

L2

# The Ultraviolet Catastrophe

Welcome to the UV catastrophe...  
Blackbody radiation

A 'black body' is an idealised object that

- Does not reflect any light at any wavelength
- Absorbs internally all incident light (none shines out the other side)

→ ie is perfectly black ('vantablack' is close)

All bodies emit electromagnetic energy.  
Paul Hollywood (and other humans) emits at 300K → infrared  
→ army night vision goggles

For the blackbody, the emission spectrum is only from this thermal emission (no reflection etc)

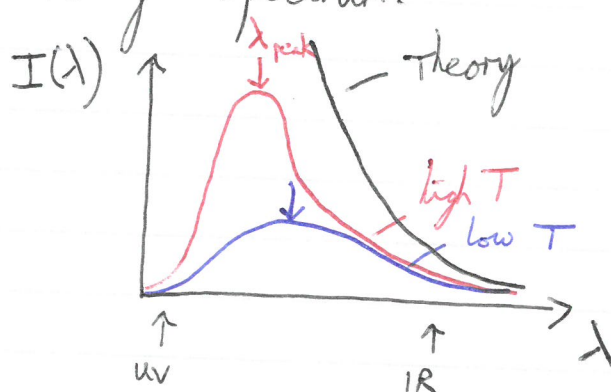
Hotter things are brighter and bluer (high energy, short wavelength)  
↳ Think of hotter flames going blue

Blackbody  
Radiation  
Spectra

Classical thermodynamics gets the blackbody spectrum totally wrong - especially at small wavelengths  
→ "UV catastrophe"

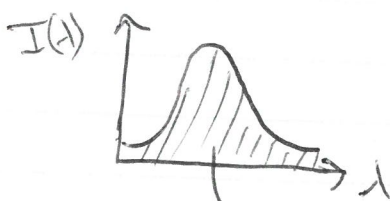
L2

- Blackbody spectrum:



$I(\lambda)$  is the intensity per wavelength, at a wavelength  $\lambda$

$I$ , intensity, is the total energy per time,  $I = \int_0^{\infty} I(\lambda) d\lambda$   
 in  $\text{Wm}^{-2}$  ie Power per area, adding up all colours



$I$  is area under curve  $\rightarrow$  total light

Empirical results:

- Stefan - Boltzmann Law:

$$I = \sigma T^4$$

$\uparrow$   
 Stefan - Boltzmann constant

$$\sigma = 5.67 \times 10^{-8} \text{ Wm}^2 \text{ K}^{-4}$$

- Wien's Displacement Law:

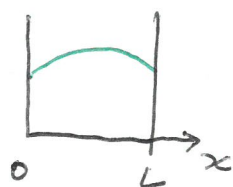
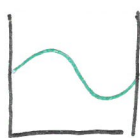
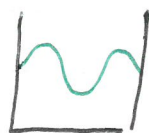
$$\lambda_{\text{peak}} = \frac{b}{T} \leftarrow 2.898 \times 10^{-3} \text{ Km}$$

L2

- Why does classical thermodynamics break?

Model  $I(\lambda)$  spectrum by slotting standing waves into a cavity

Consider a 1D cavity of length  $L$

 $n=1$  $n=2$  $n=3$  $n = (\text{Big number})$ 

"Cavity modes"  $\rightarrow$  the only allowed waves in there

Amplitude:  $a(x) = \sin\left(\frac{n\pi x}{L}\right)$ ,  $n = 1, 2, 3, \dots, \infty$   
 $\uparrow$  'mode number'

- Relationship between  $n$  and  $\lambda$ :

$$\lambda = \frac{2L}{n} \quad (\text{look at the pictures above!})$$

$\Rightarrow$  number of modes per wavelength,

$$n(\lambda) = \frac{2L}{\lambda}$$

L2

4

classically:

$$I(\lambda) \propto$$

$$\frac{n(\lambda)}{\lambda} k_B T$$

$$\propto \frac{1}{\lambda^2}$$

↑  
density of  
modes at  $\lambda$

↑  
average energy of modes

$$n(\lambda) \propto \frac{1}{\lambda}$$

$$I(\lambda) \propto \frac{1}{\lambda^2} \rightarrow \text{for } \lambda \rightarrow 0, I(\lambda) \rightarrow \infty$$

$\Rightarrow$  CATASTROPHE

- In 3D, same argument gives  $n(\lambda) \propto \frac{1}{\lambda^3}$

$$\Rightarrow I(\lambda) = \frac{2\pi c}{\lambda^4} k_B T$$

(Rayleigh-Jeans Law)

CATASTROPHE - works at large  $\lambda$  though

$\rightarrow$  The problem is the ' $k_B T$ ' bit. The assumption is that all cavity modes have an average energy  $k_B T$  - "Equipartition Theorem" - you will meet this in later courses.

In brief: Probability distribution of energies is Boltzmann distribution:

$$P(E) = \frac{e^{-E/k_B T}}{k_B T}$$

$k_B$  is Boltzmann constant

Average energy:

$$\bar{E} = \int_0^{\infty} E P(E) dE = k_B T$$

(see problem set)





- Planck's hypothesis (desperate, crazy!):

'Energy comes in discrete packets, 'quanta', that are proportional to frequency

$$\Delta E = hf = \frac{hc}{\lambda}$$

$$\left[ h = 6.626 \times 10^{-34} \text{ Js} \right]$$

Planck's constant

Striking this into the Partition Function of statistical mechanics (will learn this in later courses, don't worry now!), we get an average energy:

$$\bar{E}(\lambda) = \frac{hc/\lambda}{e^{\frac{hc}{\lambda k_B T}} - 1}$$

instead of  $k_B T$

Looking at limits:

$$\rightarrow \bar{E}(\lambda \rightarrow \infty): \quad \frac{hc}{\lambda k_B T} \ll 1$$

$$\therefore e^{\frac{hc}{\lambda k_B T}} \approx 1 + \frac{hc}{\lambda k_B T} + \dots \quad \left[ \text{Taylor series of } e^x \right]$$

$$\text{so } \bar{E}(\lambda) \approx \frac{\frac{hc}{\lambda}}{1 + \frac{hc}{\lambda k_B T} - 1} = k_B T$$

↓  
classical behaviour  
recovered at low energy

$$\rightarrow \bar{E}(\lambda \rightarrow 0): e^{\frac{1}{0}} = e^{\infty} \rightarrow \infty \text{ (fast)}$$

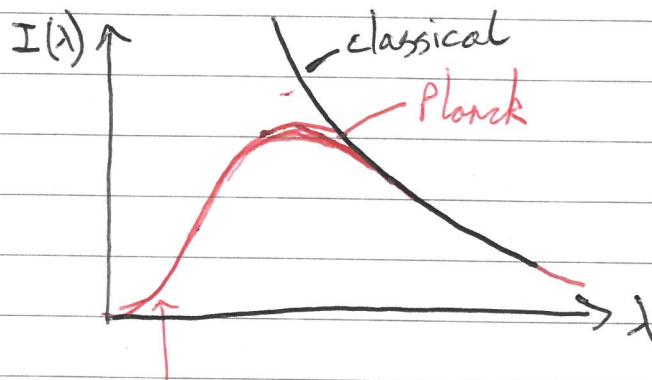
$$\therefore \bar{E}(\lambda \rightarrow 0) \approx \frac{\frac{hc}{0}}{e^{\frac{1}{0}} - 1} \rightarrow \text{small}$$

Exp. term wins

$$\bar{E}(\lambda \rightarrow 0) \rightarrow 0$$

[Can do this rigorously via L'Hopital's Rule]

$\Rightarrow$  No catastrophe!



Energy per mode,  $\bar{E}$ , goes to 0  $\rightarrow$  brings curve down

|| Science Works!

Conclusion  $\rightarrow$  This weird hypothesis fits the data

- Quantising energy means average energy of cavity modes is wavelength dependent
- Solves UV catastrophe
- Introduces idea of quantised 'pockets' of energy