

# Homework # 2

PHY-905-003, Computational Astrophysics and Astrostatistics  
Spring 2017

**This assignment is due the evening of Friday March 17, 2017.** Turn in all materials via GitHub. Include your code, plots, and answers to any questions asked in your assignment. Your code must (1) be easily readable, with good use of whitespace, clear variable names, and adequate commenting (which documents design and purpose, not mechanics) and (2) use functions to break up code into logical components. The solutions to individual problems should be saved in separate, clearly-named source files or Jupyter notebooks. Plots should have easily readable and logical axis labels and titles, and the source code and data used to generate the plots should be included. Questions should be answered in the file `ANSWERS.md` or in a L<sup>A</sup>T<sub>E</sub>X-created PDF document of a similar name (e.g., `ANSWERS.pdf`).

**Part 1:** The homework repository contains a file named `time_varying_object.txt`, which contains observations of the flux measured from a variable compact object. There are several regularly-spaced objects per day over the course of the year, and the flux magnitude is in arbitrary units. Perform a Fourier transform of this data and determine the **frequency, amplitude, and phase** of all cycles observed in this object. (Note: consult your lecture notes for information on how to determine the phase.) Then, do/answer the following:

1. Make a plot of the amplitude as a function of period, using appropriate limits in both the x and y dimensions of the plot. In other words, I'm not interested in the large range of frequencies/periods where nothing is happening - zoom in on the good stuff.
2. What are the periods (in days), relative amplitudes, and relative phases of the cycles observed in this object?
3. Create a fitting function composed of a sum of sinusoidal functions with the appropriate amplitudes, frequencies, and phases, and plot it for the duration of the observation period *on a plot that includes the original data*. Does this fitting function look like it adequately represents the dataset, or is something missing? If so, speculate about what that might be, and why you might be missing it in your fitting function.

**Part 2:** The [Lax-Wendroff method](#) for solving hyperbolic PDEs can be derived by Taylor expansion in a similar manner to the methods done in the books and in class. By keeping terms to  $\mathcal{O}(\Delta t^2)$  and using centered, second-order spatial derivatives, the update equation for the linear advection equation ( $a_t + ua_x = 0$ ) becomes:

$$a_i^{n+1} = a_i^n - \frac{\Delta t}{2\Delta x} u [a_{i+1}^n - a_{i-1}^n] + \frac{\Delta t^2}{2\Delta x^2} u [a_{i+1}^n - 2a_i^n + a_{i-1}^n] \quad (1)$$

The first two terms on the right hand side are just the FTCS update (which was shown to be unstable in a pre-class assignment) and the second term is a diffusion term. This method is stable - the explicit diffusion term counteracts the numerical diffusion that made the method unstable.

Implement this method on a cell-centered finite-volume grid. For the initial conditions, choose a Gaussian:

$$a(x, t=0) = e^{-(x-0.5)^2/0.1^2} \quad (2)$$

with a velocity of  $u = 1$  in a domain  $x = [0, 1]$  and using periodic boundary conditions. Evolve it for one period,  $T = 1/u = 1$ , and demonstrate that you have implemented it correctly by verifying that you get the same solution for  $u = -1$  and  $u = 1$  (and show this in a plot).

We are going to measure the convergence of this method as you change the resolution of the grid. Testing convergence requires that you change your timestep as you change your grid resolution, which is handled automatically if you use the same Courant number at each resolution. To get a single number from the spatially discretized solution that represents the numerical error, we're going to use the L2 norm (as defined in Section 1.2.4 of Zingale's lecture notes). We will define our error as:

$$\epsilon_{\Delta x} = \|a(x, t = T) - a(x, T = 0)\|_2 \quad (3)$$

In other words, this is the pointwise RMS error of the final solution compared to the initial solution, normalized by  $\Delta x$  (so, scaling the error with the number of grid points). **Plot  $\epsilon_{\Delta x}$  vs.  $\Delta x$  for several grid resolutions on a log-log plot and estimate the convergence rate from the slope of the line.** Note: if you don't see a convergence rate of  $\mathcal{O}(\Delta t^2)$ , you almost certainly have a bug somewhere!

**Part 3:** In this problem, we're going to be using an iterative method to calculate the gravitational potential of two 3D physical setups. The file `elliptical_equation_arrays.ipynb` contains code to generate two three-dimensional numpy arrays, `single_star_array` and `many_star_array`, which contain a single  $1 M_\odot$  star at the center of the grid and a large number of stars scattered randomly about the grid, respectively. The grid has a size that is determined by a variable in the file, and is assumed to be a cubic volume that is 1 parsec across. Do the following:

1. Write a function that solves Poisson's equation,  $\nabla^2 \phi = 4\pi G\rho$ , using your choice of either Jacobi or Gauss-Seidel iteration. This function should take as inputs a user-provided cubic 3D grid of densities of arbitrary grid size (in CGS units) and a desired level of error (to determine when convergence has been achieved), as well as any other parameters that you deem necessary. The function should return an array containing the gravitational potential in CGS units, the number of iterations required to achieve convergence, and the actual error (calculated as the L2 norm of the residual for the grid, excluding the ghost zones). Assume that the grid is cell-centered (finite volume) and using Neumann boundary conditions that ensure that the derivative of the slope across the boundary of the domain (i.e., between the ghost zone and the first "real" zone in your 3D array) is identical to the slope of the cells immediately inside the boundary. This is effectively an "isolated" boundary condition. How might you implement a periodic boundary condition in this circumstance?
2. Test the function that you have written using the `single_star_array`. Plot a slice of the gravitational potential through the center of the array (i.e., at the same 2D slice of cells containing the star) and verify that you get a sensible answer – in other words, that the potential falls off as  $1/r$  and that it has the correct units. You should do this by making a line plot that includes calculated values of the potential from the cell containing the star to the edge of the domain, as well as the analytic solution that you would expect. You may also want to make an image of a slice of the potential array (Hint: you may wish to take the log of the absolute magnitude of a slice of the array to be able to see dynamic range!) How many iterations does it take for your solution to converge to an accuracy in the residual of  $10^{-2}$ ,  $10^{-3}$ , and  $10^{-4}$ , respectively?
3. Now that you've verified that your code works correctly in a simple situation, calculate the potential for `many_star_array`. Plot a 2D image of the potential through several different slices of the domain (again using the log of the absolute value so you can see dynamic range). Does the result look sensible? Does the number of iterations required to achieve comparable accuracy change between this part of the problem and the single-star version?
4. Congratulations! You've written a gravity solver. If you wanted to turn this into a simple simulation of the evolution of a cluster of stars, how would you do that? You don't need to write the actual code, just outline the steps you'd take in `ANSWERS.md`.

**Part 4:** In this problem we're going to make a model of the conduction of heat through the Earth's crust using a simple 1D model. This problem has a time-dependent boundary condition – the heating of the surface of the sun throughout the seasons.

The mean temperature at the Earth's surface can be approximated as:

$$T_0(t) = A + B \sin\left(\frac{2\pi t}{\tau}\right) \quad (4)$$

where  $\tau = 365$  days,  $A = 10^\circ \text{ C}$  is the average surface temperature and  $B = 12^\circ \text{ C}$  captures the seasonal variation.

At a depth of 20 m the temperature of the crust is relatively constant, and is set by the heat flux coming from the interior of the Earth. We assume it to be  $11^\circ \text{ C}$ . The thermal diffusivity of the Earth's crust is roughly  $D = 0.1 \text{ m}^2 \text{ day}^{-1}$ , and we assume it to be constant throughout the domain. As an initial condition, take the temperature to be  $T = 10^\circ \text{ C}$  throughout the interior zones of the domain.

Solve the diffusion equation on a domain that goes from the surface to a depth of 20 m using the [Crank-Nicolson method](#) (as discussed in class and in Zingale's book). Use Dirichlet boundary conditions with the time-dependent values given above. Evolve the simulation for 10 years – after a few years, the initial guess at the temperature will have relaxed out (i.e., will have disappeared and been replaced by the “true” values). Experiment with the number of zones and the timestep to get a result that is converged and runs in a sensible amount of time.

For the final year, plot  $T(x)$  as a function of depth at 4 different time zones that correspond to the four seasons. By how much does the temperature of the ground vary at 1 m, 5 m, 10 m, and 20 m below ground throughout the year?