Dipolar stability in spherical simulations: the impact of an inner stable zone

Bonnie R. Zaire¹ and Laurene Jouve²

IRAP, Université de Toulouse, CNRS / UMR 5277, CNES, UPS, 14 avenue E. Belin, Toulouse, F-31400 France

¹email: bzaire@irap.omp.eu; ²email: ljouve@irap.omp.eu

Abstract.

Keywords. Dynamo, Stably stratified, Numerical simulation, etc.

1. Introduction

2. Governing equations

We perform 3D magnetohydrodynamic simulations of a stratified fluid in a spherical shell of outer radius r_0 and inner radius r_i . We use the anelastic version of the code MagIC to solve the non-dimensional equations that govern convective motions and magnetic field generation.

$$\frac{\partial \vec{\mathbf{u}}}{\partial t} + \vec{\mathbf{u}} \cdot \nabla \vec{\mathbf{u}} + \frac{2}{E} \hat{\mathbf{e}}_z \times \vec{\mathbf{u}} = -\nabla p^* + \frac{Ra}{Pr} g s' \hat{\mathbf{e}}_r + \frac{1}{Pm E \bar{\rho}} (\nabla \times \vec{\mathbf{B}}) \times \vec{\mathbf{B}} + \vec{\mathbf{F}}_{\nu}, \quad (2.1)$$

$$\nabla \cdot (\bar{\rho}\vec{\mathbf{u}}) = 0, \tag{2.2}$$

$$\bar{\rho}\bar{T}\left(\frac{\partial s'}{\partial t} + \vec{\mathbf{u}}\cdot\boldsymbol{\nabla}s'\right) + u_{r}\frac{\mathrm{d}\bar{s}}{\mathrm{d}r} = \frac{1}{\mathrm{Pr}}\boldsymbol{\nabla}\cdot\left(\bar{\kappa}\bar{\rho}\bar{T}\boldsymbol{\nabla}s'\right) + \frac{\mathrm{Pr}\,\mathrm{Di}}{\mathrm{Ra}}Q_{\nu} + \frac{\mathrm{Pr}\,\mathrm{Di}}{\mathrm{Pm}^{2}\,\mathrm{E}\,\mathrm{Ra}}\lambda(\boldsymbol{\nabla}\times\vec{\mathbf{B}})^{2}, \tag{2.3}$$

$$\frac{\partial \vec{\mathbf{B}}}{\partial t} = \nabla \times \left(\vec{\mathbf{u}} \times \vec{\mathbf{B}} \right) - \frac{1}{Pm} \nabla \times \left(\lambda \nabla \times \vec{\mathbf{B}} \right), \tag{2.4}$$

$$\nabla \cdot \vec{\mathbf{B}} = 0, \tag{2.5}$$

Our reference state is represented by an ideal gas nearly adiabatic, given by

$$\frac{1}{\overline{T}}\frac{\partial \overline{T}}{\partial r} = \epsilon_{\rm s} \frac{\mathrm{d}\overline{s}}{\mathrm{d}r} - \operatorname{Di}\overline{\alpha}g(r), \tag{2.6}$$

$$\frac{1}{\bar{\rho}}\frac{\partial \bar{\rho}}{\partial r} = \epsilon_{\rm s}\frac{\mathrm{d}\bar{s}}{\mathrm{d}r} - \frac{\mathrm{Di}\,\bar{\alpha}}{\Gamma}g(r) \tag{2.7}$$

3. Implications

References

Amari, S., Hoppe, P., Zinner, E., & Lewis R.S. 1995, Meteoritics, 30, 490