Electromagnetism

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Lecture 13
Magnetic Fields from Currents
Biot-Savart Law
Week 7

Last Lecture

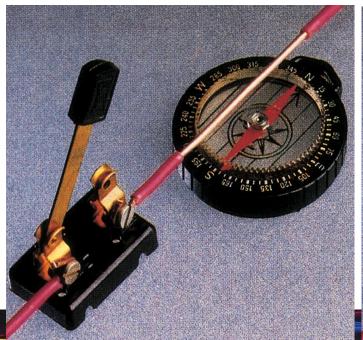
- Special cases of magnetic force
- Force on current carrying conductor
 - $\underline{F} = I \underline{l} \wedge \underline{B}$
- Current Loops and Magnetic Dipoles
 - $\underline{\mu} = I\underline{A}$
- Torque on magnetic dipole in B-field
 - $\underline{\tau} = \underline{\mu} \wedge \underline{B}$
- Potential energy of magnetic dipole in B-field
 - $U = -\boldsymbol{\mu} \cdot \underline{\boldsymbol{B}}$

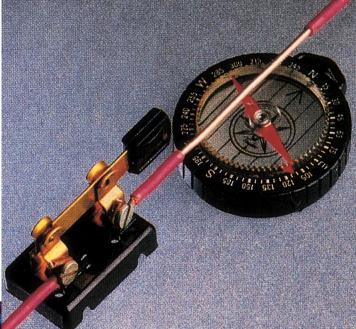
This Lecture

- Magnetic field from moving charge
- Magnetic field from current element
- Biot-Savart Law
 - B-Field at centre of current loop (magnetic dipole)
 - B-field from line of current
 - B-field from infinite line of current
 - B-field along axis of current loop (magnetic dipole)
 - B-field and E-field from dipoles

Magnetic Field From Current

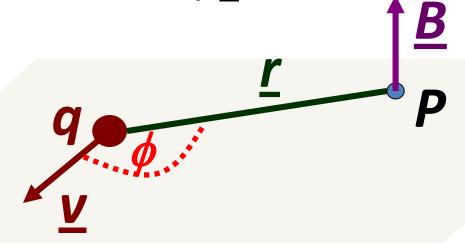
- Hans Christian Oersted, Danish Physicist (1777-1851)
- In 1820 Oersted demonstrated that a magnetic field exists near a current-carrying wire - first connection between electric and magnetic phenomena.





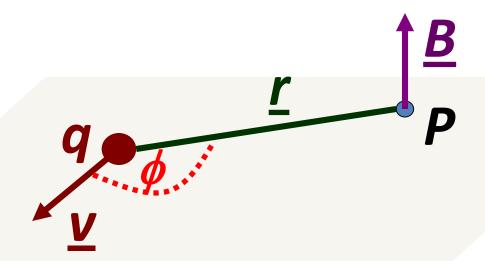
Magnetic Field from Moving Charge

• Consider a B-field at a point \underline{r} from a charge q moving with velocity \underline{v} .



• $B \propto \frac{qv \sin \phi}{r^2}$ perpendicular to \underline{v} and \underline{r} .

Magnetic Field from Moving Charge



• In vector form: $\underline{\boldsymbol{B}} \propto \frac{q}{r^2} \underline{\boldsymbol{v}} \wedge \hat{\underline{\boldsymbol{r}}}$

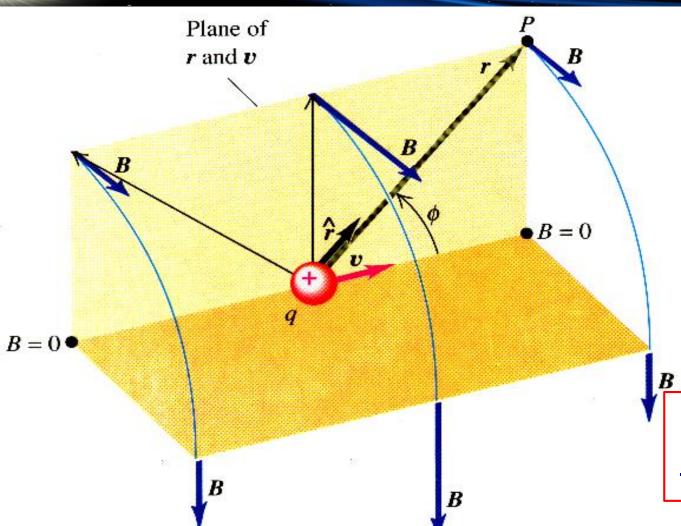
Magnetic Field from Moving Charge

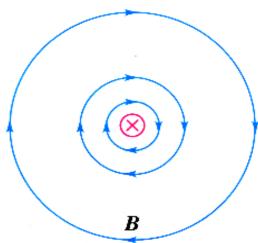
- In SI units
- In vector form: $\underline{\underline{B}} = \frac{\mu_0}{4\pi} \frac{q}{r^2} \underline{\underline{v}} \wedge \hat{\underline{r}}$
- μ_0 is called the **Permeability of free space**.

•
$$\mu_0 = 4\pi \times 10^{-7} T m A^{-1}$$

$$P$$

Magnetic Field from Moving Charge

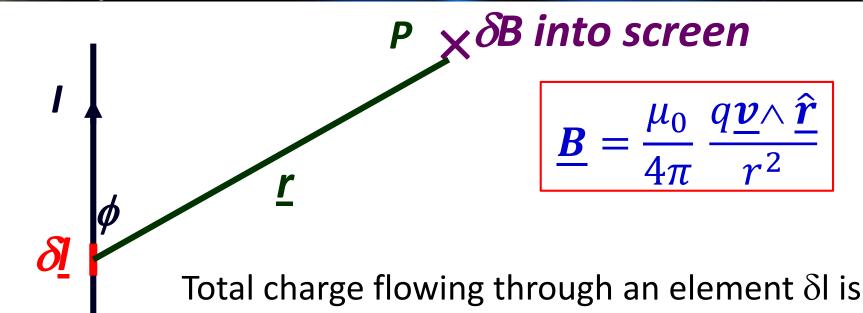




Direction of B same a screw

$$\underline{\boldsymbol{B}} = \frac{\mu_0}{4\pi} \, \frac{q\underline{\boldsymbol{v}} \wedge \hat{\underline{\boldsymbol{r}}}}{r^2}$$

Magnetie Field due to Current Element



So, $q\underline{\boldsymbol{v}} = \frac{q\boldsymbol{v}}{\delta l}\delta\underline{\boldsymbol{l}} = I\delta\underline{\boldsymbol{l}}$ where $\delta\underline{\boldsymbol{l}}$ is defined to be in the same direction as the current.

Magnetie Field due to Current Element

Plugging this in, gives us the Biot Savart Law:

Biot Savart Law

$$\delta \mathbf{B} = \frac{\mu_0}{4\pi} \frac{I \, \delta \mathbf{l} \wedge \hat{\mathbf{r}}}{r^2}$$

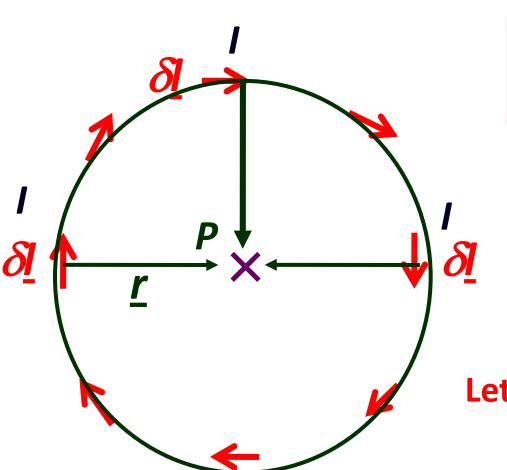
I'm sorry, this is just one of those few equations you need to learn.

The Magnetic Fields from Circuits

Procedure

- Write down dB in terms of a single variable
- Integrate between the limits applicable to the problem
- Be careful about the directions of the vector quantities
- Use symmetry to simplify the problem

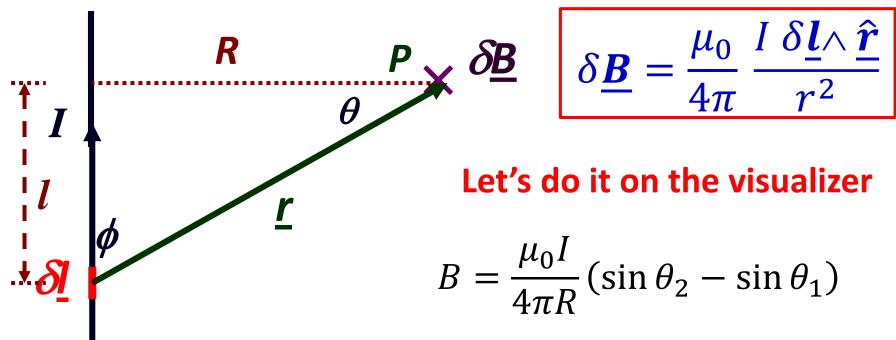
B-field in Centre of Current Loop (magnetic dipole)



$$\delta \underline{\mathbf{B}} = \frac{\mu_0}{4\pi} \, \frac{I \, \delta \underline{\mathbf{l}} \wedge \hat{\underline{\mathbf{r}}}}{r^2}$$

Let's do it on the visualizer

B-field from a Line of Gurrent



$$\delta \underline{\boldsymbol{B}} = \frac{\mu_0}{4\pi} \, \frac{I \, \delta \underline{\boldsymbol{l}} \wedge \hat{\boldsymbol{r}}}{r^2}$$

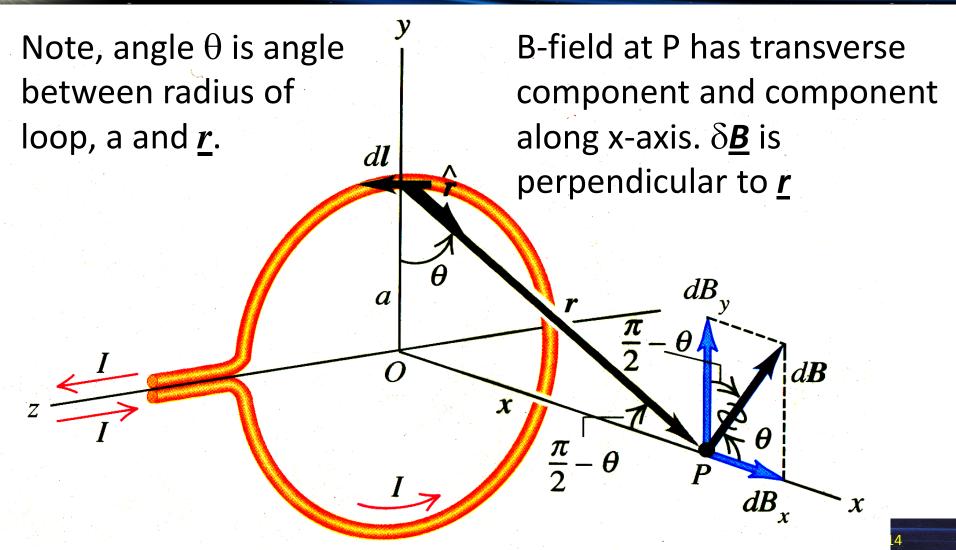
Let's do it on the visualizer

$$B = \frac{\mu_0 I}{4\pi R} (\sin \theta_2 - \sin \theta_1)$$

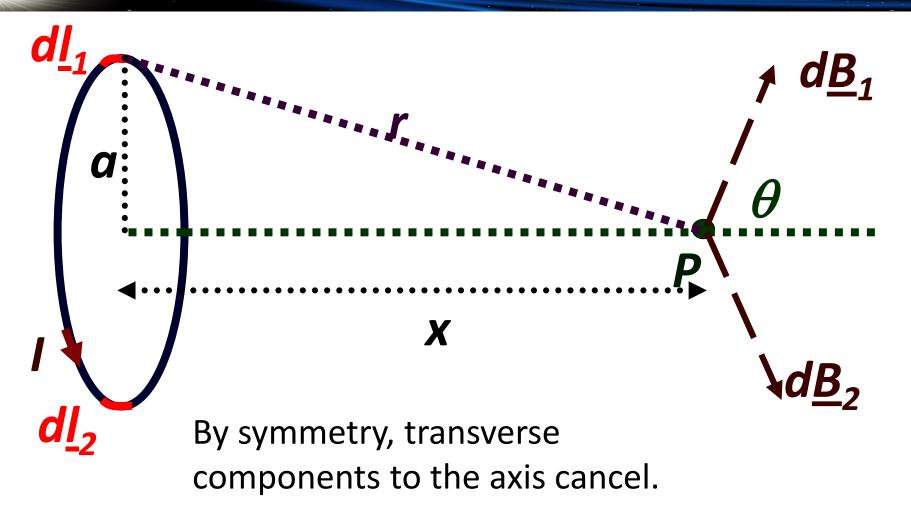
For infinite line of current

$$B = \frac{\mu_0 I}{2\pi R}$$

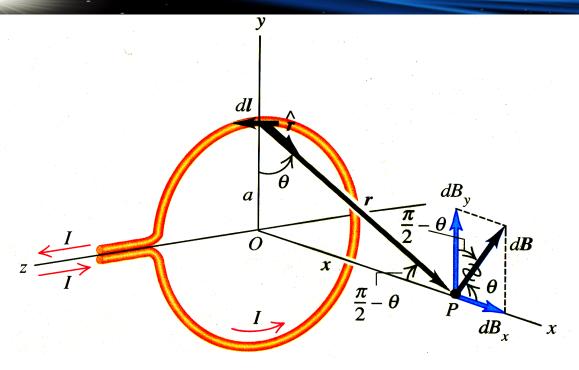
B at any Point on the Axis of a Single Current Loop



B at any Point on the Axis of a Single Gurrent Loop



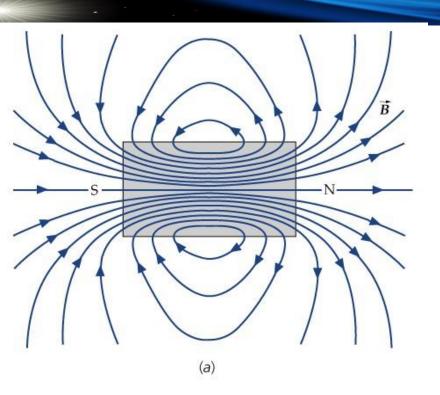
B at any Point on the Axis of a Single Gurrent Loop



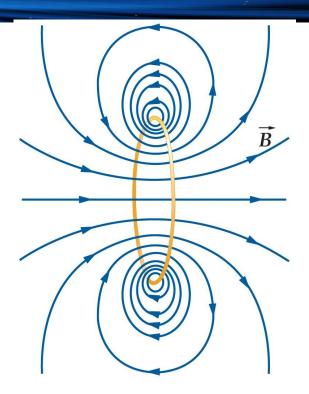
Do problem on visualizer.

Result:
$$B_{\chi} = \frac{\mu_0 I \, a^2}{2(\chi^2 + a^2)^{3/2}}$$

B-field from Current Loop



B-field from bar magnet



B-field from current loop (magnetic Dipole)

Exercise Time

- Obtain an expression for the magnetic field
- (i) At the centre of the loop

Use previous result:
$$B_{\chi} = \frac{\mu_0 I a^2}{2(\chi^2 + a^2)^{3/2}}$$

Hence
$$B_0 = \frac{\mu_0 I}{2a}$$
 (as shown in Ex 13.1)

(ii) At x >> a

Here we see:
$$B_x \to \frac{\mu_0 I \ a^2}{2x^3}$$
 (x >> a)

B-field from Magnetic Dipole (x >> a)

For x >> a (usually the case for atoms)

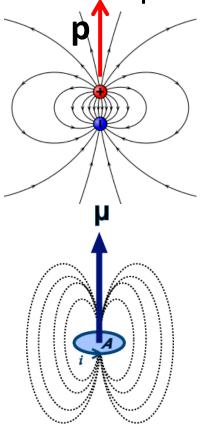
$$\bullet \ B_{\chi} = \frac{\mu_0 I \ a^2}{2\chi^3}$$

- But magnetic dipole moment
- $\mu = I \times area = I \pi a^2$ Hence

•
$$B_{\chi} = \frac{\mu_0}{4\pi} \frac{2\mu}{\chi^3}$$
 (don't get μ_0 and μ confused)

Magnetic and Electric Dipole Fields

Compare B-field from magnetic dipole to E-field from electric dipole



Electric Dipole (Lecture 7)

$$E_{\chi} = \frac{1}{4\pi\varepsilon_0} \frac{2p}{\chi^3}$$

Magnetic Dipole

$$B_{\chi} = \frac{\mu_0}{4\pi} \frac{2\mu}{\chi^3}$$

Review - Biot-Savart Law

• The magnetic field set up by a current-carrying conductor can be found from the Biot-Savart law. This law asserts that the contribution $\delta \underline{\boldsymbol{B}}$ to the field set up by a current element $I \delta \underline{\boldsymbol{l}}$ at a point P, a distance $\underline{\boldsymbol{r}}$ from the current element, is:

dB = 0

$$\delta \underline{\boldsymbol{B}} = \frac{\mu_0}{4\pi} \frac{I \, \delta \underline{\boldsymbol{l}} \wedge \underline{\boldsymbol{\hat{r}}}}{r^2}$$

Axis of dl

Next Eecture

- Ampere's Law
 - B-fields inside and outside current carrying wires
 - B-fields inside solenoids
 - B-field from Toroidal Solenoid

Force between two long parallel currents