## Caleb Logemann AER E 546 Fluid Mechanics and Heat Transfer I Homework 2

1. The heat fin equation is the linear o.d.e.

$$\frac{\mathrm{d}^2 T}{\mathrm{d}x^2} = MT$$

where M is a sort of thermal mass. First write the finite difference equation in terms of a tridiagonal matrix. Solve that equation using the Thomas algorithm (Gaussian elimination) for:

(a) Compute a solution with the boundary conditions T(0) = 1 and T(1) = 0. This corresponds to a fin that is between a hot and a cold reservoir. In non-dimensional terms, the heat flux into the cold reservoir is  $-\frac{\mathrm{d}T}{\mathrm{d}x}$  at x = 1. Obtain the heat flux as x = 1 for M = 1, 5, 9. Use enough grid points to obtain 1% accuracy. Provide your three numerical values of the heat flux. Provide a single graph with curves of T(x) for the 3 values of M.

First I will establish a some notation. I will discretize the fin into N+1 points. Let  $x_i = \frac{i}{N}$ , then  $x_0 = 0$  and  $x_N = 1$ . Let the approximate solution at  $x_i$  be represented by  $T_i$ . Then a numerical solution consists of a set of values  $T_i$  for  $i \in \mathbb{N}$ ,  $0 \le i \le N$ .

Next I will discretize the partial differential equation into a discrete equation. The second order central finite difference for the second derivative is

$$\frac{T_{i-1} - 2T_i + T_{i+1}}{\Delta x^2}$$

Plugging this into the partial differential equation gives the following difference equation

$$\frac{T_{i-1} - 2T_i + T_{i+1}}{\Delta x^2} = MT_i$$

Simplifying this gives

$$T_{i-1} + \left(-2 - M\Delta x^2\right)T_i + T_{i+1} = 0$$

For this problem the boundary conditions are T(0) = 1 and T(1) = 0. This can be encoded into the numerical solution at  $T_0 = 1$  and  $T_N = 0$ . Now only the values for  $T_i$  for  $1 \le i \le N-1$  need to be found. These can be found by solving the equations

$$T_{i-1} + \left(-2 - M\Delta x^2\right)T_i + T_{i+1} = 0$$

for i = 1, this become

$$\left(-2 - M\Delta x^2\right)T_1 + T_2 = -1$$

and for i = N - 1 the equation is

$$T_{N-2} + \left(-2 - M\Delta x^2\right)T_{N-1} = 0$$

These equation can be written in matrix form as

$$\begin{bmatrix} -2 - \Delta x^2 & 1 & & & 0 \\ 1 & -2 - \Delta x^2 & 1 & & & \\ & \ddots & \ddots & \ddots & & \\ & & 1 & -2 - \Delta x^2 & 1 \\ 0 & & & 1 & -2 - \Delta x^2 \end{bmatrix} \begin{bmatrix} T_1 \\ T_2 \\ \vdots \\ T_{N-2} \\ T_{N-1} \end{bmatrix} = \begin{bmatrix} -1 \\ 0 \\ \vdots \\ 0 \\ 0 \end{bmatrix}$$

This is a tridiagonal system which can be solved easily with the Thomas algorithm. The following function evaluates the Thomas algorithm on a tridiagonal system.

1

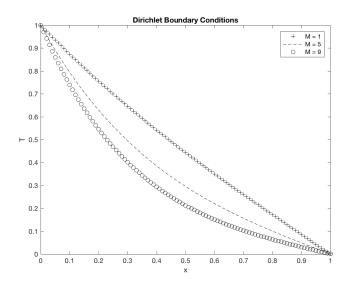
```
function [y] = tridiag(n, a, b, c, v);
   % solve a tridiagonal system with Gaussian elimination
   % Also known as Thomas Algorithm
   % (a_1 c_1 0
                  0
                                      0 ) (
                                            y_1
                                                   ) (
                                                          v_1
                                                                )
                                            у_2
   % (b_1 a_2 c_2 0
                                      0
                                        ) (
                                                  ) ( v_2 )
   % ( 0 b_2 a_3 c_3
                                      0 )(
                                                  ) = ( v_3 )
                                             у_3
   % ( 0
                                      0 ) (
                                                   ) = (
   % ( .
                          a_{n-1} c_{n-1} (y_{n-1}) (v_{n-1})
   % ( 0
                                    a_n)( y_n ) ( v_n )
   % create array zero to store solutions
   y = zeros(size(v));
   % eliminate b_i's
   for (i=1:n-1)
       a(i+1) = a(i+1) + c(i) * (-b(i)/a(i));
       v(i+1) = v(i+1) + v(i) * (-b(i)/a(i));
   end
   % solve for y_n
   y(n) = v(n)/a(n);
   for (i=(n-1):-1:1)
       y(i) = (v(i) - c(i)*y(i+1))/a(i);
end
```

The following script uses the previous function to run the Thomas algorithm on the tridiagonal system that was found earlier.

```
%% Problem 2a
n = 100;
deltaX = 1/n;
nM = 3;
sol = zeros(nM, n+1);
sol(:,1) = ones(nM,1)';
for iter = [1, 5, 9; 1:nM]
    M = iter(1);
    i = iter(2);
    mainDiagonal = (-M*deltaX^2 - 2)*ones(n-1,1);
    lowerDiagonal = ones(n-2,1);
    upperDiagonal = ones (n-2,1);
    RHS = zeros(n-1,1);
    % boundary conditions
    RHS (1) = -1;
    sol(i,2:end-1) = tridiag(n-1, mainDiagonal, lowerDiagonal, upperDiagonal, RHS);
end
x = linspace(0, 1, n+1);
plot(x, sol(1,:), 'k+', x, sol(2,:), 'k--', x, sol(3,:), 'ko');
legend('M = 1', 'M = 5', 'M = 9');
title('Dirichlet Boundary Conditions');
xlabel('x');
ylabel('T');
saveas(gcf, 'Figures/02_01.png', 'png');
% -dT/dx at x = 1 can be approximated by -dT/dx = (T_(n-1) - T_n)/deltaX
% T_n = sol(i, end) = sol(i, n+1)
heatFlux = zeros(nM, 1);
```

```
for i = 1:nM
    heatFlux(i) = (sol(i, end-1) - sol(i,end))/deltaX;
end
disp(heatFlux);
```

This script outputs the following image, which shows the three plots for M = 1, 5, 9.



The values of the heat flux at x = 1 are given to be

$$M = 1 \qquad \left( -\frac{\mathrm{d}T}{\mathrm{d}x} \right) \Big|_{x=1} = 0.850933420029442$$

$$M = 5 \qquad \left( -\frac{\mathrm{d}T}{\mathrm{d}x} \right) \Big|_{x=1} = 0.483548793162072$$

$$M = 9 \qquad \left( -\frac{\mathrm{d}T}{\mathrm{d}x} \right) \Big|_{x=1} = 0.299532256416841$$

Note that these values are found by evaluating

$$\left. \left( -\frac{\mathrm{d}T}{\mathrm{d}x} \right) \right|_{x=1} \approx \frac{T_{n-1} - T_n}{\Delta x}$$

(b) Compute a solution with the boundary conditions T(0) = 1,  $\frac{\mathrm{d}T(1)}{\mathrm{d}x} = 0$ . This corresponds to a fin that is insulated at one end. Solve for the temperature, T(1), at the insulated end for M = 1, 5, 9. Provide your three numerical values of T(1). Also plot T(x) for M = 9 with each pair of boundary conditions and compare to the exact solution.

For this problem we begin with the same difference equation as in part (a).

$$T_{i-1} + \left(-2 - M\Delta x^2\right)T_i + T_{i+1} = 0$$

Again we have the boundary condition that T(0) = 1, in results in the same modified equation for i = 1.

$$\left(-2 - M\Delta x^2\right)T_1 + T_2 = -1$$

However we have a different boundary condition at x=1. In this case we wish to enforce  $\frac{\mathrm{d}T(1)}{\mathrm{d}x}=0$ . In order to enforce this consdition we will consider an imaginary point on the fin  $T_{N+1}$ , this point doesn't actually exist on the fin, but if it did exist and the first derivative was zero then

$$\frac{T_{N+1} - T_N}{\Delta x} = 0.$$

In other words the finite difference for the first derivative should be zero, this simplifies to  $T_N = T_{N+1}$ . This should make intuitive sense, as if the derivative is zero the value shouldn't change past then endpoint. Using this condition in the difference equation for i = N gives

$$T_{N-1} + \left(-2 - M\Delta x^2\right)T_N + T_{N+1} = 0$$
  

$$T_{N-1} + \left(-2 - M\Delta x^2\right)T_N + T_N = 0$$
  

$$T_{N-1} + \left(-1 - M\Delta x^2\right)T_N = 0$$

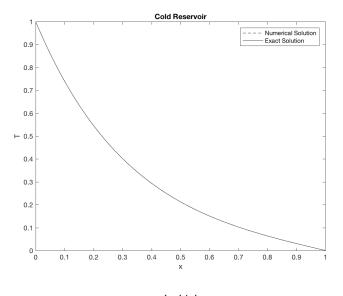
We now have N equations for N unknowns. This forms the following matrix equation.

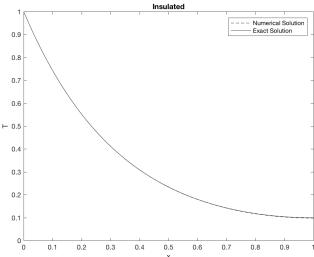
$$\begin{bmatrix} -2 - \Delta x^2 & 1 & & & 0 \\ 1 & -2 - \Delta x^2 & 1 & & & \\ & \ddots & \ddots & \ddots & & \\ & & 1 & -2 - \Delta x^2 & 1 \\ 0 & & & 1 & -1 - \Delta x^2 \end{bmatrix} \begin{bmatrix} T_1 \\ T_2 \\ \vdots \\ T_{N-1} \\ T_N \end{bmatrix} = \begin{bmatrix} -1 \\ 0 \\ \vdots \\ 0 \\ 0 \end{bmatrix}$$

Note that this system is one equation larger than in part (a) as the value of  $T_N$  must be found. Still this is a tridiagonal system that can be solved using the Thomas algorithm, which was shown in part (a). The following script using the Thomas algorithm to solve this system for M = 1, 5, 9.

```
%% Problem 2b
n = 100;
deltaX = 1/n;
nM = 3;
sol2 = zeros(nM, n+1);
sol2(:,1) = ones(nM,1)';
for iter = [1, 5, 9; 1:nM]
   M = iter(1);
    i = iter(2);
    mainDiagonal = (-M*deltaX^2 - 2)*ones(n,1);
    mainDiagonal(end) = (-M*deltaX^2 - 1);
    lowerDiagonal = ones(n-1,1);
    upperDiagonal = ones (n-1, 1);
    RHS = zeros(n, 1);
    % boundary conditions
    RHS (1) = -1;
    sol2(i,2:end) = tridiag(n, mainDiagonal, lowerDiagonal, upperDiagonal, RHS);
exactSol1 = @(x) (exp(-3)/(exp(-3) - exp(3))) * exp(3*x) + (-exp(3)/(exp(-3) - exp(-3)))
   \hookrightarrow (3))) * exp(-3*x)
exactSol2 = @(x) (exp(-3)/(exp(-3) + exp(3))) * exp(3*x) + (exp(3)/(exp(-3) + exp(3)))
   \hookrightarrow (3))) * exp(-3*x)
x = linspace(0, 1, n+1);
plot(x, sol(3,:), 'k--', x, exactSol1(x), 'k-');
legend('Numerical Solution', 'Exact Solution ');
title('Cold Reservoir');
xlabel('x');
ylabel('T');
saveas(gcf, 'Figures/02_02.png', 'png');
```

The script output the following images. Note that the numerical solutions and the exact solutions are almost indistinguishable.





The following are the numerical values found for T(1) for M=1,5,9.

$$M = 1$$
  $T(1) = 0.645597948372887$ 

$$M = 5$$
  $T(1) = 0.209066847236553$ 

$$M = 9$$
  $T(1) = 0.097878298010251$ 

(c) Add a distributed heat source: Compute and plot a solution of the non-homogeneous equation

$$\frac{\mathrm{d}^2 T}{\mathrm{d}x^2} = MT - 100x^2(1-x)^2$$

with M = 9, T(0) = 1 and  $\frac{dT(1)}{dx} = 0$ .

2. (i) What type of p.d.e. is

$$\frac{\partial^2 \phi}{\partial x \partial y} + \phi = 25?$$

(ii) What type of p.d.e. does the velocity potential,  $\phi$ , satisfy if

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

with

$$u = \frac{\partial \phi}{\partial x} \qquad v = \frac{\partial \phi}{\partial y}?$$

(iii) The boundary layer momentum equation is

$$u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y} = \frac{1}{Re}\frac{\partial^2 u}{\partial y^2}$$

where Re is the Reynolds number. What type is this equation?