

Sharp Corner Singularity of the White–Metzner Model

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Viscoelastic?

- Many materials we encounter have **viscoelastic** properties.
 - Viscosity: measure of resistance to flow.
 - Elasticity: ability of a material to return to its original state after removal of deforming forces.
- Characterised by a *relaxation time* — quantifies ‘memory’ of deformation history.



iStock, (2021)



Cabrera, X., (2020)



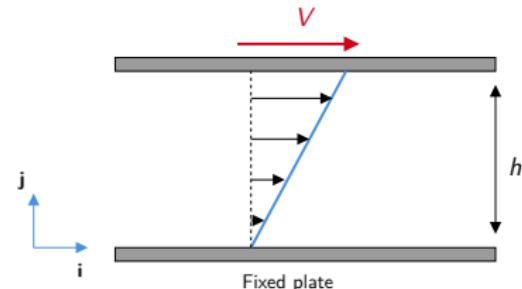
Janine, (2009)

Properties in Shearing Flow

- In (simple) shear, viscosity μ should vary with shear rate $\dot{\gamma}$.

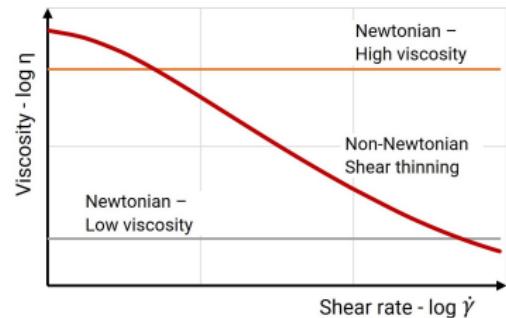
$$\mathbf{u} = (\dot{\gamma}y, 0, 0)^T, \quad \dot{\gamma} = \frac{V}{h},$$

$$\psi = \frac{1}{2}\dot{\gamma}y^2.$$



- Material is called:

- Shear thinning: μ **decreases** with increasing $\dot{\gamma}$.
- Shear thickening: μ **increases** with increasing $\dot{\gamma}$.



Østergård, A. L., 2020.

Viscoelastic Modelling

Overview

So, how do we model these materials?

Conservation of mass and momentum:

$$\nabla \cdot \mathbf{u} = 0, \quad \rho \left(\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right) = -\nabla p + \nabla \cdot \boldsymbol{\tau}, \quad \boldsymbol{\tau}^T = \boldsymbol{\tau}.$$

How do we specify $\boldsymbol{\tau}$?

Viscoelastic Modelling

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How do we specify $\boldsymbol{\tau}$?

$$\boldsymbol{\tau} + \lambda(\dot{\gamma}) \overset{\nabla}{\boldsymbol{\tau}} = 2\mu(\dot{\gamma}) \boldsymbol{D}$$

This is the *White-Metzner* model!

White and Metzner, (1963)

Viscoelastic Modelling

The White–Metzner (WM) Model

$$\overset{\nabla}{T} + \lambda(\dot{\gamma}) \overset{\nabla}{T} = 2\mu(\dot{\gamma}) D$$

Involves an *upper convected derivative*:

$$\overset{\nabla}{T} := \underbrace{\frac{\partial T}{\partial t}}_{\text{change at } \mathbf{x} \in \mathbb{R}^3} + \underbrace{(\mathbf{u} \cdot \nabla) T}_{\text{motion of fluid}} - \underbrace{(\nabla \mathbf{u}) T - T(\nabla \mathbf{u})^T}_{\text{rotation and stretching of frame}} .$$

Viscoelastic Modelling

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Introduce shear rate dependence *phenomenologically* into relaxation time and viscosity.

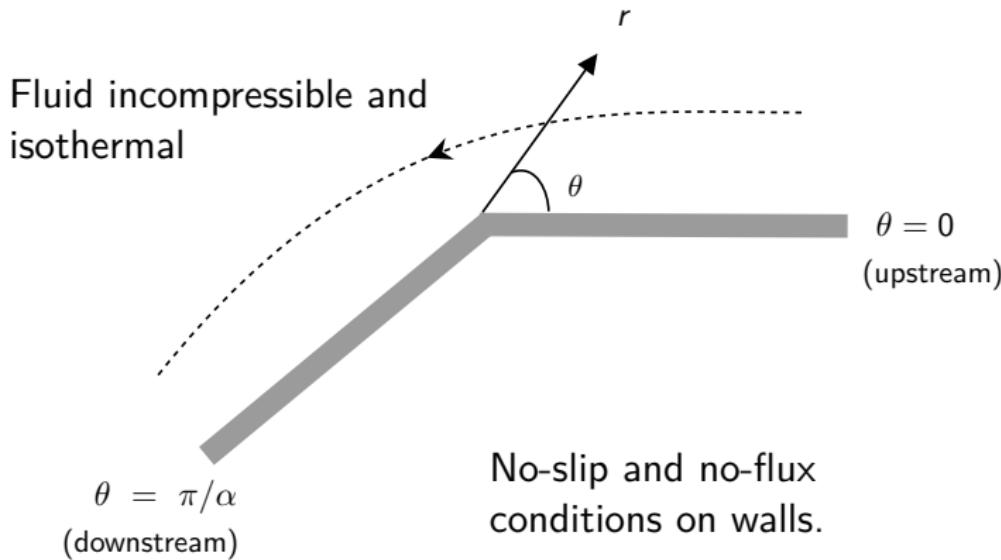
Suppose dependence is of power-law form:

$$\lambda(\dot{\gamma}) = \lambda_0 \dot{\gamma}^{q-1}, \quad \mu(\dot{\gamma}) = \mu_0 \dot{\gamma}^{n-1},$$

with $0 < \lambda_0, \mu_0, n, q < \infty$

Problem Setup

Planar, steady flow around corner of angle $\theta = \pi/\alpha$. ($1/2 \leq \alpha < 1$)



Aim: Apply asymptotic techniques to study WM fluids in region of corner apex.

Mathematics taken from Chaffin et. al. (2021), Evans and Jones (2024)

Governing Equations

Dimensionless equations:

$$\begin{aligned}\nabla \cdot \mathbf{u} &= 0, \\ \text{Re}(\mathbf{u} \cdot \nabla) \mathbf{u} &= -\nabla p + \nabla \cdot \boldsymbol{T}, \\ \boldsymbol{T} + \text{Wi} \dot{\gamma}^{q-1} \overset{\nabla}{\boldsymbol{T}} &= 2\dot{\gamma}^{n-1} \boldsymbol{D}.\end{aligned}$$

Obtain Reynolds (Re) and Weissenberg (Wi) numbers:

$$\text{Re} = \frac{[\text{inertial forces}]}{[\text{viscous forces}]}; \quad \text{Wi} = \frac{[\text{elastic forces}]}{[\text{viscous forces}]}.$$

Decompose \boldsymbol{T} into either *Cartesian* or *natural stress variables*:

$$\begin{aligned}\boldsymbol{T} &= T_{11} \mathbf{i} \mathbf{i}^T + T_{12} (\mathbf{i} \mathbf{j}^T + \mathbf{j} \mathbf{i}^T) + T_{22} \mathbf{j} \mathbf{j}^T, \\ &= -\dot{\gamma}^{n-q} \mathbf{I} + T_{uu} \mathbf{u} \mathbf{u}^T + T_{uw} (\mathbf{u} \mathbf{w}^T + \mathbf{w} \mathbf{u}^T) + T_{ww} \mathbf{w} \mathbf{w}^T.\end{aligned}$$

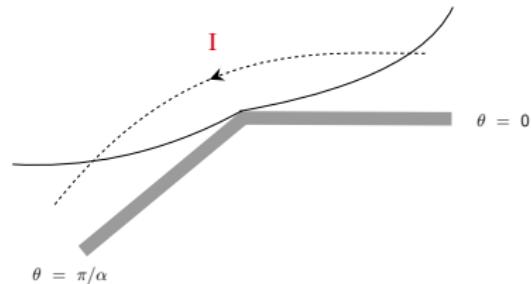
Asymptotics

Region I: Outer

Assume elasticity dominates ($\text{Wi} = \infty$):

- Leads to $\nabla \cdot \mathbf{T} = 0$ and $\nabla p = \nabla \cdot \mathbf{T}$.
- Obtain a *stretching solution*:

$$\mathbf{T} = T_{uu} \mathbf{u} \mathbf{u}^T.$$
- ‘Velocity’ $\mathbf{v} := T_{uu}^{1/2} \mathbf{u}$ governed by Euler equations



$$(\mathbf{v} \cdot \nabla) \mathbf{v} = \nabla p.$$

- Find leading order streamfunction ψ and stresses as

$$\psi \sim \frac{C_0}{\alpha^m} r^{m\alpha} \sin^m(\alpha\theta) \quad T_{uu, uw, ww} \sim d_i \left(\frac{\psi}{C_0} \right)^{mm_i}$$

with m, m_i, C_0, d_i to be determined.

Asymptotics

Viscometric Flow

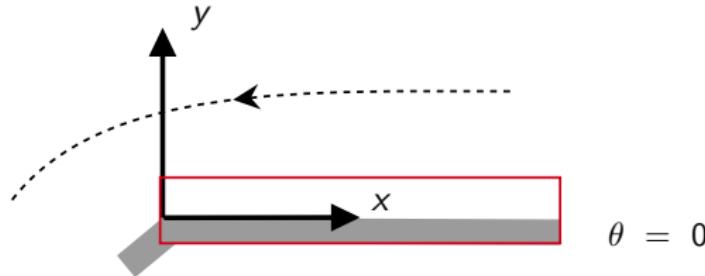
- Flow is simple shear to leading order near the walls

No-slip/no-flux $\implies \psi = O(y^2)$ as $y \rightarrow 0$.

- Cannot match outer stresses to wall behaviour:

Outer: $T_{uw} = O(y^{mm_2})$, At wall: $T_{uw} = O(1)$.

- Matching fails because we neglected viscosity — need *boundary layers* to recapture these effects.



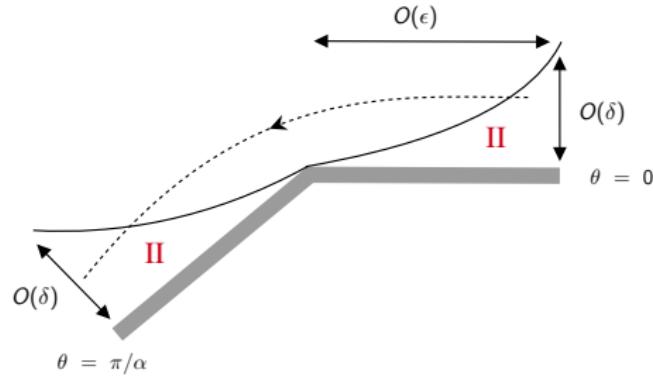
Asymptotics

Region II: Boundary Layers

- Focus on upstream layer ($x, y > 0$).
- Rescale to look close to the corner: $x = \epsilon X$, $y = \delta Y$, with $0 < \delta \ll \epsilon \ll 1$.
- Dominant balance in WM equations *determines m!* Also gives

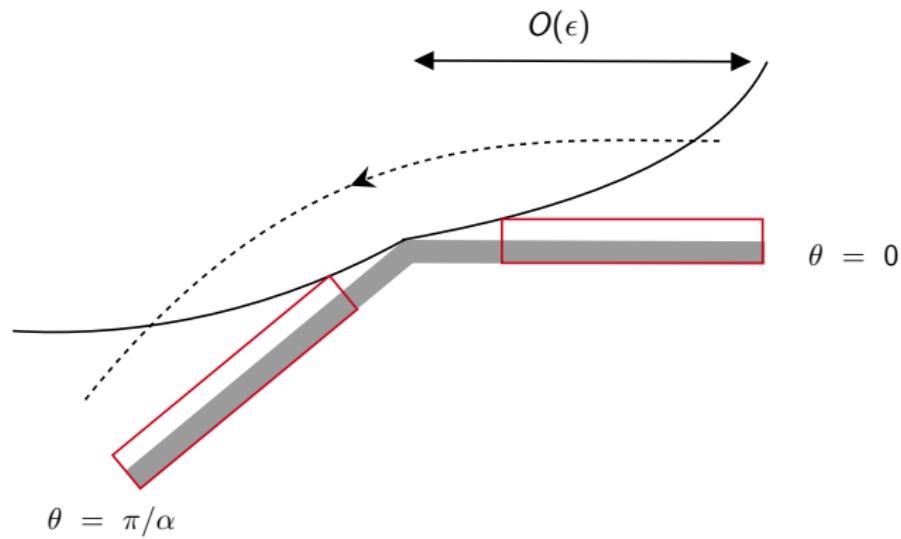
$$\delta = \epsilon^{2-\alpha + \frac{(q-n)(1-\alpha)}{n+q}}.$$

- Same results apply to the downstream layer by reorientation of axes!



Asymptotics

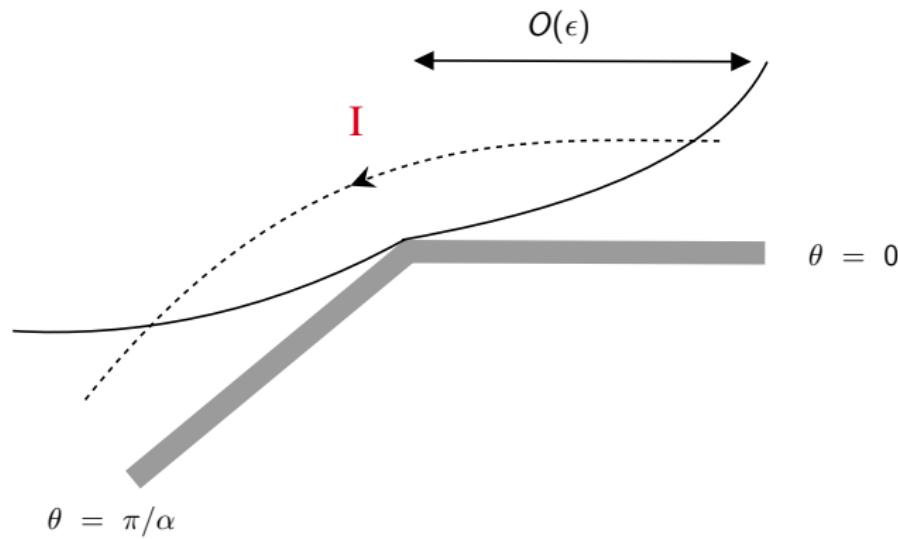
Solution Structure



1. Flow near walls is (to leading order) simple shear.

Asymptotics

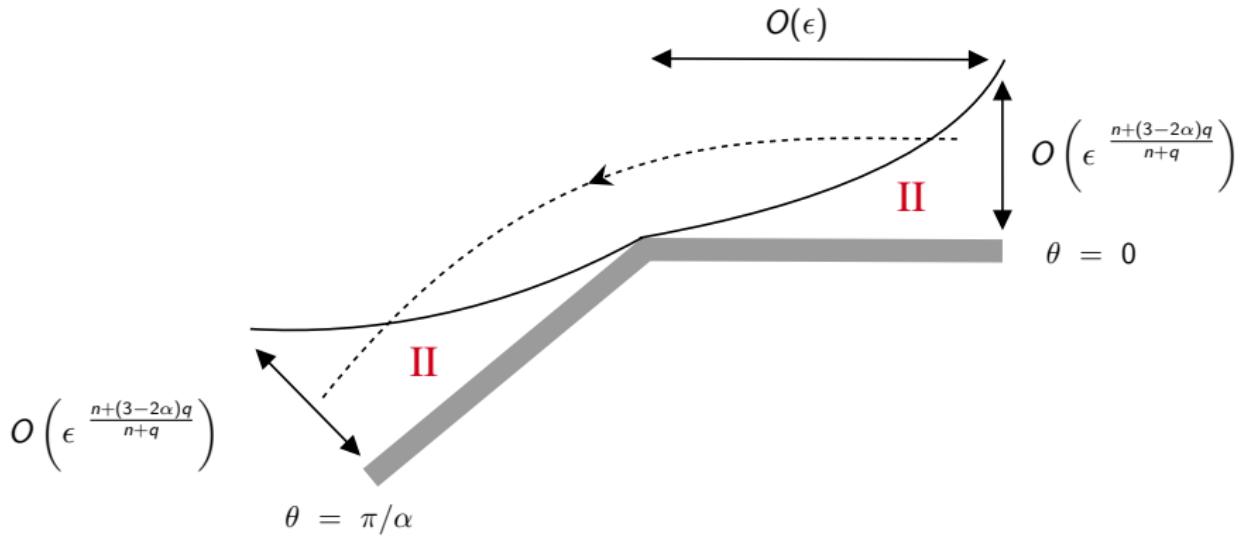
Solution Structure



2. Fluid behaves elastically as it flows round the corner (Region I).

Asymptotics

Solution Structure



- 3. Viscous effects remain important in Regions II. Reduced PDEs determine velocities/stresses.

Similarity Solutions

- BL scalings suggest layer similarity solutions in the variable

$$\chi = X^{-a} Y, \quad a = \frac{n + (3 - 2\alpha)q}{n + q}.$$

- System of PDEs in $(\psi, T_{uu}, T_{uw}, T_{ww})$ reduces to four ODEs in $(f, t_{uu}, t_{uw}, t_{ww})$!
- Boundary + Matching conditions:

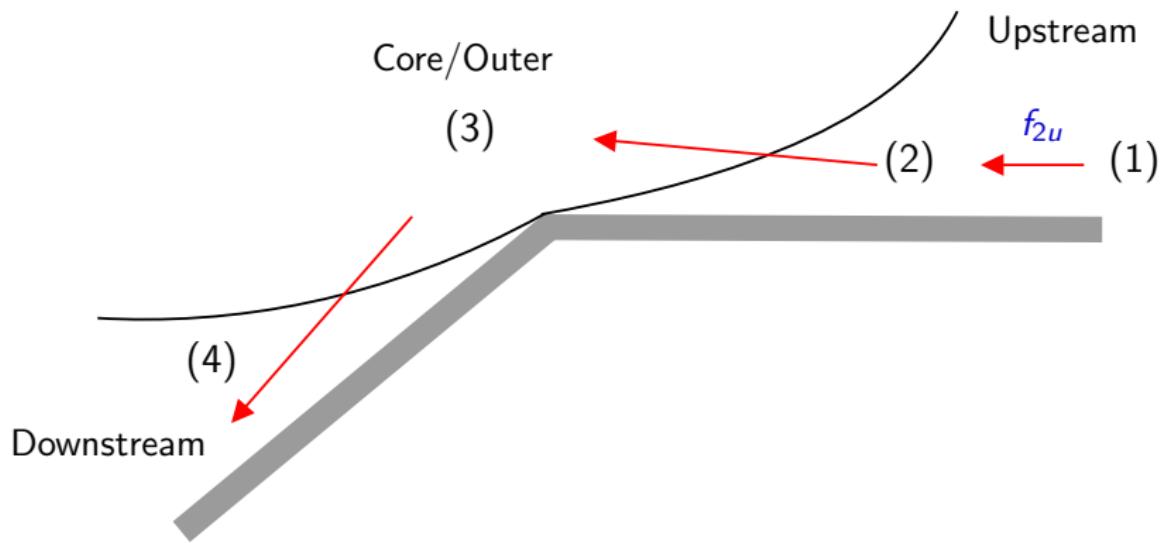
at $\chi = 0$: $f = f' = 0,$

as $\chi \rightarrow \infty$: $f \sim C_0 \chi^m, \quad t_{uu} \sim d_1 \chi^{mm_1},$
 $t_{uw} \sim d_2 \chi^{mm_2}, \quad t_{ww} \sim d_3 \chi^{mm_3}.$

- Note far-field contains four unknown constants! (m, m_i known)

Solving the Equations

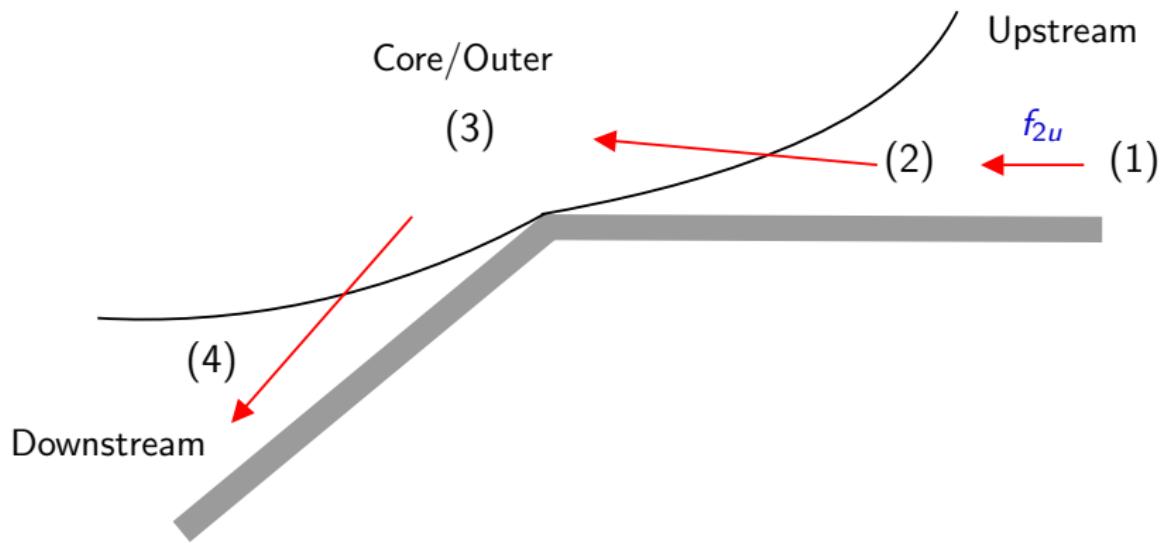
Numerical Procedure



1. Specify flow parameters (wall shear rate f_{2u} and n, q, Wi, α, p_{0u}).

Solving the Equations

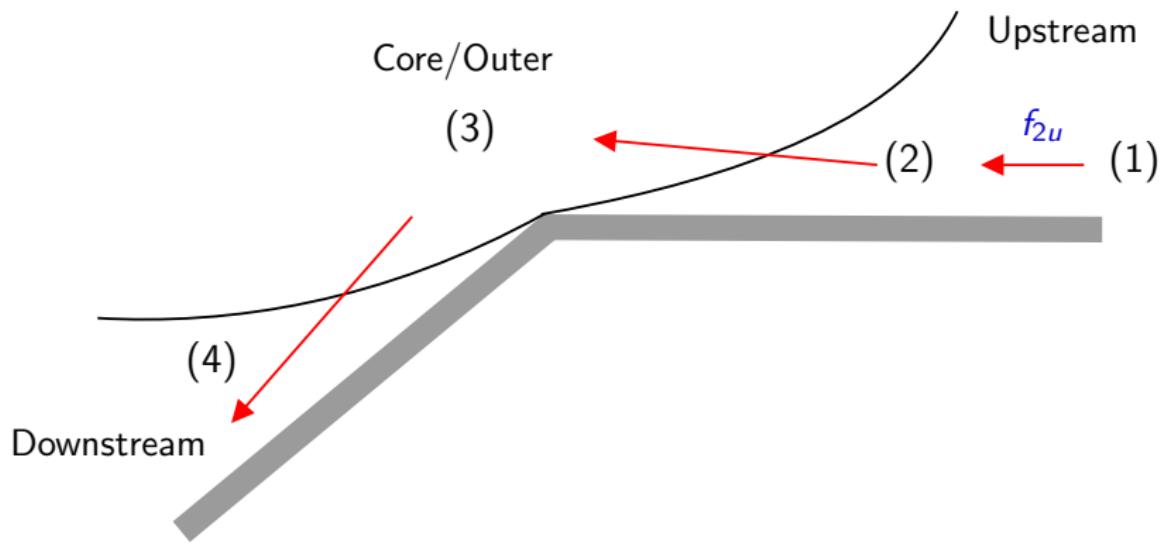
Numerical Procedure



2. Solve similarity equations in upstream region as an IVP (f_{2u} determines initial conditions).

Solving the Equations

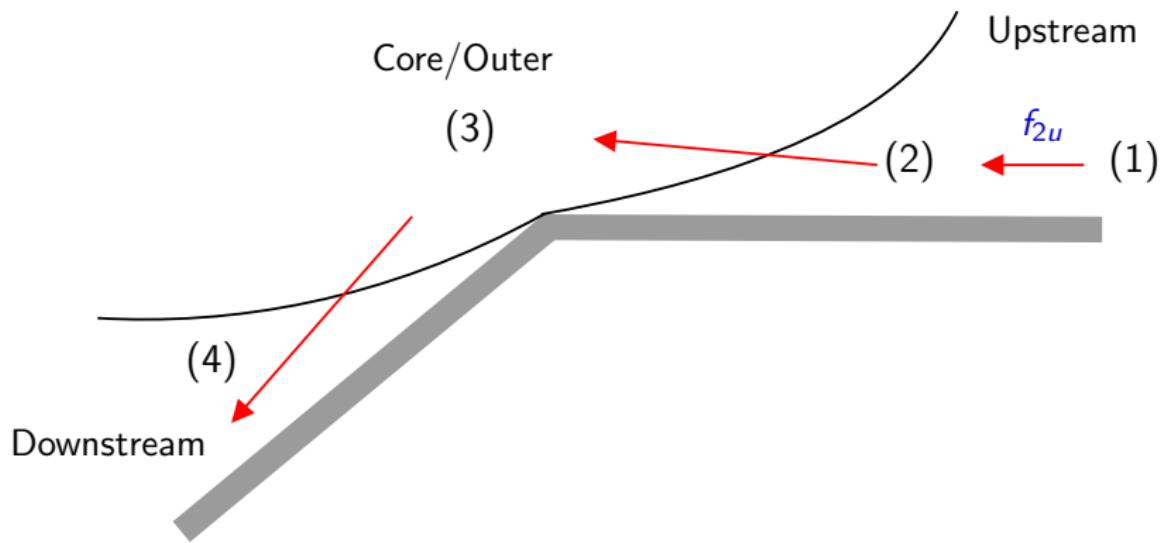
Numerical Procedure



3. Scale IVP solution by far-field behaviour to find C_0, d_1, d_2, d_3 (determines outer solution).

Solving the Equations

Numerical Procedure

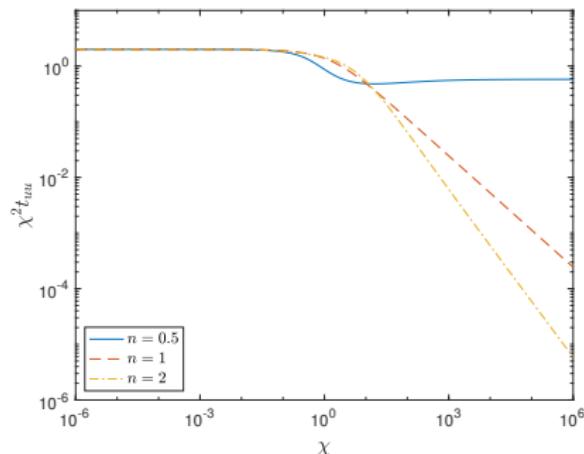
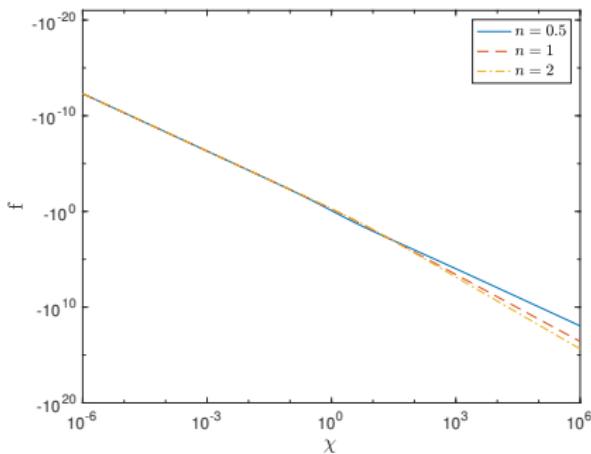


4. Solve in downstream region via a BVP (specify outer solution and wall shear as conditions)

Solving the Equations

Upstream Numerics

Streamfunction (f) and stress (t_{uu}) profiles.



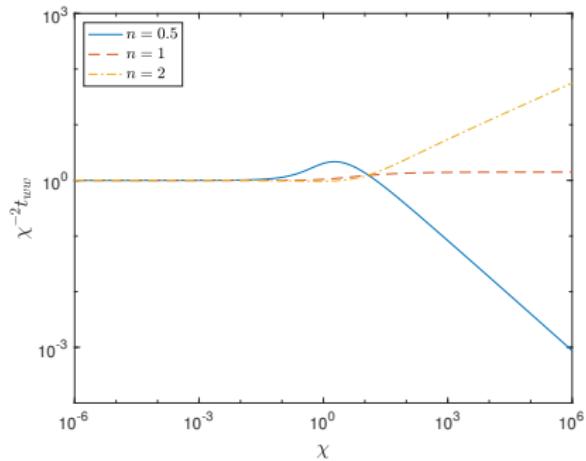
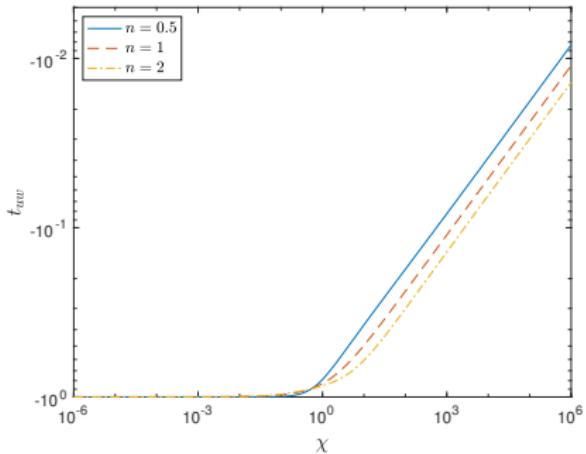
Parameters: $n = q$, $\text{Wi} = 1$, $p_{0u} = 1$, $f_{2u} = -1$, $\alpha = 2/3$.

Solved using `ode15s` on domain $[10^{-6}, 10^6]$.

Solving the Equations

Upstream Numerics

Stress component (t_{uw} , t_{ww}) profiles.

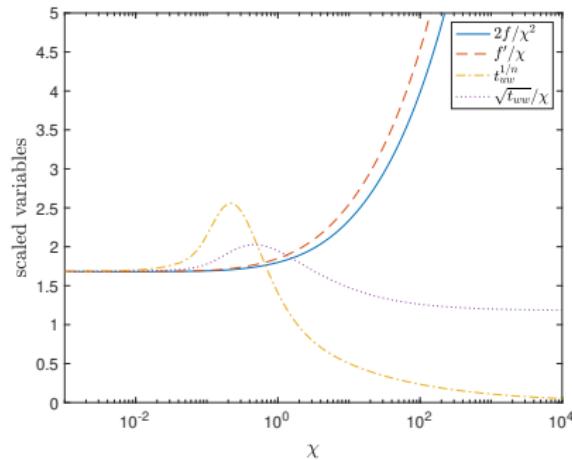
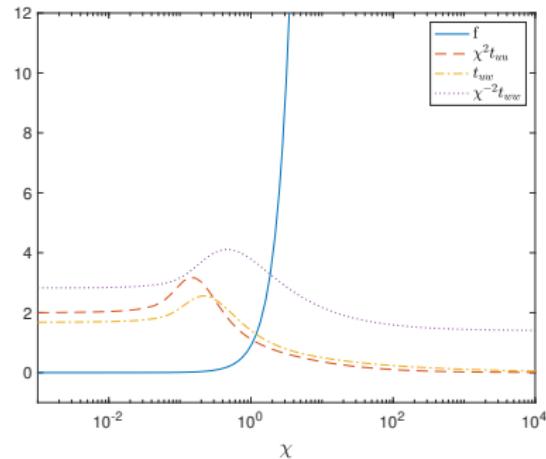


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Solving the Equations

Downstream Numerics

Downstream variables (f , t_{uu} , t_{uw} , t_{ww}) and estimation of wall shear rate (right).

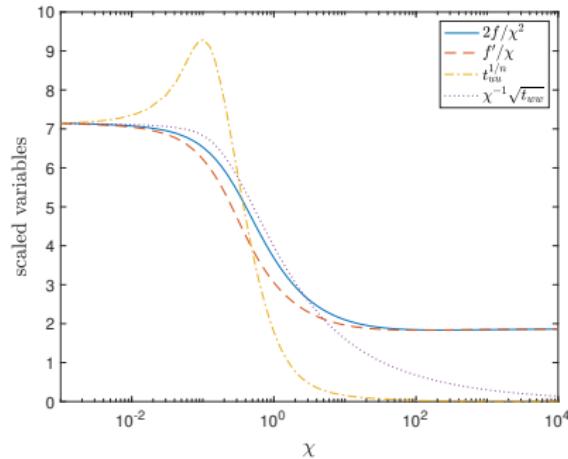
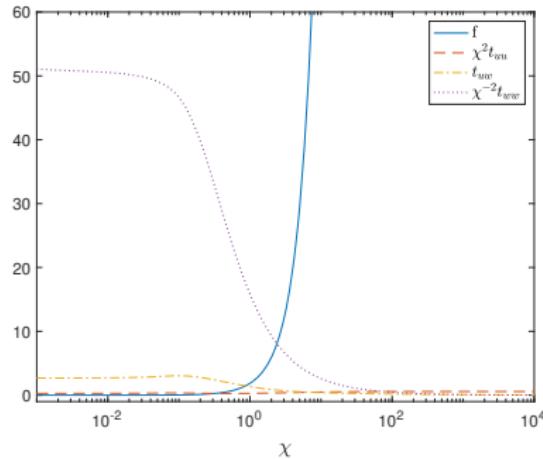


Parameters: $n = q = 1$, $p_{0u} = 1$, $f_{2u} = -1$, $\alpha = 2/3$. Solved using `bvp4c` on domain $[10^{-3}, 10^4]$.

Solving the Equations

Downstream Numerics

Downstream variables (f , t_{uu} , t_{uw} , t_{ww}) and estimation of wall shear rate (right).



Parameters: $n = q = \frac{1}{2}$, $p_{0u} = 1$, $f_{2u} = -1$, $\alpha = 2/3$. Solved using bvp4c on domain $[10^{-3}, 10^4]$.

Summary

For the WM model, we have

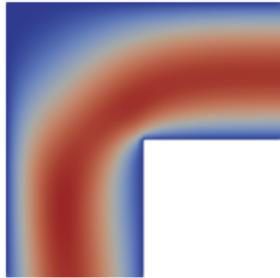
- (Re)formulated the equations in natural stress variables.
- Shown the existence of a three-region solution structure for re-entrant corner flow in natural stresses.
- Completed the flow description near to the apex via solvable similarity solutions in the upstream and downstream regions.

Future Work

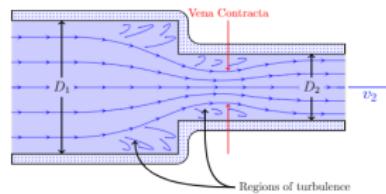
Full Numerical Verification

- Perform *volume-of-fluid* simulations for WM in OpenFOAM (+rheoTool).
- Simulate WM in re-entrant corner geometries:

L-shaped domain

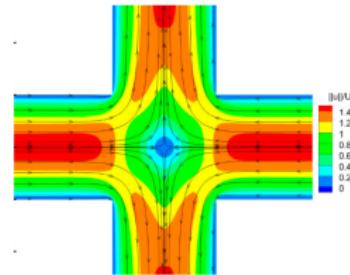


Contraction flow



Dusling, K. (n.d.)

Cross slot



Yuan et. al. (2020)

Thank you for listening!



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