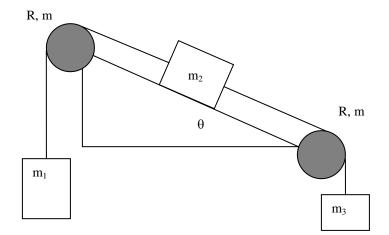
1) A metal detector is made from a single horizontal circular loop of radius R in which a constant current I_0 is maintained using an ideal current generator. A metal sphere of radius a << R is buried a depth d >> R beneath the soil (soil has $\mu = \mu_0$). Consider the sphere to be made of linear paramagnetic material of permeability μ . In the following retain only the leading order in small parameters a/R and R/d.

- a. What dipole moment is induced in the sphere when the coil is at ground level directly above it?
- b. By how much does the flux inside the current loop change as a result of moving it to the point above the sphere from somewhere far away?
- c. Denote the self-inductance of the loop, with no magnetic material around, $L_o = \alpha \mu_o R$ where α is a constant. What is the *relative change* in its self-inductance, $\Delta L/L_0$, when it is moved to the point over the metal sphere? Note whether the change is an increase or decrease.

Three masses m_1 , m_2 and m_3 are connected by massless ropes over two pulleys as shown. The solid disc pulleys have a mass m and a radius R. The slope has a coefficient of kinetic friction μ_K and a coefficient of static friction μ_S . The incline is at an angle θ to the horizontal.

- a) Find the minimum mass m₁
 for which the mass m₂ will
 begin to move up the incline.
- b) Once the mass m₂ starts to slide, find the steady acceleration 'a' of m₂.
- c) Find the power lost to heat by the friction of the mass m_2 on the incline as a function of time t.



- The tritium 3H atom is a metastable isotope of the 1H atom with a half-life of 12.33 years. It undergoes a nuclear transition and transforms into a ${}^3He^+$ ion through a *beta* (e^-) and an *antineutrino* (\bar{v}) decay, ${}^3H \rightarrow {}^3He^+ + e^- + \bar{v}$. One of the neutrons in the nucleus of 3H is converted into a proton. The energetic electron generated by the beta decay escapes the atom, leaving the He in a positively charged state. The electron that was in the ground state orbiting the 3H atom suddenly finds itself orbiting a ${}^3He^+$ ion. Assume that the *beta* decay is too fast for the orbiting electron to adjust to the new configuration until after the decay is complete. Answer the following questions:
- a. Assuming that a 3H atom starts out in the ϕ_{100} ground state, prove that the probability of the electron orbiting the ${}^3He^+$ ion in the $\phi_{nlm} = R_{nl}(r)Y_{lm}(\theta,\phi) = \phi_{210}$ state of the ${}^3He^+$ ion is zero. Is this expected? Explain.
- b. The result found in (a) would suggest that finding the ${}^3He^+$ ion in any excited state as a result of nuclear transition is zero. Rather than rushing to this conclusion, now determine the probability that the ${}^3He^+$ ion is in the $\phi_{nlm} = \phi_{200}$ state. From this result draw a general conclusion on the symmetry of the excited state wave functions.

(Hint. Useful functions:
$$\phi_{100} = \frac{1}{\sqrt{\pi}} \left(\frac{Z}{a_0}\right)^{3/2} e^{-Zr/a_o}, \ \phi_{210} = \frac{1}{4\sqrt{2\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \frac{Zr}{a_0} e^{-Zr/2a_o} \cos\theta$$

$$\phi_{200} = \frac{1}{4\sqrt{2\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \left(2 - \frac{Zr}{a_0}\right) e^{-Zr/2a_o}, \ Y_{00}(\theta, \phi) = \frac{1}{\sqrt{4\pi}})$$

A spin-1/2 particle interacts with a magnetic field $\vec{B} = B_o \hat{z}$ through the Pauli interaction $H = \mu \vec{\sigma} \cdot \vec{B}$, where μ is the magnetic moment and $\vec{\sigma} = (\sigma_x, \sigma_y, \sigma_z)$ are the Pauli spin matrices. At t = 0, a measurement determines that the spin is pointing along the positive x-axis. Calculate the probability that it will be pointing along the negative y-axis at a later time t.

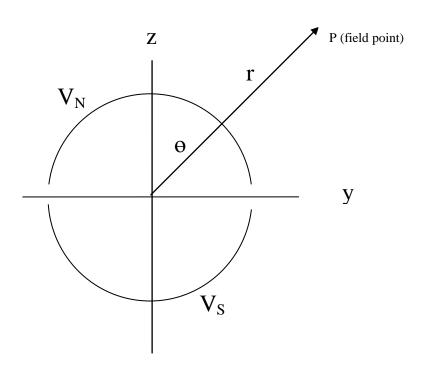
NOTE: The Pauli spin matrices are given by

$$\sigma_{x} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \qquad \sigma_{y} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \qquad \sigma_{z} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

5) Let a conducting sphere of radius $\bf a$ be split into hemispheres, held at potentials V_N and $V_{S,}$ respectively (see Figure 1). Find the potential Φ (r, θ) outside the sphere. Let Φ (r= ∞ , θ)=0.

Hint:
$$\int_0^1 P^{\ell}(x) dx = \frac{(-1)^{(\ell-1)/2} (\ell-1)!}{2^{\ell} \{(\ell+1)/2\}! \{(\ell-1)/2\}!}, \quad \text{if } \ell = \text{odd}$$
$$= 0 \quad \text{if } \ell = \text{even}$$
$$= 1 \quad \text{if } \ell = 0$$

Figure 1

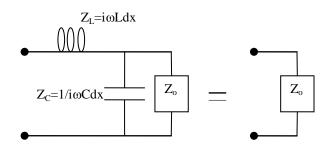


The figure below shows a spring pendulum consisting of a massless spring of spring constant k and unstretched length r_o . One end of the spring is pivoted from the ceiling while the other is attached to a point mass m. The pendulum is restricted to movement in an x-y plane as shown in the figure. The pendulum is pulled off its equilibrium and set into motion. Answer the following questions:

- a. Write down the Lagrangian and determine the equations of motion for m using polar coordinates r and θ , shown in the figure.
- b. Identify each component in the equations found in part (a) in terms of Newton's laws of motion.
- c. Now make small amplitude approximations in the Lagrangian equations associated with the $r (= r_e + u)$ and θ motions but keep the terms in the **second order** in the amplitudes of motion (i.e. in θ and u/r_e and in their derivatives) and write down the equations of motion for $\theta(t)$ and u(t); here r_e is the equilibrium position around which m oscillates.
- d. Now discuss the coupling between the u and θ motions implied in (c) and compare this motion with the motions in u and θ , but this time keeping only the **first-order** terms in the amplitudes of the motion.

m

- 7)
 A long coaxial cylinder transmission line is made of an inner conductor of radius a surrounded by an outer conductor of inner radius b. Assume that the region between a and b is filled with vacuum and that the conductors are perfectly conducting.
 - a) Find an expression for the capacitance per unit length C of the coaxial cylinder in terms of a and b.
 - b) Find an expression for the inductance per unit length L of the coaxial cylinder in terms of a and b.
 - c) Find the characteristic impedance Z_o for the coaxial cylinder transmission line. Hint: The two circuits to the right are equivalent. Use this fact to find the impedance Z_o for the transmission line. For the circuit on the left side, the impedances for the inductance and capacitance of a short length dx



of the line are also shown. You may assume dx is small and keep only appropriate terms.

- 8) Initially n moles of an ideal, monatomic gas are at pressure P_0 , volume V_0 , and temperature T_0 . The gas undergoes an isothermal expansion that doubles its volume. What are the final values for each of the following quantities in terms of the initial values of these quantities?
 - a. pressure
 - b. temperature
 - c. heat absorbed
 - d. work done
 - e. change in internal energy
 - f. change in entropy

Find these same quantities if the gas undergoes an adiabatic **free** expansion from volume V_0 to $2\ V_0$.

A two-dimensional simple harmonic oscillator has oscillation frequency Ω_0 . It is prepared in the state

$$\psi(x,y) = \left[\frac{\sqrt{3}}{2}\varphi_1(x) - \frac{1}{2}\varphi_2(x)\right] \left[\frac{1}{\sqrt{2}}\varphi_0(y) - \frac{i}{\sqrt{2}}\varphi_1(y)\right]$$
(1)

where φ_n is the n^{th} energy eigenfunction of the *one-dimensional* simple harmonic oscillator - φ_0 is the ground state.

- a. If the energy of this state were measured, what possible values could be found? (Express them in terms of Ω_0 .) With what probability will each energy be observed?
- b. What is the expected energy $\langle E \rangle$?
- c. The angular momentum of the particle is represented by the operator

$$\hat{\mathbf{L}}_{\mathbf{z}} = \mathrm{i}\hbar \left(\hat{\mathbf{a}}_{\mathbf{x}} \hat{\mathbf{a}}_{\mathbf{y}}^{\dagger} - \hat{\mathbf{a}}_{\mathbf{x}}^{\dagger} \hat{\mathbf{a}}_{\mathbf{y}} \right)$$

written in terms of raising and lower operators for the 1d SHO satisfying the following relationships

$$\hat{a}_{x}^{\dagger} \, \phi_{n}(x) = \sqrt{n+1} \, \phi_{n+1}(x)$$
 , $\hat{a}_{x} \, \phi_{n}(x) = \sqrt{n} \, \phi_{n-1}(x)$

and similarly for \hat{a}_y^{\dagger} and \hat{a}_y acting on $\varphi_n(y)$ (You do not need to show this -- it is true.) What is the expectation $\langle L_z \rangle$ for a particle in state given by eq. (1).

- 10)
- (a) Starting with the equation for mechanical pressure P:

$$P = \frac{1}{3} \int_0^\infty p \, \mathbf{v} \, \mathbf{n}(\mathbf{p}) \, \mathbf{dp}, \tag{1}$$

where p is the momentum, v is the velocity of the constituent particles and n is the momentum distribution of the particles.

Utilizing the momentum distribution in thermal equilibrium:

$$n (p) dp = \frac{4 \pi N p^2 dp}{V(2 \pi m k T)^{3/2}} exp \{-p^2/(2 m k T)\},$$
 (2)

show that the equation of state of a perfect nonrelativistic gas is expressed as:

$$PV = N k T, (3)$$

where T is temperature, N is the number of particles in a volume V and m is the mass of a particle.

(b) From the expression for force balance, show that the structure of a stationary spherical object, e.g., a star, in hydrostatic equilibrium is found from:

$$\frac{dP(r)}{dr} = -\frac{\rho(r) G M(r)}{r^2},$$

$$\frac{dM(r)}{dr} = \rho(r) 4\pi r^2$$
(4)

where P(r) and $\rho(r)$ are pressure and density at distance r from the center, M(r) is the total mass within a sphere of radius r, and G is the gravitational constant.

Consider the inhomogeneous equation with the differential operator \mathcal{L}

$$\mathcal{L}\psi = \frac{d}{dx}\left(x^2 \frac{d\psi}{dx}\right) - 2\psi = f(x)$$

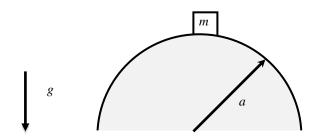
- a. Find a complete set of **homogeneous** solutions by proposing $\psi = Ax^{\lambda}$
- b. For the boundary value problem

$$\mathcal{L}\psi = f(x)$$
, $\psi(1) = 1$, $\psi \to 0$ as $x \to \infty$ (1)

find the Green's function for the operator \mathcal{L} within the region 1 < x.

c. Use the Green's function from b. to solve eq. (1) for the choice of RHS $f(x)=x^{-2}$

A smooth hemispherical surface of radius a is placed with its flat side down and fastened to the Earth whose gravitational acceleration is g. A small object of mass m is placed on top of the hemisphere and given a very small push $(v_{in} \approx 0)$. If the surface is frictionless, use the method of Lagrange multipliers to calculate at what point the small object leaves the surface of the hemisphere.

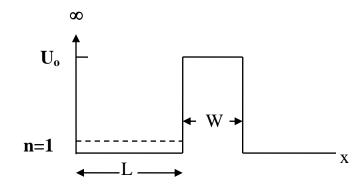


A paramagnetic material is subjected to a magnetic field, $\vec{B} = B_o \hat{z}$, and is in thermal equilibrium with a reservoir at temperature T. Assume that the paramagnetic material is made up of N distinguishable magnetic particles, with each particle having a net angular momentum of j=1 and a magnetic dipole moment of $\vec{\mu} = \mu_o \vec{j}$. Answer the following questions.

- a. Determine the entropy of the paramagnetic material under the conditions described above.
- b. Now assume that the paramagnetic material is isolated from the reservoir and the magnetic field is reduced adiabatically to $\vec{B}_f = (B_o/100)\hat{z}$. Determine the final temperature of the system in terms of T as a result of this adiabatic demagnetization.

(Hint: $F = -kT \ln Z$ and $dF = -SdT - pdV + \mu dN$, where F is the Helmholtz free energy and Z is the canonical partition function.)

An electron is in the lowest state (n=1) of a 1-D quantum well of width L as shown in the figure. The potential wall at x=0 is infinitely high, while the potential wall on the right at x=L is finite with a height of U_o and width W forming a barrier. Estimate the initial rate of escape of the electron through the barrier.



Assume that
$$L \gg \frac{\hbar}{\sqrt{mU_o}}$$
 so that the

initial state can be approximated by the ground state of an infinite 1-D square well.

Also assume that $W \gg \frac{\hbar}{\sqrt{mU_o}}$ when finding the initial rate of escape through the barrier.

- (a) Calculate the multipole moment $q_{\ell m}$ of the charge distributions shown in Figure 1 (positive charge q on the z axis distance **a** above and below the origin, and negative 2q at the origin). Obtain results for nonvanishing moments valid for all ℓ .
- (b) For the charge distribution shown in Figure 1 write down the multipole expansion for the potential ϕ , for r > a.
- (c) What is the potential if you keep only the lowest order term in the expansion (valid for r >> a)?
- (d) How does the potential in (c) behave as function of distance r on the x-y plane?
- (e) Calculate directly from Coulomb's law the exact potential for (a) in the x-y plane, and compare the behavior with the solution in (d) as $r \to \infty$.

$$\begin{split} \text{Hints:} \qquad q_{\ell m} &= \int Y_{\ell m}^*(\theta,\phi) \; r^\ell \, \rho(\boldsymbol{x}) \; d^3 \boldsymbol{x}. \\ \varphi &= \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{+\ell} \frac{4\pi}{2\ell+1} r^{-\ell-1} Y_{\ell m} \left(\theta,\phi\right) q_{\ell m} \end{split}$$

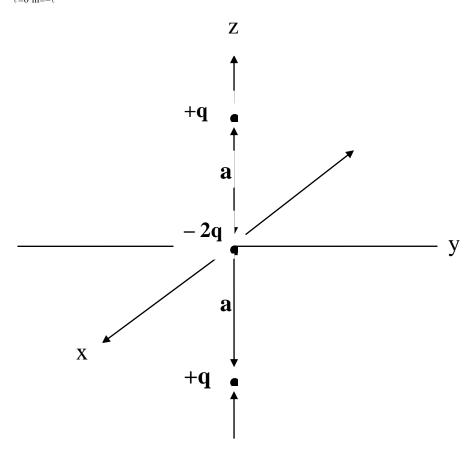


Figure 1