3. In MKS units, choose  $k = Mo/4\pi$  valuer than  $k = \frac{1}{c}(CGS)$ , so the MKS formulas for B are gotten by replacing  $\frac{1}{c}(L_{MKS}) = \frac{1}{c}(2I/r)[CGS] \rightarrow \frac{\mu_0}{4\pi}(2I/r) = \mu_0 I/2\pi r [MKS]$ .

3) Eq. (11) [ ~ Amperés Taw] shows how a source Idl generates a magnetic field dB, but we still need to know how Idl complex to an already existing field B. Answer is ... | dF f I

$$dF = \frac{1}{c} I(d1 \times B) [Corentz' Caw].$$
 (13)

I is due to the motion of a single charge  $\Delta q$ , then Idl=( $\Delta q$ ) V, and :  $\Delta F = (\Delta q/c) V \times IB$ , which is Torentz' Law.

 $d^2 F_{2m1} = \frac{1}{c} I_1 (d I_1 \times d B_2)$ 

i.e./ 
$$d^2 \mathbf{F}_{21} = \frac{\mathbf{I}_1 \mathbf{I}_2}{c^2} \left[ \frac{d \mathbf{I}_1 \times (d \mathbf{I}_2 \times \mathbf{F}_{21})}{\tau_{21}^3} \right]$$

$$\xrightarrow{\alpha_{4/}} d^2 F_{21} = \frac{I_1 I_1}{c^2 \gamma_{21}^2} \left[ (d \mathbf{l}_1 \cdot \hat{\mathbf{r}}_{21}) d \mathbf{l}_2 - (d \mathbf{l}_1 \cdot d \mathbf{l}_2) \hat{\mathbf{r}}_{21} \right]. \tag{14}$$

This is the force by Irda on Irda. Reversing the voles ...

$$d^{2}\mathbf{F}_{12} = \frac{\mathbf{I}_{2}\mathbf{I}_{1}}{c^{2}\gamma_{12}^{2}} \left[ (d\mathbf{I}_{2} \cdot \hat{\mathbf{r}}_{12}) d\mathbf{I}_{1} - (d\mathbf{I}_{2} \cdot d\mathbf{I}_{1}) \hat{\mathbf{r}}_{12} \right] \int_{\gamma_{12} = \gamma_{21}}^{\hat{\gamma}_{12}} \left[ (d\mathbf{I}_{2} \cdot \hat{\mathbf{r}}_{12}) d\mathbf{I}_{1} - (d\mathbf{I}_{2} \cdot d\mathbf{I}_{1}) \hat{\mathbf{r}}_{12} \right]$$

$$\int_{0}^{\infty} \left[ d^{2} \mathbf{F}_{11} + d^{2} \mathbf{F}_{21} = \frac{\mathbf{I}_{1} \mathbf{I}_{2}}{c^{2} \gamma_{21}^{2}} \left[ (d \mathbf{1}_{1} \cdot \hat{\gamma}_{21}) d \mathbf{1}_{2} - (d \mathbf{1}_{2} \cdot \hat{\gamma}_{21}) d \mathbf{1}_{1} \right] + 0 \right] (15)$$

A seeming Disaster in by Newton III, should have: d2 Frz + d2 Frz = 0. What has been left out is that both Ik d1k are parts of current loops.

## Magneto statics (cont'd)

1) What must be done to recover from this disaster is to include the loops...

$$dF_{21} = \frac{1}{c} I_1 (dI_1 \times B_2),$$

We Bz =  $\frac{I_z}{c} \oint \frac{d\mathbf{l}_2 \times \hat{\mathbf{r}}_{21}}{r_{21}^2} \int \frac{\text{loop # z at site of } I_1 d\mathbf{l}_1.$ 

Entire force by loop # 1 } F\_{2m1} = \( \phi \) dF\_{21} = \( \frac{\text{I\_1 I\_2}}{c^2} \phi \) \( \phi \) \( \frac{1}{\gamma\_{21}^2} \) d\( \text{I\_1} \times \frac{\text{I\_2}}{\gamma\_{21}^2} \)

$$\mathbb{F}_{2\text{oni}} = \frac{\mathbb{I}_{1}\mathbb{I}_{2}}{c^{2}} \oint \oint \frac{1}{\gamma_{21}^{2}} \left[ \left( \frac{d\mathbb{I}_{1} \cdot \hat{\gamma}_{21}}{r_{21}} \right) d\mathbb{I}_{2} - \left( d\mathbb{I}_{1} \cdot d\mathbb{I}_{2} \right) \hat{\gamma}_{21} \right]. \tag{17}$$
Contributes zero on integration

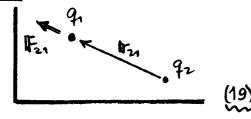
The indicated term contributes zero because it is  $\alpha \oint dR_2 \oint dR_1 \cdot \frac{r_{21}}{r_{21}^2}$ , and the  $\alpha \oint dr \cdot \frac{r}{r^3} = \frac{1}{2} \oint \frac{dr}{r^2} = -\frac{1}{2} \oint d(1/r) = 0$ . The total force law is then...

$$F_{2m1} = (-) \frac{I_1 I_2}{c^2} \oint \oint \frac{1}{r_{21}^2} (dl_1 \cdot dl_2) \mathring{r}_{21} \left[ Jk^2 E_g. (5.10) \right]. \tag{18}$$

Now, clearly, From = (-) From, and Newton III is obeyed, as must be.

Contrast Fzm (magnetic) with Fzm (electric)...

Sy [F2 mm (electric) = 9192 
$$\frac{\hat{\tau}_{21}}{\tau_{21}^2}$$
]



The sources of B, refusing to be monopoles, really complicate the fire law. About the only simple application of Eq. (18) is to two II long wires...

$$I_1$$
 $F$ 
 $I_2$ 

$$\frac{dF}{d\ell} = 2I_1I_2/c^2d$$

(20)

F is attractive if In & Iz flow 11, repulsive if In & Iz anti-11.