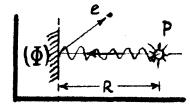
## \$506 Problems

A lightbulb, radiating total power P=10W, is placed at distance R=1m from a large metal plate. The metal has a

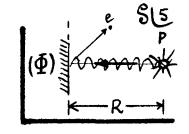


work function  $\Phi$ =3eV. An electron in the metal can absorb the incident radiation over an area adjacent to it -- suppose that over is a circle of radius  $\tau$ ~ several atomic radii (Say 5). Treat this problem <u>classically</u>.

- (A) Calculate the time Dt required for the electron to absorb enough energy to be liected from the surface. It is observed that Dt < 1ns. Comment.
- (B) With all else the same, what lightbulb power is needed to eject the electron in 1ns? Compare this with the output of a typical commercial power plant.
- (5/A) Show that a photon cannot transfer all of its energy to a free electron.
  - (B) Suppose a photon, incident on a free electron, has energy  $E_i = Nmc^2$ , where m is the electron rest mass and N is a numerical factor. What value of N is required for the photon to transfer 99% of its energy? What is  $\lambda$  for this photon?
- ⑥ Apply the Correspondence Principle to radiation from a transition  $n \rightarrow n-1$  in a hydrogenlike Bohr atom. The average power radiated is:  $\frac{P_n = \frac{1}{T_n}}{\Delta E(n \rightarrow n-1)}$ ,  $\frac{M_p}{\Delta E} = 0$  energy smitted, and  $\frac{P_n = \frac{1}{T_n}}{\Delta E(n \rightarrow n-1)}$ ,  $\frac{M_p}{\Delta E} = 0$  the classical Larmor radiation rate:  $\frac{P_n = (2e^2/3c^3)|a_n|^2}{2c^3}$ , for a charge ε undergoing (centripetal) acceleration  $\frac{P_n = (2e^2/3c^3)|a_n|^2}{2c^3}$ , and define the dimensionless constant:  $\frac{P_n = e^2}{2c^2}$ . α is the fine-structure const.
- (A) Let  $\Gamma_n = \frac{1}{T_n}$  be the "transition probability / unit time" for  $n \to n-1$ . Show that:  $\Gamma_n = \Gamma_1/n^5$ , where:  $\Gamma_1 = \frac{2}{3} Z^4 \alpha^5 \cdot (mc^2/\hbar)$ .
- (B) For hydrogen (Z=1), compare the result in part (A) for the  $n^{th}$  state lifetime  $T_n = \frac{1}{T_n}$  with known values:  $T(n=2 \rightarrow 1) = 1.6 \text{ ns}$ ,  $T(n=4 \rightarrow 3) = 73 \text{ ns}$ , and  $T(n=6 \rightarrow 5) = 610 \text{ ns}$ . Comment on the <u>trend</u> of the comparisons.

## \$ 506 Solutions

4 Time & power scales for the classical photo-electric effect.



(A) The radiation incident on the plate delivers a flux of size

If a given electron gathers energy from an area = circle of radius na, 2/2 = atomic radius and n= numerical factor, it absorbs energy at a rate...

$$\rightarrow \frac{\ln c r g n}{time} = (P/4\pi R^2) \cdot \pi (na)^2 = P \cdot (na/2R)^2.$$

In order to be ejected in time Dt, the Mectron must gather an energy exceeding the work function barrier, i.e.

$$\rightarrow \left(\frac{\ln e r q u}{t i m u}\right) \Delta t = P \Delta t \cdot (na/2R)^2 \geqslant \Phi \Rightarrow \Delta t \gg \frac{\Phi}{P} \cdot (2R/na)^2.$$
 (3)

For the given numbers ...

$$\Phi = 3eV = 4.81 \times 10^{-19} \text{ J}, P = 10W = 10 \text{ J/sec}$$

$$R = 1m = 100 \text{ cm}; \quad a = 0.53 \times 10^{-8} \text{ cm} \left(\frac{80 \text{ hr}}{\text{radin}}\right) \neq n = 5...$$

$$\Delta t = \frac{4.81 \times 10^{-19}}{10} \left(2 \times 100 / 5 \times 0.53 \times 10^{-8}\right)^2 = 2.74 \text{ sec}$$
(4)

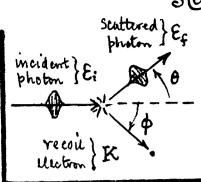
The observed Dt < 1ns. Clearly the classical picture of the electron <u>continu</u>-<u>ously</u> absorbing the vadiant energy is monumentally WRONG. Enter Ein-Stein, center stage: the "absorption" is actually a collision.

(B) If  $\Delta t = 2.74$  sec is to be decreased to 1ns, then -- per Eq.(3) -- the lightbulb power P must be increased by a factor 2.74/10-9. This gives a required power...

$$\rightarrow$$
 for  $\Delta t = 1$ ns:  $P = 2.74 \times 10^9 \times 10W = 27,400 MW$ . (5)

A typical commercial power plant puts ont ~ 100-300 MW. So the power required in Eq. (5) is equivalent to the ontput of ~100 such plants.

- 6 Most energetic photon -> electron collision.
- (A) Ref. CLASS, p. Intro. 12, Eqs. (26) (27). There, for a photon -> free electron collision, we have written energy & momentum conservation extres, and have found

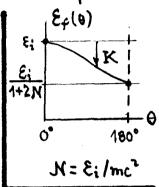


the final photon frequency. Since energy E= hv for the photon, Eq. (27) is:

$$\frac{\mathcal{E}_{f} = \mathcal{E}_{i} / [1 + (\mathcal{E}_{i} / mc^{2})(1 - \cos \theta)]}{\text{for photon}}, \theta = \frac{1}{\text{for photon}}. (1)$$

The photon transfers a recoil (kinetic) anagy K to the e, of an amount...

$$\rightarrow K = \varepsilon_i - \varepsilon_f = \varepsilon_i \left[ \frac{N(1-\cos\theta)}{1+N(1-\cos\theta)} \right] \langle \varepsilon_i, \mathcal{N} = \frac{\varepsilon_i}{mc^2}, (2)$$



R increases monotonically with photon scattering 4 θ, reaching its max. value @ θ = 180° (photon backscattering), whence

$$\rightarrow K_{\text{MAX}} = K(\theta = 160^{\circ}) = \mathcal{E}_{i} \left[ \frac{2N}{1+2N} \right] < \mathcal{E}_{i}$$

(3)

For any finite photon energy N, Kmax (Ei, and-with this max. energy transfer-the photon still has energy Ef = Ei/(1+2N). So, videed, the photon cannot transfer all its energy to the free E... that transaction would violate conservation of momentum.

(B) If the transfer E; → K is to be 99% complete, then the least value of N necessary is -- from Eq. (3)...

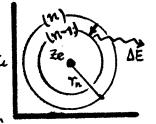
$$\rightarrow \mathbb{K}_{MRX} = \left. \mathcal{E}_{i} \left[ \frac{2N}{1+2N} \right] = 0.99 \, \mathcal{E}_{i} \Rightarrow \underline{N} = \frac{1}{2} \left( \frac{1}{\epsilon} - 1 \right) \right|_{\epsilon=0.01} = \underline{49.5}. \tag{4}$$

This is a very robust photon, with incident energy  $E_i = Nmc^2 = 25.3 \text{ MeV}.$  Its wavelength is...

$$\Delta = hc/\epsilon_i = \frac{1}{N}(h/mc) = \frac{1}{49.5} \times 2.43 \times 10^{-10} \text{ cm} = \frac{4.91 \times 10^{-4} \text{ Å}}{10^{-10} \text{ cm}} = \frac{4.91 \times 10^{-10} \text{ Å}}{10^{-10} \text{ M}}$$

6 H-atom transition votes per the Correspondence Principle.

(A)  $P_n = \Gamma_n \Delta E(n \rightarrow n-1) \rightarrow \frac{2}{3} (e^2/c^3) |a_n|^2$ , so we want to calculate  $\left[ \Gamma_n = \frac{2}{3} (e^2/c^3) |a_n|^2 / \Delta E(n \rightarrow n-1) \right] \int_{\Delta E=n-n-1}^{\Delta n=accel^n} \frac{m}{n} \frac{m}{n} \frac{m}{n} \frac{m}{n} \frac{m}{n}$ In every



Assume circular orbits of radii r. Bohr model gives [Eq. (16), p. Duality 7]:

$$\begin{cases} \text{Brbit} \\ \text{energy} \end{cases} E_{n} = -\frac{1}{2} (2\alpha)^{2} mc^{2} / n^{2} \Rightarrow \underbrace{\Delta E(n \rightarrow n-1)}_{\text{energy}} = E_{n} - E_{n-1} \approx \frac{(2\alpha)^{2} mc^{2}}{n^{3}}; \quad (2)$$

$$\left\{\begin{array}{l} \text{orbit} \\ \text{radius} \end{array}\right\} = \frac{n^2 a_0}{2}, \quad \text{if} \quad a_0 = \frac{\hbar^2}{me^2} = \text{Bohr radius}. \quad \text{NOTE: } \alpha = e^2/\kappa c. \quad (3)$$

The electron's centripetal accel an is found from:  $ma_n = -\frac{Ze^2}{r_n^2}$ , i.e.

$$\rightarrow a_n = -\frac{1}{m} z e^2 / r_n^2 = -z^3 e^6 m / n^4 t_n^4, \text{ } m = -z^3 e^2 \alpha^2 m c^2 / n^4 t_n^2.$$
 (4)

Sor Larmor power: 
$$P_n = \frac{2}{3}(e^2/c^3)|a_n|^2 = \frac{2}{3} Z^6 \alpha^2 (mc^2)^2/n^8 k$$
. (5)

Then Eq. (1) yields...

$$\rightarrow$$
 Γn = Pn [Eq.(5)]/ΔΕ[Eq.(2)] =  $\frac{2}{3}$   $Z^4$  α 5 mc²/n5 t,

$$\Gamma_n = \Gamma_1/n^5$$
,  $\omega / \Gamma_1 = \frac{2}{3} Z^4 \alpha^5 m c^2 / \hbar$ . QED (6)

(B) The difetime for 
$$n \to n-1$$
 is:  $\frac{T_n = \frac{1}{T_n} = n^5 T_1}{T_1}$ , where: (7)
$$T_1 = \frac{1}{7^4} \cdot \left(\frac{3}{2} t / \alpha^5 mc^2\right) \int \frac{\alpha}{k/mc^2} = \frac{1}{1.24 \times 10^{-21}} sec.$$

soy numerically: 
$$\underline{\tau}_1 = \frac{1}{24} \cdot 0.0932 \text{ ns}$$
.

For Z=1, we calculate the #5 at right by use of Tn=n<sup>5</sup>T<sub>1</sub> in Eq. (7). The agreement with fenous values improves as n increases; this is in accord with

TRANSMION	THIS CALCH	KNOWN Tn, ns	ratio
n=2-1	3.0	1.6	1,88
n=4+3	95	73	1.30
n=6+5	725	610	1.19

the Correspondence Principle feature that the classical result becomes more accurate as the quantum system becomes larger.