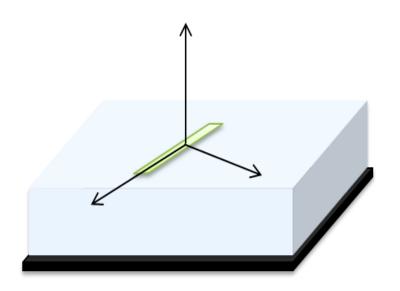
Question 1

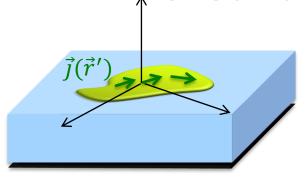
Write a Matlab routine that finds the propagation constant of the first two surface waves (TM0 and TE1) supported by a grounded slab.

To check the program, provide a plot showing the propagation constant, normalized to k_0 , of the TM0 and TE1 surface waves for a frequency range from 1 to 20 GHz. Consider h = 2 mm, $\varepsilon_r = 10$ and $\varepsilon_r = 5$.





Grounded Dielectric Substrate

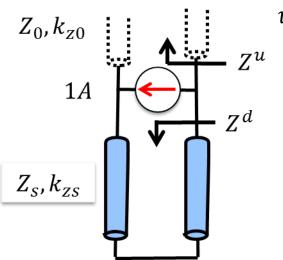


$$Z_{in} = \frac{Z_u Z_d}{Z_u + Z_d}$$

$$Z_{in} = \frac{Z_u Z_d}{Z_u + Z_d}$$

$$Z_u = Z_0$$

$$Z_d = jZ_s \tan(k_{zs}h)$$



$$v(z=z_s)=Z_{in}$$

All the voltage and current solutions at any z-quote are expressed as a function of this voltage



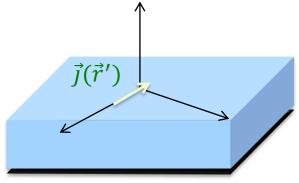
The poles in the transversal complex plane can be found solving the following dispersion equation

$$D(k_{\rho}) = Z_u + Z_d = 0$$

for TE and TM



Grounded Dielectric Substrate



All the voltage and current solutions at any z-quote are expressed as a function of the same denominator

Voltage in the slab:

$$V_{S} = \frac{Z_{u}Z_{d}}{Z_{u} + Z_{d}} \frac{\sin(k_{zS}z)}{\sin(k_{zS}h)} = \frac{Z_{u}Z_{d}}{D(k_{\rho})} \frac{\sin(k_{zS}z)}{\sin(k_{zS}h)}$$

Current in the slab:

$$I_{s} = \frac{1}{Z_{s}} \frac{Z_{u}Z_{d}}{D(k_{\rho})} \frac{j\cos(k_{zs}z)}{\sin(k_{zs}h)}$$

Voltage in the air:

$$V_0 = \frac{Z_u Z_d}{D(k_o)} e^{jk_{z0}h} e^{-jk_{z0}z}$$

Current in the air:

$$I_0 = \frac{1}{Z_0} \frac{Z_u Z_d}{D(k_0)} e^{jk_{z0}h} e^{-jk_{z0}z}$$



Solving Dispersion Equations

If we know a good guess of the solution, $k_{ ho}^g$

1) One can expand the denominator around this point using the Taylor's series:

$$D(k_{\rho}) \approx D(k_{\rho}^{g}) + D'(k_{\rho}^{g})(k_{\rho} - k_{\rho}^{g})$$

2) Evaluating this expansion in the actual zero, $k_{\rho 0}$:

$$D(k_{\rho 0}) \approx D(k_{\rho}^g) + D'(k_{\rho}^g)(k_{\rho 0} - k_{\rho}^g) = 0$$



The derivative can be done numerically

$$k_{\rho 0} = k_{\rho}^g - \frac{D(k_{\rho}^g)}{D'(k_{\rho}^g)}$$

$$D'(k_{\rho}^g) \approx \frac{D(k_{\rho}^g + \Delta k/2) - D(k_{\rho}^g - \Delta k/2)}{\Delta k} \quad \Delta k = k_0/500$$

- Write a separate routine to calculate the dispersion equation: $[D] = Den_GroundSlab(k0, er, h, kro, 'TE/TM')$

$$D(k_{\rho}) = Z_u + Z_d$$

- Write a separate routine to calculate the propagation constant: $[krho] = findprop(k0, er, h, \mathbf{k}_{\rho}^{g}, 'TE/TM')$

$$k_{\rho} = k_{\rho}^{g} - \frac{D(k_{\rho}^{g})}{D'(k_{\rho}^{g})}$$

$$D'(k_{\rho}^{g}) \approx \frac{D(k_{\rho}^{g} + \Delta k/2) - D(k_{\rho}^{g} - \Delta k/2)}{\Delta k}$$
$$\Delta k = k_{0}/500$$



Iterative Frequency Loop for solving the dispersion equation

- 1) Plot the Denominator vs k_{ρ} / k_{0} at the highest frequency to find a good initial guess: k_{ρ}^{g}
- 2) Use this guess to find the propagation constant at the highest frequency

$$k_{\rho} (f_i) = k_{\rho}^g - \frac{D(k_{\rho}^g)}{D'(k_{\rho}^g)}$$

3) Do a loop from the highest to the lowest frequency where you use the *normalized* propagation constant, $k_{\rho}^{gn} = k_{\rho} (f_i)/k_0(f_i)$, of the previous step as a guess to the next frequency k_{ρ}^0 :

$$k_{\rho} (f_i + 1) = k_{\rho}^g - \frac{D(k_{\rho}^g)}{D'(k_{\rho}^g)}$$

with $k_{\rho}^g = k_{\rho}^{gn} k_0 (f_i + 1)$

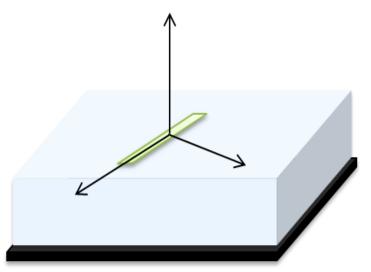


Question 2

Write a Matlab routine to calculate the electric field contribution, in cylindrical coordinates, associated to the TM0 surface wave excited by a half-wavelength dipole on top of a grounded slab.

Provide the following plots (h = 2 mm, $\varepsilon_r = 10$, f = 10 GHz):

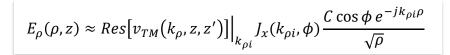
- Real and imaginary part variation of the electric field as a function of the radial distance.
- Amplitude variation of the electric field as a function of *z*.
- Amplitude variation of the electric field as a function of ϕ .





TM surface wave

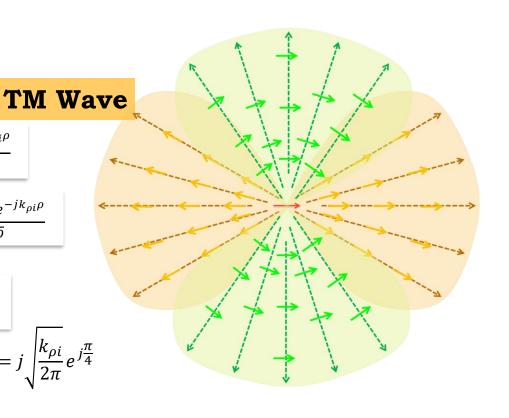
Expressions for an electric current oriented along x



$$E_z(\rho,z) \approx -\frac{\varsigma k_{\rho i}}{k} Res \big[i_{TM} \big(k_\rho, z, z' \big) \big] \bigg|_{k_{\rho i}} J_x(k_{\rho i}, \phi) \frac{C \cos \phi \, e^{-j k_{\rho i} \rho}}{\sqrt{\rho}}$$

$$H_{\phi}(\rho,z) \approx Res \left[i_{TM}(k_{\rho},z,z')\right] \Big|_{k_{\rho i}} J_{x}(k_{\rho i},\phi) \frac{C\cos\phi e^{-jk_{\rho i}\rho}}{\sqrt{\rho}}$$

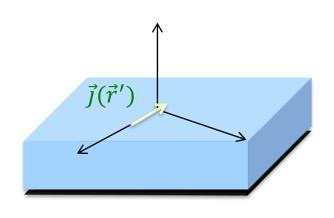
$$C = j \sqrt{\frac{k_{\rho i}}{2\pi}} e^{j\frac{\pi}{4}}$$



 $k_{
ho i}$: pole location in the spectrum







$$E_{\rho}(\rho, z) \approx Res \left[v_{TM}(k_{\rho}, z, z') \right] \Big|_{k_{\rho i}} J_{x}(k_{\rho i}, \phi) \frac{C \cos \phi \, e^{-jk_{\rho i}\rho}}{\sqrt{\rho}}$$

$$E_z(\rho,z) \approx -\frac{\varsigma k_{\rho i}}{k} Res \big[i_{TM} \big(k_\rho, z, z' \big) \big] \Big|_{k_{\rho i}} \boldsymbol{J}_x(k_{\rho i}, \phi) \frac{C \cos \phi \, e^{-jk_{\rho i}\rho}}{\sqrt{\rho}}$$

$$H_{\phi}(\rho, z) \approx Res[i_{TM}(k_{\rho}, z, z')]\Big|_{k_{\rho i}} J_{x}(k_{\rho i}, \phi) \frac{C\cos\phi e^{-jk_{\rho i}\rho}}{\sqrt{\rho}}$$

Fields inside the slab:

$$\begin{split} Res\big[v_{TE/TM}\big(k_{\rho},z,z'\big)\big]\Big|_{k_{\rho i}} &= \frac{Z_{u}Z_{d}}{D'(k_{\rho i})}\frac{sin(k_{zs}z)}{sin(k_{zs}h)} \\ Res\big[i_{TE/TM}\big(k_{\rho},z,z'\big)\big]\Big|_{k_{\rho i}} &= \frac{Z_{u}Z_{d}}{D'(k_{\rho i})}\frac{1}{Z_{s}}\frac{j\cos(k_{zs}z)}{sin(k_{zs}h)} \end{split}$$

Fields inside the air:

$$\begin{split} Res \big[v_{TE/TM} \big(k_{\rho}, z, z' \big) \big] \Big|_{k_{\rho i}} &= \frac{Z_{u} Z_{d}}{D'(k_{\rho i})} e^{j k_{z0} h} e^{-j k_{z0} z} \\ Res \big[i_{TE/TM} \big(k_{\rho}, z, z' \big) \big] \Big|_{k_{\rho i}} &= \frac{Z_{u} Z_{d}}{D'(k_{\rho i})} \frac{1}{Z_{0}} e^{j k_{z0} h} e^{-j k_{z0} z} \end{split}$$

$$D'(k_{\rho}^{g}) \approx \frac{D(k_{\rho}^{g} + \Delta k/2) - D(k_{\rho}^{g} - \Delta k/2)}{\Delta k}$$



- Write a separate routine to calculate the residues: $[VtmR, ItmR] = Residue_GroundSlab(k0, er, h, ksw, z)$

Fields inside the slab:

$$\begin{split} Res \big[v_{TE/TM} \big(k_{\rho}, z, z' \big) \big] \Big|_{k_{\rho i}} &= \frac{Z_u Z_d}{D'(k_{\rho i})} \frac{sin(k_{zs} z)}{sin(k_{zs} h)} \\ Res \big[i_{TE/TM} \big(k_{\rho}, z, z' \big) \big] \Big|_{k_{\rho i}} &= \frac{Z_u Z_d}{D'(k_{\rho i})} \frac{1}{Z_s} \frac{j \cos(k_{zs} z)}{sin(k_{zs} h)} \end{split}$$

Fields inside the air:

$$\begin{split} Res \big[v_{TE/TM} \big(k_{\rho}, z, z' \big) \big] \Big|_{k_{\rho i}} &= \frac{Z_{u} Z_{d}}{D'(k_{\rho i})} e^{j k_{z0} h} e^{-j k_{z0} z} \\ Res \big[i_{TE/TM} \big(k_{\rho}, z, z' \big) \big] \Big|_{k_{\rho i}} &= \frac{Z_{u} Z_{d}}{D'(k_{\rho i})} \frac{1}{Z_{0}} e^{j k_{z0} h} e^{-j k_{z0} z} \end{split}$$

$$k_{\rho i} = ksw$$

- Write a separate routine to calculate the field components: $[Erho, Ez, Hphi] = TMSwFields(k0, ksw, er, VtmR, ItmR, \rho, \phi, z)$

$$E_{\rho}(\rho,z) \approx Res \left[v_{TM}(k_{\rho},z,z') \right] \Big|_{k_{\rho i}} \boldsymbol{J}_{x}(k_{\rho i},\phi) \frac{C\cos\phi \, e^{-jk_{\rho i}\rho}}{\sqrt{\rho}}$$

$$E_z(\rho,z) \approx -\frac{\varsigma k_{\rho i}}{k} Res \big[i_{TM} \big(k_\rho, z, z' \big) \big] \Big|_{k_{\rho i}} \boldsymbol{J}_x(k_{\rho i}, \phi) \frac{C \cos \phi \, e^{-jk_{\rho i}\rho}}{\sqrt{\rho}}$$

$$H_{\phi}(\rho, z) \approx Res \left[i_{TM}(k_{\rho}, z, z')\right] \Big|_{k_{\rho i}} J_{x}(k_{\rho i}, \phi) \frac{C \cos \phi \, e^{-jk_{\rho i}\rho}}{\sqrt{\rho}}$$

FT of the current distribution (PWS): Jx = FTcurrent(k0, kx, ky, l, w)

$$k_{x} = k_{sw} cos \phi$$
$$k_{y} = k_{sw} sin \phi$$

$$C = j \sqrt{\frac{k_{\rho i}}{2\pi}} e^{j\frac{\pi}{4}}$$

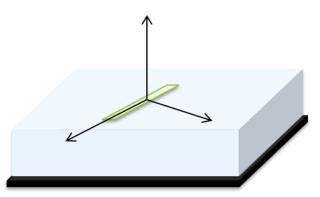


Question 3

Write a Matlab routine to calculate the power radiated into the TM0 surface wave of a dipole antenna printed on top of a grounded slab.

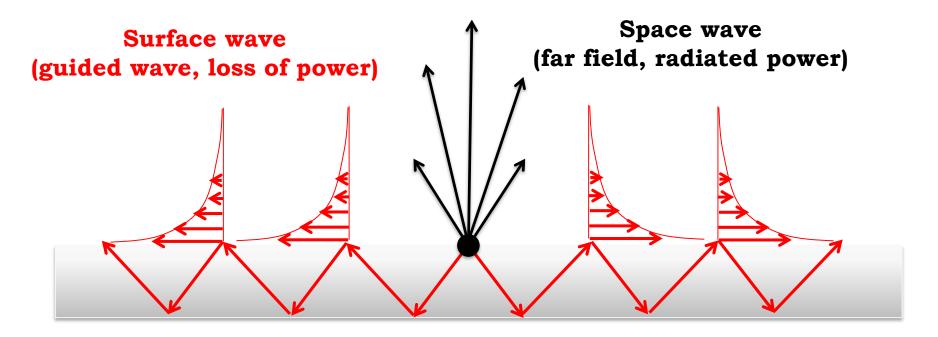
Provide the following plots (h = 2 mm, $\varepsilon_r = 10$, f = 1 - 15 GHz):

- The far field power and surface wave power radiated by an elementary dipole (J = 1) normalized to the power radiated by the same dipole in free space.
- The TM0 surface wave efficiency for the elementary dipole and for a dipole with l = 5.3 mm and w = 0.5 mm (assume a PWS distribution).
- The TM0 surface efficiency for a *uniform* current distribution in both x and y directions with a dimension of l = w = 25 mm.





Antenna SW Efficiency



We are going to characterize the loss of efficiency into SW:

- 1. Find the SW propagation constant (dispersion equation)
- 2. Calculate the SW fields (Residue)
- 3. Power into SW vs radiated power

$$\eta_{sw} = \frac{P_{rad}}{P_{rad} + P_{sw}}$$

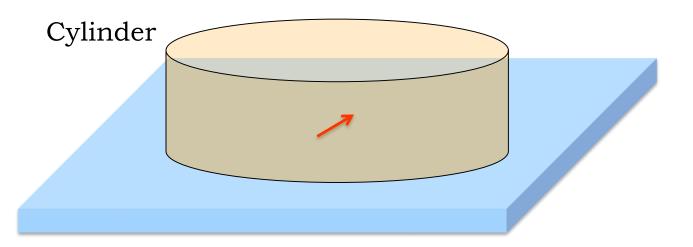


SW Power

We need to characterize the power launched into surface waves

By integrating the Poynting vector around a closed surface surrounding the antenna

$$P_{sw} = \frac{1}{2} \sum_{i} \iint_{S} Re[\overrightarrow{E}_{sw_{i}} \times \overrightarrow{H}_{sw_{i}}^{*}] d\overrightarrow{s}$$



Since the fields at the top and bottom of the cylinder are zero

$$P_{sw}^{TM} = \frac{1}{2} \int_{0}^{2\pi} \int_{0}^{\infty} Re\left[-E_z H_{\phi}^*\right] \rho_s dz d\phi$$



TM SW Power

$$P_{sw}^{TM} = \frac{1}{2} \int_{0}^{2\pi} \int_{0}^{\infty} Re\left[-E_z H_{\phi}^*\right] \rho_s dz d\phi$$

$$E_z(\rho,z) \approx -\frac{\zeta k_{\rho i}}{k} Res \left[i_{TM} \left(k_{\rho}, z, z' \right) \right] \Big|_{k_{\rho i}} C_x(k_{\rho i}, \phi) \frac{C \cos \phi \, e^{-jk_{\rho i}\rho}}{\sqrt{\rho}}$$

$$H_{\phi}(\rho,z) \approx Res \left[i_{TM}(k_{\rho},z,z')\right] \Big|_{k_{\rho i}} C_{x}(k_{\rho i},\phi) \frac{C\cos\phi \, e^{-jk_{\rho i}\rho}}{\sqrt{\rho}}$$

 $C = j \sqrt{\frac{k_{\rho i}}{2\pi}} e^{j\frac{\pi}{4}}$

$$P_{sw}^{TM} = \frac{1}{2} \frac{k_{\rho i}^{2}}{2\pi} \int_{0}^{\infty} \frac{\zeta}{k} \left| Res \left[i_{TM}(k_{\rho}, z, z') \right] \right|_{k_{\rho i}} \right|^{2} dz \int_{0}^{2\pi} \left| C_{x}(k_{\rho i}, \phi) \right|^{2} \cos^{2} \phi \, d\phi$$

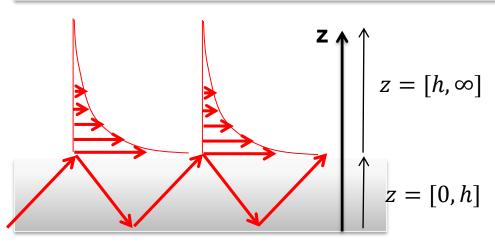
The integral in z and phi are not coupled They can be performed independently for faster computation

$$P_{sw}^{TM} = \frac{1}{2} \frac{k_{\rho i}^2}{2\pi} I_{\mathbf{z}} I_{\phi}$$



Z-integration

$$I_{z} = \int_{0}^{\infty} \frac{\zeta}{k} \left| Res[i_{TM}(k_{\rho}, z, z')] \right|_{k_{\rho i}} \left|^{2} dz = \frac{\zeta_{0}}{k_{0}} \int_{0}^{h} \frac{1}{\epsilon_{r}} \left| Res[i_{TM}(k_{\rho}, z, z')] \right|_{k_{\rho i}} \left|^{2} dz + \frac{\zeta_{0}}{k_{0}} \int_{h}^{\infty} \left| Res[i_{TM}(k_{\rho}, z, z')] \right|_{k_{\rho i}} \right|^{2} dz$$



The z-integration domain is divided in

The z-integration domain is divided in two regions associated to the different dielectrics
$$z = [h, \infty]$$

$$Res[i_{TM}(k_{\rho}, z, z')] = \begin{cases} \frac{Z_u Z_d}{D'(k_{\rho i})} \frac{1}{Z_s} \frac{j \cos(k_{zs} z)}{\sin(k_{zs} h)} \\ \frac{Z_u Z_d}{D'(k_{\rho i})} \frac{1}{Z_0} e^{jk_{z0}h} e^{-jk_{z0}z} \end{cases}$$

$$k_{z0} = \sqrt{k_0^2 - k_{\rho i}^2} = -j\sqrt{k_{\rho i}^2 - k_0^2} = -j\alpha_{z0}$$

It is possible doing the z-integration in analytical form

$$\int_{0}^{h} \cos^{2}(k_{zs}z) dz = \frac{h}{2}(1 + sinc(2k_{zs}h))$$

$$\int_{h}^{\infty} e^{-2\alpha_{z0}(z-h)} dz = \frac{1}{2\sqrt{k_{\rho i}^{2} - k_{0}^{2}}}$$



phi-integration: numerical

$$I_{\phi} = \int_{0}^{2\pi} \left| \boldsymbol{J}_{x}(k_{\rho i}, \phi) \right|^{2} \cos^{2} \phi \, d\phi$$

The integral in phi can be done numerically for any current distribution

However if we consider a short dipole
$$(J_x(k_{\rho i}, \phi) = I_0 l)$$

 $I_{\phi} = \pi$



- Write a separate routine to calculate the power radiated into the surface wave for an elementary dipole:

[Psw] = PswTMelem(k0, er, h, ksw)

$$\left| P_{sw}^{TM} \right|_{\delta} = \frac{1}{2} \frac{k_{\rho i}^{2}}{2\pi} I_{\mathbf{z}} I_{\phi}^{\delta} \qquad I_{\phi}^{\delta} = \pi$$

$$I_{z} = \int_{0}^{\infty} \frac{\zeta}{k} \left| Res[i_{TM}(k_{\rho}, z, z')] \right|_{k_{\rho i}} \right|^{2} dz \quad Numeric \ or \ analytic$$

- Write a separate routine to calculate the FT of a uniform current distribution

$$\tilde{J}_{x}(k_{x}, k_{y}) = Lsinc(\frac{k_{x}L}{2})sinc(\frac{k_{y}w}{2})$$

- Write a separate routine to calculate the phi integral for the specific current distribution:

$$[I_{\phi}] = PhiInt_PWS(k0, ksw, l, w)$$

$$[I_{\phi}] = PhiInt_Uniform(k0, ksw, l, w)$$

$$I_{\phi} = \int_{0}^{2\pi} \left| \boldsymbol{J}_{x}(k_{\rho i}, \phi) \right|^{2} \cos^{2} \phi \, d\phi$$

The power can be evaluated using the power of the elementary current $(P_{sw}^{TM}|_{\delta})$ via



$$P_{SW}^{TM} = P_{SW}^{TM} \Big|_{\delta} \frac{I_{\phi}}{I_{\phi}^{\delta}}$$