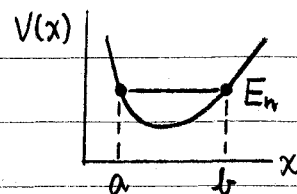


Exam Questions for Phys. 506 (Hom Test)

12/29/70

- 12/29/70 ① Consider a particle bound in an arbitrary attractive 1D potential $V(x)$ as shown. Show that the spacing between adjacent quantized energy levels E_n is given, in WKB approximation, by $\Delta E_n \approx \hbar \omega_n$, where ω_n is the natural vibration frequency of the n^{th} level. (Hint: differentiate the WKB energy quantization condition with respect to n)."



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$$\text{Level } E_n \text{ found from: } \int_a^b [2m(E_n - V(x))]^{\frac{1}{2}} dx = (n + \frac{1}{2})\pi\hbar$$

Differentiating w.r.t. n gives

$$\pi\hbar = \int_a^b \frac{1}{2} [2m(E_n - V(x))]^{-\frac{1}{2}} 2m \left(\frac{dE_n}{dn} \right) dx = \left(\frac{dE_n}{dn} \right) \int_a^b \frac{dx}{p_n(x)/m}$$

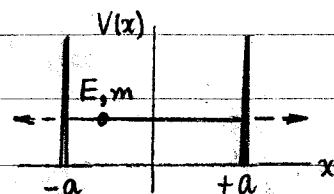
where: $p_n(x) = [2m(E_n - V(x))]^{\frac{1}{2}}$ is momentum in n^{th} level

$$\text{But natural period is: } T = \frac{2\pi}{\omega} = 2 \int_a^b \frac{dx}{p(x)/m}$$

$$\therefore \int_a^b \frac{dx}{p_n(x)/m} = \frac{\pi}{\omega_n}, \quad \omega_n = \text{natural freq. in } n^{\text{th}} \text{ level}$$

$$\text{This } \Rightarrow \frac{dE_n}{dn} = \hbar \omega_n, \text{ or } \Delta E_n \approx \hbar \omega_n \text{ for adjacent levels.}$$

- 1/5/71 ② "A particle of mass m and energy E is trapped in a 1D box of length $2a$. The walls of the box (at $\pm a$) may be represented by δ -funs of strength C , i.e. the potential is: $V(x) = C [\delta(x+a) + \delta(x-a)]$. Estimate the lifetime of the particle in the box; i.e., how long before it gets out?"



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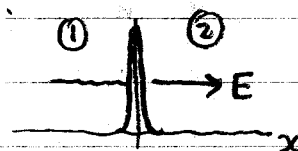
Decay rate : $\Gamma = \left(\frac{1}{\tau/2}\right) T$ $\left\{ \begin{array}{l} \tau = \text{natural period inside box} \\ T = \text{trans. coeff. at one barrier} \end{array} \right.$

Particle is free inside box $\Rightarrow \tau = \frac{4a}{v}$, $v = \sqrt{2E/m}$

$\therefore \left(\frac{1}{\tau/2}\right) = \frac{1}{2a} \sqrt{\frac{2E}{m}}$

Must now calculate T for δ -fun barrier...

$V = C\delta(x)$. Let $\hbar k = \sqrt{2mE}$.



$\psi_1 = e^{+ikx} + Ae^{-ikx}$, $\psi_2 = Be^{+ikx}$, Want $T = |B|^2$

ψ continuous at $x=0$: $1+A=B \leftarrow \textcircled{\text{I}}$

ψ' discontinuous at $x=0$...

$\psi_2'(0+) - \psi_1'(0-) = \frac{2mC}{\hbar^2} \psi_2(0)$

$\Rightarrow ik[B - (1+A)] = \frac{2mC}{\hbar^2} B \Rightarrow 1-A = \left(1 - \frac{2mC}{ik\hbar^2}\right) B \leftarrow \textcircled{\text{II}}$

Add $\textcircled{\text{I}}$ & $\textcircled{\text{II}} \Rightarrow 2 = \left(2 + i \frac{2mC}{\hbar^2 k}\right) B \Rightarrow B = 1 / \left(1 + i \frac{mC}{\hbar^2 k}\right)$

$\therefore T = |B|^2 = 1 / \left(1 + \frac{m^2 C^2}{\hbar^2 (\hbar k)^2}\right) = 1 / \left(1 + \frac{mC^2}{2E\hbar^2}\right)$

$\Rightarrow \Gamma = \frac{1}{2a} \sqrt{\frac{2E}{m}} / \left(1 + \frac{mC^2}{2E\hbar^2}\right)$

Lifetime is : $\Delta t = \frac{1}{\Gamma} = 2a \sqrt{\frac{m}{2E}} \left(1 + \frac{mC^2}{2E\hbar^2}\right)$

1/27/71 ③ "A QM system in state $\psi(x)$ at time $t=0$ is subjected to an interaction H which generates two discrete eigenstates ϕ_n and eigenenergies E_n , with $E_2 - E_1 = \hbar\Omega$. The energy spectrum for H is therefore discrete, with values

$$W_n = \left| \int \phi_n^*(x) \psi(x) dx \right|^2, \quad n = 1 \neq 2.$$

Assume $\sum_{n=1}^2 W_n = 1$ for convenience. Derive an expression for the probability of finding the state $\psi(x)$ at time $t > 0$. What is the oscillation period between times of maximum probability?"

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Desired probability is : $P(t) = \left| \sum_E W(E) e^{-\frac{i}{\hbar} E t} \right|^2$

$$\text{or } P(t) = \left| W_1 e^{-\frac{i}{\hbar} E_1 t} + W_2 e^{-\frac{i}{\hbar} E_2 t} \right|^2$$

$$= \left| e^{-\frac{i}{\hbar} E_1 t} (W_1 + W_2 e^{-i\Omega t}) \right|^2$$

$$= (W_1 + W_2 e^{+i\Omega t})(W_1 + W_2 e^{-i\Omega t})$$

$$= W_1^2 + W_2^2 + W_1 W_2 (e^{i\Omega t} + e^{-i\Omega t})$$

$$= W_1^2 + W_2^2 + 2W_1 W_2 \cos \Omega t$$

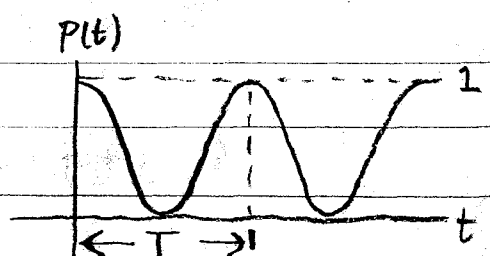
$$\nearrow 2 \sin^2 \frac{1}{2} \Omega t$$

$$= W_1^2 + W_2^2 + 2W_1 W_2 - 2W_1 W_2 (1 - \cos \Omega t)$$

$$= \underbrace{(W_1 + W_2)^2}_{=1} - 4W_1 W_2 \sin^2 \frac{1}{2} \Omega t$$

$$\therefore P(t) = 1 - 4W_1 W_2 \sin^2 \frac{1}{2} \Omega t$$

$$\text{Period } T : \frac{1}{2} \Omega T = \pi \Rightarrow T = 2\pi / \Omega$$



2/3/71 ④ "Start from the definition of the S-matrix in the form

$$\Psi_\alpha(x', t') = \sum_\beta S_{\beta\alpha} \Phi_\beta(x', t'),$$

which describes the evolution of a free particle state $\Phi_\alpha(x, t)$ in the distant past to the state $\Psi_\alpha(x', t')$ in the distant future. Suppose the Φ_β are orthonormal, and that the total interaction is at all times Hermitian. Then the normalization and orthogonalization of the Ψ_α must be time-independent. Use this fact to show that the S-matrix is unitary, i.e.

$$S^\dagger S = 1, \quad \text{or} \quad (S^\dagger S)_{ij} = \sum_\beta S_{i\beta}^\dagger S_{\beta j} = \sum_\beta S_{\beta i}^* S_{\beta j} = \delta_{ij}.$$

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Since the Ψ normalization is t' indpt, we can write

$$\begin{aligned} \int dx' \Psi_i^*(x', t') \Psi_j(x', t') &= \lim_{t' \rightarrow -\infty} \int dx' \Psi_i^*(x', t') \Psi_j(x', t') \\ &= \int dx' \Phi_i^*(x', t') \Phi_j(x', t') = \delta_{ij} \end{aligned}$$

Plugging in the expansions for Ψ on the LHS, we get

$$\sum_{\beta, \gamma} S_{\beta i}^* S_{\gamma j} \underbrace{\int dx' \Phi_\beta^*(x', t') \Phi_\gamma(x', t')}_{\delta_{\beta\gamma}} = \delta_{ij}$$

$$\text{or} \quad \sum_\beta S_{\beta i}^* S_{\beta j} = \delta_{ij} \quad \underline{\underline{\text{QED}}}$$

An alternate, somewhat more physical, proof is the following...

$$\psi_\alpha(x', t') = \sum_\beta S_{\beta\alpha} \phi_\beta(x', t') \Rightarrow S_{\beta\alpha} = \int dx' \phi_\beta^*(x', t') \psi_\alpha(x', t')$$

Plugging the matrix element directly into the sum, we get

$$\begin{aligned} \sum_\beta S_{\beta i}^* S_{\beta j} &= \sum_\beta \left(\int dx' \phi_\beta^*(x', t') \psi_i(x', t') \right)^* \left(\int dx \phi_\beta^*(x, t) \psi_j(x, t) \right) \\ &= \int dx' \psi_i^*(x', t') \int dx \left[\sum_\beta \phi_\beta^*(x, t) \phi_\beta(x', t') \right] \psi_j(x, t) \end{aligned}$$

But the $[] = i G_0(x', t'; x, t)$, the free particle propagator. So

$$\begin{aligned} \sum_\beta S_{\beta i}^* S_{\beta j} &= \int dx' \psi_i^*(x', t') \underbrace{i \int dx G_0(x', t'; x, t) \psi_j(x, t)}_{= \psi_j(x', t'), \text{ since } G_0 \Rightarrow \text{propagation without any interactions to distort } \psi_j.} \\ &= \psi_j(x', t') \end{aligned}$$

$$\begin{aligned} \therefore \sum_\beta S_{\beta i}^* S_{\beta j} &= \int dx' \psi_i^*(x', t') \psi_j(x', t') \leftarrow \text{this is the orthonormalization} \\ &= \delta_{ij}, \quad \underline{\underline{\text{QED}}} \end{aligned}$$

integral for ψ , which can be assumed to be δ_{ij} as $t' \rightarrow -\infty$.