## Cálculo

# Al futuro, cuando al fin usaremos esto con fluidez.

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## Índice de cuadros

# Índice de figuras

### Prefacio

Algo.

#### Agradecimientos

A las personas que hicieron sencillo el publicar libros.

ASG

#### 0

## Quién escribe?

Por ahora, solo ASG. Más tarde espero que más personas.

### 1

### Introducción

Si esto funciona, el sitio se actualizará automáticamente. Y la escritura se simplificará. Además, cada capítulo tendrá el enlace a su propio video sin problemas.

#### Ejemplo 01

Esta es la primera prueba con LaTeX. En esta sesión veremos una consecuencia importante del Teorema de Green: si \$F\$ es un campo de clase \$C^1\$, el rotacional \$  $\dim(A) = \sup \left\{ \left( \right) - \left( \right) \right\} \right\}$ \] \begin{prop} \label{prop1} Sean  $U\subset\R^2\$  una regi\'on,  $F = \left(P, Q\right): U\to \R^2\$  de clase  $C^1\$  en  $U\$ ,  $\$  voerline  $X\$  in  $U\$  y  $C^2\$  de clase  $C^1\$  en  $U\$  $\rot\eft(F(\overline{x})\right) = \lim\limits_{\epsilon\to 0^{+}} \frac{\displaystyle\int\limits_{\Gamma_{\overline{x}}}}$ donde \$\gamma\_{\epsilon}\$ es una parametrizaci\'on de \$\Gamma\_{\epsilon}\$ que la recorre en el sentido contrario al de \end{prop} \begin{proof} Notamos que para cada \$0< \epsilon<c\$ se cumplen las hip\'otesis del Teorema de Green, por lo cual \displaystyle \int\limits\_{\Gamma\_{\epsilon}} F\cdot d\gamma\_{\epsilon} = \int\limits\_{\Omega\_{\epsilon}} \left( Luego, para cada \$\epsilon\$, por el Teorema del Valor Promedio para integrales sobre conjuntos Jordan--medibles (\text \begin{align\*}  $\int_{\infty} \left(\frac{y}{\pi x} - \frac{y}{\pi x} \right) dx - \frac{y}{\pi x} - \frac{y}{\pi x} dx -$ & = \left( \frac{\partial Q}{\partial x}\left(\overline{\xi}\_{\epsilon}\right) - \frac{\partial P}{\partial y}\left \end{align\*}  $Ya \ que \ \overline \ \{x\}\in \climate{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\omega_{\climate{left(\climate{left(\omega_{\climate{left(\omega_{\climate{left(\climate{left(\climate{left(\climate{left(\$ 1 \frac{\displaystyle\int\limits\_{\Gamma\_{\epsilon}} F\cdot d\gamma\_{\epsilon}}{\area\left(\Omega\_{\epsilon}\right)} \] Para concluir, notamos que las funciones  $\frac{partial Q}{partial x}$  y  $\frac{partial P}{partial y}$  son continuas 1/  $\lim \int_{\left(x\right)^{+}} \operatorname{in}\left(x\right) = \operatorname{in}\left(x\right)$ \]  $porque $\langle x\rangle \in (0.000) \ porque $\langle x\rangle \in (0.000) \ para toda $0<\exp\sin(x) \ c$ y $\langle x\rangle \in (0.000) \ para toda $0<\exp\sin(x) \ para toda $0<\exp(x) \ para$ \begin{align\*}  $&= \frac{partial Q}{partial x}\left( | vartial x \right) - \frac{partial P}{partial y}\left( | vartial x \right) - \frac{partial y}\left( | vartial x \right) - \frac{partial y}{partial y}\left( | vartial x \right) - \frac{partial x}{partial y}\left( | vartial x \right) - \frac{partial x}{partial x}\left( | va$ & = \rot\left( F(\overline{x})\right). \end{align\*} Esto termina la prueba. \end{proof}

Aunque no lo parezca, la proposici\'on anterior se puede interpretar como la \emph{rotaci\'on promedio} generada por el ca

4 2 Ejemplo 01

Para concluir esta sesi\'on presentamos una aplicaci\'on directa del Teorema de Green: el c\'alculo de una integral.

```
\begin{ejer}
       Calcule la integral \star \sin x^3 - y^3, x^3 + 2y^3 = 1
\end{ejer}
\begin{proof}[Solución.]
       Tenemos que F es de clase C^1 en R^2 y si 0 en B_1(0,0), entonces 0 es un conjunto Jordan--medible y
                \displaystyle\int\limits_{\Gamma=\partial \Omega} F \cdot d\gamma = \int\limits_{\Omega} \rot\left( F \right),
           \]
       donde $\gamma$ es la parametrizaci\'on usual de $\Gamma$ que la recorre una vez en el sentido contrario al de las maneci
           Notamos que
           \ [
                        \frac{partial Q}{partial x}(x,y) = 3x^2
           \]
           y tambi\'en
           \ [
                         \frac{partial P}{partial y}(x,y) = -3y^2,
           \]
           por lo cual
           1
                        \operatorname{\mathsf{Tot}}(F)(x,y) = 3x^2 + 3y^2 = 3(x^2 + y^2).
           \]
           En virtud de lo anterior, tenemos que
                   \]
           donde f(x,y) = x^2+y^2.
           Notamos que al aplicar el cambio de variable a coordenadas polares obtenemos que
           \begin{align*}
               \int_{B_1(0,0)} f &= \int_{0}^{2\pi}\left(\int_{0}^{2\pi}\left(\int_{0}^{2\pi}\right) dx\right) dx = \int_{0}^{2\pi}\left(\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{2\pi}\int_{0}^{
                        & = \int_{0}^{2\pi} \frac{1}{4} d\theta \
                        & = \frac{\pi}{2}.
           \end{align*}
           Por lo tanto,
                        \label{limits_{Gamma=partial 0mega} F \ d d gamma = \frac{3\pi}{2}.
           \]
\end{proof}
\textcolor{blue}{?`Obtiene el mismo resultado si hace los c\'alculos directamente?}
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#### A

## Por aprender

Este sería el final. Ver la Proposición ??