

Review of G-S in residual form ( $A = L + D + U$ ) ( $\tilde{A} = L+D$ )

$$(L+D) \cdot \psi^{n+1} = \omega - U \cdot \psi^n \text{ let } \psi^{n+1} = \psi^n + \Delta\psi \text{ then}$$

$$(L+D) \cdot \Delta\psi = \omega - (L + D + U) \cdot \psi^n = \omega - A \cdot \psi^n \equiv R^n ; \quad D \cdot \Delta\psi = R^n - L \cdot \Delta\psi$$

If  $\psi^n$  is a solution, then  $R=0$ . Iteration = drive residual to zero.

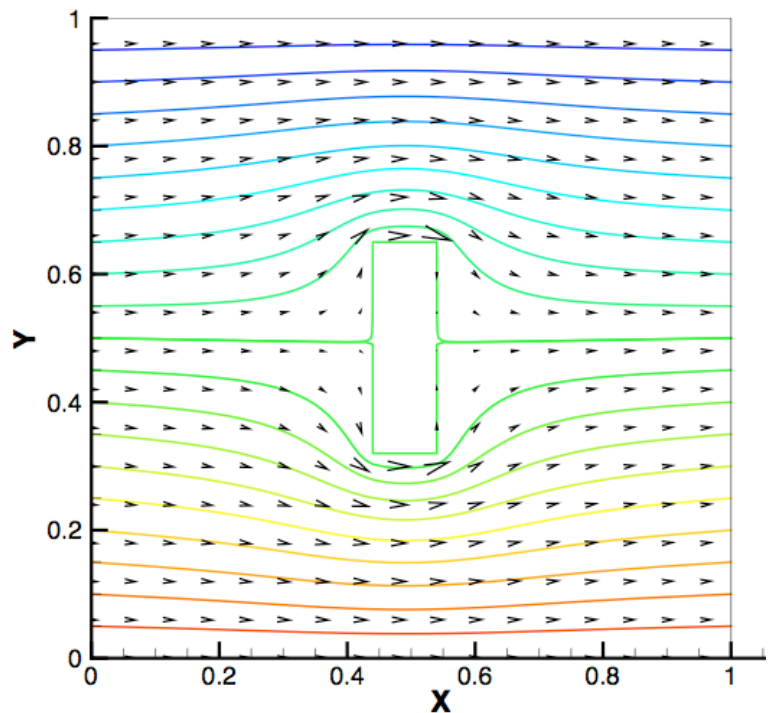
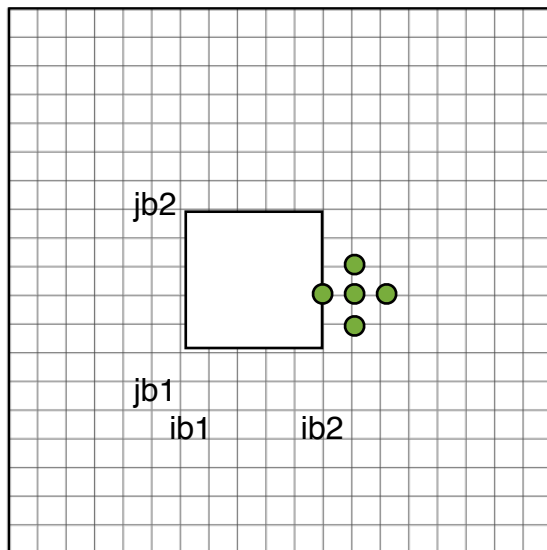
In G-S, accelerated by SOR, this is solved as

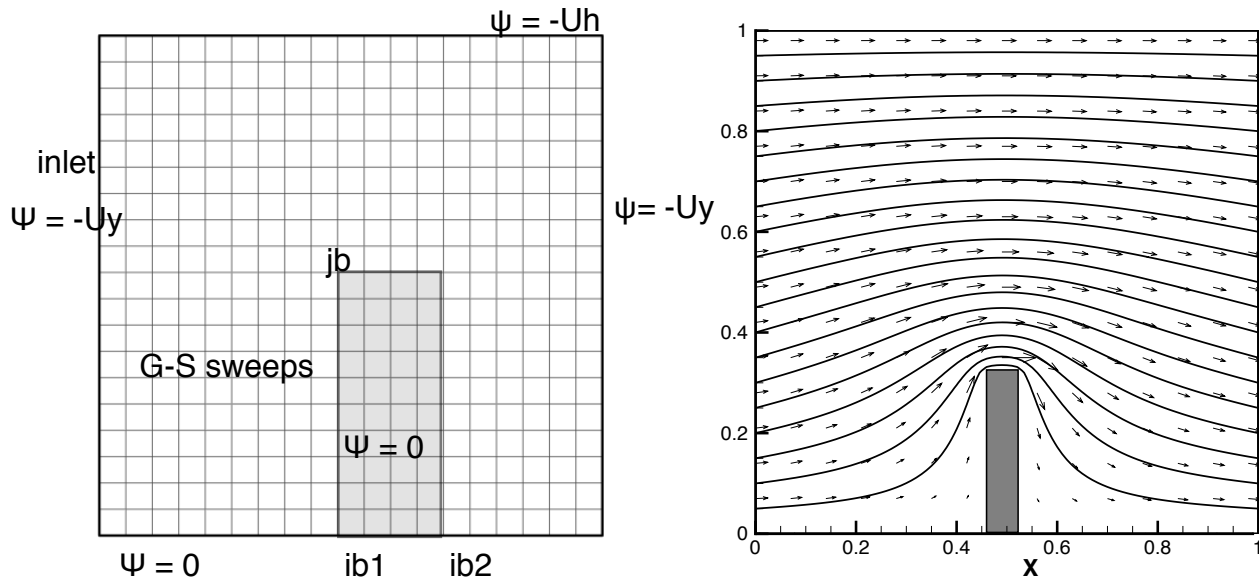
$$D \cdot \Delta\psi \equiv R^n - \lambda L \cdot \Delta\psi \text{ and } \psi^{n+1} = \psi^n + \lambda \Delta\psi$$

show Grids.pptx

### Toward geometry: i-blanking

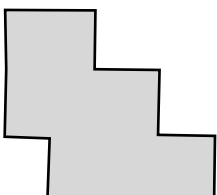
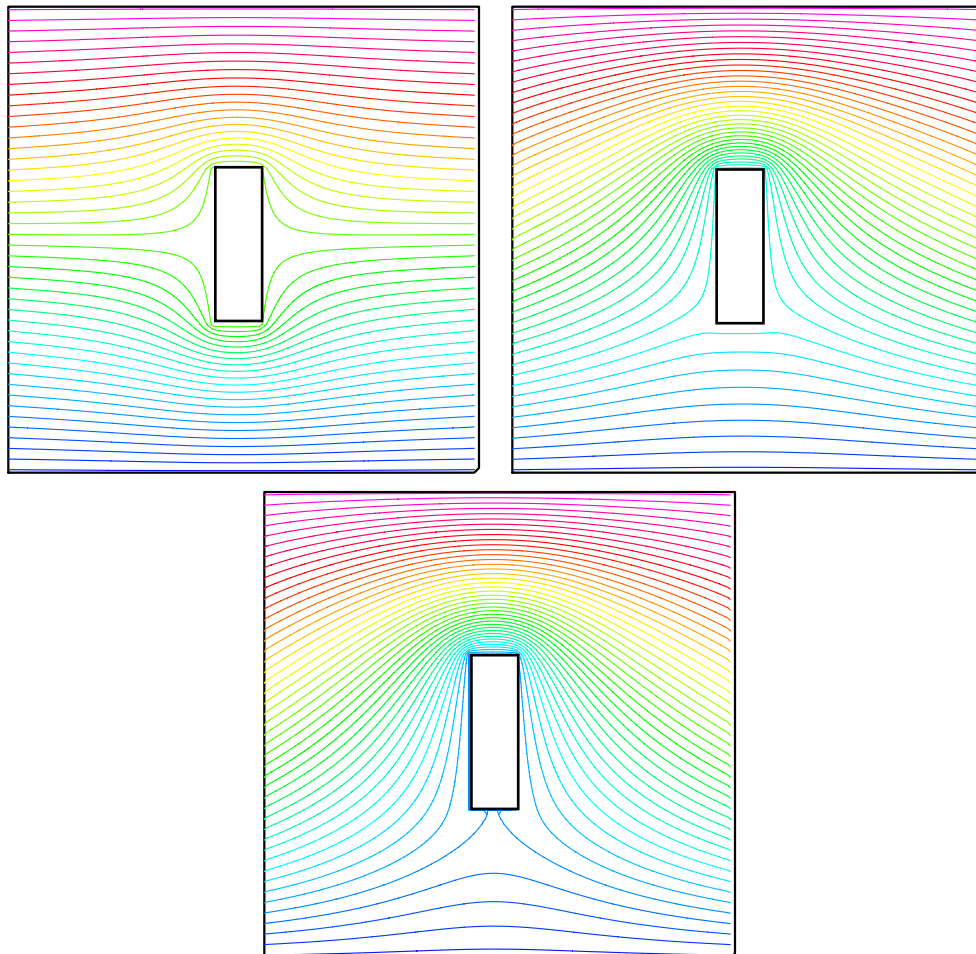
- Consider a rectangular geometry and Cartesian grid -- nodes at intersections
- Inside rectangle solve  $\Psi_{ij} = \text{const}$ , so row of matrix is 1 on diagonal and b is 0. NB:  $\Psi$  can be any constant on rectangle -- e.g. circulation; or temperature
- Linear algebra:  $A \cdot \Psi = \omega$ . In rib  $A = I$  so equation is  $\Psi = \omega = \Psi_B$ .  
Or  $A \cdot \Delta\Psi = R^n = 0$ ; actually, G-S is  $D \cdot \Delta\Psi = R^n - L \cdot \Delta\Psi$  and  $\Delta\Psi=0$  inside rectangle.





Non-dimensionalize: uniform flow becomes  $\psi = -y$ ; lower wall,  $y=0$ ,  $\psi=0$ ; upper wall  $y=1$ ,  $\psi=-1$ .

$\psi_{\text{rib}} = -0.5, -0.2 \text{ and } -0.1$



D. Pseudo-code for blanking; 2 options: change **A** or make use of **IB = 1** in fluid, 0 in geometry

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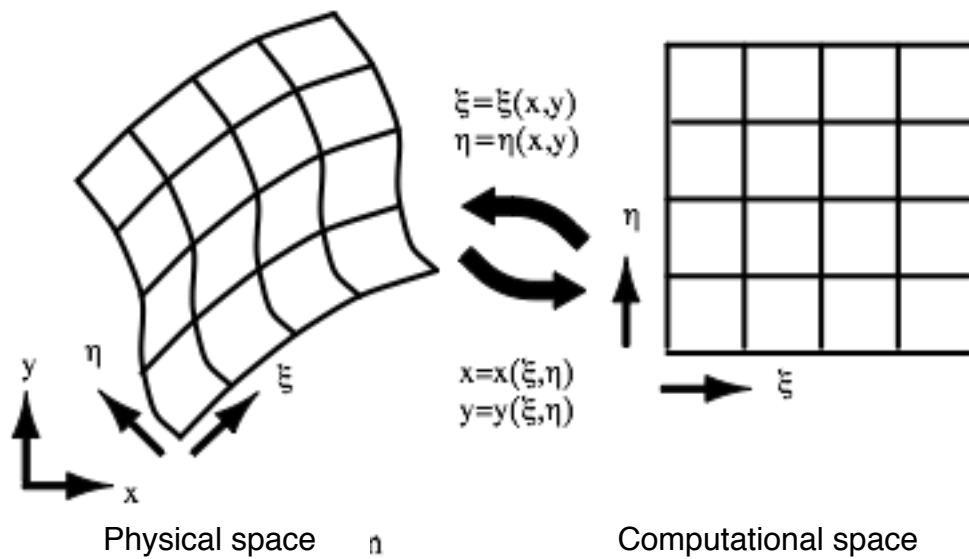
IB(:, :) = 1
 $\Psi(:, :) = 0$     ! initial guess
dy = 1/(J-1) ; dx = 1/(I-1)
DO j=1,J          ! B.C.  $\Psi = -y$ 
   $\Psi(1, j) = -(j-1)*dy$ 
   $\Psi(I, j) = -(j-1)*dy$ 
ENDDO
 $\Psi(:, J) = -1$     ! Upper boundary
 $\Psi(:, 1) = 0$      ! Lower boundary

A(:, :, 1) = A(:, :, 5) = 1./dy**2      ! 5 point stencil
A(:, :, 2) = A(:, :, 4) = 1./dx**2
A(:, :, 3) = -(A(:, :, 1)+A(:, :, 2)+A(:, :, 4)+A(:, :, 5))
!----
! inside & surface of rib: the ONLY CHANGE to Poisson2D_GS.f90
!----
DO j=jb1, jb2
  DO i=ib1, ib2
    A(i, j, [1, 2, 4, 5]) = 0
    A(i, j, 3) = 1
    IB(i, j) = 0      ! zero out in geometry
     $\omega(i, j, 3) = 0$ 
     $\Psi(i, j) = \Psi_{rib}$  (0 if on lower wall)
  ENDDO
ENDDO
! Gauss-Seidel with SOR (sweep LL-UR)
L∞ = 0; L2 = 0
DO j=2, J-1
  DO i=2, I-1
    R =  $\omega(i, j)$ 
    -( A(i, j, 1)* $\Psi(i, j-1)$ +A(i, j, 2)* $\Psi(j-1, k)$ +A(i, j, 3)* $\Psi(i, j)$  &
      +A(i, j, 4)* $\Psi(j+1, k)$ + A(i, j, 5)* $\Psi(i, j+1)$  )
     $\Delta\Psi = R/A(i, j, \mathbf{3})$ 
     $\Psi(i, j) = \Psi(i, j) + \mathbf{IB}(i, j)*\lambda*\Delta\Psi$ 

    L∞ = amax(Linf, abs[Del(i, j)])
    L2 = L2 + IB(i, j)*Del(i, j)^2
  ENDDO
ENDDO

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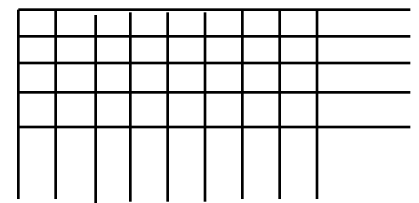
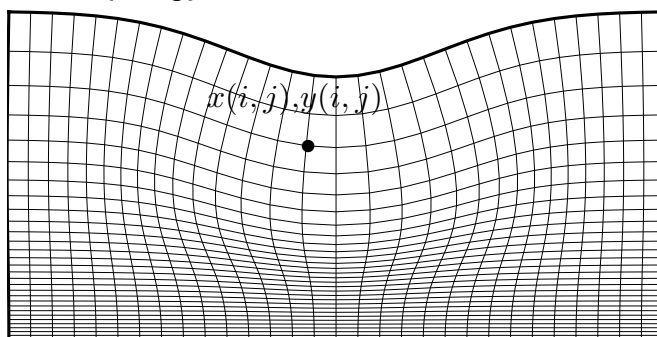
# Introduction to structured grids: Computational and physical space

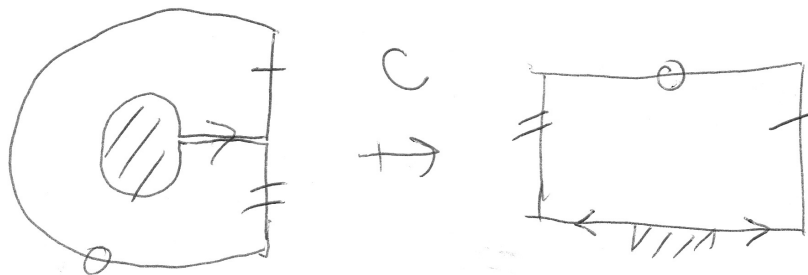
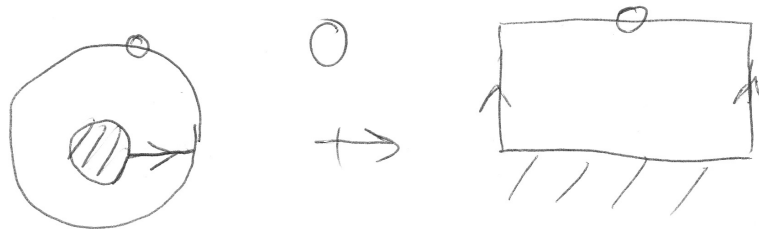
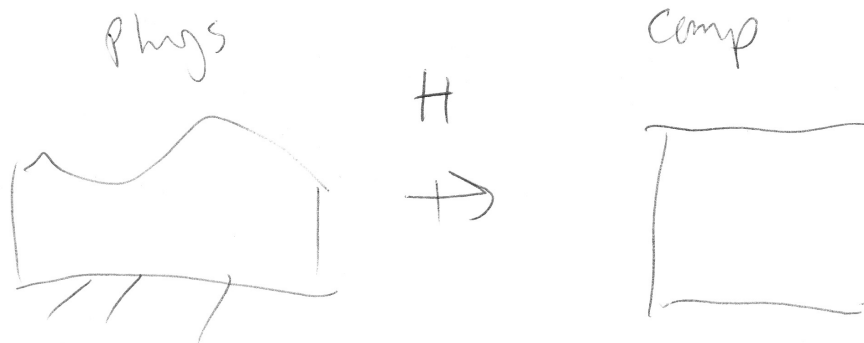


$\xi \longleftrightarrow j$  and  $\eta \longleftrightarrow k$  Note  $\Delta \eta = 1$ .  $x(i, j)$ ,  $y(i, j)$  is mapping from computational to physical space --- but just arrays of coordinates.

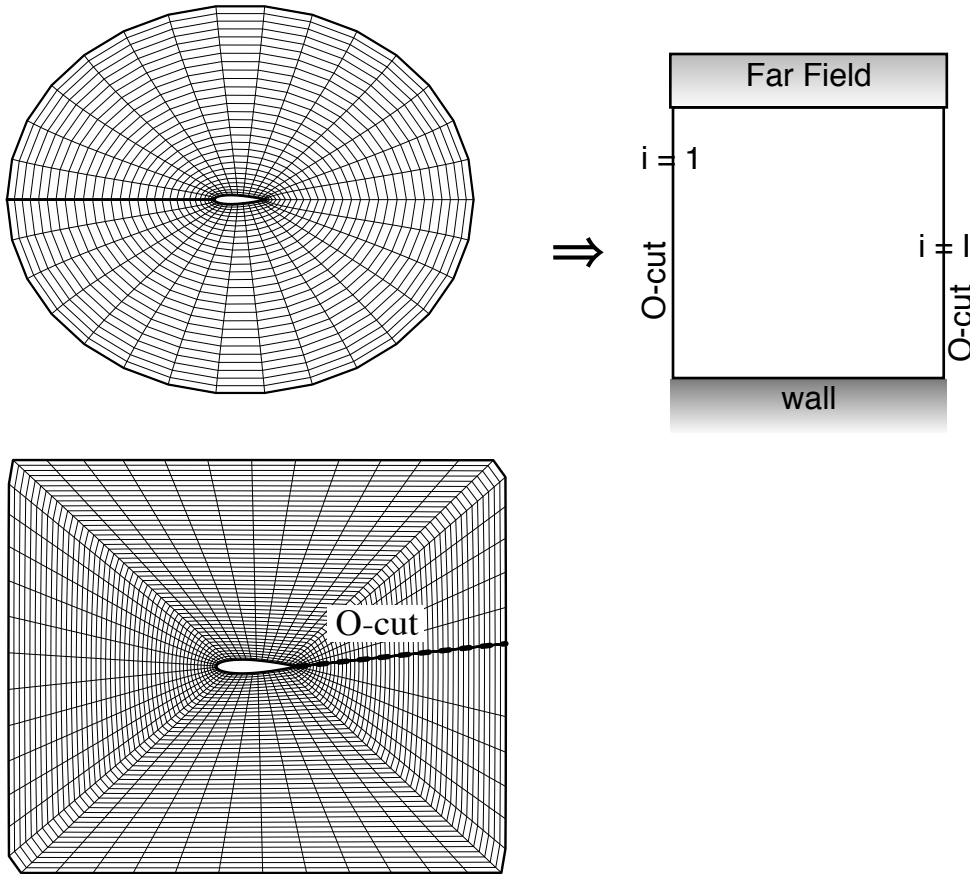
## Grid 'topologies'

A. H-topology: Lines from entrance to exit and top to bottom. Maps to rectangular grid

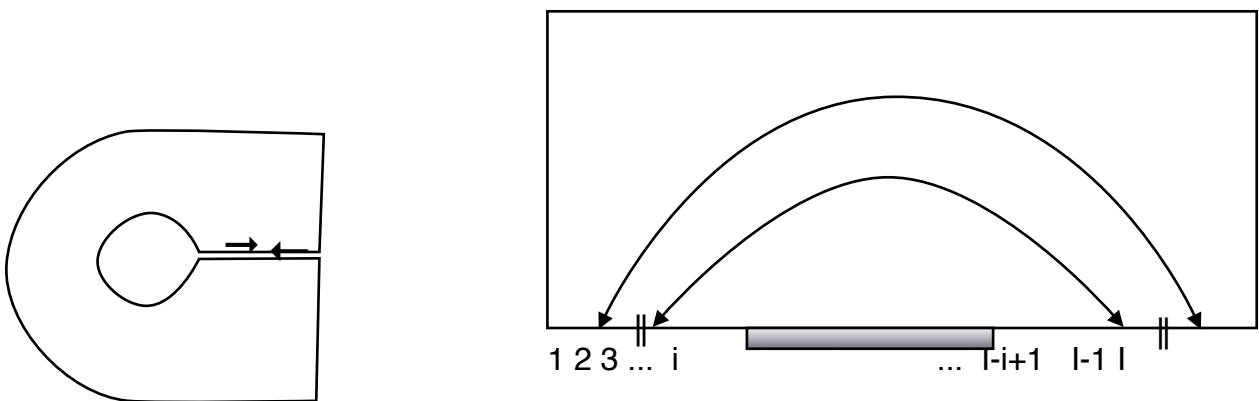


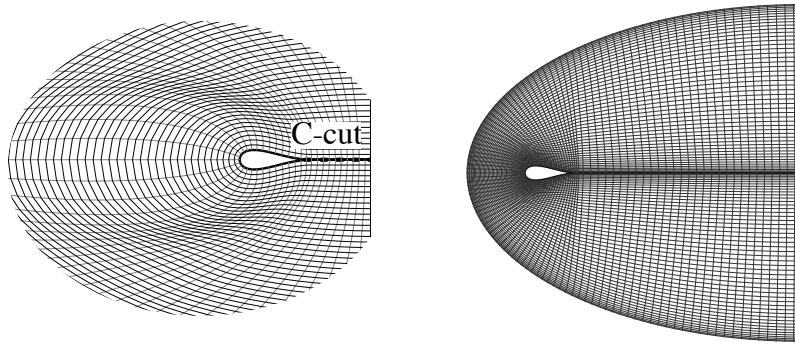


B. O-mesh topology Singly connected:  $i=1 \equiv i=I$



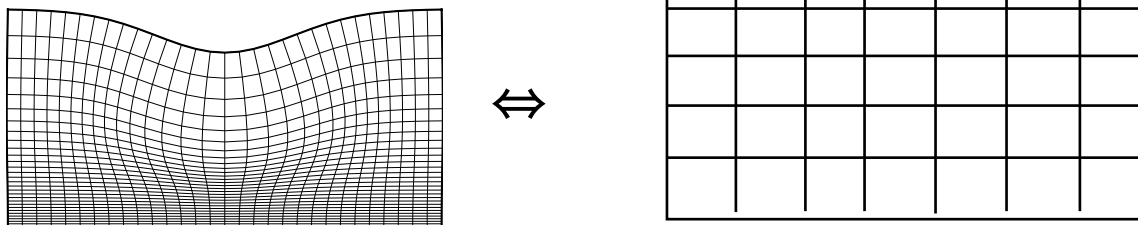
C. C-mesh





### Implications for discretization

A. H-topology: Lines from entrance to exit and top to bottom. Maps to rectangular grid

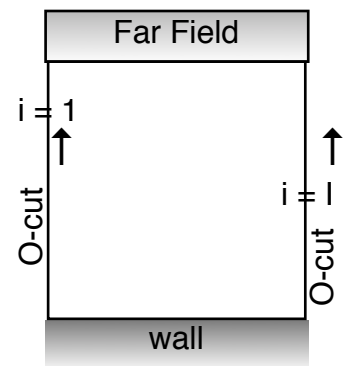
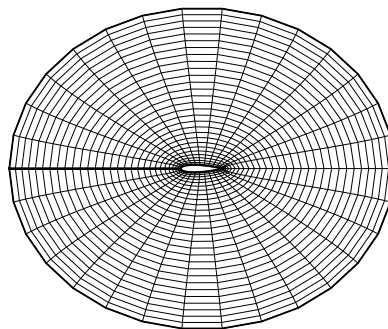
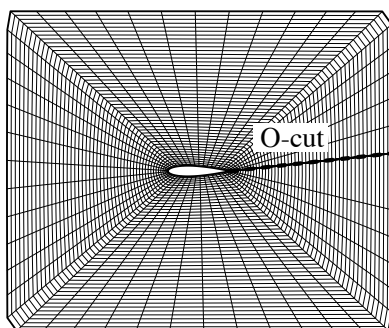


B. O-mesh topology Singly connected:  $i=1 \equiv i=I$ .

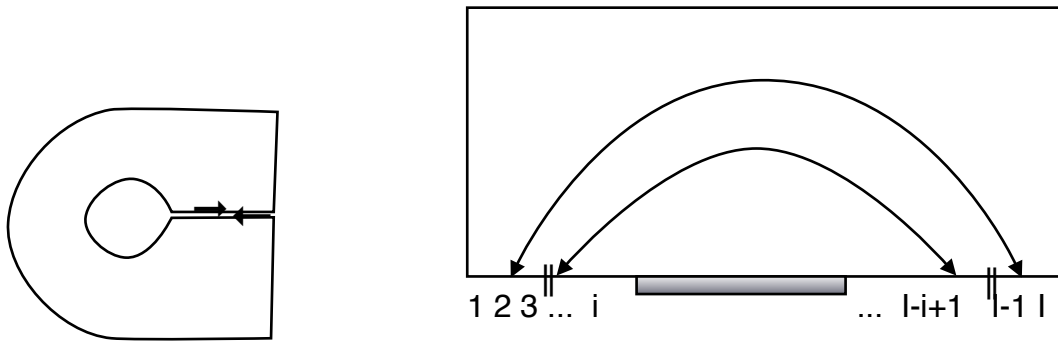
Periodic in  $i$ : Identify:  $\psi(1,j) = \psi(I,j)$

For stencil:  $\psi(0,j) = \psi(I-1,j)$

$\delta \psi / \delta x|_1 = \psi(2,j) - \psi(0,j) = \psi(2,j) - \psi(I-1,j)$



## C. C-mesh



For C-cut on  $j=1$ ,  $i=1 \leftrightarrow i=l$   $i=2 \leftrightarrow i=l-1$   
 $i \leftrightarrow l-i+1$   $\delta \psi / \delta x_{li} = \psi(i,2) - \psi(i,0) = \psi(i,2) - \psi(l-i+1,2)$

