



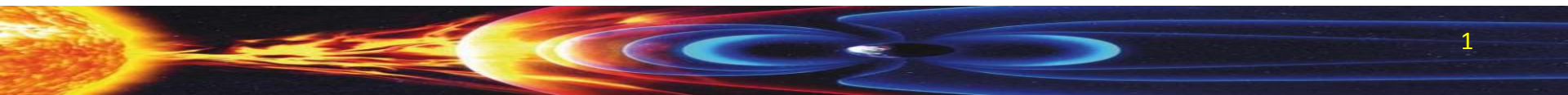
Electromagnetism

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Lecture 14

Ampere's Law

Week 7





Last Lecture

- Magnetic field from moving charge
- Magnetic field from current element
- **Biot-Savart Law**
 - B-Field at centre of current loop (magnetic dipole)
 - B-field from line of current
 - B-field from infinite line of current
 - B-field along axis of current loop (magnetic dipole)



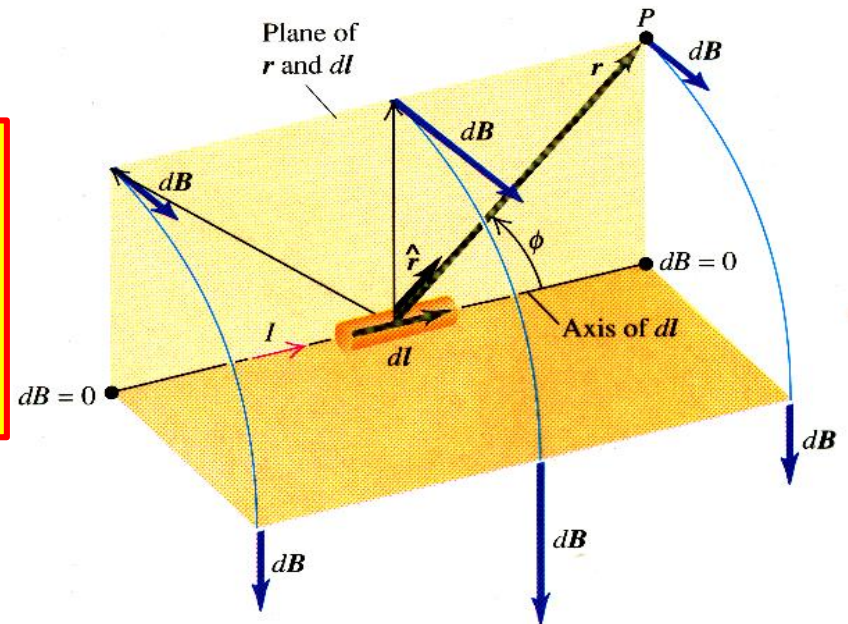
This Lecture

- Ampere's Law
 - B-fields inside and outside current carrying wires
 - B-fields inside solenoids
 - B-field from Toroidal Solenoid
- Force between two long parallel currents

Review – Biot-Savart Law

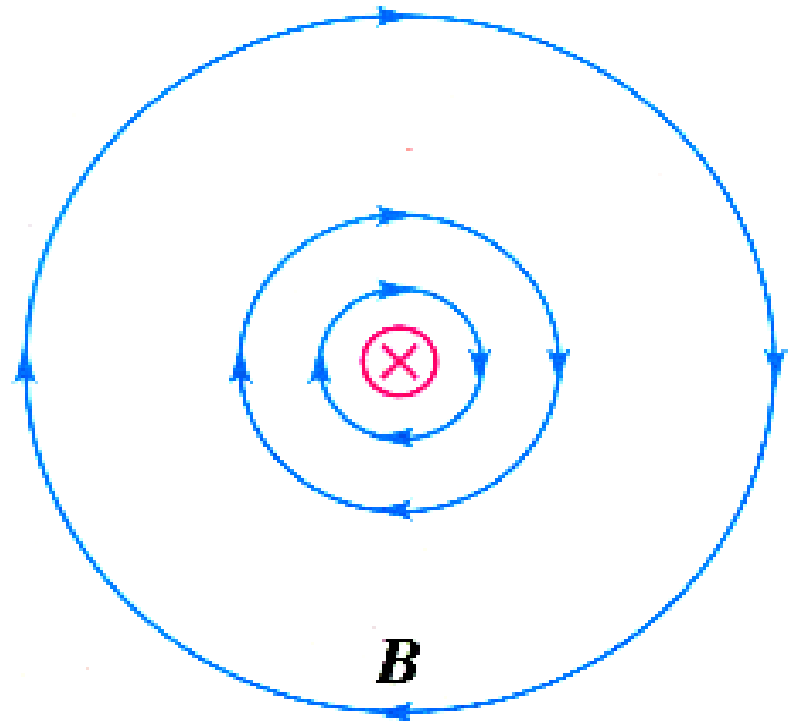
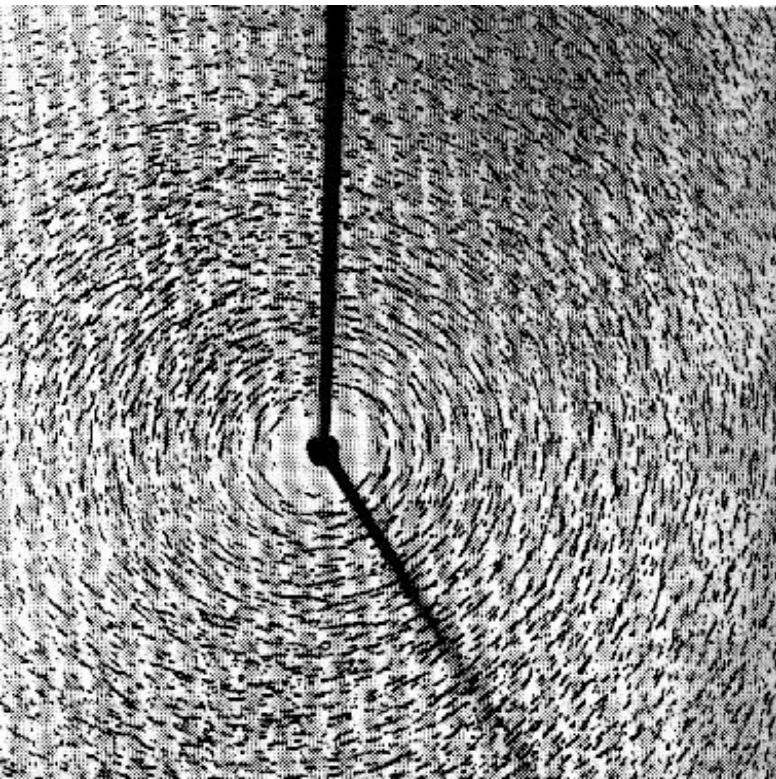
- The magnetic field set up by a current-carrying conductor can be found from the Biot-Savart law. This law asserts that the contribution $\delta \underline{B}$ to the field set up by a current element $I \delta \underline{l}$ at a point P , a distance \underline{r} from the current element, is:

$$\delta \underline{B} = \frac{\mu_0}{4\pi} \frac{I \delta \underline{l} \wedge \underline{\hat{r}}}{r^2}$$



B-field from line of Current

B-field lines *encircle* the current that acts as their source. B-field lines are continuous loops (lecture 11 - Law for Magnetism)



E- and B-Fields

- We already know the following:

Gauss's Law for E-fields

$$\int_S \underline{E} \cdot d\underline{S} = \frac{Q}{\epsilon_0}$$

Gauss's Law for B-fields

$$\int_S \underline{B} \cdot d\underline{S} = 0$$

From Lecture 5: E-field is conservative *i.e.* If a charge in an E-field returns to its original position, by any route, NO WORK IS DONE.

$$\oint \underline{E} \cdot d\underline{l} = 0$$

So, what about: $\oint \underline{B} \cdot d\underline{l}$?

Ampere's Law

Consider circular path of the B-field around an infinite line of current at a radial distance r from the line.

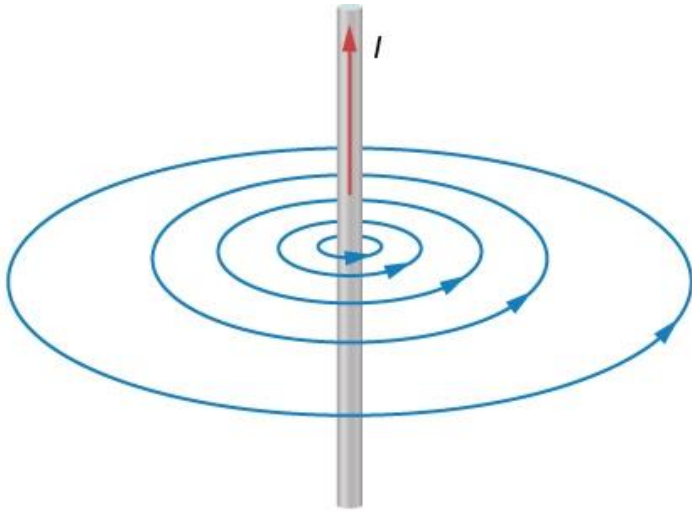
By symmetry, \underline{B} is parallel to $d\underline{l}$ and constant for fixed r . Hence

$$\oint \underline{B} \cdot d\underline{l} = \oint B dl = B \oint dl = B 2\pi r$$

From Lecture 13 (ex 13.2, Eq 13.2) B-field from infinite line of current is:

$$B = \frac{\mu_0 I}{2\pi r}$$

Ampere's Law



$$\oint \underline{B} \cdot d\underline{l} = B 2\pi r$$

But from Lecture 13

$$B = \frac{\mu_0 I}{2\pi r}$$

Hence :

$$\oint \underline{B} \cdot d\underline{l} = B 2\pi r = \frac{\mu_0 I}{2\pi r} 2\pi r = \mu_0 I$$

Ampere's Law

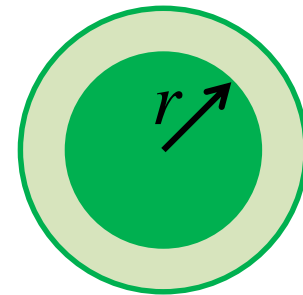
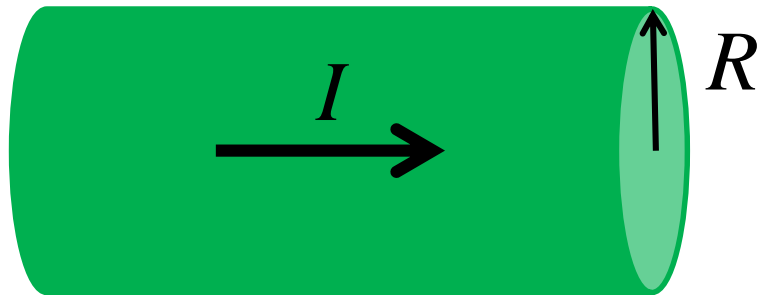
This is true for B-fields in general and is known as Ampere's Law:

$$\oint \underline{B} \cdot d\underline{l} = \mu_0 I_{enc}$$

Where I is the current enclosed in the integration loop

First Example (*Ex 14.1*)

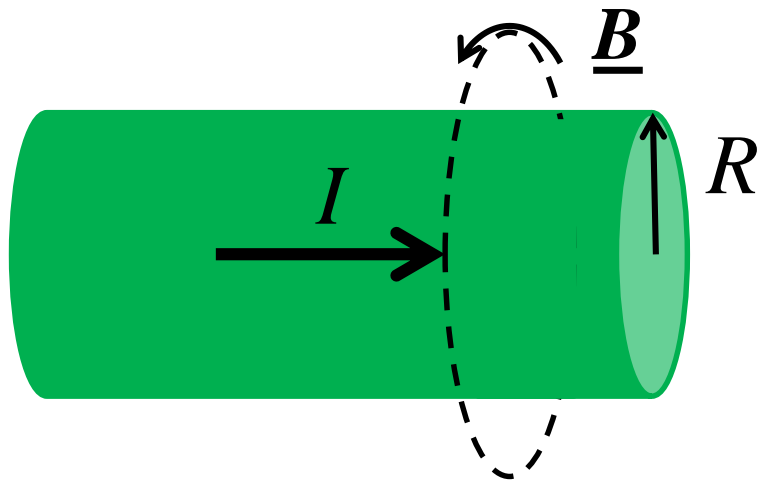
- B-Field Outside and Inside a Long Solid Cylindrical Conductor Carrying Uniformly Distributed Current



- Use Ampere's Law: $\oint \underline{B} \cdot d\underline{l} = \mu_0 I_{enc}$

First Example (Ex 14.1)

- Outside: $\oint \underline{B} \cdot d\underline{l} = \mu_0 I$
- By symmetry: \underline{B} is parallel to $d\underline{l}$ and constant for fixed r .



$$\text{LHS: } B 2\pi r$$

$$\text{RHS: } \mu_0 I$$

$$\Rightarrow B = \frac{\mu_0 I}{2\pi r}$$

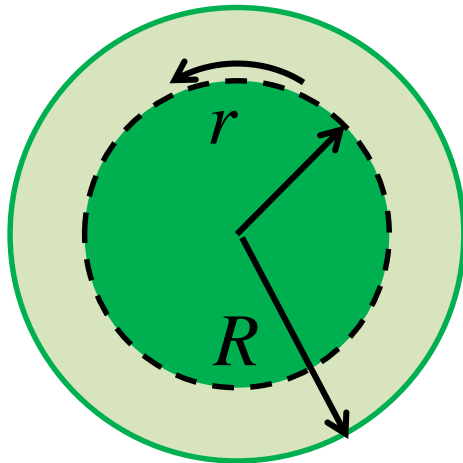
First Example (Ex 14.1)

- Inside: $\oint \underline{B} \cdot d\underline{l} = \mu_0 I_{enc}$
- By symmetry: \underline{B} is parallel to $d\underline{l}$ and constant for fixed r .

LHS: $B 2\pi r$

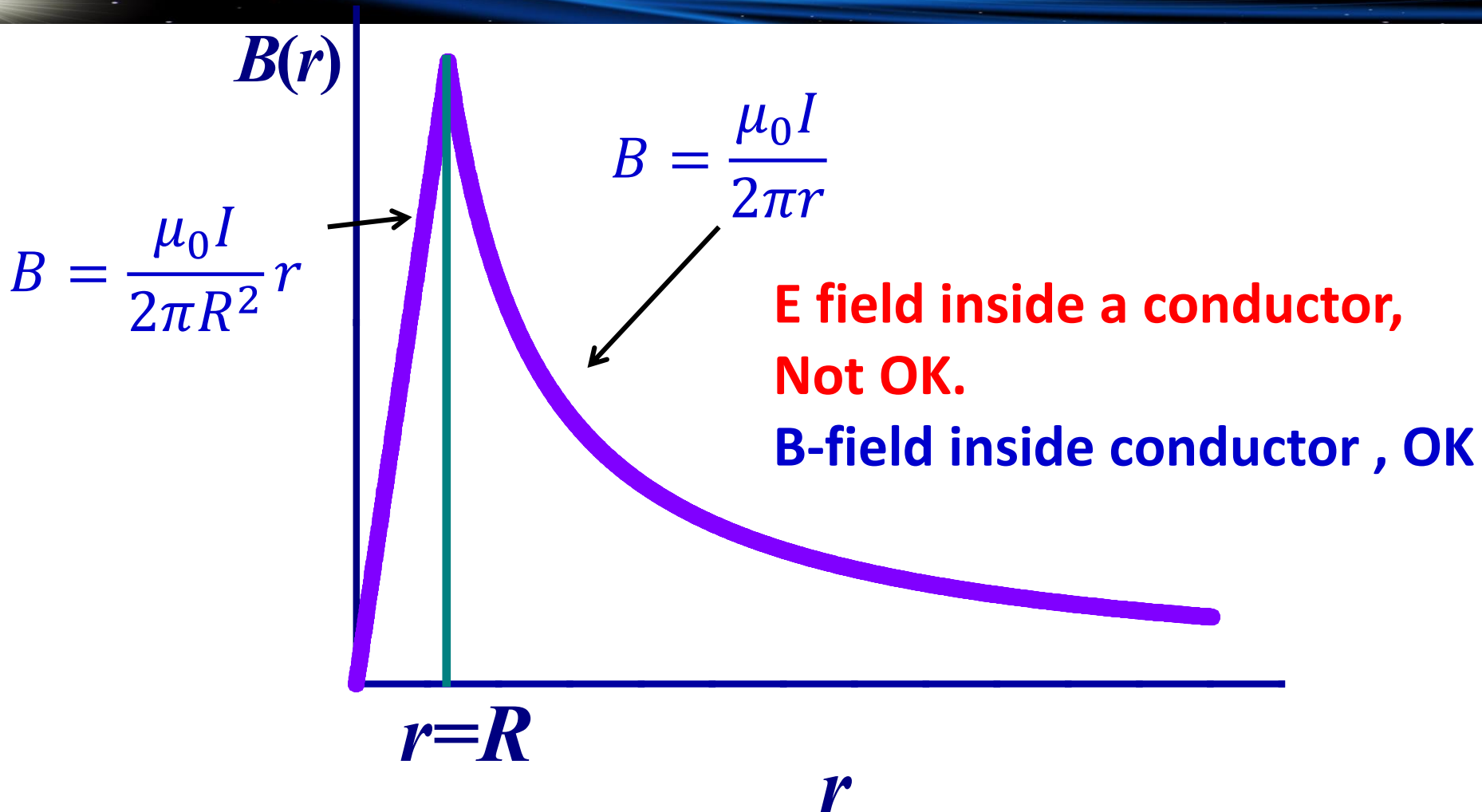
RHS: $\mu_0 I_{enc}$

For uniform current $I_{enc} = \frac{\pi r^2}{\pi R^2} I$



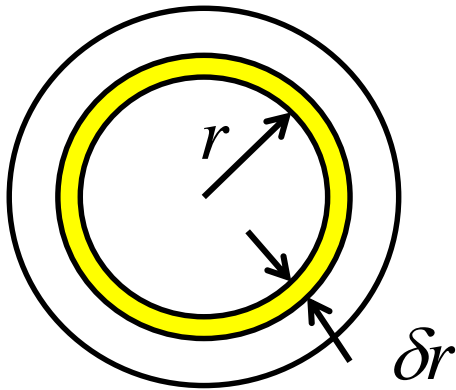
$$\Rightarrow B = \frac{\mu_0 I}{2\pi R^2} r$$

First Example: Long Solid Cylindrical Conductor



Example 14.2

- B-Field Inside a Long Solid Cylindrical Conductor Carrying Non-Uniformly Current.
- Current density $J = J_0 \frac{r^2}{R^2}$
- Ampere's Law becomes: $\oint \underline{B} \cdot d\underline{l} = \mu_0 \int_0^r \underline{J} \cdot d\underline{S}$

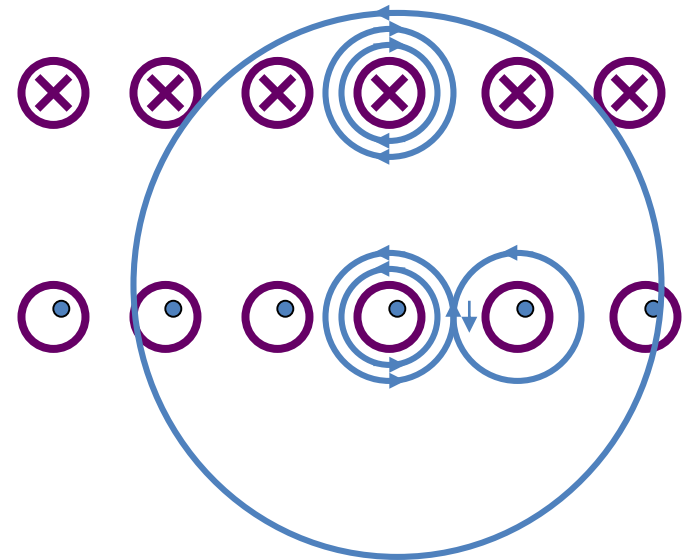
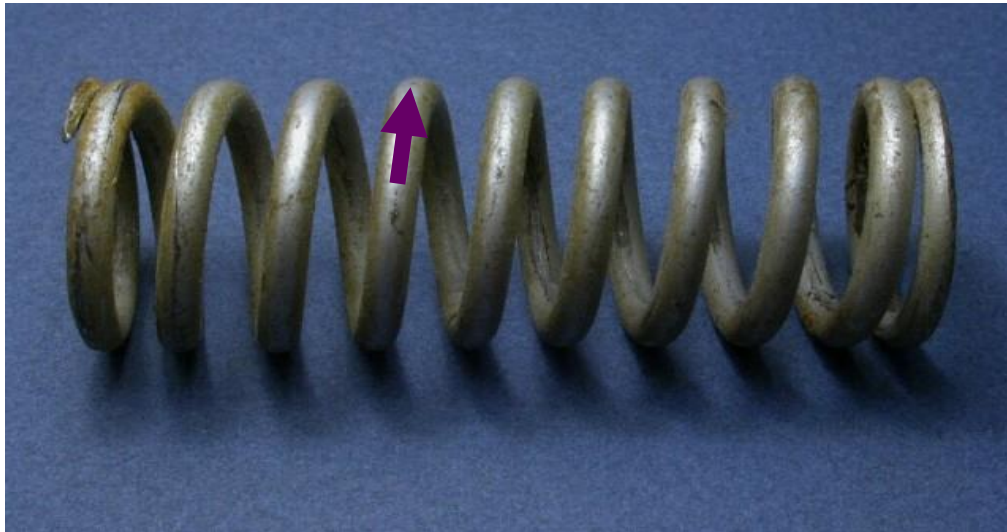


Element of area, $\delta S = 2\pi r \delta r$

Let's do it on the visualizer

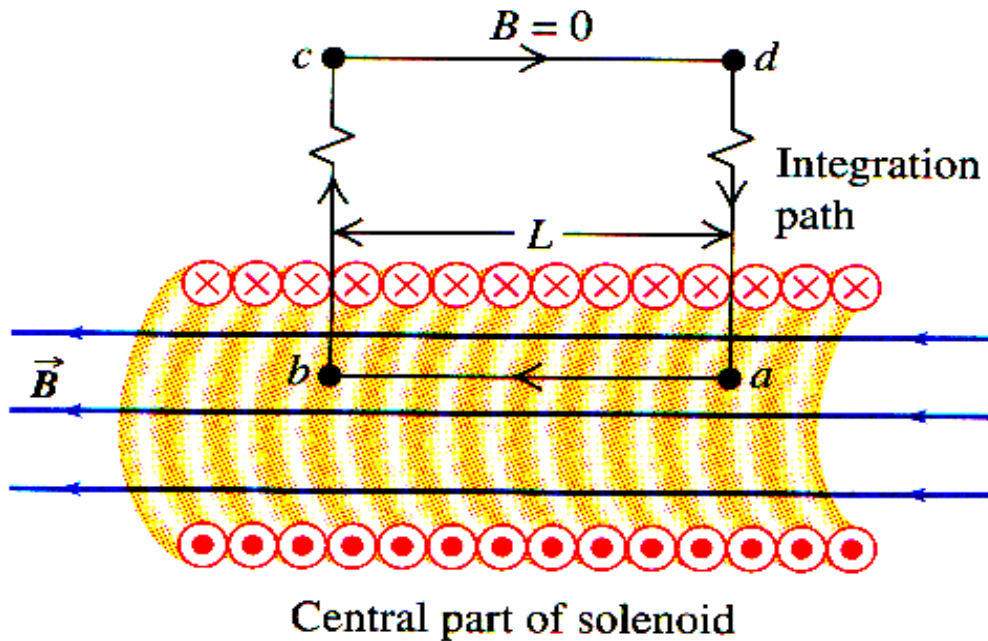
Example 14.3

- B-Field Inside a Long Solenoid



- At first sight, this looks complicated – Don't Panic!

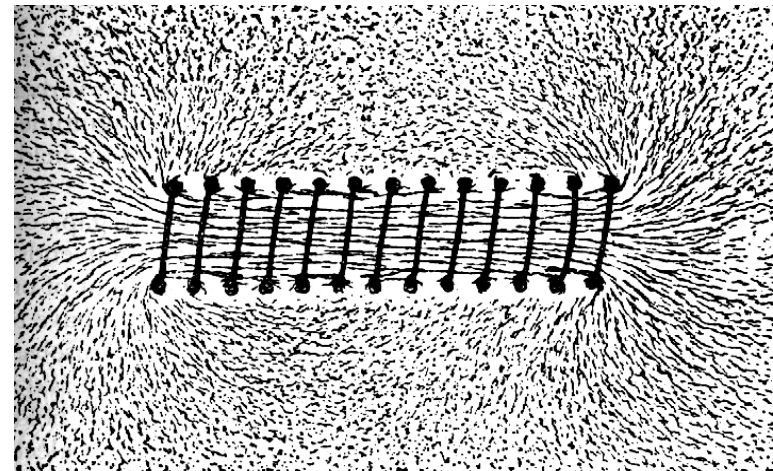
Example 14.3



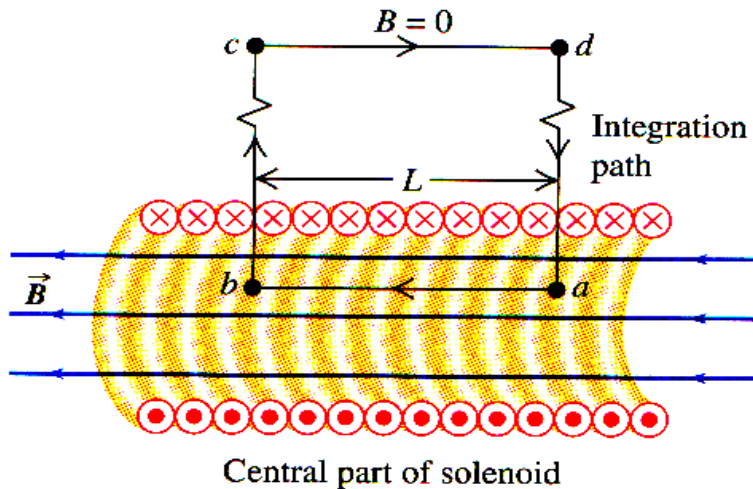
Choose integration path as shown

$$\oint \underline{B} \cdot d\underline{l} = BL$$

n = number of turns per unit length
So current enclosed in integration loop: $I_{enc} = nLI$



Example 14.3



$$\oint \underline{B} \cdot d\underline{l} = \mu_0 I_{enc}$$

$$\text{LHS} = B L$$

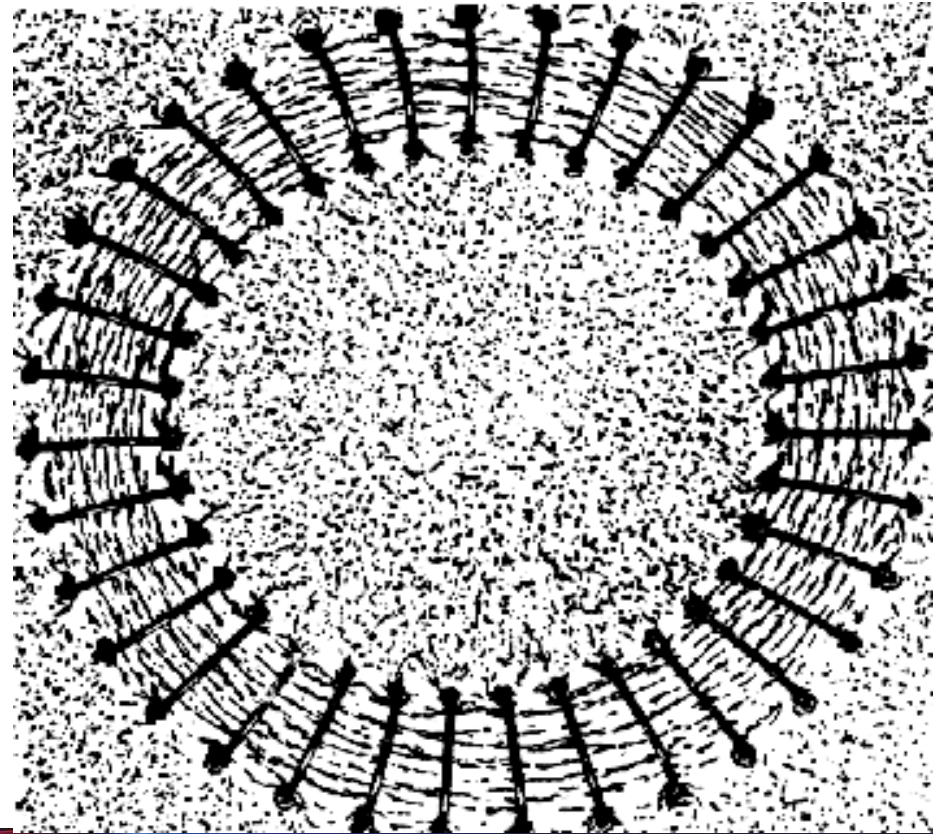
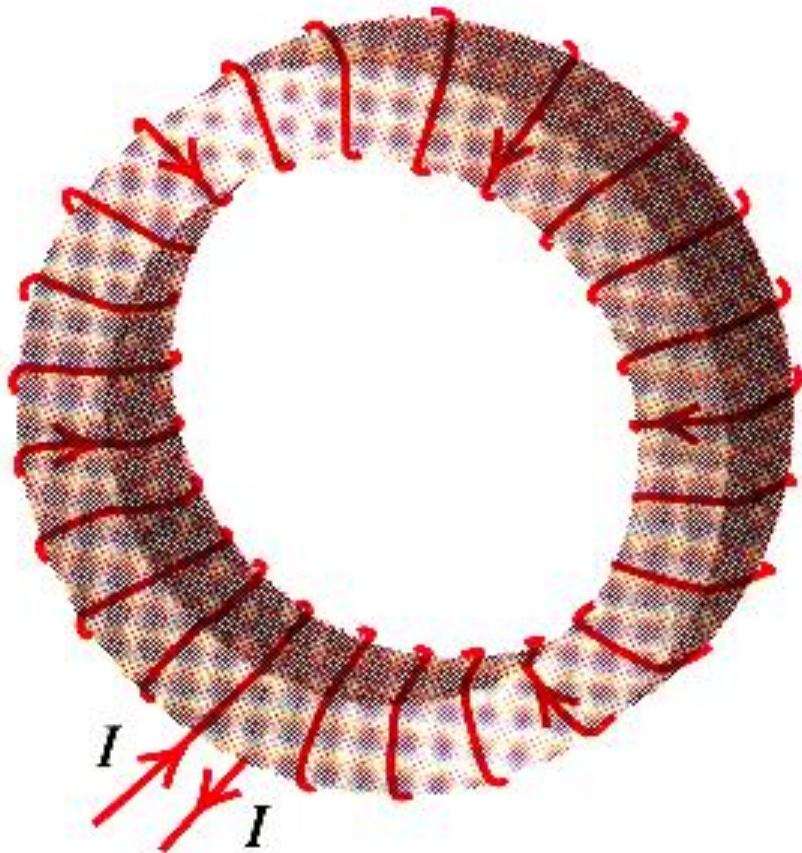
$$\text{RHS} = \mu_0 n L I$$

B-field inside long (i.e. neglecting end effects) Solenoid:

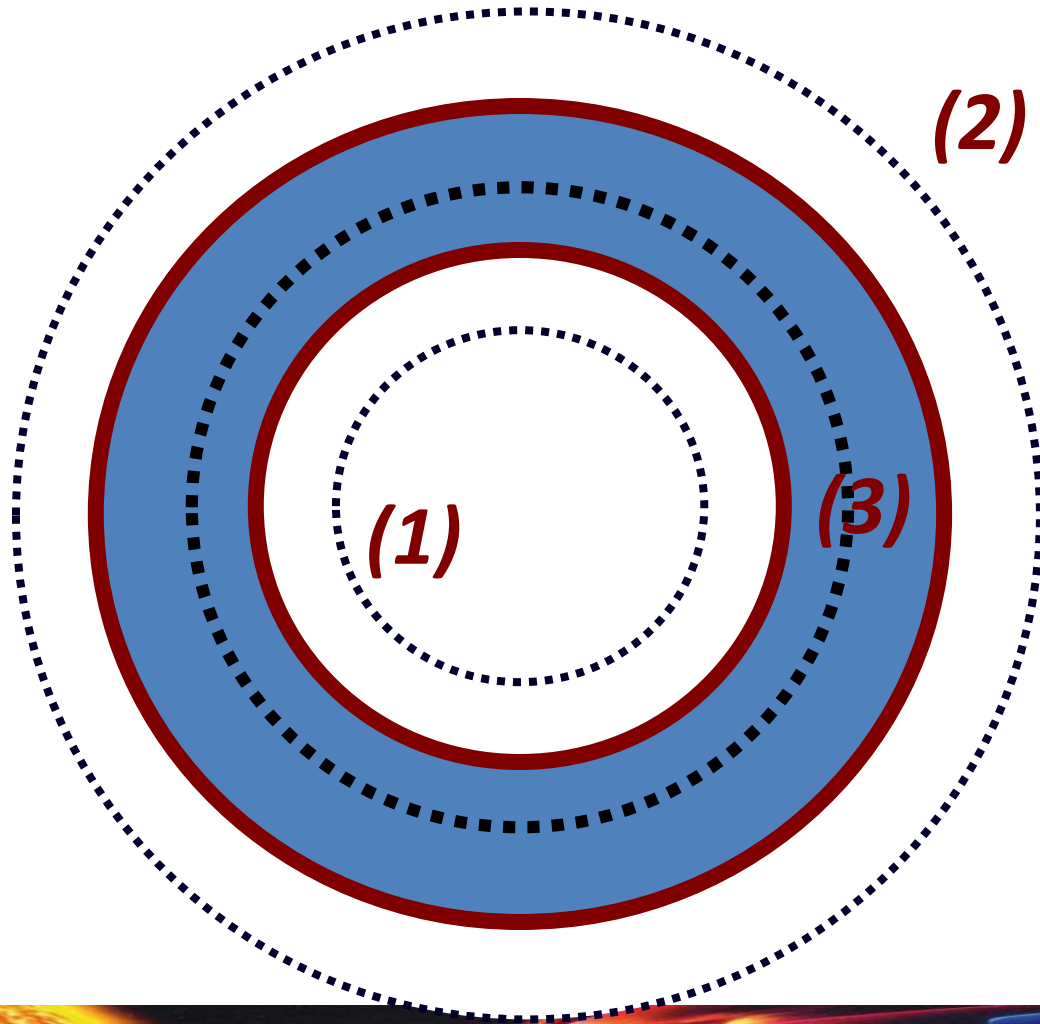
$$B = \mu_0 n I$$

Example 14.4

- Field of a Toroidal Solenoid



Field of a Toroidal Solenoid with N Turns



Path 1 – no current enclosed: $\underline{B} = 0$

Path 2 – no current enclosed: $\underline{B} = 0$

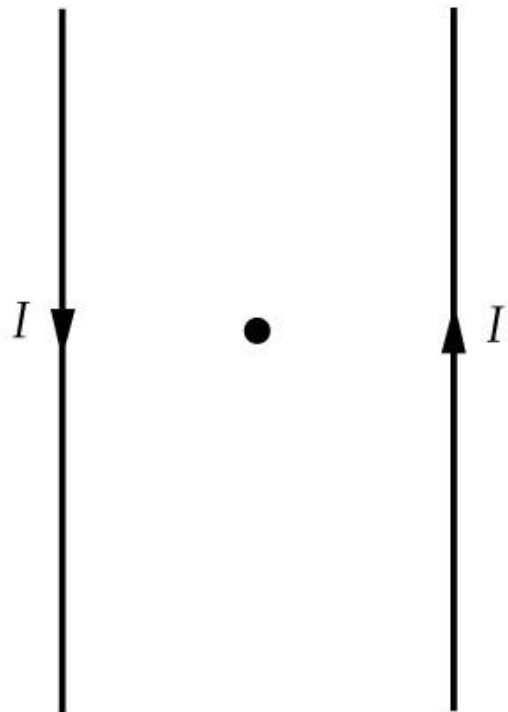
Path 3 – net current enclosed = NI

$$\oint \underline{B} \cdot d\underline{l} = B2\pi r = \mu_0 NI$$

$$B = \frac{\mu_0 NI}{2\pi r}$$

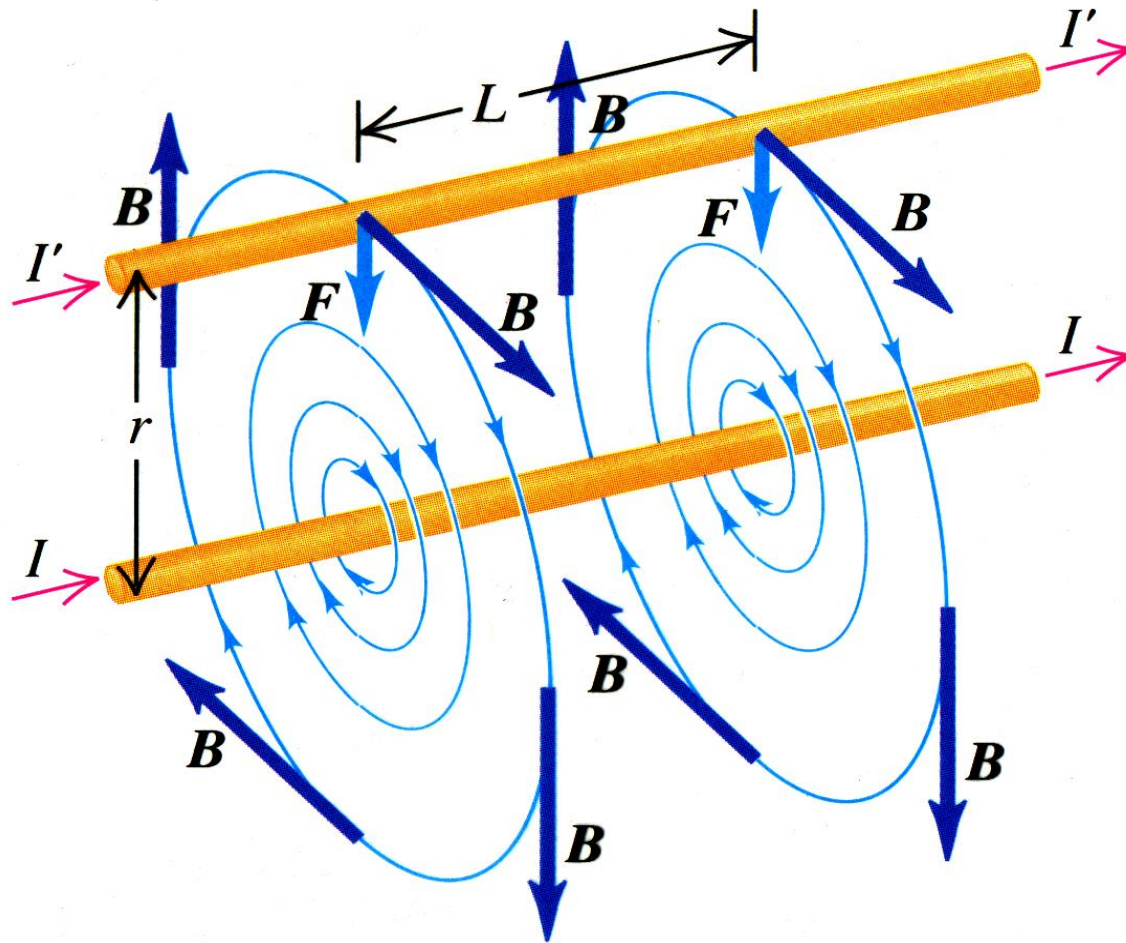
Quiz Time

- Two wires lie in the plane of the screen and carry equal currents in opposite directions. At a point midway between the wires, the magnetic field is

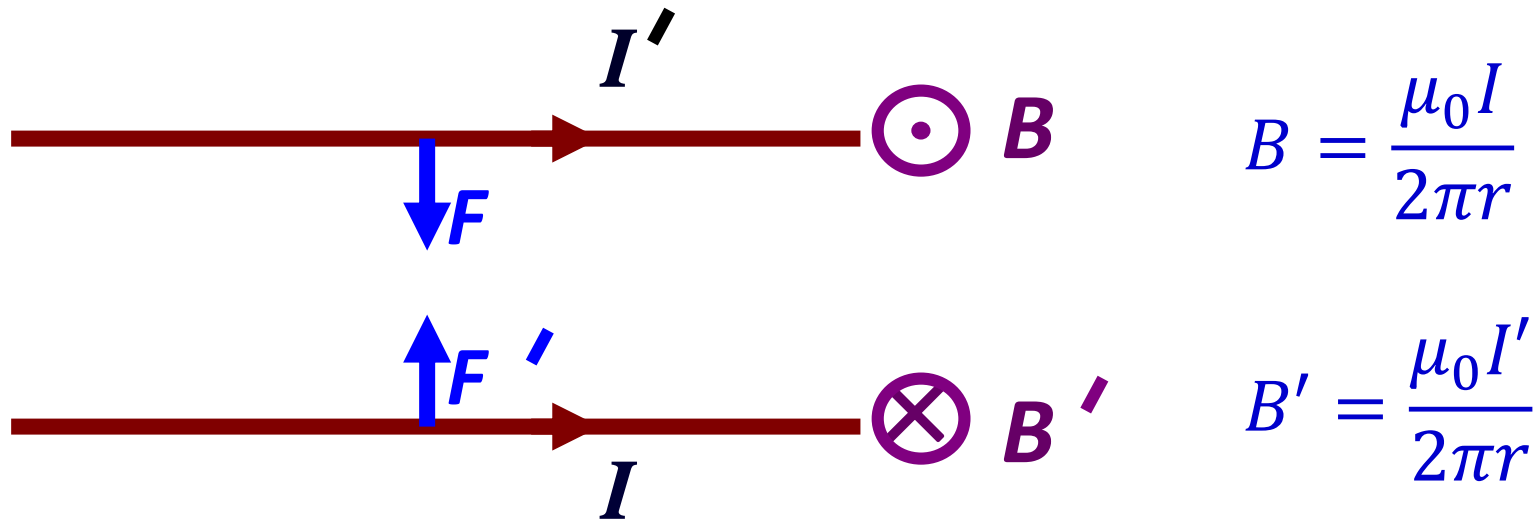


- (a) zero
- (b) into the screen
- (c) out of the screen
- (d) toward the top or bottom of the screen
- (e) toward one of the wires

Force between Two Long Parallel Currents



Force between parallel Currents



Force F on a length L of the upper conductor is:

$$F = I' L B = \frac{\mu_0 I I' L}{2\pi r}$$

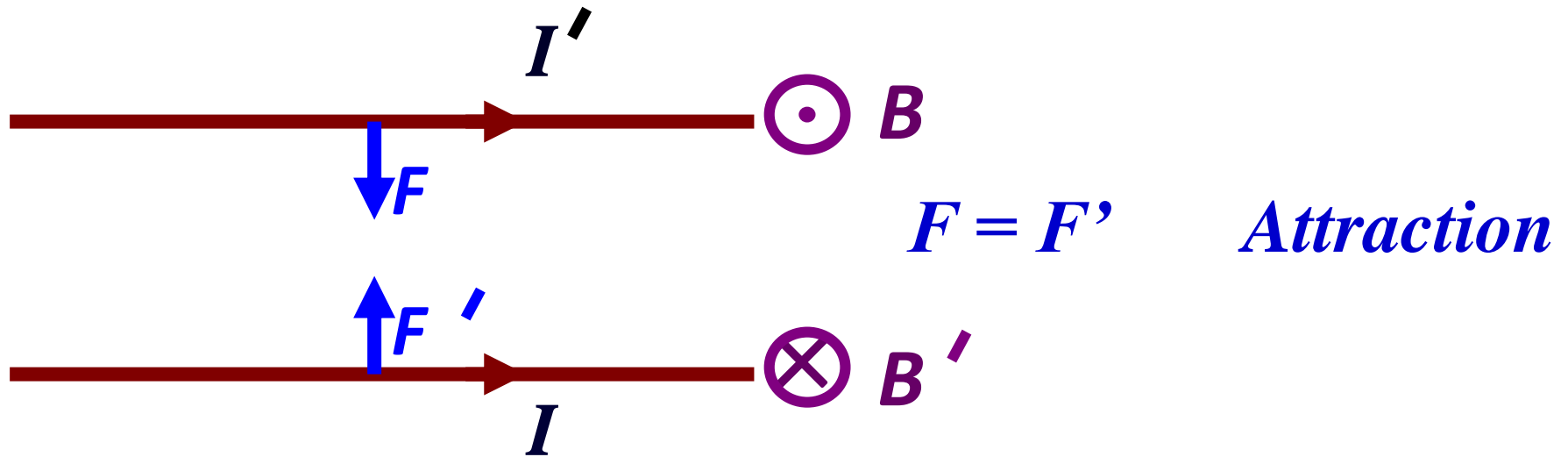
and

$$F' = I L B' = \frac{\mu_0 I I' L}{2\pi r}$$

$$F = F'$$

Attraction

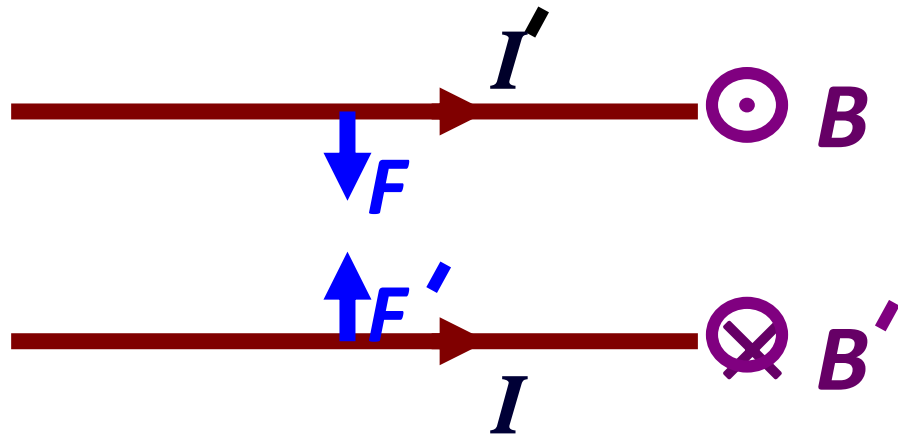
Force between parallel Currents



What happens when the currents are in opposite directions? (Ans. Repulsion)

$$\text{Force per unit length: } \frac{F}{L} = \frac{\mu_0 I I'}{2\pi r}$$

Force between parallel Currents



$F = F'$ *Attraction*

Force per unit length:
$$\frac{F}{L} = \frac{\mu_0 I I'}{2\pi r}$$

This *fundamental magnetic effect* was first studied by Ampere (1822)



Definition of the Ampere

- The ampere is that steady current which, flowing in two infinitely long straight parallel conductors of negligible cross-sectional area placed 1 m apart in a vacuum, causes each wire to exert a force of 2×10^{-7} N on each metre of the other wire.
- $$\frac{F}{L} = \frac{\mu_0 II'}{2\pi r} = \frac{4\pi \times 10^{-7} \times 1 \times 1}{2\pi \times 1} = 2 \times 10^{-7} \text{ Nm}^{-1}$$
- **Definition of the Coulomb:** A current of one ampere carries a charge of one coulomb per second

Equations of Static Electric and Magnetic Fields

- For E- and B-fields that don't vary with time

Laws of Electrostatics

Laws of Magnetostatics

$$\int_S \underline{E} \cdot d\underline{S} = \frac{Q_{enc}}{\epsilon_0}$$

$$\int_S \underline{B} \cdot d\underline{S} = 0$$

(Integrals over the closed surface)

$$\oint \underline{E} \cdot d\underline{l} = 0$$

$$\oint \underline{B} \cdot d\underline{l} = \mu_0 I_{enc}$$

Quiz Time

- Two parallel wires carry currents I_1 and $I_2 (= 2I_1)$ in the same direction. The forces F_1 and F_2 on the wires are related by:

(a) $F_1 = F_2$

(b) $F_1 = 2F_2$

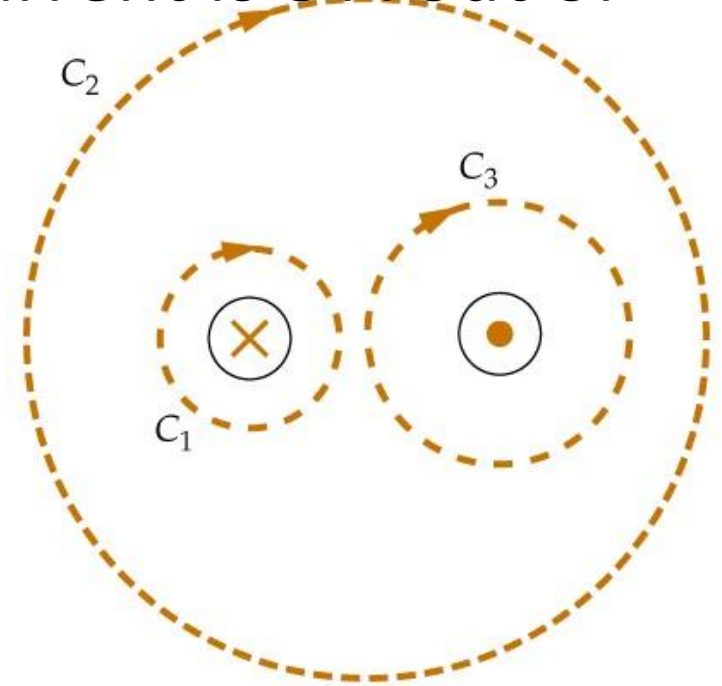
(c) $2F_1 = F_2$

(d) $F_1 = 4F_2$

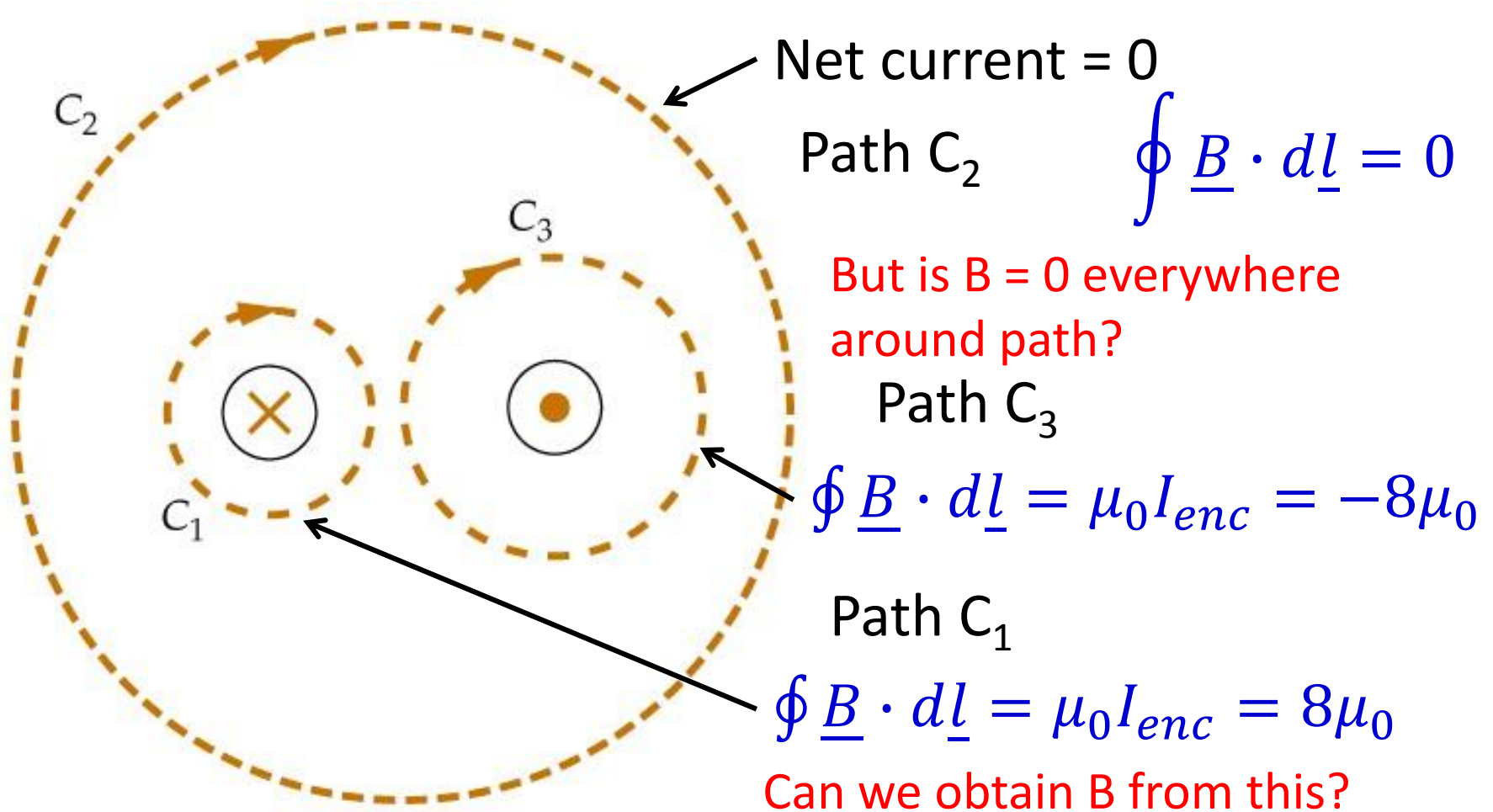
(e) $4F_1 = F_2$

Short Exercise

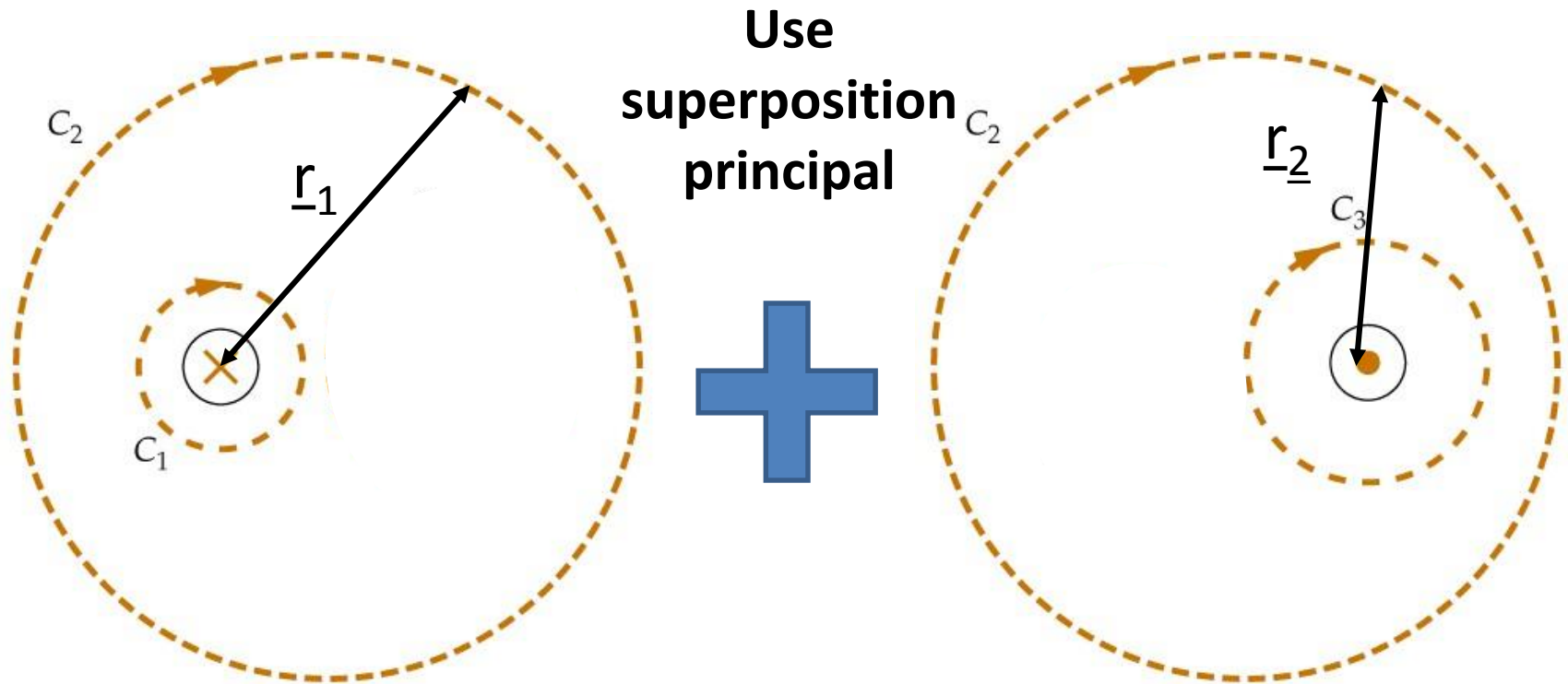
- The diagram shows two currents associated with infinitely long wires, one current of 8 A into the screen, the other current is 8 A out of the screen. Find
- $\oint \underline{B} \cdot d\underline{l} = \mu_0 I_{enc}$ for each path indicated.



Short Exercise



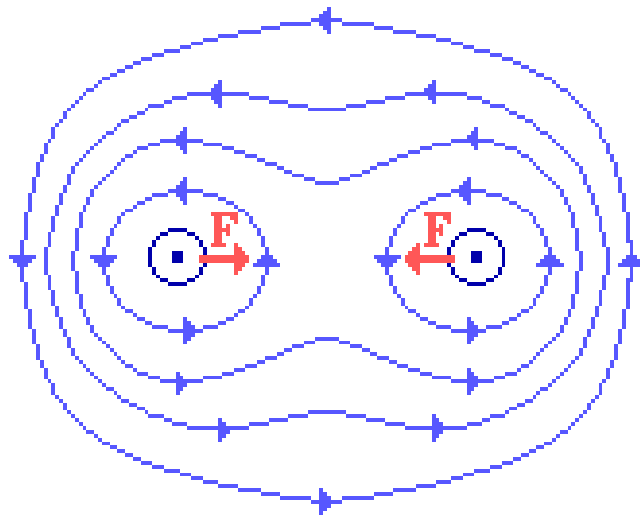
Short Exercise



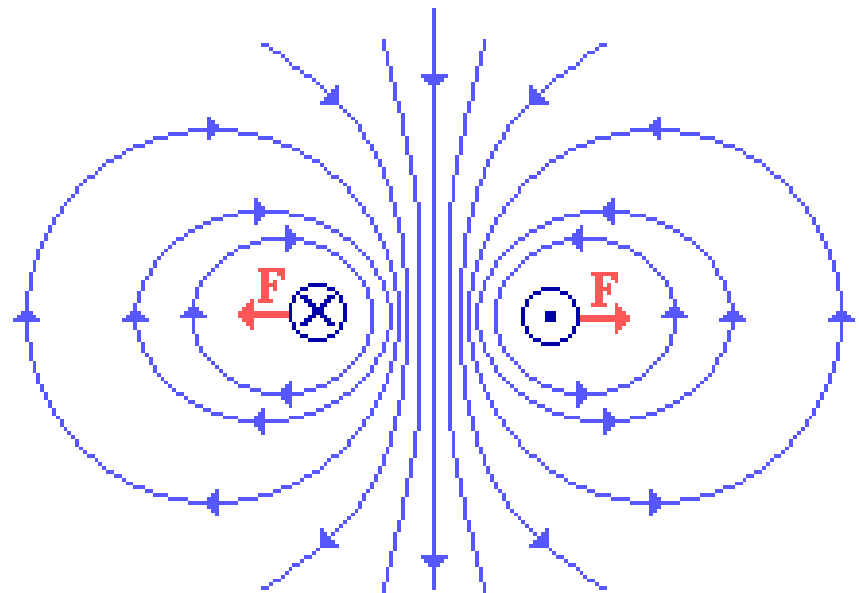
$$\mathbf{B} = \mathbf{B}_1 + \mathbf{B}_2 = \frac{\mu_0}{2\pi} \left\{ \frac{I \wedge \hat{\mathbf{r}}_1}{r_1} - \frac{I \wedge \hat{\mathbf{r}}_2}{r_2} \right\}$$

Cross products
to get direction
of B-field

B-field from two Parallel Currents



Same direction



Opposite direction

Example

- Two straight rods 50 cm long and 1.5 mm apart carry a current of 15 A in opposite directions. One rod lies vertically above the other. What mass must be placed on the upper rod to balance the magnetic force of repulsion?
- $mg = I'LB = \frac{\mu_0 II' L}{2\pi r} = \frac{4\pi \times 10^{-7} \times 15 \times 15 \times 0.5}{2\pi \times 1.5 \times 10^{-3}} = 0.015N$
- **Mass = 1.53 grams**
- **Note:** the magnetic force between two current-carrying wires is relatively small, even for currents as large as 15 A separated by only 1.5 mm.

Summary of Magnetostatics

$$\underline{F}_m = q \underline{v} \wedge \underline{B}$$

$$\underline{F} = I \underline{l} \wedge \underline{B}$$

$$\underline{\mu} = I \underline{A}$$

$$U = -\underline{\mu} \cdot \underline{B}$$

$$\underline{\tau} = \underline{\mu} \wedge \underline{B}$$

$$\underline{B} = \frac{\mu_0}{4\pi} \frac{q}{r^2} \underline{v} \wedge \underline{\hat{r}}$$

$$\phi_m = \int_S \underline{B} \cdot d\underline{S} = 0$$

$$\delta \underline{B} = \frac{\mu_0}{4\pi} \frac{I \delta \underline{l} \wedge \underline{\hat{r}}}{r^2}$$

$$\oint \underline{B} \cdot d\underline{l} = \mu_0 I_{enc}$$