

18-660: Numerical Methods for Engineering Design and Optimization

Xin Li

Department of ECE
Carnegie Mellon University
Pittsburgh, PA 15213



Slide 1

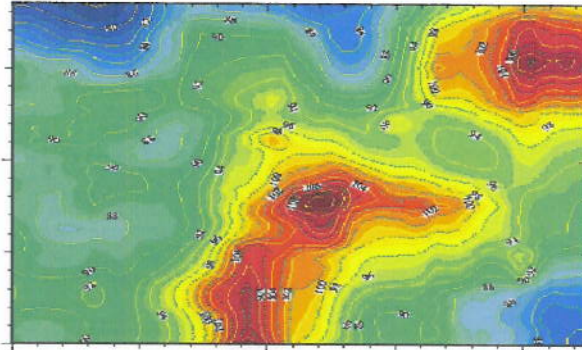
Overview

- Partial Differential Equation (PDE)
 - ▼ Heat equation
 - ▼ Boundary condition

Slide 2

Partial Differential Equation (PDE)

- Partial differential equation is often used to describe a physical system or process
- Example: heat equation is an PDE for thermal analysis
 - ▼ We will derive heat equation step by step in this lecture



Slide 3

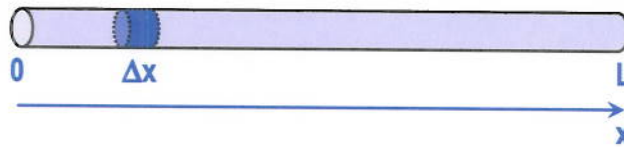
Heat Equation

- Heat equation can be derived from several fundamental physics laws
 - ▼ Fourier's law
 - ▼ Conservation of heat
 - ▼ Etc.
- We will use a 1-D example to illustrate heat equation
 - ▼ Help to get many insights about heat transfer process

Slide 4

1-D Heat Model

- We consider a 1-D rod of length L



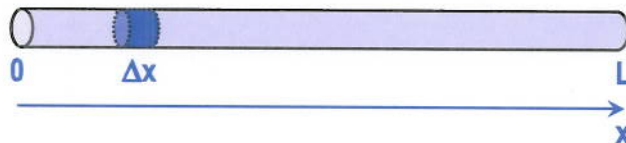
- Three major assumptions

- ▼ Rod is made of a single homogenous conducting material
- ▼ Rod is laterally insulated (heat flows only in the x-direction)
- ▼ Rod is thin (constant temperature at all points of a cross section)

Slide 5

Conservation of Heat

- Apply conservation of heat to the segment $[x, x+\Delta x]$



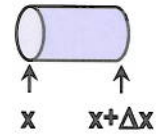
$$\text{Net change of heat inside } [x, x+\Delta x] = \text{Net flux of heat across boundaries} + \text{Total heat generated inside } [x, x+\Delta x]$$

We will look at each of these three components in detail

Slide 6

Conservation of Heat

$$\text{Net change of heat inside } [x, x+\Delta x] = \text{Net flux of heat across boundaries} + \text{Total heat generated inside } [x, x+\Delta x]$$



- Total heat inside $[x, x+\Delta x]$ is equal to:

$$\int_x^{x+\Delta x} \rho \cdot C_p \cdot A \cdot T(s, t) ds$$

ρ : density

C_p : thermal Capacity

A : cross-section Area

T : temperature

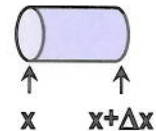
s : x-Coordinate

t : time

Slide 7

Conservation of Heat

$$\text{Net change of heat inside } [x, x+\Delta x] = \text{Net flux of heat across boundaries} + \text{Total heat generated inside } [x, x+\Delta x]$$



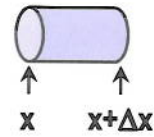
- Net change of heat is equal to:

$$\begin{aligned} & \frac{d}{dt} \left[\int_x^{x+\Delta x} \rho \cdot C_p \cdot A \cdot T(s, t) ds \right] \\ &= \frac{d}{dt} \left[\rho \cdot C_p \cdot A \cdot T(x, t) \cdot \Delta x \right] \quad \Delta x \rightarrow 0 \\ &= \rho \cdot C_p \cdot A \cdot \Delta x \cdot \frac{dT(x, t)}{dt} \end{aligned}$$

Slide 8

Conservation of Heat

$$\text{Net change of heat inside } [x, x+\Delta x] = \text{Net flux of heat across boundaries} + \text{Total heat generated inside } [x, x+\Delta x]$$



- Fourier's law: heat flux across a boundary is proportional to the temperature gradient across the boundary:

$$K \cdot A \cdot \frac{\partial T(x,t)}{\partial x} \quad \leftarrow \text{cylinder} \quad \rightarrow K \cdot A \cdot \frac{\partial T(x+\Delta x, t)}{\partial x}$$

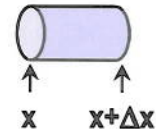
$x \quad x+\Delta x$

K : thermal conductivity

Slide 9

Conservation of Heat

$$\text{Net change of heat inside } [x, x+\Delta x] = \text{Net flux of heat across boundaries} + \text{Total heat generated inside } [x, x+\Delta x]$$



- Net flux of heat across boundaries is equal to:

$$\kappa \cdot A \cdot \frac{\partial T(x,t)}{\partial x} \quad \leftarrow \text{cylinder} \quad \rightarrow \kappa \cdot A \cdot \frac{\partial T(x+\Delta x, t)}{\partial x}$$

$x \quad x+\Delta x$

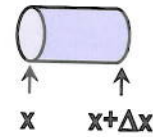
Heat flux at x Heat flux at $x+\Delta x$

$$K \cdot A \cdot \frac{\partial T(x+\Delta x, t)}{\partial x} - K \cdot A \cdot \frac{\partial T(x, t)}{\partial x}$$

Slide 10

Conservation of Heat

$$\text{Net change of heat inside } [x, x+\Delta x] = \text{Net flux of heat across boundaries} + \text{Total heat generated inside } [x, x+\Delta x]$$



- Total heat generated inside $[x, x+\Delta x]$ is equal to:

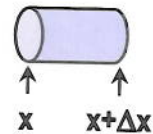
$$\int_x^{x+\Delta x} A \cdot f(s, t) ds$$

f : heat source

Slide 11

1-D Heat Equation

$$\text{Net change of heat inside } [x, x+\Delta x] = \text{Net flux of heat across boundaries} + \text{Total heat generated inside } [x, x+\Delta x]$$



- Overall, we have:

$$\rho \cdot C_p \cdot \cancel{A} \cdot \Delta x \cdot \frac{\partial T(x, t)}{\partial t} = \kappa \cdot \cancel{A} \cdot \left[\frac{\partial T(x + \Delta x, t)}{\partial x} - \frac{\partial T(x, t)}{\partial x} \right] + \int_x^{x+\Delta x} \cancel{A} \cdot f(s, t) \cdot ds$$

Slide 12

1-D Heat Equation

$$\rho \cdot C_p \cdot \frac{\partial T(x,t)}{\partial t} = \kappa \cdot \frac{1}{\Delta x} \left[\frac{\partial T(x+\Delta x, t)}{\partial x} - \frac{\partial T(x, t)}{\partial x} \right] + \frac{1}{\Delta x} \cdot \int_x^{x+\Delta x} f(s, t) \cdot ds$$

$\Delta x \rightarrow 0$

$$\rho \cdot C_p \cdot \frac{dT(x, t)}{dt} = \kappa \cdot \frac{d^2 T(x, t)}{dx^2} + f(x, t)$$

$$\rho \cdot C_p \cdot T_t(x, t) = \kappa \cdot T_{xx}(x, t) + f(x, t)$$

Slide 13

Partial Differential Equation (PDE)

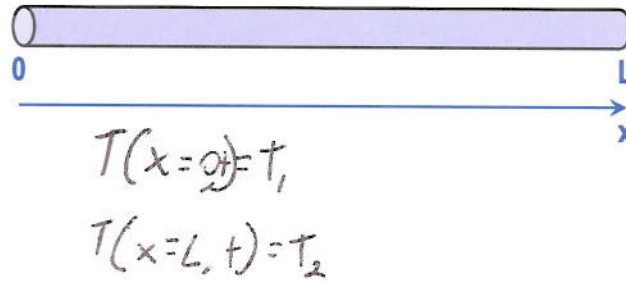
$$\rho \cdot C_p \cdot T_t(x, t) = \kappa \cdot T_{xx}(x, t) + f(x, t)$$

- T (which we differentiate) is called the **dependent variable**
- t and x (which we differentiate with respect to) are called **independent variables**
- In addition to this PDE, we further need to know the boundary and initial conditions to uniquely determine $T(x, t)$

Slide 14

Boundary Conditions

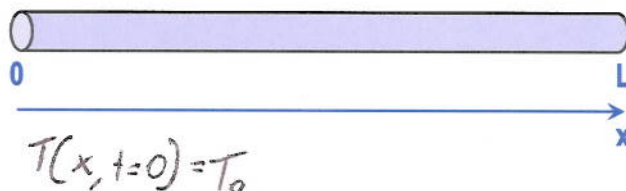
- Boundary conditions describe the physical nature of our problem on the boundaries
- A simple example
 - ▼ Rod temperature is fixed at the two ends



Slide 15

Initial Conditions

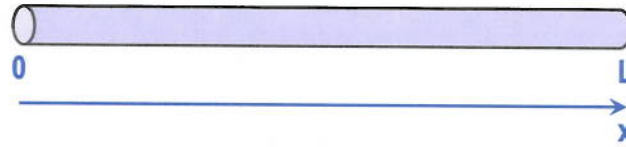
- Initial conditions describe the physical phenomenon at the beginning of the thermal transfer process
- A simple example
 - ▼ Rod is initially at an equilibrium point – constant temperature



Slide 16

1-D Heat Equation

- The complete PDE with boundary and initial conditions



PDE $\rho \cdot C_p \cdot T_t(x,t) = \kappa \cdot T_{xx}(x,t) + f(x,t) \quad (0 \leq x \leq L \quad 0 \leq t \leq \infty)$

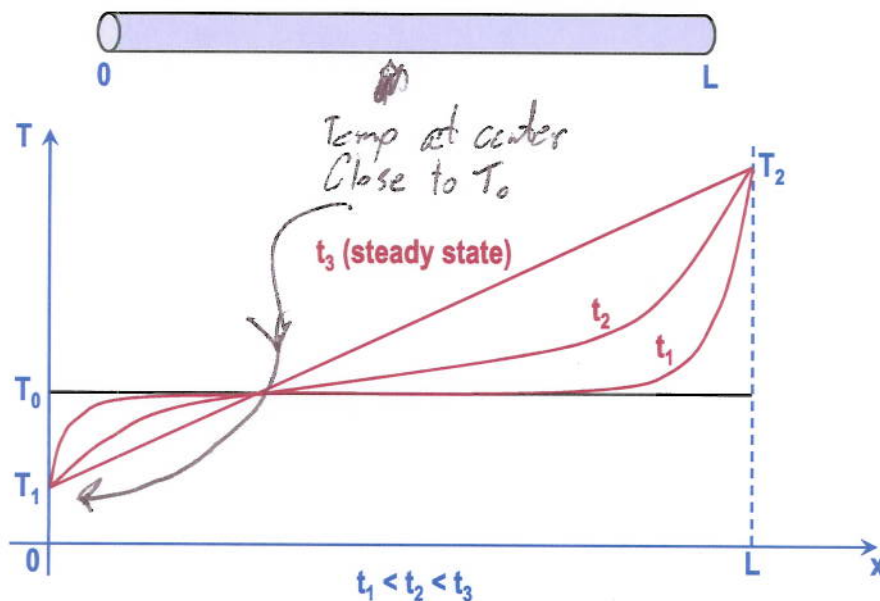
BCs $\begin{cases} T(x=0,t) = T_1 \\ T(x=L,t) = T_2 \end{cases} \quad (0 < t \leq \infty)$

IC $T(x,t=0) = T_0 \quad (0 \leq x \leq L)$

Slide 17

1-D Heat Equation

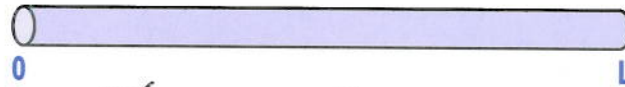
- The solution of this PDE can be analytically calculated



Slide 18

Boundary Conditions

- There are many ways to specify boundary conditions
- Option 1: temperature is specified on boundaries



$$T(x=0, t) = g_1(t)$$
$$T(x=L, t) = g_2(t)$$

Slide 19

Boundary Conditions

- The BCs used for our 1-D rod belongs to this category



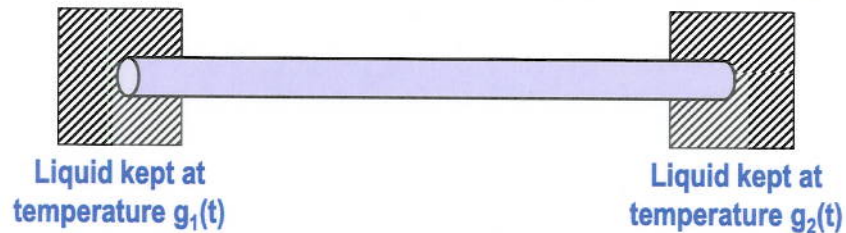
$$T(x=0, t) = T_1$$
$$T(x=L, t) = T_2$$

- Practically useful to force the temperature to behave in a suitable manner
 - ▼ E.g., boundary control in steel industry

Slide 20

Boundary Conditions

- Option 2: temperature of the surrounding medium is specified



- Newton's law: heat flux across the boundary is proportional to the temperature difference

$$h \cdot A \cdot [T(x=0^+) - g_1(t)]$$

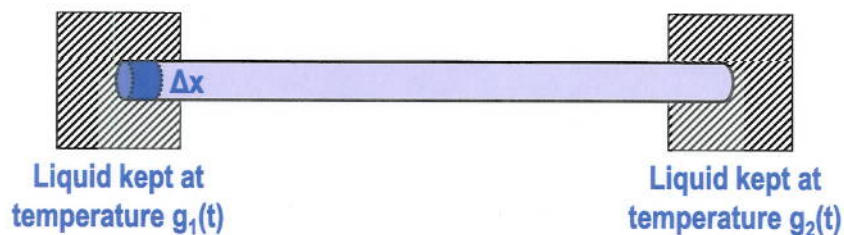
h: heat exchange coefficient

A diagram shows the right end of a rod inserted into a shaded block labeled "Liquid kept at temperature $g_1(t)$ ". A red arrow points from the rod into the liquid, representing heat flux.

Slide 21

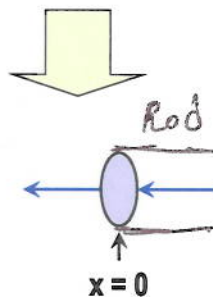
Boundary Conditions

- Consider Newton's and Fourier's laws at $x = 0$



$$h \cdot A \cdot [T(x=0) - g_1(t)]$$

Newton's Law



$$k \cdot A \cdot \frac{dT(x,t)}{dx} \bigg|_{x=0}$$

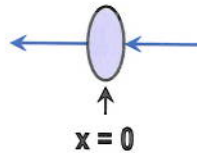
Fourier's Law

Slide 22

Boundary Conditions

$$h \cdot A \cdot [T(x=0, t) - g_1(t)]$$

Heat flux at $x = 0$
(Newton's law)



$$\kappa \cdot A \cdot \left. \frac{\partial T(x, t)}{\partial x} \right|_{x=0}$$

Heat flux at $x = 0$
(Fourier's law)

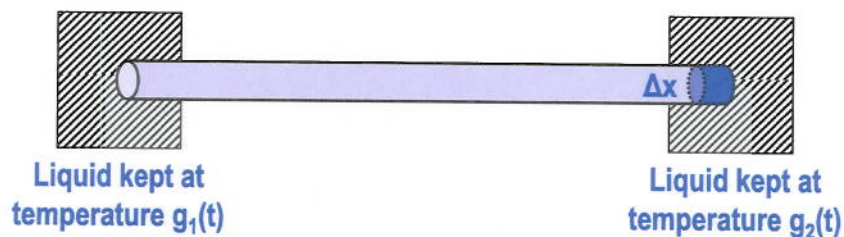
- Apply conservation of heat flux to the boundary $x = 0$

$$\cancel{\kappa \cdot A} \cdot \left. \frac{\partial T(x, t)}{\partial x} \right|_{x=0} = h \cdot \cancel{A} \cdot [T(x=0, t) - g_1(t)]$$

Slide 23

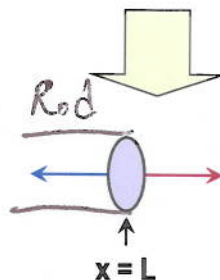
Boundary Conditions

- Similarly, we can derive the boundary condition at $x = L$



$$\kappa \cdot A \cdot \left. \frac{\partial T(x, t)}{\partial x} \right|_{x=L}$$

Fourier's Law

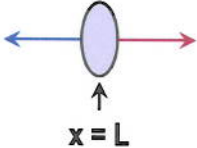


$$h \cdot A \cdot [T(x=L, t) - g_2(t)]$$

Newton's Law

Slide 24

Boundary Conditions

$$\kappa \cdot A \cdot \frac{\partial T(x,t)}{\partial x} \Big|_{x=L} \quad \leftarrow \text{Heat flux at } x = L \text{ (Fourier's law)}$$


$$h \cdot A \cdot [T(x=L,t) - g_2(t)] \quad \text{Heat flux at } x = L \text{ (Newton's law)}$$

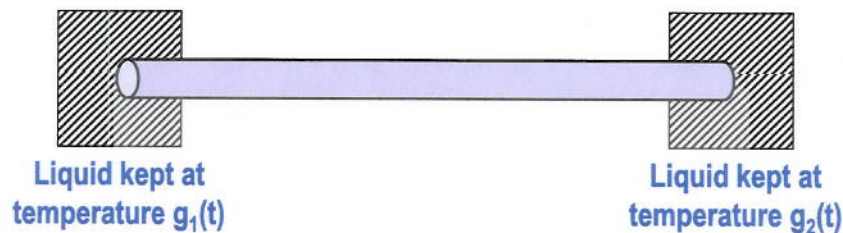
- Apply conservation of heat flux to the boundary $x = L$

$$\cancel{\kappa \cdot A} \cdot \frac{\partial T(x,t)}{\partial x} \Big|_{x=L} = -\cancel{h \cdot A} \cdot [T(x=L,t) - g_2(t)]$$

Slide 25

Boundary Conditions

- Option 2: temperature of the surrounding medium is specified



$$\kappa \cdot \frac{\partial T(x,t)}{\partial x} \Big|_{x=0} = h \cdot [T(x=0,t) - g_1(t)]$$

$$\kappa \cdot \frac{\partial T(x,t)}{\partial x} \Big|_{x=L} = -h \cdot [T(x=L,t) - g_2(t)]$$

Slide 26

Boundary Conditions

- Option 3: heat flow across the boundaries is specified



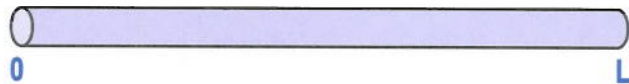
$$\left. \frac{dT(x,t)}{dx} \right|_{x=0} = g_1(t)$$

$$\left. \frac{dT(x,t)}{dx} \right|_{x=L} = g_2(t)$$

Slide 27

Boundary Conditions

- Insulated boundaries (also referred to as reflective boundaries)
 - ▼ No heat passes through boundaries



$$\left. \frac{\partial T(x,t)}{\partial x} \right|_{x=0} = 0$$

$$\left. \frac{\partial T(x,t)}{\partial x} \right|_{x=L} = 0$$

Slide 28

Summary

- **Partial differential equation (PDE)**

- ▼ Heat equation
- ▼ Boundary condition

18-660: Numerical Methods for Engineering Design and Optimization

Xin Li
Department of ECE
Carnegie Mellon University
Pittsburgh, PA 15213



Slide 1

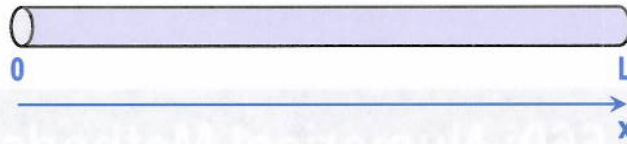
Overview

- Thermal Analysis
 - ▼ 2-D / 3-D heat equation
 - ▼ Finite difference

Slide 2

1-D Heat Equation

- The complete PDE with boundary and initial conditions



PDE $\rho \cdot C_p \cdot T_t(x, t) = \kappa \cdot T_{xx}(x, t) + f(x, t) \quad (0 \leq x \leq L \quad 0 \leq t \leq \infty)$

BCs $\begin{cases} T(x=0, t) = T_1 \\ T(x=L, t) = T_2 \end{cases} \quad (0 < t \leq \infty)$

IC $T(x, t=0) = T_0 \quad (0 \leq x \leq L)$

Slide 3

2-D / 3-D Heat Equation

- 2-D heat equation

$$\rho \cdot C_p \cdot \frac{\partial T(x, y, t)}{\partial t} = \kappa \cdot \nabla^2 T(x, y, t) + f(x, y, t)$$

- 3-D heat equation

$$\begin{array}{ccccc} \text{Density} & & \text{Laplace operator} & & \\ \downarrow & & \downarrow & & \\ \rho \cdot C_p \cdot \frac{\partial T(x, y, z, t)}{\partial t} = \kappa \cdot \nabla^2 T(x, y, z, t) + f(x, y, z, t) & & & & \\ \uparrow & & \uparrow & & \uparrow \\ \text{Thermal capacity} & \text{Thermal conductivity} & & \text{Heat source} & \end{array}$$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

Slide 4

2-D / 3-D Heat Equation

- Heat equation is a 2nd-order linear PDE

$$\rho \cdot C_p \cdot \frac{\partial T(x, y, z, t)}{\partial t} = \kappa \cdot \nabla^2 T(x, y, z, t) + \underline{f(x, y, z, t)}$$

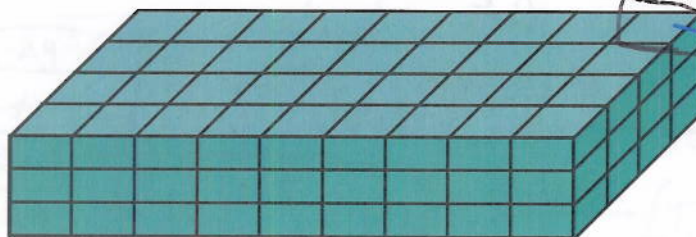
- Order of PDE – the order of the highest partial derivative
- Linearity – the dependent variable T and all its derivatives appear in a linear fashion
- Homogeneity
 - ▼ Homogenous if $f(x, y, z, t) = 0$
 - ▼ Non-homogenous if $f(x, y, z, t) \neq 0$

Slide 5

Finite Difference Method

- PDE can be numerically solved using finite difference method
 - ▼ Discretize 3-D space into a number of small control volumes

$$\rho \cdot C_p \cdot \frac{\partial T(x, y, z, t)}{\partial t} = \kappa \cdot \nabla^2 T(x, y, z, t) + f(x, y, z, t)$$



why? partial derivative of algebraic equations

A control volume

Slide 6

Finite Difference Method

■ We have:

$$I_{i,j,k} = C \cdot \frac{\partial T_{i,j,k}}{\partial t} + G_x \cdot [T_{i,j,k} - T_{i+1,j,k}] + G_x \cdot [T_{i,j,k} - T_{i-1,j,k}] \\ G_y \cdot [T_{i,j,k} - T_{i,j+1,k}] + G_y \cdot [T_{i,j,k} - T_{i,j-1,k}] + G_z \cdot [T_{i,j,k} - T_{i,j,k+1}] + G_z \cdot [T_{i,j,k} - T_{i,j,k-1}]$$

▼ where

$$G_x = \frac{\kappa \cdot \Delta y \cdot \Delta z}{\Delta x} \quad G_y = \frac{\kappa \cdot \Delta x \cdot \Delta z}{\Delta y} \quad G_z = \frac{\kappa \cdot \Delta x \cdot \Delta y}{\Delta z}$$

$$C = \rho \cdot C_p \cdot \Delta x \cdot \Delta y \cdot \Delta z \quad I_{i,j,k} = f_{i,j,k} \cdot \Delta x \cdot \Delta y \cdot \Delta z$$

The discretized thermal equation has a form similar to a circuit equation

Slide 11

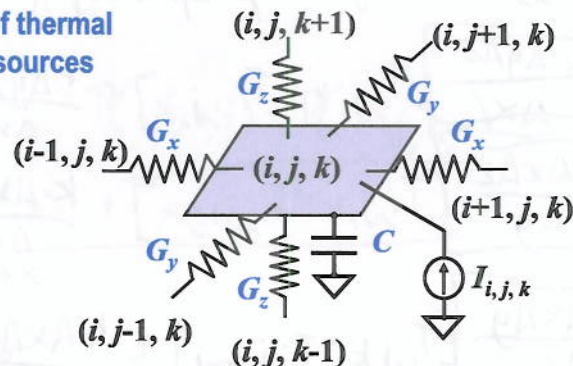
Finite Difference Method

$$I_{i,j,k} = C \cdot \frac{\partial T_{i,j,k}}{\partial t} + G_x \cdot [T_{i,j,k} - T_{i+1,j,k}] + G_x \cdot [T_{i,j,k} - T_{i-1,j,k}] \\ G_y \cdot [T_{i,j,k} - T_{i,j+1,k}] + G_y \cdot [T_{i,j,k} - T_{i,j-1,k}] + G_z \cdot [T_{i,j,k} - T_{i,j,k+1}] + G_z \cdot [T_{i,j,k} - T_{i,j,k-1}]$$

Equivalent circuit consisting of thermal resistors/capacitors and heat sources

T == nodal voltage

I == branch current



Slide 12

Finite Difference Method

$$I_{i,j,k} = C \cdot \frac{\partial T_{i,j,k}}{\partial t} + G_x \cdot [T_{i,j,k} - T_{i+1,j,k}] + G_x \cdot [T_{i,j,k} - T_{i-1,j,k}] \\ G_y \cdot [T_{i,j,k} - T_{i,j+1,k}] + G_y \cdot [T_{i,j,k} - T_{i,j-1,k}] + G_z \cdot [T_{i,j,k} - T_{i,j,k+1}] + G_z \cdot [T_{i,j,k} - T_{i,j,k-1}]$$

- The operator $\partial/\partial t$ can be handled by numerical integration

- ▼ We need to solve a large-scale linear equation to find $T_{i,j,k}(t_n)$ at each time point t_n

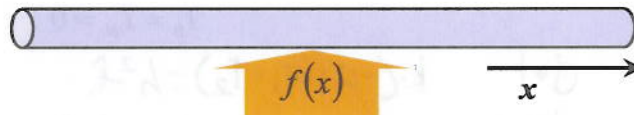
- Generally interested only in steady state – thermal capacitance is not considered

Slide 13

1-D Thermal Analysis Example

- 1-D PDE to describe the steady-state temperature distribution along a uniform rod at $[0, 1]$

$$T(x, 0) = T_{init} \\ T(x = 0, t) = T(x = 1, t) = 0$$



$$\rho \cdot C_p \cdot \frac{\partial T(x, t)}{\partial t} = \kappa \cdot \frac{\partial^2 T(x, t)}{\partial x^2} + f(x, t)$$

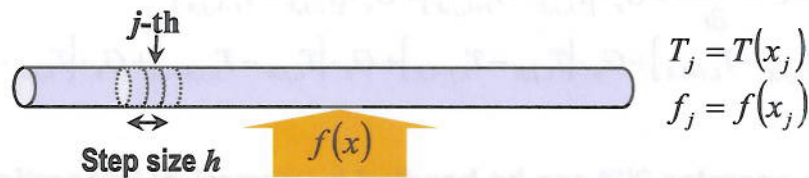
steady state $\frac{dT(x, t)}{dt} = 0$

$$-k \cdot T_{xx}(x, t) = f(x, t) \quad (0 < x < 1)$$

Slide 14

1-D Thermal Analysis Example

- Approximate 2nd order derivative using finite difference



$$T(x=0,t) = T(x=1,t) = 0 \quad -\kappa \cdot T_{xx}(x) = f(x) \quad T_{xx}(x_j) = \frac{T_{j+1} + T_{j-1} - 2T_j}{h^2}$$

$$-\kappa \frac{T_{j+1} + T_{j-1} - 2T_j}{h^2} = f_j \quad (1 \leq j \leq N-1)$$

$$-\kappa \cdot [T_{j+1} + T_{j-1} - 2T_j] = h^2 \cdot f_j$$

$$T_0 = T_N = 0$$

Slide 15

1-D Thermal Analysis Example

- The linear system is:

$$\kappa \cdot (-T_{j-1} + 2T_j - T_{j+1}) = h^2 \cdot f_j \quad (1 \leq j \leq N-1)$$

$$T_0 = T_N = 0$$

$$j=1 \quad \kappa \cdot (-0 + 2T_1 - T_2) = h^2 \cdot f_1$$

$$j=2 \quad \kappa \cdot (-T_1 + 2T_2 - T_3) = h^2 \cdot f_2$$

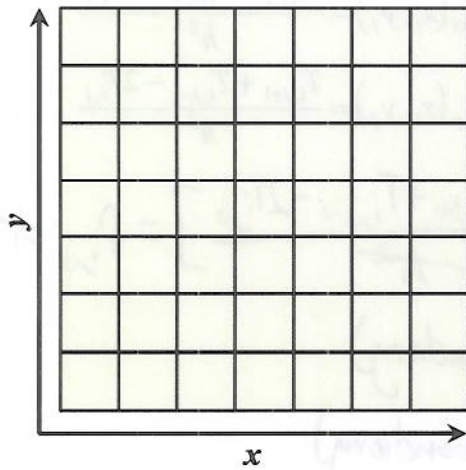
$$j=N-1 \quad \kappa \cdot (-T_{N-2} + 2T_{N-1} - 0) = h^2 \cdot f_{N-1}$$

$$T_1, T_2, T_3, \dots, T_{N-1}$$

Slide 16

2-D Thermal Analysis Example

- 2-D PDE to describe the steady-state temperature distribution over a uniform plane $x, y \in [0, 1]$



$$T(x, y, 0) = T_{init}$$

$$T(x=0, t) = T(x=1, t) = 0$$

$$T(y=0, t) = T(y=1, t) = 0$$

$$\rho \cdot C_p \cdot \frac{\partial T}{\partial t} = \kappa \cdot \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right) + f(x, y, t)$$

Steady state $\frac{dT}{dt} = 0$

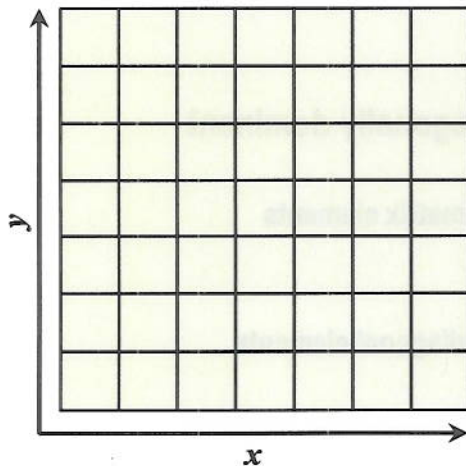
$$-\kappa \cdot (T_{xx} + T_{yy}) = f(x, y)$$

$$(0 < x, y < 1)$$

Slide 17

2-D Thermal Analysis Example

- Approximate 2nd order derivative using finite difference



$$T(x=0, t) = T(x=1, t) = 0$$

$$T(y=0, t) = T(y=1, t) = 0$$

$$-\kappa \cdot (T_{xx} + T_{yy}) = f(x, y)$$

$$T_{i,j} = T(x_i, y_j)$$

$$f_{i,j} = f(x_i, y_j)$$

$$T_{xx}(x_i, y_j) = \frac{T_{i+1,j} + T_{i-1,j} - 2T_{i,j}}{h^2}$$

$$T_{yy}(x_i, y_j) = \frac{T_{i,j+1} + T_{i,j-1} - 2T_{i,j}}{h^2}$$

Slide 18

2-D Thermal Analysis Example

- The linear system is:

$$T(x=0,t) = T(x=1,t) = 0 \quad T_{xx}(x_i, y_j) = \frac{T_{i+1,j} + T_{i-1,j} - 2T_{i,j}}{h^2}$$

$$T(y=0,t) = T(y=1,t) = 0$$

$$-\kappa \cdot (T_{xx} + T_{yy}) = f(x, y) \quad T_{yy}(x_i, y_j) = \frac{T_{i,j+1} + T_{i,j-1} - 2T_{i,j}}{h^2}$$

$$-\kappa \left[\frac{T_{i+1,j} + T_{i-1,j} - 2T_{i,j}}{h^2} + \frac{T_{i,j+1} + T_{i,j-1} - 2T_{i,j}}{h^2} \right] = f_{i,j} \cdot h^2$$

(Not at boundary)

$T_{i,j} = 0$ (at boundary)

Slide 19

Thermal Analysis

- Thermal analysis generally requires to solve a large-scale linear equation

$$A \cdot X = B$$

- The matrix A is **symmetric** and **diagonally dominant**

$$A_{ij} = A_{ji} \quad \text{For all matrix elements}$$

$$|A_{ii}| \geq \sum_{j=1, j \neq i}^N |A_{ij}| \quad \text{For all diagonal elements}$$

Slide 20

Thermal Analysis

■ 1-D thermal analysis example

$$\kappa \cdot (-T_{j-1} + 2T_j - T_{j+1}) = h^2 \cdot f_j \quad (1 \leq j \leq N-1)$$

$$T_0 = T_N = 0$$

$j=1$
 $j=2$

$$\begin{bmatrix}
 2 & -1 & & & \\
 -1 & 2 & -1 & & \\
 & \ddots & \ddots & \ddots & \\
 & & -1 & 2 & -1 \\
 & & & -1 & 2
 \end{bmatrix}
 \begin{bmatrix}
 T_1 \\
 T_2 \\
 T_3 \\
 \vdots \\
 T_{N-1}
 \end{bmatrix}
 = \frac{h^2}{\kappa}
 \begin{bmatrix}
 f_1 \\
 f_2 \\
 \vdots \\
 f_{N-1}
 \end{bmatrix}$$

Symmetric
 diagonally dominant

Slide 21

Thermal Analysis

■ A matrix "A" is **positive definite**, if

- ▼ A is symmetric and
- ▼ A is diagonally dominant and
- ▼ All diagonal elements of A are non-negative and
- ▼ A is not singular
- ▼ Sufficient but NOT necessary condition

■ Definition of positive definite matrix

$$P^T \cdot A \cdot P > 0 \quad \text{for any real-valued vector } P \neq 0$$

All eigenvalues of A are positive

Slide 22

Thermal Analysis

- Positive definite linear equation $AX = B$ can be solved by efficient numerical algorithms
 - ▼ Cholesky decomposition
 - ▼ Conjugate gradient method
 - ▼ Etc.
- We will try to cover some of these efficient algorithms in future lectures

Slide 23

Summary

- Thermal analysis
 - ▼ 2-D / 3-D heat equation
 - ▼ Finite difference

Slide 24