

Notes: cubic couplings in $3d$

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Scalar fields

The scalar Kaluza-Klein fluctuations of the $\mathrm{SO}(8, 4 + m)$ exceptional field theory [1] are labelled by pair of indices $\Phi^{\alpha, \Sigma}$ [2], with α the index of non-compact generators of $\mathrm{SO}(8, 4 + m)$ (or equivalently the fields within the $3d$ truncation) and Σ labelling the scalar harmonic \mathcal{Y}^Σ on S^3 . They transform in the symmetric vector representation $(n/2, n/2; 0, 0)$ of $\mathrm{SO}(4)_{\text{gauge}} \times \mathrm{SO}(4)_{\text{global}} \times \mathrm{SO}(m)$. Under

$$\mathrm{SO}(8, 4 + m) \longrightarrow \mathrm{SO}(4)_{\text{gauge}} \times \mathrm{SO}(4)_{\text{global}} \times \mathrm{SO}(m), \quad (0.1)$$

the non-compact generators t_α of $\mathrm{SO}(8, 4 + m)$ decompose as

$$\begin{aligned} t_\alpha \longrightarrow & (0, 0; 0, 0) \oplus (1, 0; 0, 0) \oplus (0, 1; 0, 0) \oplus (1, 1; 0, 0) \\ & \oplus (1/2, 1/2; 1/2, 1/2) \oplus (1/2, 1/2; 0, 0)^{(m)} \oplus (0, 0; 1/2, 1/2)^{(m)}, \end{aligned} \quad (0.2)$$

where the exponent $^{(m)}$ indicates representations that transforms as vectors under $\text{SO}(m)$. The other representations are $\text{SO}(m)$ scalars. Then, the fluctuations $\Phi^{\alpha, \Sigma}$ are made of the following representations:

$$\begin{aligned}
\alpha \otimes \Sigma \longrightarrow & \left[(0, 0; 0, 0) \oplus (1, 0; 0, 0) \oplus (0, 1; 0, 0) \oplus (1, 1; 0, 0) \right. \\
& \left. \oplus (1/2, 1/2; 1/2, 1/2) \oplus (1/2, 1/2; 0, 0)^{(m)} \oplus (0, 0; 1/2, 1/2)^{(m)} \right] \otimes (n/2, n/2; 0, 0) \\
\longrightarrow & \left[(n/2, n/2; 0, 0) \right] \\
& \oplus \left[((n-2)/2, n/2; 0, 0) \oplus (n/2, n/2; 0, 0) \oplus ((n+2)/2, n/2; 0, 0) \right] \\
& \oplus \left[(n/2, (n-2)/2; 0, 0) \oplus (n/2, n/2; 0, 0) \oplus (n/2, (n+2)/2; 0, 0) \right] \\
& \oplus \left[((n-2)/2, (n-2)/2; 0, 0) \oplus ((n-2)/2, n/2; 0, 0) \oplus ((n-2)/2, (n+2)/2; 0, 0) \right. \\
& \quad \oplus (n/2, (n-2)/2; 0, 0) \oplus (n/2, n/2; 0, 0) \oplus (n/2, (n+2)/2; 0, 0) \\
& \quad \left. \oplus ((n+2)/2, (n-2)/2; 0, 0) \oplus ((n+2)/2, n/2; 0, 0) \oplus ((n+2)/2, (n+2)/2; 0, 0) \right] \\
& \oplus \left[((n-1)/2, (n-1)/2; 1/2, 1/2) \oplus ((n-1)/2, (n+1)/2; 1/2, 1/2) \right. \\
& \quad \left. \oplus ((n+1)/2, (n-1)/2; 1/2, 1/2) \oplus ((n+1)/2, (n+1)/2; 1/2, 1/2) \right] \\
& \oplus \left[((n-1)/2, (n-1)/2; 0, 0)^{(m)} \oplus ((n-1)/2, (n+1)/2; 0, 0)^{(m)} \right. \\
& \quad \left. \oplus ((n+1)/2, (n-1)/2; 0, 0)^{(m)} \oplus ((n+1)/2, (n+1)/2; 0, 0)^{(m)} \right] \\
& \oplus \left[(n/2, n/2; 1/2, 1/2)^{(m)} \right].
\end{aligned} \tag{0.3}$$

This corresponds to the scalars fields and Goldstone modes belonging to the level n in the Kaluza-Klein tower. The supermultiplets at level n are $[3, 4, 2]$ (see tab. 1 for the notations)

$$\mathcal{S}_{(2,0)}^{(n)} = [\overset{\circ}{\mathbf{n} + \mathbf{1}, \mathbf{n} + \mathbf{1}}]_{\text{s}} + [\overset{\square}{\mathbf{n} + \mathbf{3}, \mathbf{n} + \mathbf{3}}]_{\text{s}} + [\overset{\diamond}{\mathbf{n} + \mathbf{2}, \mathbf{n} + \mathbf{2}}]_{\text{s}}^{(m)} + [\overset{\triangle}{\mathbf{n} + \mathbf{1}, \mathbf{n} + \mathbf{3}}]_{\text{s}} + [\overset{\nabla}{\mathbf{n} + \mathbf{3}, \mathbf{n} + \mathbf{1}}]_{\text{s}}. \tag{0.4}$$

In this equation, we have assigned a geometric shape to each of these supermultiplets. We use in the following these shapes to identify the multiplet to which each representation in eq. (0.3) belongs:

$$\begin{aligned}
\alpha \otimes \Sigma \longrightarrow & \left[\underline{(n/2, n/2; 0, 0)} \right] \\
& \oplus \left[\underline{((n-2)/2, n/2; 0, 0)} \oplus \underline{(n/2, n/2; 0, 0)} \oplus \underline{((n+2)/2, n/2; 0, 0)} \right] \\
& \oplus \left[\underline{(n/2, (n-2)/2; 0, 0)} \oplus \underline{(n/2, n/2; 0, 0)} \oplus \underline{(n/2, (n+2)/2; 0, 0)} \right] \\
& \oplus \left[\underline{((n-2)/2, \overset{\circ}{(n-2)/2}; 0, 0)} \oplus \underline{((n-2)/2, n/2; 0, 0)} \oplus \underline{((n-2)/2, \overset{\Delta}{(n+2)/2}; 0, 0)} \right. \\
& \quad \left. \oplus \underline{(n/2, (n-2)/2; 0, 0)} \oplus \underline{(n/2, n/2; 0, 0)} \oplus \underline{(n/2, (n+2)/2; 0, 0)} \right. \\
& \quad \left. \oplus \underline{((n+2)/2, \overset{\nabla}{(n-2)/2}; 0, 0)} \oplus \underline{((n+2)/2, n/2; 0, 0)} \oplus \underline{((n+2)/2, \overset{\square}{(n+2)/2}; 0, 0)} \right] \quad (0.5) \\
& \oplus \left[\underline{((n-1)/2, \overset{\circ}{(n-1)/2}; 1/2, 1/2)} \oplus \underline{((n-1)/2, \overset{\Delta}{(n+1)/2}; 1/2, 1/2)} \right. \\
& \quad \left. \oplus \underline{((n+1)/2, \overset{\nabla}{(n-1)/2}; 1/2, 1/2)} \oplus \underline{((n+1)/2, \overset{\square}{(n+1)/2}; 1/2, 1/2)} \right] \\
& \oplus \left[\underline{((n-1)/2, \overset{\diamond}{(n-1)/2}; 0, 0)}^{(m)} \oplus \underline{((n-1)/2, \overset{\diamond}{(n+1)/2}; 0, 0)}^{(m)} \right. \\
& \quad \left. \oplus \underline{((n+1)/2, \overset{\diamond}{(n-1)/2}; 0, 0)}^{(m)} \oplus \underline{((n+1)/2, \overset{\diamond}{(n+1)/2}; 0, 0)}^{(m)} \right] \\
& \oplus \left[\underline{(n/2, n/2; 1/2, 1/2)}^{(m)} \right].
\end{aligned}$$

Representations to which no shape has been assigned may belong to multiple multiplets (**ce: This needs to be fixed**). Underlined representations are Goldstone modes (or potential Goldstone modes if the line is dashed and not solid).

ExFT cubic Couplings TT term obtained from Cadabra + hand massaging

$$\begin{aligned}
TT = & -\frac{5}{32} T_{ST}^{\Sigma\Omega} T_{UV}^{\Lambda\Delta} c^{\Omega\Delta\Gamma} \delta^{TV} \phi^{SU\Sigma} \phi_W^{W\Gamma\Lambda} + T_{ST}^{\Sigma\Omega} T_{UV}^{\Lambda\Delta} c^{\Omega\Delta\Gamma} \delta^{TV} \phi^{SU\Lambda\Sigma\Gamma} \\
& + \frac{3}{2} T_{ST}^{\Sigma\Omega} T_{UV}^{\Lambda\Sigma} c^{\Omega\Delta\Gamma} \delta^{TV} \phi^{US\Delta\Gamma\Lambda} + \frac{5}{2} T_{ST}^{\Sigma\Omega} T_{UV}^{\Lambda\Sigma} \phi^{TV\Delta} \phi^{SU\Lambda\Gamma} c^{\Omega\Delta\Gamma} - T_{ST}^{\Sigma\Omega} T_{UV}^{\Lambda\Delta} \phi^{SV\Gamma} \phi^{UT\Lambda\Sigma} c^{\Omega\Delta\Gamma} \\
& - T_{ST}^{\Sigma\Omega} T_{UV}^{\Lambda\Sigma} \phi^{TU\Delta} \phi^{SV\Gamma\Lambda} c^{\Omega\Delta\Gamma} - \frac{1}{2} T_{ST}^{\Sigma\Omega} T_{UV}^{\Lambda\Sigma} \phi_W^{U\Delta} \varphi^{ST\Lambda} \varphi^{VW\Gamma} c^{\Omega\Delta\Gamma} \\
& + \frac{1}{2} T_{ST}^{\Sigma\Omega} T_{UV}^{\Lambda\Delta} \phi_W^{S\Sigma} \varphi^{TW\Lambda} \varphi^{UV\Gamma} c^{\Omega\Delta\Gamma} + \frac{1}{4} T_{ST}^{\Sigma\Omega} T_{UV}^{\Lambda\Sigma} \phi_W^{S\Delta} \varphi^{TW\Lambda} \varphi^{UV\Gamma} c^{\Omega\Delta\Gamma} \quad (0.6)
\end{aligned}$$

$$\begin{aligned}
TTCamille := & c^{\Sigma\Omega\Gamma} (\gamma T_{XY}^{\Delta\Omega} T_{WV}^{\Lambda\Sigma} (\frac{3}{4} \Delta^{XW} j^{ST\Lambda} j_{ST}^{\Delta} j^{YV\Gamma} + \Delta^{XW} j^{YSA} j_{SU}^{\Gamma} j^{UV\Delta} \\
& - j^{XU\Lambda} \Delta_U^Y j^{WS\Delta} j_S^{V\Gamma} + \frac{3}{2} j^{XU\Lambda} \Delta_U^Y j^{WS\Gamma} j_S^{V\Delta} - j^{YV\Gamma} j^{XSA} \Delta_{SU} j^{UV\Delta}) \\
& + \gamma T_{XY}^{\Delta\Lambda} T_{WV}^{\Lambda\Sigma} (\Delta^{XW} j^{Y\Sigma\Gamma} j_{SU}^{\Delta} j^{UV\Omega} - \Delta^{XW} j^{YSA} j_{SU}^{\Omega} j^{UV\Gamma} \\
& - 2 \Delta_{SU} j^{YV\Omega} j^{XSA} j^{UV\Delta} + 2 \Delta^{XW} j^{YSA} j_{SU}^{\Gamma} j^{UV\Delta} - 2 \Delta_{SU} j^{YV\Delta} j^{XSA} j^{UV\Gamma})); \quad (0.7)
\end{aligned}$$

Δ_L	Δ_R	Δ	s	$\text{SO}(4)_{\text{gauge}}$	$\text{SO}(4)_{\text{global}}$
Spin-1 multiplet $[\mathbf{k} + \mathbf{1}, \mathbf{k} + \mathbf{1}]_s$					
$k/2$	$k/2$	k	0	$(k/2, k/2)$	$(0, 0)$
$k/2$	$(k+1)/2$	$k+1/2$	1/2	$(k/2, (k-1)/2)$	$(0, 1/2)$
$(k+1)/2$	$k/2$	$k+1/2$	-1/2	$((k-1)/2, k/2)$	$(1/2, 0)$
$(k+1)/2$	$(k+1)/2$	$k+1$	0	$((k-1)/2, (k-1)/2)$	$(1/2, 1/2)$
$k/2$	$(k+2)/2$	$k+1$	1	$(k/2, (k-2)/2)$	$(0, 0)$
$(k+2)/2$	$k/2$	$k+1$	-1	$((k-2)/2, k/2)$	$(0, 0)$
$(k+1)/2$	$(k+2)/2$	$k+3/2$	1/2	$((k-1)/2, (k-2)/2)$	$(1/2, 0)$
$(k+2)/2$	$(k+1)/2$	$k+3/2$	-1/2	$((k-2)/2, (k-1)/2)$	$(0, 1/2)$
$(k+2)/2$	$(k+2)/2$	$k+2$	0	$((k-2)/2, (k-2)/2)$	$(0, 0)$
Spin-2 multiplet $[\mathbf{p}, \mathbf{p} + \mathbf{2}]_s$					
$(p-1)/2$	$(p+1)/2$	p	1	$((p-1)/2, (p+1)/2)$	$(0, 0)$
$(p-1)/2$	$(p+2)/2$	$p+1/2$	3/2	$((p-1)/2, p/2)$	$(0, 1/2)$
$p/2$	$(p+1)/2$	$p+1/2$	1/2	$((p-2)/2, (p+1)/2)$	$(1/2, 0)$
$p/2$	$(p+2)/2$	$p+1$	1	$((p-2)/2, p/2)$	$(1/2, 1/2)$
$(p-1)/2$	$(p+3)/2$	$p+1$	2	$((p-1)/2, (p-1)/2)$	$(0, 0)$
$(p+1)/2$	$(p+1)/2$	$p+1$	0	$((p-3)/2, (p+1)/2)$	$(0, 0)$
$p/2$	$(p+3)/2$	$p+3/2$	3/2	$((p-2)/2, (p-1)/2)$	$(1/2, 0)$
$(p+1)/2$	$(p+2)/2$	$p+3/2$	1/2	$((p-3)/2, p/2)$	$(0, 1/2)$
$(p+1)/2$	$(p+3)/2$	$p+2$	1	$((p-3)/2, (p-1)/2)$	$(0, 0)$

Tab. 1 Spin-1 $[\mathbf{k} + \mathbf{1}, \mathbf{k} + \mathbf{1}]_s$ and spin-2 $[\mathbf{p}, \mathbf{p} + \mathbf{2}]_s$ multiplets of $\text{SU}(2|1, 1)_L \times \text{SU}(2|1, 1)_R$, for $k \geq 2$ and $p \geq 3$ [4]. The $\text{SO}(4)$ representations are given by a couple of $\text{SU}(2)$ spins. The conjugate spin-2 multiplet $[\mathbf{p} + \mathbf{2}, \mathbf{p}]_s$ is obtained by inverting Δ_L with Δ_R , taking the opposite spin $-s$ and exchanging the $\text{SU}(2)$ spins inside each $\text{SO}(4)_{\text{gauged}}$ and $\text{SO}(4)_{\text{global}}$ representations. Taken from ref. [2].

XT terms

$$\begin{aligned}
X\text{Theta} = & c^{\Sigma\Omega\Gamma} (-3\Theta_{KLMN}\Delta^{KP}T_{PQ}\Delta^{\Sigma}j^{MU\Delta}\Delta_U{}^Nj^{LR\Omega}\Delta_{RS}j^{SQ\Gamma} + 6\Theta_{KLMN}\Delta^{KP}T_{PQ}\Delta^{\Sigma}j^{MU\Gamma}j_U{}^N\Delta j^{LQ\Omega} \\
& + 2\Theta_{KLMN}T_{PQ}\Delta^{\Sigma}j^{MU\Gamma}\Delta_U{}^Nj^{KP\Delta}j^{LQ\Omega} - 2\Theta_{KLMN}T_{PQ}\Delta^{\Sigma}j^{MU\Delta}\Delta_U{}^Nj^{KP\Gamma}j^{LQ\Omega} \\
& - 3(\Theta_{KL} + \Theta_{\eta KL})T_{PQ}\Delta^{\Sigma}j^{LQ\Delta}j^{KR\Omega}\Delta_{RS}j^{S\Gamma} \\
& + 3\Theta_{KLMN}(\eta^{KP}\eta^{LQ} - \Delta^{KP}\Delta^{LQ})T_{PQ}\Delta^{\Sigma}\Delta_{RS}j^{MU\Delta}j_U{}^{R\Omega}j^{S\Gamma});
\end{aligned} \tag{0.8}$$

XX terms ($\Theta\Theta$)

Notes on the symmetries of different tensors and usefull identities

- $j_{\bar{M}N} = j_{\bar{M}}^{\bar{N}} \delta_{A\bar{N}}$ is symmetric in $\bar{M} \leftrightarrow \bar{N}$
- $j_{\bar{M}N} \delta^{\bar{N}K} \eta_{\bar{K}L}$ is anti-symmetric in $\bar{M} \leftrightarrow \bar{L}$
- $T.C = 0$ or $T_{..}^{\Sigma\Omega} c^{\Omega\Delta\Gamma} + T_{..}^{\Delta\Omega} c^{\Omega\Sigma\Gamma} + T_{..}^{\Gamma\Omega} c^{\Omega\Sigma\Delta} = 0$
- $f_{\alpha\beta\gamma} = t_{\alpha M}^L t_{\beta L}^K t_{\gamma K}^M = 0$

References

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