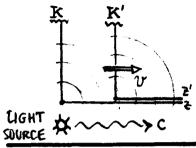
3) If c=cost, then the GT does not work for EM. The transformation which does work -- called the Loventz Transformation (LT) -- guaranteeing that C= same cost in both systems K & K', can be found by the following analysis ...



Z & Z' axes of systems K & K' are 11. K' moves down Z-axis of K at velocity v. At t=t'=0, when K & K' origins coincide, a light source flashes a pulse. The pulse must propagate as a spherical wave in both K & K', in order to be independent of the motion. So:

Assume the K&K' coordinates at pulse front, viz. (t; x, y, z) & (t'; x', y', z'), are related by a <u>linear</u> transformation. Where y is an as-yet-unknown for of 1v1:

$$\begin{cases} Z' = \gamma(z - vt), & y' = y, \quad \chi' = x \leftarrow \text{ just a modified GT;} \\ t' = \gamma(t - \frac{vz}{c^2}) \leftarrow \text{ gives up idea of absolute time (!): } t' \neq t. \end{cases}$$

Civing up the idea of (Newton's) absolute time is the BIG STEP. Now plug Egs. (12) into Eq. (11) and equate like powers of (t; x, y, z) to get:

$$\chi'_{0} = \gamma(x_{0} - \beta x_{1});$$

$$\chi'_{1} = \gamma(x_{1} - \beta x_{0});$$

$$\chi'_{1} = \gamma(x_{1} - \beta x_{0});$$

$$\chi'_{2} = \chi_{2}, \chi'_{3} = \chi_{3}$$

$$\chi'_{2} = \chi_{2}, \chi'_{3} = \chi_{3}$$

$$\chi'_{3} = \chi_{3}$$

$$\chi'_{4} = \gamma(x_{1} - \beta x_{0});$$

$$\chi'_{5} = \chi'_{5} = \chi'_{5} = \chi'_{5}$$

$$\chi'_{5} = \chi'_{5} = \chi'_{5} = \chi'_{5}$$

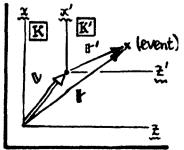
$$\chi'_{5} = \chi'_{5} = \chi'_{5} = \chi'_{5}$$

$$\chi'_{5} = \chi'_{5} = \chi$$

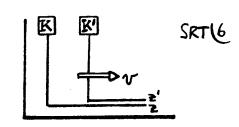
The inverse transform K+K is found by interchanging {unprimed} variables, and B->(-)B.

The IT can be generalized to the case where the K&K' axes are TEST Still parallel, but v is not necessarily along the Z-axis:

$$x' = \gamma(x_0 - \beta \cdot r);$$
 $y' = \gamma(r - \beta x_0) - (\gamma - 1)[r - \frac{1}{\beta^2}(\beta \cdot r)\beta].$
[Jkⁿ Eq. 11.19)]



SRT Introd 1 Consequences of LT.



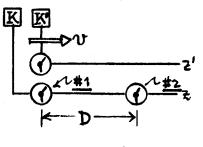
REMARKS on Loventz Transform (LT).

1. The immediate consequences of the LT are to change the absolute "character of measurements of length and time for KEK'. Some elementary results are:

Plength of Stick ΔL... ΔL(at rest) \rightarrow Δ L'(moving) = Δ L(at rest) $\sqrt{1-\beta^2}$ Scontraction (15)

B Duration of event ΔT... ΔT (same pt.) - ΔT'(two pts) = ΔT (same pt.) / 1-β2 Stime (15)

D'Simultaneity: events simultaneons in K are not simultaneons in K', and vice versa.



Two clocks, located on Z-axis of K, are synchronized by K.

K', passing by at velocity v, checks the synchrony by one clock situated at K' origin. K' finds the clocks in K are not synchronized: to K', clock #2 shows a later time than #1 by an amount: $\Delta t = vD/c^2$ (in K-time). (17)

K clocks synchronized by K.

2. The "warping" of ΔL > ΔΙ' & ΔΤ > ΔΤ' measurements for K& K'... more precisely the dependence of lengths & times on the relative motion of K& K'... is demanded by the constancy of lightspeed C. ΔΙ' & ΔΤ's are no longer absolate quantities independent of the relative motion of the observer. Only the ratio ΔL/ΔΤ = C, in the tracking of a light pulse, must be constant.

It is worth noting that Newtonian mechanics contains no intrinsic scale constants... it scales arbitrarily with M(mass), L(length), and T(time). By contrast, Maxwell's E&M generates the constant C = lightspeed, intrinsic to the whole theory of E&M. Then L&T scales cannot be fixed independently... you always have to adjust L=CT for phenomena involving EM radiation.

C=cnst <00 => E&M obeys the LT; th=cnst>0 => QMobeys quantization.

Do all fund = ensts (e, me, etc.) => some radical condition on the "right" theory?

REMARKS on IT (cont'd)

3. The essence of the LT is that it preserves the "space-time interval, i.e. (DS) = (CDt) - (DX) , for all possible events measured by K & K'... not just observation of a light pulse (per Eg. 191) above), but all events. Consider...

Starts at finisher at duration separation

[EVENT] { Viewed by
$$K: (x_1,t_1) \quad (x_2,t_2) \quad \Delta x_0 = c(t_2-t_1) \quad \Delta x = (x_2-x_1) }$$

[EVENT] { Viewed by $K': (x_1',t_1') \quad (x_2',t_2') \quad \Delta x_0' = c(t_2'-t_1') \quad \Delta x' = (x_2'-x_1') }$

[18)

Now, construct: $(\Delta S')^2 = (\Delta X')^2 - (\Delta X')^2 \leftarrow \text{luent spacetime interval for } K'$.

Plug in LT intervals: $\Delta X'_0 = \gamma (\Delta X_0 - \beta \Delta X)$, $\Delta X' = \gamma (\Delta X - \beta \Delta X_0)$.

Find (2) Some algebra): $C^2(\Delta t')^2 - (\Delta x')^2 \equiv C^2(\Delta t)^2 - (\Delta x)^2 \leftarrow \text{interval for } K$.

$$(\Delta s)^2 = c^2(\Delta t)^2 - (\Delta x)^2, \text{ is invariant under a LT.}$$
(19)

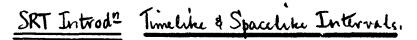
This invariance contains most of the physics of the LT. E.g. we can derive the time dilation formula quoted in Eq. (16) above. As follows...

Let K measure an event of duration Dt, on one clock fixed at one place in his system; $\Delta x=0$. For the same event, the relatively moving observer K' measures $\Delta t'$ on his (one) clock, but the start of finish of the event are separated in space by $\Delta x'=v\Delta t'$. Invariance of the spacetime interval permits us to write:

$$\longrightarrow (C\Delta t')^2 - (v\Delta t')^2 = (c\Delta t)^2 - (\Delta x)^2, \quad \Delta t' = \Delta t / \sqrt{1-\beta^2} \int_{\text{dilation}}^{\text{time}} \frac{(20)}{\sqrt{1-\beta^2}} \frac{(20)}{\sqrt{1-\beta^2}} \int_{\text{dilation}}^{\text{time}} \frac{(20)}{\sqrt{1-\beta^2}} \frac{(20)}{\sqrt{1-\beta^2}} \frac{(20)}{\sqrt{1-\beta^2}} \frac{(20)}{\sqrt{1-\beta^2}} \frac{(20)}{\sqrt{1-\beta^2}} \frac{(20)}{\sqrt{1-\beta^2}} \frac{(20)}{\sqrt{1-\beta^2}} \frac{(20)}{\sqrt{1-\beta^2}} \frac{(20)}{\sqrt{1-\beta^2}} \frac{(20)}{$$

EXERCISE: Deric length contraction [Eq. (15)] by considering invariance of (DS)?

We note that $(\Delta S)^2 = (C\Delta t)^2 - (\Delta x)^2$ can be $(H)ve_1(-)ve_1$ or zero... depending on the relative sizes of the space-interval $\Delta x = 0$ time-interval between the start of finish of the event. Further, the sign of $(\Delta S)^2$ is a Lorentz invariant.



The Torentz invariance of the Size of sign of the Spacetime interval (AS)² provides a scheme for Classifying events according to whether (AS)²

Is Hive, (-) ve, or zero. This classification puts

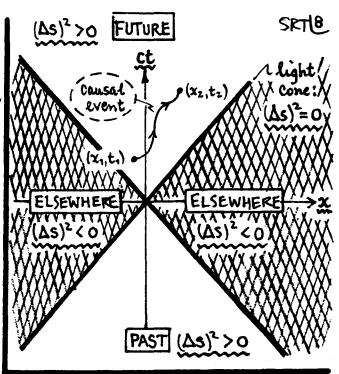
Causality on a quantitative basis. For EVENTS

Specified by: (x₁, t₁)[start] of (x₂, t₂)[finish],

We consider the interval:

$$|(\Delta s)^2 = c^2 (t_2 - t_1)^2 - (x_2 - x_1)^2 = cnst$$
 (21)

Three cases are possible:



$$\frac{(\Delta S)^2 = 0}{(\Delta S)^2 = 0}, \frac{x_2 - x_1}{(\Delta z - z_1)} = \pm c \cdot \begin{cases} \text{EVENTS } 1 \neq 2 \text{ can only concern broadcast } \Rightarrow \text{ Teception} \\ \text{of a light Signal; this defines "light Cone" in above dam.} \end{cases}$$

Por 2 (Δ5)²>0, ^{SQ} | x₂-x₁ | < C. {EVENTS 1 & 2 can be connected cansally by a light signal; events ²/(Δ5)²>0 lie <u>inside</u> the light cone, and they define the PAST of FUTURE as indicated.

frame @ $v = \frac{x_2 - x_1}{t_2 - t_1}$ => each $x' = \gamma(x - vt) = 0$, this $(\Delta s)^2 > 0$ is called TIME-LIKE.)

(\(\D s)^2 = c^2(t_2' - t_1')^2\) \(\((\D s)^2 > 0\) is called TIME-LIKE.)

(22b)

OHUNNE (ΔS)²<0, | | | | | | > C. { EVENTS 1 & 2 <u>Cannot</u> be connected <u>Cansally</u> by a light Signal; events | (ΔS)²<0 lie <u>outside</u> the light cone, in a place called (with engaging whimsy) ELSEWHERE.

frame @
$$\frac{v}{c} = \frac{c(t_2-t_1)}{x_2-x_1} \Rightarrow \frac{each \ t'=\gamma \ (t-\frac{vx}{c^2})=0}{(\Delta s)^2 = -(x_2'-x_1')^2} \int \frac{this}{(\Delta s)^2 \langle 0 \ is \ SPACE-LIKE.} (EVENT' is at same t', different x's).$$

In the real world, causally connected wents occur inside or on the light cone, (22c) So long as C= const is the limiting signal velocity in the universe. ELSEWHERE is the realm of coincidence, chance, and the current home of John Sumunu.