (Bottom Left) All misidentified Λ and K^+ excluded. (Right) The difference of (Top Left) - (Bottom Left), which reveals the contamination in our Λ collection. The structure which appears around $k_{\text{Rec}}^* = k_{\text{True}}^* - 0.15 \text{ GeV}/c$ is mainly caused by K_S^0 contamination in our $\Lambda(\bar{\Lambda})$ sample.



(b) (Top Left) All misidentified Λ excluded. (Bottom Left) All misidentified Λ and K^+ excluded. (Right) The difference of (Top Left) - (Bottom Left), which reveals the contamination in our K^+ collection. The structure not distributed about $k_{\text{Rec}}^* = k_{\text{True}}^*$ is due to π and e^- contamination in our K^\pm sample.

Fig. 17: In the figure, the y-axis = k_{Rec}^* , and the x-axis = k_{True}^* . (Left) k_{Rec}^* vs. k_{True}^* plots for a small sample of the ΛK^+ 0-10% central analysis, MC truth was used to eliminate misidentified particles in the K^+ (a) and Λ (b) collections. (Right) The difference of the top left and bottom left plots. Contaminations in our particle collections are clearly visible. Figure (a) demonstrates a K_S^0 contamination in our Λ collection; Figure (b) demonstrates a π and e^- contamination in our K^\pm collection.

443 energy and momentum conservation, etc). Therefore, one may build the numerator distribution using
 444 mixed-event pairs. This corresponds, more or less, to simply running the weight generator through the
 445 detector framework.

446 A second approach (“Matrix” method) is to use information gained from plots like those in Figure 16,

447 which can be considered response matrices. The response matrix describes quantitatively how each
 448 k_{Rec}^* bin receives contributions from multiple k_{True}^* bins, and can be used to account for the effects of
 449 finite momentum resolution. With this approach, the resolution correction is applied on-the-fly during
 450 the fitting process by propagating the theoretical correlation function (fit) through the response matrix,
 451 according to:

$$C_{\text{Fit}}(k_{\text{Rec}}^*) = \frac{\sum_{k_{\text{True}}^*} M_{k_{\text{Rec}}^*, k_{\text{True}}^*} C_{\text{Fit}}(k_{\text{True}}^*)}{\sum_{k_{\text{True}}^*} M_{k_{\text{Rec}}^*, k_{\text{True}}^*}} \quad (14)$$

452 where $M_{k_{\text{Rec}}^*, k_{\text{True}}^*}$ is the response matrix (Figure 16), $C_{\text{Fit}}(k_{\text{True}}^*)$ is the fit binned in k_{True}^* , and the denominator normalizes the result.

454 Equation 14 describes that, for a given k_{Rec}^* bin, the observed value of $C(k_{\text{Rec}}^*)$ is a weighted average of
 455 all $C(k_{\text{True}}^*)$ values, where the weights are the normalized number of counts in the $[k_{\text{Rec}}^*, k_{\text{True}}^*]$ bin. As
 456 seen in Figure 16, overwhelmingly the main contributions comes from the $k_{\text{Rec}}^* = k_{\text{True}}^*$ bins. Although
 457 the correction is small, it is non-negligible for the low- k^* region of the correlation function.

458 Here, the momentum resolution correction is applied to the fit, not the data. In other words, during
 459 fitting, the theoretical correlation function is smeared just as real data would be, instead of unsmearing
 460 the data. This may not be ideal for the theorist attempting to compare a model to experimental data, but
 461 it leaves the experimental data unadulterated. The current analyses use this second approach to applying
 462 momentum resolution corrections because of two major advantages. First, the MC data must be analyzed
 463 only once, and no assumptions about the fit are needed. Secondly, the momentum resolution correction
 464 is applied on-the-fly by the fitter, delegating the iterative process to a computer instead of the user.

465 The two methods described above, Ratio and Matrix, should reproduce the same results at the parameter
 466 set used to generate the $C_{\text{True}}/C_{\text{Rec}}$ needed for the Ratio method. Figure 18 shows that the two methods
 467 converge as the binning size is decreased.

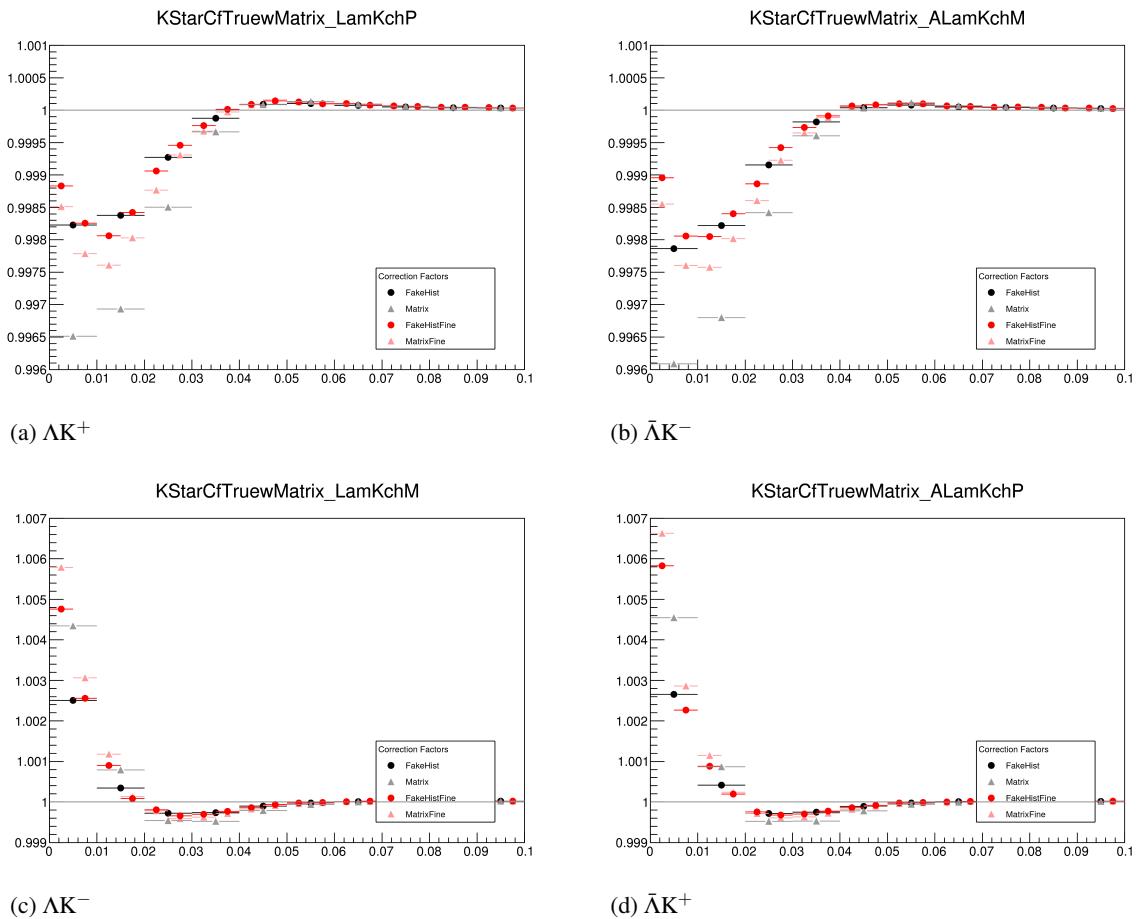


Fig. 18: Comparison of the two methods, Ratio and Matrix, for accounting for momentum resolution effects with HIJING. The Ratio method corresponds to the “FakeHist” histograms (circles), while the Matrix method corresponds to the “Matrix” histograms (triangles). Black shows a course binning, while red shows a finer binning.

468 **5.4 Residual Correlations**

469 The purpose of this analysis is to study both the interaction, and the scale of the emitting source, of
 470 the pairs. In order to obtain correct results, it is important for our particle collections to consist of
 471 primary particles. In practice, this is difficult to achieve for our Λ and $\bar{\Lambda}$ collections. Many of our Λ
 472 particles are not primary, but originate as decay products from other hyperons, including Σ^0 , Ξ^0 , Ξ^- and
 473 $\Sigma^{*(+,-,0)}(1385)$. Additionally, many of our K particles are not primary, but decay from $K^{*(+,-,0)}(892)$
 474 parents. In these decays, the daughter Λ or K_S^0 carries away a momentum very similar to that of its parent.
 475 As a result, the correlation function between a secondary Λ and, for instance, a K^+ will be sensitive to,
 476 and dependent upon, the interaction between the parent of the Λ and the K^+ . In effect, the correlation
 477 between the parent of the Λ and the K^+ (ex. $\Sigma^0 K^+$) will be visible, although smeared out, in the ΛK^+
 478 data; we call this a residual correlation resulting from feed-down. The contributions from the primary
 479 correlation, residual correlations, and fake pairs on the finally measure data is shown schematically in
 480 Figure 19. Residual correlations are important in an analysis when three criteria are met [6]: i) the parent
 481 correlation signal is large, ii) a large fraction of pairs in the sample originate from the particular parent
 482 system, and iii) the decay momenta are comparable to the expected correlation width in k^* .

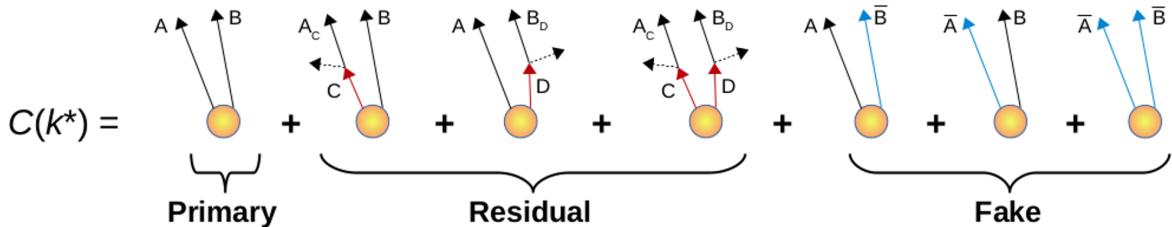


Fig. 19: A schematic representation of the contributions to the finally measured data from the primary correlation, residual correlations, and fake pairs.

483 As it is difficult for us to eliminate these residual correlations in our analyses, we must attempt to account
 484 for them in our fit. To achieve this, we will simultaneously fit the data for both the primary correlation
 485 function and the residual correlations. For example, in the simple case of a ΛK^+ analysis with residuals
 486 arising solely from $\Sigma^0 K^+$ feed-down:

$$C_{measured}(k_{\Lambda K^+}^*) = 1 + \lambda_{\Lambda K^+}[C_{\Lambda K^+}(k_{\Lambda K^+}^*) - 1] + \lambda_{\Sigma^0 K^+}[C_{\Sigma^0 K^+}(k_{\Lambda K^+}^*) - 1]$$

$$C_{\Sigma^0 K^+}(k_{\Lambda K^+}^*) \equiv \frac{\sum_{k_{\Sigma^0 K^+}^*} C_{\Sigma^0 K^+}(k_{\Sigma^0 K^+}^*) T(k_{\Sigma^0 K^+}^*, k_{\Lambda K^+}^*)}{\sum_{k_{\Sigma^0 K^+}^*} T(k_{\Sigma^0 K^+}^*, k_{\Lambda K^+}^*)} \quad (15)$$

487 $C_{\Sigma^0 K^+}(k_{\Sigma^0 K^+}^*)$ is the $\Sigma^0 K^+$ correlation function from, for instance, Equation 5, and T is the transform
 488 matrix generated with THERMINATOR. The transform matrix is formed for a given parent pair, AB,
 489 by taking all ΛK pairs originating from AB, calculating the relative momentum of the parents (k_{AB}^*)
 490 and daughters ($k_{\Lambda K}^*$), and filling a two-dimensional histogram with the values. The transform matrix
 491 is essentially an unnormalized probability distribution mapping the k^* of the parent pair to that of the
 492 daughter pair when one or both parents decay. An example of such transform matrices can be found in
 493 Figures 20 and 21.

494 The above equation can be easily extended to include feed-down from more sources:

$$\begin{aligned}
C_{measured}(k_{\Lambda K}^*) &= 1 + \lambda_{\Lambda K}[C_{\Lambda K}(k_{\Lambda K}^*) - 1] + \lambda_{\Sigma^0 K}[C_{\Sigma^0 K}(k_{\Lambda K}^*) - 1] + \dots \\
&\quad + \lambda_{P_1 P_2}[C_{P_1 P_2}(k_{\Lambda K}^*) - 1] + \lambda_{other}[C_{other}(k_{\Lambda K}^*) - 1]
\end{aligned}
\tag{16}$$

$$C_{P_1 P_2}(k_{\Lambda K}^*) \equiv \frac{\sum_{k_{P_1 P_2}^*} C_{P_1 P_2}(k_{P_1 P_2}^*) T(k_{P_1 P_2}^*, k_{\Lambda K}^*)}{\sum_{k_{P_1 P_2}^*} T(k_{P_1 P_2}^*, k_{\Lambda K}^*)}$$

495 Or, more compactly:

$$C_{measured}(k_{\Lambda K}^*) = 1 + \sum_i \lambda_i [C_i(k_{\Lambda K}^*) - 1] \tag{17}$$

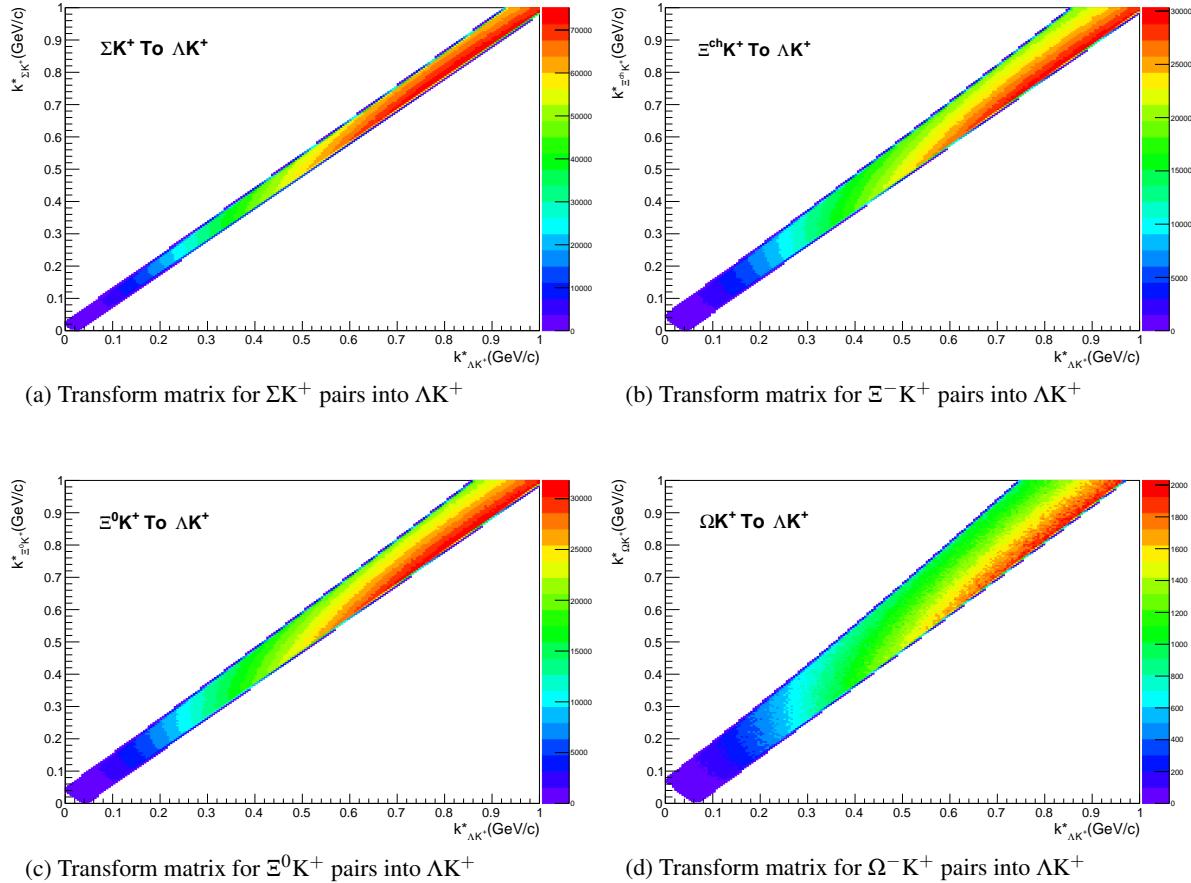


Fig. 20: Sample Transform Matrices generated with THERMINATOR for ΛK^+ Analysis

496 So, in practice, we model the correlation function of the parents (ex. $\Sigma^0 K^+$), and run the correlation
497 function through the appropriate transform matrix to determine the contribution to the daughter correla-
498 tion function (ex. ΛK^+). A few questions still remain. First, what λ values should be used in the above
499 equation? One option would be to leave all of these λ -parameters free during the fit process. However,
500 this would introduce a huge number of new parameters into the fitter, and would make the fit results less
501 trustworthy. The λ parameters roughly dictate the strength of the parent contribution to the daughter
502 pair. Additionally, as found in [7], the reconstruction efficiency for primary Λ particles is nearly equal
503 to that of Λ particles originating from Σ , Σ^* , Ξ^0 , Ξ^- , and Ω hyperons. Therefore, the λ parameter for

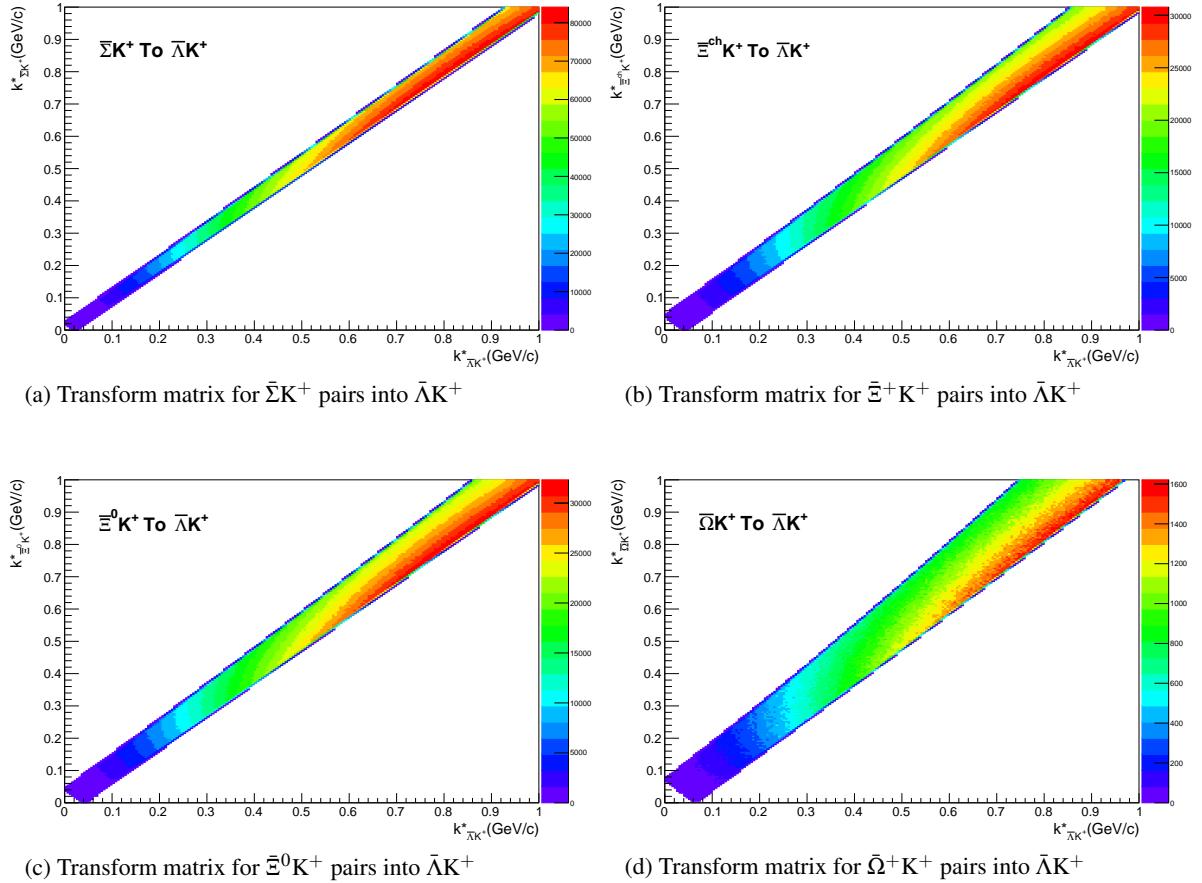
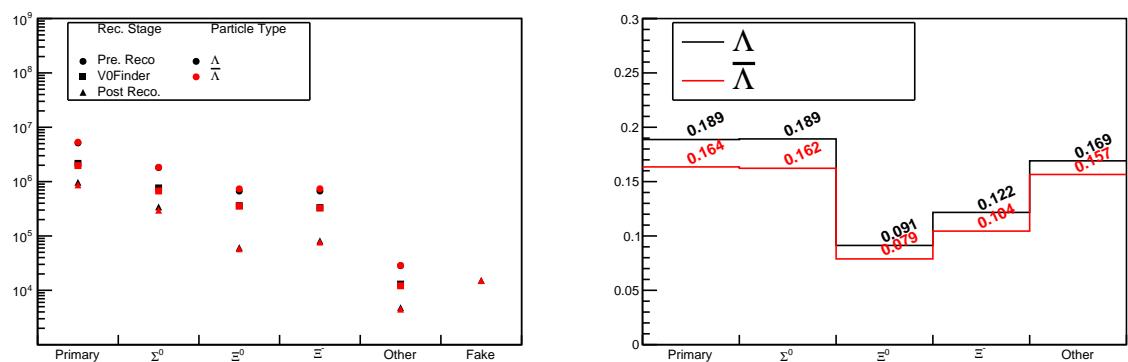


Fig. 21: Sample Transform Matrices generated with THERMINATOR for $\bar{\Lambda}K^+$ Analysis

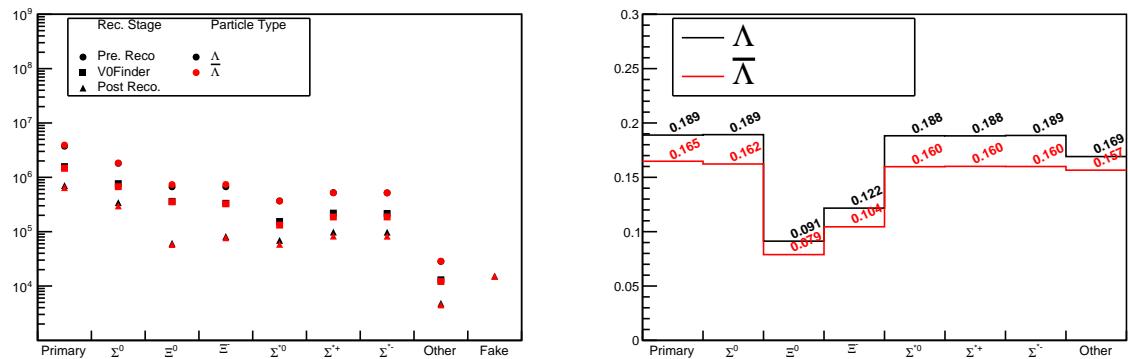
parent system AB can be estimated using THERMINATOR as the total number of ΛK pairs originating from AB (N_{AB}) divided by the total number of ΛK pairs (N_{Total}):

$$\lambda_{AB} = \frac{N_{AB}}{N_{Total}} \quad (18)$$

Note, for our study, we consider a particle to be primary if its parent has a proper decay length of < 4 fm. The λ values used can be found in Table 4, for the case of both three and ten residual contributors. In the table, we also list the λ values used for “Other” and “Fakes”. The “Other” category contains pairs which are not primary, and which do not originate from the (3 or 10) residual pairs included in the fit. The “Fakes” category represents pairs that are mistakenly identified as ΛK . To estimate this λ_{Fakes} value, we assumed that the number of fake pairs was equal to the total number of pairs multiplied by the Λ purity (i.e. assuming perfect purity for the kaons); or, more simply, $\lambda_{Fakes} = 1.0 - \text{Purity}(\Lambda)$. For both of these contributors (“Other” and “Fakes”), we assume that these correlations average to unity, and therefore do not contribute to the final correlation function.



(a) Reconstruction Efficiencies (3 Residuals)



(b) Reconstruction Efficiencies (10 Residuals)

Fig. 22: Reconstruction Efficiencies

| ΛK^+ residuals | | $\bar{\Lambda}K^-$ residuals | | ΛK^- residuals | | $\bar{\Lambda}K^+$ residuals | | ΛK_S^0 residuals | | $\bar{\Lambda}K_S^0$ residuals | |
|-------------------------|-----------------|-------------------------------|-----------------|-------------------------|-----------------|------------------------------|-----------------|---------------------------|-----------------|--------------------------------|-----------------|
| Pair System | λ value | Pair System | λ value | Pair System | λ value | Pair System | λ value | Pair System | λ value | Pair System | λ value |
| 3 Residuals | | | | | | | | | | | |
| ΛK^+ | 0.488 | $\bar{\Lambda}K^-$ | 0.488 | ΛK^- | 0.487 | $\bar{\Lambda}K^+$ | 0.489 | ΛK_S^0 | 0.505 | $\bar{\Lambda}K_S^0$ | 0.507 |
| $\Sigma^0 K^+$ | 0.103 | $\bar{\Sigma}^0 K^-$ | 0.102 | $\Sigma^0 K^-$ | 0.102 | $\bar{\Sigma}^0 K^+$ | 0.102 | $\Sigma^0 K_S^0$ | 0.112 | $\bar{\Sigma}^0 K_S^0$ | 0.111 |
| $\Xi^0 K^+$ | 0.046 | $\bar{\Xi}^0 K^-$ | 0.040 | $\Xi^0 K^-$ | 0.046 | $\bar{\Xi}^0 K^+$ | 0.041 | $\Xi^0 K_S^0$ | 0.050 | $\bar{\Xi}^0 K_S^0$ | 0.044 |
| $\Xi^- K^+$ | 0.063 | $\bar{\Xi}^+ K^-$ | 0.056 | $\Xi^- K^-$ | 0.062 | $\bar{\Xi}^+ K^+$ | 0.057 | $\Xi^- K_S^0$ | 0.068 | $\bar{\Xi}^+ K_S^0$ | 0.062 |
| Other | 0.252 | Other | 0.261 | Other | 0.255 | Other | 0.259 | Other | 0.217 | Other | 0.222 |
| Fakes | 0.048 | Fakes | 0.048 | Fakes | 0.048 | Fakes | 0.048 | Fakes | 0.048 | Fakes | 0.048 |
| 10 Residuals | | | | | | | | | | | |
| ΛK^+ | 0.317 | $\bar{\Lambda}K^-$ | 0.319 | ΛK^- | 0.317 | $\bar{\Lambda}K^+$ | 0.318 | ΛK_S^0 | 0.342 | $\bar{\Lambda}K_S^0$ | 0.344 |
| $\Sigma^0 K^+$ | 0.105 | $\bar{\Sigma}^0 K^-$ | 0.105 | $\Sigma^0 K^-$ | 0.104 | $\bar{\Sigma}^0 K^+$ | 0.106 | $\Sigma^0 K_S^0$ | 0.113 | $\bar{\Sigma}^0 K_S^0$ | 0.114 |
| $\Xi^0 K^+$ | 0.047 | $\bar{\Xi}^0 K^-$ | 0.042 | $\Xi^0 K^-$ | 0.047 | $\bar{\Xi}^0 K^+$ | 0.042 | $\Xi^0 K_S^0$ | 0.051 | $\bar{\Xi}^0 K_S^0$ | 0.045 |
| $\Xi^- K^+$ | 0.064 | $\bar{\Xi}^+ K^-$ | 0.058 | $\Xi^- K^-$ | 0.063 | $\bar{\Xi}^+ K^+$ | 0.059 | $\Xi^- K_S^0$ | 0.069 | $\bar{\Xi}^+ K_S^0$ | 0.063 |
| $\Sigma^{*+} K^+$ | 0.049 | $\bar{\Sigma}^{*-} K^-$ | 0.047 | $\Sigma^{*+} K^-$ | 0.048 | $\bar{\Sigma}^{*-} K^+$ | 0.048 | $\Sigma^{*+} K_S^0$ | 0.052 | $\bar{\Sigma}^{*-} K_S^0$ | 0.051 |
| $\Sigma^{*-} K^+$ | 0.044 | $\bar{\Sigma}^{*+} K^-$ | 0.046 | $\Sigma^{*-} K^-$ | 0.044 | $\bar{\Sigma}^{*+} K^+$ | 0.047 | $\Sigma^{*-} K_S^0$ | 0.047 | $\bar{\Sigma}^{*+} K_S^0$ | 0.050 |
| $\Sigma^{*0} K^+$ | 0.044 | $\bar{\Sigma}^{*0} K^-$ | 0.041 | $\Sigma^{*0} K^-$ | 0.044 | $\bar{\Sigma}^{*0} K^+$ | 0.042 | $\Sigma^{*0} K_S^0$ | 0.047 | $\bar{\Sigma}^{*0} K_S^0$ | 0.045 |
| ΛK^{*0} | 0.041 | $\bar{\Lambda}K^{*0}$ | 0.042 | $\Lambda \bar{K}^{*0}$ | 0.041 | $\bar{\Lambda}K^{*0}$ | 0.042 | ΛK^{*0} | 0.020 | $\bar{\Lambda}K^{*0}$ | 0.020 |
| $\Sigma^0 K^{*0}$ | 0.037 | $\bar{\Sigma}^0 \bar{K}^{*0}$ | 0.037 | $\Sigma^0 \bar{K}^{*0}$ | 0.037 | $\bar{\Sigma}^0 K^{*0}$ | 0.037 | $\Sigma^0 K^{*0}$ | 0.017 | $\bar{\Sigma}^0 K^{*0}$ | 0.018 |
| $\Xi^0 K^{*0}$ | 0.017 | $\bar{\Xi}^0 \bar{K}^{*0}$ | 0.015 | $\Xi^0 \bar{K}^{*0}$ | 0.017 | $\bar{\Xi}^0 K^{*0}$ | 0.015 | $\Xi^0 K^{*0}$ | 0.008 | $\bar{\Xi}^0 K^{*0}$ | 0.007 |
| $\Xi^- K^{*0}$ | 0.022 | $\bar{\Xi}^+ \bar{K}^{*0}$ | 0.021 | $\Xi^- \bar{K}^{*0}$ | 0.022 | $\bar{\Xi}^+ K^{*0}$ | 0.021 | $\Xi^- K^{*0}$ | 0.011 | $\bar{\Xi}^+ K^{*0}$ | 0.010 |
| Other | 0.166 | Other | 0.174 | Other | 0.169 | Other | 0.171 | Other | 0.176 | Other | 0.181 |
| Fakes | 0.048 | Fakes | 0.048 | Fakes | 0.048 | Fakes | 0.048 | Fakes | 0.048 | Fakes | 0.048 |

Table 4: λ values for the individual components of the ΛK correlation functions for the case of 3 and 10 residual contributions.

515 Now, the remaining question is how do we model the parent correlation functions? In an ideal world, we
 516 would simply look up the parent interaction in some table, and input this into our Lednicky equation (for
 517 the case of one or more charge neutral particle in the pair), or run it through the CoulombFitter machinery
 518 described in Sec.5.2. Unfortunately, the world in which we live is not perfect, such a table does not exists,
 519 and little is know about the interaction between the residual pairs in this study. One solution would be
 520 to introduce a set of scattering parameters and radii for each residual system. However, as was the case
 521 of the λ -parameters above, this would introduce a large number of additional fit parameters, and would
 522 make our fitter too unconstrained and would yield untrustworthy results. The second option, which is
 523 adopted in this analysis, is to assume all residual pairs have the same source size as the daughter pair,
 524 and all Coulomb-neutral residual pairs also share the same scattering parameters as the daughter pair (the
 525 case of charged pairs, such as $\Xi^- K^\pm$ or $\Sigma^{*\pm} K^\pm$, will be described below).

526 Concerning the radii of the residual parent pairs, it was suggested that these should be set to smaller
 527 values than those of the daughter pair. In the interest of minimizing the number of parameters in the
 528 fitter, we tested this by introducing an m_T -scaling of the parents' radii. The motivation for this scaling
 529 comes from the approximate m_T -scaling of the radii observed in ???. To achieve this scaling, we assume
 530 the radii follow an inverse-square-root distribution: $R_{AB} = \alpha m_T^{-1/2}$. Then, it follows that we should scale
 531 the parent radii as:

$$R_{AB} = R_{\Lambda K} \left(\frac{m_{T,AB}}{m_{T,\Lambda K}} \right)^{-1/2} \quad (19)$$

532 The values of m_T for each pair system were taken from THERMINATOR. As the fitter dances around
 533 parameter space and selects a new radius for the ΛK system, the radii of the residuals is simply the ΛK
 534 radius scaled by the appropriate factor, given above (Eq.19). In the end, this scaling factor made no
 535 significant difference in our fit results, so this complication is excluded from our final results. Note that
 536 this is not surprising, as the most extreme scaling factor was, in the case of using 10 residual systems,
 537 between ΛK^+ with $m_{T,\Lambda K^+} \approx 1.4 \text{ GeV}/c^2$ and $\Xi^- K^{*0}$ with $m_{T,\Xi^- K^{*0}} \approx 1.8 \text{ GeV}/c^2$, resulting in a scale
 538 factor of ≈ 0.9 .

539 Now, as hinted above, accounting for charged residuals adds a complication in that they necessitate the
 540 inclusion of the CoulombFitter (described in Sec. 5.7) into the process. The complication of combining
 541 the two fitters is not troubling; however, the substantial increase in the fitting time is (the parallelization
 542 of the CoulombFitter across a large number of GPU cores, to drastically decrease run-time, is currently
 543 underway). We have two solutions to bypass such a large increase in run time. First, we can use our
 544 experimental $\Xi^{\text{ch}} K^{\text{ch}}$ data to represent all charged parent pair system. In this case, there is no need to
 545 make any assumption about scattering parameters or source sizes, as we already have the experimental
 546 data. The downside is that, especially in the 30-50% centrality bin, the statistics are low and error bars
 547 large. Alternatively, we can assume the strong interaction is negligible in the charged residual, and
 548 generate the parent correlation function given radius and λ parameters. We find in our $\Xi^{\text{ch}} K^{\text{ch}}$ study
 549 that a Coulomb-only description of the system describes, reasonably well, the broad features of the
 550 correlation. The strong interaction is necessary for the fine details. However, as these correlations are
 551 run through a transform matrix, which largely flattens out and fine details, a Coulomb-only description
 552 should be sufficient. In practice, this Coulomb-only scenario is achieved by first building a large number
 553 of Coulomb-only correlations for various radii and λ parameter values, and interpolating from this grid
 554 during the fitting process. We find consistent results between using the ΞK data and the Coulomb-only
 555 interpolation method. When quantifying the $\Xi^- K^\pm$ residual contribution, the experimental $\Xi^- K^\pm$ data
 556 is always used. When the number of residual pairs used is increased to 10, so that contributors such
 557 as $\Sigma^{*+} K^-$ enter the picture, the Coulomb-only interpolation method is used. In other words, the ΞK
 558 experimental data is only used to model the ΞK residual contribution, all other charged pairs are treated
 559 with the Coulomb-only interpolation method.

560 Two examples of how very different transform matrices can alter a correlation function are shown in
 561 Figures 23 and 24 below. These figures were taken using parameter values obtained from fits to the data.
 562 In the top left corner of the figures, the input correlation function (closed symbols) is shown together
 563 with the output, transformed, correlation function (open symbols). In the bottom left, the transformed
 564 correlation is shown by itself (with zoomed y-axis). This is especially helpful when the λ parameter is
 565 very small, in which case the contribution in the top left can look flat, but the zoomed in view in the
 566 bottom left shows the structure. The right two plots in each figure show the transform matrix without
 567 (top right) and with (bottom right) a log-scale on the z-axis. Note, more examples of these transforms
 568 can be found in Sec. 9.

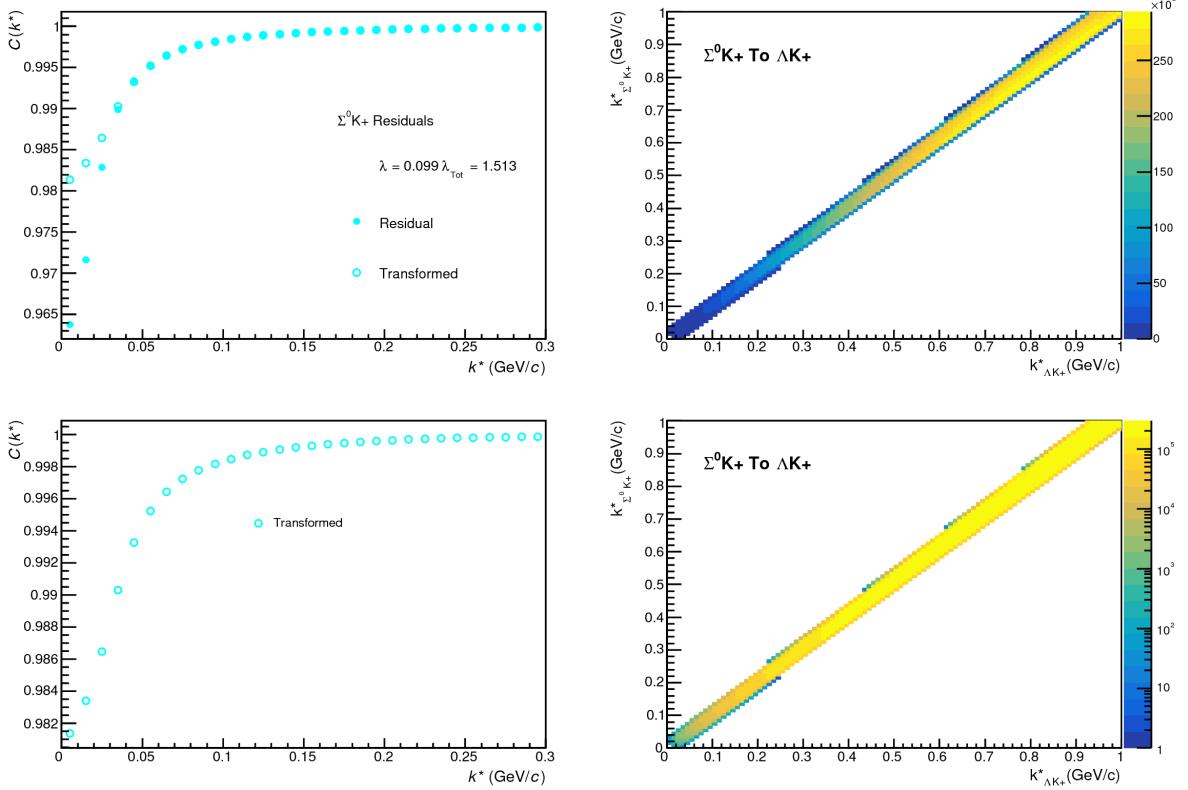


Fig. 23: $\Sigma^0 K^+$ Transform. These figures were taken using parameter values obtained from fits to the data. In the top left corner of the figures, the input correlation function (closed symbols) is shown together with the output, transformed, correlation function (open symbols). In the bottom left, the transformed correlation is shown by itself. The right two plots in each figure show the transform matrix without (top right) and with (bottom right) a log-scale on the z-axis.

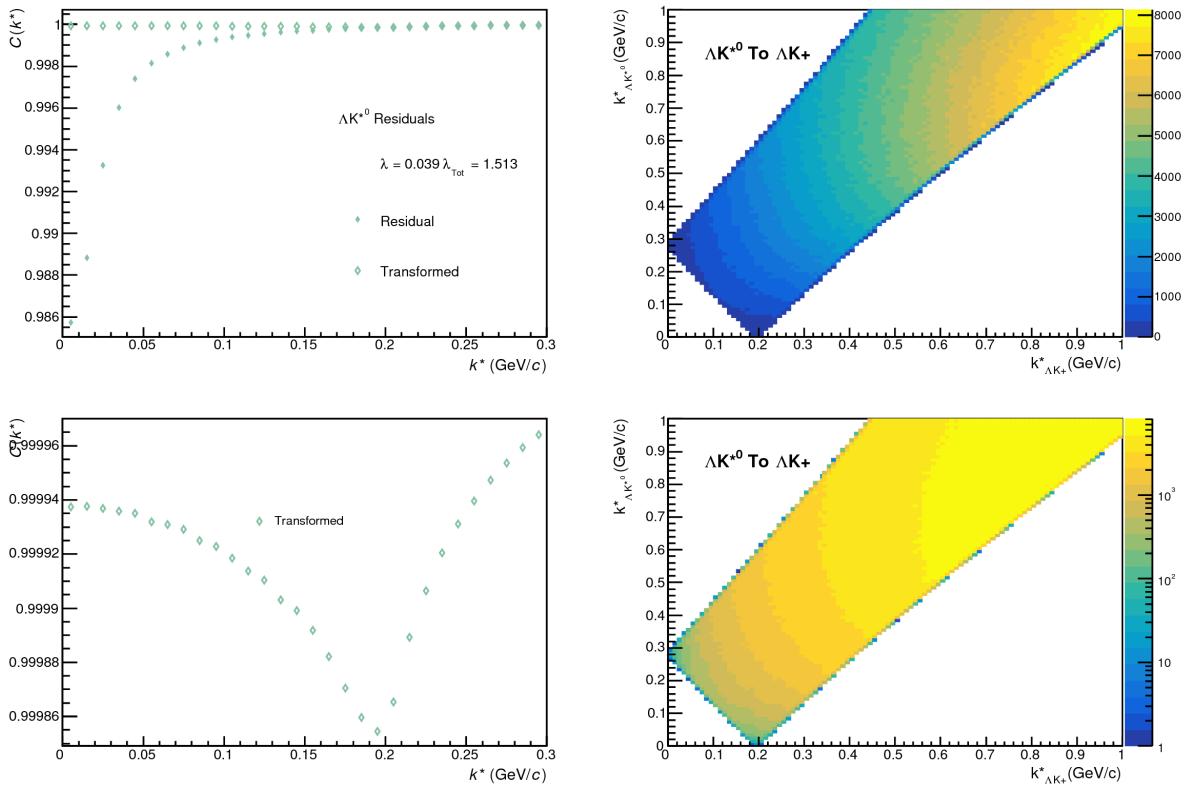


Fig. 24: ΛK^0 Transform. These figures were taken using parameter values obtained from fits to the data. In the top left corner of the figures, the input correlation function (closed symbols) is shown together with the output, transformed, correlation function (open symbols). In the bottom left, the transformed correlation is shown by itself. The right two plots in each figure show the transform matrix without (top right) and with (bottom right) a log-scale on the z-axis.

569 **5.5 Non-Flat Background**

570 We observe a significant non-femtoscopic, non-flat, background in all of our correlations at large k^* .
 571 This background increases with decreasing centrality, is the same amongst all ΛK^\pm pairs, and is more
 572 pronounced in the ΛK_S^0 system, as can be seen in Fig. 25. Figure 26a shows that THERMINATOR 2
 573 simulation does a good job of describing the difference in backgrounds between ΛK^\pm and ΛK_S^0 .

574 Before beginning, it is important to note that the difference in ΛK^\pm and ΛK_S^0 backgrounds is due mainly
 575 to the difference in kinematic cuts, not due to any interesting physics. Figure 26b shows that, for THER-
 576 MINATOR simulation, when restrictions are imposed on the p_T of the K_S^0 to more closely match the
 577 K^\pm cuts, the backgrounds align much better. Therefore, we conclude that the difference in background
 578 between ΛK^\pm and ΛK_S^0 observed in our experimental data is simply due to a difference in kinematic cuts
 579 between K^\pm and K_S^0 particles.

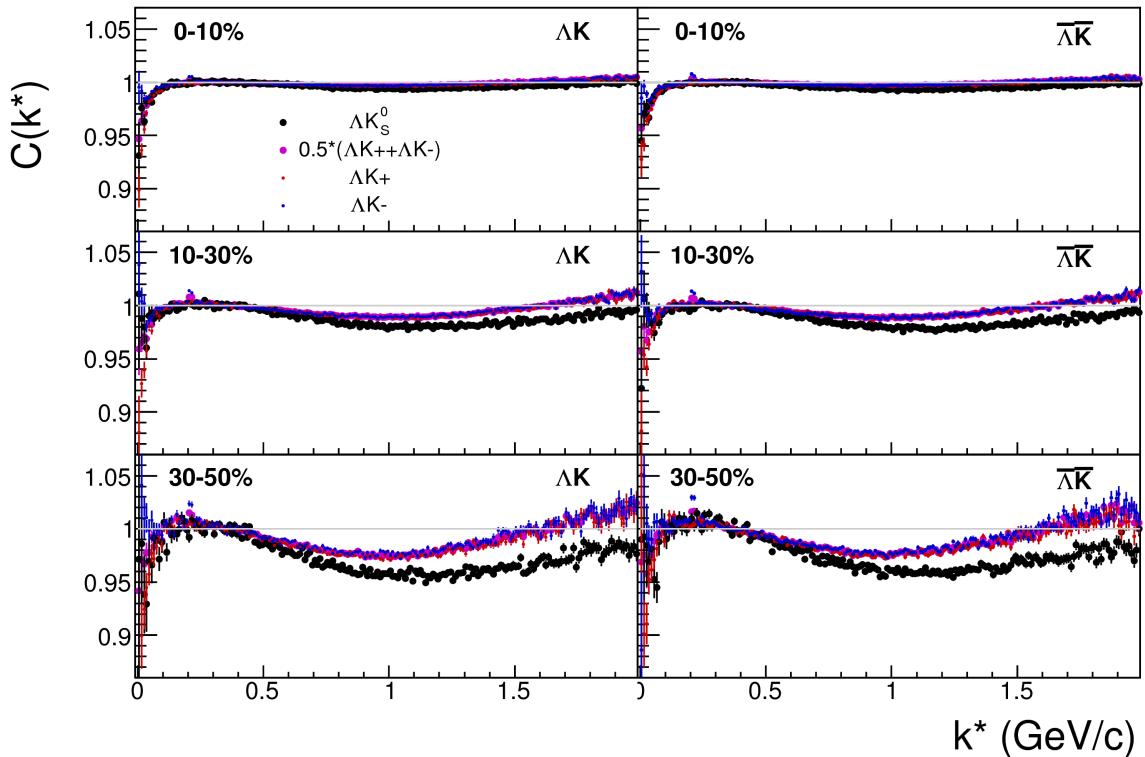


Fig. 25: A comparison on the non-femtoscopic backgrounds observed in our our ΛK experimental data.

580 It is suggested that this background effect is due primarily to particle collimation associated with elliptic
 581 flow [8]. More specifically, these backgrounds result from mixing events with unlike event-plane angles
 582 (Ψ_{EP}). As explained in [8], when elliptic flow is present, all particles are more likely to be emitted
 583 in a specific direction (in-plane), as opposed to a perpendicular direction. Therefore, the difference in
 584 momenta for pairs of particles tends to be smaller, compared to the case of no flow. In the case of mixed-
 585 event pairs, the two events used do not share an event-plane, and therefore there is no collimation effect
 586 in the pairs from flow. As a result, pairs with larger momentum are more likely when mixed-events are
 587 used, causing the correlation function to be observed below unity. In general, a dip below unity, at a given
 588 k^* , means it is more probable to find a pair at that k^* when the daughters are taken from mixed-events, as
 589 compared to when they are taken from the same event.

590 This same reasoning suggests that the background should lead to an enhancement at low- k^* . The en-
 591 hancement at high- k^* ($k^* \gtrsim 1.5 \text{ GeV}/c$) does not result from the collective flow of the system. We are not

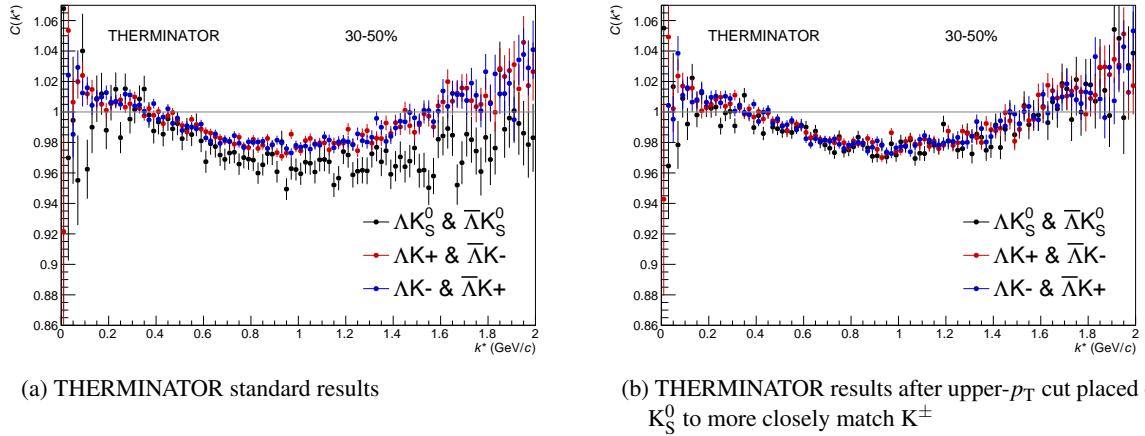


Fig. 26: THERMINATOR 2 simulation for ΛK^+ (red), ΛK^- (blue), and ΛK_S^0 (black). In 26a, we show the standard THERMINATOR 2 results. THERMINATOR 2 does a good job describing qualitatively the different between the ΛK^\pm and ΛK_S^0 backgrounds. In 26b, an upper- p_T cut was placed on the K_S^0 particles to more closely match the K^\pm kinematic cuts. After this tweak, the ΛK^\pm and ΛK_S^0 backgrounds agree much better.

592 certain was causes this enhancement, but typical suspects are jet-like correlations and resonance decays.
 593 We can split our correlation functions into three main regions. First, the low- k^* region ($k^* \lesssim 0.3$ GeV/c)
 594 contains the femtoscopic correlations, as well as a likely enhancement from the background. The
 595 intermediate- k^* region ($0.3 \lesssim k^* \gtrsim 1.5$ GeV/c) contains a suppression from the background. Finally,
 596 the high- k^* region ($k^* \gtrsim 1.5$ GeV/c) contains an enhancement with unknown origin.
 597 THERMINATOR 2 simulation has been shown to reproduce the background features in a πK analysis
 598 [8]. As the background effect can be attributed mainly to elliptic flow, which is a global feature of the
 599 system, we suspected THERMINATOR 2 could also, at least qualitatively, describe our backgrounds.
 600 After ensuring each simulated event received a random event-plane angle (Ψ_{EP})¹, we found THERMI-
 601 NATOR 2 did a good job of describing our data qualitatively, and, in many cases, quantitatively. Figure
 602 27 shows the THERMINATOR 2 simulation (gold) together with experimental data (red, blue, or black).
 603 The figure also shows a 6th-order polynomial fit to the simulation (gold), as well as the fit polynomial
 604 scaled to match the data (red, blue, black).

605 Figure 28 shows three different correlation function generated using THERMINATOR 2 simulation (“Cf
 606 w/o Bgd (A)”, “Cf w. Bgd (B)”, “Bgd(C)”), as well as two histograms describing the relation between
 607 the three (“Ratio (B/C)”, “1+Diff(B-C)”). “Cf w/o Bgd (A)” shows a correlation function with a femto-
 608 scopic correlation, but without background. When THERMINATOR 2 is run without randomizing event
 609 planes, and therefore having all events share a common event plane, no background is observed, as ex-
 610 pected. The femtoscopic correlation effect was introduced by assuming a set of scattering parameters for
 611 the system, and weighting the numerators appropriately. The second correlation, ”Cf w. Bgd (B)”, shows
 612 a correlation function with both a femtoscopic correlation and a background (most closely matches our
 613 situation in experiment). To generate the background, each event was given a random event-plane angle,
 614 as is given to us in experiment. To generate the femtoscopic correlation, the same numerator weighting
 615 procedure was used. Finally, ”Bgd (C)”, shows a correlation function with a non-femtoscopic back-
 616 ground, but no femtoscopic correlation, i.e. background only. This is generated just as ”Cf w. Bgd (B)”,
 617 with randomized event planes, but unit weights are used when filling the numerators, so no femtoscopic
 618 effects are included.

¹default was for all events to share a common event plane

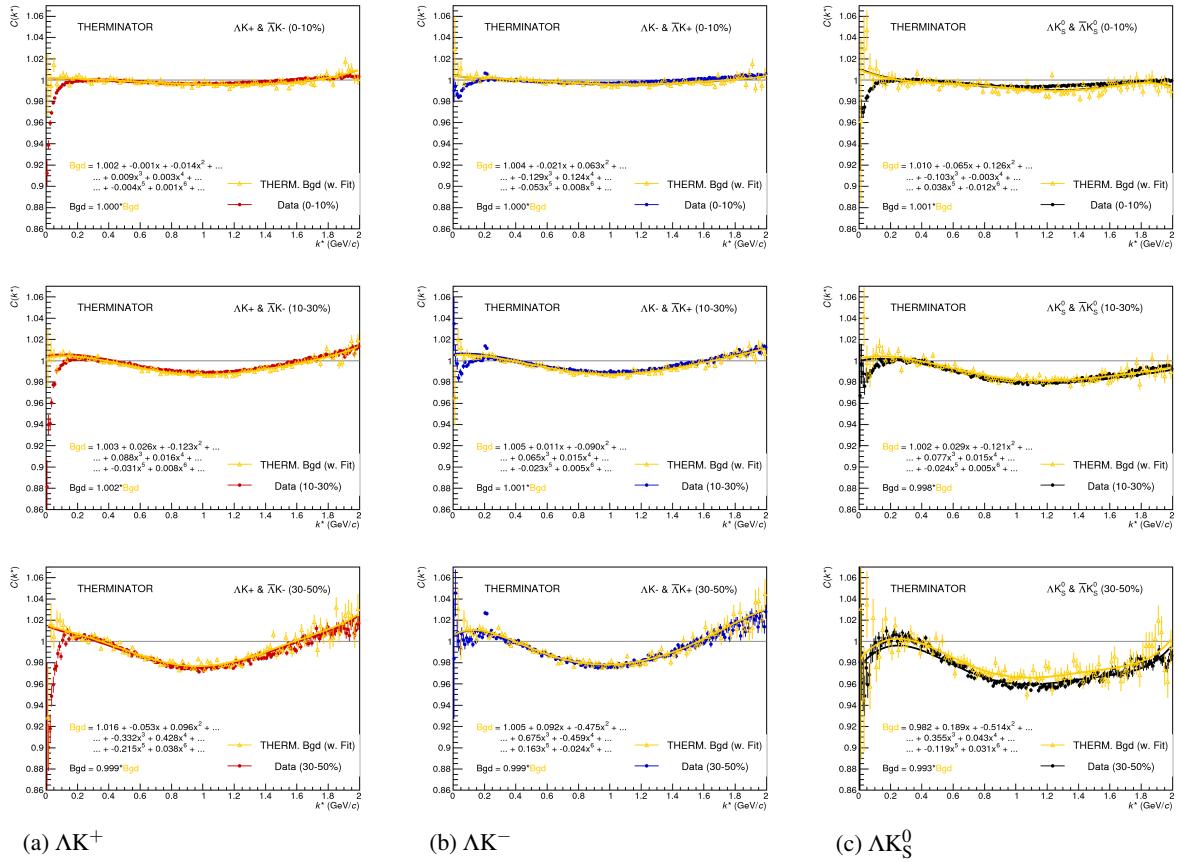


Fig. 27: THERMINATOR 2 simulation (gold) together with experimental data (red, blue, or black). The left column shows results for ΛK^+ (27a), middle for ΛK^- (27b), and right for ΛK_S^0 (27c). A 6th-order polynomial fit to the simulation is shown as a solid gold line, and whose fit parameters are printed on the lower left of each plot. This polynomial is scaled to match the experimental data; the value of this scale is printed in the lower left corner of each plot. The polynomial fit with scale factor applied is drawn in a color matching the experimental data (red, blue, black).

619 The main point of Fig. 28 is that the black points match the blue (and purple) points; or, equivalently:

$$C_{fw/oBgd} = \frac{C_{fw.Bgd}}{Bgd} \rightarrow C_{theory} = \frac{C_{exp}}{F_{Bgd}} \rightarrow C_{exp} = C_{theory} \cdot F_{Bgd} \quad (20)$$

620 As shown and described in Fig 28 and Eq. 20, THERMINATOR 2 simulation shows the non-femtosscopic
621 background affects the correlation function as a separable scale factor. We expect this behavior to be
622 roughly the same in the experimental data.

623 Figure 29 demonstrates the use of the Stavinsky method with THERMINATOR 2. In the figure, unit
624 weights were used for all numerators, so no femtoscopic signal is included, only background effects.
625 The black points show an ideal, experimentally unreachable, situation of aligning all of the event-plane
626 angles. With THERMINATOR 2, when the event-planes are aligned, the background signal is killed.
627 The green points show the case of random event-plane angles, a situation more closely matching that
628 of experiment. The purple points show the affect of applying the Stavinsky method to the case of
629 random event-planes. The figure shows that this method effectively kills the non-flat background (i.e.
630 the procedure takes the green points to the purple). Finally, the blue points show the effect of applying
631 the Stavinsky method when all of the event-planes are aligned. This shows that the Stavinsky method

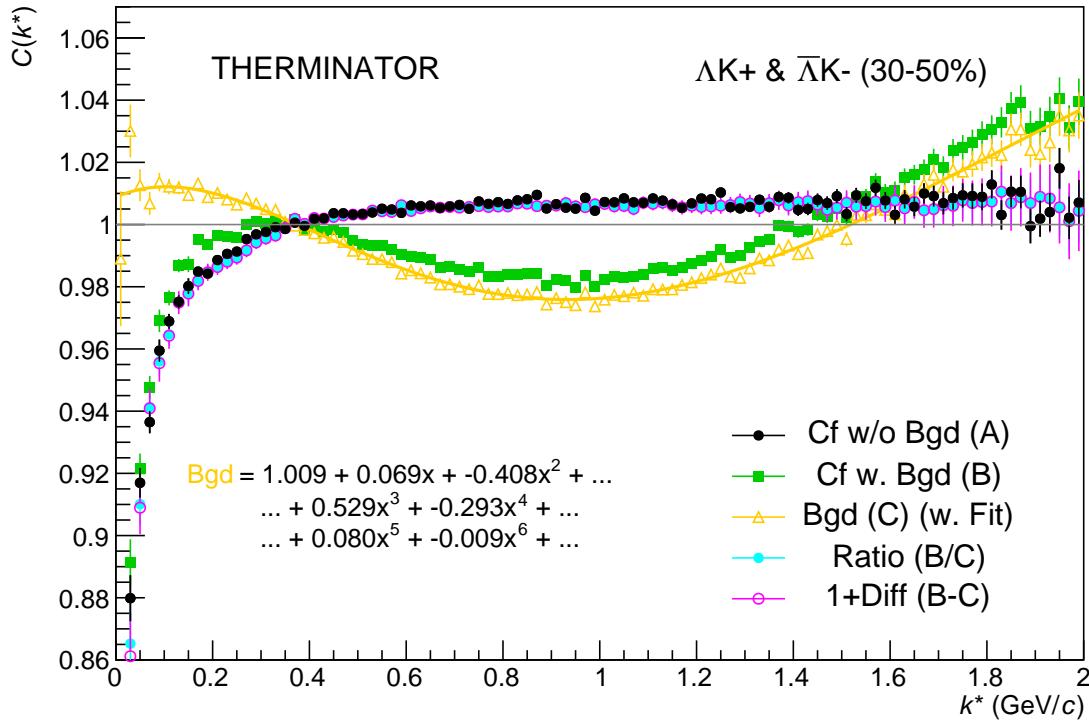


Fig. 28: Correlation with background decomposition with THERMINATOR. "Cf w/o Bgd (A)" shows a correlation function with a femtoscopic correlation, but without background. "Cf w. Bgd (B)", shows a correlation function with both a femtoscopic correlation and a background (most closely matches our situation in experiment). "Bgd (C)", shows a correlation function with a non-femtoscopic background, but no femtoscopic correlation, i.e. background only.

632 does not introduce any signal to an already flat background.

633 5.6 LednickyFitter

634 The code developed to fit the data is called "LednickyFitter", and utilizes the ROOT TMinuit implemen-
 635 tation of the MINUIT fitting package. In short, given a function with a number of parameters, the fitter
 636 explores the parameter space searching for the minimum of the function. In this implementation, the
 637 function to be minimized should represent the difference between the measured and theoretical corre-
 638 lation functions. However, a simple χ^2 test is inappropriate for fitting correlation functions, as the ratio of
 639 two Poisson distributions does not result in a Poisson distribution. Instead, a log-likelihood fit function
 640 of the following form is used [2]:

$$\chi_{PML}^2 = -2 \left[A \ln \left(\frac{C(A+B)}{A(C+1)} \right) + B \ln \left(\frac{A+B}{B(C+1)} \right) \right] \quad (21)$$

641 where A is the experimental signal distribution (numerator), B is the experimental background distribu-
 642 tion (denominator), and C is the theoretical fit correlation function.

643 The LednickyFitter uses Equations 5 – 7 to build the theoretical fit, and Equation 21 as the statistic
 644 quantifying the quality of the fit. The parameters to be varied by MINUIT are: λ , R , f_0 ($\mathbb{R}f_0$ and $\mathbb{I}f_0$
 645 separately), d_0 , and normalization N . The fitter currently includes methods to correct for momentum
 646 resolution and a non-flat background. These corrections are applied to the fit function, the data is never

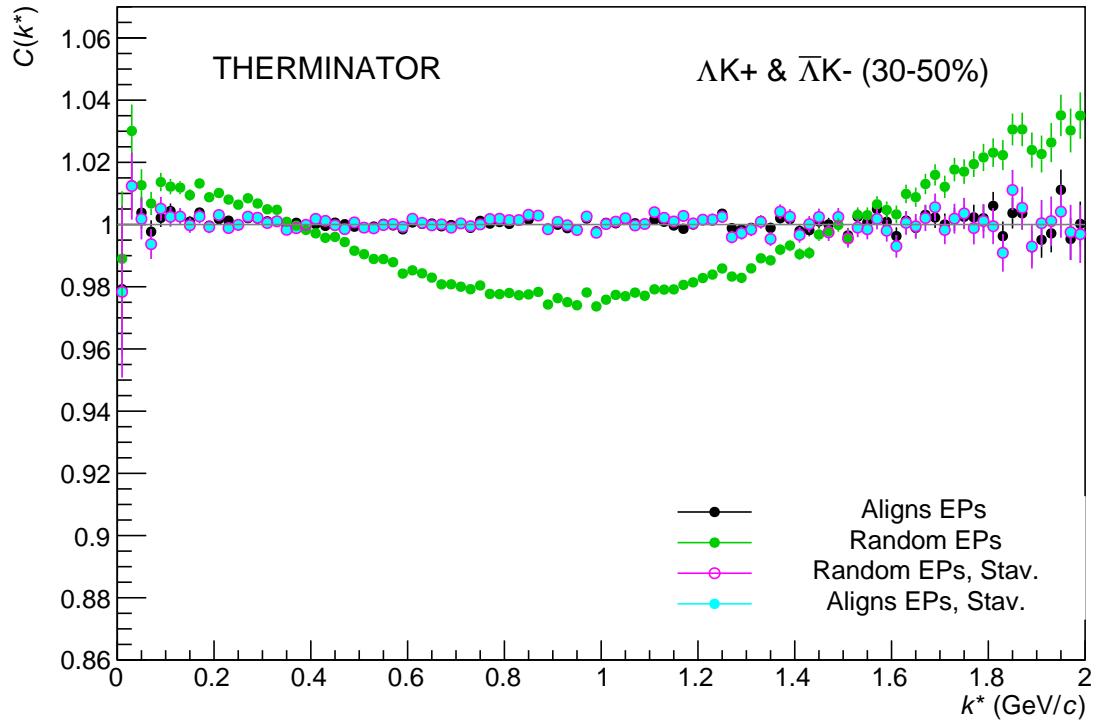


Fig. 29: The use of the Stavinsky method with THERMINATOR 2. Unit weights were used for all numerators, so no femtoscopic signal is included, only background effects. The black points show an ideal, experimentally unreachable, situation of aligning all of the event-plane angles. The green points show the experimental situation of random event-plane angles. The purple points shown the affect of applying the Stavinsky method to the case of random event-planes. Finally, the blue points show the effect of applying the Stavinsky method when all of the event-planes are aligned.

- 647 touched. The fitter is able to share parameters between different analyses and fit all simultaneously.
- 648 In a typical fit, a given pair is fit with its conjugate (ex. ΛK^+ with $\bar{\Lambda} K^-$) across all centralities (0-10%,
649 10-30%, 30-50%), for a total of 6 simultaneous analyses. Each analysis has a unique λ and normalization
650 parameter. The radii are shared between analyses of like centrality, as these should have similar source
651 sizes. The scattering parameters ($\mathbb{R}f_0, \mathbb{I}f_0, d_0$) are shared amongst all.
- 652 Initially, we left open the possibility for the $\Lambda K^+(\bar{\Lambda} K^-)$ and $\Lambda K^-(\bar{\Lambda} K^+)$ systems to have different source
653 radii. After always finding these two to be consistent, we decided to join the radii parameters between these
654 systems. So, now, in a typical fit of our ΛK^\pm data, all ΛK^\pm analyses ($\Lambda K^+, \bar{\Lambda} K^-, \Lambda K^-, \bar{\Lambda} K^+$) are fit
655 simultaneously across all centralities. Scattering parameters are shared between pair-conjugate systems
656 (i.e. a parameter set describing the $\Lambda K^+ & \bar{\Lambda} K^-$ system, and a separate set describing the $\Lambda K^- & \bar{\Lambda} K^+$
657 system). For each centrality, a radius and λ parameters are shared amongst all pairs. Each analysis has a
658 unique normalization parameter.
- 659 In the case of fitting with residuals, the λ_{Fit} parameter serves as an overall normalization shared by all
660 contributors, such that Eqn 17 becomes:

$$\begin{aligned}
C_{measured}(k_{\Lambda K}^*) &= 1 + \sum_i \lambda'_i [C_i(k_{\Lambda K}^*) - 1] \\
\lambda'_i &= \lambda_{Fit} \lambda_i \\
\sum_i \lambda'_i &= \lambda_{Fit} \sum_i \lambda_i = \lambda_{Fit}
\end{aligned} \tag{22}$$

where λ_i is obtained from THERMINATOR, as explained in Section 5.4, and whose values are presented in Tables ?? through ???. For Coulomb-neutral pairs, such as ΛK , $\Sigma^0 K$, and $\Xi^0 K$, $C_i(k_{\Lambda K}^*)$ is calculated from Eqn. 5, with the help of Eqn. 7. For those residual pairs which include a Coulomb interaction, $C_i(k_{\Lambda K}^*)$ is either calculated using the CoulombFitter method (Sections 5.2 and 5.7) with no strong interaction, or by using the $\Xi^{ch} K^{ch}$ data directly. Unless otherwise stated, the $\Xi^{ch} K^{ch}$ residual contribution is modeled using the experimental $\Xi^{ch} K^{ch}$ data, and all other charged contributors (ex. $\Sigma^{*ch} K^{ch}$) are modeled using the CoulombFitter technique with no strong interaction contribution.

To summarize, the complete fit function is constructed as follows:

1. The uncorrected, primary, correlation function, $C_{\Lambda K}(k_{True}^*)$, is constructed using Eq. 22 (with the help of Eqns. 5 and 7)
2. If residuals are included:
 - the parent correlation functions are obtained using:
 - Eq. 22 (with the help of Eqns. 5 and 7) for the case of Coulomb-neutral pairs
 - $\Xi^- K^\pm$ experimental data for $\Xi^- K^\pm$ contributions
 - a Coulomb-only curve, with the help of Secs. 5.2 and 5.7, for pairs including the Coulomb interaction
 - the contribution to the ΛK correlation function is found by running the parent correlation function through the appropriate transform, via Eq.16
3. The primary and residual correlations are combined, via Eq.17, to form $C'_{Fit}(k_{True}^*)$
 - in the case of no residual contributions included in the fit, $\lambda_i = \lambda_{\Lambda K}$ in Eq. 22 is set equal to 1. Then, the extracted λ_{Fit} parameter should be roughly equal to the pair purity
 - when residuals are included, the λ_i values are presented in Table 4
4. The correlation function is corrected to account for momentum resolution effects using Eq. 14

$$C'_{fit}(k_{Rec}^*) = \frac{\sum_{k_{True}^*} M_{k_{Rec}^*, k_{True}^*} C'_{fit}(k_{True}^*)}{\sum_{k_{True}^*} M_{k_{Rec}^*, k_{True}^*}}$$

5. Finally, the non-flat background correction is applied, and the final fit function is obtained

$$C_{Fit}(k_{Rec}^*) = C'_{Fit}(k_{Rec}^*) * F_{Bgd}(k_{Rec}^*)$$

Figures 47, 48, and 49 (33, 34, and 35, or 40, 41, and 42), in Section 7, show experimental data with fits for all studied centralities for $\Lambda K_S^0 (\bar{\Lambda} K_S^0)$, $\Lambda K^+ (\bar{\Lambda} K^-)$, and $\Lambda K^- (\bar{\Lambda} K^+)$, respectively. In the figures, the black solid line represents the “raw” fit, i.e. not corrected for momentum resolution effects nor non-flat background. The green line shows the fit to the non-flat background. The purple points show the fit after momentum resolution, non-flat background, and residual correlations (if applicable) corrections have been applied. The extracted fit values with uncertainties are also printed on the figures.

693 **5.7 Coulomb Fitter**

694 When fitting the $\Xi^-(\bar{\Xi}^+)K^\pm$ results, it is necessary to include both strong and Coulomb effects. In this
 695 case, Equation 5 is no longer valid, and, in fact, there is no analytical form with which to fit. We therefore
 696 must take a more basic approach, and integrate out Eq.4 by hand. To achieve this, one has two options.
 697 The first option is to numerically integrate Eq.4. The second option is to simulate a large sample of
 698 particle pairs, calculate the wave function describing the interaction, and average to obtain the integral.
 699 Having no experience with either of these options, we elected the latter of simulating pairs. The code
 700 developed to achieve this functionality is called “CoulombFitter”. Currently, in order to generate the
 701 statistics needed for a stable fit, we find that $\sim 10^4$ simulated pairs per 10 MeV bin are necessary. The
 702 nature of this process means that the “CoulombFitter” takes much longer to run than the “LednickyFitter”
 703 of Section 5.1.

704 Unfortunately, with this analysis, we are not sensitive to, and therefore not able to distinguish between,
 705 the iso-spin singlet and triplet states. We proceed with our analysis, but the results must be interpreted
 706 as iso-spin averaged scattering parameters.

707 As stated before, to generate a fit correlation function, we must simulate a large number of pairs, calculate
 708 the wave-function, and average Ψ^2 over all pairs in a given k^* bin. Essentially, we calculate Equation 9
 709 by hand:

$$\begin{aligned} C(\mathbf{k}^*) &= \sum_S \rho_S \int S(\mathbf{r}^*) |\Psi_{\mathbf{k}^*}^S(\mathbf{r}^*)|^2 d^3\mathbf{r}^* \\ &\longrightarrow C(|\mathbf{k}^*|) \equiv C(k^*) = \sum_S \rho_S \langle |\Psi^S(\mathbf{k}_i^*, \mathbf{r}_i^*)|^2 \rangle_i \\ &\longrightarrow C(k^*) = \lambda \sum_S \rho_S \langle |\Psi^S(\mathbf{k}_i^*, \mathbf{r}_i^*)|^2 \rangle_i + (1 - \lambda) \end{aligned} \quad (23)$$

710 where $\langle \rangle_i$ represents an average over all pairs in a given k^* bin.

711 In summary, for a given k^* bin, we must draw $N_{pairs} \sim 10^4$ pairs, and for each pair:

712 1. Draw a random \mathbf{r}^* vector according to our Gaussian source distribution $S(\mathbf{r}^*)$

713 2. Draw a random \mathbf{k}^* vector satisfying the $|\mathbf{k}^*|$ restriction of the bin

714 – We draw from real k^* vectors obtained from the data

715 – However, we find that drawing from a distribution flat in k^* gives similar results

716 3. Construct the wave-function Ψ

717 After all pairs for a given k^* bin are simulated and wave-functions obtained, the results are averaged to
 718 give the fit result.

719 Construction of the wave-functions, Equation 10, involves a number of complex functions not included
 720 in standard C++ or ROOT libraries (namely, $h(\eta)$, $\tilde{G}(\rho, \eta)$), and $F(-i\eta, 1, i\xi)$. These functions were
 721 even difficult to find and implement from elsewhere. Our solution was to embed a Mathematica kernel
 722 into our C++ code to evaluate these functions. However, having Mathematica work on-the-fly with the
 723 fitter was far too time consuming (fitter would have taken days, maybe weeks to finish). Our solution
 724 was to use Mathematica to create matrices representing these functions for different parameter values.
 725 During fitting, these matrices were then interpolated and the results used to build the wave-functions.
 726 This method decreased the running time dramatically, and we are not able to generate results in under \sim
 727 1 hour. This process will be explained in more detail in future versions of the note.

728 **6 Systematic Errors**

729 In order to understand my systematic uncertainties, the analysis code was run many times using slightly
 730 different values for a number of important cuts, and the results were compared. To quantify the systematic
 731 errors on the data, all correlation functions built using all varied cut values were bin-by-bin averaged,
 732 and the resulting variance of each bin was taken as the systematic error. The cuts which were utilized in
 733 this study are presented in Sections 6.1.1 (ΔK_S^0) and 6.2.1 (ΔK^\pm).

734 Similarly, the fit parameters extracted from all of these correlation functions were averaged, and the
 735 resulting variances were taken as the systematic errors for the fit parameters. As with the systematic
 736 errors on the data, this was performed for all varied cut values. Additionally, a systematic analysis was
 737 done on our fit method (which, for now, just includes our choice of fit range and functional form used to
 738 model the non-femtoscopic background). These two sources of uncertainty were combined in quadrature
 739 to obtain the final systematic uncertainties on the extracted fit parameters.

740 **6.1 Systematic Errors: ΔK_S^0**

741 **6.1.1 Particle and Pair Cuts**

742 The cuts included in the systematic study, as well as the values used in the variations, are listed below.
 743 Note, the central value corresponds to that used in the analysis.

| ΔK_S^0 systematics | |
|--|------------------------|
| DCA $\Lambda(\bar{\Lambda})$ | 4, 5, 6 mm |
| DCA K_S^0 | 2, 3, 4 mm |
| DCA $\Lambda(\bar{\Lambda})$ Daughters | 3, 4, 5 mm |
| DCA K_S^0 Daughters | 2, 3, 4 mm |
| $\Lambda(\bar{\Lambda})$ Cosine of Pointing Angle | 0.9992, 0.9993, 0.9994 |
| K_S^0 Cosine of Pointing Angle | 0.9992, 0.9993, 0.9994 |
| DCA to Primary Vertex of $p(\bar{p})$ Daughter of $\Lambda(\bar{\Lambda})$ | 0.5, 1, 2 mm |
| DCA to Primary Vertex of $\pi^-(\pi^+)$ Daughter of $\Lambda(\bar{\Lambda})$ | 2, 3, 4 mm |
| DCA to Primary Vertex of π^+ Daughter of K_S^0 | 2, 3, 4 mm |
| DCA to Primary Vertex of π^- Daughter of K_S^0 | 2, 3, 4 mm |
| Average Separation of Like-Charge Daughters | 5, 6, 7 cm |

Table 5: ΔK_S^0 systematics

744 **6.1.2 Non-Flat Background**

745 We fit our non-flat background with a linear function. To study the contribution of this choice to our sys-
 746 tematic errors, we also fit with a quadratic and Gaussian form. The resulting uncertainties are combined
 747 with the uncertainties arising from our particle cuts.

748 **6.1.3 Fit Range**

749 Our choice of k^* fit range was varied by $\pm 25\%$. The resulting uncertainties in the extracted parameter
 750 sets were combined with our uncertainties arising from our particle and pair cuts.

751 **6.2 Systematic Errors: ΛK^\pm**

752 **6.2.1 Particle and Pair Cuts**

753 The cuts included in the systematic study, as well as the values used in the variations, are listed below.
 754 Note, the central value corresponds to that used in the analysis.

| ΛK^\pm systematics | |
|---|------------------------|
| DCA $\Lambda(\bar{\Lambda})$ | 4, 5, 6 mm |
| DCA $\Lambda(\bar{\Lambda})$ Daughters | 3, 4, 5 mm |
| $\Lambda(\bar{\Lambda})$ Cosine of Pointing Angle | 0.9992, 0.9993, 0.9994 |
| DCA to Primary Vertex of $p(\bar{p})$ Daughter of $\Lambda(\bar{\Lambda})$ | 0.5, 1, 2 mm |
| DCA to Primary Vertex of $\pi^-(\pi^+)$ Daughter of $\Lambda(\bar{\Lambda})$ | 2, 3, 4 mm |
| Average Separation of $\Lambda(\bar{\Lambda})$ Daughter with Same Charge as K^\pm | 7, 8, 9 cm |
| Max. DCA to Primary Vertex in Transverse Plane of K^\pm | 1.92, 2.4, 2.88 |
| Max. DCA to Primary Vertex in Longitudinal Direction of K^\pm | 2.4, 3.0, 3.6 |

Table 6: ΛK^\pm systematics

755 **6.2.2 Non-Flat Background**

756 We fit our non-flat background with a linear function. To study the contribution of this choice to our sys-
 757 tematic errors, we also fit with a quadratic and Gaussian form. The resulting uncertainties are combined
 758 with the uncertainties arising from our particle cuts.

759 **6.2.3 Fit Range**

760 Our choice of k^* fit range was varied by $\pm 25\%$. The resulting uncertainties in the extracted parameter
 761 sets were combined with our uncertainties arising from our particle and pair cuts.

762 **6.3 Systematic Errors: ΞK^\pm**

763 **6.3.1 Particle and Pair Cuts**

764 The cuts included in the systematic study, as well as the values used in the variations, are listed below.
 765 Note, the central value corresponds to that used in the analysis.

$\Xi^- K^\pm$ systematics

| | |
|---|---------------------------|
| | $\Xi^- K^\pm$ systematics |
| Max. DCA $\Xi(\bar{\Xi})$ | 2, 3, 4 mm |
| Max. DCA $\Xi(\bar{\Xi})$ Daughters | 2, 3, 4 mm |
| Min. $\Xi(\bar{\Xi})$ Cosine of Pointing Angle to Primary Vertex | 0.9991, 0.9992, 0.9993 |
| Min. $\Lambda(\bar{\Lambda})$ Cosine of Pointing Angle to $\Xi(\bar{\Xi})$ Decay Vertex | 0.9992, 0.9993, 0.9994 |
| Min. DCA Bachelor π | 0.5, 1, 2 mm |
| Min. DCA $\Lambda(\bar{\Lambda})$ | 1, 2, 3 mm |
| Max. DCA $\Lambda(\bar{\Lambda})$ Daughters | 3, 4, 5 mm |
| Min. DCA to Primary Vertex of $p(\bar{p})$ Daughter of $\Lambda(\bar{\Lambda})$ | 0.5, 1, 2 mm |
| Min. DCA to Primary Vertex of $\pi^-(\pi^+)$ Daughter of $\Lambda(\bar{\Lambda})$ | 2, 3, 4 mm |
| Min. Average Separation of $\Lambda(\bar{\Lambda})$ Daughter and K^\pm with like charge | 7, 8, 9 cm |
| Min. Average Separation of Bachelor π and K^\pm with like charge | 7, 8, 9 cm |
| Max. DCA to Primary Vertex in Transverse Plane of K^\pm | 1.92, 2.4, 2.88 |
| Max. DCA to Primary Vertex in Longitudinal Direction of K^\pm | 2.4, 3.0, 3.6 |

Table 7: $\Xi^- K^\pm$ systematics

766 **7 Results and Discussion**

767 **7.1 Results: ΛK_S^0 and ΛK^\pm**

768 In the following sections, we present results assuming (i) no residual correlations (Sec. 7.1.3), (ii) three
 769 residual contributors (Sec. 7.1.1), and (iii) ten residual contributors (Sec. 7.1.2). We find the case of
 770 three and ten contributors to be consistent; therefore, for simplicity, we will quote the result utilizing
 771 three residuals as our final result.

772 For the results shown, unless otherwise noted, the following hold true: All correlation functions were
 773 normalized in the range $0.32 < k^* < 0.40$ GeV/c, and fit in the range $0.0 < k^* < 0.30$ GeV/c. For the
 774 ΛK^- and $\bar{\Lambda} K^+$ analyses, the region $0.19 < k^* < 0.23$ GeV/c was excluded from the fit to exclude the
 775 bump caused by the Ω^- resonance. The non-femtoscopic background was modeled by a (6th-order
 776 polynomial fit to THERMINATOR simulation. The $\Lambda K^+(\bar{\Lambda} K^-)$ radii are shared with $\Lambda K^-(\bar{\Lambda} K^+)$, while
 777 the $\Lambda K_S^0(\bar{\Lambda} K_S^0)$ radii are unique. In the figures showing experimental correlation functions with fits, the
 778 black solid line represents the primary (ΛK) correlation's contribution to the fit. The green line shows
 779 the fit to the non-flat background. The purple points show the fit after all residuals' contributions have
 780 been included, and momentum resolution and non-flat background corrections have been applied.

781 Before beginning, I first collect a summary of my final results in Figure 30. In the summary plot, we
 782 show the extracted scattering parameters in the form of a $\text{Im}[f_0]$ vs $\text{Re}[f_0]$ plot, which includes the d_0
 783 values to the right side. We also show the λ vs. radius parameters for all three of our studied centrality
 784 bins. In Fig. 30, three residual contributors were used, and the background was modeled by a (6th-
 785 order polynomial fit to THERMINATOR simulation. For the ΛK_S^0 results shown in the figure, the ΛK_S^0
 786 and $\bar{\Lambda} K_S^0$ analyses were fit simultaneously across all centralities (0-10%, 10-30%, 30-50%); scattering
 787 parameters and a single λ parameter were shared amongst all, the radii were shared amongst results
 788 of like-centrality, and each has a unique normalization parameter. For the ΛK^\pm results shown, all four
 789 pair combinations were fit simultaneously ($\Lambda K^+, \bar{\Lambda} K^-, \Lambda K^-, \bar{\Lambda} K^+$) across all centralities. Scattering
 790 parameters were shared between pair-conjugate systems (i.e. a parameter set describing ΛK^+ & $\bar{\Lambda} K^-$,
 791 and a separate set describing ΛK^- & $\bar{\Lambda} K^+$). For each centrality, a radius and λ parameters were shared
 792 between all pairs. Each analysis has a unique normalization parameter.

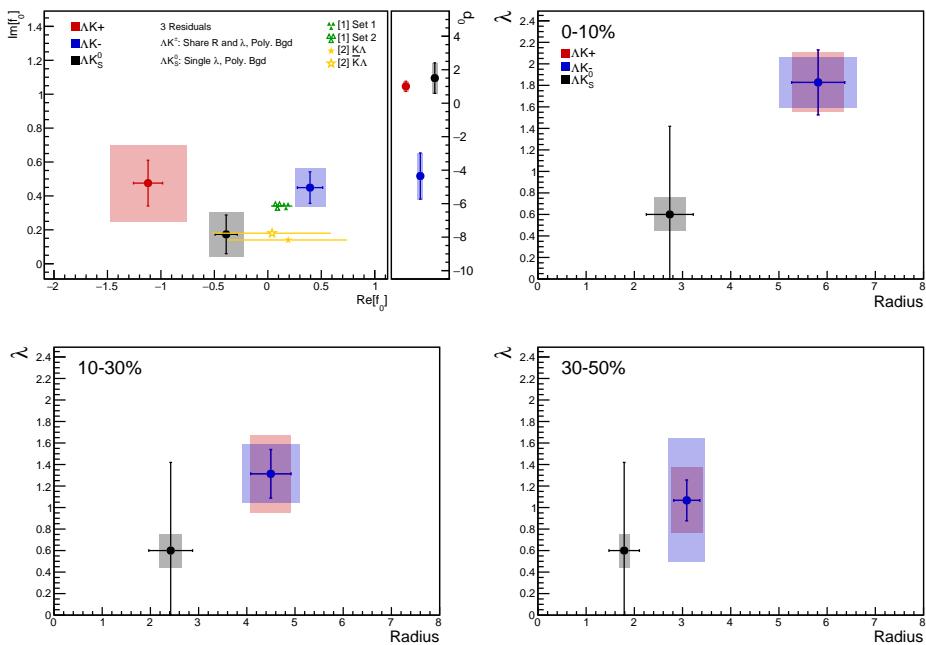


Fig. 30: Extracted scattering parameters for the case of 3 residual contributors for all of our ΛK systems. [Top Left]: $\text{Im}[f_0]$ vs. $\text{Re}[f_0]$, together with d_0 to the right. [Top Right (Bottom Left, Bottom Right)]: λ vs. Radius for the 0-10% (10-30%, 30-50%) bin. The green [9] and yellow [10] points show theoretical predictions made using chiral perturbation theory.

793 **7.1.1 3 Residual Correlations Included in Fit**

794 Figure 31 nicely collects and summarizes all of our extracted fit parameters for the case of 3 included
 795 residual contributors. Figure 32 presents our extracted fit radii, along with those of other systems pre-
 796 viously analyzed by ALICE [11], as a function of pair transverse mass (m_T). Figures 33, 34, and 35
 797 show the experimental correlation functions with fits, assuming 3 residual contributors, for all studied
 798 centralities for ΛK_S^0 with $\bar{\Lambda} K_S^0$, ΛK^+ with $\bar{\Lambda} K^-$, and ΛK^- with $\bar{\Lambda} K^+$, respectively. The parameter sets
 799 extracted from the fits can be found in Tables 12 and 13.

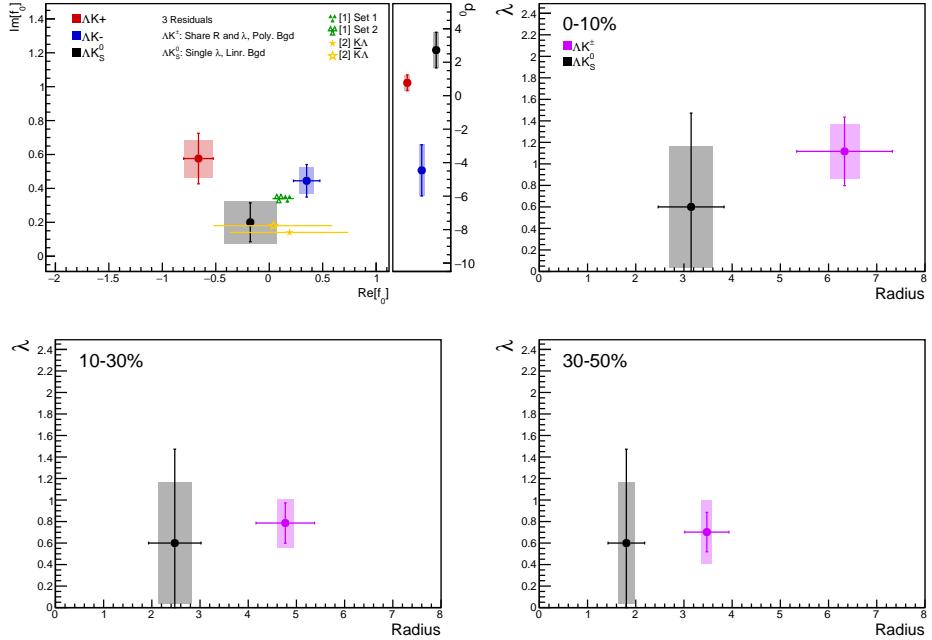


Fig. 31: Extracted scattering parameters for the case of 3 residual contributors for all of our ΛK systems. [Top Left]: $\text{Im}(f_0)$ vs. $\text{Re}(f_0)$, together with d_0 to the right. [Top Right (Bottom Left, Bottom Right)]: λ vs. Radius for the 0-10% (10-30%, 30-50%) bin. The green [9] and yellow [10] points show theoretical predictions made using chiral perturbation theory.

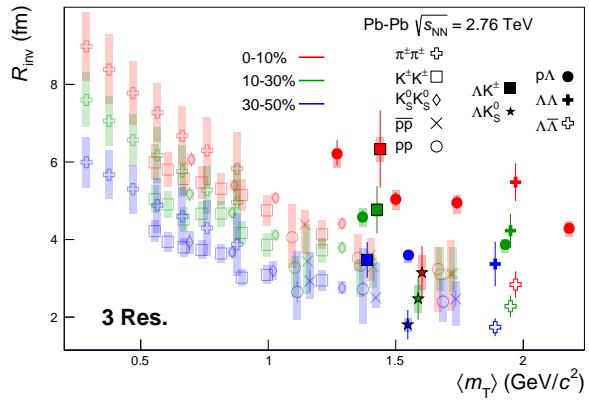
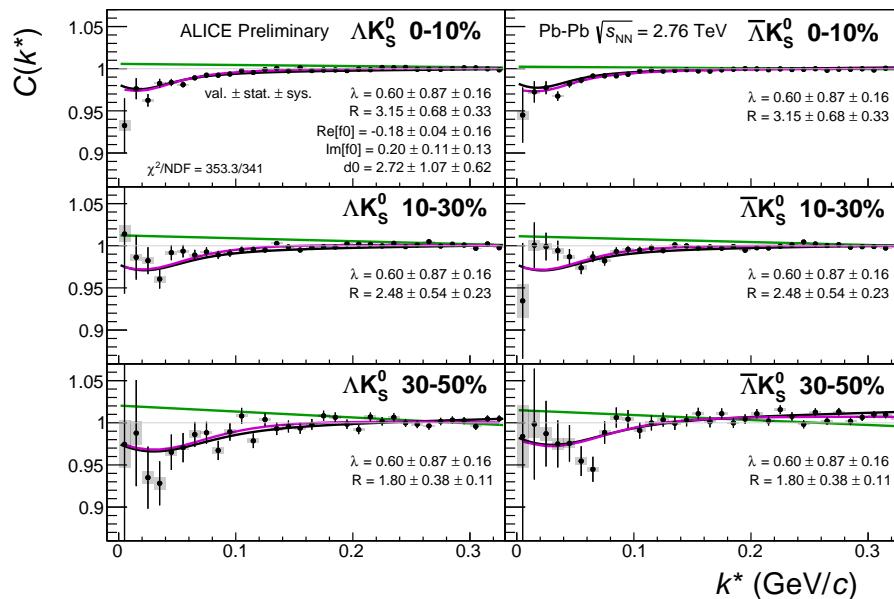
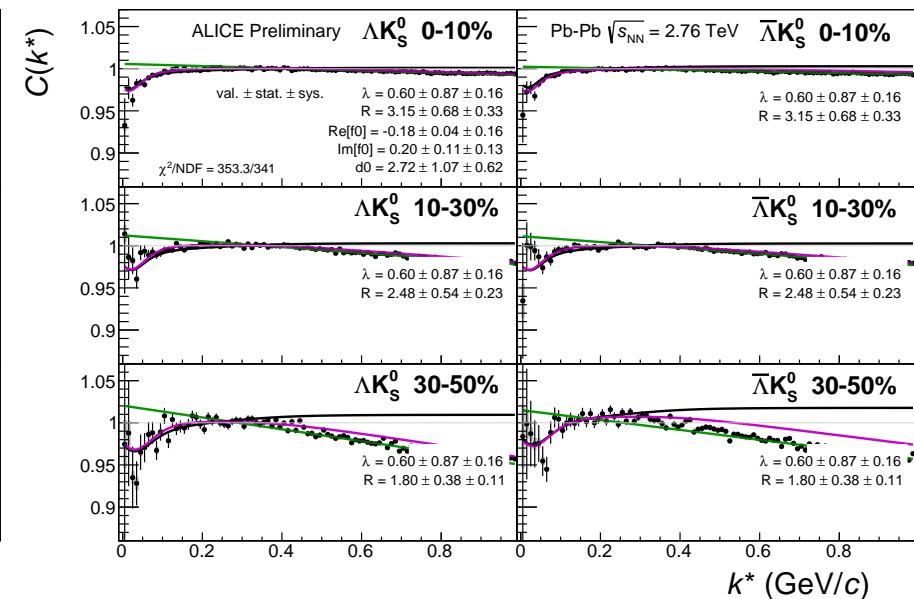


Fig. 32: 3 residual correlations in ΛK fits. Extracted fit R_{inv} parameters as a function of pair transverse mass (m_T) for various pair systems over several centralities. The ALICE published data [11] is shown with transparent, open symbols. The new ΛK results are shown with opaque, filled symbols. In the left, the ΛK^+ (with its conjugate pair) results are shown separately from the ΛK^- (with its conjugate pair) results. In the right, all ΛK^\pm results are averaged.

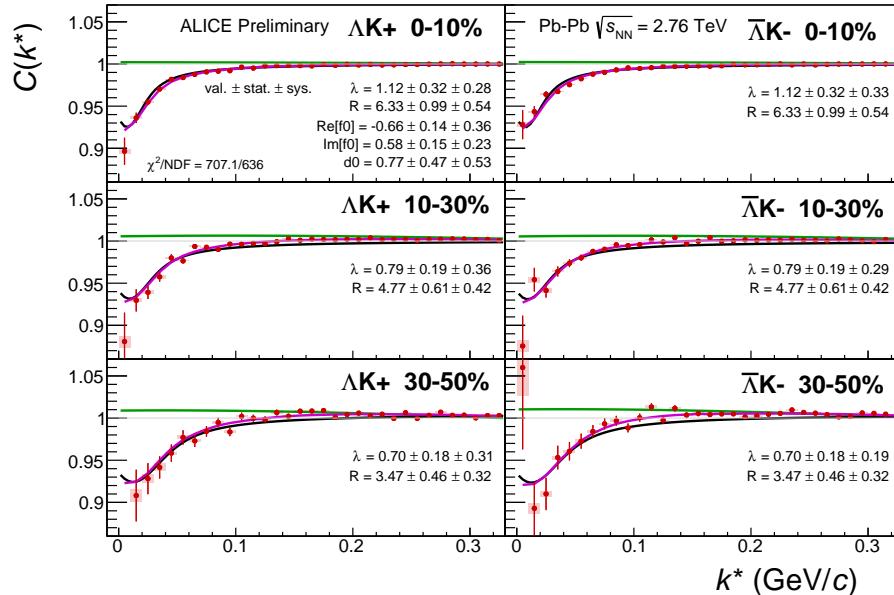


(a) Signal region view ($k^* \lesssim 0.3 \text{ GeV}/c$)

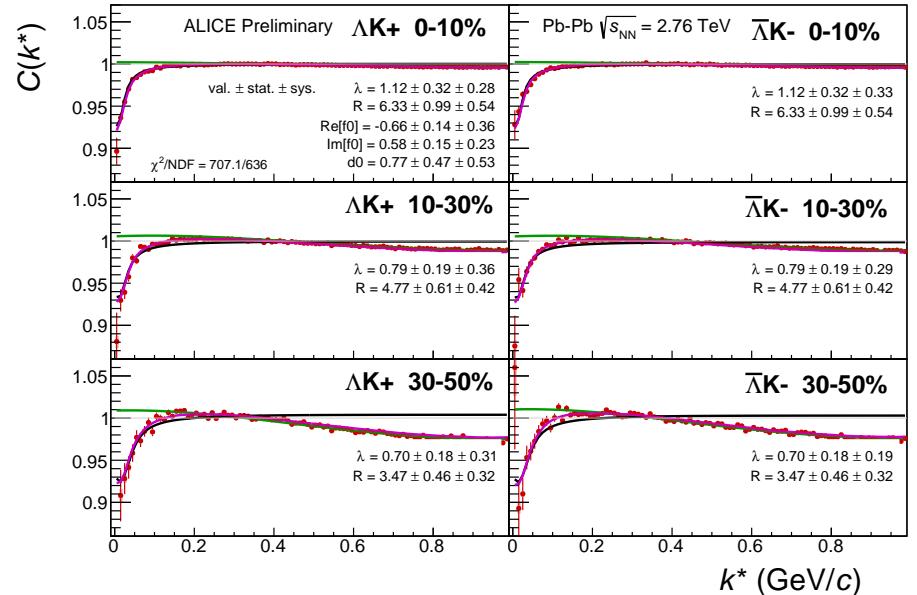


(b) Wide view ($k^* \lesssim 1.0 \text{ GeV}/c$)

Fig. 33: Fits, with 3 residual correlations included, to the ΛK_s^0 (left) and $\bar{\Lambda} K_s^0$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The lines represent the statistical errors, while the boxes represent the systematic errors. A single λ parameter is shared amongst all. Each analysis has a unique normalization parameter. The radii are shared between analyses of like centrality, as these should have similar source sizes. The scattering parameters ($\mathbb{R} f_0$, $\mathbb{I} f_0$, d_0) are shared amongst all. The background is modeled by a (6th-degree polynomial fit to THERMINATOR simulation. The black solid line represents the primary (ΛK) correlation's contribution to the fit. The green line shows the fit to the non-flat background. The purple points show the fit after all residuals' contributions have been included, and momentum resolution and non-flat background corrections have been applied. The extracted fit values with uncertainties are printed.

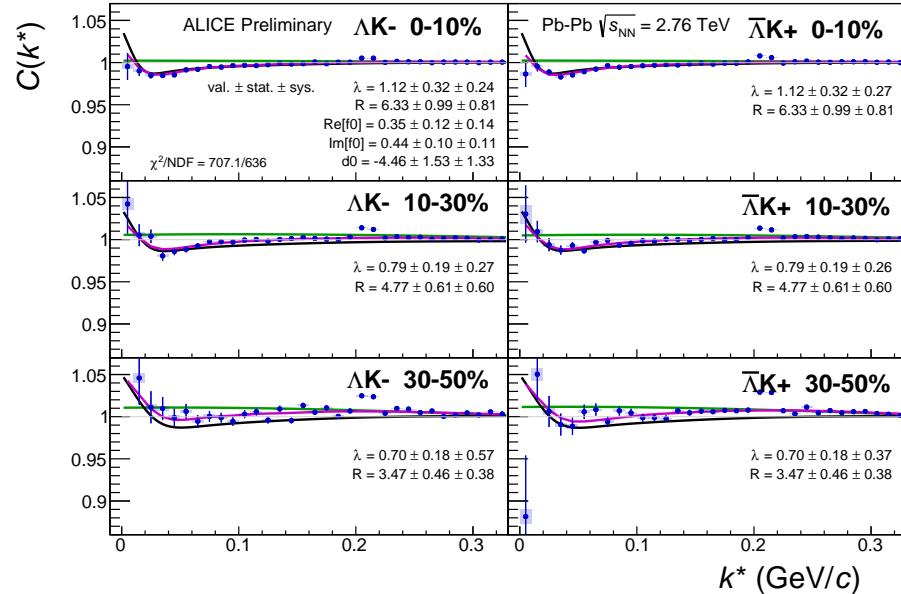


(a) Signal region view ($k^* \lesssim 0.3 \text{ GeV}/c$)

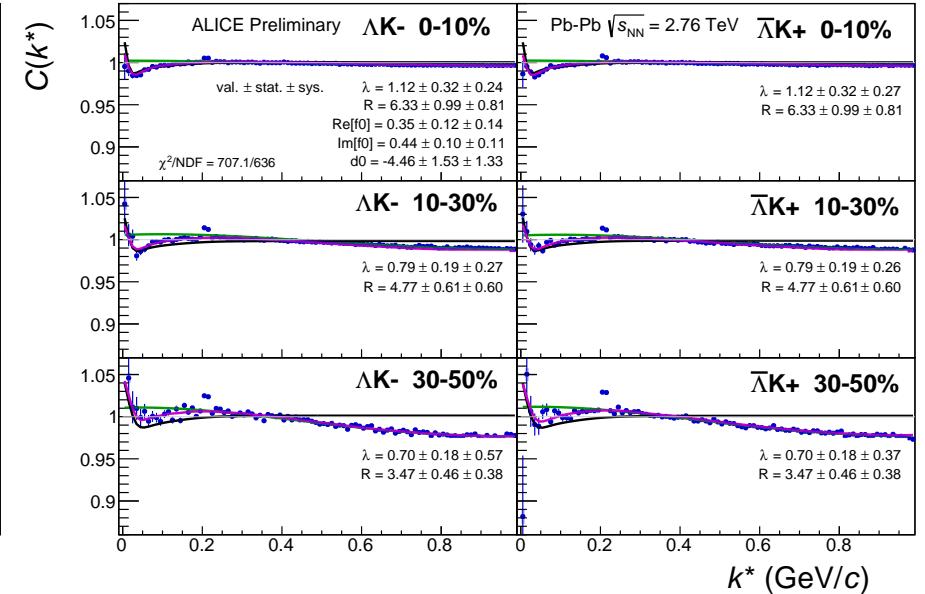


(b) Wide view ($k^* \lesssim 1.0 \text{ GeV}/c$)

Fig. 34: Fits, with 3 residual correlations included, to the ΛK^+ (left) and $\bar{\Lambda} K^-$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The lines represent the statistical errors, while the boxes represent the systematic errors. All ΛK^\pm analyses are fit simultaneously across all centralities (0-10%, 10-30%, 30-50%). Scattering parameters ($\text{Re}[f_0]$, $\text{Im}[f_0]$, d_0) are shared between pair-conjugate systems (i.e. a parameter set describing the ΛK^+ & $\bar{\Lambda} K^-$ system, and a separate set describing the ΛK^- & $\bar{\Lambda} K^+$ system). For each centrality, a radius and λ parameters are shared between all pairs (ΛK^+ , $\bar{\Lambda} K^-$, ΛK^- , $\bar{\Lambda} K^+$). Each analysis has a unique normalization parameter. The background is modeled by a (6th-degree polynomial fit to THERMINATOR simulation. The black solid line represents the primary (ΛK) correlation's contribution to the fit. The green line shows the fit to the non-flat background. The purple points show the fit after all residuals' contributions have been included, and momentum resolution and non-flat background corrections have been applied. The extracted fit values with uncertainties are printed.



(a) Signal region view ($k^* \lesssim 0.3 \text{ GeV}/c$)



(b) Wide view ($k^* \lesssim 1.0 \text{ GeV}/c$)

Fig. 35: Fits, with 3 residual correlations included, to the ΛK^- (left) with $\bar{\Lambda} K^+$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The lines represent the statistical errors, while the boxes represent the systematic errors. All ΛK^\pm analyses are fit simultaneously across all centralities (0-10%, 10-30%, 30-50%). Scattering parameters ($\mathbb{R}f_0$, $\mathbb{I}f_0$, d_0) are shared between pair-conjugate systems (i.e. a parameter set describing the ΛK^+ & $\bar{\Lambda} K^-$ system, and a separate set describing the ΛK^- & $\bar{\Lambda} K^+$ system). For each centrality, a radius and λ parameters are shared between all pairs (ΛK^+ , $\bar{\Lambda} K^-$, ΛK^- , $\bar{\Lambda} K^+$). Each analysis has a unique normalization parameter. The background is modeled by a (6th-)degree polynomial fit to THERMINATOR simulation. The black solid line represents the “raw” fit, i.e. not corrected for momentum resolution effects nor non-flat background. The green line shows the fit to the non-flat background. The purple points show the fit after momentum resolution and non-flat background corrections have been applied. The extracted fit values with uncertainties are printed.

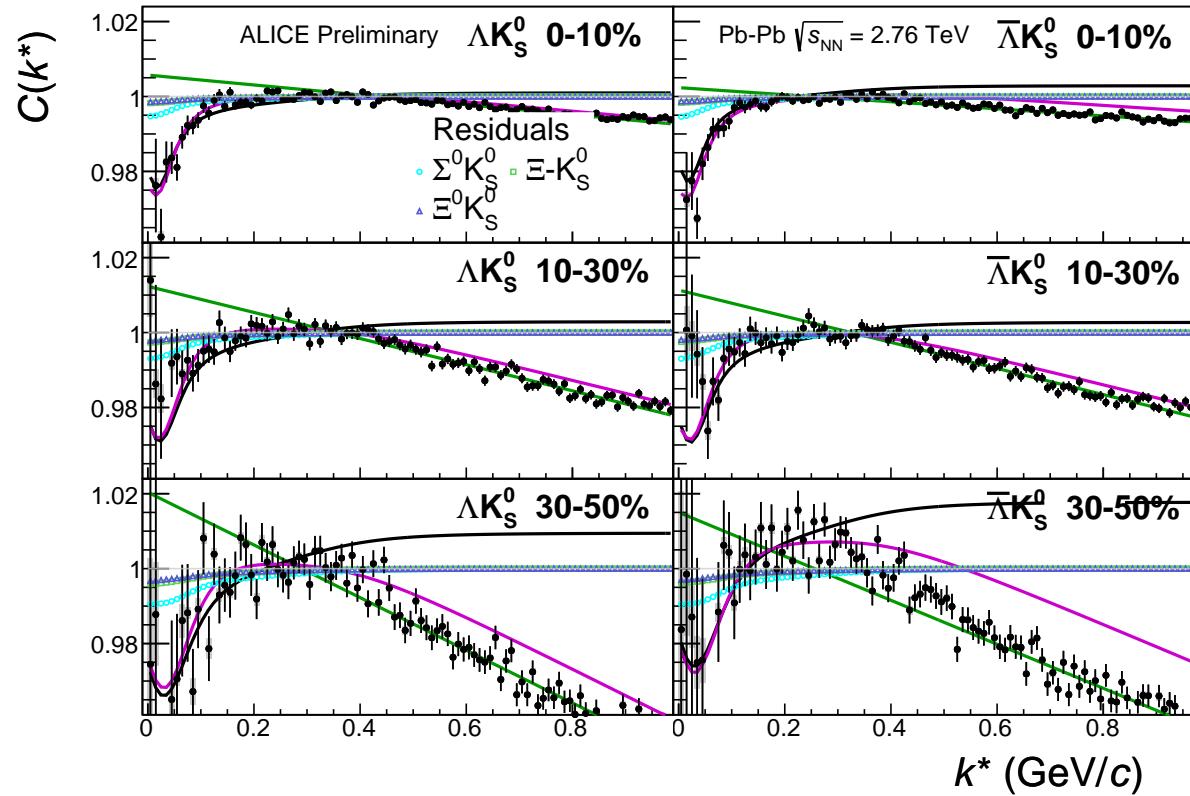
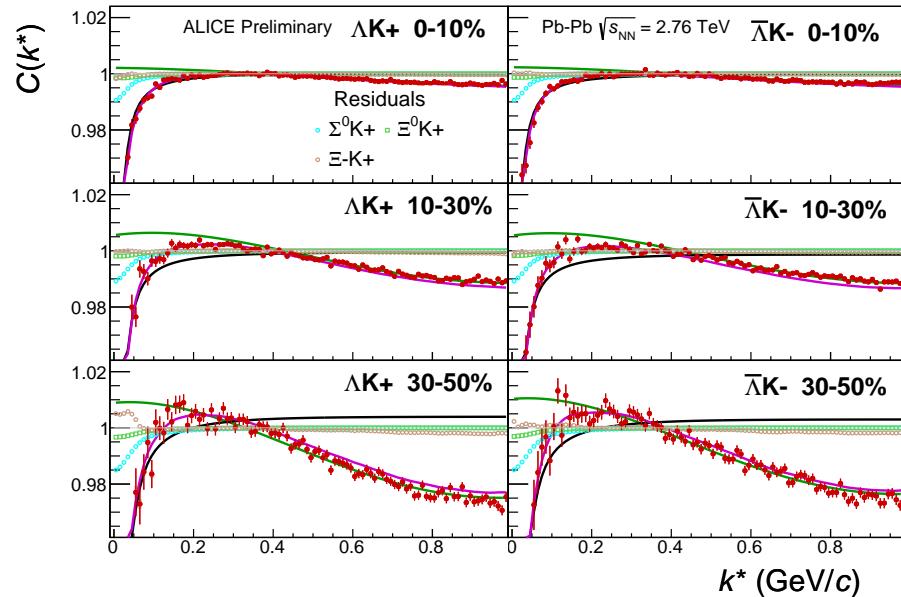
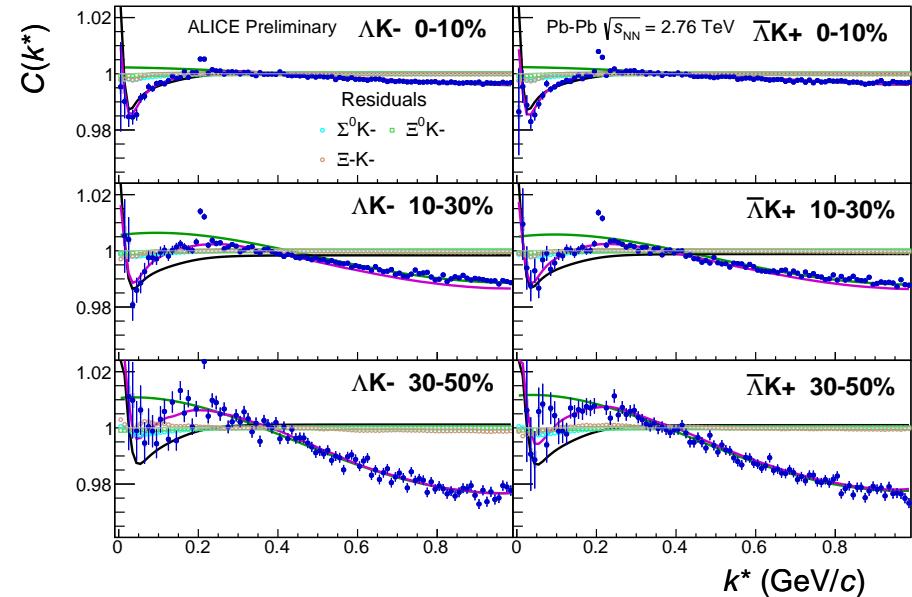


Fig. 36: Fits, with 3 residual correlations included and shown, to the ΛK_S^0 (left) and $\bar{\Lambda} K_S^0$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The three parent pairs used for the residual correction to the ΛK_S^0 ($\bar{\Lambda} K_S^0$) fit are $\Sigma^0 K_S^0$, $\Xi^0 K_S^0$, and $\Xi^- K_S^0$ ($\bar{\Sigma}^0 K_S^0$, $\bar{\Xi}^0 K_S^0$, and $\bar{\Xi}^+ K_S^0$).



(a) $\Lambda K^+(\bar{\Lambda}K^-)$ fits with residual contributions shown for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom)



(b) $\Lambda K^-(\bar{\Lambda}K^+)$ fits with residual contributions shown for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom)

Fig. 37: Fits, with 3 residual correlations included and shown, to the ΛK^+ & $\bar{\Lambda}K^-$ (left) and ΛK^- & $\bar{\Lambda}K^+$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The three parent pairs used for the residual correction to the $\Lambda K^+(\bar{\Lambda}K^-)$ fit are $\Sigma^0 K^+$, $\Xi^0 K^+$, and $\Xi^- K^-$ ($\bar{\Sigma}^0 K^-$, $\bar{\Xi}^0 K^-$, and $\bar{\Xi}^+ K^-$).

| Fit Results $\Lambda(\bar{\Lambda})K_S^0$ | | | | | | |
|--|------------|---|---|--|---|---|
| System | Centrality | Fit Parameters | | | | |
| | | λ | R | $\mathbb{R}f_0$ | $\mathbb{I}f_0$ | d_0 |
| $\Lambda K_S^0 \text{ & } \bar{\Lambda} K_S^0$ | 0-10% | | | $3.15 \pm 0.68 \text{ (stat.)} \pm 0.45 \text{ (sys.)}$ | | |
| | 10-30% | $0.60 \pm 0.87 \text{ (stat.)} \pm 0.57 \text{ (sys.)}$ | $2.48 \pm 0.54 \text{ (stat.)} \pm 0.35 \text{ (sys.)}$ | $-0.18 \pm 0.04 \text{ (stat.)} \pm 0.25 \text{ (sys.)}$ | $0.20 \pm 0.11 \text{ (stat.)} \pm 0.13 \text{ (sys.)}$ | $2.72 \pm 1.07 \text{ (stat.)} \pm 2.12 \text{ (sys.)}$ |
| | 30-50% | | | $1.80 \pm 0.38 \text{ (stat.)} \pm 0.17 \text{ (sys.)}$ | | |

Table 8: Fit Results $\Lambda(\bar{\Lambda})K_S^0$, with 3 residual correlations included. Each pair is fit simultaneously with its conjugate (ie. ΛK_S^0 with $\bar{\Lambda} K_S^0$) across all centralities (0-10%, 10-30%, 30-50%), for a total of 6 simultaneous analyses in the fit. A single λ parameter is shared amongst all. Each analysis has a unique normalization parameter. The radii are shared between analyses of like centrality, as these should have similar source sizes. The scattering parameters ($\mathbb{R}f_0$, $\mathbb{I}f_0$, d_0) are shared amongst all. The background is fit with a linear form in the range $0.6 < k^* < 0.9 \text{ GeV}/c$. The fit is done on the data with only statistical error bars. The errors marked as “stat.” are those returned by MINUIT. The errors marked as “sys.” are those which result from my systematic analysis (as outlined in Section 6).

| Fit Results $\Lambda(\bar{\Lambda})K^\pm$ | | | | | | |
|--|------------|---|---|--|---|--|
| System | Centrality | Fit Parameters | | | | |
| | | λ | R | $\mathbb{R}f_0$ | $\mathbb{I}f_0$ | d_0 |
| $\Lambda K^+ \text{ & } \bar{\Lambda} K^-$ | 0-10% | $1.12 \pm 0.32 \text{ (stat.)} \pm 0.25 \text{ (sys.)}$ | $6.33 \pm 0.99 \text{ (stat.)} \pm 0.31 \text{ (sys.)}$ | $-0.66 \pm 0.14 \text{ (stat.)} \pm 0.13 \text{ (sys.)}$ | $0.58 \pm 0.15 \text{ (stat.)} \pm 0.11 \text{ (sys.)}$ | $0.77 \pm 0.47 \text{ (stat.)} \pm 1.66 \text{ (sys.)}$ |
| | 10-30% | $0.79 \pm 0.19 \text{ (stat.)} \pm 0.23 \text{ (sys.)}$ | $4.77 \pm 0.61 \text{ (stat.)} \pm 0.17 \text{ (sys.)}$ | | | |
| $\Lambda K^+ \text{ & } \bar{\Lambda} K^-$ | 30-50% | $0.70 \pm 0.18 \text{ (stat.)} \pm 0.30 \text{ (sys.)}$ | $3.47 \pm 0.46 \text{ (stat.)} \pm 0.10 \text{ (sys.)}$ | $0.35 \pm 0.12 \text{ (stat.)} \pm 0.07 \text{ (sys.)}$ | $0.44 \pm 0.10 \text{ (stat.)} \pm 0.08 \text{ (sys.)}$ | $-4.46 \pm 1.53 \text{ (stat.)} \pm 1.36 \text{ (sys.)}$ |

Table 9: Fit Results $\Lambda(\bar{\Lambda})K^\pm$, with 3 residual correlations included. All ΛK^\pm analyses are fit simultaneously across all centralities (0-10%, 10-30%, 30-50%). Scattering parameters ($\mathbb{R}f_0$, $\mathbb{I}f_0$, d_0) are shared between pair-conjugate systems (i.e. a parameter set describing the $\Lambda K^+ \text{ & } \bar{\Lambda} K^-$ system, and a separate set describing the $\Lambda K^- \text{ & } \bar{\Lambda} K^+$ system). For each centrality, a radius and λ parameters are shared between all pairs ($\Lambda K^+, \bar{\Lambda} K^-, \Lambda K^-, \bar{\Lambda} K^+$). Each analysis has a unique normalization parameter. The background is modeled by a (6th-degree polynomial fit to THERMINATOR simulation. The fit is done on the data with only statistical error bars. The errors marked as “stat.” are those returned by MINUIT. The errors marked as “sys.” are those which result from my systematic analysis (as outlined in Section 6).

800 7.1.2 10 Residual Correlations Included in Fit

801 Figure 38 nicely collects and summarizes all of our extracted fit parameters for the case of 10 included
 802 residual contributors. Figure 39 presents our extracted fit radii, along with those of other systems pre-
 803 viously analyzed by ALICE [11], as a function of pair transverse mass (m_T). Figures 40, 41, and 42
 804 show the experimental correlation functions with fits, assuming 10 residual contributors, for all studied
 805 centralities for ΛK_S^0 with $\bar{\Lambda} K_S^0$, ΛK^+ with $\bar{\Lambda} K^-$, and ΛK^- with $\bar{\Lambda} K^+$, respectively. The parameter sets
 806 extracted from the fits can be found in Tables ?? and ??.

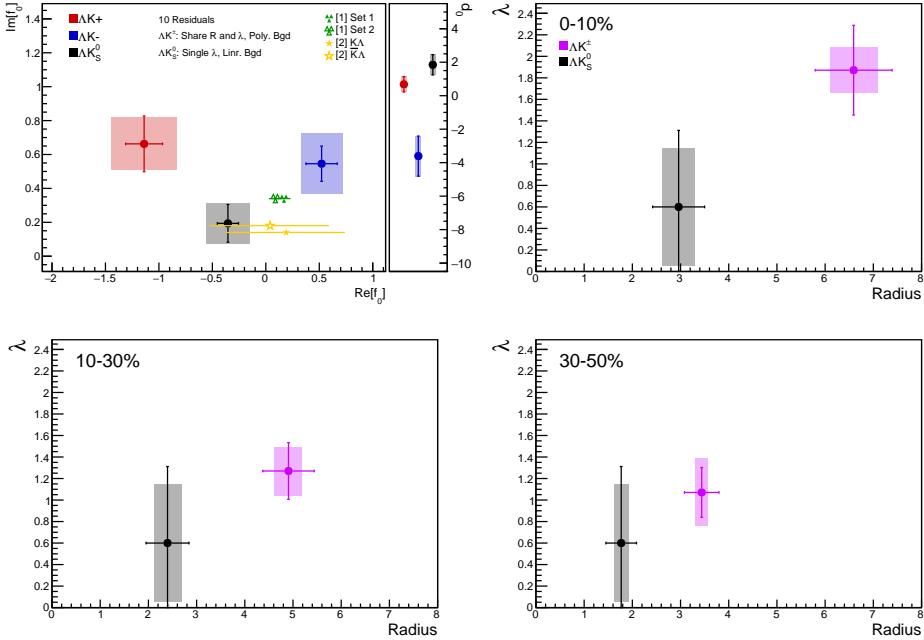


Fig. 38: Extracted scattering parameters for the case of 10 residual contributors for all of our AK systems. [Top Left]: $\text{Im}[f_0]$ vs. $\text{Re}[f_0]$, together with d_0 to the right. [Top Right (Bottom Left, Bottom Right)]: λ vs. Radius for the 0-10% (10-30%, 30-50%) bin. The green [9] and yellow [10] points show theoretical predictions made using chiral perturbation theory.

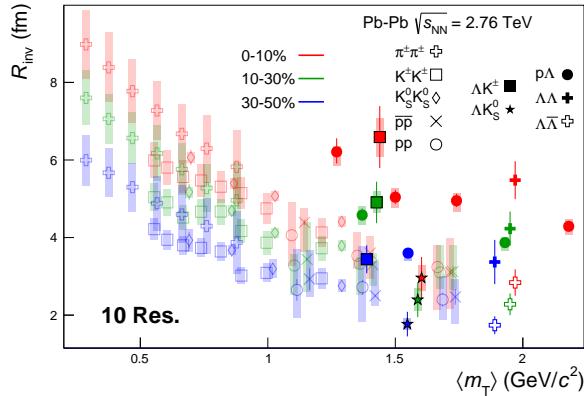
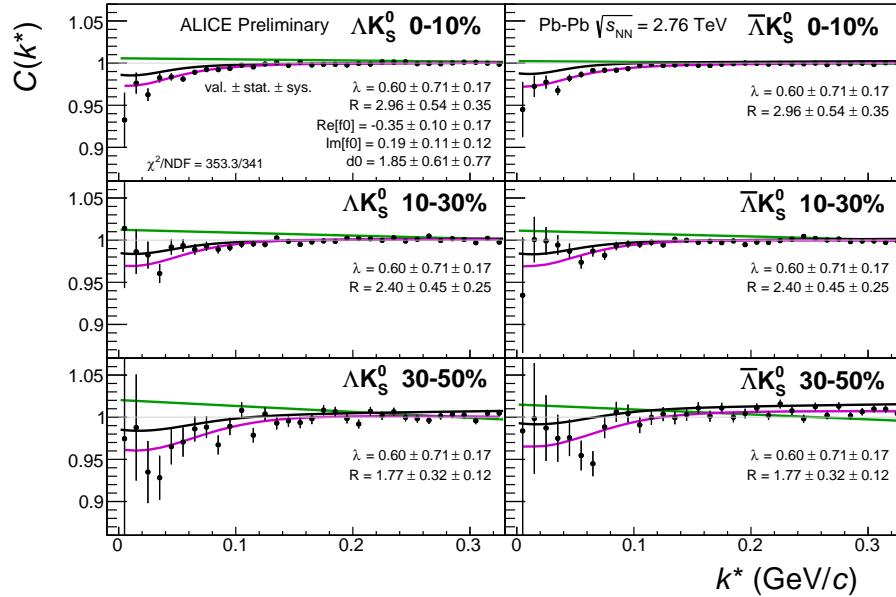
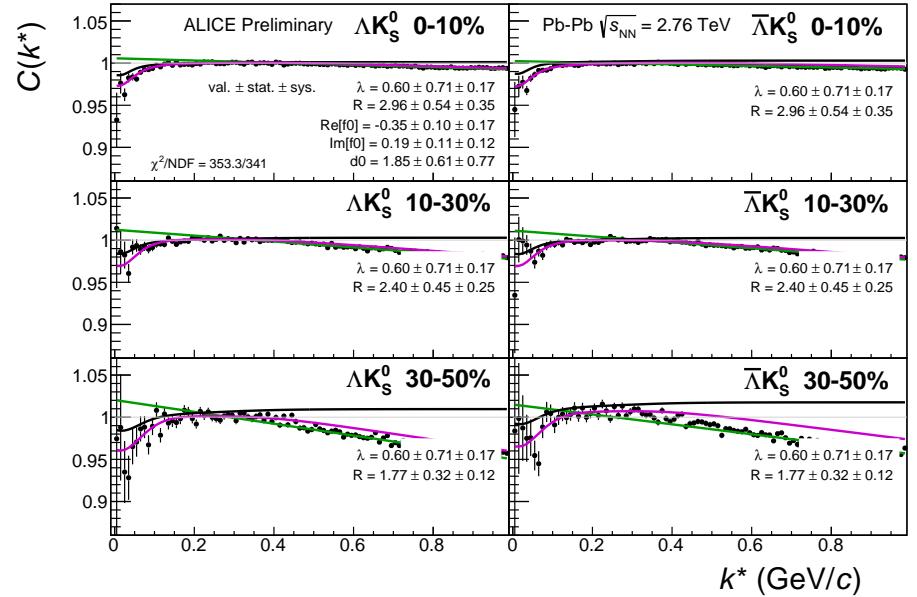


Fig. 39: 10 residual correlations in AK fits. Extracted fit R_{inv} parameters as a function of pair transverse mass (m_T) for various pair systems over several centralities. The ALICE published data [11] is shown with transparent, open symbols. The new AK results are shown with opaque, filled symbols. In the left, the ΛK^+ (with its conjugate pair) results are shown separately from the ΛK^- (with its conjugate pair) results. In the right, all ΛK^\pm results are averaged.

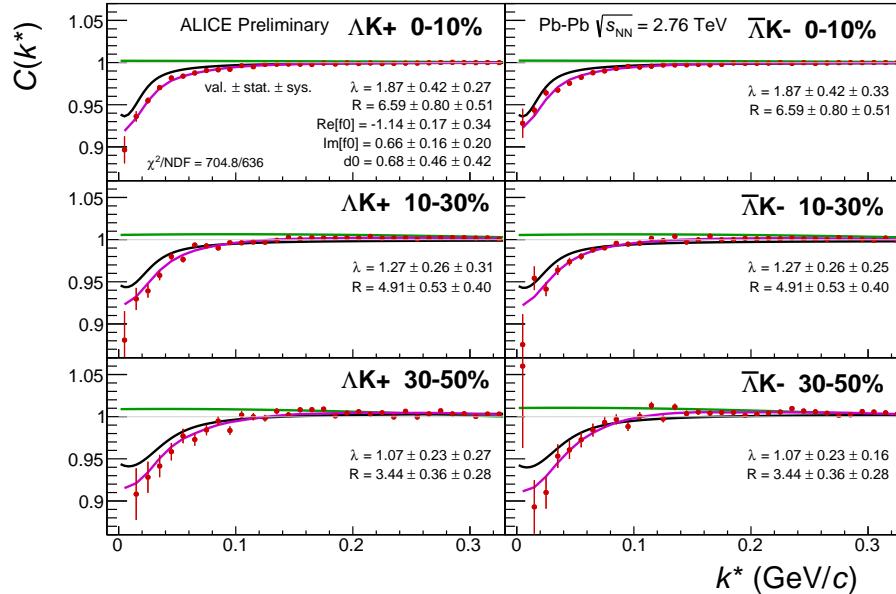


(a) Signal region view ($k^* \lesssim 0.3 \text{ GeV}/c$)

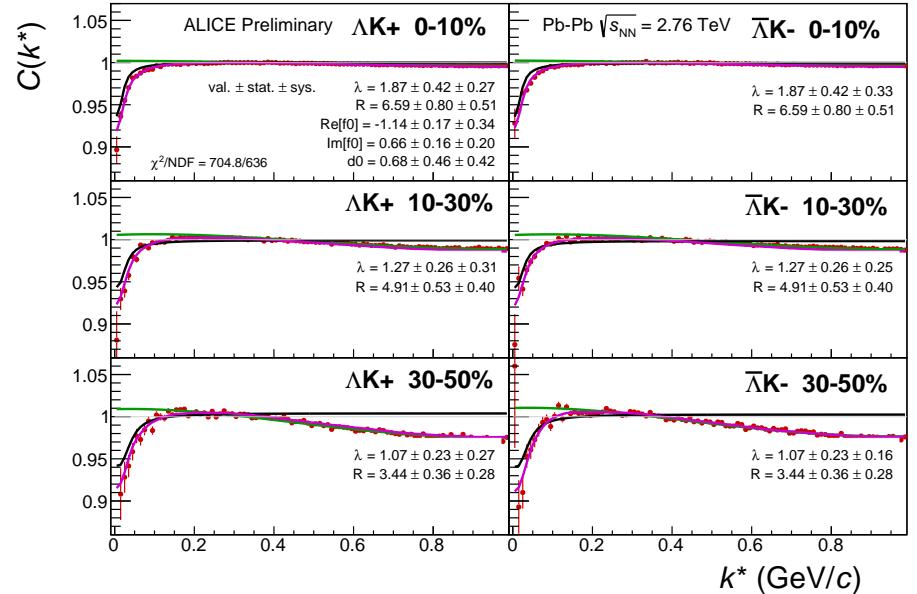


(b) Wide view ($k^* \lesssim 1.0 \text{ GeV}/c$)

Fig. 40: Fits, with 10 residual correlations included, to the ΛK_s^0 (left) and $\bar{\Lambda} K_s^0$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The lines represent the statistical errors, while the boxes represent the systematic errors. A single λ parameter is shared amongst all. Each analysis has a unique normalization parameter. The radii are shared between analyses of like centrality, as these should have similar source sizes. The scattering parameters ($\mathbb{R} f_0$, $\mathbb{I} f_0$, d_0) are shared amongst all. The background is modeled by a (6th-degree polynomial fit to THERMINATOR simulation. The black solid line represents the primary (ΛK) correlation's contribution to the fit. The green line shows the fit to the non-flat background. The purple points show the fit after all residuals' contributions have been included, and momentum resolution and non-flat background corrections have been applied. The extracted fit values with uncertainties are printed.

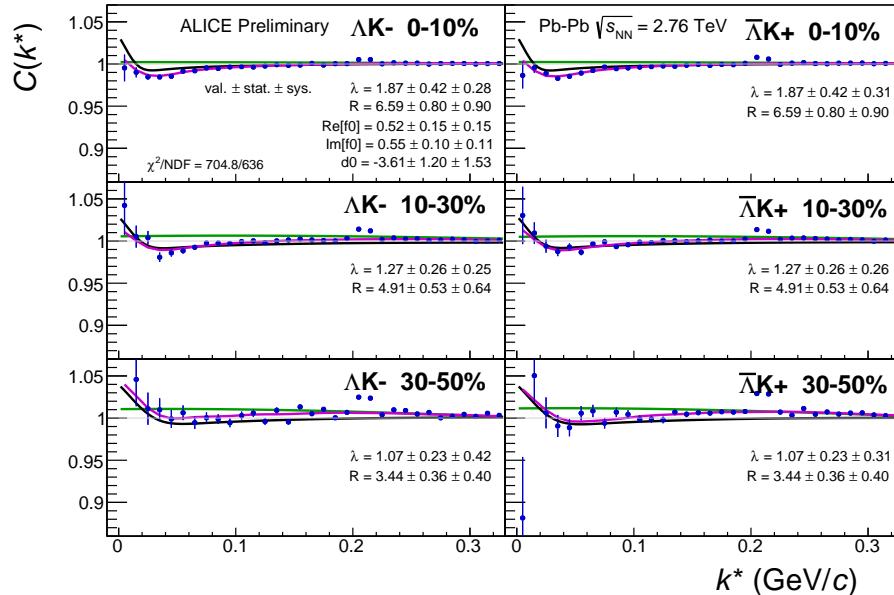


(a) Signal region view ($k^* \lesssim 0.3 \text{ GeV}/c$)

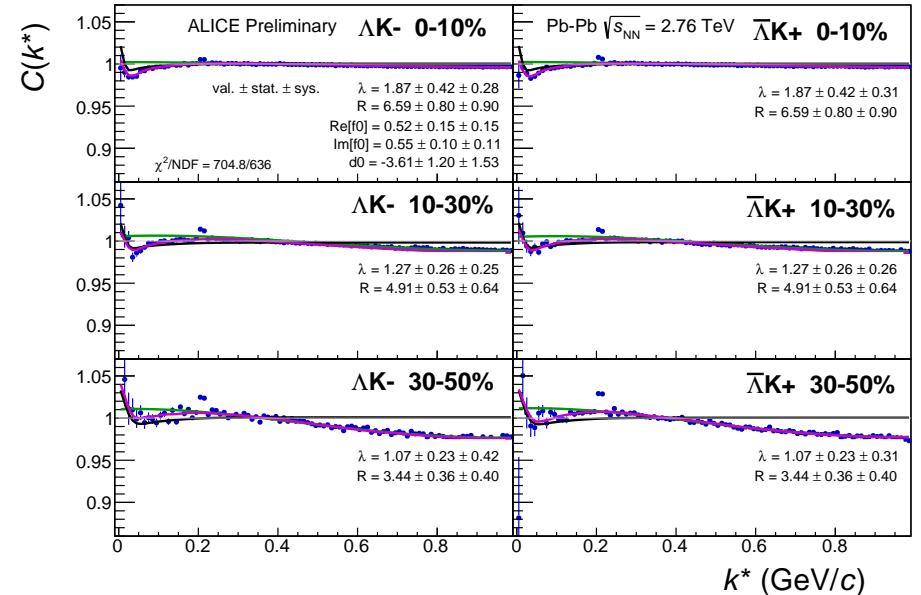


(b) Wide view ($k^* \lesssim 1.0 \text{ GeV}/c$)

Fig. 41: Fits, with 10 residual correlations included, to the ΔK^+ (left) and $\bar{\Delta}K^-$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The lines represent the statistical errors, while the boxes represent the systematic errors. All ΔK^\pm analyses are fit simultaneously across all centralities (0-10%, 10-30%, 30-50%). Scattering parameters ($\text{Re}[f_0]$, $\text{Im}[f_0]$, d_0) are shared between pair-conjugate systems (i.e. a parameter set describing the ΔK^+ & $\bar{\Delta}K^-$ system, and a separate set describing the ΔK^- & $\bar{\Delta}K^+$ system). For each centrality, a radius and λ parameters are shared between all pairs (ΔK^+ , $\bar{\Delta}K^-$, ΔK^- , $\bar{\Delta}K^+$). Each analysis has a unique normalization parameter. The background is modeled by a (6th-degree polynomial fit to THERMINATOR simulation. The black solid line represents the primary (ΔK) correlation's contribution to the fit. The green line shows the fit to the non-flat background. The purple points show the fit after all residuals' contributions have been included, and momentum resolution and non-flat background corrections have been applied. The extracted fit values with uncertainties are printed.



(a) Signal region view ($k^* \lesssim 0.3 \text{ GeV}/c$)



(b) Wide view ($k^* \lesssim 1.0 \text{ GeV}/c$)

Fig. 42: Fits, with 10 residual correlations included, to the ΛK^- (left) with $\bar{\Lambda} K^+$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The lines represent the statistical errors, while the boxes represent the systematic errors. All ΛK^\pm analyses are fit simultaneously across all centralities (0-10%, 10-30%, 30-50%). Scattering parameters ($R f_0$, $I f_0$, d_0) are shared between pair-conjugate systems (i.e. a parameter set describing the ΛK^+ & $\bar{\Lambda} K^-$ system, and a separate set describing the ΛK^- & $\bar{\Lambda} K^+$ system). For each centrality, a radius and λ parameters are shared between all pairs (ΛK^+ , $\bar{\Lambda} K^-$, ΛK^- , $\bar{\Lambda} K^+$). Each analysis has a unique normalization parameter. The background is modeled by a (6th-degree polynomial fit to THERMINATOR simulation. The black solid line represents the primary (ΛK) correlation's contribution to the fit. The green line shows the fit to the non-flat background. The purple points show the fit after all residuals' contributions have been included, and momentum resolution and non-flat background corrections have been applied. The extracted fit values with uncertainties are printed.

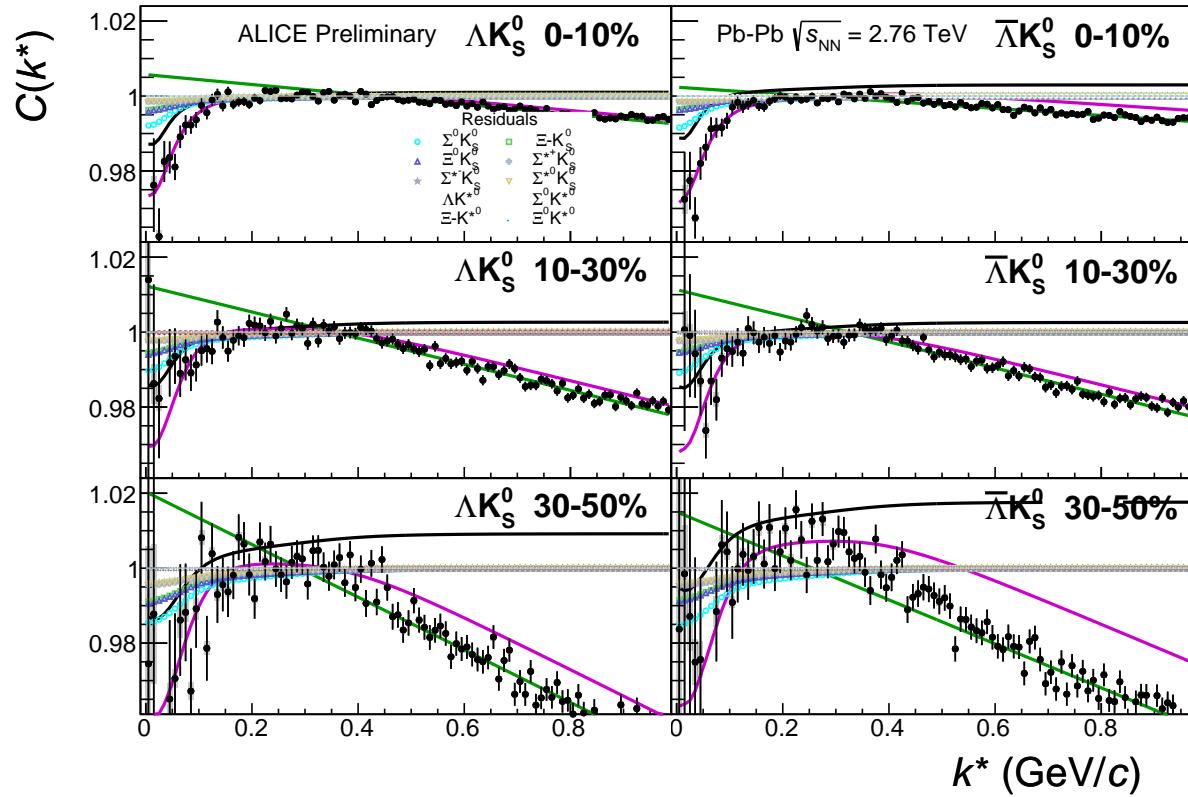
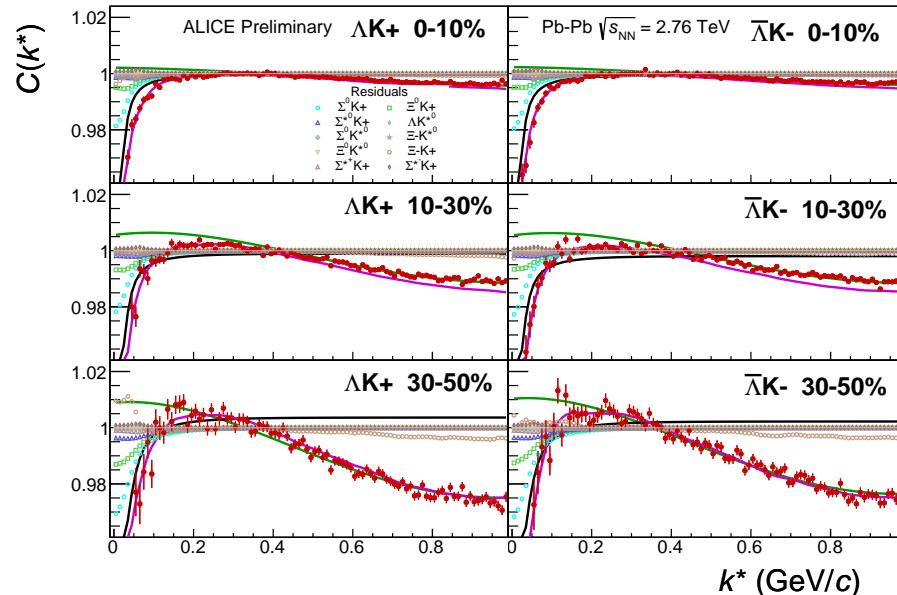
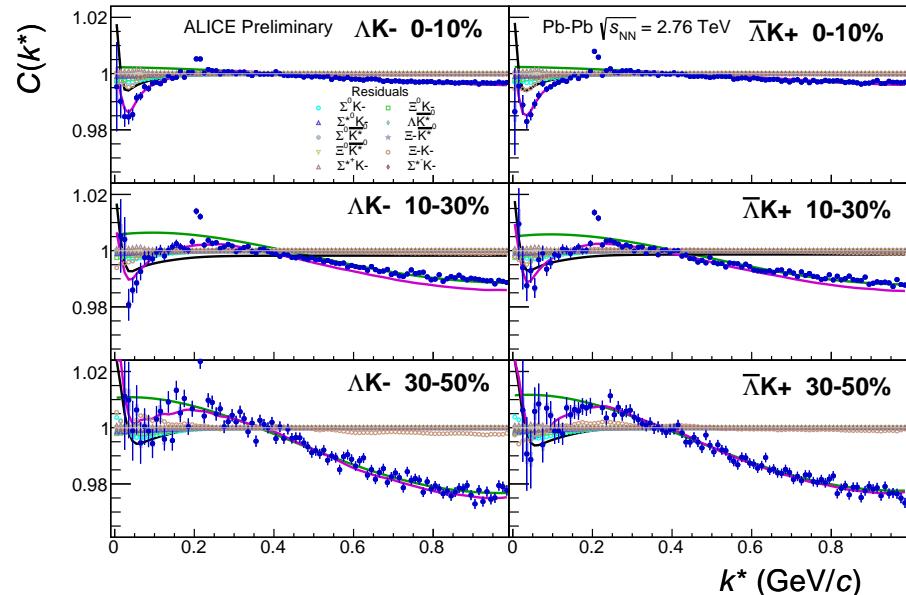


Fig. 43: Fits, with 10 residual correlations included and shown, to the ΛK_s^0 (left) and $\bar{\Lambda} K_s^0$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The ten parent pairs used for the residual correction to the ΛK_s^0 ($\bar{\Lambda} K_s^0$) fit are $\Sigma^0 K_S^0$, $\Xi^0 K_S^0$, $\Xi^- K_S^0$, $\Sigma^{*(+,-,0)} K_S^0$, ΛK^{*0} , $\Sigma^0 K^{*0}$, $\Xi^0 K^{*0}$, and $\Xi^- K^{*0}$ ($\bar{\Sigma}^0 K_S^0$, $\bar{\Xi}^0 K_S^0$, $\Xi^+ K_S^0$, $\Sigma^{*(+,-,0)} \bar{K}_S^0$, $\bar{\Lambda} K^{*0}$, $\bar{\Sigma}^0 \bar{K}^{*0}$, $\bar{\Xi}^+ \bar{K}^{*0}$).



(a) $\Lambda K^+(\bar{\Lambda} K^-)$ fits with residual contributions shown for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom)



(b) $\Lambda K^- (\bar{\Lambda} K^+)$ fits with residual contributions shown for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom)

Fig. 44: Fits, with 10 residual correlations included and shown, to the ΛK^+ & $\bar{\Lambda} K^-$ (left) and ΛK^- & $\bar{\Lambda} K^+$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The ten parent pairs used for the residual correction to the ΛK^+ ($\bar{\Lambda} K^-$) fit are $\Sigma^0 K^+$, $\Xi^0 K^+$, $\Xi^- K^+$, $\Sigma^{*(+,-,0)} K^+$, ΛK^{*0} , $\Sigma^0 K^{*0}$, $\Xi^0 K^{*0}$, and $\Xi^- K^{*0}$ ($\bar{\Sigma}^0 K^-$, $\bar{\Xi}^0 K^-$, $\bar{\Xi}^+ K^-$, $\bar{\Sigma}^{*(+,-,0)} K^-$, $\bar{\Lambda} K^{*0}$, $\bar{\Sigma}^0 \bar{K}^{*0}$, $\bar{\Xi}^0 \bar{K}^{*0}$, and $\bar{\Xi}^- \bar{K}^{*0}$).

| Fit Results $\Lambda(\bar{\Lambda})K_S^0$ | | | | | | |
|---|------------|---|---|--|---|---|
| System | Centrality | Fit Parameters | | | | |
| | | λ | R | $\mathbb{R}f_0$ | $\mathbb{I}f_0$ | d_0 |
| $\Lambda K_S^0 \& \bar{\Lambda} K_S^0$ | 0-10% | | | 2.96 ± 0.54 (stat.) ± 0.33 (sys.) | | |
| | 10-30% | 0.60 ± 0.71 (stat.) ± 0.54 (sys.) | 2.40 ± 0.45 (stat.) ± 0.29 (sys.) | -0.35 ± 0.10 (stat.) ± 0.21 (sys.) | 0.19 ± 0.11 (stat.) ± 0.12 (sys.) | 1.85 ± 0.61 (stat.) ± 2.68 (sys.) |
| | 30-50% | | | 1.77 ± 0.32 (stat.) ± 0.15 (sys.) | | |

Table 10: Fit Results $\Lambda(\bar{\Lambda})K_S^0$, with 10 residual correlations included. Each pair is fit simultaneously with its conjugate (ie. ΛK_S^0 with $\bar{\Lambda} K_S^0$) across all centralities (0-10%, 10-30%, 30-50%), for a total of 6 simultaneous analyses in the fit. A single λ parameter is shared amongst all. Each analysis has a unique normalization parameter. The radii are shared between analyses of like centrality, as these should have similar source sizes. The scattering parameters ($\mathbb{R}f_0$, $\mathbb{I}f_0$, d_0) are shared amongst all. The background is fit with a linear form in the range $0.6 < k^* < 0.9$ GeV/c. The fit is done on the data with only statistical error bars. The errors marked as “stat.” are those returned by MINUIT. The errors marked as “sys.” are those which result from my systematic analysis (as outlined in Section 6).

| Fit Results $\Lambda(\bar{\Lambda})K^\pm$ | | | | | | |
|---|------------|---|---|--|---|--|
| System | Centrality | Fit Parameters | | | | |
| | | λ | R | $\mathbb{R}f_0$ | $\mathbb{I}f_0$ | d_0 |
| $\Lambda K^+ \& \bar{\Lambda} K^-$ | 0-10% | 1.87 ± 0.42 (stat.) ± 0.21 (sys.) | 6.59 ± 0.80 (stat.) ± 0.49 (sys.) | | | |
| | 10-30% | 1.27 ± 0.26 (stat.) ± 0.23 (sys.) | 4.91 ± 0.53 (stat.) ± 0.28 (sys.) | -1.14 ± 0.17 (stat.) ± 0.31 (sys.) | 0.66 ± 0.16 (stat.) ± 0.15 (sys.) | 0.68 ± 0.46 (stat.) ± 0.53 (sys.) |
| $\Lambda K^+ \& \bar{\Lambda} K^-$ | 30-50% | 1.07 ± 0.23 (stat.) ± 0.32 (sys.) | 3.44 ± 0.36 (stat.) ± 0.13 (sys.) | 0.52 ± 0.15 (stat.) ± 0.19 (sys.) | 0.55 ± 0.10 (stat.) ± 0.18 (sys.) | -3.61 ± 1.20 (stat.) ± 1.02 (sys.) |

Table 11: Fit Results $\Lambda(\bar{\Lambda})K^\pm$, with 10 residual correlations included. All ΛK^\pm analyses are fit simultaneously across all centralities (0-10%, 10-30%, 30-50%). Scattering parameters ($\mathbb{R}f_0$, $\mathbb{I}f_0$, d_0) are shared between pair-conjugate systems (i.e. a parameter set describing the ΛK^+ & $\bar{\Lambda} K^-$ system, and a separate set describing the ΛK^- & $\bar{\Lambda} K^+$ system). For each centrality, a radius and λ parameters are shared between all pairs (ΛK^+ , $\bar{\Lambda} K^-$, ΛK^- , $\bar{\Lambda} K^+$). Each analysis has a unique normalization parameter. The background is modeled by a (6th-)degree polynomial fit to THERMINATOR simulation. The fit is done on the data with only statistical error bars. The errors marked as “stat.” are those returned by MINUIT. The errors marked as “sys.” are those which result from my systematic analysis (as outlined in Section 6).

807 **7.1.3 No Residual Correlations Included in Fit**

808 Figure 45 nicely collects and summarizes all of our extracted fit parameters for the case of no included
 809 residual contributors. Figure 46 presents our extracted fit radii, along with those of other systems pre-
 810 viously analyzed by ALICE [11], as a function of pair transverse mass (m_T). Figures 47, 48, and 49
 811 show the experimental correlation functions with fits, assuming no residual contributors, for all studied
 812 centralities for ΛK_S^0 with $\bar{\Lambda} K_S^0$, ΛK^+ with $\bar{\Lambda} K^-$, and ΛK^- with $\bar{\Lambda} K^+$, respectively. The parameter sets
 813 extracted from the fits can be found in Tables 12 and 13.

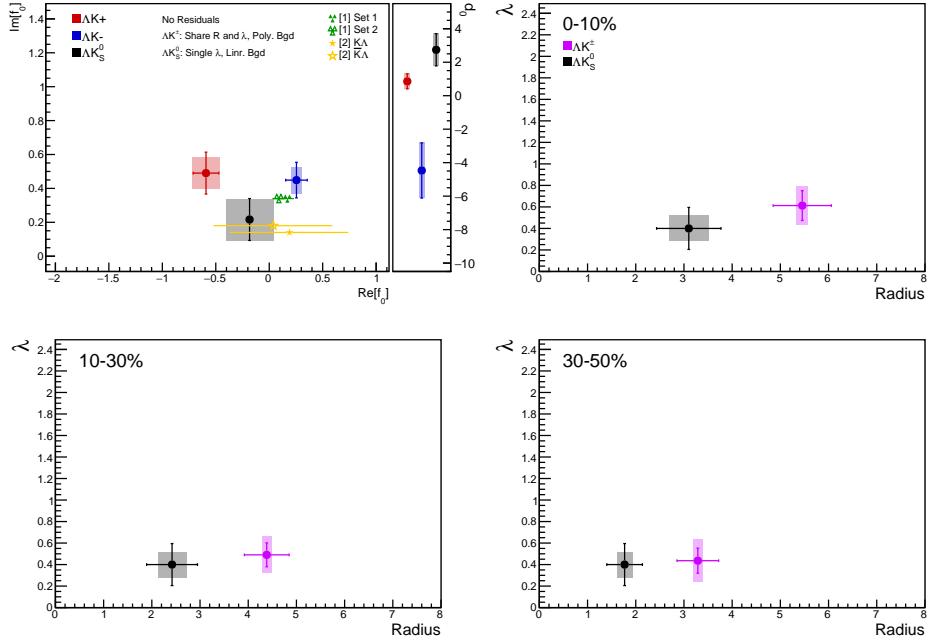


Fig. 45: Extracted scattering parameters for the case of NO residual contributors for all of our ΛK systems. [Top Left]: $\text{Im}[f_0]$ vs. $\text{Re}[f_0]$, together with d_0 to the right. [Top Right (Bottom Left, Bottom Right)]: λ vs. Radius for the 0-10% (10-30%, 30-50%) bin. The green [9] and yellow [10] points show theoretical predictions made using chiral perturbation theory.

814 Figure 46 shows extracted R_{inv} parameters as a function of tranverse mass (m_T) for various pair systems
 815 over several centralities. The published ALICE data [11] is shown with transparent, open symbols. The
 816 new ΛK results are shown with opaque, filled symbols. The radii shown an increasing size with increas-
 817 ing centrality, as is expected from the simple geometric picture of the collisions. The radii decrease
 818 in size with increasing m_T , and we see an approximate scaling of the radii with transverse mass, as is
 819 expected in the presence of collective flow in the system.

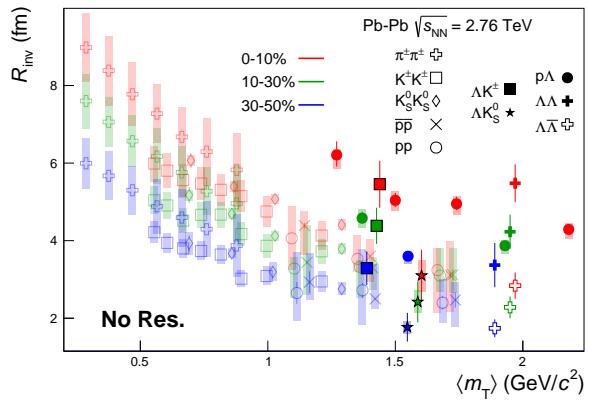
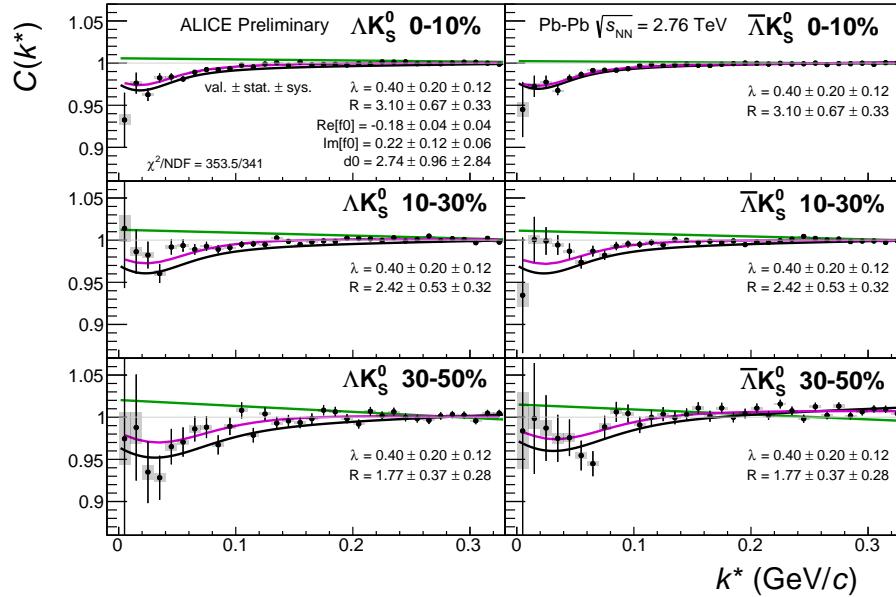
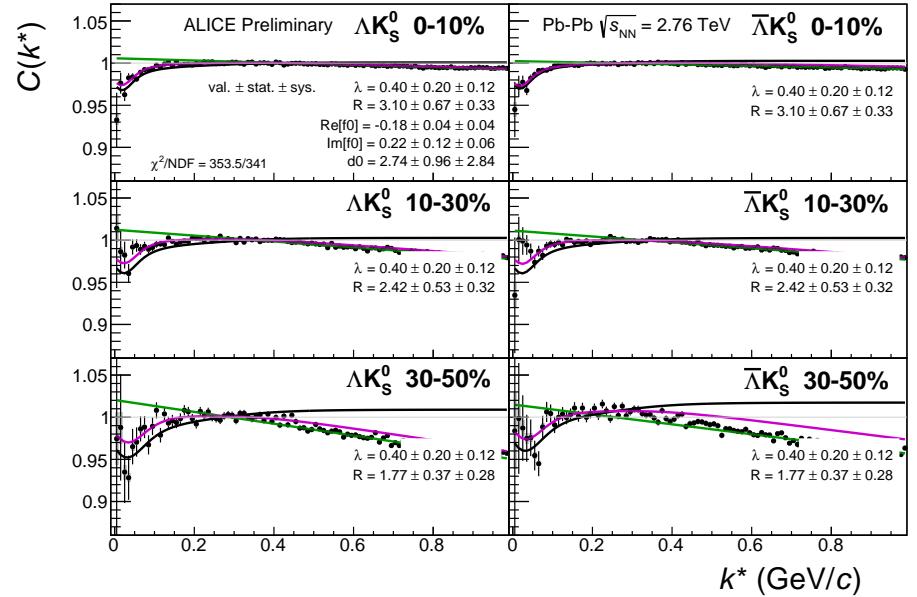


Fig. 46: No residual correlations in ΛK fits. Extracted fit R_{inv} parameters as a function of pair transverse mass (m_T) for various pair systems over several centralities. The ALICE published data [11] is shown with transparent, open symbols. The new ΛK results are shown with opaque, filled symbols. In the left, the ΛK^+ (with its conjugate pair) results are shown separately from the ΛK^- (with its conjugate pair) results. In the right, all ΛK^\pm results are averaged.

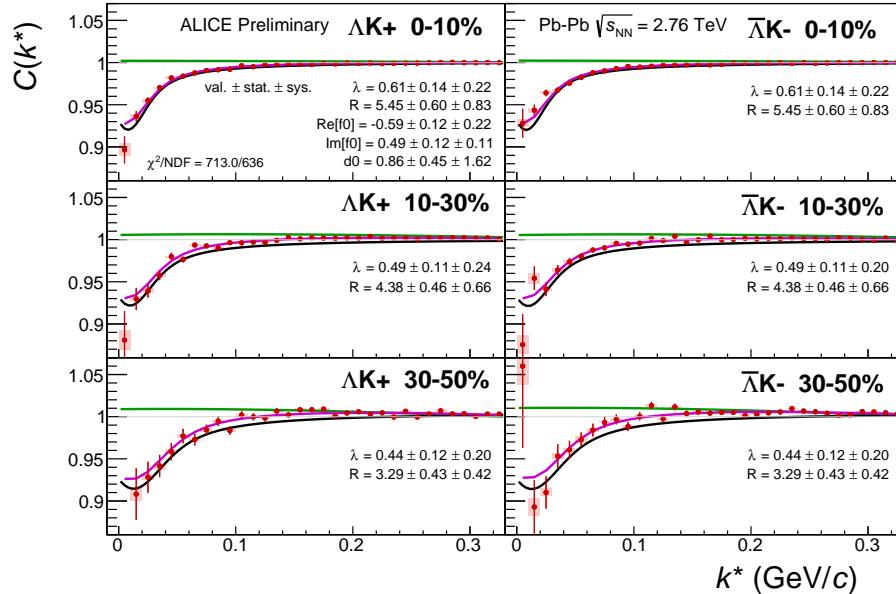


(a) Signal region view ($k^* \lesssim 0.3 \text{ GeV}/c$)

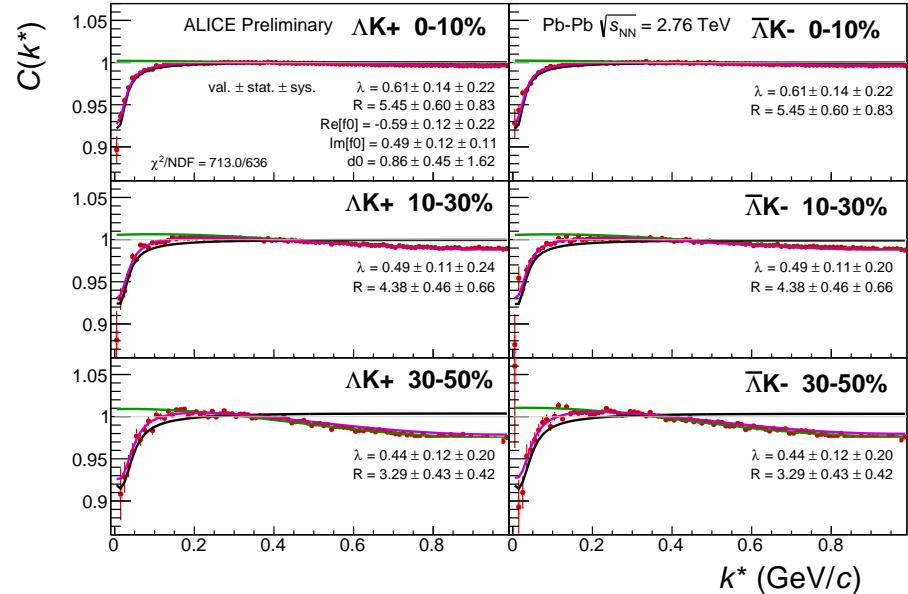


(b) Wide view ($k^* \lesssim 1.0 \text{ GeV}/c$)

Fig. 47: Fits, with NO residual correlations included, to the ΛK_s^0 (left) and $\bar{\Lambda} K_s^0$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The lines represent the statistical errors, while the boxes represent the systematic errors. A single λ parameter is shared amongst all. Each analysis has a unique normalization parameter. The radii are shared between analyses of like centrality, as these should have similar source sizes. The scattering parameters ($\mathbb{R} f_0$, $\mathbb{I} f_0$, d_0) are shared amongst all. The background is modeled by a (6th-degree polynomial fit to THERMINATOR simulation. The black solid line represents the “raw” primary fit, i.e. not corrected for momentum resolution effects nor non-flat background. The green line shows the fit to the non-flat background. The purple points show the fit after momentum resolution and non-flat background corrections have been applied. The extracted fit values with uncertainties are printed.

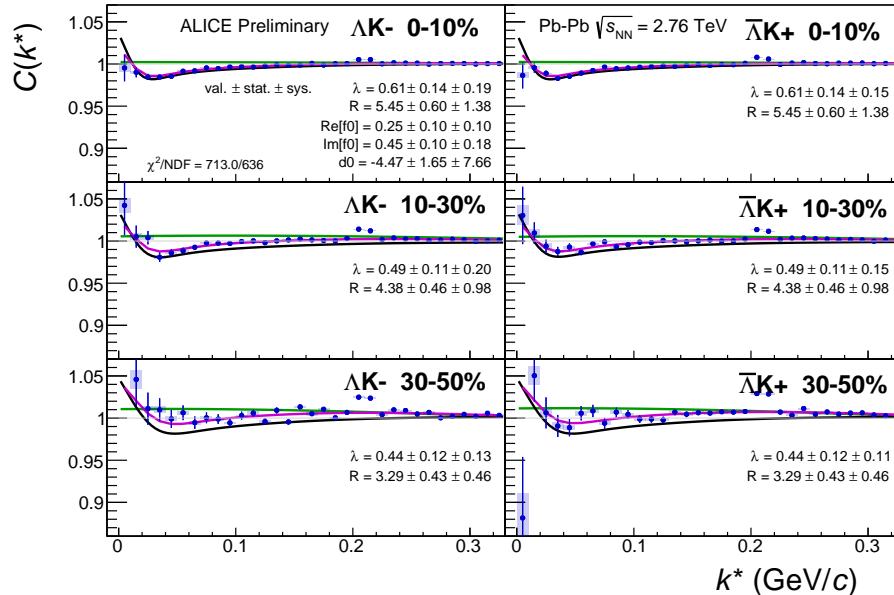


(a) Signal region view ($k^* \lesssim 0.3 \text{ GeV}/c$)

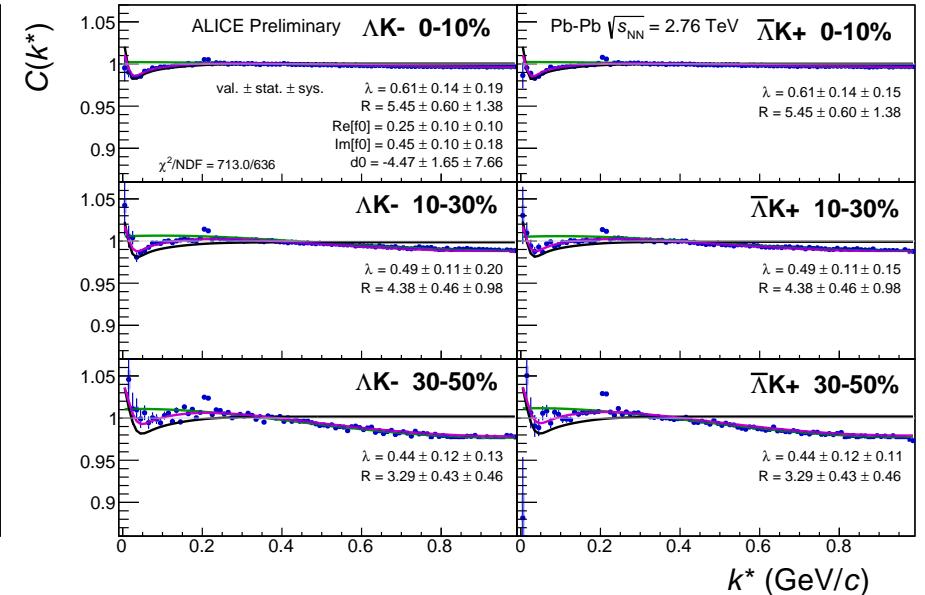


(b) Wide view ($k^* \lesssim 1.0 \text{ GeV}/c$)

Fig. 48: Fits to the ΛK^+ (left) and $\bar{\Lambda} K^-$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The lines represent the statistical errors, while the boxes represent the systematic errors. All ΛK^\pm analyses are fit simultaneously across all centralities (0-10%, 10-30%, 30-50%). Scattering parameters ($\mathbb{R}f_0$, $\mathbb{I}f_0$, d_0) are shared between pair-conjugate systems (i.e. a parameter set describing the ΛK^+ & $\bar{\Lambda} K^-$ system, and a separate set describing the ΛK^- & $\bar{\Lambda} K^+$ system). For each centrality, a radius and λ parameters are shared between all pairs (ΛK^+ , $\bar{\Lambda} K^-$, ΛK^- , $\bar{\Lambda} K^+$). Each analysis has a unique normalization parameter. The background is modeled by a (6th-)degree polynomial fit to THERMINATOR simulation. The black solid line represents the “raw” primary fit, i.e. not corrected for momentum resolution effects nor non-flat background. The green line shows the fit to the non-flat background. The purple points show the fit after momentum resolution and non-flat background corrections have been applied. The extracted fit values with uncertainties are printed.



(a) Signal region view ($k^* \lesssim 0.3 \text{ GeV}/c$)



(b) Wide view ($k^* \lesssim 1.0 \text{ GeV}/c$)

Fig. 49: Fits, with NO residual correlations included, to the ΛK^- (left) with $\bar{\Lambda} K^+$ (right) data for the centralities 0-10% (top), 10-30% (middle), and 30-50% (bottom). The lines represent the statistical errors, while the boxes represent the systematic errors. All ΛK^\pm analyses are fit simultaneously across all centralities (0-10%, 10-30%, 30-50%). Scattering parameters ($\Re f_0$, $\Im f_0$, d_0) are shared between pair-conjugate systems (i.e. a parameter set describing the ΛK^+ & $\bar{\Lambda} K^-$ system, and a separate set describing the ΛK^- & $\bar{\Lambda} K^+$ system). For each centrality, a radius and λ parameters are shared between all pairs (ΛK^+ , $\bar{\Lambda} K^-$, ΛK^- , $\bar{\Lambda} K^+$). Each analysis has a unique normalization parameter. The background is modeled by a (6th-degree polynomial fit to THERMINATOR simulation. The black solid line represents the “raw” primary fit, i.e. not corrected for momentum resolution effects nor non-flat background. The green line shows the fit to the non-flat background. The purple points show the fit after momentum resolution and non-flat background corrections have been applied. The extracted fit values with uncertainties are printed.

| Fit Results $\Lambda(\bar{\Lambda})K_S^0$ | | | | | | |
|---|------------|---|---|--|---|---|
| System | Centrality | Fit Parameters | | | | |
| | | λ | R | $\mathbb{R}f_0$ | $\mathbb{I}f_0$ | d_0 |
| $\Lambda K_S^0 \& \bar{\Lambda} K_S^0$ | 0-10% | | | 3.10 ± 0.67 (stat.) ± 0.41 (sys.) | | |
| | 10-30% | 0.40 ± 0.20 (stat.) ± 0.12 (sys.) | 2.42 ± 0.53 (stat.) ± 0.29 (sys.) | -0.18 ± 0.04 (stat.) ± 0.22 (sys.) | 0.22 ± 0.12 (stat.) ± 0.12 (sys.) | 2.74 ± 0.96 (stat.) ± 1.28 (sys.) |
| | 30-50% | | | 1.77 ± 0.37 (stat.) ± 0.16 (sys.) | | |

Table 12: Fit Results $\Lambda(\bar{\Lambda})K_S^0$, with no residual correlations included. Each pair is fit simultaneously with its conjugate (ie. ΛK_S^0 with $\bar{\Lambda} K_S^0$) across all centralities (0-10%, 10-30%, 30-50%), for a total of 6 simultaneous analyses in the fit. A single λ parameter is shared amongst all. Each analysis has a unique normalization parameter. The radii are shared between analyses of like centrality, as these should have similar source sizes. The scattering parameters ($\mathbb{R}f_0$, $\mathbb{I}f_0$, d_0) are shared amongst all. The background is fit with a linear form in the range $0.6 < k^* < 0.9$ GeV/c. The fit is done on the data with only statistical error bars. The errors marked as “stat.” are those returned by MINUIT. The errors marked as “sys.” are those which result from my systematic analysis (as outlined in Section 6).

| Fit Results $\Lambda(\bar{\Lambda})K^\pm$ | | | | | | |
|---|------------|---|---|--|---|--|
| System | Centrality | Fit Parameters | | | | |
| | | λ | R | $\mathbb{R}f_0$ | $\mathbb{I}f_0$ | d_0 |
| $\Lambda K^+ \& \bar{\Lambda} K^-$ | 0-10% | 0.61 ± 0.14 (stat.) ± 0.18 (sys.) | 5.45 ± 0.60 (stat.) ± 0.12 (sys.) | -0.59 ± 0.12 (stat.) ± 0.13 (sys.) | 0.49 ± 0.12 (stat.) ± 0.09 (sys.) | 0.86 ± 0.45 (stat.) ± 1.63 (sys.) |
| | 10-30% | 0.49 ± 0.11 (stat.) ± 0.17 (sys.) | 4.38 ± 0.46 (stat.) ± 0.10 (sys.) | | | |
| $\Lambda K^+ \& \bar{\Lambda} K^+$ | 30-50% | 0.44 ± 0.12 (stat.) ± 0.20 (sys.) | 3.29 ± 0.43 (stat.) ± 0.10 (sys.) | 0.25 ± 0.10 (stat.) ± 0.05 (sys.) | 0.45 ± 0.10 (stat.) ± 0.08 (sys.) | -4.47 ± 1.65 (stat.) ± 1.60 (sys.) |

Table 13: Fit Results $\Lambda(\bar{\Lambda})K^\pm$, with no residual correlations included. All ΛK^\pm analyses are fit simultaneously across all centralities (0-10%, 10-30%, 30-50%). Scattering parameters ($\mathbb{R}f_0$, $\mathbb{I}f_0$, d_0) are shared between pair-conjugate systems (i.e. a parameter set describing the $\Lambda K^+ \& \bar{\Lambda} K^-$ system, and a separate set describing the $\Lambda K^- \& \bar{\Lambda} K^+$ system). For each centrality, a radius and λ parameters are shared between all pairs ($\Lambda K^+, \bar{\Lambda} K^-, \Lambda K^-, \bar{\Lambda} K^+$). Each analysis has a unique normalization parameter. The background is modeled by a (6th-)degree polynomial fit to THERMINATOR simulation. The fit is done on the data with only statistical error bars. The errors marked as “stat.” are those returned by MINUIT. The errors marked as “sys.” are those which result from my systematic analysis (as outlined in Section 6).

820 7.1.4 Fit Method Comparisons

821 In Figure 50, we show extracted fit parameters for the case of $\Lambda K^+(\bar{\Lambda} K^-)$ sharing radii with $\Lambda K^-(\bar{\Lambda} K^+)$.
 822 The figure shows results for three different treatments of the non-femtoscopic background: a polynomial
 823 fit to THERMINATOR 2 simulation to model the background (circles), a linear fit to the data to model
 824 the background (squares), and the Stavinsky method (crosses).

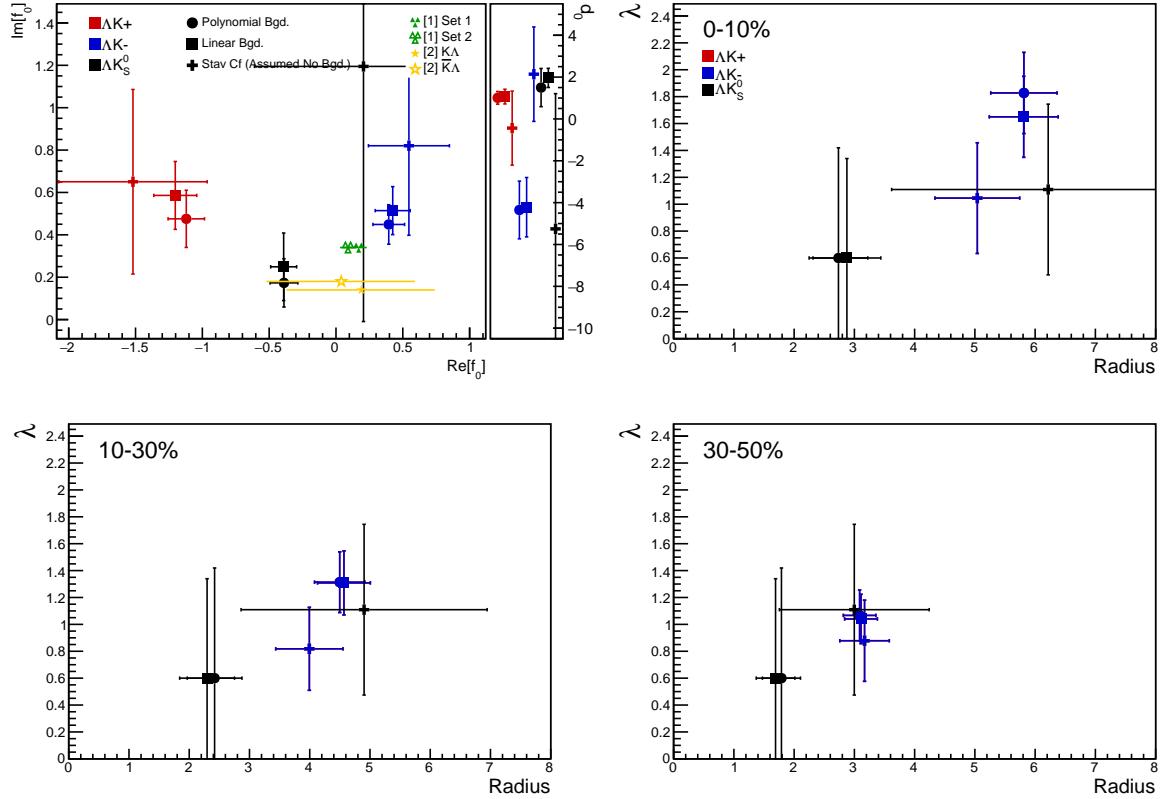


Fig. 50: Compare Fit Parameters: Background treatment: Extracted fit results for all of our $\Lambda(\bar{\Lambda})K^\pm$ systems across all studied centrality bins (0-10%, 10-30%, 30-50%). The $\Lambda K^+(\bar{\Lambda} K^-)$ and $\Lambda K^-(\bar{\Lambda} K^+)$ systems share both a radius and a λ parameter for each centrality bin (i.e. 3 total radius parameters, 3 total λ parameters). The figure shows results for three different treatments of the non-femtoscopic background: a polynomial fit to THERMINATOR 2 simulation to model the background (circles), a linear fit to the data to model the background (squares), and the Stavinsky method (crosses). The green [9] and yellow [10] points show theoretical predictions made using chiral perturbation theory.

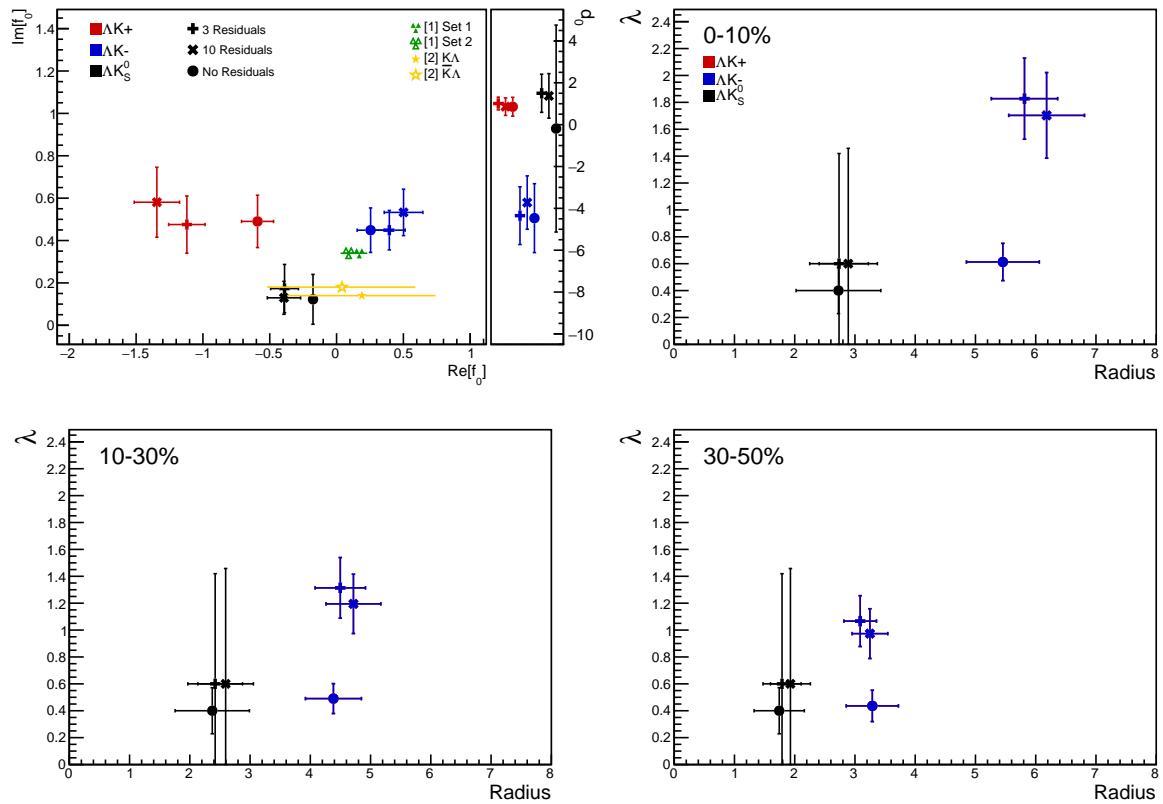


Fig. 51: Compare Fit Parameters: Number of residuals: Results shown for the case of 3 (+), 10 (X), and no (circles) residual contributors.

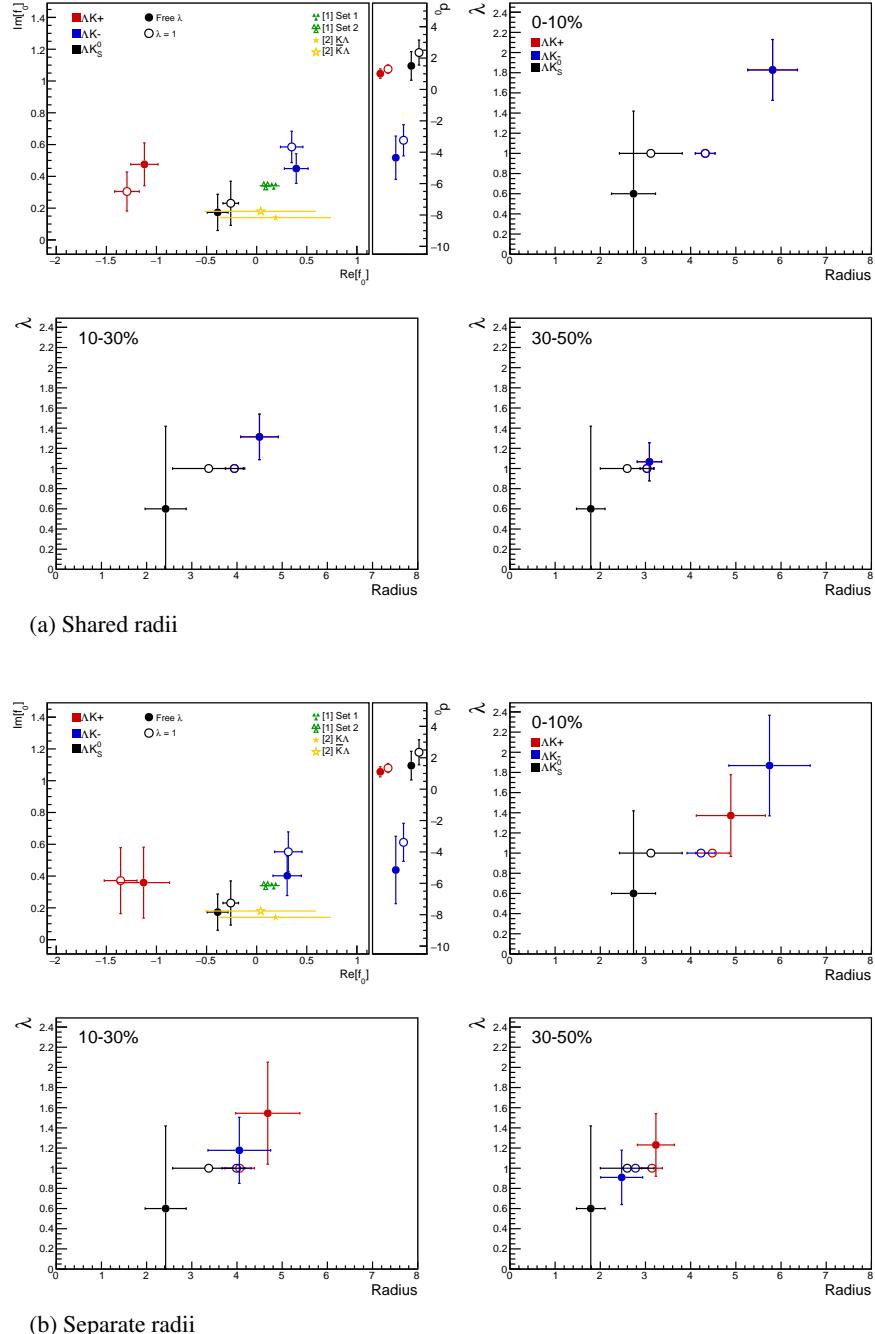
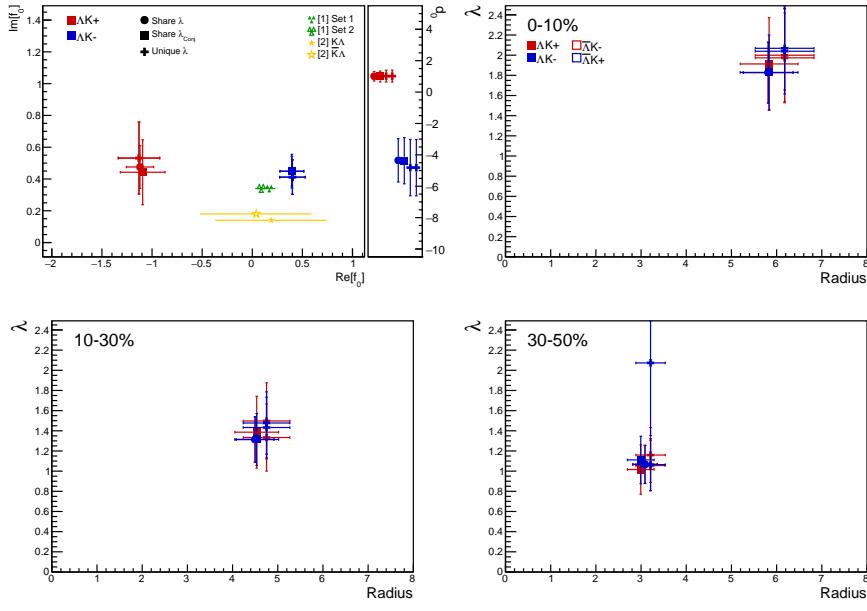
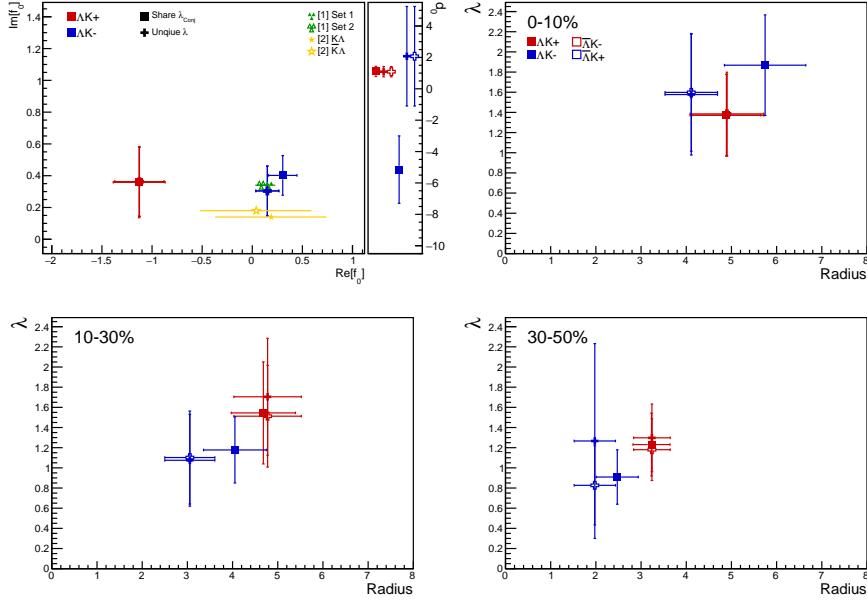


Fig. 52: Compare Fit Parameters: Free vs fixed λ : Results shown for λ parameters left free (filled symbols) and fixed to 1 (open symbols). In the top plot (52a), the ΛK^+ and ΛK^- analyses share radii, whereas in the bottom (52b) they have unique radii.



(a) Shared radii



(b) Separate radii

Fig. 53: Compare Fit Parameters: Shared vs unique λ : Results shown for different sharing of the λ parameters between analyses and systems. In the top (53a), the ΛK^+ and $\bar{\Lambda} K^-$ analyses share radii, whereas in the bottom (53b), they do not. “Share λ ” (circles) is the case where a single λ is shared amongst all analyses for a given centrality bin (i.e., in 53a, 3 radius parameters and 3 λ parameters). “Share λ_{Conj} ” (squares) means that conjugate pairs (ex. ΛK^+ and $\bar{\Lambda} K^-$) share a λ parameter for each centrality. This corresponds to 6 total λ parameters (for each of the 3 centrality bins, the $\Lambda K^+(\bar{\Lambda} K^-)$ receives a unique λ , as does $\Lambda K^-(\bar{\Lambda} K^+)$). Finally, in “Unique λ ” (+), each analysis received its own unique λ parameter. This corresponds to 12 λ parameters (for each of the 3 centrality bins, each ΛK^+ , $\bar{\Lambda} K^-$, ΛK^- , and $\bar{\Lambda} K^+$ receives a unique λ).

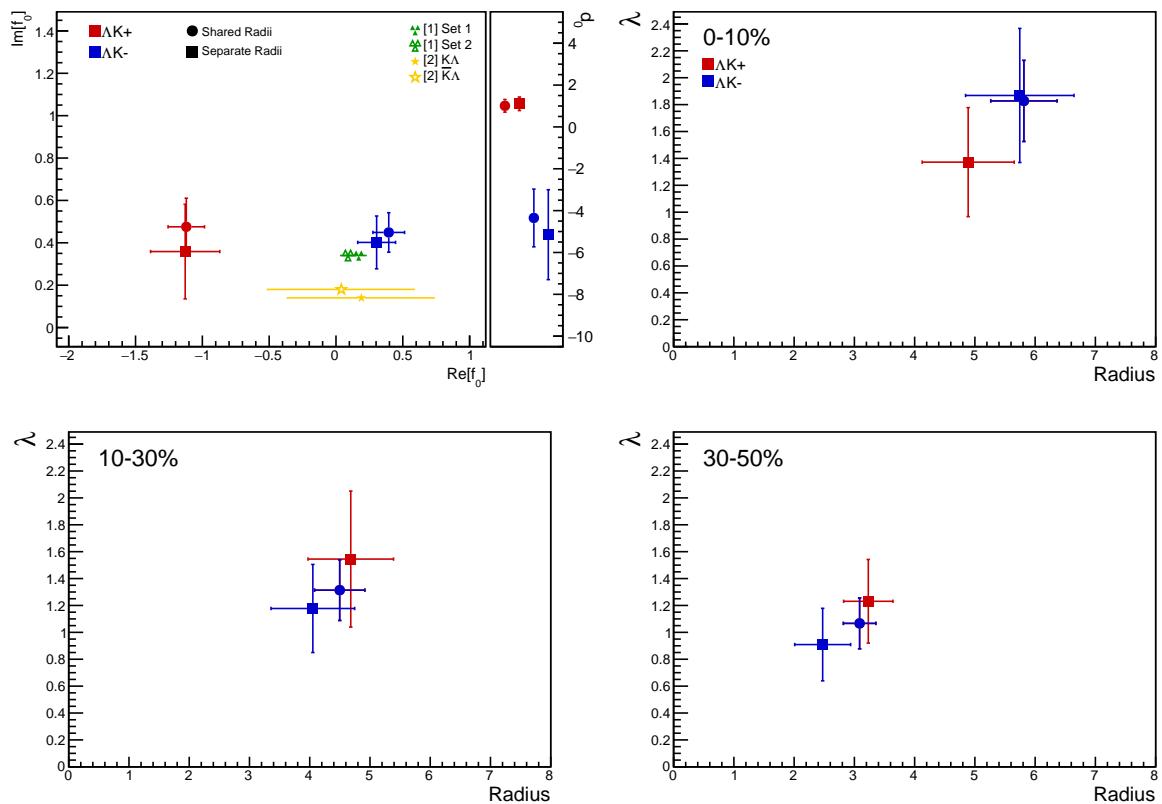


Fig. 54: Compare Fit Parameters: Shared vs. Separate Radii: Results shown for the case of radii being shared between $\Lambda K^+(\bar{\Lambda} K^-)$ and $\Lambda K^-(\bar{\Lambda} K^+)$ (circles) vs not shared (squares).

825 **7.1.5 Discussion of m_T -Scaling**

826 It is clear from the results presented in the previous sections, that the ΛK systems do not conform to the
 827 approximate m_T -scaling of the pair source sizes. At first thought, this may appear to be a troubling result;
 828 the approximate scaling is an observed consequence of the collective behavior of the soft ($low-p_T$) sector
 829 of the produced system. The Λ and K particles certainly participate in the collective expansion of the
 830 QGP medium, so why do their extracted femtoscopic radii not behave as expected? To get straight to the
 831 point: the ΛK systems are (obviously) comprised on non-identical particles, each with its own and unique
 832 single particle source. Each source is, in general, unique in both its overall size, and in its space-time po-
 833 sition within the produced medium. The hydrodynamic nature of the medium produces the approximate
 834 m_T -scaling with respect to these single-particle sources, not the pair sources. The combination of these
 835 effects, when probing correlations between non-identical particle pairs, leads to extracted radii falling
 836 outside of the (identical particle femtoscopy) m_T -scaling trend. Figure 55 (which contains the same data
 837 as Fig.32), shows again the R_{inv} vs m_T plot, but also highlights (with arrows) the approximate individual
 838 $\langle m_T \rangle$ values of the single particle distributions. The grey circles show how to single particle sizes change
 839 with m_T .

840 Taking a close look at Fig. 55, one can see that the previously published data (transparent points) are
 841 for identical particle analyses only. For these cases, the pair source, probed through femtoscopy, is
 842 comprised of two identical sources laying on top of each other. The extracted femtoscopic radii are
 843 related to the single particle source sizes by a factor of $\sqrt{2}$, and of course follow the m_T -scaling trend.
 844 The other (unpublished) non-identical particle femtoscopy study ($p\Lambda$) included in the figure, also shows
 845 radii deviating from the m_T -scaling band. Drawing a comparison with the $\Lambda\bar{\Lambda}$ study shown in Fig. 32
 846 is a bit more complicated; the $\Lambda\bar{\Lambda}$ system, although containing non-identical particles, does contain a
 847 particle with its antiparticle, for which annihilation could conceivably alter the pair source distribution.
 848 It would be more surprising if the non-identical analyses did happen to conform to the scaling; although,
 849 this could occur for a non-identical analysis in which the particles have similar masses as well as similar
 850 m_T distributions. For the case presented here, the result differing from m_T -scaling is not surprising.

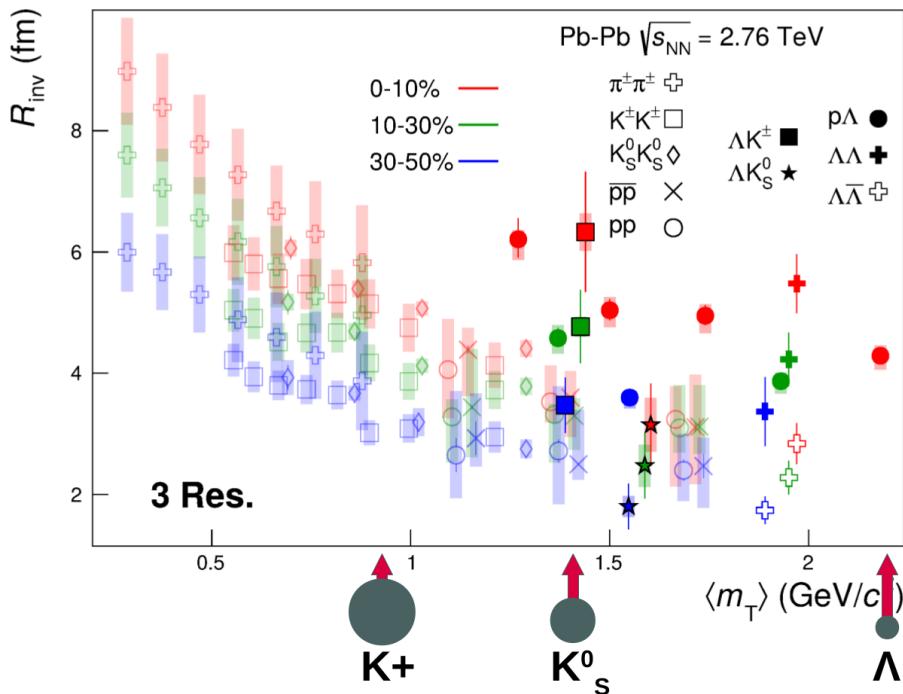


Fig. 55: Same as Fig. 32, but with the individual m_T values for the single particle distributions identified. The grey circles show how the single particle sizes are expected to change with m_T .

851 I will also briefly point out that it is not automatically clear where a non-identical study should be placed
 852 on such a R_{inv} vs m_T plot. Each single particle distribution has a well-defined $\langle m_T \rangle$, which, to a large
 853 extent, determines the single particle region of homogeneity. When combining two sources with different
 854 spatio-temporal characteristics, originating from particles of different m_T , how should one define the pair
 855 m_T ? A simple mathematical expression for the pair m_T is easy to come up with, but that's not exactly
 856 what I'm hinting at here. With respect to this m_T -scaling picture, the m_T value dictates the source size,
 857 and one desires the same for non-identical particles. However, do the two unequal sized sources both
 858 contribute equally to the extracted femtoscopic size? Or does the larger (smaller) source more closely
 859 dictate the femtoscopic signal? If the contribution is equal, then it seems natural to simply more-or-less
 860 average the two, single particle, m_T values. If the contribution is unequal, then there should be introduced
 861 some sort of weighting in the pair m_T calculation reflecting this fact. In any case, in our study we use
 862 the most straightforward definition of pair m_T , defined as:

$$m_{T,\text{pair}}^2 = \left(\frac{m_{\text{inv}}}{2}\right)^2 + \left(\frac{1}{2}|\mathbf{p}_{T,1} + \mathbf{p}_{T,2}|\right)^2 \quad (24)$$

863 Many times, the equation for non-identical particle pair m_T is defined with the average mass replacing
 864 $m_{\text{inv}}/2$. However, the above Eq. 24 is more directly analogous to the single particle m_T :

$$m_T^2 = m^2 + \mathbf{p}_T^2 = (p^0)^2 - (p^3)^2 \quad (25)$$

865 as, Eq. 24 may be rewritten as:

$$\begin{aligned} m_{T,\text{pair}}^2 &= (K^0)^2 - (K^3)^2 \\ K^\mu &\equiv \frac{1}{2}(p_1^\mu + p_2^\mu) \end{aligned} \quad (26)$$

866 Identical particle femtoscopic studies are able to probe only the size of the emitting region, or, more
 867 precisely, the second moments of the emission function. In addition to this, non-identical particle studies
 868 are able to measure the relative emission shifts, the first moments of the emission function. One method
 869 to extract information about the emission asymmetries in the system is via a spherical decomposition
 870 of the correlation function. With this method, one can draw a wealth of information from just a few
 871 components of the decomposition. More specifically, the C_{00} component is similar to the 1D correlation
 872 functions typically studied, and probes the overall size of the source. The $\Im C_{11}$ component probes the
 873 asymmetry in the system; a non-zero value reveals the asymmetry.

874 In Fig. 56 we show results for the C_{00} and $\Im C_{11}$ components from the spherical decomposition of our
 875 ΛK^+ system in the 0-10% centrality bin (red circles). Results from a number of other components
 876 within the decomposition, as well as for our ΛK_S^0 and ΛK^- systems, are contained in 9.2. Along with
 877 the experimental data in Fig. 56, we have also included results from THERMINATOR simulation for an
 878 impact parameter of $b = 2$ fm (gold stars). As THERMINATOR does not include any final state effects,
 879 we assumed scattering parameters $(\Re f_0, \Im f_0, d_0) = (-1.16, 0.51, 1.08)$ and weighted the numerator pairs
 880 with $|\Psi|^2$, as discussed previously. As seen in the figure, the C_{00} signal is similar to that observed in our
 881 one-dimensional study. The $\Im C_{11}$ component shows a clear deviation from zero, and the negative value
 882 signifies that the Λ particles are, on average, emitted further out and/or earlier than the K mesons.

883 Fig. 57 shows a closer look at the THERMINATOR simulation, whose spherical harmonic decompositon
 884 was shown along with the data in Fig. 56. The top left of Fig. 57a shows a fit to the one-dimensional
 885 correlation function from THERMINATOR. The scattering parameters are known precisely here, as they
 886 served as the weights used in the simulation, and are kept constant in the fit. We are interested at looking

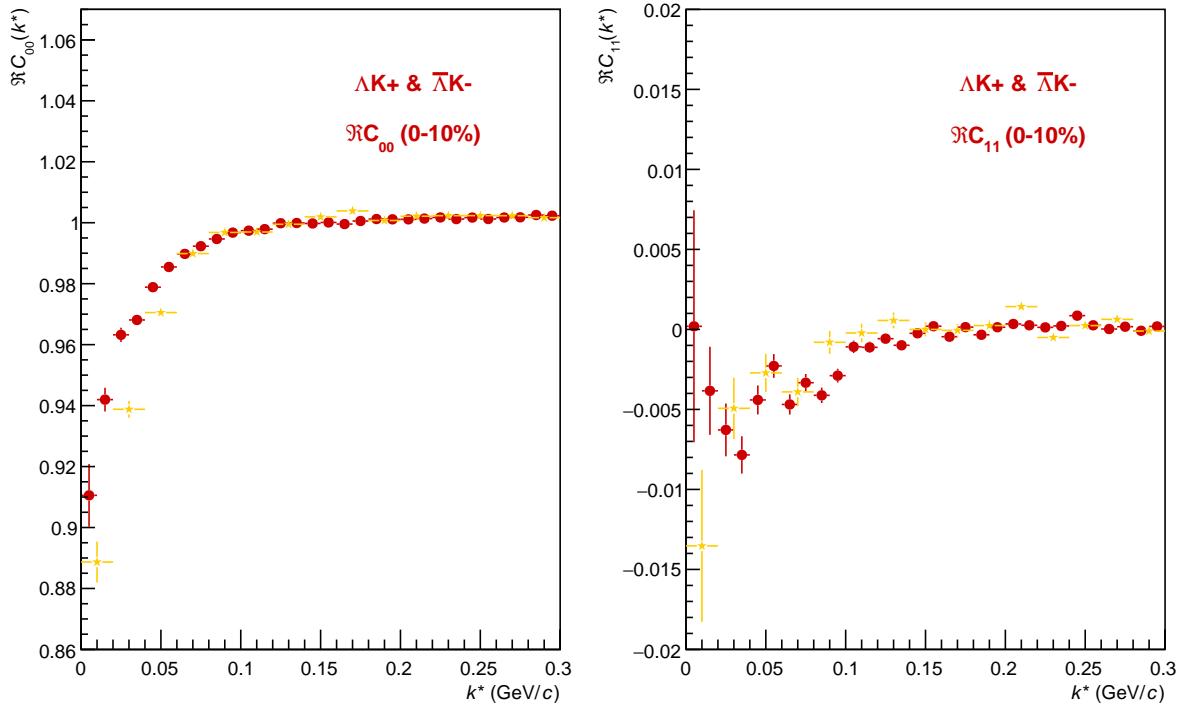
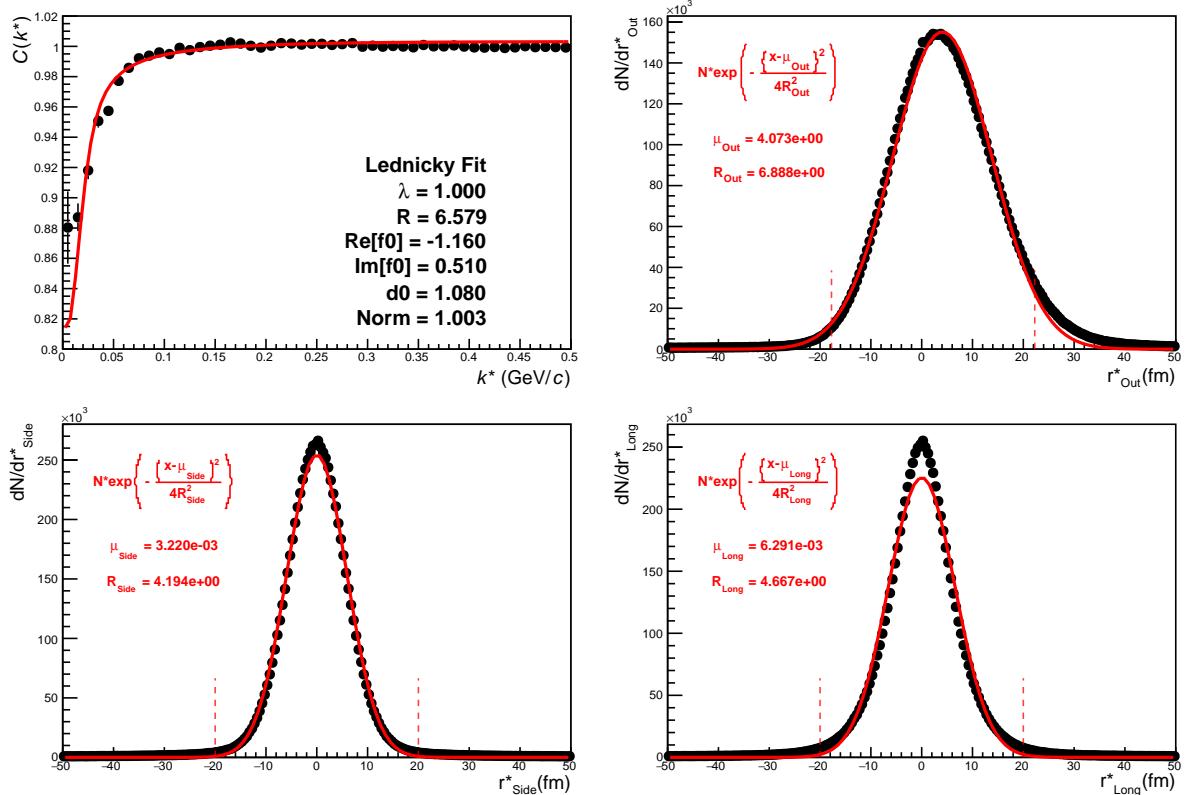


Fig. 56: C_{00} (left) and $\Re C_{11}$ (right) components of a spherical harmonic decomposition of the ΛK^+ correlation function for the 0-10% centrality bin. The C_{00} component is similar to the 1D correlation functions typically studied, and probes the overall size of the source. The $\Re C_{11}$ component probes the asymmetry in the system; a non-zero value reveals the asymmetry

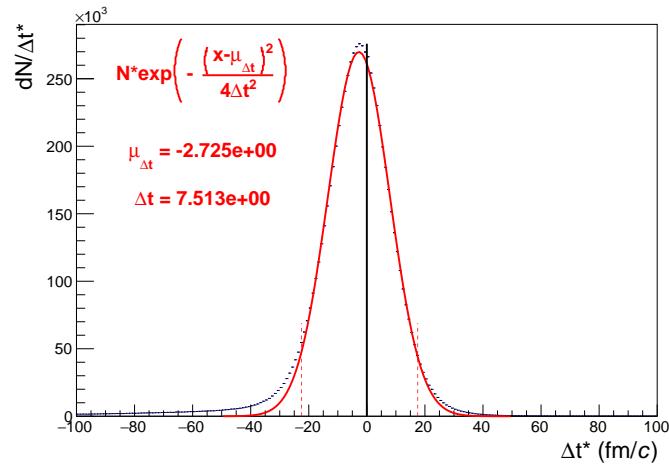
at the extracted one-dimensional source size here, so the λ parameter is also fixed at unity. The other three plots in Fig. 57a show the source distribution in the out (top right), side (bottom left), and long (bottom right) directions (all in the PRF). The source distributions have all been fitted with a Gaussian form, the result of which is printed within the respective plot. One immediately sees a significant shift in the out direction, $\mu_{\text{out}} \approx 4 \text{ fm}$, and negligible shift in the other two directions, $\mu_{\text{side}} \approx \mu_{\text{long}} \approx 0 \text{ fm}$. The figure demonstrates that, within the THERMINATOR model, the Λ is, on average, emitted further out than its K partner. Finally, Fig. 57b shows the distribution of the relative time of emittance, again in the PRF. The figure shows that the Λ is, on average, emitted earlier than its K partner.

We end this section with a brief look at how a spatial separation of the single particle sources affects the radii extracted from a femtoscopic analysis. To achieve this, we use THERMINATOR in a similar fashion as described above, but with one important difference. Instead of taking the source information from THERMINATOR, we instead draw the source from a pre-determined Gaussian distribution. In all cases, we take $R_{\text{out}} = R_{\text{side}} = R_{\text{long}} = 5 \text{ fm}$, and $\mu_{\text{side}} = \mu_{\text{long}} = 0 \text{ fm}$. Figure 58 shows an example of results obtained from THERMINATOR following this procedure, where $\mu_{\text{out}} = 3 \text{ fm}$.

In Figure 59, we show results for the case of $\mu_{\text{out}} = 1 \text{ fm}$, $\mu_{\text{out}} = 3 \text{ fm}$, and $\mu_{\text{out}} = 6 \text{ fm}$. In this figure, we do not show the side and long distributions, as they appear identical to those shown in Fig. 58. The figure demonstrates that as the separation μ_{out} increases, so do the extracted femtoscopic radii.



(a) Caption 1



(b) Caption 2

Fig. 57: Long Overall

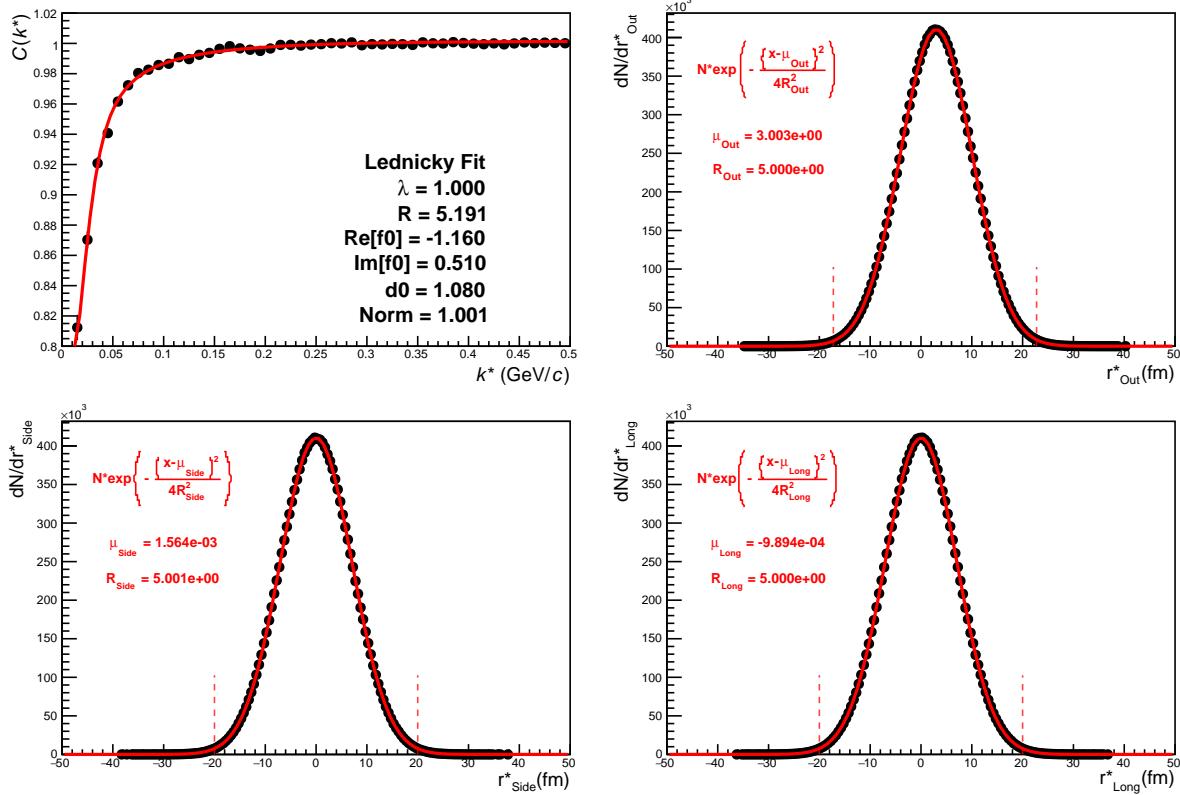


Fig. 58: Long Caption

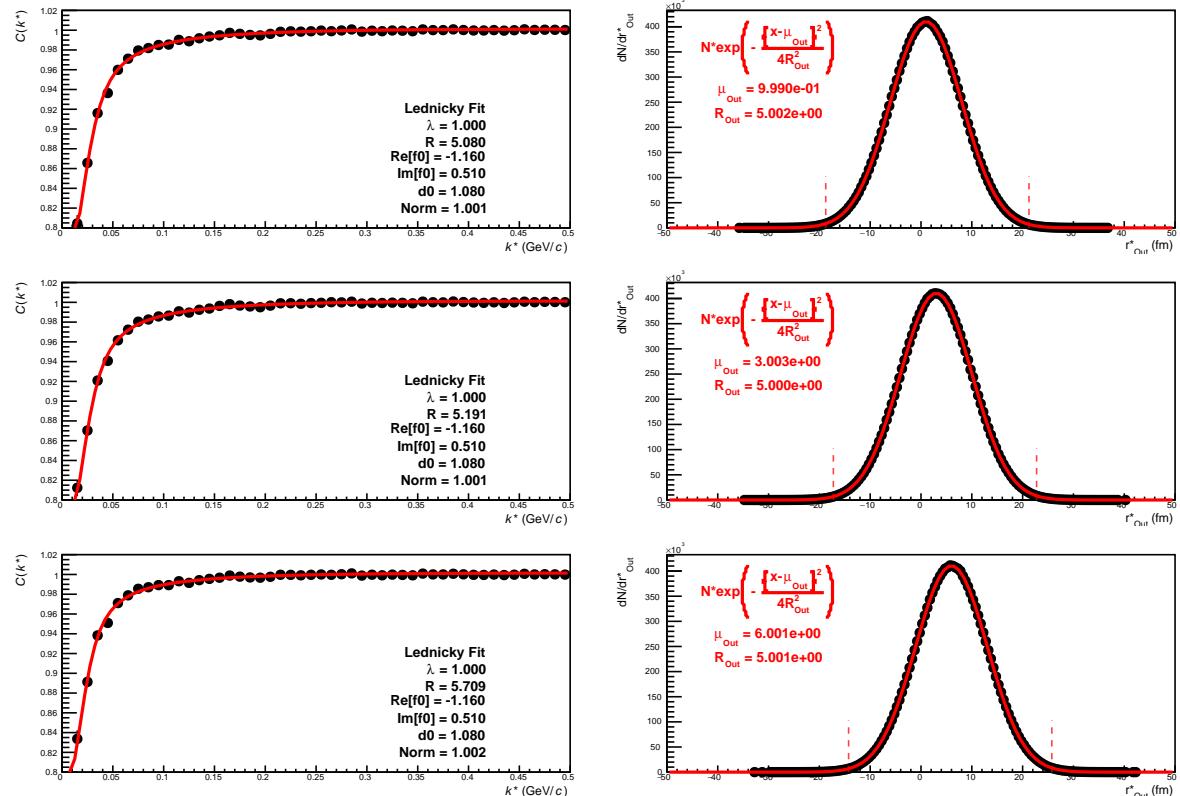


Fig. 59: Long Caption

904 **7.2 Results: ΞK^\pm**

905 Even without any fits to the data, the fact that the $\Xi^- K^+$ data dips below unity (Fig. 60) is exciting, as
 906 this cannot occur purely from a Coulomb interaction. We hope that this dip signifies that we are able to
 907 peer through the overwhelming contribution from the Coulomb interaction to see the effects arising from
 908 the strong interaction.

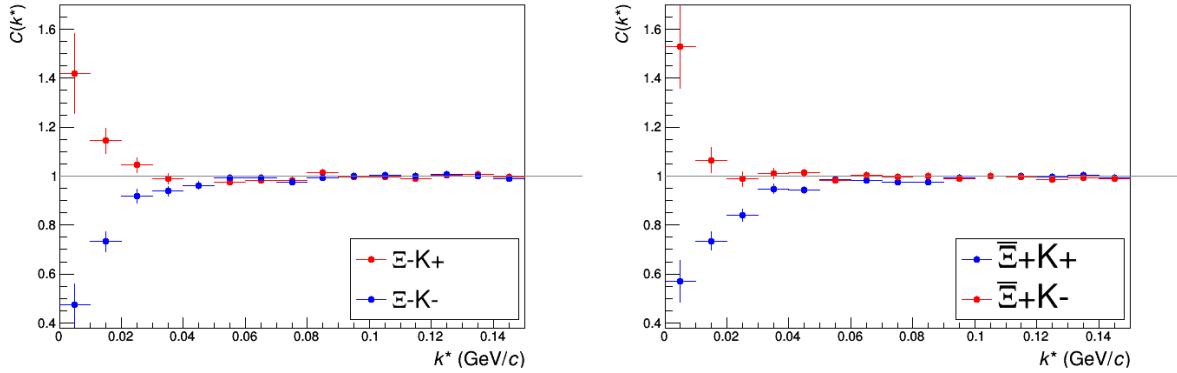


Fig. 60: ΞK^\pm Results for 0-10% Centrality. (Left) $\Xi^- K^+$ and $\Xi^- K^-$ (Right) $\Xi^+ K^+$ and $\Xi^+ K^-$

909 Figure 61 demonstrates graphically, that the $\Xi^- K^+$ results cannot be described by solely the Coulomb
 910 interaction. In this figure, we present the data along with a Coulomb-only band. The Coulomb-only
 911 band is spanned by two Coulomb-only curves, whose parameters are given in the figure. The Coulomb-only
 912 curves were generated using a technique identical to the generation of the fit function, described
 913 in Sec. 5.2, except, of course, with the nuclear scattering parameters all set to zero. The Coulomb-only
 914 curves change monotonically with varying λ or varyin radius parametres, therefore, any curves built with
 915 parameter sets intermediate to those use in the Coulomb-only band will be contained in the band.

916 Including the strong interaction into the simulation can change, sometimes dramatically, the resulting
 917 correlation function, as shown in Figure 62. In the figure, the solid line represents a Coulomb-only curve,
 918 i.e. a simulated correlation function with the strong interaction turned off. The dashed lines represent a
 919 full simulation, including both the strong and Coulomb interactions. The two dashed lines differ only in
 920 the real part of the assumed scattering length: positive in Set 1, and negative in Set 2. In the top figure,
 921 for the $\Xi^- K^+$ simulation, we see that parameter set 2, with a negative real part of the scattering length,
 922 causes the simulated curve to dip below unity, as is seen in the data. If there is a parallel to be drawn
 923 between this analysis and the ΛK analysis, we expect to see similar effects in the ΛK^+ system and the
 924 $\Xi^- K^+$ systems. In these systems, we could have an $s\bar{s}$ annihilation picture. Or, another possible way of
 925 thinking about these systems is in terms of net strangeness. The ΛK^+ system has $S=0$, while the ΛK^-
 926 has $S=-2$. The $\Xi^- K^+$ has $S=-1$, while the $\Xi^- K^-$ has $S=-3$.

927 The author was asked to perform a global Coulomb-only fit to the data, to ensure that the system truly
 928 could not be described simply by the Coulomb interaction. In order words, in the fit, the strong force was
 929 turned off, and the $\Xi^- K^+$, $\Xi^+ K^-$, $\Xi^- K^-$, $\Xi^+ K^+$ systems all share one sinlge radius parameter, while the
 930 pair and conjugate pair systems share a λ parameter. The results of this fit are shown in Figures 63 and
 931 64. In Fig. 63, there was a lower limit of 0.1 fm placed on the radius parameter, and the radius parameter
 932 was initialized to 3 fm (as seems reasonable, when considering the transverse mass of the system and
 933 looking at Fig. ??). As is shown in the results, the radius parameter reached this unrealistic lower bound
 934 of 0.1 fm. In Fig. 64, the parameters were all unbounded, and the radius parameter was initialized to 10
 935 fm. In this case, the radius parameters reamins high, and ends at an unrealistic value of 10.84 fm. In both
 936 cases, the λ parameters are too low. From these figures, we conclude that a global Coulomb-only fit is
 937 not suitable for the data.

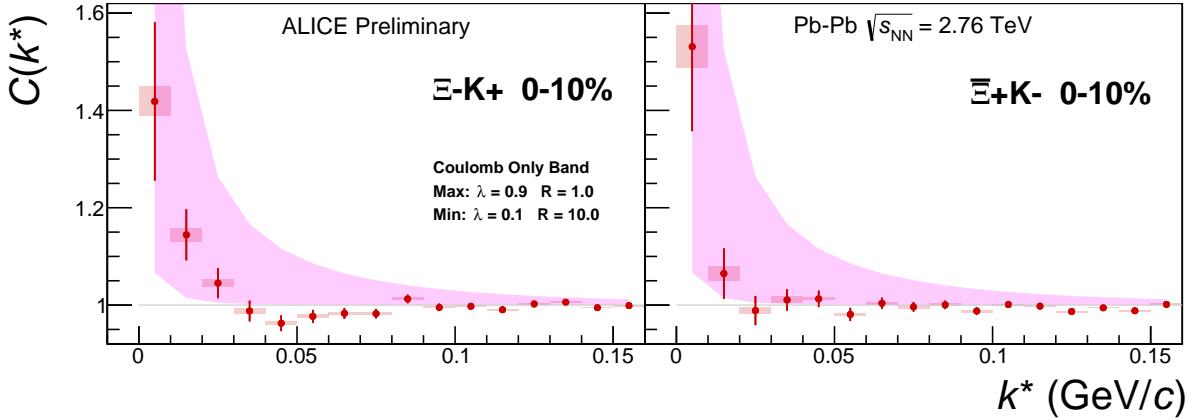
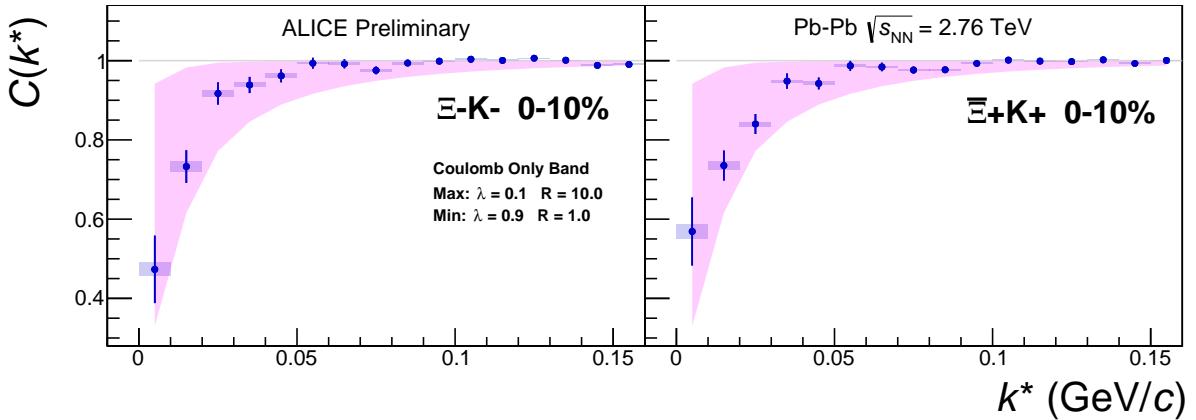
(a) (Left) ΞK^+ and (Right) $\Xi\text{-K}^-$ (b) (Left) ΞK^- and (Right) $\Xi\text{-K}^+$

Fig. 61: ΞK^\pm data with Coulomb-only bands for the 0-10% centrality bin. The Coulomb-only bands span two sets of Coulomb-only curves: (1) $\lambda = 0.9$, $R = 1.0$ fm and (2) $\lambda = 0.1$, $R = 10.0$ fm. The Coulomb-only curves are simulated correlation functions for the respective pair system assuming only a Coulomb interaction, i.e. ignoring the strong interaction. The Coulomb-only curves change monotonically with varying λ and varying R , therefore, any intermediate parameter set will fall within this Coulomb-only band.

938 Although the global Coulomb-only fit failed, it is possible that a Coulomb-only fit performed on $\Xi^- \text{K}^+$
 939 and $\Xi^+ \text{K}^-$ separately from $\Xi^- \text{K}^-$ and $\Xi^+ \text{K}^+$ could be suitable. The result of such fits are shown in
 940 Figures 65 and 66. Figure 65, shows that the fit is not able to describe the dip in the $\Xi^- \text{K}^+$ data below
 941 unity. Of course, this is obviously true for an attractive Coulomb-only fit. The radius parameter of
 942 8.43 fm extracted from this fit is unrealistically large. In Figure 66 shows the Coulomb-only fit can
 943 described the $\Xi^- \text{K}^-$ data reasonable well; although the extracted radius of 3.73 fm is somewhat larger
 944 than expected.

945 8 To Do

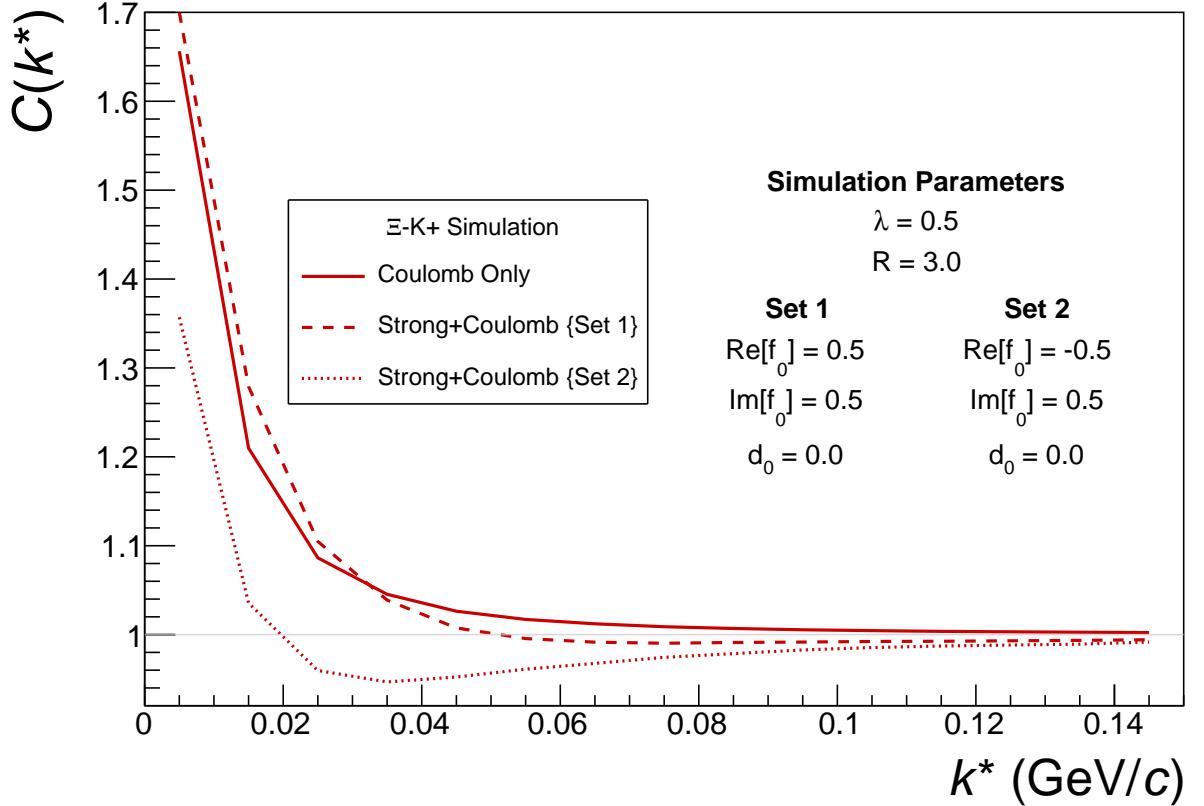
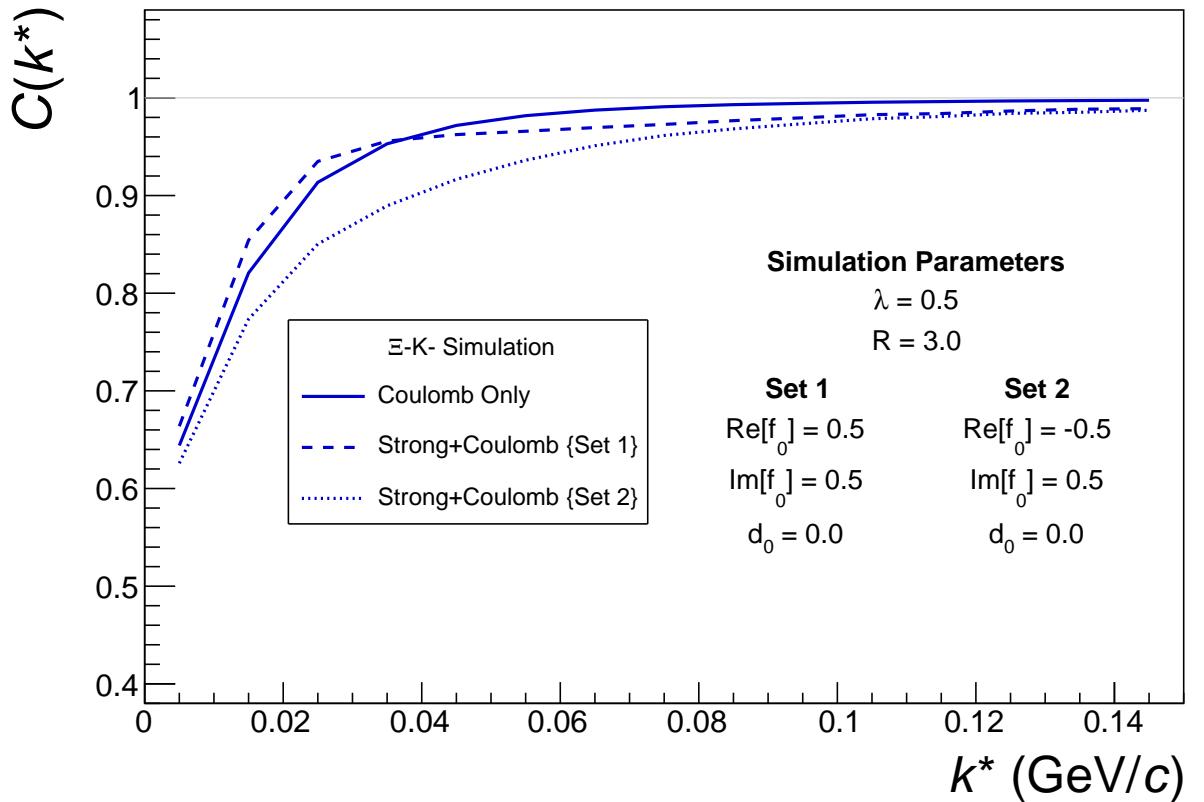

 (a) ΞK^+ and $\bar{\Xi}\text{K}^-$ simulation

 (b) ΞK^- and $\bar{\Xi}\text{K}^+$ simulation

Fig. 62: Effect on the Coulomb-only curve of including the strong interaction for ΞK^\pm systems. The solid line represents a Coulomb-only curve, i.e. a simulated correlation function with the strong interaction turned off. The dashed lines represent a full simulation, including both the strong and Coulomb interactions. The two dashed lines differ only in the real part of the assumed scattering length: positive in Set 1, and negative in Set 2.

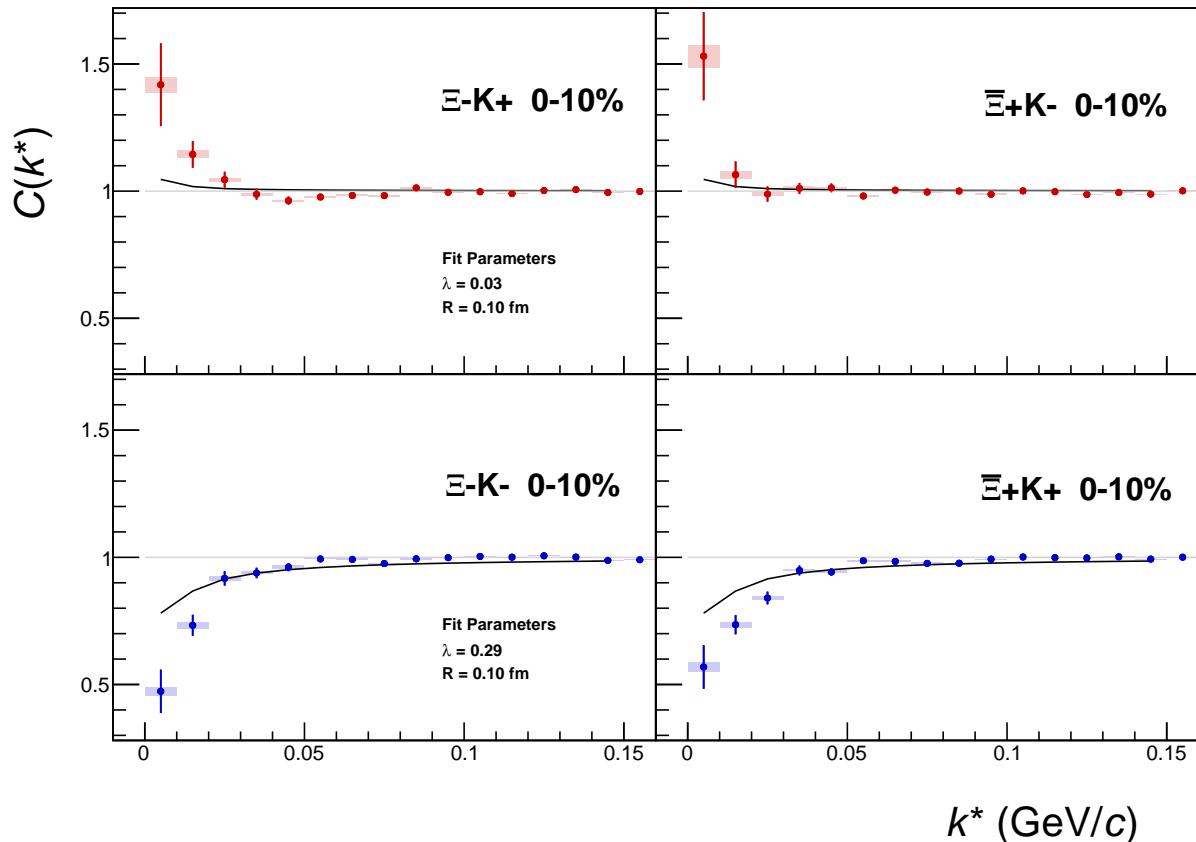


Fig. 63: ΞK^\pm Global Coulomb-only fit (Set 1) for 0-10% centrality. In this fit, there was a lower limit of 0.1 fm placed on the radius parameter, and the radius parameter was initialized to 3 fm (as seems reasonable, when considering the transverse mass of the system and looking at Fig. ??). As is shown in the results, the radius parameter reached this unrealistic lower bound of 0.1 fm. Also, the extracted λ parameters are too low.

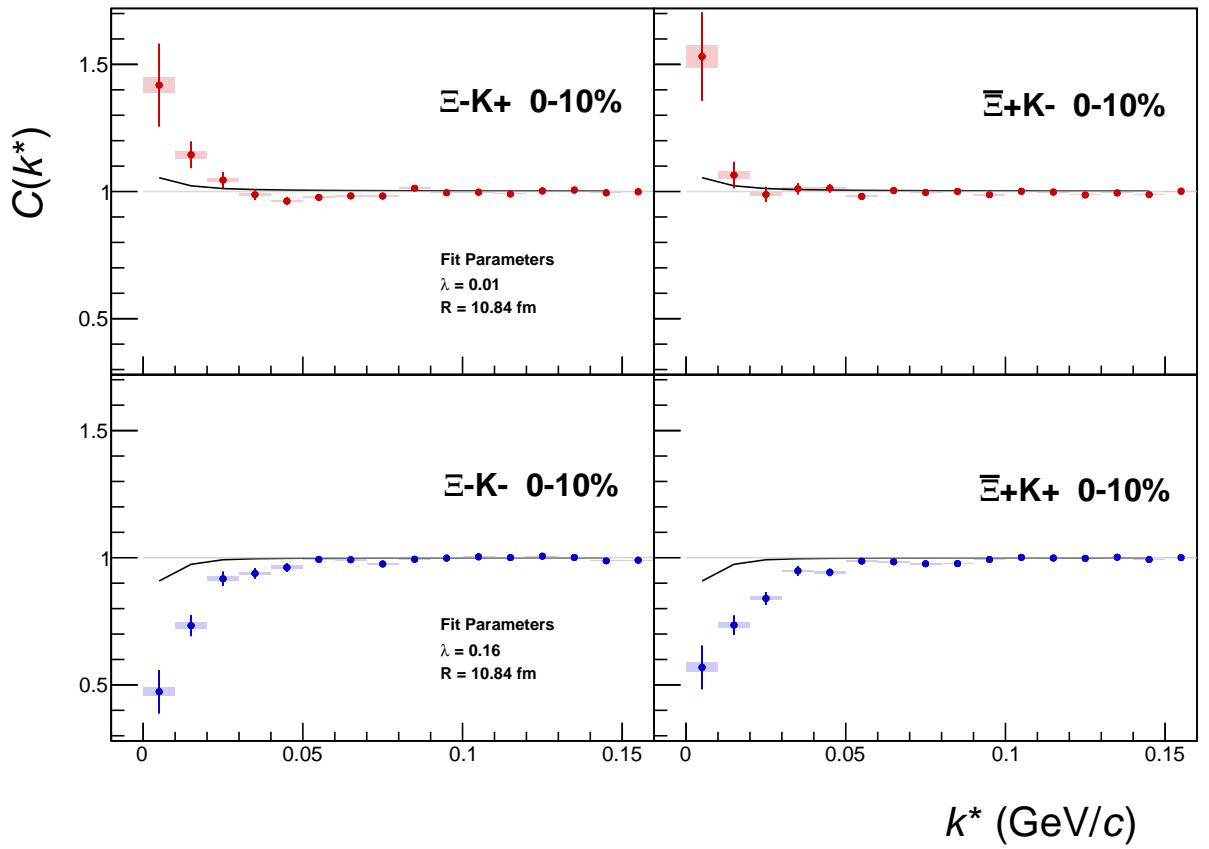


Fig. 64: ΞK^\pm Global Coulomb-only fit (Set 2) for 0-10% centrality. In this fit, the parameters were all unbounded, and the radius parameter was initialized to 10 fm. In this case, the radius parameters remain high, and ends at an unrealistic value of 10.84 fm. Also, the extracted λ parameters are too low.

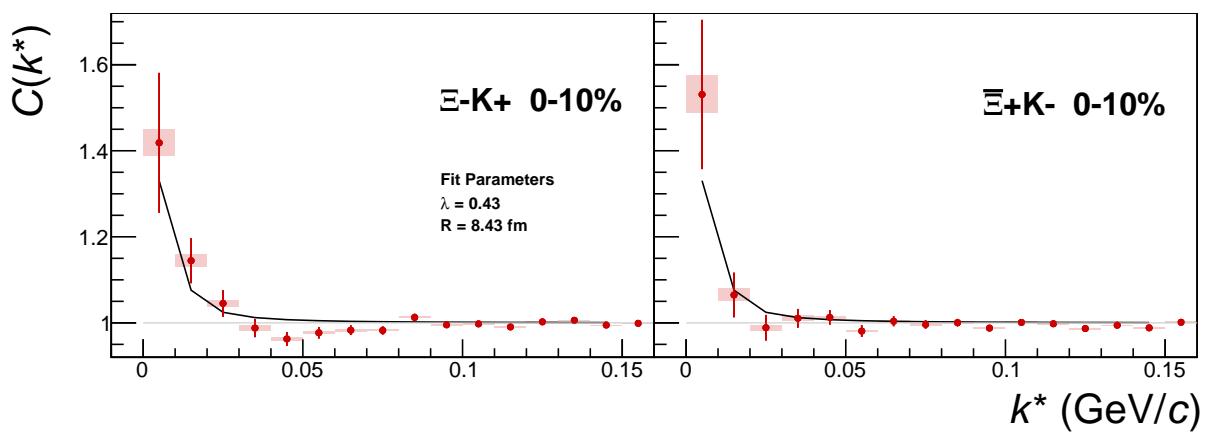


Fig. 65: $\Xi^- K^+$ Coulomb-only fit for 0-10% centrality

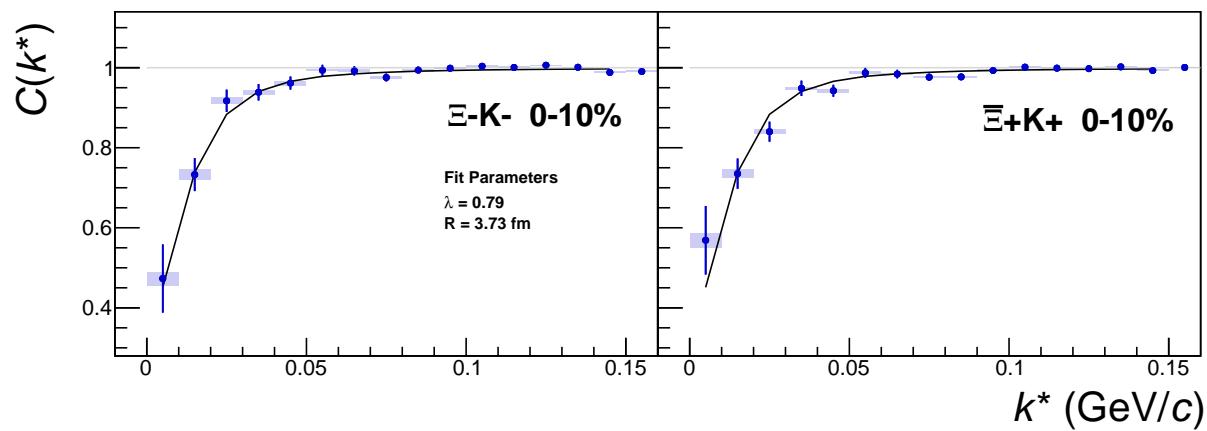


Fig. 66: $\Xi^- K^-$ Coulomb-only fit for 0-10% centrality

9 Additional Figures

9.1 Residuals

9.1.1 ΛK^+ Residuals

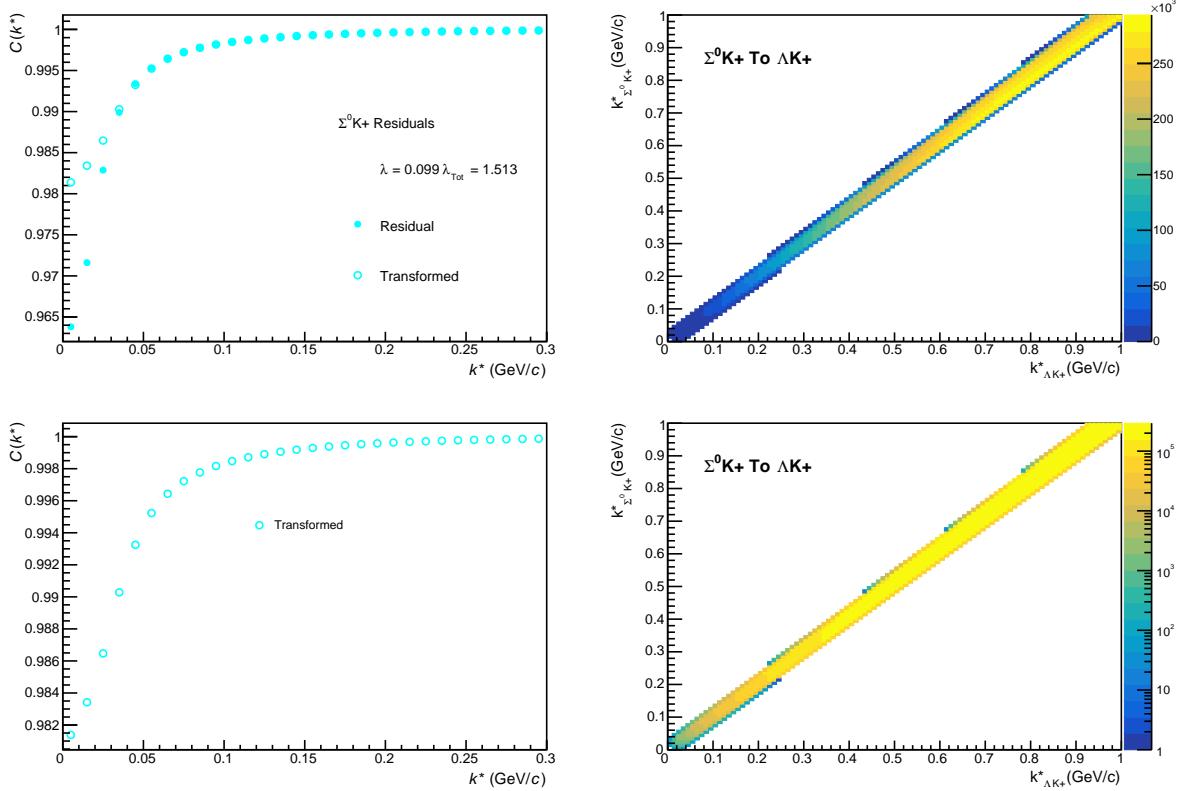


Fig. 67: Residuals: $\Sigma^0 K^+$ to ΛK^+ (0-10% Centrality)

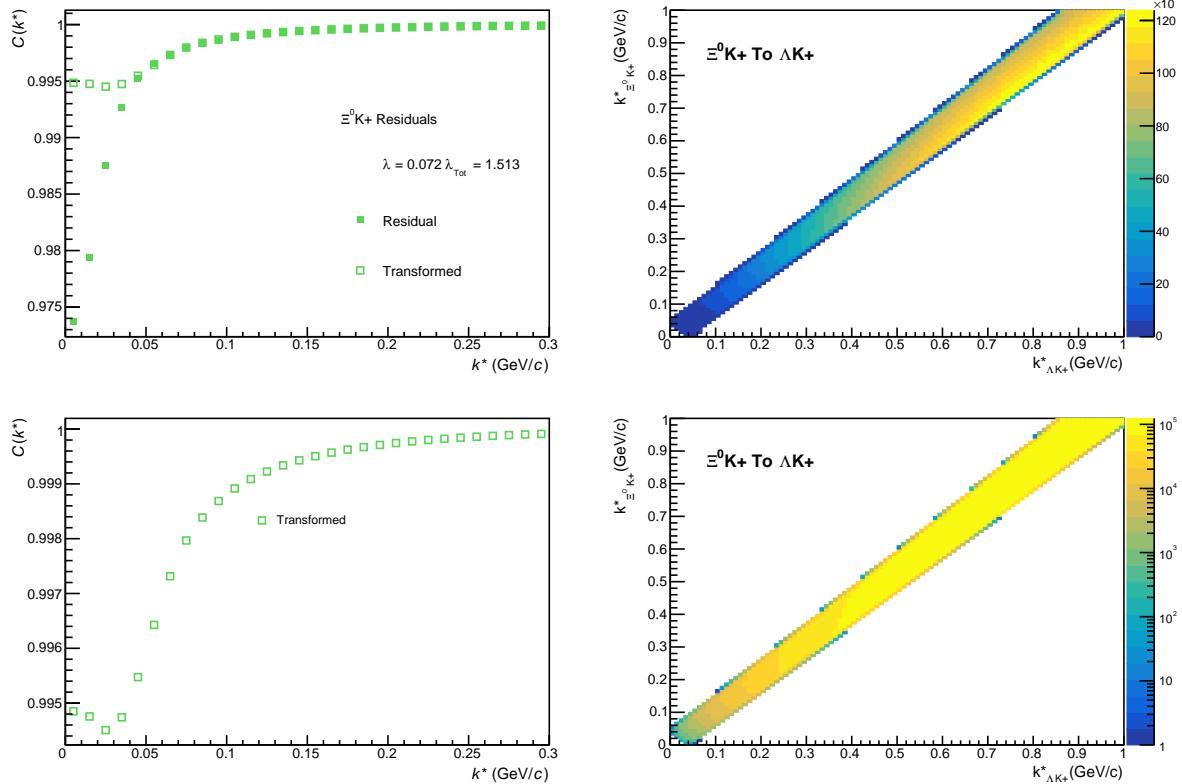


Fig. 68: Residuals: $\Xi^0 \text{K}^+$ to ΛK^+ (0-10% Centrality)

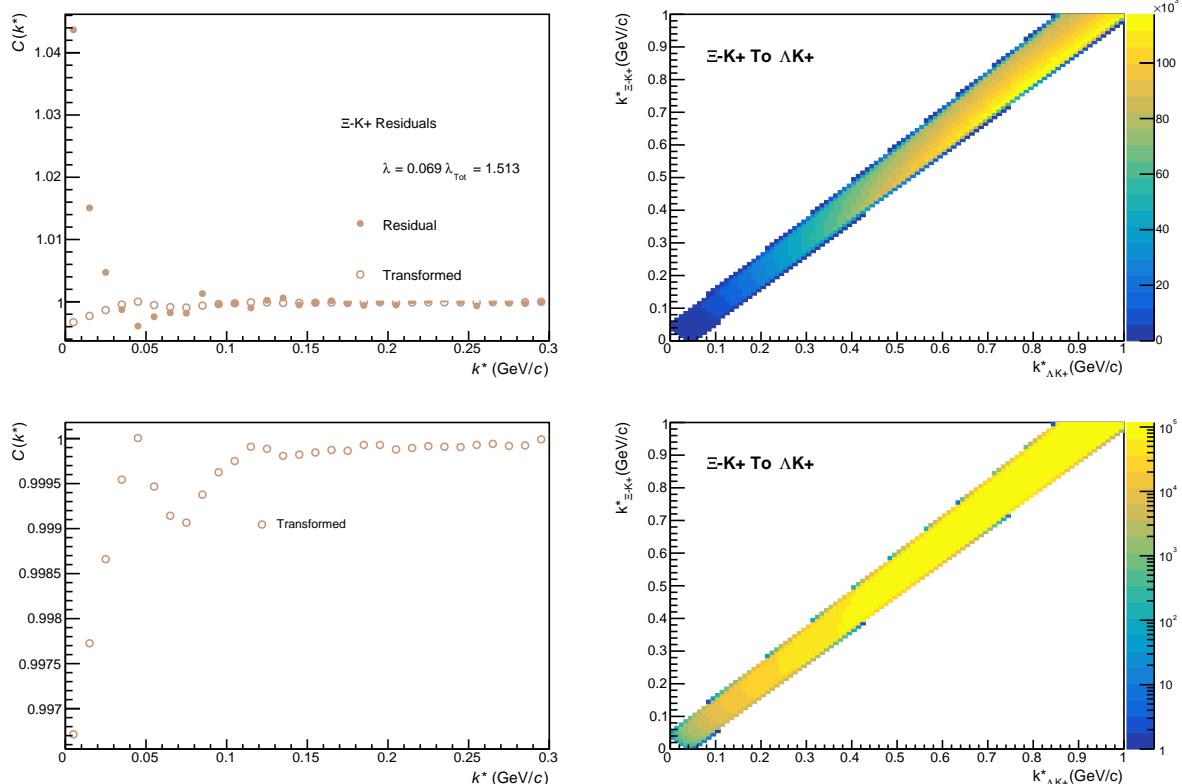
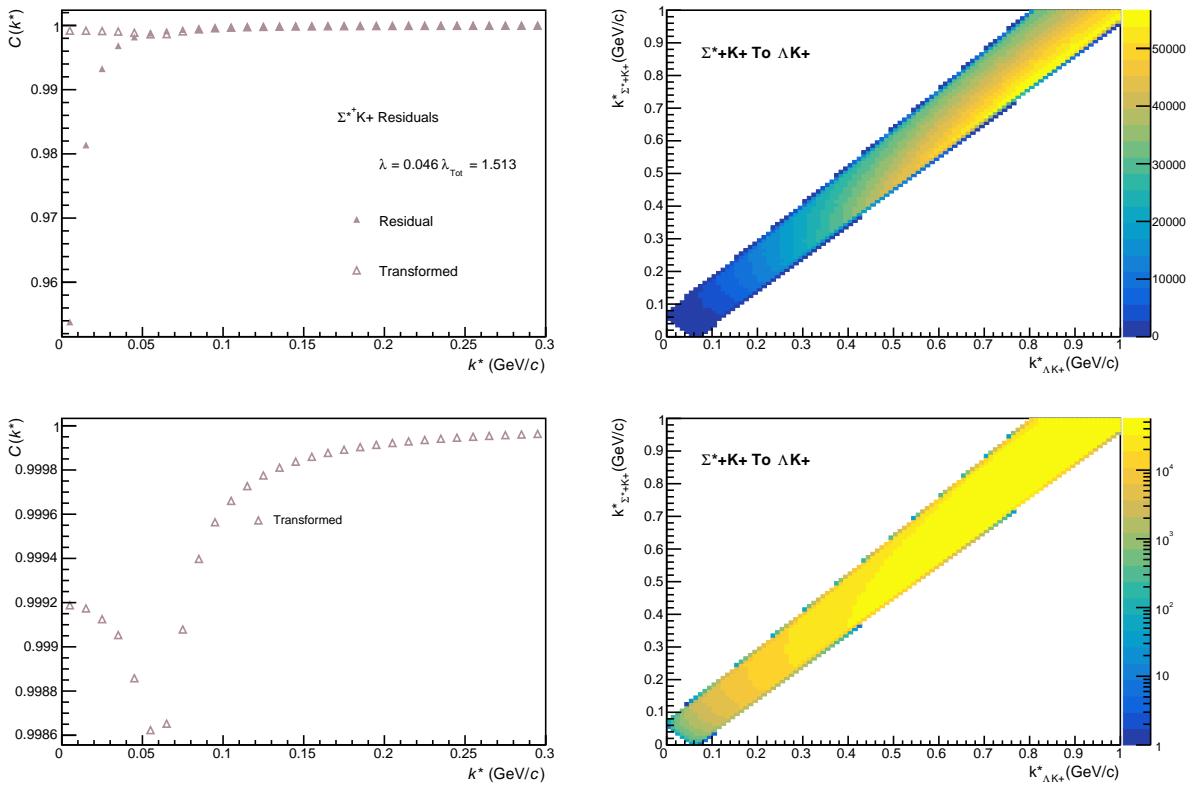
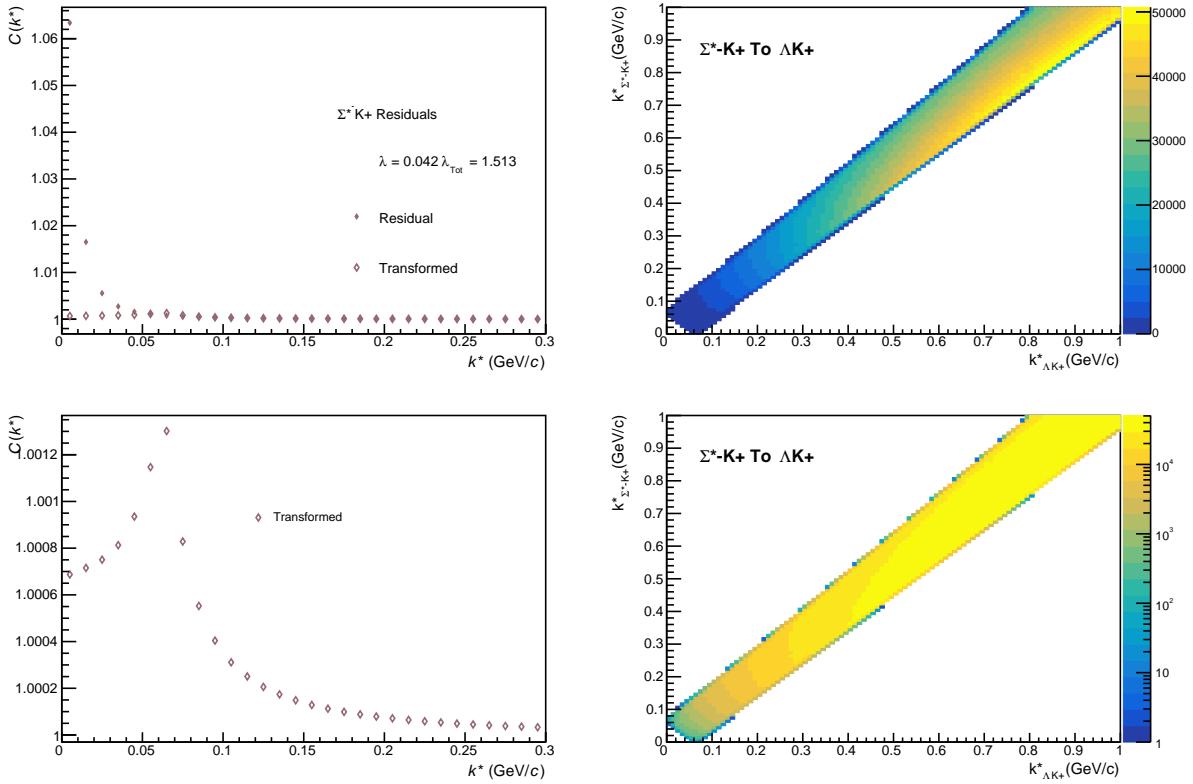


Fig. 69: Residuals: $\Xi^- \text{K}^+$ to ΛK^+ (0-10% Centrality)


 Fig. 70: Residuals: $\Sigma^{*+}K^+$ to ΛK^+ (0-10% Centrality)

 Fig. 71: Residuals: $\Sigma^{*-}K^+$ to ΛK^+ (0-10% Centrality)

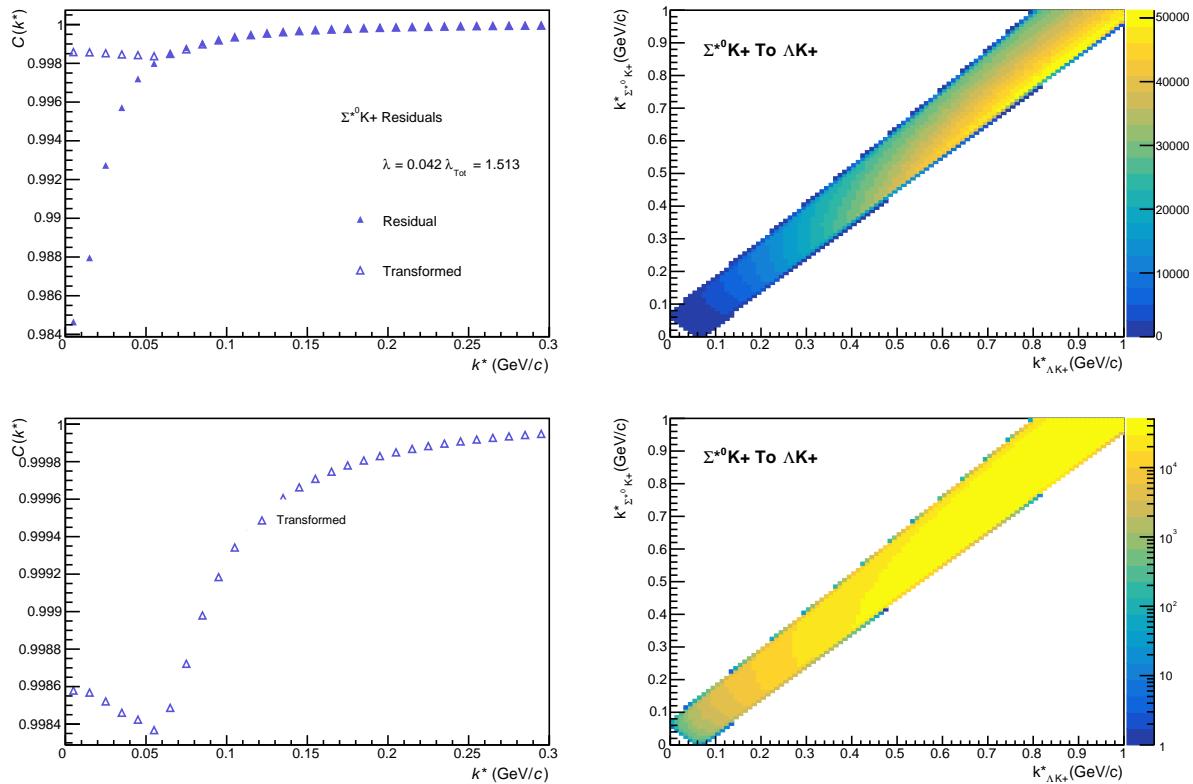


Fig. 72: Residuals: $\Sigma^{*0} \text{K}^+$ to ΛK^+ (0-10% Centrality)

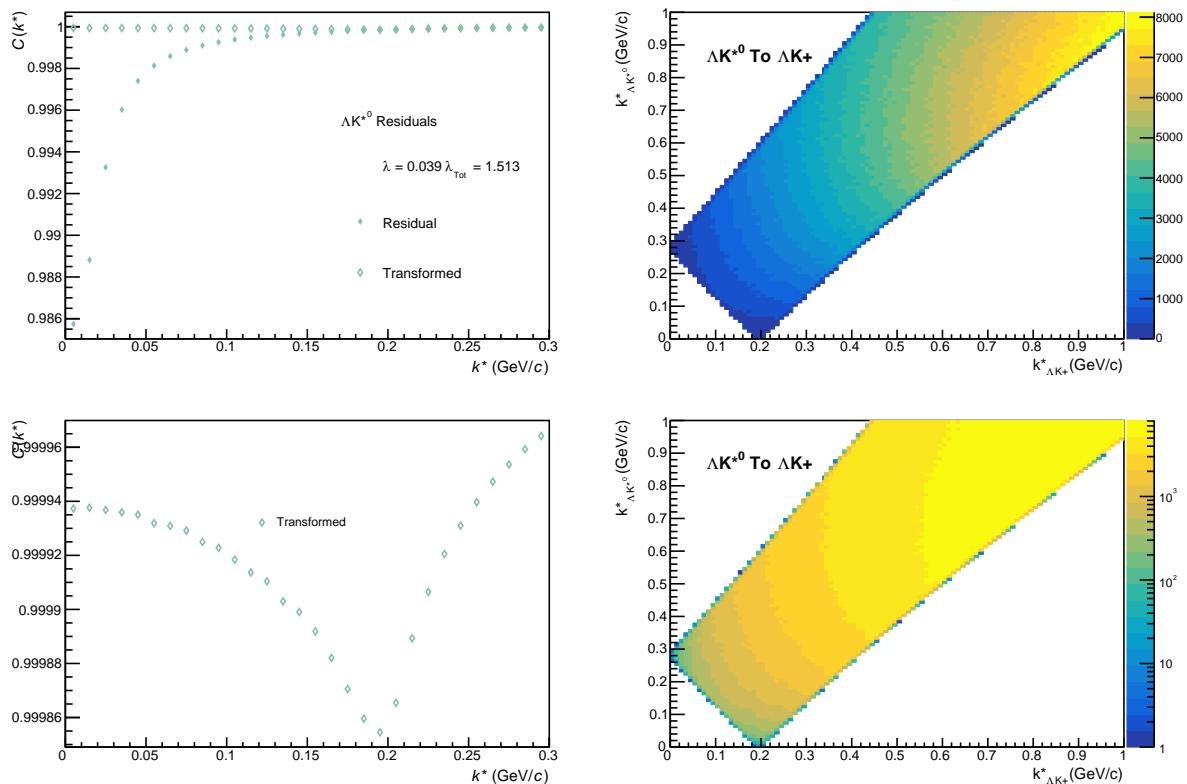


Fig. 73: Residuals: ΛK^{*0} to ΛK^+ (0-10% Centrality)

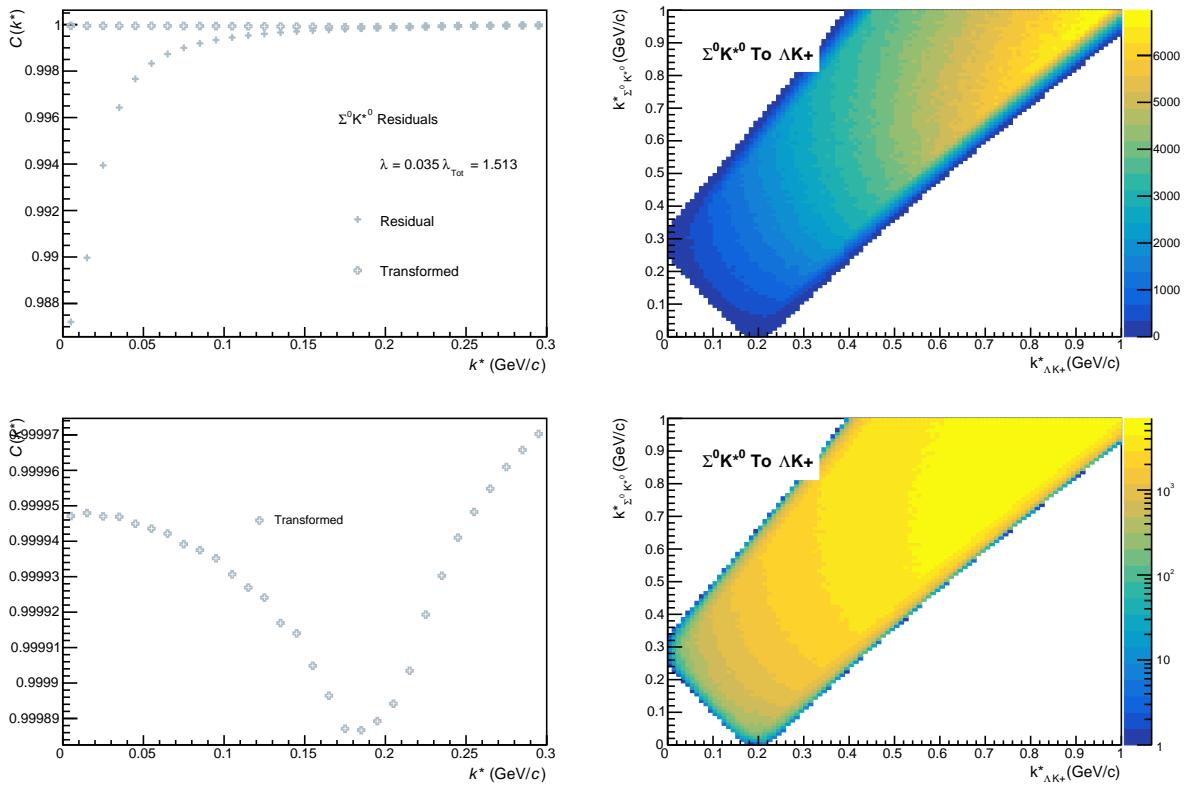


Fig. 74: Residuals: $\Sigma^0 K^{*0}$ to ΛK^+ (0-10% Centrality)

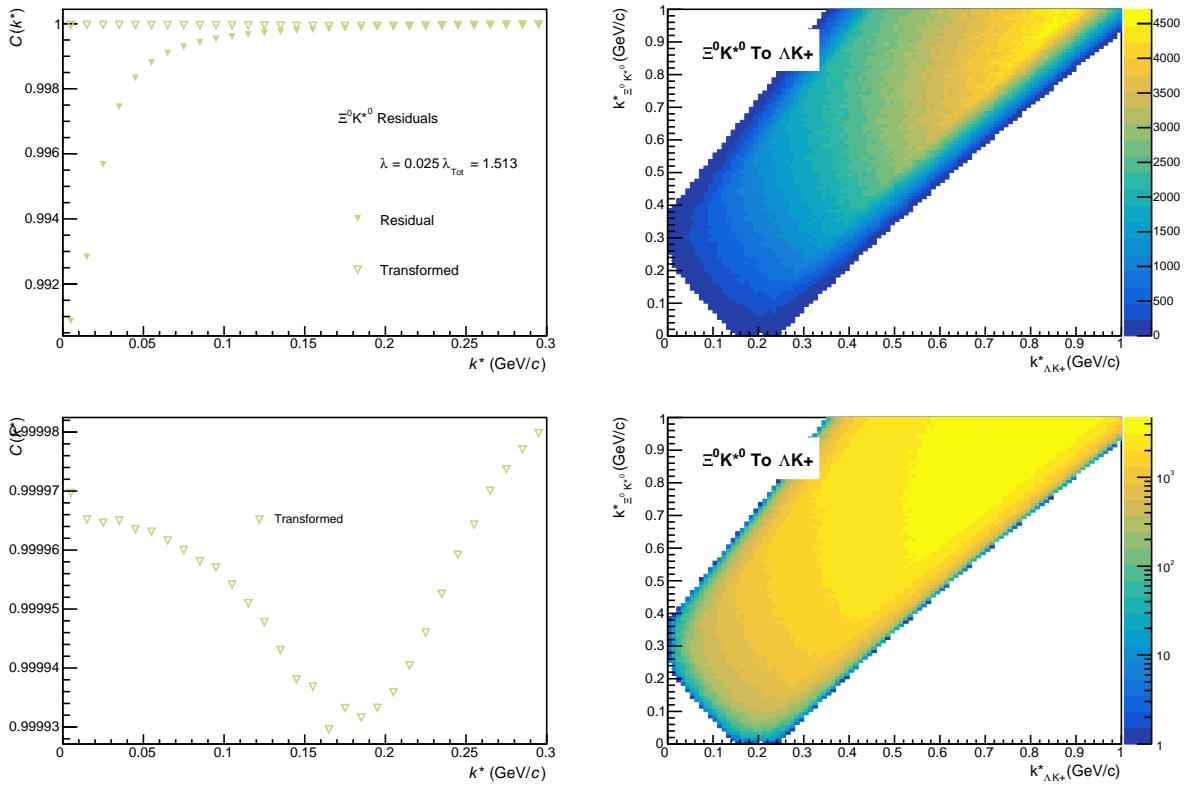


Fig. 75: Residuals: $\Xi^0 K^{*0}$ to ΛK^+ (0-10% Centrality)

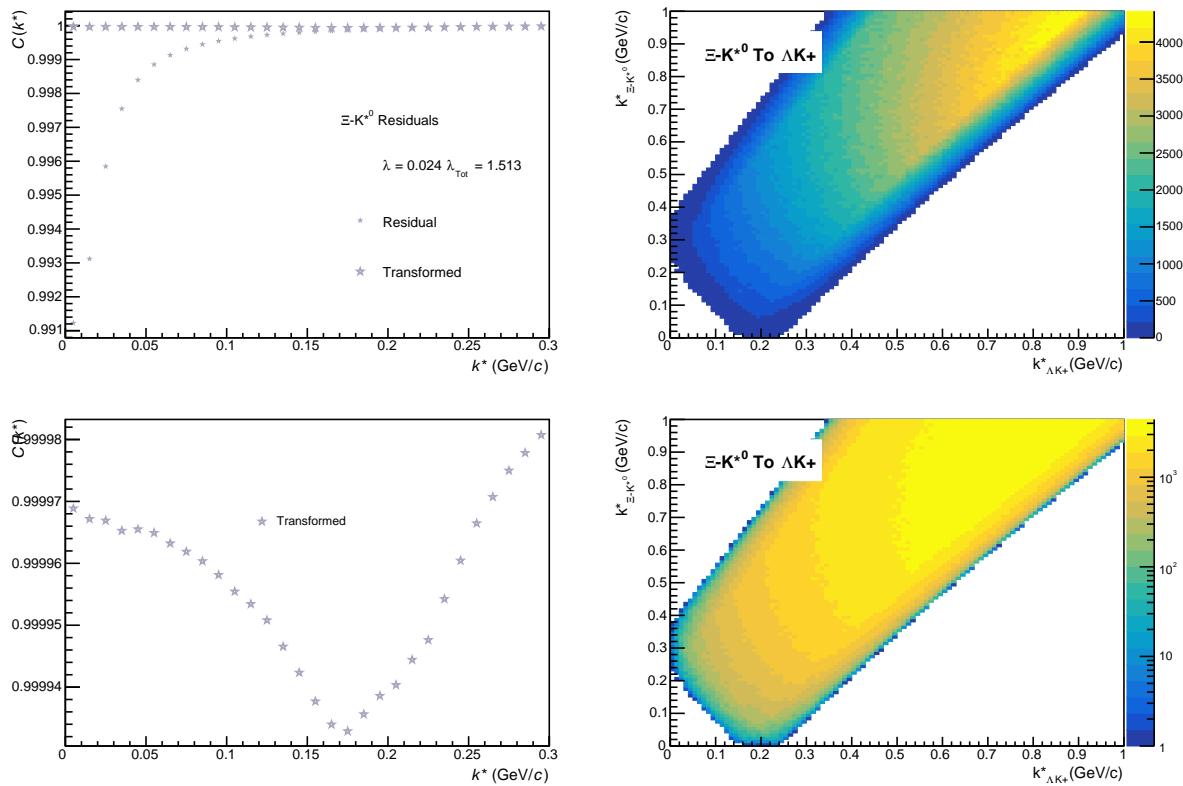


Fig. 76: Residuals: $\Xi^- K^{*0}$ to ΛK^+ (0-10% Centrality)

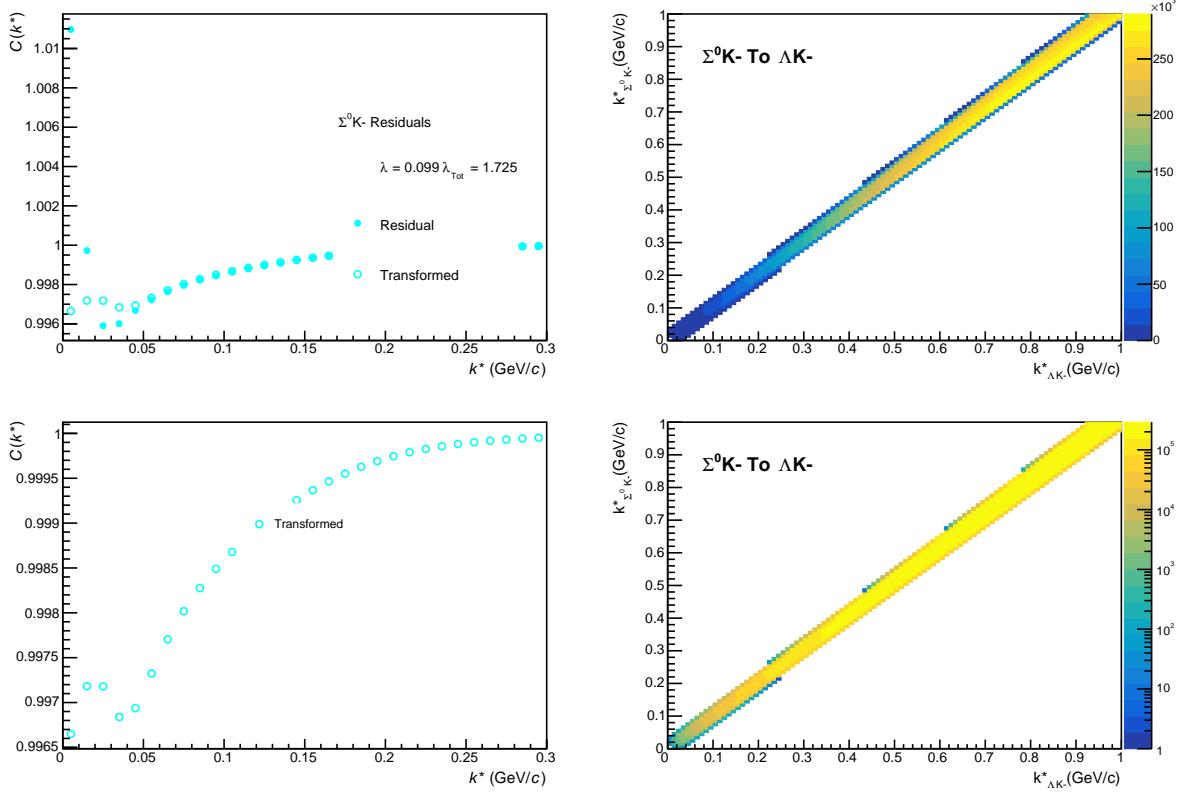
9.1.2 ΛK^- Residuals


Fig. 77: Residuals: $\Sigma^0 K^-$ to ΛK^- (0-10% Centrality)

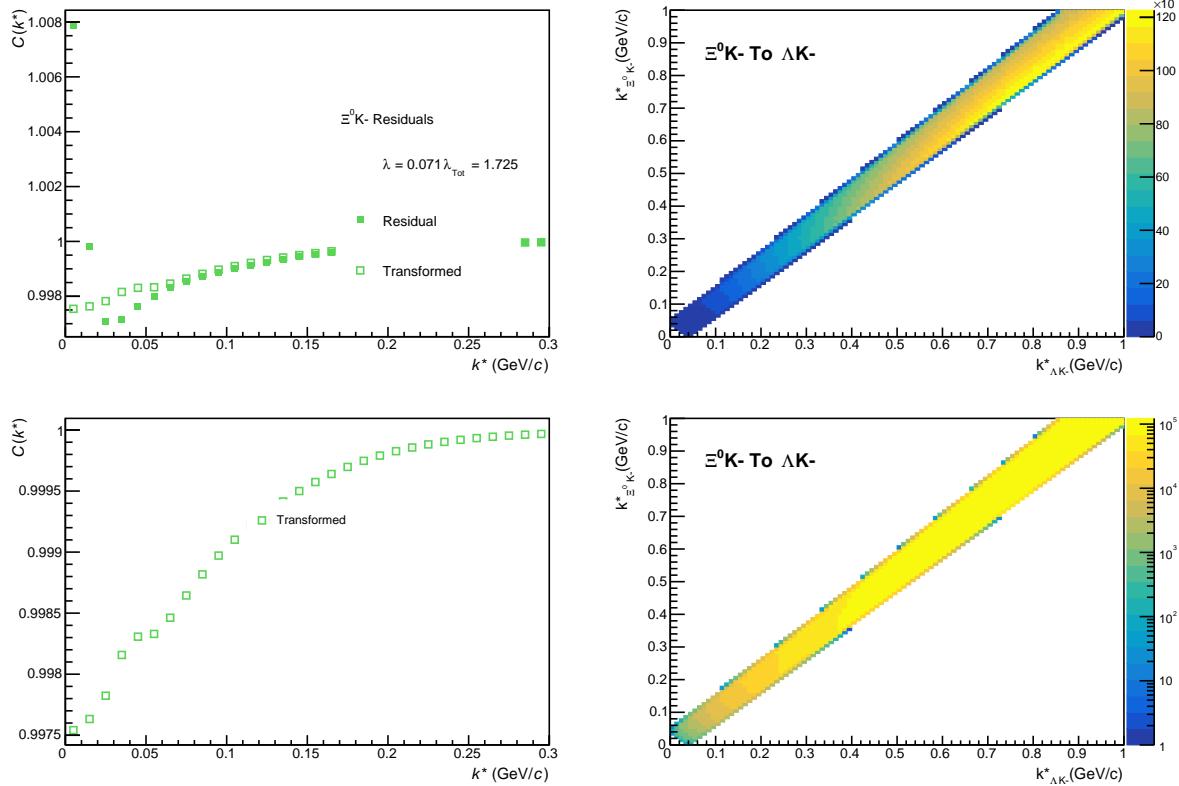


Fig. 78: Residuals: $\Xi^0 \text{K}^-$ to ΛK^- (0-10% Centrality)

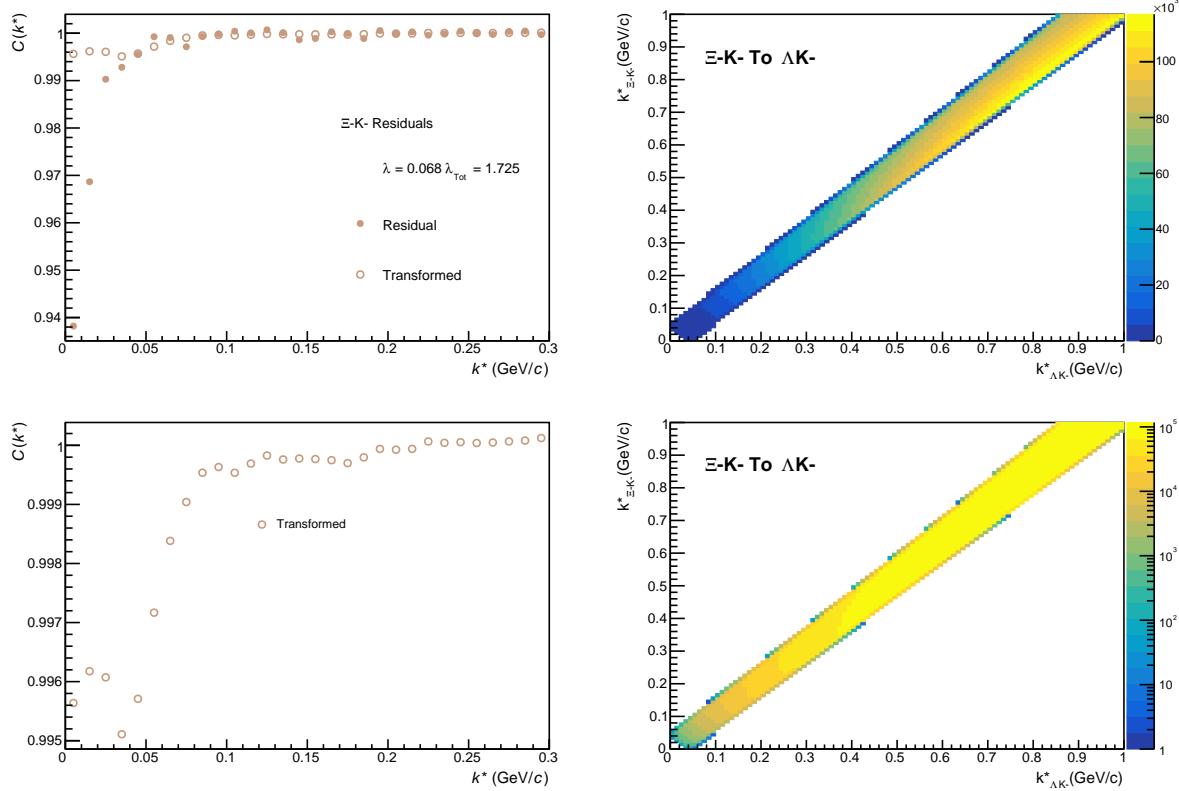
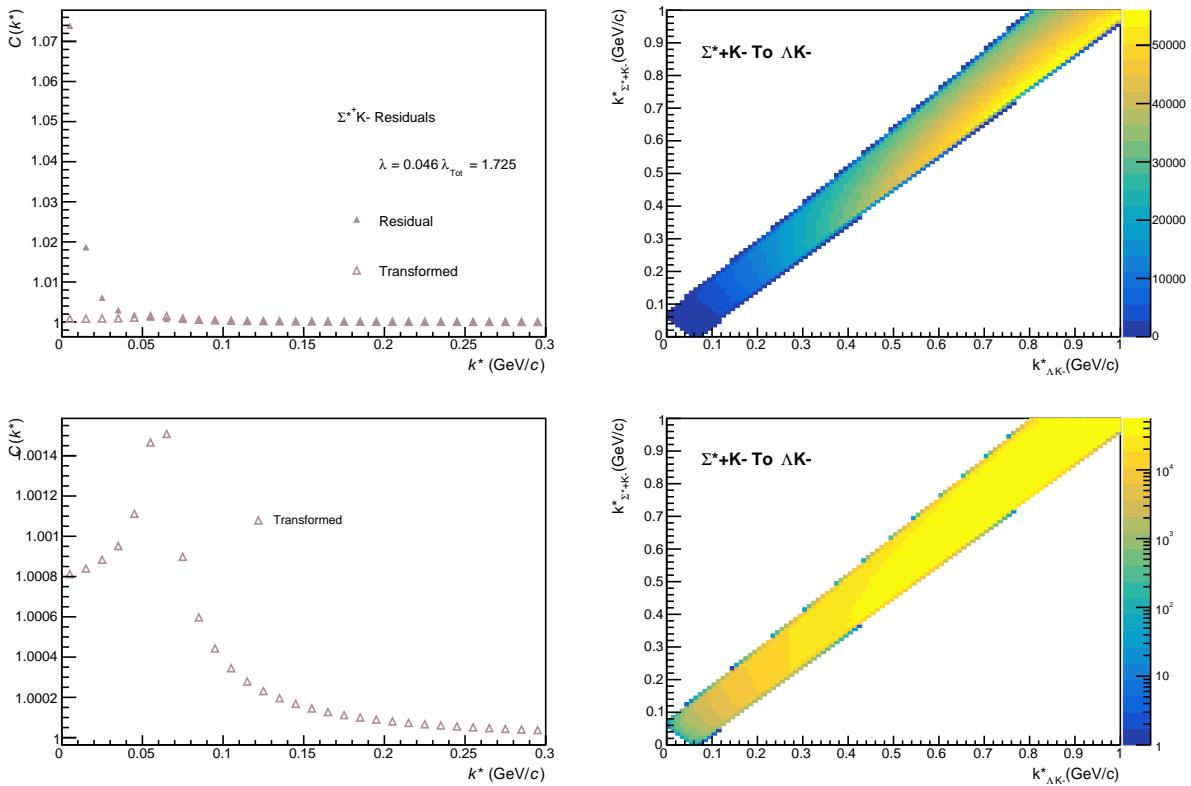
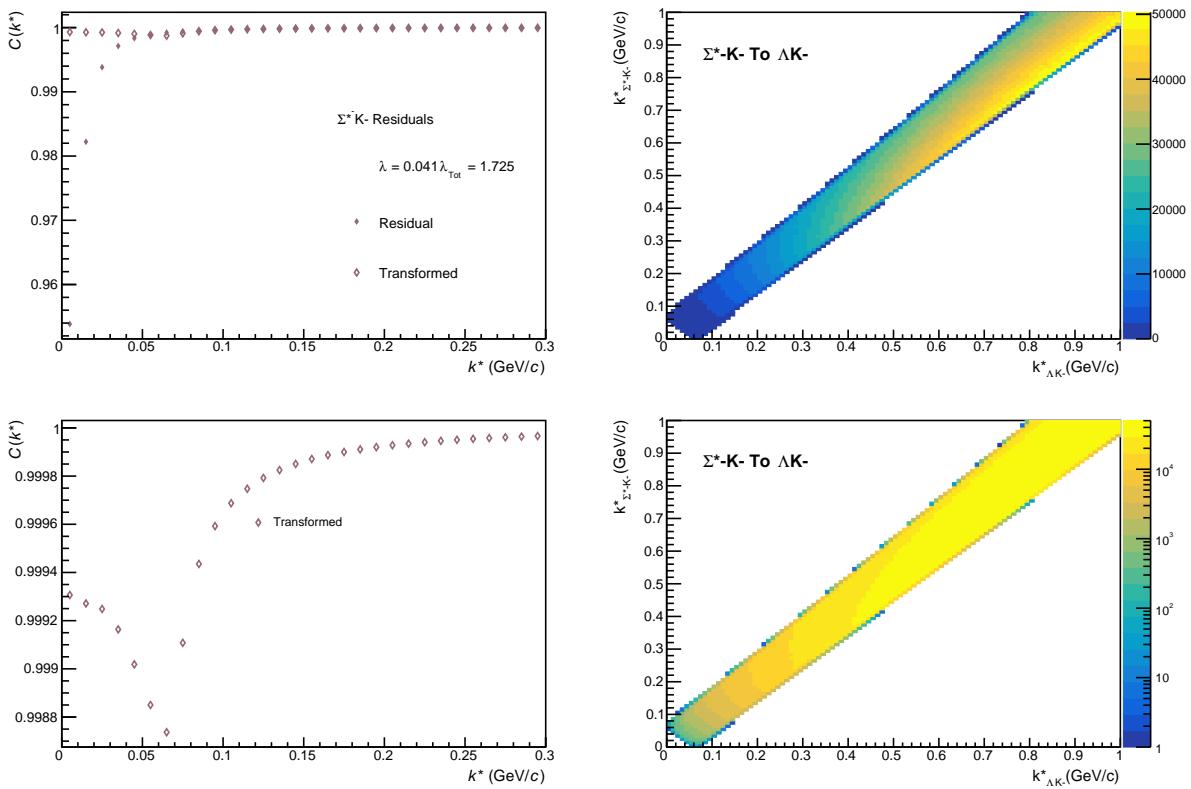


Fig. 79: Residuals: $\Xi^- \text{K}^-$ to ΛK^- (0-10% Centrality)


Fig. 80: Residuals: $\Sigma^{+} K^{-}$ to ΛK^{-} (0-10% Centrality)

Fig. 81: Residuals: $\Sigma^- K^-$ to ΛK^- (0-10% Centrality)

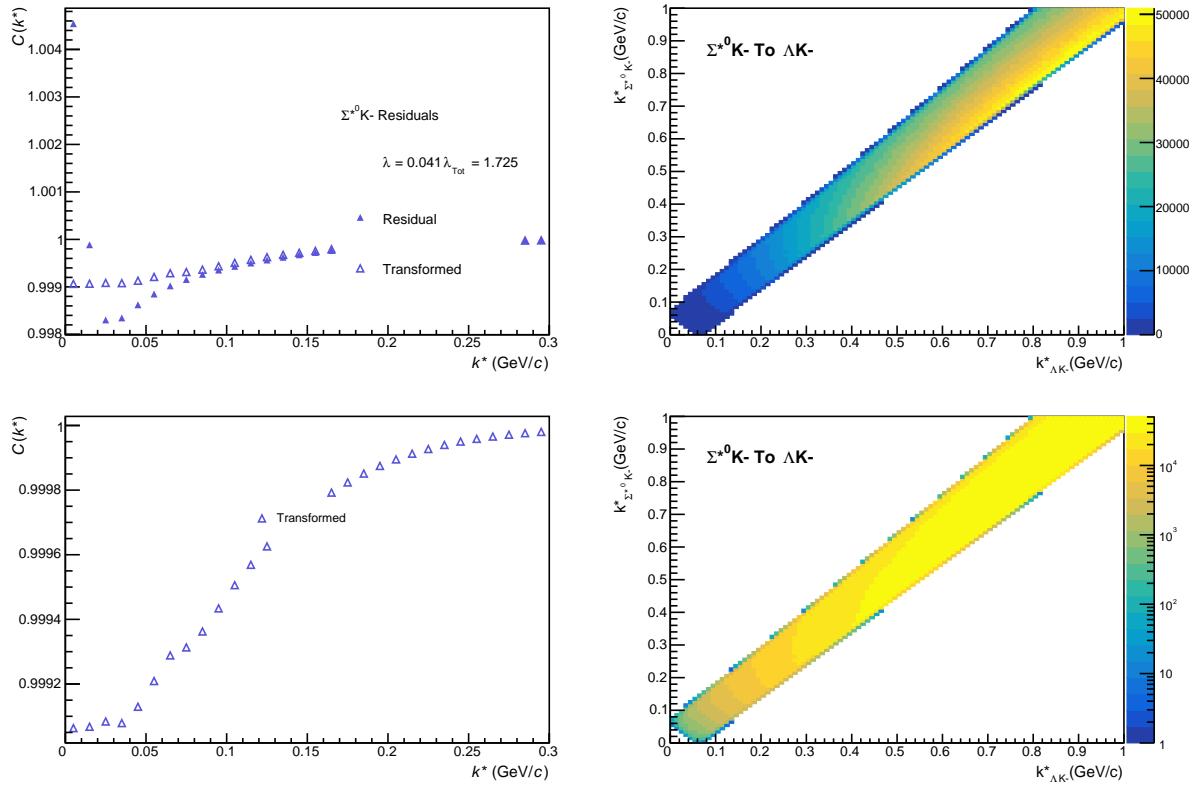


Fig. 82: Residuals: $\Sigma^0 \text{K}^-$ to ΛK^- (0-10% Centrality)

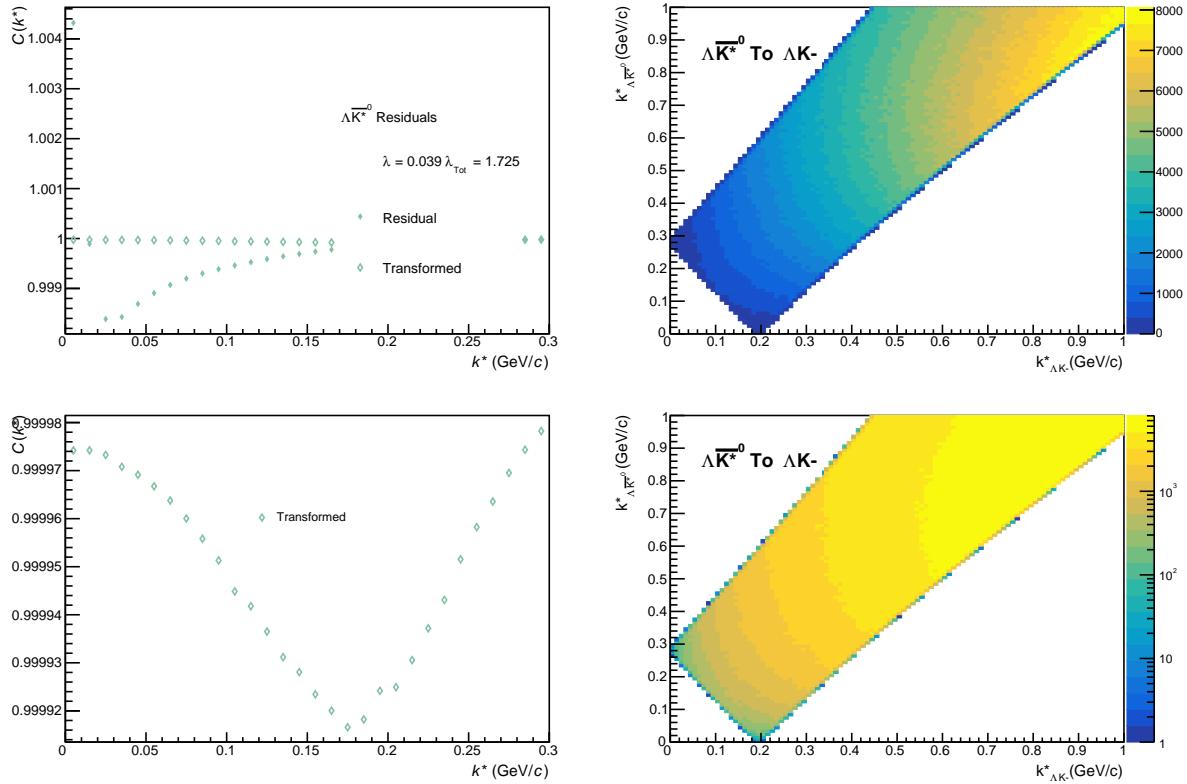


Fig. 83: Residuals: $\Lambda \bar{K}^0$ to ΛK^- (0-10% Centrality)

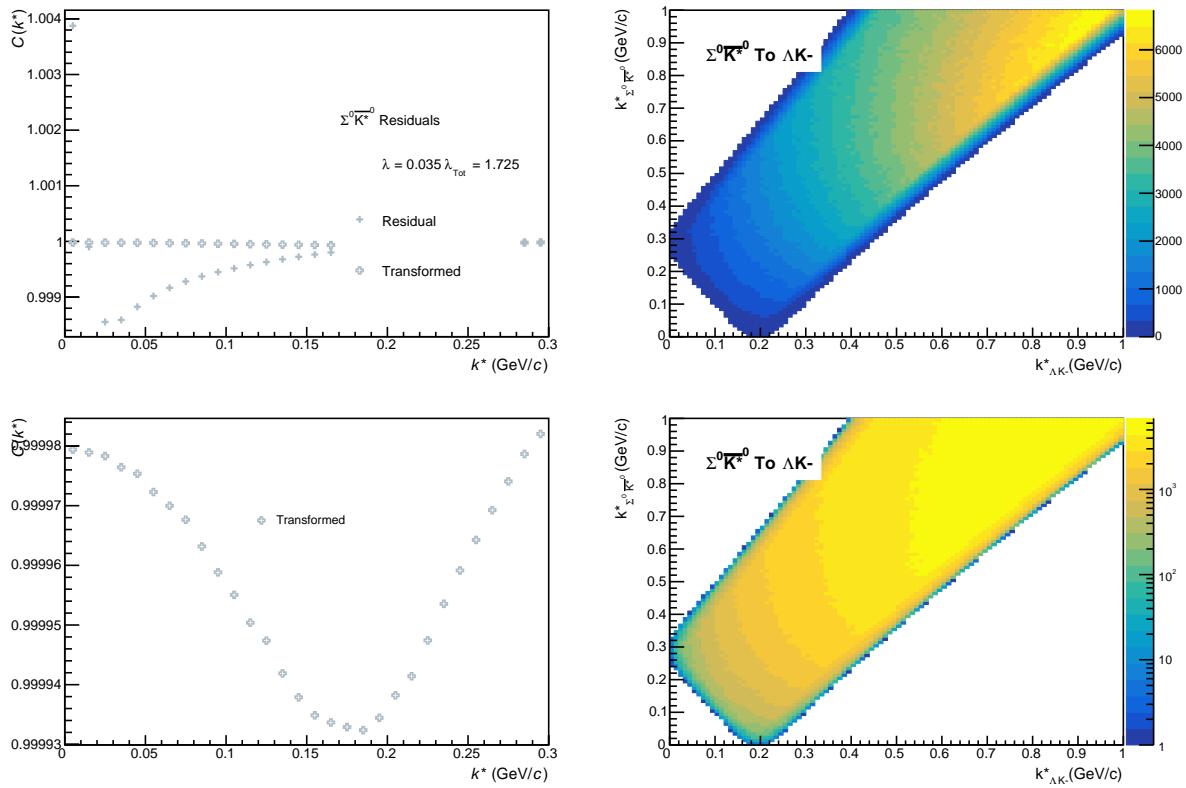


Fig. 84: Residuals: $\Sigma^0 \bar{K}^{*0}$ to ΛK^- (0-10% Centrality)

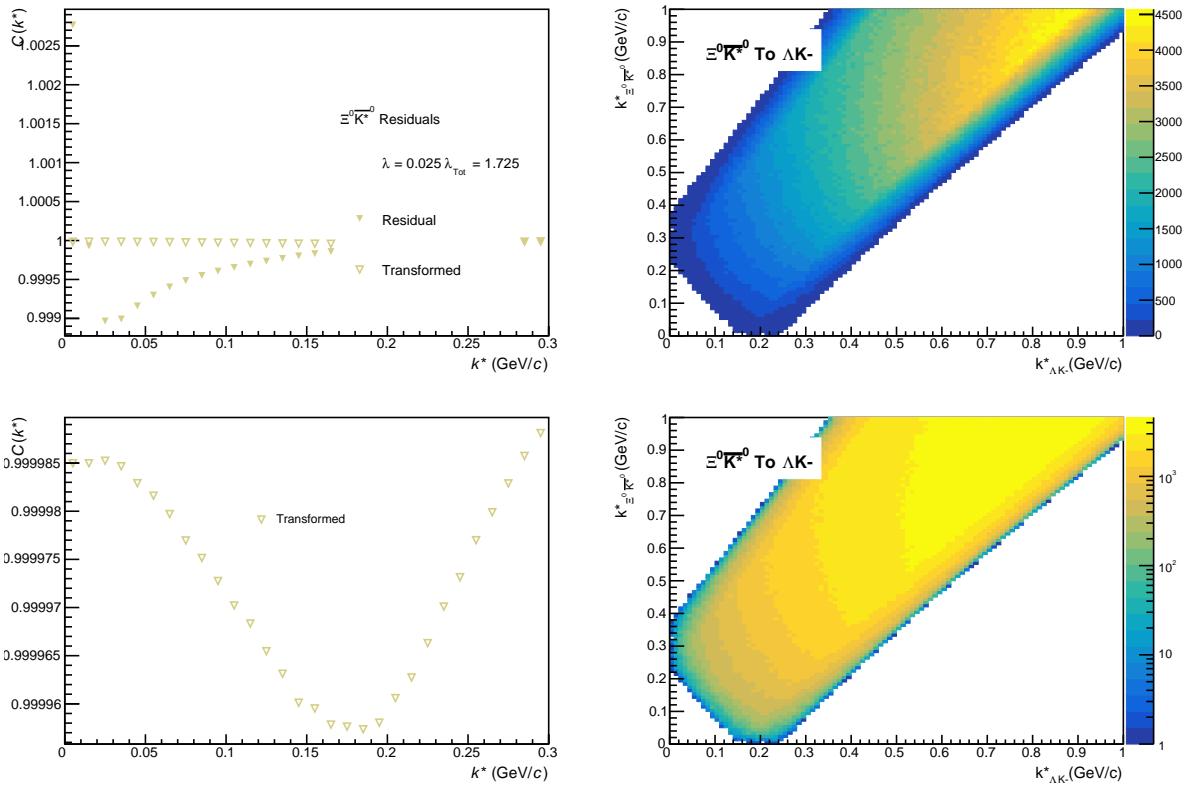


Fig. 85: Residuals: $\Xi^0 \bar{K}^{*0}$ to ΛK^- (0-10% Centrality)

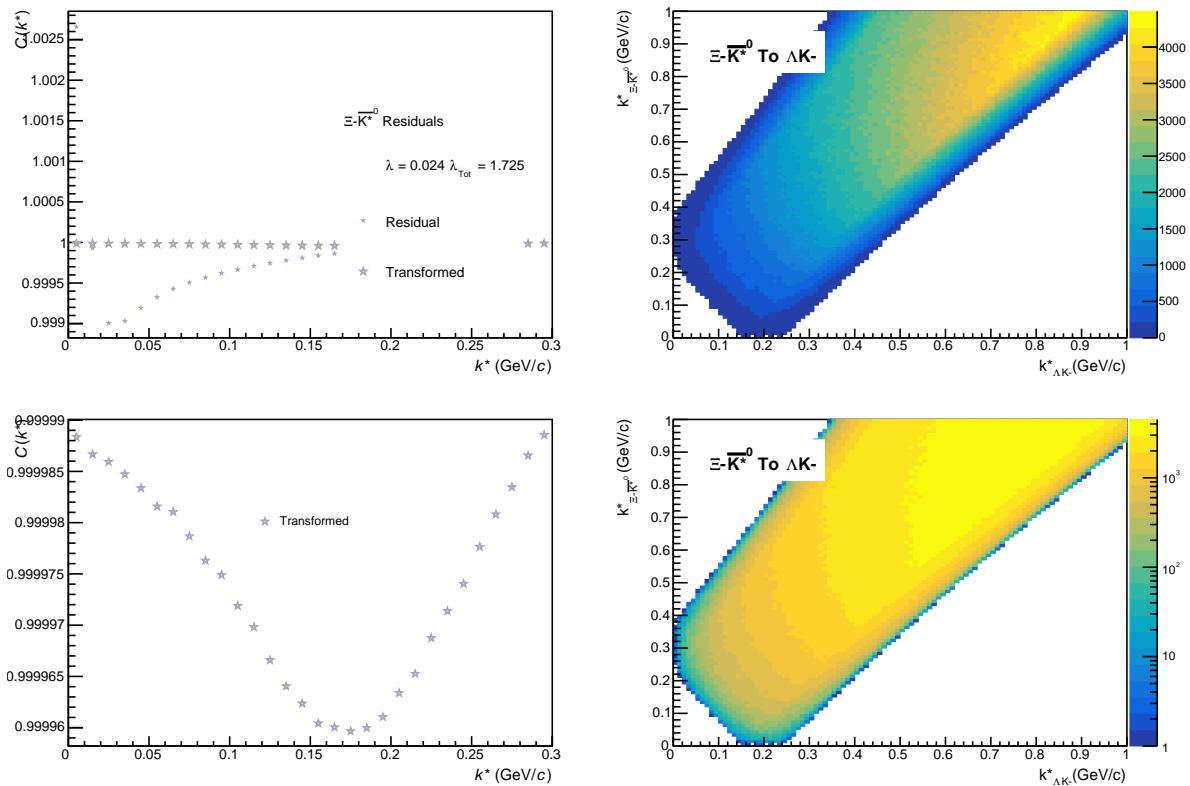


Fig. 86: Residuals: $\Xi^- \bar{K}^{*0}$ to ΛK^- (0-10% Centrality)

950 **9.1.3 ΛK_S^0 Residuals**

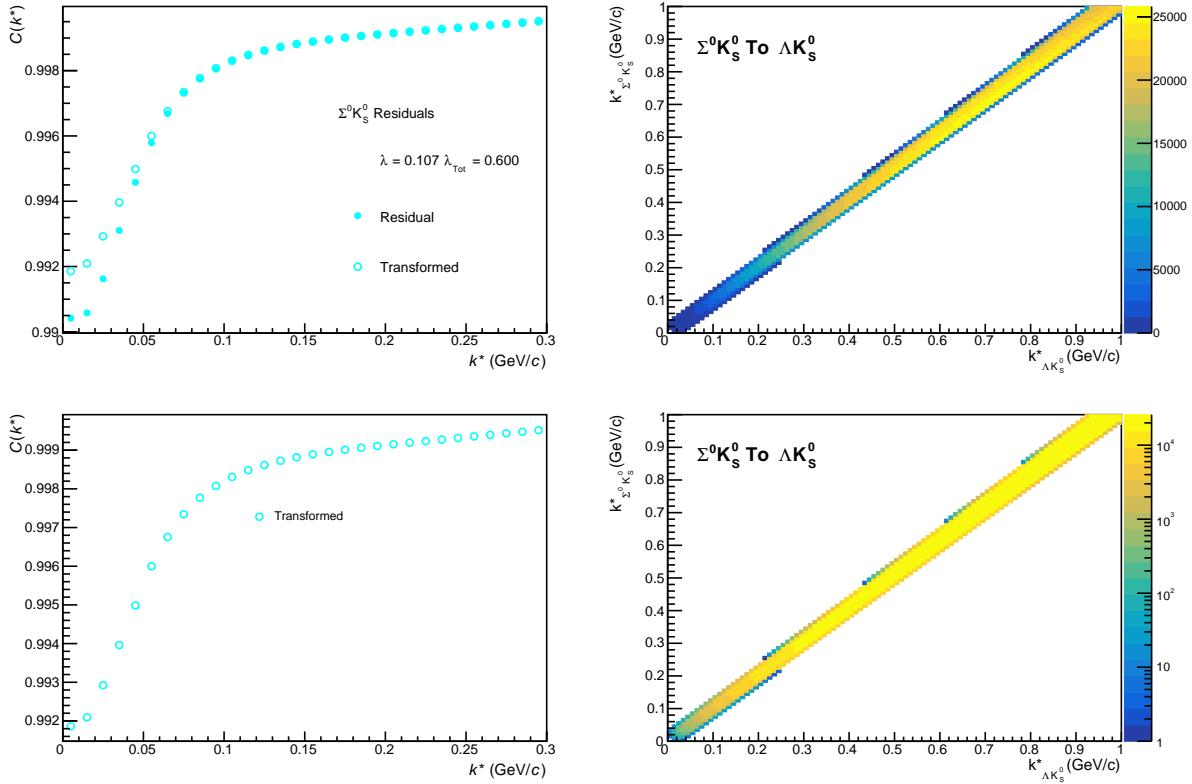


Fig. 87: Residuals: $\Sigma^0 K_S^0$ to ΛK_S^0 (0-10% Centrality)

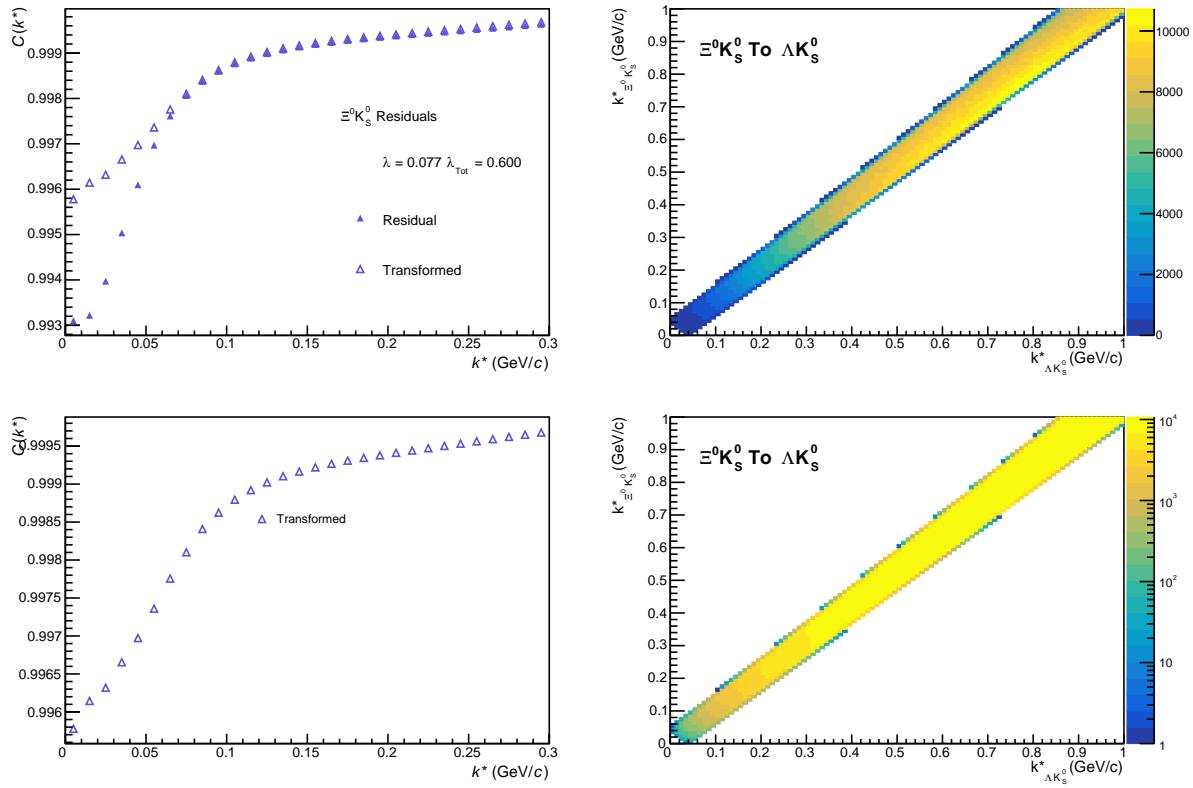


Fig. 88: Residuals: $\Xi^0 K_S^0$ to ΛK_S^0 (0-10% Centrality)

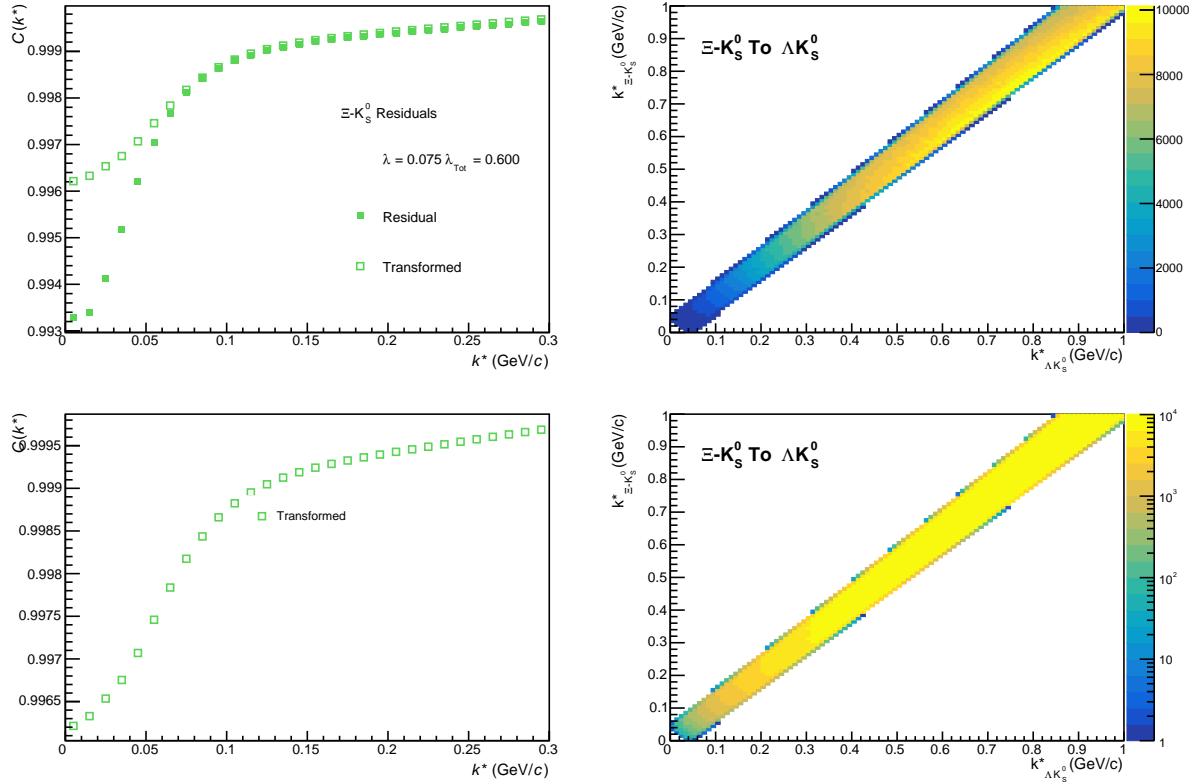
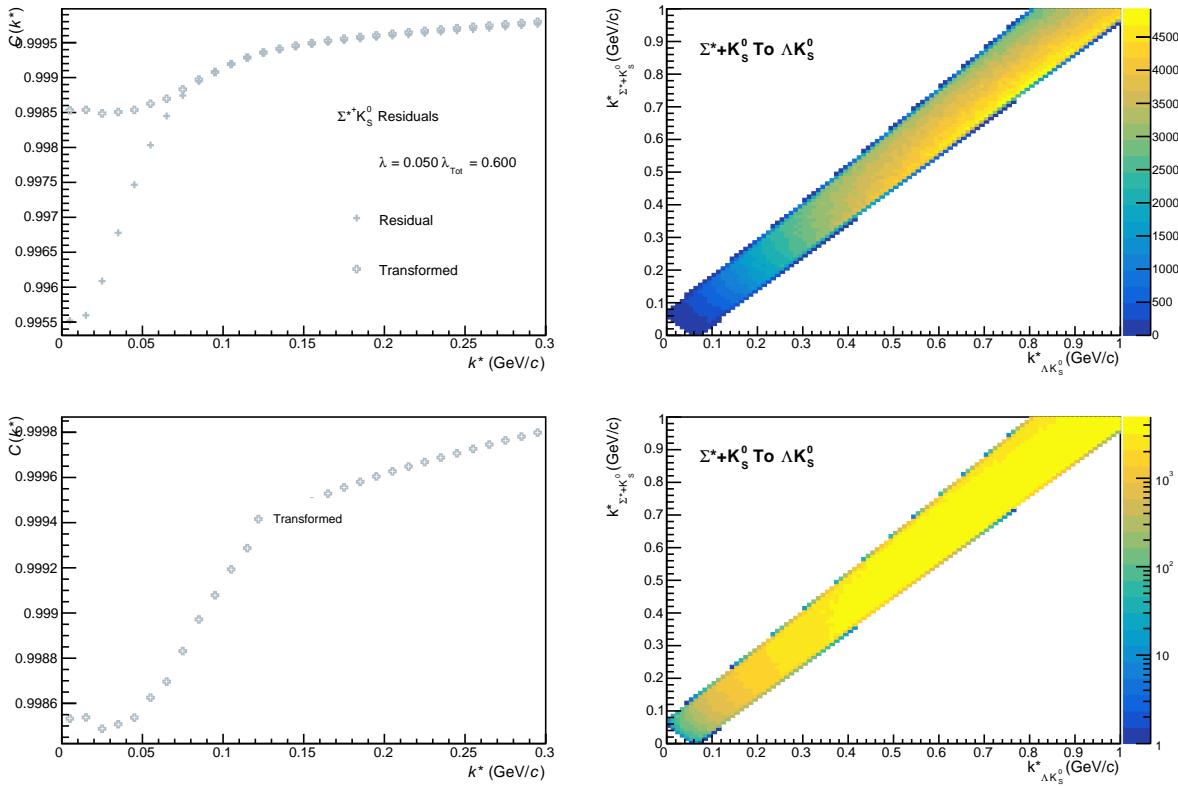
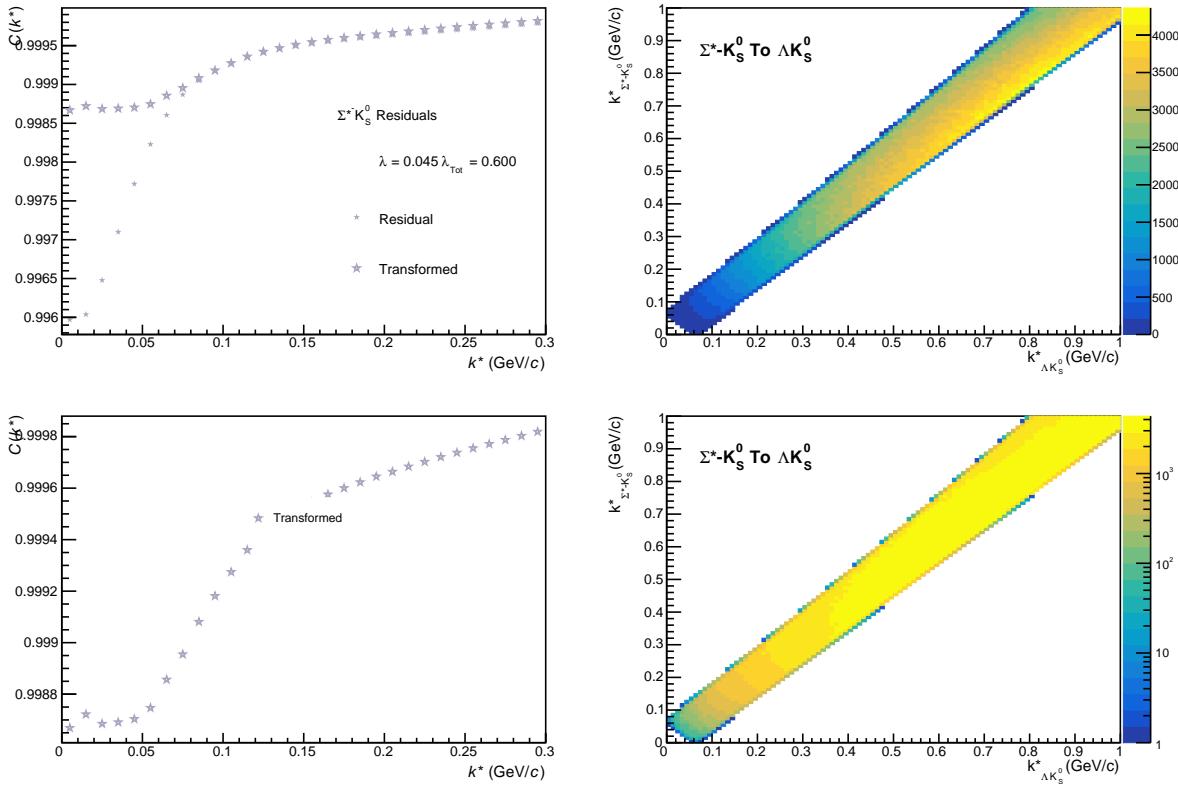


Fig. 89: Residuals: $\Xi^- K_S^0$ to ΛK_S^0 (0-10% Centrality)


Fig. 90: Residuals: $\Sigma^* + K_S^0 \rightarrow \Lambda K_S^0$ (0-10% Centrality)

Fig. 91: Residuals: $\Sigma^* - K_S^0 \rightarrow \Lambda K_S^0$ (0-10% Centrality)

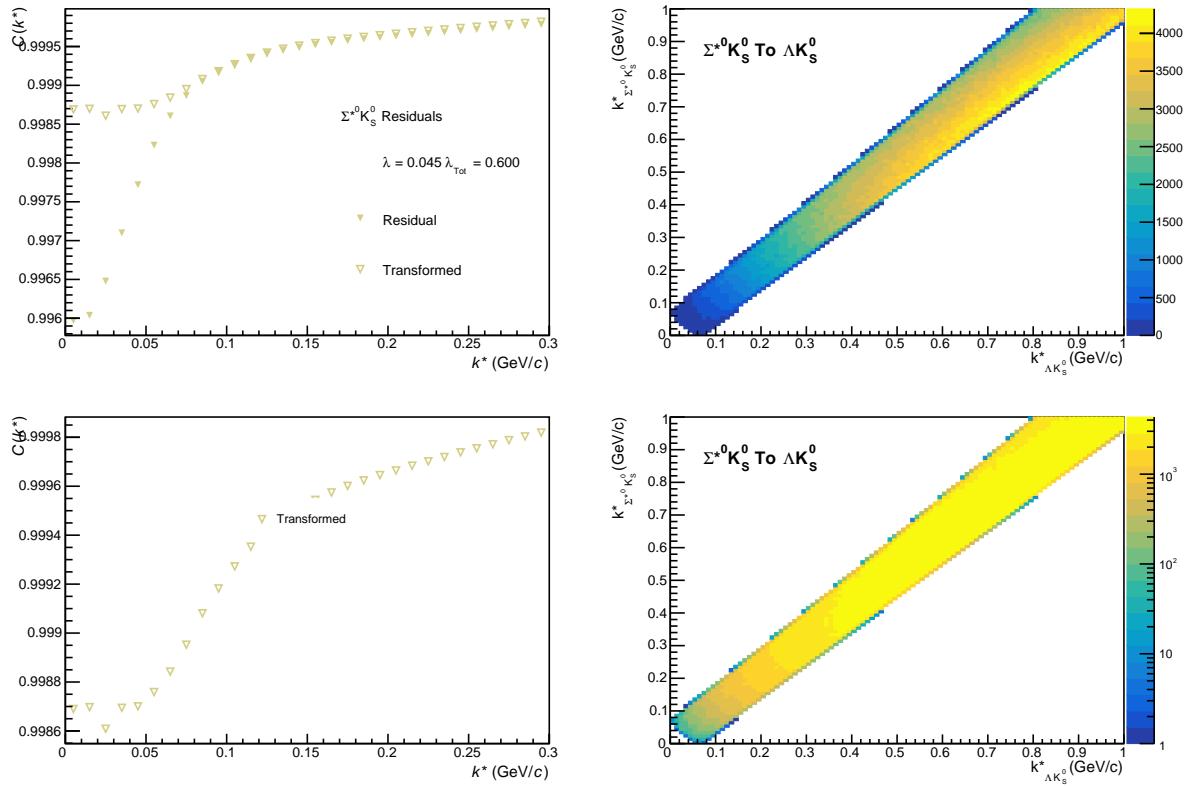


Fig. 92: Residuals: $\Sigma^{*0} K_S^0$ to ΛK_S^0 (0-10% Centrality)

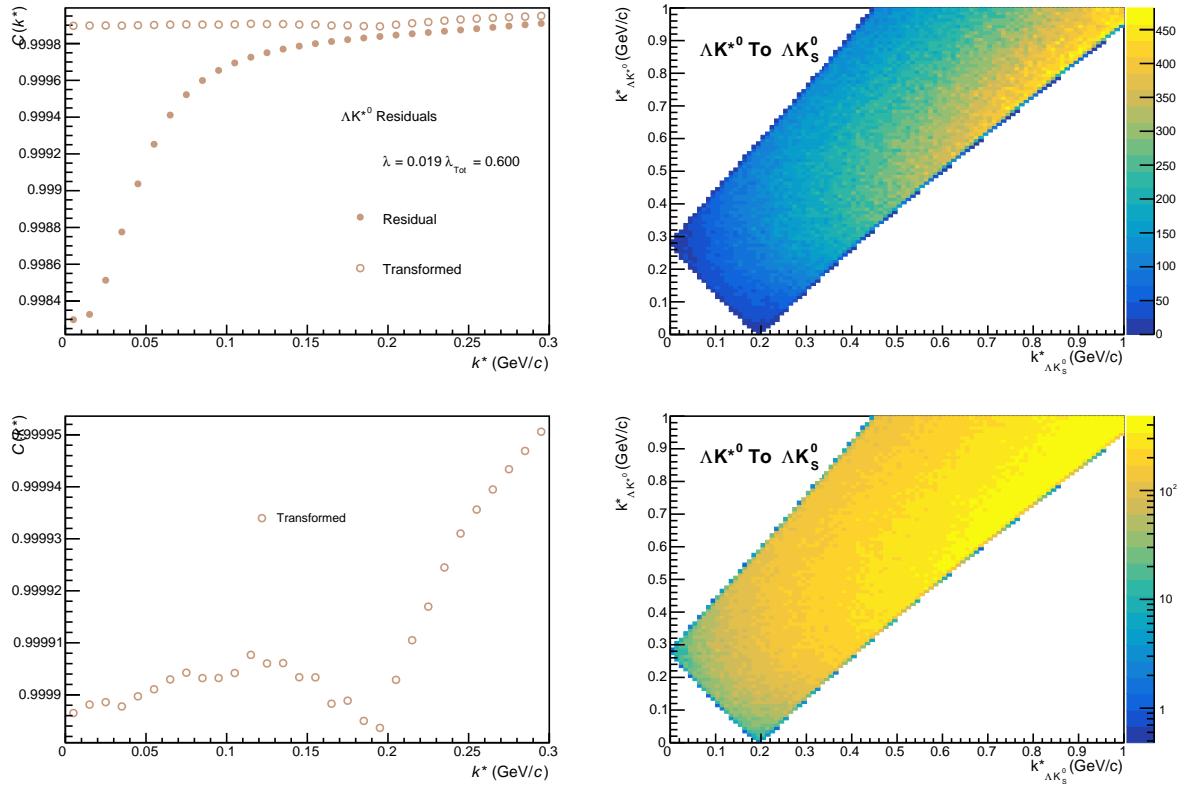
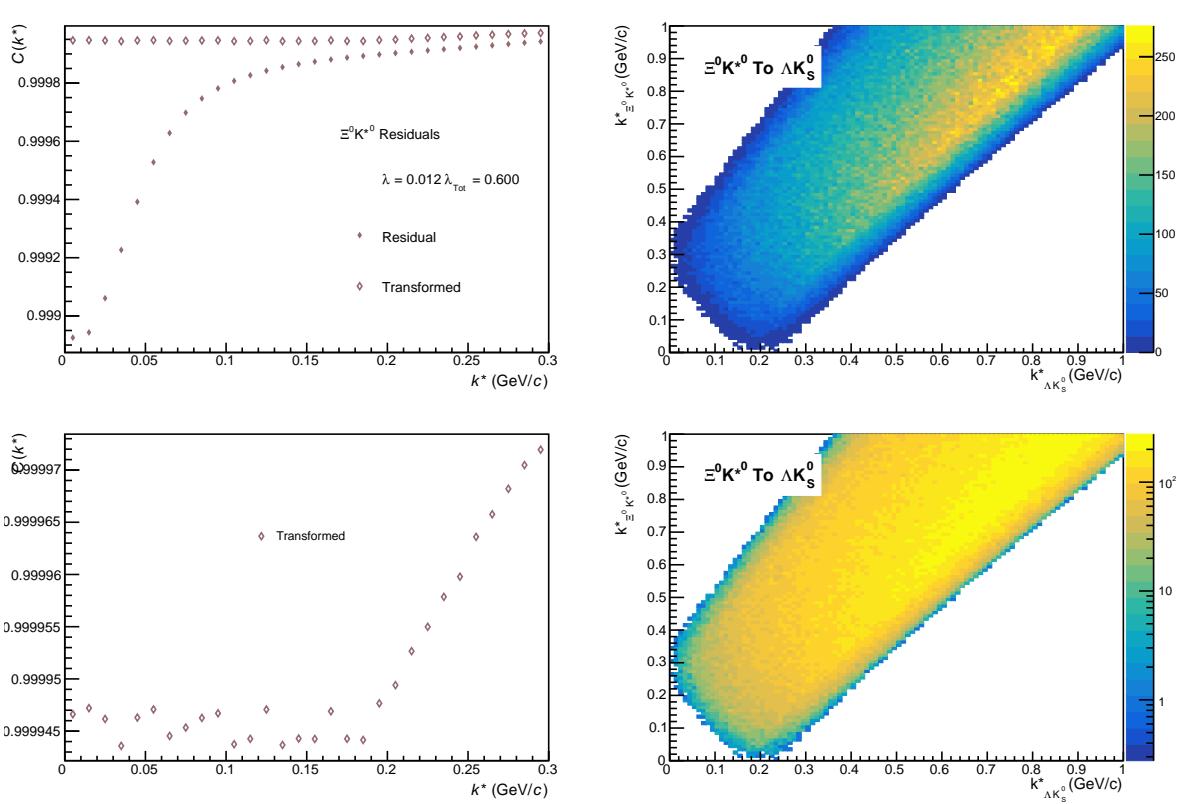
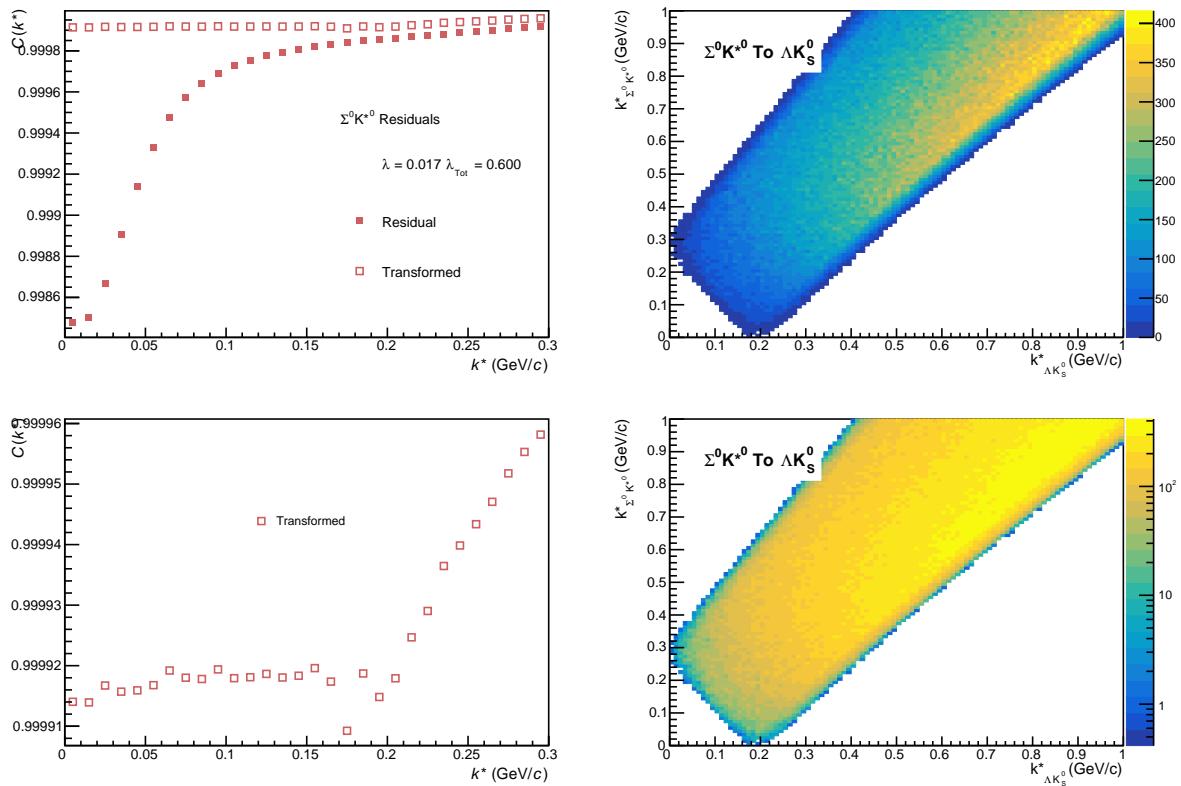


Fig. 93: Residuals: ΔK^{*0} to ΛK_S^0 (0-10% Centrality)



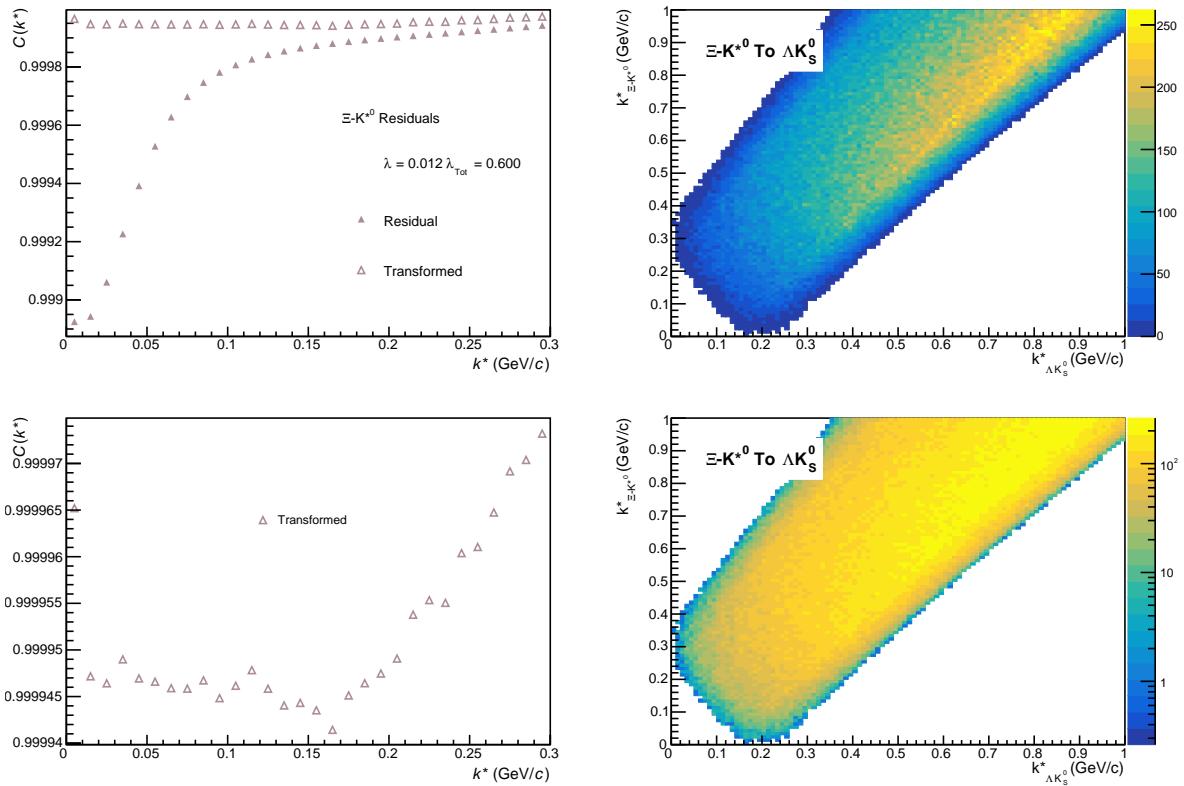


Fig. 96: Residuals: $\Xi^- K^{*0}$ to ΛK_s^0 (0-10% Centrality)

951 9.2 Spherical Harmonics

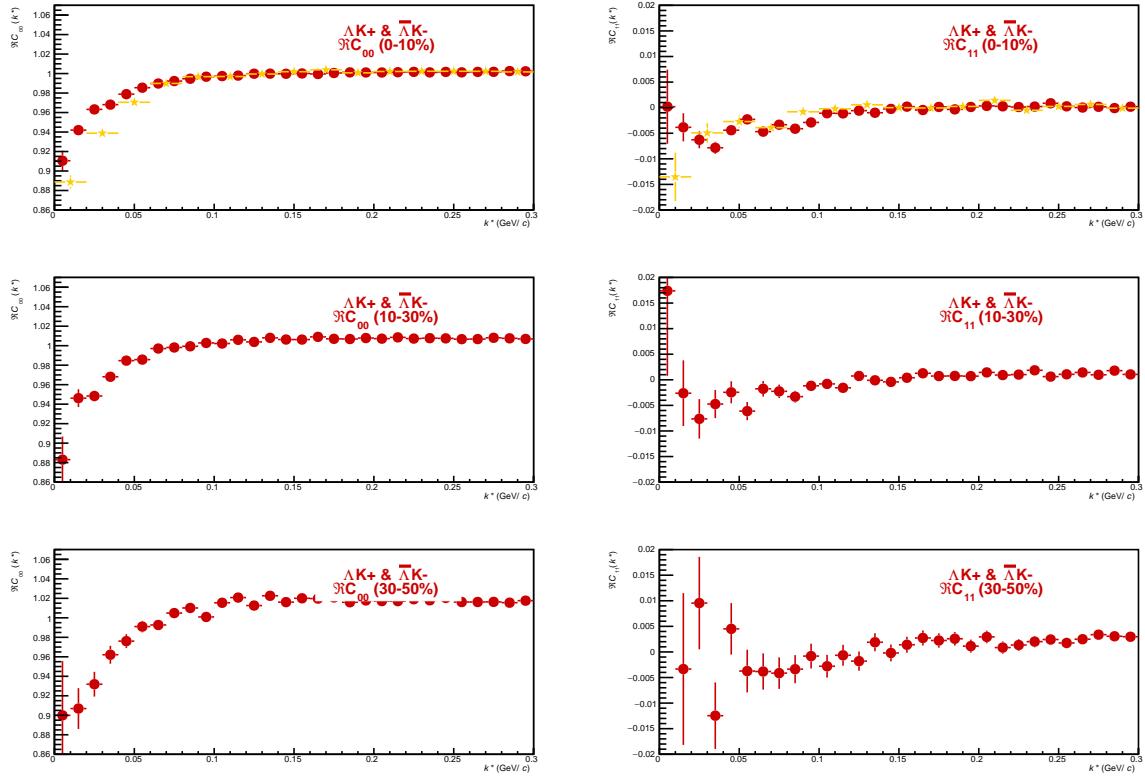


Fig. 97: C_{00} (left) and $\Im C_{11}$ (right) components of a spherical harmonic decomposition of the ΛK^+ correlation function for the 0-10% (top), 10-30% (middle), and 30-50% (bottom) centrality bins

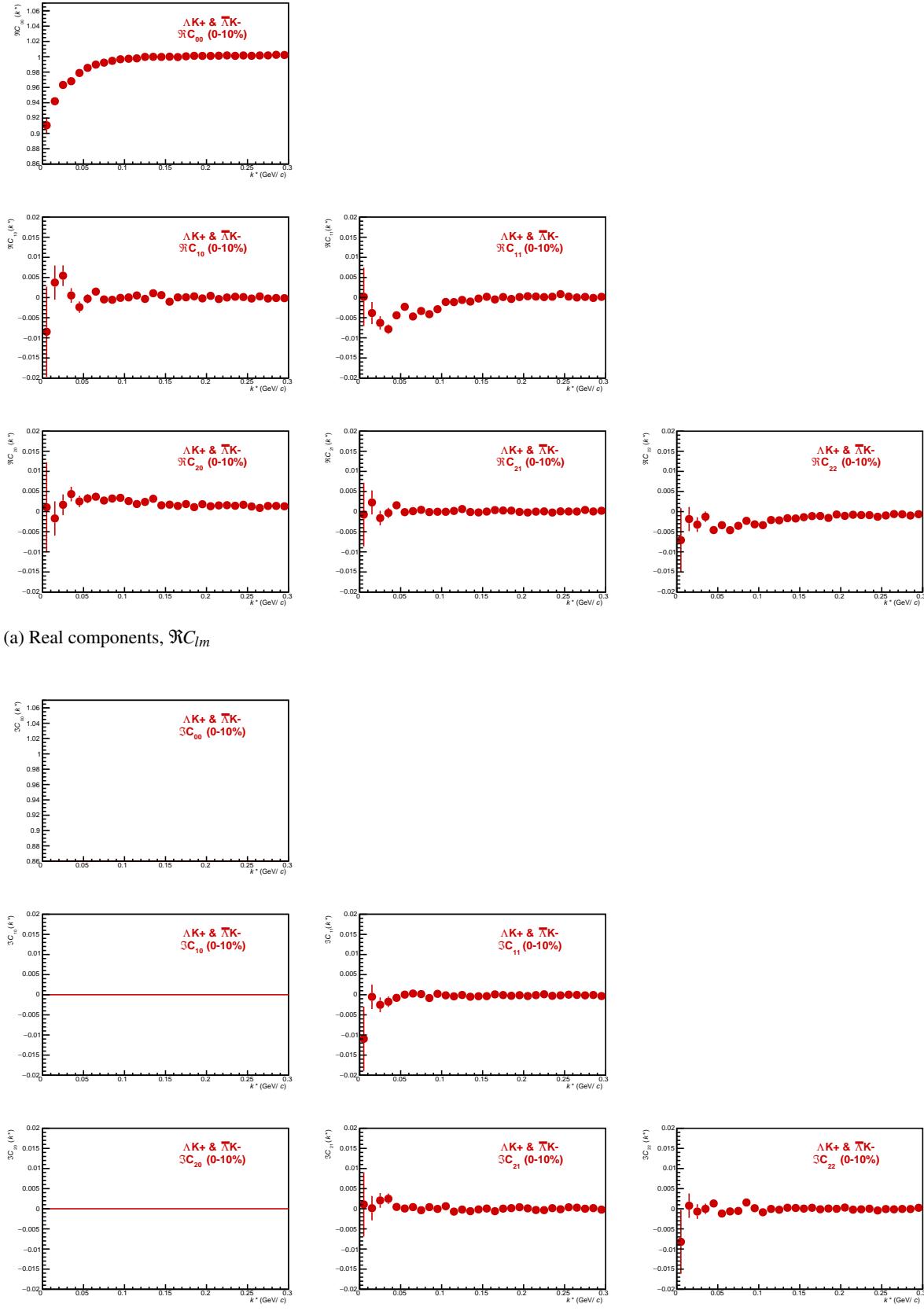
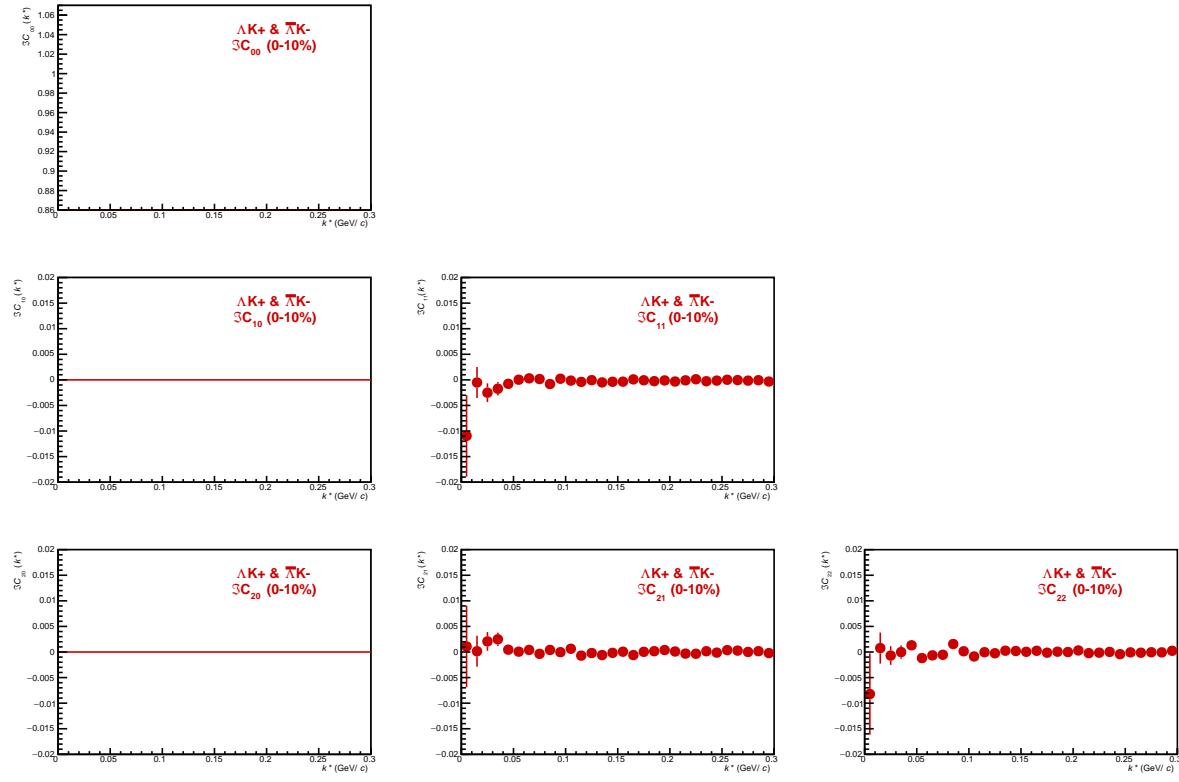
(a) Real components, $\Re C_{lm}$ (b) Imaginary components, $\Im C_{lm}$

Fig. 98: First six components ($C_{00}, C_{10}, C_{11}, C_{20}, C_{21}, C_{22}$) of the spherical harmonic decomposition of the ΛK^+ correlation function for the 0-10% centrality bin. Note, $\Im C_{00}$, $\Im C_{10}$, and $\Im C_{20}$ are zero by definition.

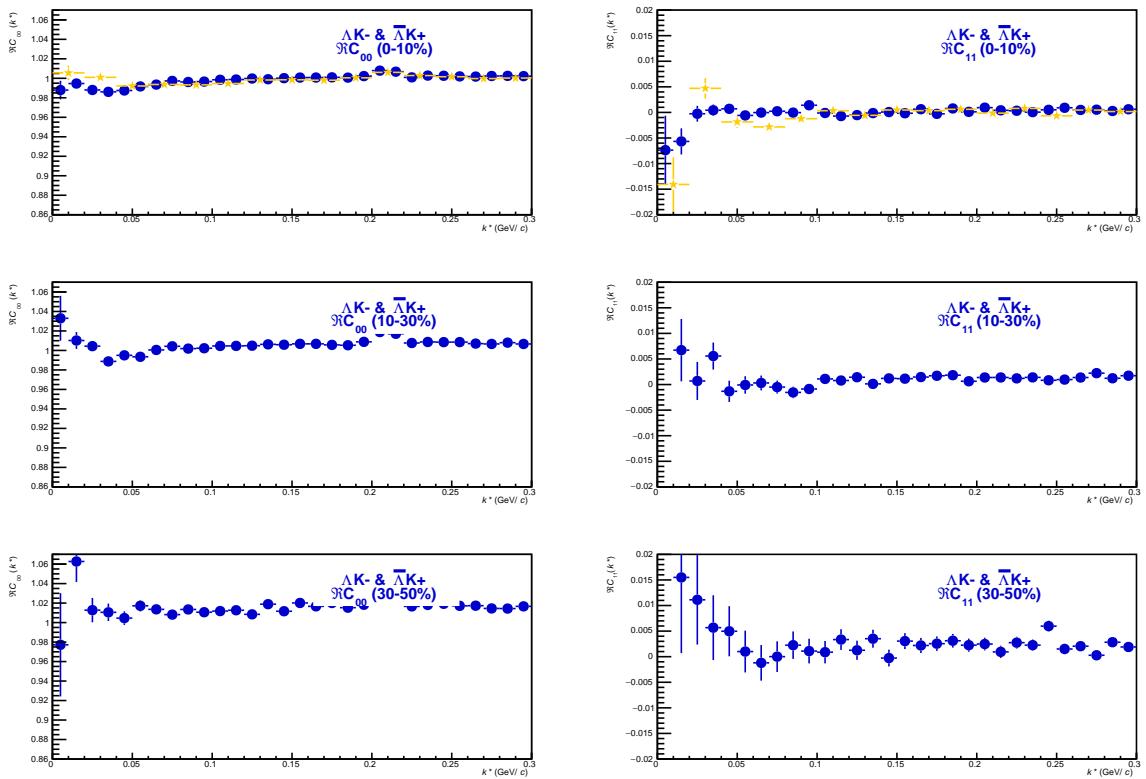


Fig. 99: C_{00} (left) and $\Im C_{11}$ (right) components of a spherical harmonic decomposition of the ΛK^- correlation function for the 0-10% (top), 10-30% (middle), and 30-50% (bottom) centrality bins

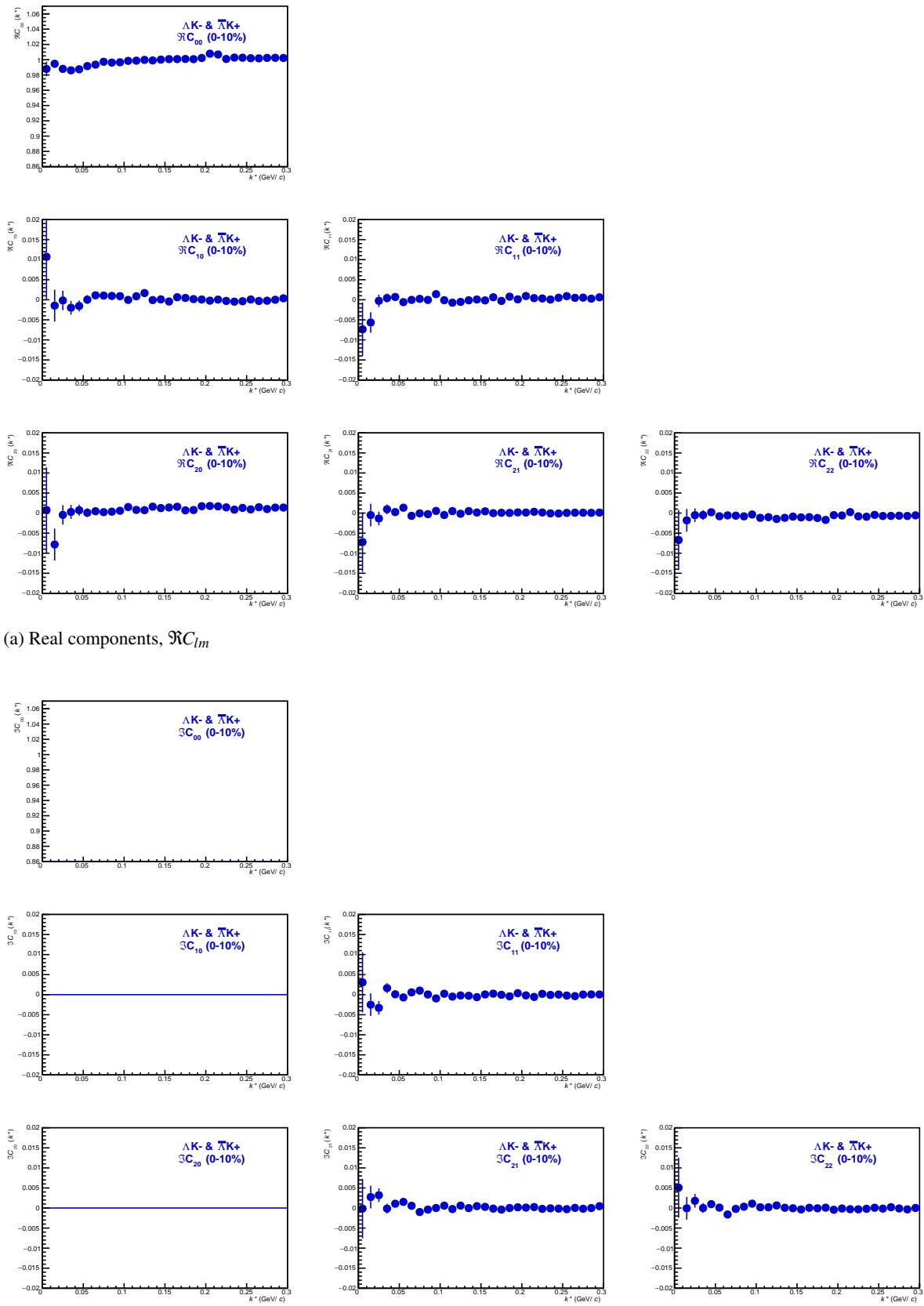


Fig. 100: First six components ($C_{00}, C_{10}, C_{11}, C_{20}, C_{21}, C_{22}$) of the spherical harmonic decomposition of the ΛK^- correlation function for the 0-10% centrality bin. Note, $\Im C_{00}$, $\Im C_{10}$, and $\Im C_{20}$ are zero by definition.

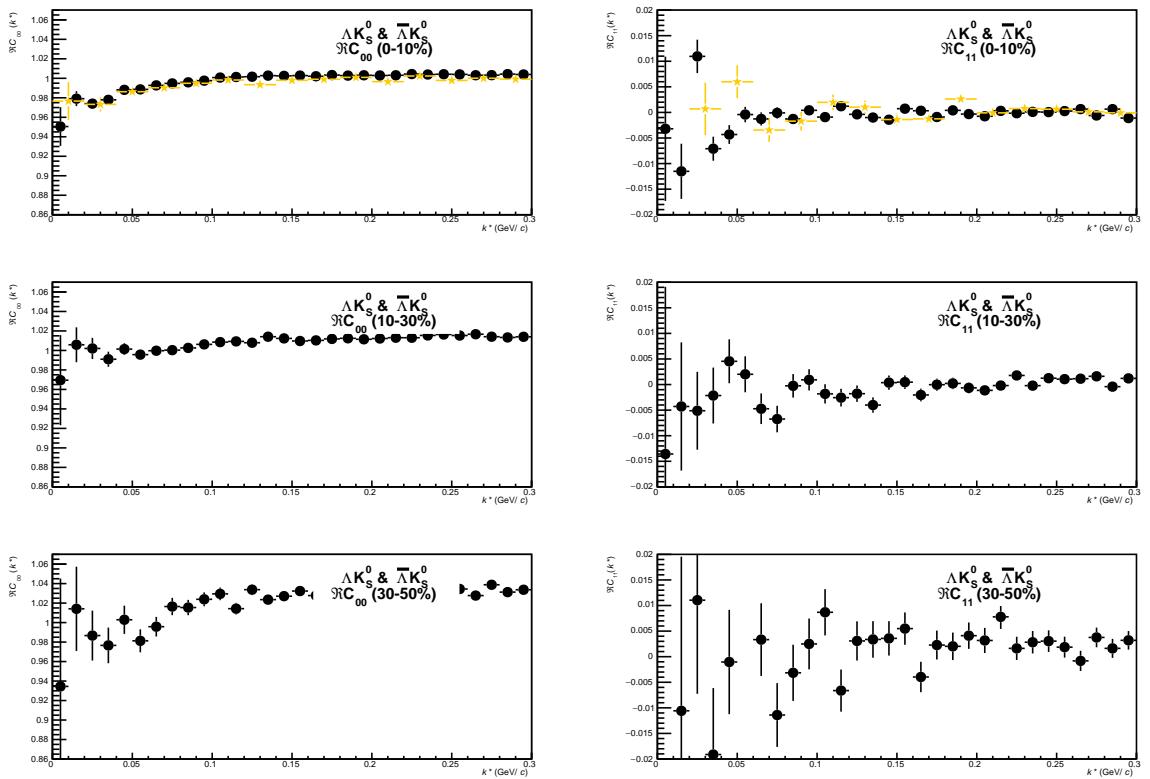


Fig. 101: C_{00} (left) and $\Im C_{11}$ (right) components of a spherical harmonic decomposition of the ΛK_s^0 correlation function for the 0-10% (top), 10-30% (middle), and 30-50% (bottom) centrality bins

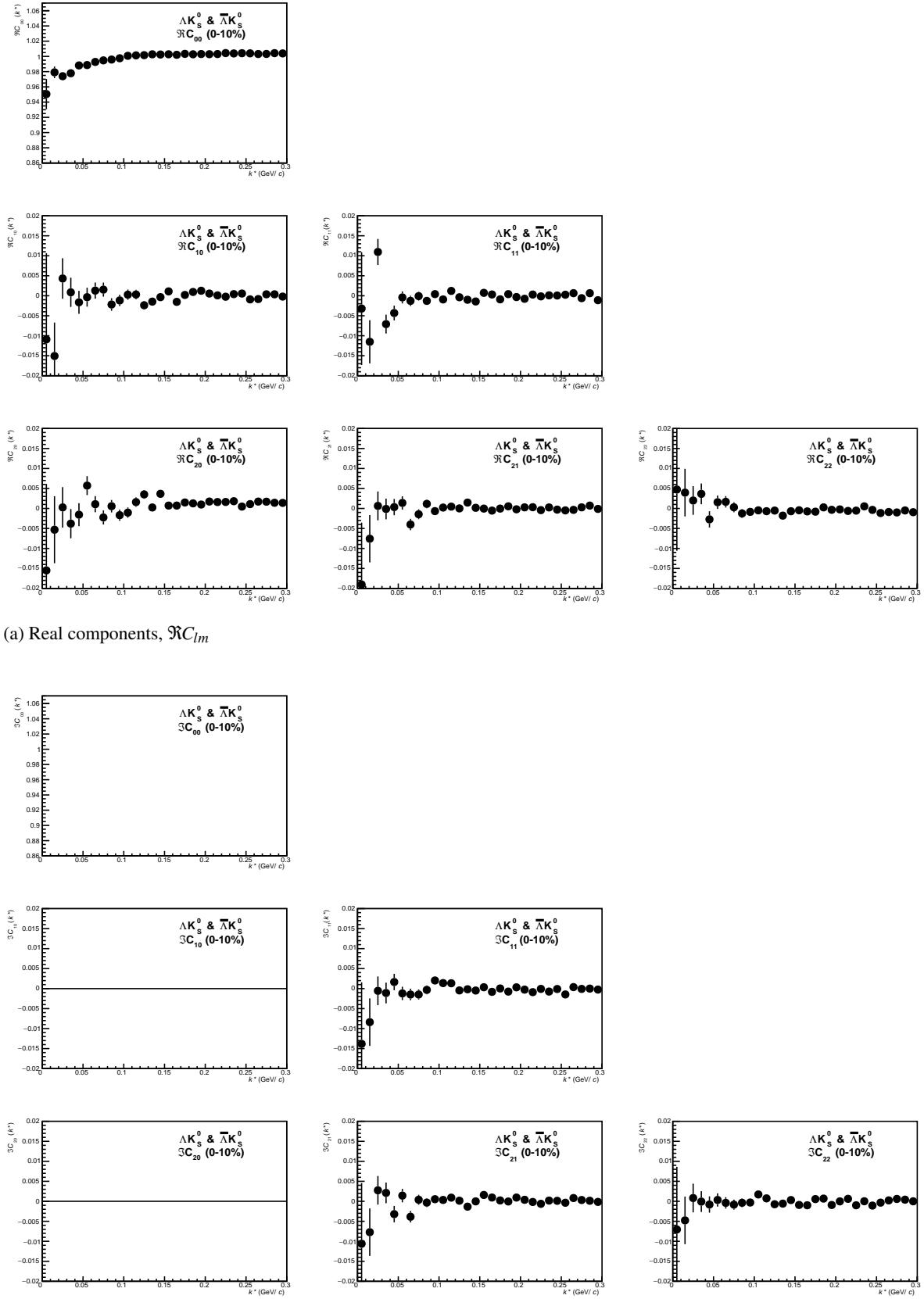


Fig. 102: First six components ($C_{00}, C_{10}, C_{11}, C_{20}, C_{21}, C_{22}$) of the spherical harmonic decomposition of the ΛK_s^0 correlation function for the 0-10% centrality bin. Note, $\Im C_{00}$, $\Im C_{10}$, and $\Im C_{20}$ are zero by definition.

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