

# Symmetries in Spin Systems and Physical Aware Equivariant Neural Network

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## I. INTRODUCTION

Representing a spin system in a material efficiently is a challenge due to the sheer number of parameters involved, to name a few: atomic types and their physical descriptors, atomic positions, spins, coupling constants, etc. Each of them also obeys different sets of rules which always make representing the system complicated especially if one want to numerically represent it for computational purposes, e.g. machine learning.

A class of machine learning that would greatly benefit from efficient representation of the system is neural network, which contains a large amount of training parameters and require significant amount of distinct data points with good (sometimes crafty) augmentations. With the advance in recent development of equivariant neural network, it provides a great method to reduce both training parameters and augmentation for the same level of model complexity due to the fact that the models are aware of the symmetries of the data points which limit the possible choice of outcome, and drive the necessity of the augmentations related to said symmetries obsolete.

The goal of this work is to serve as the comprehensive mathematical guide for developing equivariant neural networks for systems with known physical symmetries, particularly, in this case, spin system.

## II. SYMMETRIES IN SPIN SYSTEM

For a general spin system, one can always divided the representation of the system into the physical space information, and the spin space information which is the internal degree of freedom of the quantum angular momentum. Hence, in addition to the normal symmetries in  $E(3)$  for physical information, we need to consider the  $SU(2)$  properties of the spin as well. Furthermore, we will also assume the non-relativistic treatment of the system which entails the consideration of Time reversal symmetry as well.

## III. SYMMETRIES AND GROUP REPRESENTATION

## IV. IRREDUCIBLE REPRESENTATION AND EQUIVARIANT

## V. IRREDUCIBLE REPRESENTATION OF SPIN SYSTEMS

In order to fully utilize the symmetries of spin systems, one need to find the irreducible representations

of the combined symmetries. However, not all symmetries commute with each other which make the action of one symmetry group mix different irreducible representations of the other non-commuting group. This defeats the purposes of using irreducible representation for efficient training of equivariant network. Hence, in these cases, we will restrict the model to only use invariant information of one of the symmetries.

## VI. PHYSICAL TRANSLATION ( $T$ )

The physical translation of the system doesn't commute with the physical rotation which we will discuss in more detail in the following sections. This means that the translation will mix the different irreducible representation of the rotation. Fortunately, the spin interaction directly depends on the relative distance between the involved objects rather than their absolute position. Hence, it should be the case that we will restrict the model to only use the relative positions for positional inputs since it is invariant to any physical translation.

## VII. PHYSICAL PERMUTATION ( $P$ )

The physical permutation of the system's objects also doesn't commute with the physical rotation. However, the quantum mechanical laws are applied to each and every objects in the same way regardless of the label we put on each object. Hence, it is also trivial to choose to restrict the model to be invariant of the object labeling, e.g. use sets as the key object containers.

## VIII. PHYSICAL ROTATION ( $R$ )

Any physical rotation of the system in 3D can be fully described by a real 3 by 3 orthogonal matrix with unit determinant which means that the action form the  $SO(3)$  group. One popular irreducible representation of such group is by representing the system with spherical harmonic basis, and the actions with Wigner-D matrices. Since one of the main challenge in deep neural network for learning 3D objects is in the amount of rotation augmentation which grow cubically with the accuracy instead of linearly in the case of 2D objects, it is mostly beneficial to prioritize the model to be physical rotation equivariant.

## IX. PHYSICAL INVERSION ( $I$ )

The physical inversion or parity is a very important physical symmetry that determine the possibility of certain interactions. Having this symmetry equivariance allows us to create physical aware machine learning model that can be fine tuned or separate the interested interactions. Fortunately, physical inversion commutes with physical rotation. This means that the combined irreducible representations are just the tensor products between their respective irreducible representations. The intuitive irreducible representation of the physical inversion is the symmetric-antisymmetric functions with the action to be multiplication by 1 for symmetric and by -1 for antisymmetric functions.

## X. SPIN ROTATION ( $R_s$ )

Similar to the physical rotation, any spin rotation in spin space can be fully described by a complex 2 by 2 unitary matrix with unit determinant which is the  $SU(2)$  group. The popular irreducible representation of this group is the spinor representation which is representing the system with spinors, and the actions with Wigner-D matrices for spinors. We can detect that obvious that the choice for irreducible representation for spin rotation is similar to physical rotation since  $SU(2)$  is a double cover of  $SO(3)$  with the same Lie algebra. Hence, the irreducible representation of physical rotation is redundant with spin rotation. Since each irreducible representation of spin rotation can be indexed by a non-negative integer or half-integer  $j$ , and those integer indexed representations are redundant with physical rotation. Hence, spin rotation equivariance model can also be used as physical equivariance model. One can also incorporate the physical inversion symmetry similar to the physical rotation case, however, the representation for the actual spin is

on spin space. Therefore, in the actual system, we need to restrict the parity of 1 to all spin representations.

## XI. TIME REVERSAL ( $\Theta$ )

At first glance, the time reversal seems to be a direct copy of physical inversion, but for temporal axis, since it commutes with all other symmetries excluding the one we already only use the invariances. However, in order for the action to preserve the non-relativistic Schrödinger equation, not only is the sign of time changed, but also the wavefunction conjugated. This fact will make the representation of the group action vary vastly depending on the choice of other symmetric representations whether the field of the representation is real or complex, and the behavior of those bases under time reversal actions.

For a concrete example, consider two irreducible representations of physical rotation where both of them are representing system with spherical harmonic basis, but one use real spherical harmonic while other use complex spherical harmonic. The first representation also use real number field which make time reversal actions the same as physical inversion. On the other hand, the complex spherical harmonic representation require the representations to be of complex field which make time reversal actions change the basis (flip sign of projected angular momentum number, i.e.  $m$ ), and perform conjugation on the representations. The situation is also different for the spin rotation representations where the real spinor bases representation can be of complex field.

## XII. EQUIVARIANT NEURAL NETWORK OF SPIN SYSTEM

## XIII. CONCLUSION