

**DEPARTMENT OF PHYSICS**

**M.S. COMPREHENSIVE / PH. D. QUALIFYING EXAMINATION**

**SATURDAY APRIL 2, 1986**

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SATURDAY, APRIL 2, 8-12 AM 1986

Answer each of the following questions, all of which carry equal weight. Begin your answer to each question on a new sheet of paper; solutions to different questions must not appear on the same sheet. Each sheet of paper must be labeled with your name and the problem number, in the upper right hand corner. When more than one sheet is submitted for a problem, be sure the pages are ordered properly. [Problems 1-8].

1. The Coriolis force is given by  $\vec{F} = -2m\vec{\omega} \times \vec{v}$ .

- a) Discuss physically what this expression means.
- b) The water in a river flowing to the south in the northern hemisphere will have the water level higher on one side than the other. Which side, east or west, will be higher?
- c) For the Mississippi river, estimate that height.

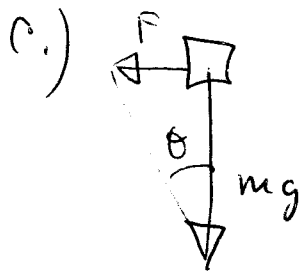
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-

Sol'n

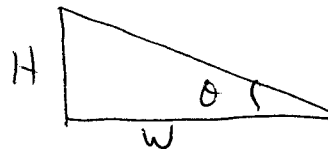
a.) The Coriolis force occurs in a non-inertial reference frame.  $\omega$  is the angular velocity with respect to the inertial frame and  $v$  is the velocity within the rotating frame.

It can be thought of as conservation of angular momentum.

b.) To the west. The west bank will be higher.



$$F = -2m|\vec{\omega} \times \vec{v}| = 2m\omega v \sin \lambda$$



$$\tan \theta = \frac{2m\omega v \sin \lambda}{mg} = \frac{2\omega v \sin \lambda}{g}$$

$$H = w \tan \theta = w \frac{2\omega v \sin \lambda}{g}$$

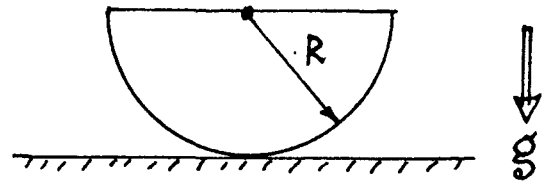
$$\lambda = 45^\circ$$

$$= \frac{2(10^3) \text{ m} (2) \left( 2\pi \frac{\text{rad}}{\text{day } 24 \text{ hr}} \right) \left( \frac{1 \text{ hr}}{3600 \text{ sec}} \right) \left( \frac{1}{\sqrt{2}} \right) \left( \frac{\text{m}}{\text{s}} \right)}{10 \frac{\text{m}}{\text{s}^2}}$$

$$= 0.02 \text{ m} = 2 \text{ cm.}$$

2. A solid glass hemisphere sits on a horizontal surface. Find its frequency for small oscillations.

Note:  $I_{\text{cm}} = \frac{2}{5} MR^2$  for solid sphere.

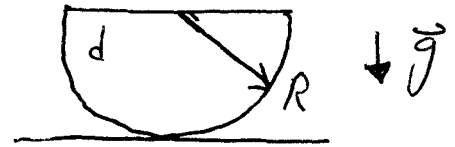


hemisphere has uniform density  $d$ .

## Classical Mechanics - Problem ZHS

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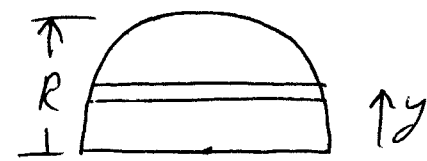
# Classical Mechanics - solution P.1

Method

- 1) Find C.M.
- 2) Find P.E. for small displacement from equilibrium.
- 3) Equate this to KE when passing thru equilibrium position.
- 4) Find frequency from KE.

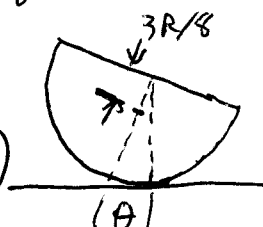
$$1) y_{cm} = \int y dV / \int dV$$

$$= \int_0^R y \pi (R^2 - y^2) dy / \frac{2}{3} \pi R^3$$

$$= (-\pi \frac{R^4}{4} + \pi \frac{R^4}{2}) / \frac{2}{3} \pi R^3 = \frac{3}{8} R$$


$$2) PE = mg(h - h_0)$$

$$= \frac{2}{3} \pi R^3 \rho g \left( -\frac{5}{8} R + R - \frac{3}{8} R \cos \theta \right)$$

$$= mg \frac{3}{16} R \theta^2$$


$$3) KE = \frac{1}{2} I_{cm} \dot{\theta}^2 + \frac{1}{2} m v_{cm}^2 \quad v_{cm} = \frac{5}{8} R \dot{\theta}$$

$$I_{center} = \frac{1}{5} m R^2 = I_{cm} + m \left( \frac{3}{8} R \right)^2 = I_{cm} + \frac{9}{64} m R^2$$

$$I_{cm} = \frac{64 - 45}{320} m R^2 = \frac{19}{320} m R^2$$

$$KE = \frac{19}{640} m R^2 \dot{\theta}^2 + \frac{25}{128} m R^2 \dot{\theta}^2 = \frac{144}{640} m R^2 \dot{\theta}^2 = \frac{9}{40} m R^2 \dot{\theta}^2$$



Classical Mechanics Solution P. 2

$$4) \theta = \theta_0 \sin \omega t \quad \dot{\theta} = \omega \theta_0 \cos \omega t$$

$\Rightarrow \omega \theta_0$  at equilibrium

$$PE = KE; \quad \frac{3}{16} mgR \theta_0^2 = \frac{9}{40} mR^2 \omega^2 \theta_0^2$$

$$\omega = \sqrt{\frac{\frac{3}{16} g}{\frac{9}{40} R}} = \sqrt{\frac{120}{144} \frac{g}{R}} = \sqrt{\frac{5}{6} \frac{g}{R}}$$

Or, using Lagrangian approach,

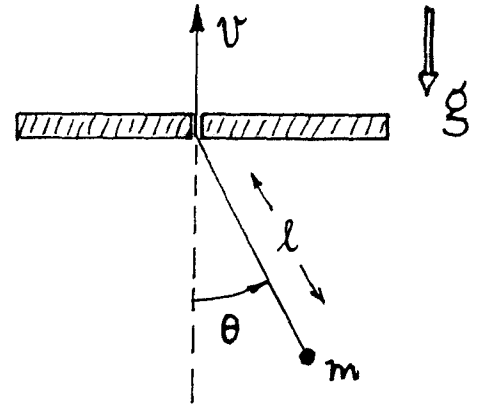
$$L = T - V = \frac{9}{40} mR^2 \dot{\theta}^2 - \frac{3}{16} mgR \theta^2$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = \frac{9}{20} mR^2 \ddot{\theta} + \frac{3}{8} mgR \theta = 0$$

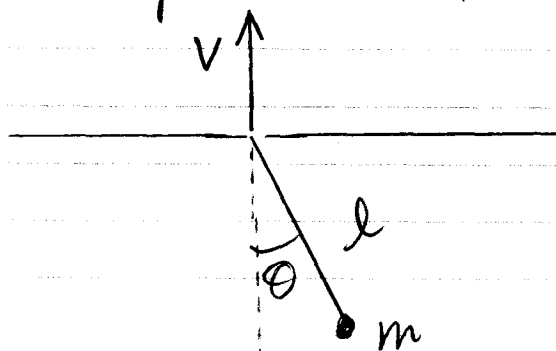
$$\ddot{\theta} + \frac{5}{6} \frac{g}{R} \theta = 0 \quad \theta = \theta_0 \sin \omega t$$

$$\omega = \sqrt{\frac{5}{6} \frac{g}{R}}$$

3. A point particle of mass  $m$  is attached to the end of a massless string which is pulled at a constant velocity  $v$  through a small hole in a table. The mass moves in a vertical plane, with an oscillating angular displacement  $\theta \ll 1$  with respect to the vertical. At  $t=0$  the maximum displacement is  $\theta_0$ , and the length of string below the hole is  $l$ . Determine the maximum displacement for small positive  $t$ .



3. A point particle of mass  $m$  is attached to the end of a massless string which is pulled at a constant velocity  $v$  through a small hole in a table. The mass moves in a vertical plane, with an oscillating angular displacement  $\theta \ll 1$  with respect to the vertical. At  $t=0$  the maximum displacement is  $\theta_0$ , and the length of string below the hole is  $l$ . Determine the maximum displacement for small positive  $t$ .



Soln:  $L = \frac{1}{2} m (v^2 + (l-vt)^2 \dot{\theta}^2) + mg(l-vt) \cos \theta$

Lagrange's Eqn's  $\Rightarrow \frac{d^2}{dt^2} [(l-vt) \theta] + g \sin \theta = 0$

For small  $\theta$ ,  $\frac{d^2}{dt^2} [(l-vt) \theta] + g \theta = 0$

or  $(l-vt) \ddot{\theta} - 2v \dot{\theta} + g \theta = 0$

For small  $t$ ,  $l-vt \approx l$ , and

$$\ddot{\theta} - \frac{2v}{l} \dot{\theta} + \frac{g}{l} \theta = 0$$

or  $\boxed{\ddot{\theta} - b \dot{\theta} + \omega_0^2 \theta = 0}$   $b = \frac{2v}{l}$   
 $\omega_0 = \sqrt{g/l}$   
 ↑ negative damping!

Assume  $\theta \sim e^{(a+i\omega)t}$

$$(a+i\omega)^2 - b(a+i\omega) + \omega_0^2 = 0$$

$$a^2 + 2ia\omega - \omega^2 - ba - ib\omega + \omega_0^2 = 0$$

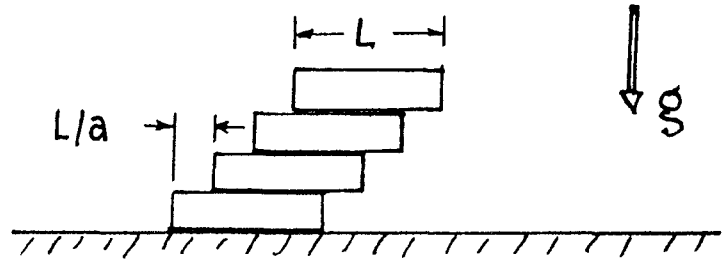
$$\text{Re, Im} = 0 \begin{cases} 2a\omega - b\omega = 0 \Rightarrow \boxed{a = \frac{b}{2}} \\ a^2 - \omega^2 + ab + \omega_0^2 = 0 \Rightarrow \omega = \sqrt{\omega_0^2 - \frac{b^2}{4}} \end{cases}$$

Thus  $\theta_{\text{max}}$  increases exponentially -

$$\theta \sim e^{bt/2} e^{i\omega t}$$

$$\boxed{\theta_{\text{max}} = \theta_0 e^{bt/2} = \theta_0 e^{vt/l}}$$

4. A uniform brick of length  $L$  is laid on a smooth horizontal surface. Other equal bricks are piled as shown, so that the ends are offset at each brick by  $L/a$ , where  $a$  is an integer (the other vertical sides form continuous planes). How many bricks can be stacked in this manner before the pile collapses?



The old Frank Lloyd Wright problem. Do students know enough architecture? RTR

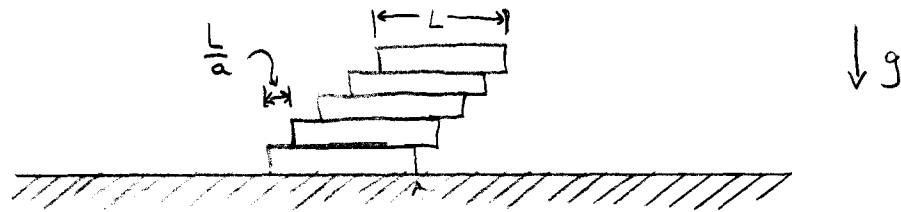
Classical Mechanics

Bill Hirsch

### A Pile of Bricks

Too easy?  
yes.  
Hugo

4. A uniform brick of length  $L$  is laid on a smooth horizontal surface. Other equal bricks are piled as shown, so that the ends are offset at each brick by  $L/a$ , where  $a$  is an integer (the other vertical sides form continuous planes). How many bricks can be stacked in this manner before the pile collapses?



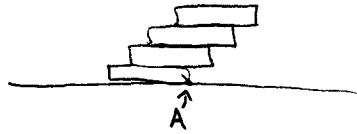
Is hint necessary?  
no

Hint:  $\sum_{i=1}^n i = \frac{n(n+1)}{2}$

Solution:

OK w/br

Evaluate the torque about the point A. Clockwise torques are negative, counterclockwise positive.



Torque of  $i$ th brick about A:

mass of single brick =  $m$

$$\tau_i = mg \left[ -\frac{L}{2} + (i-1)\frac{L}{a} \right]$$

Pile will collapse when total torque about A becomes zero

$$0 = \sum_{i=1}^n \tau_i = \sum_{i=1}^n \frac{mgL}{a} \left[ -\frac{a}{2} + (i-1) \right]$$

$\Rightarrow$

$$0 = \sum_{i=1}^n -\frac{a}{2} + i - 1 = -n\left(\frac{a}{2} + 1\right) + \sum_{i=1}^n i$$

or

$$\sum_{i=1}^n i = n\left(\frac{a}{2} + 1\right)$$

$$\text{but } \sum_{i=1}^n i = \frac{n(n+1)}{2}$$

so

$$\frac{n(n+1)}{2} = \frac{n(a+2)}{2}$$

$n = a + 1$  if bricks are glued together

$n = a$  if bricks are "loose"

(so that bottom brick does not move)

5. A delton is a one-dimensional system in which an electron interacts with an attractive delta function potential whose strength is  $-g$ .
- a) Solve Schrodinger's equation to determine the bound states of the delton.
  - b) Discuss qualitatively under what conditions the delton "negative ion" exists, i.e. a delton having two electrons.



Deltons

5. A delton is a one-dimensional system in which an electron interacts with an attractive delta function potential whose strength is  $-g$ .

(a) Discuss ~~the eigenvalue  $E_0$  and quantum mechanical~~ bound states of the delton by solving Schrödinger's equation.

(a) Solve Schrödinger's equation ~~for the delton~~ to determine the bound states ~~spectrum~~ of the delton.

(b) Discuss qualitatively under what conditions the delton "negative ion" exists, i.e., a delton having two electrons.

Solve  $\left\{ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right\} \psi(x) = E \psi(x)$

Let  $\hbar^2/2m = 1$  and  $V(x) = -g \delta(x)$ .

(1)  $\left( \frac{d^2}{dx^2} + E \right) \psi(x) = -g \delta(x) \psi(x)$

Look for bound states. The homogeneous solns are

$$\left( \frac{d^2}{dx^2} + E \right) \psi(x) = 0$$

$$\psi(x) = \begin{cases} A e^{-\kappa x} & x > 0 \\ A e^{\kappa x} & x < 0 \end{cases}$$

Assume  $\kappa$  is real for b.s.

(2)  $\kappa^2 + E = 0 \quad \text{or} \quad E = -\kappa^2 < 0$

Inhomog. eqn.

$$\left( \frac{d^2}{dx^2} + E \right) \psi(x) = -g \delta(x) \psi(x)$$

$$\lim_{\varepsilon \rightarrow 0} \int_{-\varepsilon}^{\varepsilon} \left( \frac{d^2 \psi}{dx^2} + E \psi \right) dx = -g \psi(0)$$

LHS:  $\psi'_+ - \psi'_- + E(\psi_+ - \psi_-) = -g \psi(0)$

$$\psi'_+ = -\kappa \psi(0)$$

$$\psi'_- = \kappa \psi(0)$$

Thus,  $-2\kappa^2\psi(0) = -g^2\psi(0)$

or  $\kappa = \frac{g}{2}$

(2)  $E = -\kappa^2 = -\frac{g^2}{4} \left( = \frac{\hbar^2}{2m} \cdot \frac{1}{4} \left[ \frac{g}{\hbar^2/2m} \right]^2 \right) = -\frac{mg^2}{2\hbar^2}$

There is only one.

(b) To lowest order, the Coulomb repulsion<sup>exists</sup> of the electrons  
 (3) must be weaker than the sum of the binding energies.

6. A particle of mass  $m$  moves in one dimension in the potential

$$V(x) = \begin{cases} V_0 \cos (\pi/a)x & , 0 < x < a \\ \infty, & \text{otherwise} \end{cases}$$

Find the ground state energy to second order in  $V_0$ , and the ground state wave function to first order in  $V_0$ .

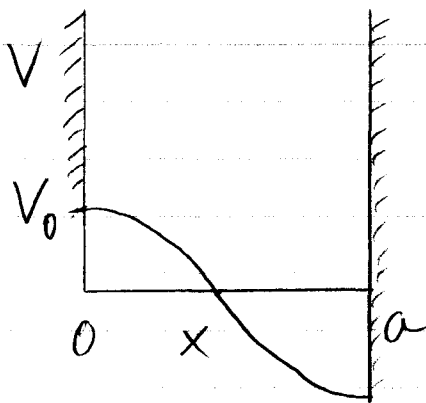
# Quantum Mechanics J. Hermanson

6. A particle of mass  $m$  moves in one dimension in the potential

$$V(x) = \begin{cases} V_0 \cos \frac{\pi}{a} x & , 0 < x < a \\ \infty & , \text{otherwise} \end{cases}$$

Find the ground state energy to second order in  $V_0$ , and the ground state wavefunction to first order in  $V_0$ .

Soln:



$$V = V_0 \cos kx, \quad k \equiv \frac{\pi}{a}$$

$$\psi_1 = \sqrt{\frac{2}{a}} \sin kx \quad \psi_2 = \sqrt{\frac{2}{a}} \sin 2kx$$

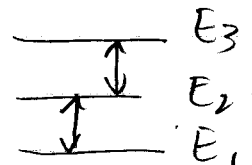
$$E_1 = \frac{\hbar^2 \pi^2}{2ma^2}$$

$$E_2 = 4E_1$$

The perturbation  $V = V_0 \frac{e^{ikx} + e^{-ikx}}{2} \Rightarrow \Delta k = \pm k$   
selection rule

The unperturbed states  $\psi_n \sim \sin nkx$   
 $= \frac{e^{inkx} - e^{-inkx}}{2i}$

Thus  $\psi_1 \leftrightarrow \psi_2 \leftrightarrow \psi_3$  etc.



The non-vanishing matrix elements are

$$V_{12} = V_{21} = V_0 \cdot \frac{2}{a} \int_0^a \sin 2kx \cos kx \sin kx dx$$

$$[V_{11} = 0 = V_{22}] = V_0 \cdot \frac{1}{a} \int_0^a \sin^2 2kx dx = \frac{1}{2} V_0$$

To 2nd order,

$$E_0 = E_1 + V_{11} + \frac{|V_{21}|^2}{E_1 - E_2} = E_1 - \frac{V_0^2}{12E_1}$$

$$E_0 = \frac{\pi^2 \hbar^2}{2ma^2} - \frac{ma^2 V_0^2}{6\pi^2 \hbar^2}$$

To 1st order,

$$\psi_0 = \psi_1 + \frac{V_{21}}{E_1 - E_2} \psi_2 = \psi_1 - \frac{V_0}{6E_1} \psi_2$$

$$\psi_0 = \sqrt{\frac{2}{a}} \sin \frac{\pi}{a} x - \sqrt{\frac{2}{a}} \frac{ma^2 V_0}{3\pi^2 \hbar^2} \sin \frac{2\pi}{a} x$$

7. Classical electromagnetic fields interact with a quantum mechanical particle according to the Schrodinger equation:

$$i\hbar \frac{\partial \Psi}{\partial t} = \left[ \frac{1}{2m} \left( \frac{\hbar}{i} \vec{\nabla} - e \vec{A}(\vec{x}, t) \right)^2 + e\phi(\vec{x}, t) \right] \Psi(\vec{x}, t)$$

Since the wave function  $\Psi$  depends on the electromagnetic potential  $(\phi, \vec{A})$ , it must change if we change the potential by a gauge transformation:

$$\begin{aligned} \vec{A}'(\vec{x}, t) &= \vec{A}(\vec{x}, t) + \vec{\nabla} f(\vec{x}, t) \\ \phi'(\vec{x}, t) &= \phi(\vec{x}, t) - \frac{1}{c} \frac{\partial f}{\partial t} \end{aligned}$$

(the electromagnetic fields,  $\vec{E}$  and  $\vec{B}$ , are of course, unchanged by this transformation), where  $f(\vec{x}, t)$  is an arbitrary function of  $\vec{x}$  and  $t$ . The wave function  $\Psi$  in the new gauge obeys

$$i\hbar \frac{\partial \Psi'}{\partial t} = \left[ \frac{1}{2m} \left( \frac{\hbar}{i} \vec{\nabla} - e \vec{A}' \right)^2 + e\phi' \right] \Psi'(\vec{x}, t),$$

$$\text{where } \Psi' = \exp \left[ \frac{ie}{\hbar c} f(\vec{x}, t) \right] \Psi$$

- Compare  $\langle \Psi | \vec{x} | \Psi \rangle$  and  $\langle \Psi' | \vec{x} | \Psi' \rangle$ . Are they equal?
- Compare  $\langle \Psi | \vec{p} | \Psi \rangle$  and  $\langle \Psi' | \vec{p} | \Psi' \rangle$ . Show that they are not equal.
- Does (b) imply that quantum mechanics is not gauge invariant? Is  $\vec{p}$  an observable? (Hint: consider the operator equation of motion for  $\vec{x}$ ).

7. Classical electromagnetic fields interact with a quantum mechanical particle according to the Schrodinger equation:

I'll do equations

$$\left\{ i\hbar \frac{\partial \Psi}{\partial t} = \left[ \frac{1}{2m} \left( \frac{\hbar}{i} \vec{\nabla} - \frac{e}{c} \vec{A}(\vec{x}, t) \right)^2 + e\phi(\vec{x}, t) \right] \Psi(\vec{x}, t) \right\}$$

Since the wave function  $\Psi$  depends on the electromagnetic potential  $(\phi, \vec{A})$ , it must change if we change the potential by a gauge transformation:

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$\Psi'(\vec{x}, t)$  is related to  $\Psi(\vec{x}, t)$  by

$$\Psi' = \exp \left[ \frac{ie}{\hbar c} f(\vec{x}, t) \right] \Psi$$

(a) Compare  $\langle \Psi | \vec{x} | \Psi \rangle$  and  $\langle \Psi' | \vec{x} | \Psi' \rangle$ . Are they equal?

(b) Compare  $\langle \Psi | \vec{p} | \Psi \rangle$  and  $\langle \Psi' | \vec{p} | \Psi' \rangle$ . Show that they are not equal.

(c) Does (b) imply that quantum mechanics is not gauge invariant?

Is  $\vec{p}$  an observable? (Hint: consider the operator equation of motion for  $\vec{x}$ ).



Solution.

$$(a) \langle \psi | \vec{x} | \psi \rangle = \int d^3x \psi^* \vec{x} \psi$$

$$\langle \psi' | \vec{x} | \psi' \rangle = \int d^3x \psi^* e^{-\frac{ie}{\hbar c} F} \vec{x} e^{\frac{ie}{\hbar c} F} \psi = \int d^3x \psi^* \vec{x} \psi = \langle \psi | \vec{x} | \psi \rangle$$

They are the same:  $\langle \vec{x} \rangle$  is gauge invariant

$$(b) \langle \psi | \vec{p} | \psi \rangle = \int d^3x \psi^* (-i\hbar \vec{\nabla}) \psi$$

$$\langle \psi' | \vec{p} | \psi' \rangle = \int d^3x \psi^* e^{-\frac{ie}{\hbar c} F} (-i\hbar \vec{\nabla}) e^{\frac{ie}{\hbar c} F} \psi$$

$$= \int d^3x \psi^* (-i\hbar \vec{\nabla}) \psi + \int d^3x \psi^* \left( \frac{e}{c} \vec{\nabla} F \right) \psi$$

$$= \langle \psi | \vec{p} | \psi \rangle + \int d^3x \psi^* \left( \frac{e}{c} \vec{\nabla} F \right) \psi$$

They are not equal!  $\langle \vec{p} \rangle$  is not gauge invariant! However, this does not mean quantum mechanics is not gauge invariant: gauge invariance only requires that all observable physical quantities must be gauge invariant. The physically observable momentum is  $m d\vec{x}/dt$ : from the equation of motion for  $\vec{x}$ :

$$i\hbar \frac{d}{dt} [\vec{x}, H] = -\frac{1}{m} \frac{\hbar}{i} \left( \frac{\hbar}{i} \vec{\nabla} - \frac{e}{c} \vec{A} \right) = +\frac{i\hbar}{m} (\vec{p} - \frac{e}{c} \vec{A})$$

$$\Rightarrow \boxed{m \frac{d\vec{x}}{dt} = (\vec{p} - \frac{e}{c} \vec{A})} \text{ as in classical mechanics}$$

Note that

$$\langle \psi | \vec{p} - \frac{e}{c} \vec{A} | \psi \rangle = \langle \psi' | \vec{p} - \frac{e}{c} \vec{A} | \psi' \rangle \text{ is the observable quantity is}$$

gauge invariant.

and

$$\langle V_e | \Psi(t, x) \rangle = \exp(-i \frac{E_0 t}{\hbar}) \left\{ \cos^2 \theta \exp\left(\frac{i \sqrt{2m_1} E_0}{\hbar} x\right) + \sin^2 \theta \exp\left(\frac{i \sqrt{2m_2} E_0}{\hbar} x\right) \right\}$$

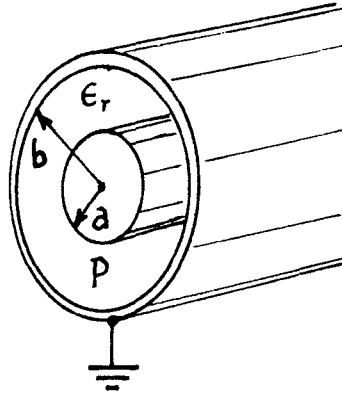
$$|\langle V_e | \Psi \rangle|^2 = \cos^4 \theta + \sin^4 \theta + \cos^2 \theta \sin^2 \theta \left\{ \exp\left[i \frac{\sqrt{2E_0}}{\hbar} (\sqrt{m_1} - \sqrt{m_2}) x\right] - \exp\left[-i \frac{\sqrt{2E_0}}{\hbar} (\sqrt{m_1} - \sqrt{m_2}) x\right] \right\}$$

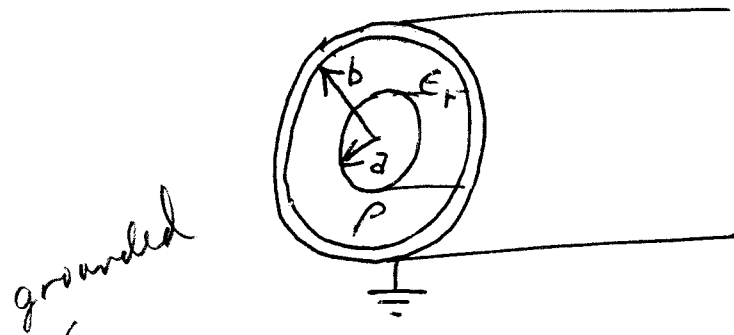
$$= \cos^4 \theta + \sin^4 \theta + 2 \cos^2 \theta \sin^2 \theta \cos\left[\frac{\sqrt{2E_0}}{\hbar} (\sqrt{m_1} - \sqrt{m_2}) x\right]$$

$$|\langle V_e | \Psi \rangle|^2 = 1 + \frac{1}{2} \sin 2\theta \left\{ \cos\left[\frac{\sqrt{2E_0}}{\hbar} (\sqrt{m_1} - \sqrt{m_2}) x\right] - 1 \right\}$$

Note that if either  $\theta \rightarrow 0$  or  $m_1 \rightarrow m_2$ , the oscillations vanish.

8. A coaxial cable of inner conductor radius  $a$ , outer conductor inside radius  $b$ , and relative dielectric permittivity  $\epsilon_r$  is open-circuited with initial voltage 0 on the inner conductor. It receives a burst of ionizing radiation which leaves the dielectric with a uniform free volume charge density  $P$ . Find the voltage on the inner conductor if the outer conductor is grounded.





8. A coaxial cable of ~~inside~~ <sup>inner</sup> conductor radius  $a$ , outer conductor inside radius  $b$ , and ~~dielectric~~ relative dielectric permittivity  $\epsilon_r$  is open-circuited with initial voltage 0 on the inner conductor. It receives a burst of ionizing radiation which leaves the dielectric with <sup>(a uniform)</sup> ~~a free~~ volume charge density  $\rho$ . Find the voltage on the inner conductor.

R

## EM solution

Use Gauss' Law:  $\int \frac{dq_{\text{enc}}}{\epsilon_0 \epsilon_r} = \int \vec{E} \cdot d\vec{A}$

$$\frac{\rho \pi (r^2 - a^2) l}{\epsilon_0 \epsilon_r} = 2\pi r l E$$

$$E = \frac{\rho (r^2 - a^2) l}{2\pi \epsilon_0 \epsilon_r r l}$$

$$\vec{E} = \frac{\rho (r^2 - a^2)}{2\epsilon_0 \epsilon_r r} \hat{r}$$

$$V = - \int_b^a \vec{E} \cdot d\vec{r} = + \int_a^b E dr$$

$$V = \frac{\rho}{2\epsilon_0 \epsilon_r} \int_a^b \left( r - \frac{a^2}{r} \right) dr = \boxed{\frac{\rho}{2\epsilon_0 \epsilon_r} \left( \frac{b^2 - a^2}{2} - a^2 \ln \frac{b}{a} \right)}$$

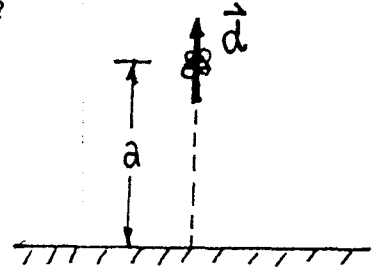
DEPARTMENT OF PHYSICS

M.S. COMPREHENSIVE/PH.D. QUALIFYING EXAMINATION

SATURDAY, APRIL 2, 1-5 PM / 1986

Answer each of the following questions, all of which carry equal weight. Begin your answer to each question on a new sheet of paper; solutions to different questions must not appear on the same sheet. Each sheet of paper must be labeled with your name and the problem number, in the upper right hand corner. When more than one sheet is submitted for a problem, be sure the pages are ordered properly. [Problems 9-16].

9. Consider a carbon monoxide molecule positioned a distance  $a$  above an ideal metal surface. The molecule has a dipole moment  $\vec{d}$ . Assuming that  $\vec{d}$  is perpendicular to the surface and parallel to the surface normal, what, if any, force exists on the molecule?



E &amp; M

good - maybe better for  
Ph.D. Comp. JED.

WJ

OK/WJ

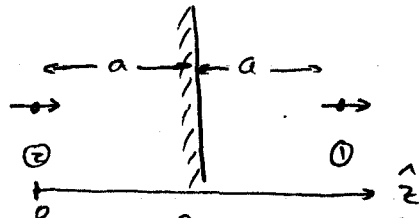
9. Consider a carbon monoxide molecule positioned a distance  $a$  above an ideal metal surface. The molecule has a dipole moment  $\vec{p}$ . Assuming that  $\vec{p}$  is perpendicular to the surface and parallel to the surface normal, what, if any, force exists on the molecule?

Soln

General dipole potl:  $\varphi = \frac{\vec{d} \cdot \vec{r}}{r^3} = \frac{d \cos \theta}{r^2}$

$$\vec{E} = (-\nabla \varphi) = -\hat{e}_r \frac{\partial \varphi}{\partial r} - \frac{\hat{e}_\theta}{r} \frac{\partial \varphi}{\partial \theta} \quad (\text{axial symmetry})$$

$$\vec{E} = \frac{2d \cos \theta}{r^3} \hat{e}_r + \frac{d \sin \theta}{r^3} \hat{e}_\theta$$



Let  $E_z$  be the image field. Then the interaction energy is

$$U_{\text{int}} = -\vec{p}_1 \cdot \vec{E}_2$$

$$E_z = \frac{2d_2}{z^3} \hat{e}_z$$

$$U_{\text{int}} = -\frac{2d_1 d_2}{z^3}$$

$$F = (-\nabla U) = + \frac{2d^2}{z^4} (-3) \hat{e}_z = -\frac{6d^2}{z^4} \hat{e}_z$$

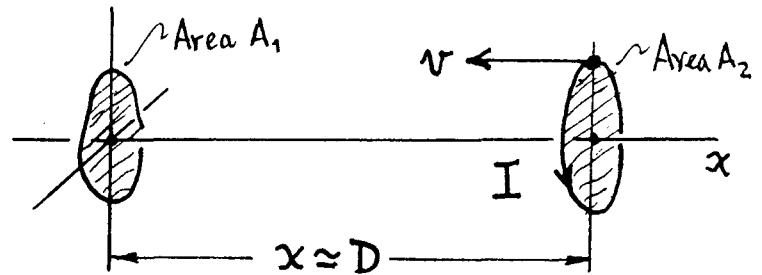
$$F = -\frac{6}{24} \frac{d^2}{a^4} \hat{e}_z = -\frac{3}{8} \frac{d^2}{a^4} \hat{e}_z$$



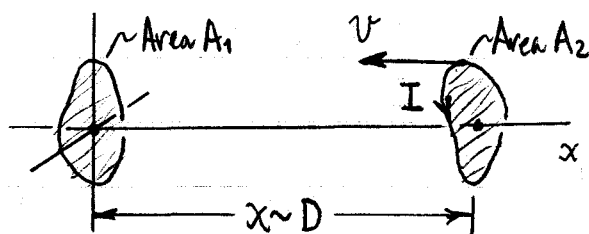
10. Two coaxial, coplanar wire loops of areas  $A_1$  &  $A_2$  are situated on the  $x$ -axis as shown, and are initially separated by a "large" distance  $D$  ( $D \gg$  linear dimension of either loop). Loop 1 remains fixed at the origin, while loop 2 carries a steady current  $I$  and is made to move at a known velocity  $v = dx/dt$  to the left (toward loop 1) along the  $x$ -axis.

- a) Calculate the emf induced in loop 1 by the movement of loop 2, so long as the loop separation  $x \sim D$  is large.
- b) If the current flow is counterclockwise in loop 2, in what direction does the induced current flow in loop 1?

Hint: Use the dipole approximation.



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A. Calculate the emf induced in loop 1 by the movement of loop 2, so long as the loop separation  $x \sim D$  is "large".

B. If the current flow is counterclockwise in loop 2, in what direction does the induced current flow in loop 1?

HINT: use the dipole approximation.

### Solution

1.  $D \rightarrow$  large  $\Rightarrow$  dipole approxn can be used for the  $B$ -fld of loop # 2:

$$\mathbf{B} = \frac{1}{x^3} [3\mathbf{m}(\mathbf{m} \cdot \hat{\mathbf{m}}) - \mathbf{m}] \quad \begin{matrix} \mathbf{m} = (-) \hat{\mathbf{x}}, \text{ unit vector from \# 2 to \# 1, (1)} \\ \mathbf{m} = \left(\frac{I}{c} A_2\right) \hat{\mathbf{x}}, \text{ mag. moment of \# 2.} \end{matrix}$$

(We are using cgs units, as any sensible person would do). For the given geometry, this gives the field at loop 1 due to loop 2...

$$\mathbf{B} = \left(\frac{2IA_2}{cx^3}\right) \hat{\mathbf{x}} \leftarrow \text{points to the right, and increases as } x \text{ decreases. (2)}$$

2. With  $x \sim D \rightarrow$  large (compared to loop # 1 dimensions),  $B$  is  $\sim$  const over the area  $A_1$ , so the flux through loop # 1 is

(over)

$$\phi = BA_1 = (2IA_1A_2/c) \frac{1}{x^3} \quad (3)$$

Faraday's law then provides the induced emf...

$$\mathcal{E} = - \frac{d\phi}{dt} = - \underbrace{\left( \frac{dx}{dt} \right)}_{v, \text{ velocity of loop \#2}} \frac{d\phi}{dx} = -v \left( \frac{2IA_1A_2}{c} \right) \frac{d}{dx} \left( \frac{1}{x^3} \right)$$

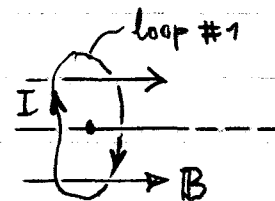
or

$$\boxed{\mathcal{E} = 6vIA_1A_2/cx^4}$$

(4)

This is the required emf in #1 due to #2's motion, per part A.

3. As noted in Eq. (2), B is to the right at loop #1, and is increasing with time. By Lenz's law, the induced I will flow in a direction so as to oppose this change... I flows to create a B' in the opposite direction. So...



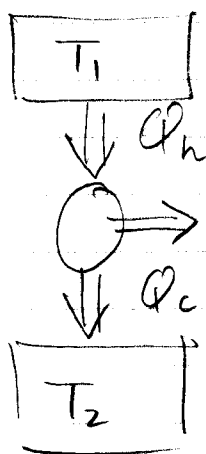
If I flows CCW in loop #2, the induced I flows CW in loop #1.

11. Two identical bodies with heat capacities  $C_p$  and  $C_v$  are at different initial temperatures  $T_1$  and  $T_2 < T_1$ . While remaining at constant pressure the bodies are used as reservoirs for a heat engine. What is the maximum amount of work obtainable?

Hint: allow the two bodies to come to a common final temperature.

11. Two identical bodies with heat capacities  $C_p$  and  $C_v$  are at different initial temperatures  $T_1$  and  $T_2 < T_1$ . While remaining at constant pressure the bodies are used as reservoirs for a heat engine. What is the maximum amount of work obtainable? Hint: allow the two bodies to come to a common final temperature  $T_f$ .

Soln:



Here  $T_1 > T_f > T_2$

$$W = Q_h - Q_c$$

$$= C_p (T_1 - T_f) - C_p (T_f - T_2)$$

$$= C_p (T_1 + T_2 - 2T_f)$$

What is  $T_f$  for max  $W$ ?

The maximum efficiency (max  $W$ ) is obtained when  $\Delta S = 0$  for the complete cycle.

$$\text{Now } \Delta S = C_p \ln \frac{T_f}{T_1} + C_p \ln \frac{T_f}{T_2} \quad \text{for a cycle}$$

$$= C_p \ln \frac{T_f^2}{T_1 T_2} = 0 \quad \text{when } T_f = \sqrt{T_1 T_2}$$

Thus 
$$W_{\max} = C_p (T_1 + T_2 - 2\sqrt{T_1 T_2})$$

12. A careless experimenter left the valve on a tank of helium gas slightly open over the weekend. The gas, originally at 100 atm. slowly escaped isothermally at 300 K. What was the change in entropy per kilogram of helium?

Boltzmann's constant  $k = 1.38 \times 10^{-23} \text{ J/K}$

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Boltzmann's constant  $k = 1.38 \times 10^{-23} \text{ J/K}$

Solution:

$$\Delta S = \int_a^b \frac{dQ}{T} = \frac{\Delta Q}{T} \text{ since isothermal}$$

isothermal expansion  $\Rightarrow dU = 0$      $\Delta Q = \int_a^b P dV = - \int_{P_1}^{P_2} V dP = - \int_{P_1}^{P_2} \frac{NkT dP}{P}$

$$\Delta Q = NkT \log\left(\frac{P_1}{P_2}\right) \Rightarrow \Delta S = Nk \log\left(\frac{P_1}{P_2}\right)$$

$$\frac{\Delta S}{\text{kg}} = N_{\text{kg}} k \log\left(\frac{P_1}{P_2}\right)$$

$$1 \text{ kg He} = 250 \text{ moles} \approx 1.5 \times 10^{26} \text{ He atoms}$$

$$\frac{\Delta S}{\text{kg}} \approx (1.5 \times 10^{26}) \left(1.38 \times 10^{-23} \frac{\text{J}}{\text{K}}\right) \log_e(100)$$

$$\boxed{\frac{\Delta S}{\text{kg}} = 9.5 \times 10^3 \text{ J/K}}$$

13.  $T(a)$  is a translation operator which converts  $\psi(x)$  to  $\psi(x+a)$ ;

$$T(a) \psi(x) = \psi(x+a) \quad .$$

Show that  $T(a) = \exp(iap_x)$ , where  $p_x$  is the quantum mechanical momentum operator,

$$p_x = -i(d/dx) \quad .$$



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$$T(a)\psi(x) = \psi(x+a)$$

Show that  $T(a) = \exp(ia\hat{p}_x)$ , where  $\hat{p}_x$  is the quantum mechanical momentum operator,

$$\hat{p}_x = -i\hbar \frac{d}{dx},$$

Solution:

$$T(a) = \exp(ia\hat{p}_x) = \exp\left(a\frac{d}{dx}\right) = 1 + a\frac{d}{dx} + \frac{1}{2}a^2\left(\frac{d}{dx}\right)^2 + \frac{1}{3!}a^3\left(\frac{d}{dx}\right)^3 + \dots$$

$$T(a)\psi(x) = \psi(x) + a\frac{d\psi}{dx} + \frac{1}{2}a^2\frac{d^2\psi}{dx^2} + \frac{1}{3!}a^3\frac{d^3\psi}{dx^3} + \dots$$

$$= \psi(x+a) \quad \blacksquare$$

14. A proton-antiproton pair may be created in the absorption of a photon ( $\gamma$ ) by a proton at rest:

$$\gamma + p \rightarrow p + p + \bar{p}$$

The threshold energy for this reaction corresponds to the three particles on the right moving off together as a single particle of rest mass  $3m_p$ . What is the threshold energy  $E_\gamma$  of the photon (in terms of  $m_p c^2$ )?

Stfwd, but ok. // RTR

## Special Relativity - Easy

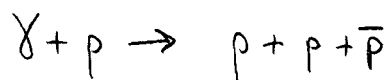
Bill Hirsch

standard:  
= e minor  
problem more.

OK RTR

OK WBS

14. A proton-antiproton pair may be created in the absorption of a photon ( $\gamma$ ) by a proton at rest:



The threshold energy for this reaction corresponds to the three particles on the right moving off together as a single particle of rest mass  $3m_p$ . What is the threshold energy  $E_\gamma$  of the photon (in terms of  $m_p c^2$ )?

Solution:

conservation of momentum

$$p_\gamma = p_{3p} \quad (1)$$

conservation of energy

$$E_\gamma + m_p c^2 = E_{3p} \quad (2)$$

but  $E_\gamma = p_\gamma c$   $p_\gamma = E_\gamma / c \xrightarrow{(1)} p_{3p} = E_\gamma / c$

and

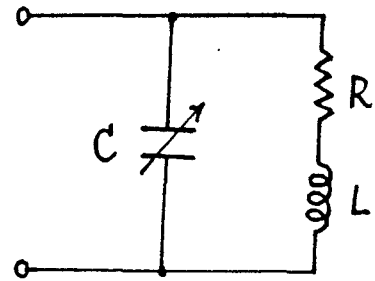
$$\begin{aligned} E_{3p}^2 &= (3m_p c^2)^2 + (p_{3p} c)^2 \xrightarrow{(2)} E_\gamma^2 + 2E_\gamma m_p c^2 + m_p^2 c^4 \\ &= 9m_p^2 c^4 + E_\gamma^2 = E_\gamma^2 + 2E_\gamma m_p c^2 + m_p^2 c^4 \\ &\Rightarrow \boxed{E_\gamma = 4m_p c^2} \end{aligned}$$

15. Describe how you would generate and measure electromagnetic radiation in all different frequency ranges.

- 
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-

|           | generator                     | vac's               |
|-----------|-------------------------------|---------------------|
| Audio     | transistors                   | tuned circuits      |
| r.f.      | "                             | "                   |
| VHF       | "                             | diodes              |
| microwave | klystrons                     | diodes, thermistors |
| far i.r.  |                               |                     |
| ir.       | heat                          | thermistors, film   |
| visible   | filament, gaseous discharge   | eye, film, "        |
| u.v.      |                               |                     |
| X-ray     | hi-voltage electrons on metal | film                |
| X-ray     | nuclear excitation            | secondary electrons |

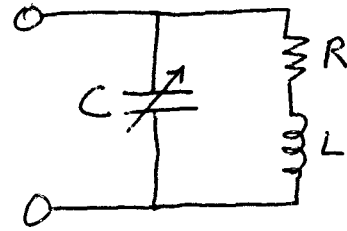
16. To what value must  $C$  be adjusted to achieve maximum impedance at angular frequency  $\omega$ ?



## Experimental Problem

VHS

16. To what value must  $C$  be adjusted to achieve maximum impedance at angular frequency  $\omega$ ?





Experimental solution

$$Z_{RL} = R + j\omega L$$

$$Y_{RL} = \frac{1}{Z_{RL}} = \frac{R - j\omega L}{R^2 + \omega^2 L^2}$$

$$Y = Y_{RL} + Y_C = \frac{R - j\omega L}{R^2 + \omega^2 L^2} + j\omega C$$

$Y$  is smallest when real, so

$$C = \frac{L}{R^2 + \omega^2 L^2}$$