Electromagnetism

Professor D. Evans d.evans@bham.ac.uk

Physics beyond the Stand Module!

Maxwell's Equations

Week 10

(Week 11 of Semester 2)

Last Lecture

- Paramagnetic materials
- Magnetic Susceptibility
- Relative permeability
- Ferromagnetic materials
- Diamagnetic materials
- Displacement Current Density

This ecture

- Ampere-Maxwell Law
- Maxwell's Equations in integral form
- Vector calculus
- Divergence and Stokes' theorems
- Maxwell's Equations in differential form
- The End

Extra Material - Towards Maxwell's Equations

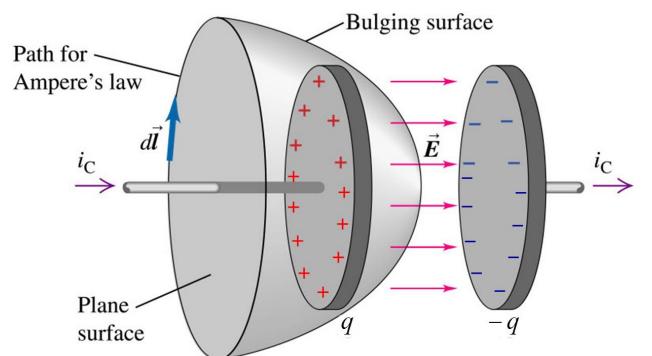
- Physics beyond the Stand Module!
- This lecture is not in the syllabus so you can relax.
- However, we've got this far so may as well get to Maxwell's Equations.

Changing B-fields induce a changing E-Field.

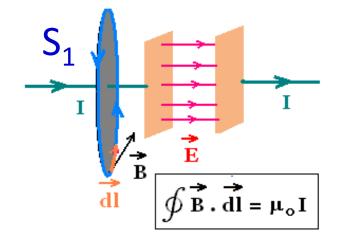
•
$$\oint \underline{\mathbf{E}} \cdot d\underline{\mathbf{l}} = -\frac{d\Phi_m}{dt}$$

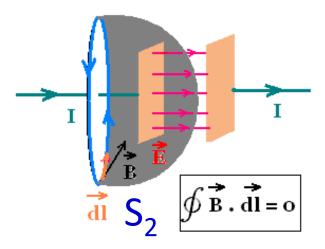
- And $\phi_m = \int_S \ \underline{\boldsymbol{B}} \cdot d\underline{\boldsymbol{S}}$
- So: $\oint \underline{E} \cdot d\underline{l} = -\int_{S} \frac{\partial \underline{B}}{\partial t} \cdot d\underline{S}$ (Maxwell's third eqⁿ)
- So, can a changing E-field induce a changing B-field?

- Consider an electrical circuit with a capacitor in.
- For an AC voltage, there will be an AC current in the circuit.



Apply
Ampere's Law
around the
wire, in front
of a capacitor
plate.





Consider plane surface, S₁.

$$\oint \underline{\mathbf{B}} \cdot d\underline{\mathbf{l}} = \int_{S_1} \underline{\mathbf{J}} \cdot d\underline{\mathbf{S}} = \mu_0 I$$

Now consider the surface, S_2 .

$$\oint \underline{\mathbf{B}} \cdot d\underline{\mathbf{l}} = \int_{S_2} \underline{\mathbf{J}} \cdot d\underline{\mathbf{S}} = 0$$

We seem to have a problem!

Displacement Current Density

- There is an AC current in the wires but not between the capacitor plates.
- However, there must be something and that is an alternative electric field.
- The plates and discharging and charging as the voltage alternates.
- Current is rate of change of charge $I = \frac{dQ}{dt}$
- And Current density, $J = \frac{d\sigma}{dt}$

Displacement Current Density

- Current density, $J_D = \frac{d\sigma}{dt}$
- But for capacitor plates, $E = \frac{\sigma}{\varepsilon_0}$
- So $J_D = \varepsilon_0 \frac{\partial \underline{E}}{\partial t}$ This is known as the Displacement Current Density
- It's not a current, it's a changing electric field but has the same effect as a current i.e. produces B-field.

 Ampere's law needs to be modified to take the displacement current into account:

Ampere-Maxwell Law

$$\oint \underline{\mathbf{B}} \cdot d\underline{\mathbf{l}} = \mu_0 \int_{\mathcal{S}} \left(\underline{\mathbf{J}} + \varepsilon_0 \frac{\partial \underline{\mathbf{E}}}{\partial t} \right) \cdot d\underline{\mathbf{s}}$$

In free space, in integral form:

•
$$\int_{S} \underline{\boldsymbol{E}} \cdot d\underline{\boldsymbol{S}} = \int_{V} \frac{\rho}{\varepsilon_{0}} dV$$

M1

•
$$\int_{S} \underline{\mathbf{B}} \cdot d\underline{\mathbf{S}} = 0$$

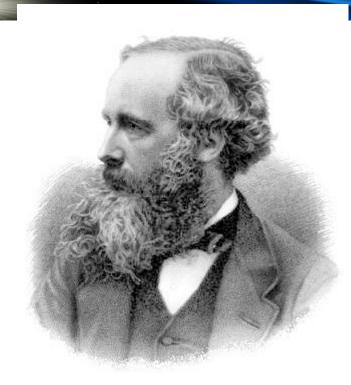
M2

•
$$\oint \underline{E} \cdot d\underline{l} = -\int_{S} \frac{\partial \underline{B}}{\partial t} \cdot d\underline{S}$$

M3

•
$$\oint \underline{\mathbf{B}} \cdot d\underline{\mathbf{l}} = \mu_0 \int_{\mathcal{S}} \left(\underline{\mathbf{J}} + \varepsilon_0 \frac{\partial \underline{\mathbf{E}}}{\partial t} \right) \cdot d\underline{\mathbf{S}}$$
 M4

James Clerk Maxwell



13 June 1831 – 5 November 1879

- British physicist
- formulate the classical theory of electromagnetic radiation (unifications of electric and magnetic forces)
- Proposed light is EM radiation
- helped develop the Maxwell– Boltzmann distribution, to describe kinetic theory of gases.

Considered the 3rd greatest physicist of all time (after Newton and Einstein)

Vector Calculus Year 2 mathematics

•
$$\nabla = \frac{\partial}{\partial x} \hat{\underline{\imath}} + \frac{\partial}{\partial y} \hat{\underline{\jmath}} + \frac{\partial}{\partial z} \hat{\underline{k}}$$

• grad(a)
$$\rightarrow \nabla a = \frac{\partial a}{\partial x} \hat{\underline{\imath}} + \frac{\partial a}{\partial y} \hat{\underline{\jmath}} + \frac{\partial a}{\partial z} \hat{\underline{k}}$$

– (a is any scalar)

•
$$\operatorname{div} \underline{A} \to \nabla \cdot \underline{A} = \frac{\partial A_{\chi}}{\partial \chi} + \frac{\partial A_{y}}{\partial y} + \frac{\partial A_{z}}{\partial z}$$

• curl
$$\underline{A} \to \nabla \wedge \underline{A} = \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z}\right) \hat{\underline{\iota}} - \left(\frac{\partial A_z}{\partial x} - \frac{\partial A_x}{\partial z}\right) \hat{\underline{\jmath}} + \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y}\right) \hat{\underline{k}}$$
 (describes the rotation of a vector field)

Two Vector Theorems (year 2 maths)

• Divergence Theorem

•
$$\int_{S} \underline{A} \cdot d\underline{S} = \int_{V} \nabla \cdot \underline{A} \, dV$$

where S is the enclosed surface around a volume V.

Stokes' Theorem

•
$$\oint \underline{A} \cdot d\underline{l} = \int_{S} (\nabla \wedge \underline{A}) \cdot d\underline{S}$$

where l is length of the boundary around the area S.

• Both true for any vector \underline{A} .

•
$$\int_{S} \underline{E} \cdot d\underline{S} = \int_{V} \frac{\rho}{\varepsilon_{0}} dV$$

Applying the divergence theorem to the LHS

•
$$\int_{S} \underline{E} \cdot d\underline{S} = \int_{V} \nabla \cdot \underline{E} \, dV$$

- Hence $\int_V \nabla \cdot \underline{\boldsymbol{E}} \, dV = \int_V \frac{\rho}{\epsilon_0} \, dV$
- If this is true for <u>any</u> volume, then:

$$\nabla \cdot \underline{\boldsymbol{E}} = \frac{\rho}{\varepsilon_0}$$

•
$$\int_{S} \underline{\mathbf{B}} \cdot d\underline{\mathbf{S}} = 0$$

Applying the divergence theorem to the LHS

•
$$\int_{S} \underline{\boldsymbol{B}} \cdot d\underline{\boldsymbol{S}} = \int_{V} \nabla \cdot \underline{\boldsymbol{B}} \, dV$$

- Hence $\int_{V} \nabla \cdot \underline{\boldsymbol{B}} \, dV = 0$
- If this is true for <u>any</u> volume, then:

$$\nabla \cdot \underline{\boldsymbol{B}} = 0$$

•
$$\oint \underline{E} \cdot d\underline{l} = -\int_{S} \frac{\partial \underline{B}}{\partial t} \cdot d\underline{S}$$

Applying Stokes' theorem to the LHS

•
$$\oint \underline{E} \cdot d\underline{l} = \int_{S} (\nabla \wedge \underline{E}) \cdot d\underline{S}$$

- Hence $\int_{S} (\nabla \wedge \underline{\mathbf{E}}) \cdot d\underline{\mathbf{S}} = -\int_{S} \frac{\partial \underline{\mathbf{B}}}{\partial t} \cdot d\underline{\mathbf{S}}$
- If this is true for <u>any</u> surface, then:

$$\nabla \wedge \underline{\boldsymbol{E}} = -\frac{\partial \underline{\boldsymbol{B}}}{\partial t}$$

•
$$\oint \underline{\mathbf{B}} \cdot d\underline{\mathbf{l}} = \mu_0 \int_{\mathcal{S}} \left(\underline{\mathbf{J}} + \varepsilon_0 \frac{\partial \underline{\mathbf{E}}}{\partial t} \right) \cdot d\underline{\mathbf{S}}$$

- Applying Stokes' theorem to the LHS
- $\oint \underline{\mathbf{B}} \cdot d\underline{\mathbf{l}} = \int_{S} (\nabla \wedge \underline{\mathbf{B}}) \cdot d\underline{\mathbf{S}}$
- Hence $\int_{S} (\nabla \wedge \underline{\mathbf{B}}) \cdot d\underline{\mathbf{S}} = \mu_0 \int_{S} (\underline{\mathbf{J}} + \varepsilon_0 \frac{\partial \underline{\mathbf{E}}}{\partial t}) \cdot d\underline{\mathbf{S}}$
- If this is true for <u>any</u> surface, then:

$$\nabla \wedge \underline{\boldsymbol{B}} = \mu_0 \left(\underline{\boldsymbol{J}} + \varepsilon_0 \frac{\partial \underline{\boldsymbol{E}}}{\partial t} \right)$$

and God said $\nabla \cdot E = \frac{\rho}{\varepsilon_0}$ $\nabla \cdot B = 0$ $\nabla x E = -\frac{\partial B}{\partial t}$ $\nabla x B = \mu_0 \left(J + \varepsilon_0 \frac{\partial E}{\partial t} \right)$ and then there was LIGHT

Electromagnetism

