

Caleb Logemann
MATH666 Finite Element Methods
Homework 1

#1

- (a) In order to recast this as a variational problem we will multiply the equation by a test function and integrate.

$$\begin{aligned} -u'' + qu &= f \\ -\int_{-\pi}^{\pi} u''v \, dx + \int_{-\pi}^{\pi} quv \, dx &= \int_{-\pi}^{\pi} fv \, dx \\ \int_{-\pi}^{\pi} u'v' \, dx - (u'v)|_{-\pi}^{\pi} + \int_{-\pi}^{\pi} quv \, dx &= \int_{-\pi}^{\pi} fv \, dx \end{aligned}$$

If we let $v(\pi) = v(-\pi)$, then the boundary term goes to zero because $u'(\pi) = u'(-\pi)$

$$\int_{-\pi}^{\pi} u'v' \, dx + \int_{-\pi}^{\pi} quv \, dx = \int_{-\pi}^{\pi} fv \, dx$$

Thus the variational problem is to find a $u \in V = \left\{ \int_{-\pi}^{\pi} ((u')^2 + qu^2) \, dx < \infty \mid u(\pi) = u(-\pi), u'(\pi) = u'(-\pi) \right\}$

$$\int_{-\pi}^{\pi} u'v' \, dx + \int_{-\pi}^{\pi} quv \, dx = \int_{-\pi}^{\pi} fv \, dx$$

for all $v \in V$. The test function v must satisfy the periodic boundary conditions on the function value, but I believe we can also choose to have v satisfy the periodic condition on the derivative as well. This makes the test and trial function spaces the same.

- (b) The energy minimization problem that is equivalent to the variational problem is to find a $u \in V$ such that

$$F(u) \leq F(w)$$

for all $w \in V$, where

$$F(w) = \frac{1}{2} \left(\int_{-\pi}^{\pi} (w')^2 \, dx + \int_{-\pi}^{\pi} qw^2 \, dx \right) - \int_{-\pi}^{\pi} fw \, dx$$

I will now prove that the energy minimization problem and the variational problem are equivalent.

Proof. Let u be a solution to the variational problem, and consider some $w \in V$. Then there exists $v \in V$ such that $u + v = w$. Now consider $F(w)$.

$$\begin{aligned} F(w) &= F(u + v) \\ &= \frac{1}{2} \left(\int_{-\pi}^{\pi} (u' + v')^2 \, dx + \int_{-\pi}^{\pi} q(u + v)^2 \, dx \right) - \int_{-\pi}^{\pi} f(u + v) \, dx \\ &= \frac{1}{2} \left(\int_{-\pi}^{\pi} (u')^2 + 2u'v' + (v')^2 \, dx + \int_{-\pi}^{\pi} qu^2 + 2quv + qv^2 \, dx \right) - \int_{-\pi}^{\pi} f(u + v) \, dx \\ &= \frac{1}{2} \left(\int_{-\pi}^{\pi} (u')^2 \, dx + \int_{-\pi}^{\pi} qu^2 \, dx \right) - \int_{-\pi}^{\pi} fu \, dx + \int_{-\pi}^{\pi} u'v' \, dx + \int_{-\pi}^{\pi} quv \, dx - \int_{-\pi}^{\pi} fv \, dx + \int_{-\pi}^{\pi} (v')^2 + qv^2 \, dx \\ &= F(u) + \int_{-\pi}^{\pi} u'v' \, dx + \int_{-\pi}^{\pi} quv \, dx - \int_{-\pi}^{\pi} fv \, dx + \int_{-\pi}^{\pi} (v')^2 + qv^2 \, dx \end{aligned}$$

Since u is a solution to the variational problem the middle three terms cancel

$$= F(u) + \int_{-\pi}^{\pi} (v')^2 + qv^2 dx$$

Also since $q(x) > 0$ for $x \in [-\pi, \pi]$, this integral is nonnegative

$$\geq F(u)$$

This shows that u is a solution to the energy minization problem when u is a solution to the variational problem.

Now let u be a solution to the energy minimization problem, then

$$F(u) \leq F(w)$$

for all $w \in V$. Let $v \in V$ and consider $w = u + \varepsilon v$, then $F(u) \leq F(u + \varepsilon v)$. Consider the function of ε ,

$$g(\varepsilon) = F(u + \varepsilon v)$$

we know that g has a minimum at $\varepsilon = 0$ thus $g'(0) = 0$.

$$\begin{aligned} g(\varepsilon) &= \frac{1}{2} \left(\int_{-\pi}^{\pi} (u' + \varepsilon v')^2 dx + \int_{-\pi}^{\pi} q(u + \varepsilon v)^2 dx \right) - \int_{-\pi}^{\pi} f(u + \varepsilon v) dx \\ &= \frac{1}{2} \left(\int_{-\pi}^{\pi} (u')^2 dx + \int_{-\pi}^{\pi} qu^2 dx \right) - \int_{-\pi}^{\pi} fu dx + \varepsilon \int_{-\pi}^{\pi} u'v' dx + \varepsilon \int_{-\pi}^{\pi} quv dx - \varepsilon \int_{-\pi}^{\pi} fv dx + \varepsilon^2 \int_{-\pi}^{\pi} (v')^2 + qv^2 dx \\ g'(\varepsilon) &= \int_{-\pi}^{\pi} u'v' dx + \int_{-\pi}^{\pi} quv dx - \int_{-\pi}^{\pi} fv dx + 2\varepsilon \int_{-\pi}^{\pi} (v')^2 + qv^2 dx \\ g'(0) &= \int_{-\pi}^{\pi} u'v' dx + \int_{-\pi}^{\pi} quv dx - \int_{-\pi}^{\pi} fv dx \\ 0 &= \int_{-\pi}^{\pi} u'v' dx + \int_{-\pi}^{\pi} quv dx - \int_{-\pi}^{\pi} fv dx \\ \int_{-\pi}^{\pi} fv dx &= \int_{-\pi}^{\pi} u'v' dx + \int_{-\pi}^{\pi} quv dx \end{aligned}$$

Since this is true for any $v \in V$, this shows that u is a solution to the variational problem. \square

(c) Consider $u, w \in V$ solutions to the variational, then for any $v \in V$ we have

$$\begin{aligned} \int_{-\pi}^{\pi} u'v' dx + \int_{-\pi}^{\pi} quv dx &= \int_{-\pi}^{\pi} fv dx \\ \int_{-\pi}^{\pi} w'v' dx + \int_{-\pi}^{\pi} qwv dx &= \int_{-\pi}^{\pi} fv dx \\ \int_{-\pi}^{\pi} (u - w)'v' dx + \int_{-\pi}^{\pi} q(u - w)v dx &= 0 \end{aligned}$$

Thus the Galerkin Orthogonality property of this problem is

$$\int_{-\pi}^{\pi} (u - w)'v' dx + \int_{-\pi}^{\pi} q(u - w)v dx = 0$$

for u, w solutions to the variational problem and for all $v \in V$.

The energy norm for this problem is found by letting $v = u - w$ and taking the squareroot.

$$\|u\|_E = \sqrt{\int_{-\pi}^{\pi} (u')^2 dx + \int_{-\pi}^{\pi} qu^2 dx}$$

Note that an energy inner product can be formed as well,

$$[u, v]_E = \int_{-\pi}^{\pi} u'v' dx + \int_{-\pi}^{\pi} quv dx \quad \|u\|_E = \sqrt{[u, u]_E}$$

This satisfies all the properties of an inner product.

$$\begin{aligned} [u, v]_E &= [v, u]_E \\ [au, v]_E &= a[u, v]_E \\ [u + w, v]_E &= [u, v]_E + [w, v]_E \\ [u, u]_E &\geq 0 \\ [u, u]_E &= 0 \Leftrightarrow u = 0 \end{aligned}$$

Note that $q > 0$ is required for the last two statements. The fact that this forms an inner product allows the Cauchy-Schwarz inequality to be applied directly to the energy norm. Also the Galerkin Orthogonality condition can be expressed as

$$[u - w, v]_E = 0$$

for u, w solutions to the variational problem and for all $v \in V$.

- (d) The cG(1) method for this problem is formulated by replacing the test and trial space V with a subspace, V_h^1 . Let $-\pi = x_0 < x_1 < \dots < x_M < x_{M+1} = \pi$ be a partition of $[-\pi, \pi]$, and define $h_j = x_j - x_{j-1}$ and let $h = \max_{1 \leq j \leq M+1} \{h_j\}$. I will also define the functions

$$\phi_j(x) = \begin{cases} \frac{x-x_{j-1}}{x_j-x_{j-1}} & x_{j-1} \leq x \leq x_j \\ \frac{x-x_{j+1}}{x_j-x_{j+1}} & x_j \leq x \leq x_{j+1} \\ 0 & \text{otherwise} \end{cases}$$

for $j = 1, 2, \dots, M$ and

$$\phi_0(x) = \phi_{M+1}(x) = \begin{cases} \frac{x-x_0}{x_1-x_0} & x_0 \leq x \leq x_1 \\ \frac{x-x_{M+1}}{x_M-x_{M+1}} & x_M \leq x \leq x_{M+1} \\ 0 & \text{otherwise} \end{cases}$$

Then the space $V_h^1 = \left\{ \sum_{j=1}^{M+1} (\xi_j \phi_j) \mid \xi_{M+1} = \frac{h_1 h_{M+1}}{h_1 + h_{M+1}} \left(\frac{\xi_1}{h_1} + \frac{\xi_M}{h_{M+1}} \right) \right\}$. The condition on ξ_{M+1} guarantees the continuity of the derivative at $x_0 = x_{M+1}$. The design of the basis function ϕ_{M+1} guarantees the continuity of the function.

- (e) The following will show that the cG(1) method gives the optimal solution in the energy norm.

Proof. Let u be the solution to the original variational and let U be the solution to the cG(1) method. Consider an arbitrary $v \in V_h^1$, and note that Galerkin orthogonality states that $[u - U, U - v]_E = 0$.

$$\begin{aligned} \|u - U\|_E^2 &= [u - U, u - U]_E \\ &= [u - U, u - U]_E + [u - U, U - v]_E \\ &= [u - U, u - v]_E \end{aligned}$$

Using Cauchy-Schwarz

$$\leq \|u - U\|_E \|u - v\|_E$$

This shows that $\|u - U\|_E \leq \|u - v\|_E$. □

Now in order to find an error estimate let $v = \pi_h u$.

$$\begin{aligned} \|u - U\|_E^2 &\leq \|u - \pi_h u\|_E^2 \\ &= \int_{-\pi}^{\pi} ((u - \pi_h u)')^2 dx + \int_{-\pi}^{\pi} q(u - \pi_h u)^2 dx \\ &\leq Ch^2 \|u''\|^2 + Ch^4 \|u''\|^2 \end{aligned}$$

Thus

$$\|u - U\|_E \leq Ch \|u''\|$$

since the h term dominates the error.

(f) Now I will find an error estimate in the L^2 norm.

First note that this problem is self adjoint, and so the dual problem is

$$\begin{aligned} -\phi'' + q\phi &= e \\ \phi(-\pi) &= \phi(\pi) \\ \phi'(-\pi) &= \phi'(\pi) \end{aligned}$$

Now consider $\|e\|_{L^2}^2$,

$$\|e\|_{L^2}^2 = \int_{-\pi}^{\pi} e^2 dx$$

Using the dual problem

$$\begin{aligned} &= \int_{-\pi}^{\pi} e(-\phi'' + q\phi) dx \\ &= - \int_{-\pi}^{\pi} e\phi'' dx + \int_{-\pi}^{\pi} qe\phi dx \\ &= \int_{-\pi}^{\pi} e'\phi' dx + \int_{-\pi}^{\pi} qe\phi dx \end{aligned}$$

Using Galerkin Orthogonality

$$\begin{aligned} &= \int_{-\pi}^{\pi} e'\phi' dx + \int_{-\pi}^{\pi} qe\phi dx - \int_{-\pi}^{\pi} e'\pi_h\phi' dx - \int_{-\pi}^{\pi} qe\pi_h\phi dx \\ &= \int_{-\pi}^{\pi} e'(\phi - \pi_h\phi)' dx + \int_{-\pi}^{\pi} qe(\phi - \pi_h\phi) dx \\ &= [e, \phi - \pi_h\phi]_E \end{aligned}$$

Now applying Cauchy-Schwarz

$$\leq \|e\|_E \|\phi - \pi_h\phi\|_E$$

From part (e)

$$\leq Ch\|u''\|\|\phi - \pi_h\phi\|_E$$

Now consider $\|\phi - \pi_h\phi\|_E$

$$\begin{aligned}\|\phi - \pi_h\phi\|_E^2 &= \int_{-\pi}^{\pi} ((\phi - \pi_h\phi)')^2 dx + \int_{-\pi}^{\pi} q(\phi - \pi_h\phi)^2 dx \\ &\leq C_1 h^2 \|\phi''\|^2 + C_2 h^4 \|\phi''\|^2\end{aligned}$$

Since the h^2 term dominates, the h^4 term can be incorporated into C_1

$$\leq C_1 h^2 \|\phi''\|^2$$

Therefore $\|\phi - \pi_h\phi\|_E \leq Ch\|\phi''\|$.

Next I will show that $\|\phi''\| \leq C\|e\|$. Consider the dual problem again.

$$\begin{aligned}-\phi'' + q\phi &= e \\ \int_{-\pi}^{\pi} (-\phi'' + q\phi)^2 dx &= \int_{-\pi}^{\pi} e^2 dx \\ \int_{-\pi}^{\pi} (\phi'')^2 - 2q\phi''\phi + q^2\phi^2 dx &= \int_{-\pi}^{\pi} e^2 dx \\ \|\phi''\|^2 + \int_{-\pi}^{\pi} -2q\phi''\phi + q^2\phi^2 dx &= \|e\|^2 \\ \|\phi''\|^2 + \int_{-\pi}^{\pi} 2q(\phi')^2 + q^2\phi^2 dx &= \|e\|^2 \\ \|\phi''\|^2 &= \|e\|^2 - \int_{-\pi}^{\pi} 2q(\phi')^2 + q^2\phi^2 dx\end{aligned}$$

Since $\int_{-\pi}^{\pi} 2q(\phi')^2 + q^2\phi^2 dx$ is strictly positive

$$\|\phi''\|^2 \leq \|e\|^2$$

Now putting this all together shows that

$$\|e\|_{L^2}^2 \leq Ch\|u''\|_{L^2}\|\phi - \pi_h\phi\|_E \leq Ch^2\|u''\|_{L^2}\|\phi''\|_{L^2} \leq Ch^2\|u''\|_{L^2}\|e\|_{L^2}$$

Dividing both side by $\|e\|_{L^2}$ gives

$$\|e\|_{L^2} \leq Ch^2\|u''\|_{L^2}$$

(g)

#2