

Maxwell Equations: t-Variation & Conservation Laws

We now attack Jackson's Chap. 6. Our order-of-battle includes...

MAXWELL EQUATIONS || $\left. \begin{array}{l} \textcircled{1} \nabla \cdot \mathbf{D} = 4\pi \rho, \text{ } \text{free} \\ \textcircled{2} \nabla \cdot \mathbf{B} = 0, \\ \textcircled{3} \nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}, \\ \textcircled{4} \nabla \times \mathbf{H} = +\frac{1}{c} \frac{\partial \mathbf{D}}{\partial t} + \frac{4\pi}{c} \mathbf{J}, \text{ } \text{Maxwell addition} \text{ } \text{free} \end{array} \right\} \textcircled{1}$

CONSERVATION of CHARGE } $\nabla \cdot \{ \text{Eq. } \textcircled{4} \}, \text{ and use of Eq. } \textcircled{1} \Rightarrow \boxed{\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{J} = 0} \textcircled{2}$

LORENTZ FORCE LAW || $\left. \begin{array}{l} \mathbf{F} = q(\mathbf{E} + \frac{1}{c} \mathbf{v} \times \mathbf{B}), \text{ on single charge } q; \\ \mathbf{F} = \rho \mathbf{E} + \frac{1}{c} \mathbf{J} \times \mathbf{B}, \text{ force/vol. on distributions } \rho \text{ \& } \mathbf{J}. \end{array} \right\} \textcircled{3}$

CONSTITUTIVE RELATIONS || $\left. \begin{array}{l} \mathbf{D} = \mathbf{E} + 4\pi \mathbf{P} \rightarrow \epsilon \mathbf{E}, \text{ } \mathbf{P} = (\text{bound}) \text{ polarization field}; \\ \mathbf{B} = \mathbf{H} + 4\pi \mathbf{M} \rightarrow \mu \mathbf{H}, \text{ } \mathbf{M} = (\text{bound}) \text{ magnetization field}. \end{array} \right\} \textcircled{4}$

NOTE... In same vein, Ohm's Law often assumed: $\mathbf{J} = \sigma \mathbf{E}$, $\sigma = \text{electrical conductivity}$ $\textcircled{5}$

NOTE... For vacuum (non-interacting sources): $\epsilon = 1$ & $\mu = 1$, $\text{so } \mathbf{D} \equiv \mathbf{E}, \mathbf{B} \equiv \mathbf{H}$. $\textcircled{6}$

FIELD ENERGY DENSITIES || $\left. \begin{array}{l} \text{Electric [linear medium: } \epsilon \neq \text{fcn}(\mathbf{E})]: u_E = \frac{1}{8\pi} \mathbf{E} \cdot \mathbf{D}; \\ \text{Magnetic [linear medium: } \mu \neq \text{fcn}(\mathbf{B})]: u_M = \frac{1}{8\pi} \mathbf{B} \cdot \mathbf{H}. \end{array} \right\} \textcircled{7}$

REPRESENTATION by POTENTIALS || $\left. \begin{array}{l} \mathbf{B} = \nabla \times \mathbf{A}, \text{ still OK by Max. Eq. } \textcircled{2}; \\ \mathbf{E} = -\nabla \phi - \frac{1}{c} \frac{\partial \mathbf{A}}{\partial t}, \text{ } \mathbf{A} \text{ term reg'd by Eq. } \textcircled{3}. \end{array} \right\} \textcircled{8}$

This is a powerful strike force. Virtually all of classical E & M appears here.

† u_M not actually derived until Jk^{II} Sec. 6.2, Eq. (6.16).